



Higgs 2022 | PISA

Constraints on anomalous Higgs boson couplings and EFT with the CMS experiment



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on behalf of the CMS Collaboration

November 8th, 2022

Motivation |

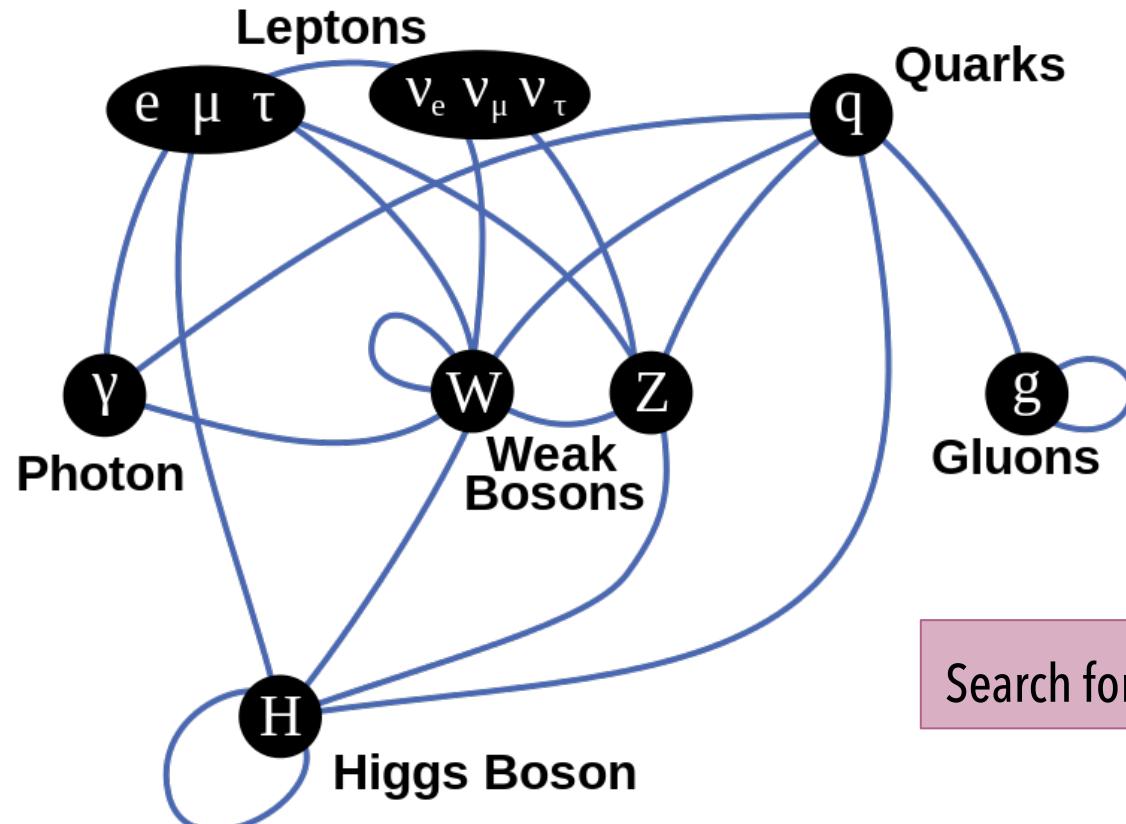
10 years studying the Higgs boson

Run 1 → LHC data compatible with an SM $H(125)$ of spin 0 and even CP

Run 2 → Precision era for Higgs physics +

JCP studies:

- Coupling to bosons
- Coupling to fermions



CP is a property of the interaction

&

Matter asymmetry
CP-violating baryon-generating interaction

Search for Higgs Anomalous Couplings (AC) to bosons and fermions

Higgs AC to bosons | HVV parametrization

$$A(HVV) \sim \left[a_1^{VV} + \frac{\kappa_1^{VV} q_1^2 + \kappa_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} \right] m_{V1}^2 \epsilon_{V1}^* \epsilon_{V2}^* + a_2^{VV} f_{\mu\nu}^{*(1)} f^{*(2)\mu\nu} + a_3^{VV} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2)\mu\nu}$$

→ If VV= **WW, ZZ, Zγ** → Tree-level SM → If VV= **gg** → 1-loop SM

AC approach	SMEFT approach
$a_i^{ZZ} = a_i^{WW}$	$SU(2) \times U(1)$
4 anomalous couplings:	$a_i^{ZZ} \neq a_i^{WW}$
a_2 (CP)	3 anomalous couplings:
a_3 (CP)	a_2 (CP)
a_{A1} (CP)	a_3 (CP)
$a_{A1}Z\gamma$ (CP)	a_{A1} (CP)

1. Amplitude parametrization can be related to the **Higgs basis EFT**.
2. **Cross-section fraction:**

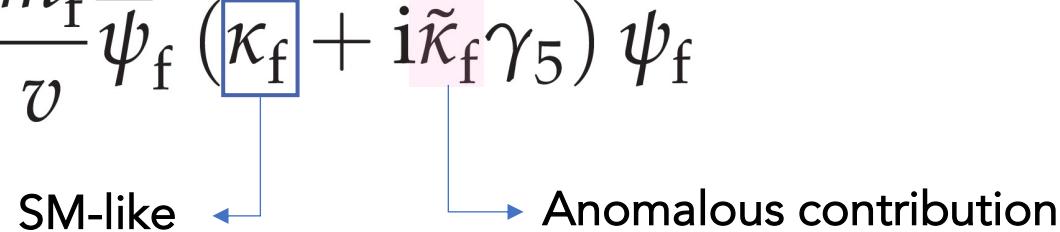
$$f_{ai} = \frac{|a_i|^2 \sigma_i}{\sum_{j=1,2,3\dots} |a_j|^2 \sigma_j} \text{sign} \left(\frac{a_i}{a_1} \right)$$

If VV= **gg** → 1 AC: **a_3 (~~CP~~)**

Higgs AC to fermions | Hff parametrization

$$A(\text{Hff}) = -\frac{m_f}{v} \bar{\psi}_f (\kappa_f + i\tilde{\kappa}_f \gamma_5) \psi_f$$

SM-like Anomalous contribution



AC approach/SMEFT approach

1 Anomalous coupling:

$$\tilde{\kappa}_f : \mathbb{CP}$$

Cross-section fraction:

$$f_{CP}^{\text{Hff}} = \frac{|\tilde{\kappa}_f|^2}{|\kappa_f|^2 + |\tilde{\kappa}_f|^2} \text{ sign} \left(\frac{\tilde{\kappa}_f}{\kappa_f} \right)$$

κ_f and $\tilde{\kappa}_f$ are Yukawa coupling strength modifiers related to the mixing angle $\alpha^{Hff} = \tan^{-1} \left(\frac{\tilde{\kappa}_f}{\kappa_f} \right)$

Higgs couplings analyses with the CMS experiment I

Since 2021:

CMS Analysis	Channel	Measurement	Combined with	REF
HIG-19-009	on-Shell $H \rightarrow ZZ$	HVV, Hgg, Htt	[Htt] $H \rightarrow \gamma\gamma$ (HIG-19-013)	PRD 104 (2021) 052004
HIG-21-013	off-Shell $H \rightarrow ZZ$	HVV	on-Shell $H \rightarrow ZZ$	NP (2022) 01682
HIG-20-007	$H \rightarrow \tau\tau$	HVV, Hgg, Htt	on-Shell $H \rightarrow ZZ + H \rightarrow \gamma\gamma$	arXiv:2205.05120 (PRD)
HIG-21-006	ttH and tH	Htt	on-Shell $H \rightarrow ZZ + H \rightarrow \gamma\gamma$	arXiv:2208.02686 (sub)
HIG-20-006	$H \rightarrow \tau\tau$	$H\tau\tau$	-	JHEP 06 (2022) 012

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“Constraints on anomalous Higgs boson couplings to vector
bosons and fermions in its production and decay using the
four-lepton final state”

On-Shell H \rightarrow ZZ analysis

Channels: 4ℓ ($4e$, 4μ , $2e2\mu$)

Productions: ggH, VBF, VH, ttH.

Dedicated measurements (detector-level)

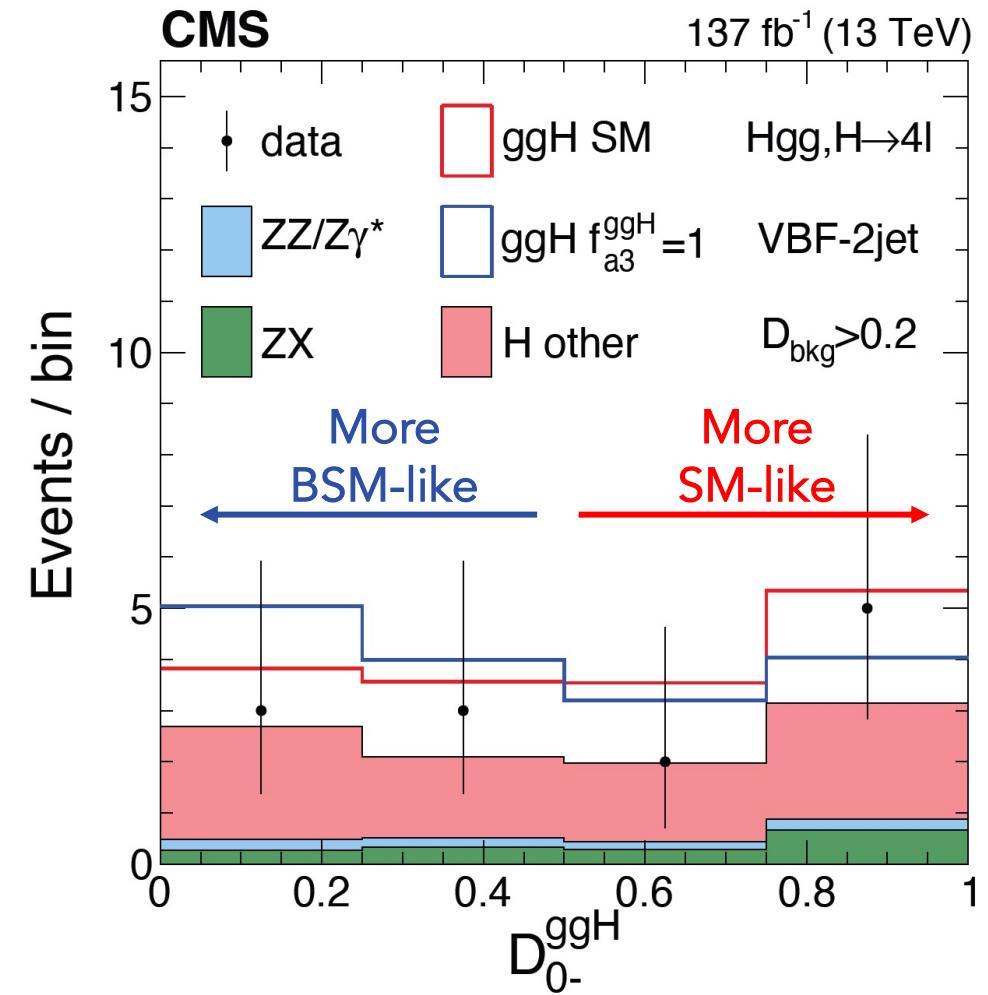
- Exploit full production and decay kinematic information with MELA-based discriminants.

$$\mathcal{D}_{\text{alt}}(\Omega) = \frac{\mathcal{P}_{\text{sig}}(\Omega)}{\mathcal{P}_{\text{sig}}(\Omega) + \mathcal{P}_{\text{alt}}(\Omega)}$$

$$\mathcal{D}_{\text{int}}(\Omega) = \frac{\mathcal{P}_{\text{int}}(\Omega)}{2 \sqrt{\mathcal{P}_{\text{sig}}(\Omega) \mathcal{P}_{\text{alt}}(\Omega)}}$$

Multi-dimensional analysis:

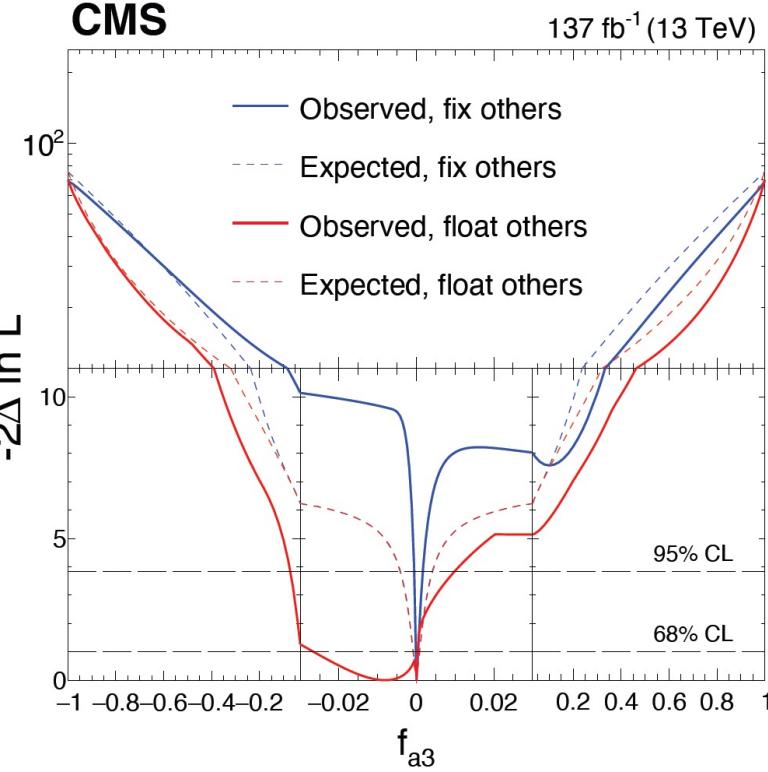
- Several parameters are simultaneously extracted.



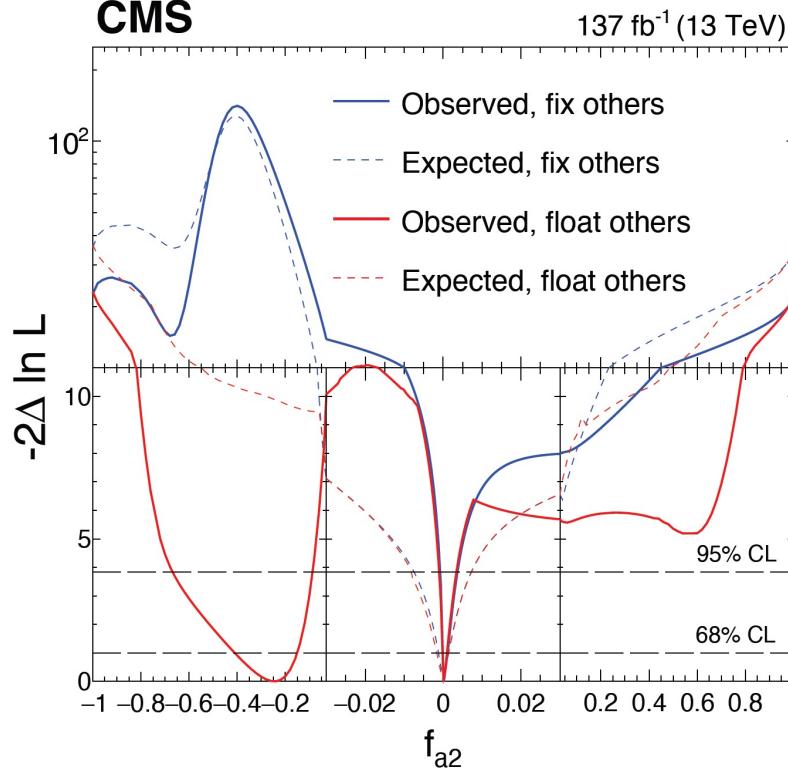
H \rightarrow ZZ Results | HVV (AC approach)

VBF & VH & H \rightarrow 4 ℓ

CMS



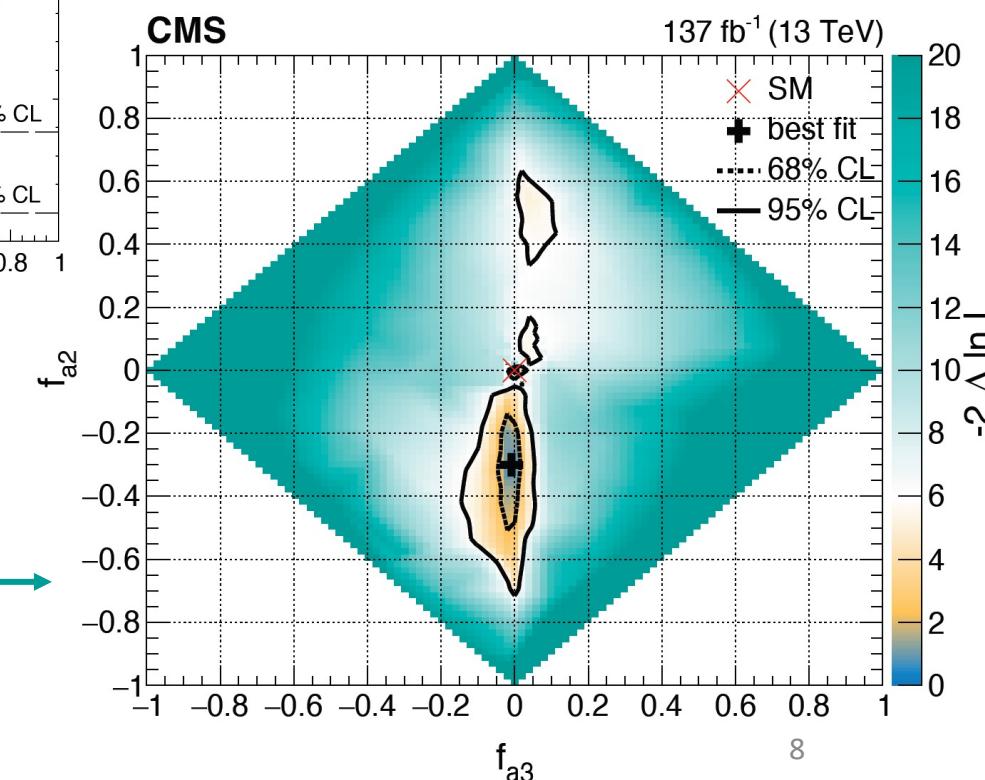
CMS



- Fix:** Other couplings are fixed to the SM expectation.
- Float:** Other couplings are profiled in the fit



f_{a1} and $f_{a1} Z\gamma$ scans are also available
Notice we can also perform f_{a3} , f_{a2} ,
and f_{a1} scans in the SMEFT approach



H \rightarrow ZZ Results | HVV (SMEFT approach)

VBF & VH & H \rightarrow 4 ℓ

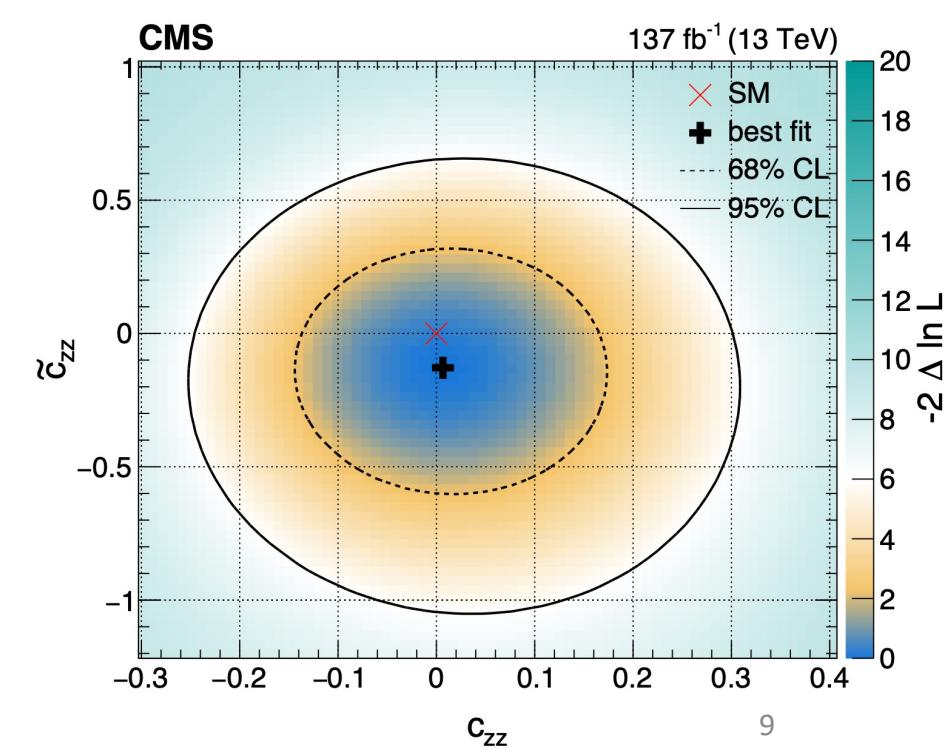
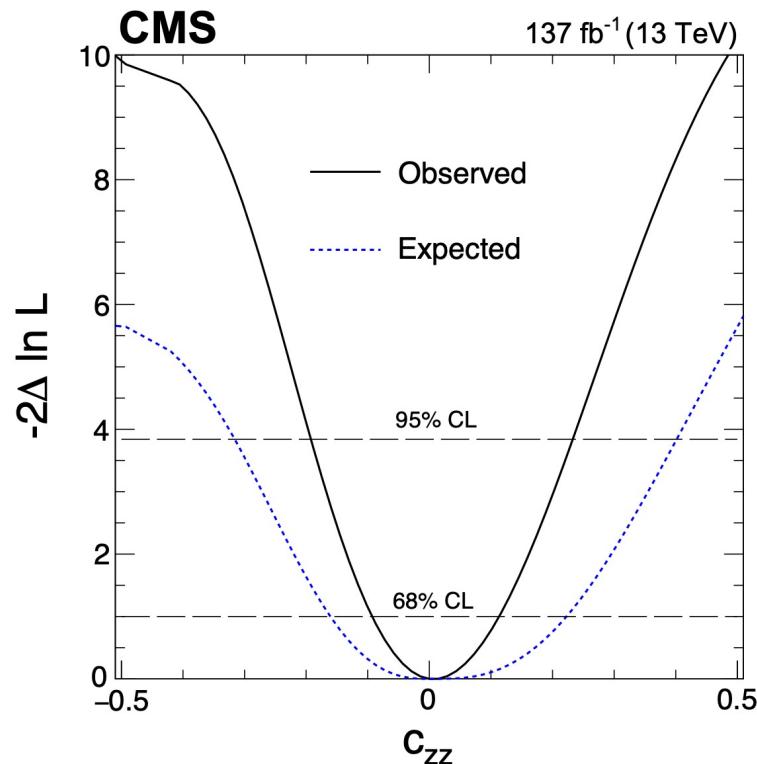
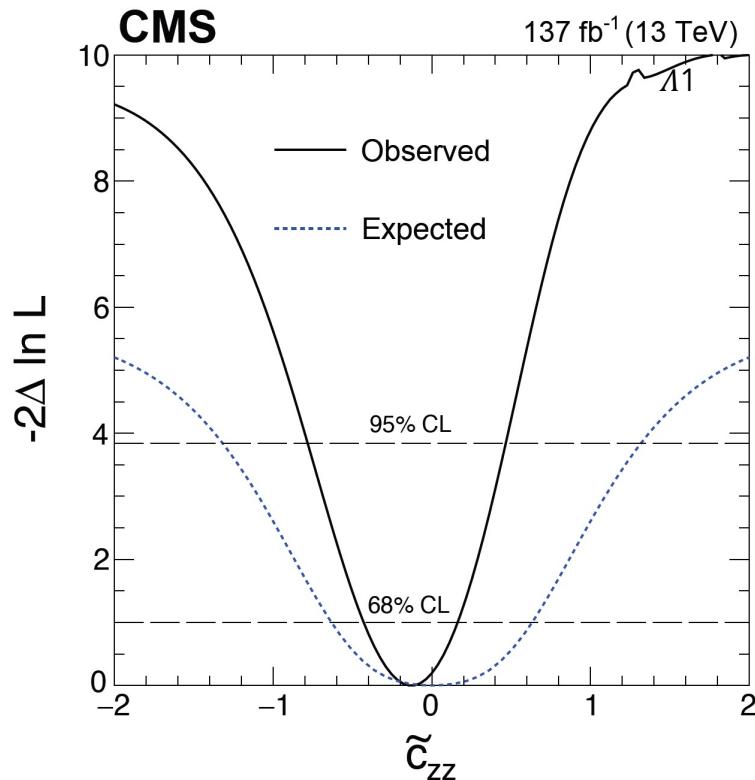
HVV parametrization can be related to the

EFT couplings using the Higgs basis.

+ Results on Warsaw Basis

$$\delta c_z = \frac{1}{2}a_1 - 1 \quad c_{zz} = -\frac{s_w^2 c_w^2}{2\pi\alpha} a_2$$

$$c_{z\square} = \frac{m_Z^2 s_w^2}{4\pi\alpha} \frac{\kappa_1}{(\Lambda_1)^2} \quad \tilde{c}_{zz} = -\frac{s_w^2 c_w^2}{2\pi\alpha} a_3$$



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“Measurement of the Higgs boson width and **evidence** of its off-shell contributions to ZZ production”

Off-Shell H \rightarrow ZZ analysis |

Channels: 2 ℓ 2 ν /4 ℓ

Productions: ggH, VBF, VH.

m_{ZZ} line shape is sensitive to the presence of anomalous HVV

- Could affect the measurement of Γ_{Higgs}

$$\sigma^{\text{off-Shell}} \propto \sigma^{\text{on-Shell}} \Gamma_{\text{Higgs}}$$

2 ℓ 2 ν \rightarrow m_T^{ZZ}>300 GeV

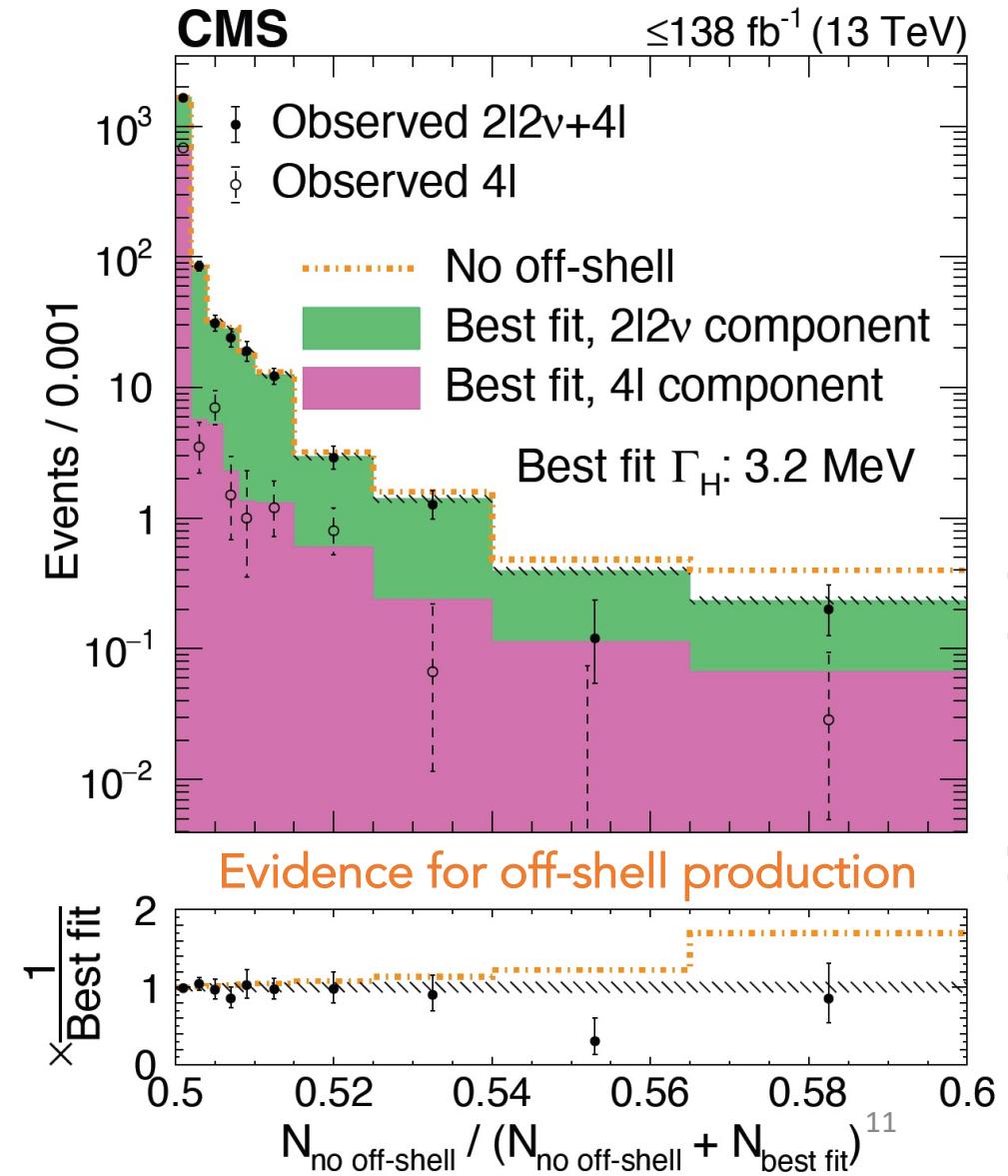
4 ℓ events \rightarrow MELA D^{VBF}_{2j} \rightarrow Sensitive to AC HVV.

Combined with:

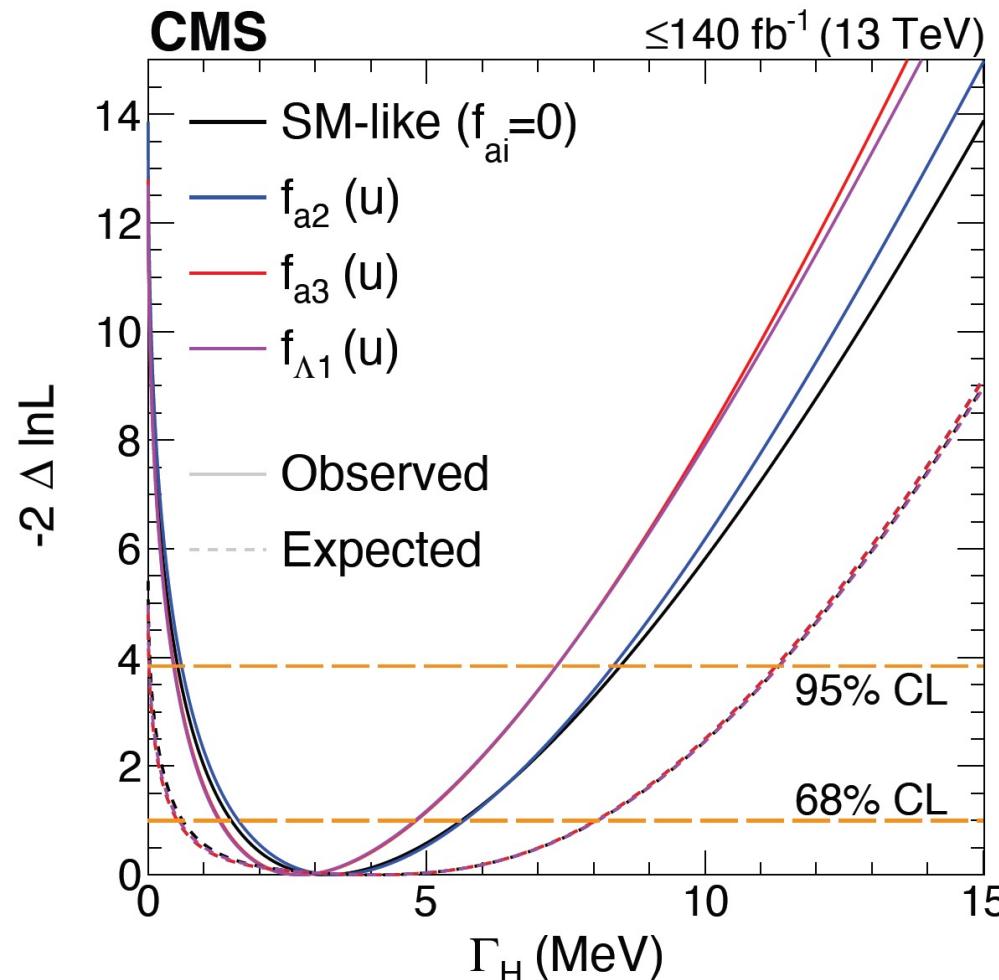
Run 2 On-Shell 4 ℓ \rightarrow PRD 104 (2021) 052004

(2016-2017) Off-Shell 4 ℓ \rightarrow Phys. Rev. D 99, 112003 (2019)

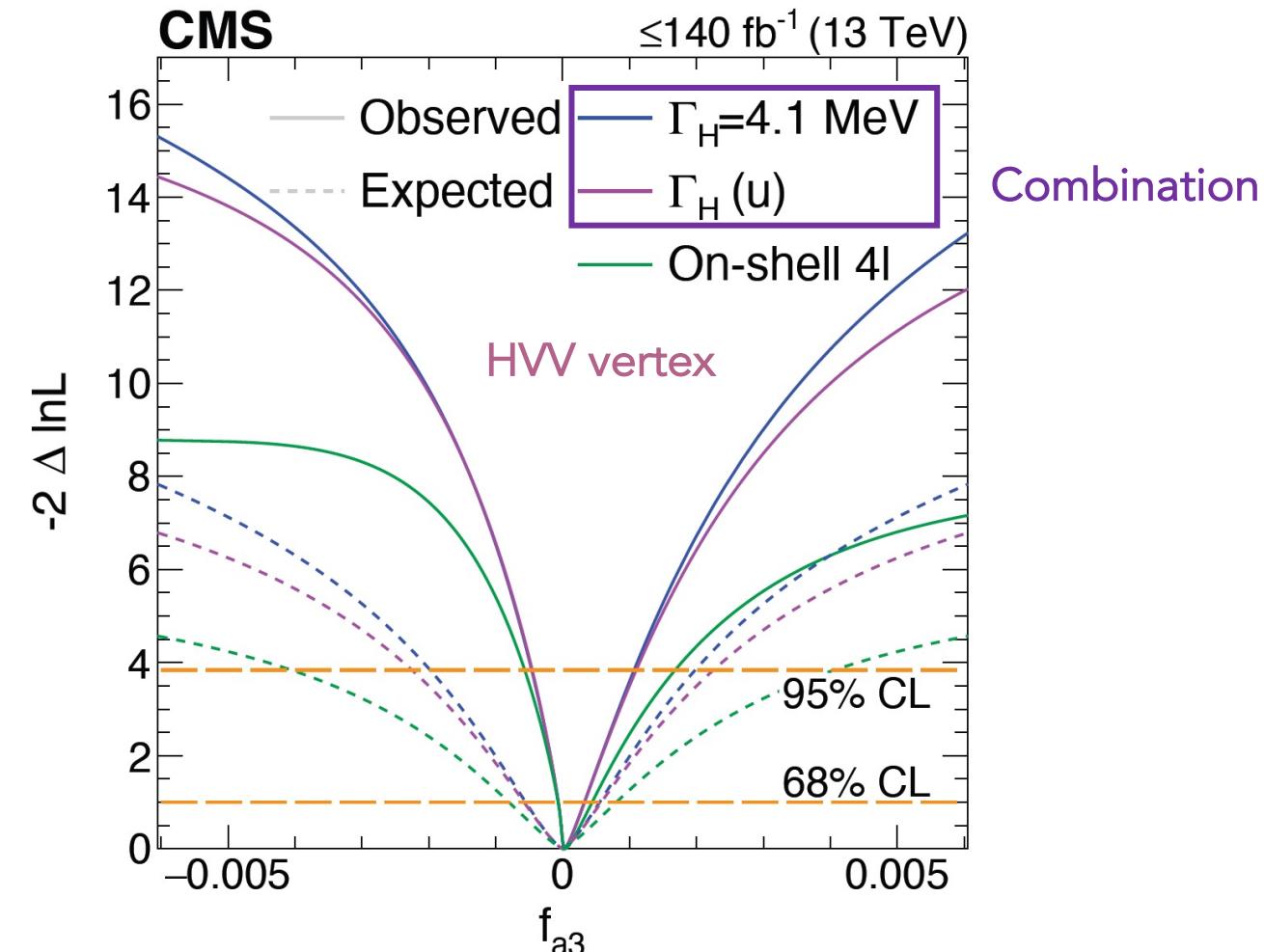
2015 On-Shell 4 ℓ \rightarrow Phys. Lett. B 775 (2017) 1



Off-Shell $H \rightarrow ZZ$ analysis Results | Combination with on-Shell $H \rightarrow ZZ$



$\Gamma_H = 3.2 [+2.4, -1.7] \text{ MeV (68\% CL)}$



AC approach

Higgs couplings analyses with the CMS experiment I

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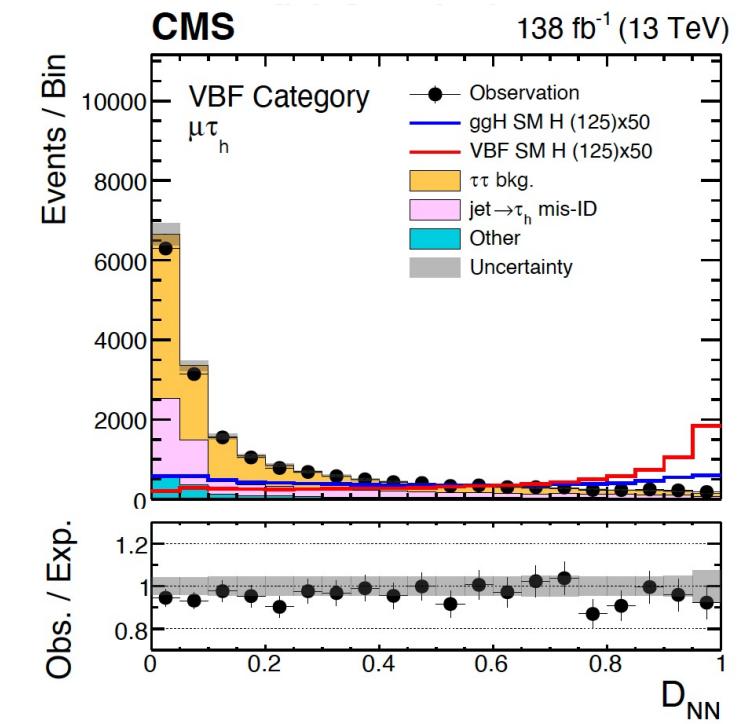
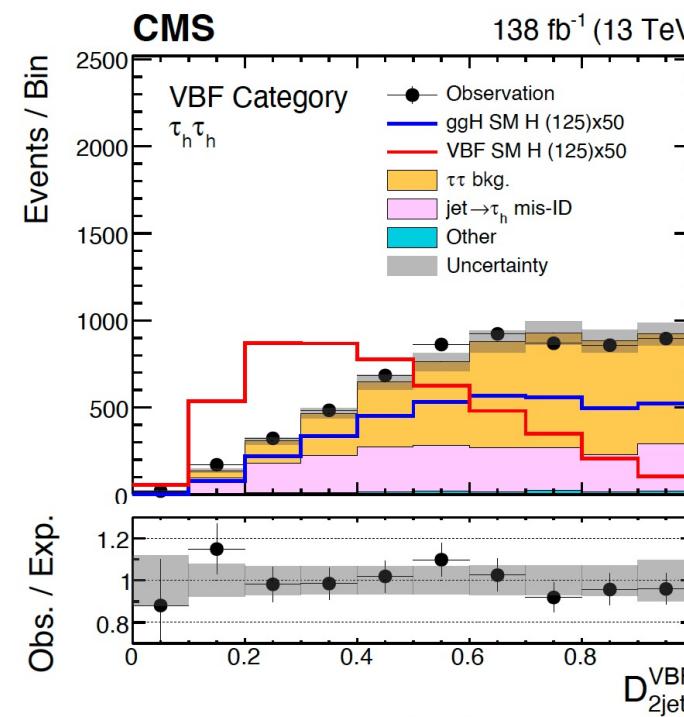
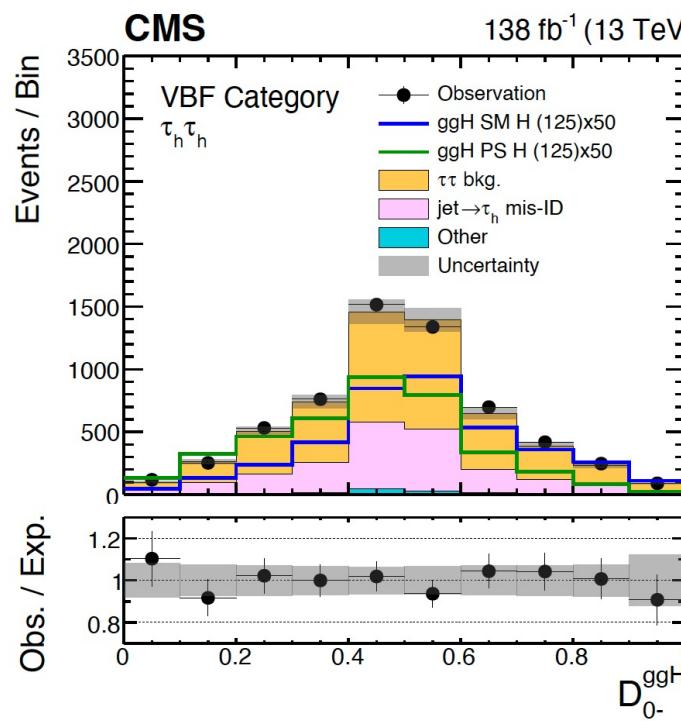
“Constraints on anomalous Higgs boson couplings to vector
bosons and fermions from the production of Higgs bosons
using the $\tau\tau$ final state”

H $\rightarrow\tau\tau$ analysis |

Channels: $\tau_h\tau_h$, $\mu\tau_h$, $e\tau_h$, $e\mu$

VBF production analysis \rightarrow HVV

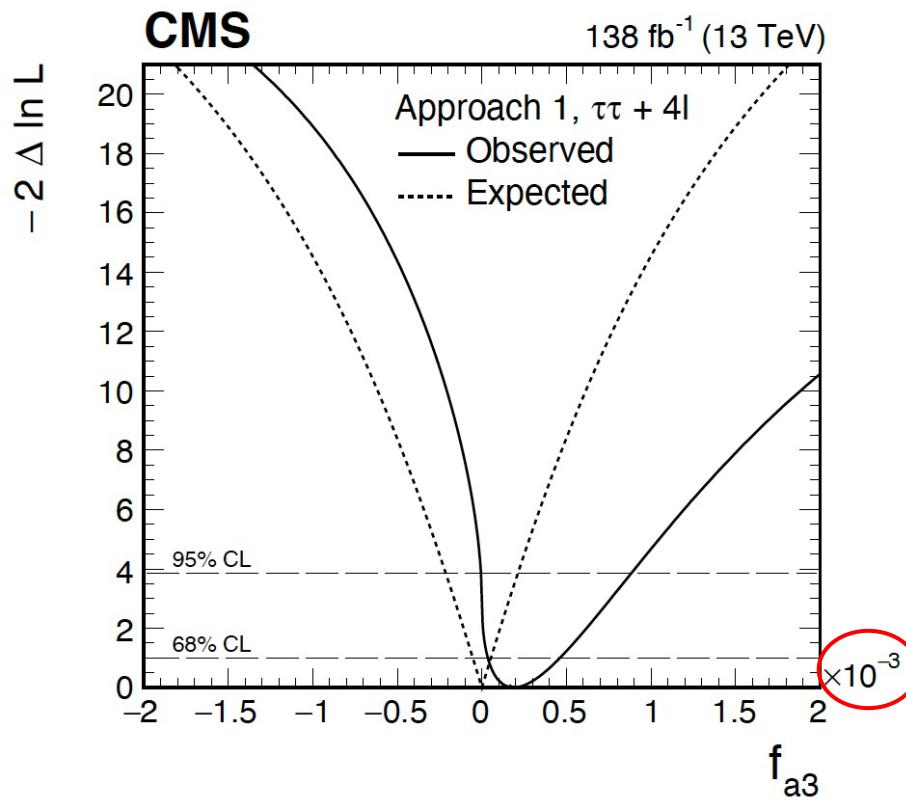
ggH production analysis \rightarrow Hgg



Combined to construct 2D and 3D [kinematic discriminants](#)

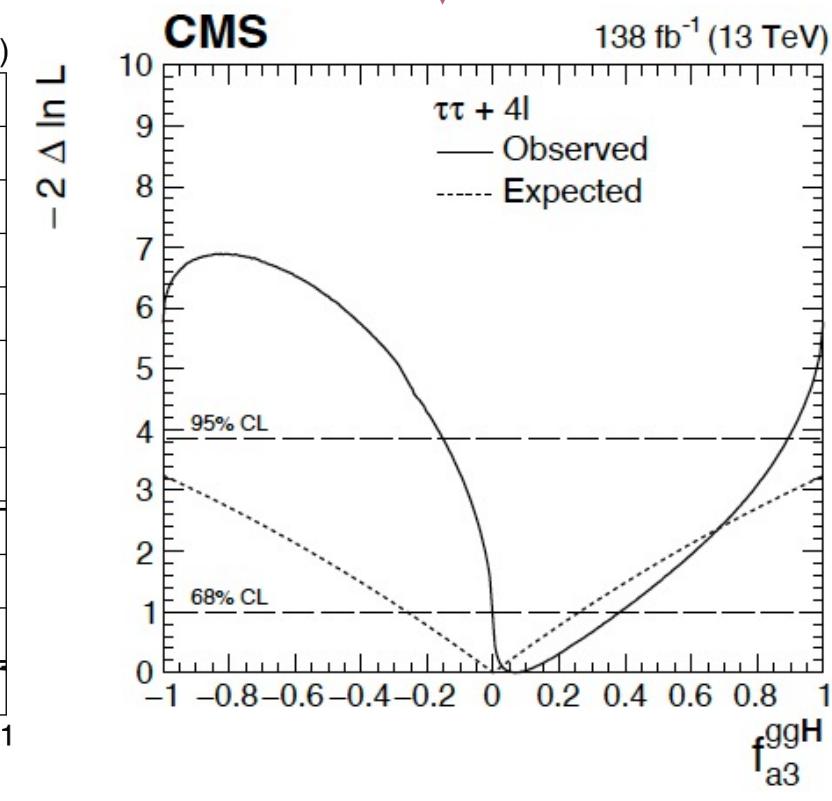
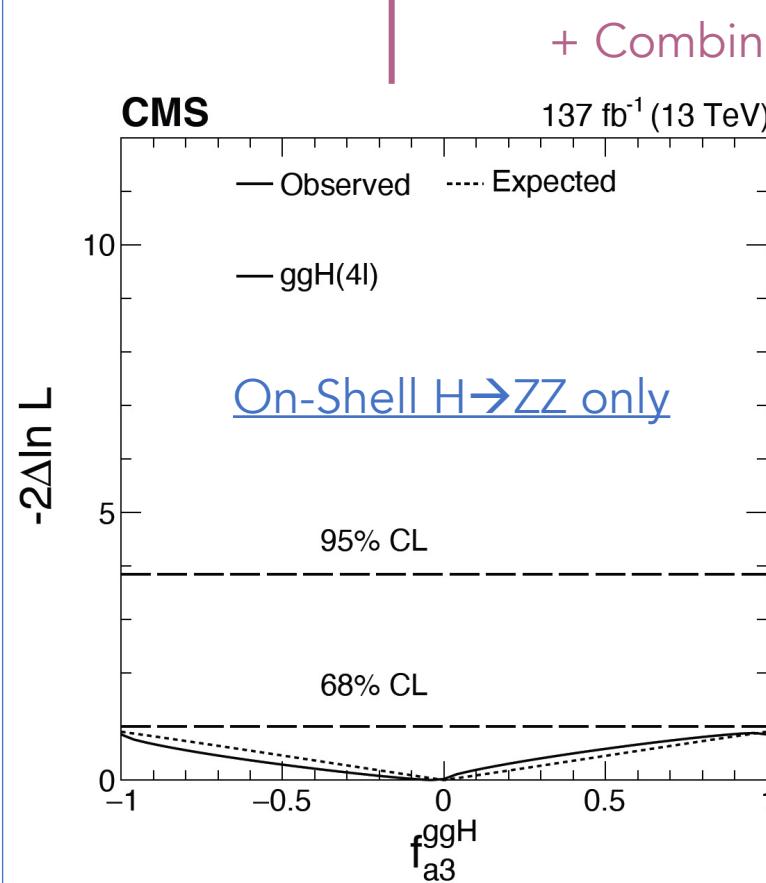
$H \rightarrow \tau\tau$ Results | HVV combining with $H \rightarrow ZZ$ (4ℓ)

HVV vertex



f_{a2} , f_{a1} and $f_{a1}z_\gamma$
also constrained (in backup)

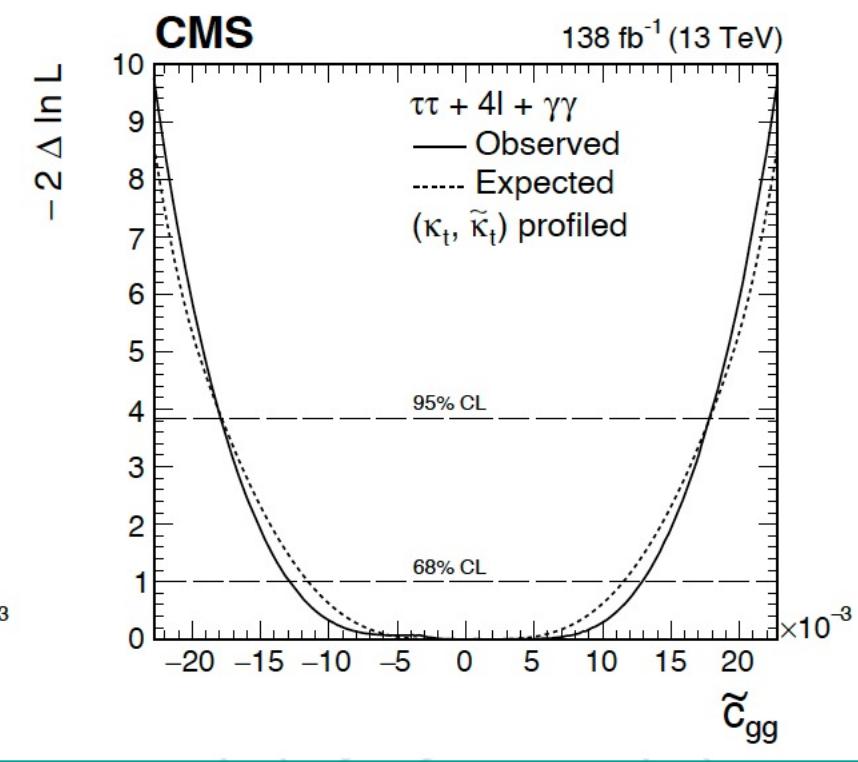
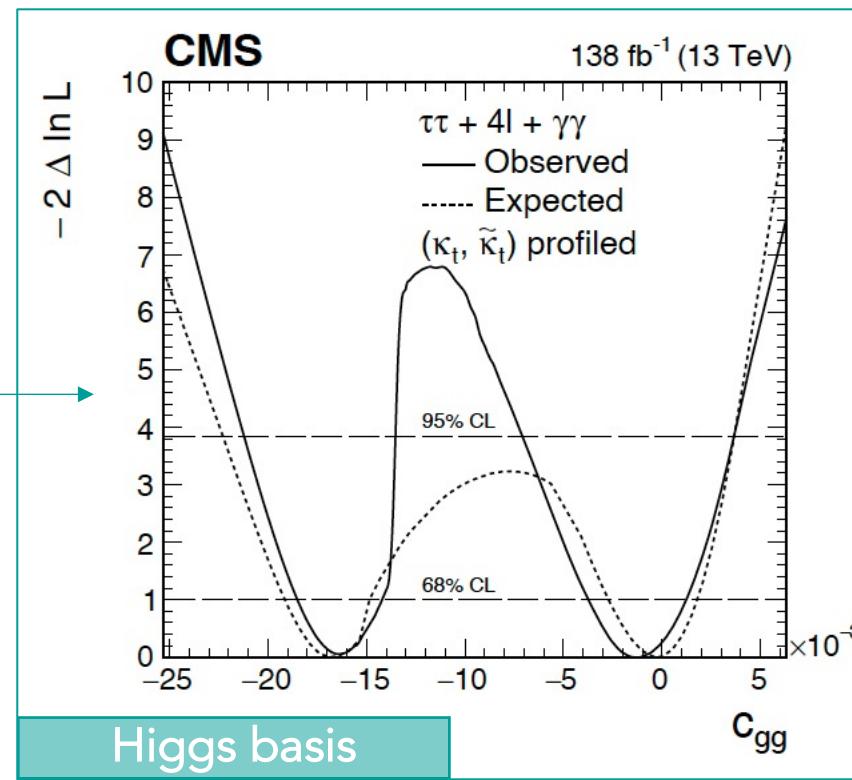
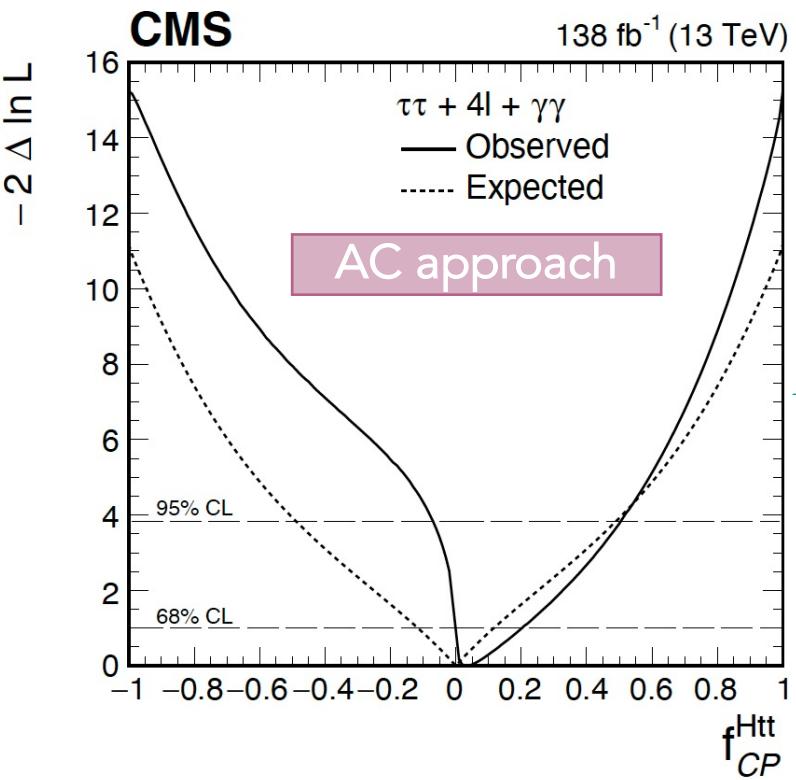
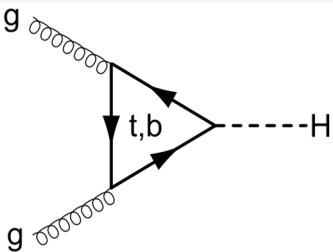
Hgg vertex



The pure CP-odd scenario of the Higgs boson coupling to gluons is excluded at 2.4σ

$H \rightarrow \tau\tau$ Results | Hff combining with $H \rightarrow ZZ$ (4ℓ) + ttH (2γ)

ggH loop is dominated by the top quark



The combination improves the limits on the anomalous coupling parameters typically by about 20–50%

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“Search for CP violation in ttH and tH production in multilepton channels in proton-proton collisions at $\sqrt{s} = 13$ TeV”

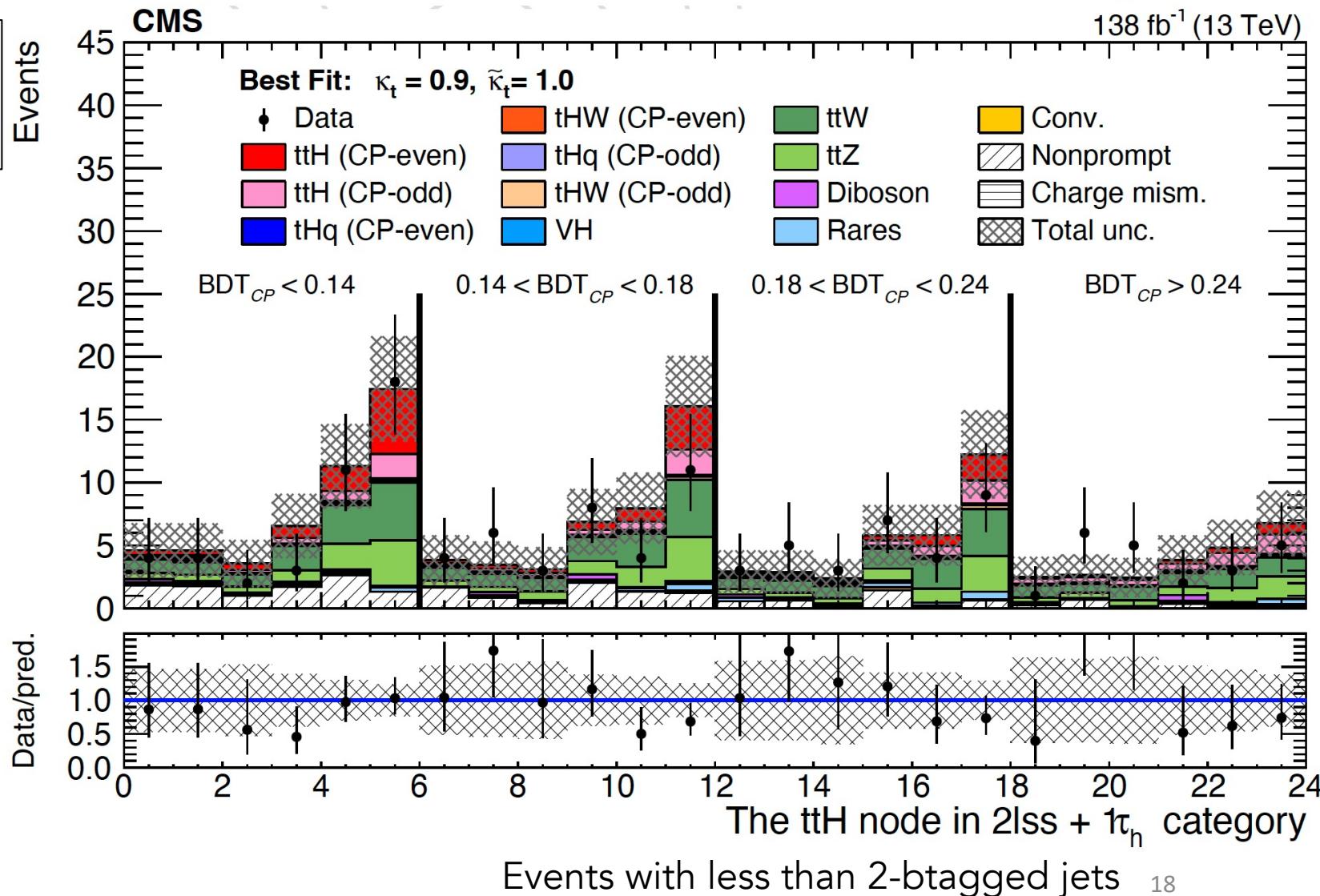
ttH & tH | H_t (Higgs coupling to fermions)

Channels: HWW/H $\tau\tau$

Productions: ttH & tH

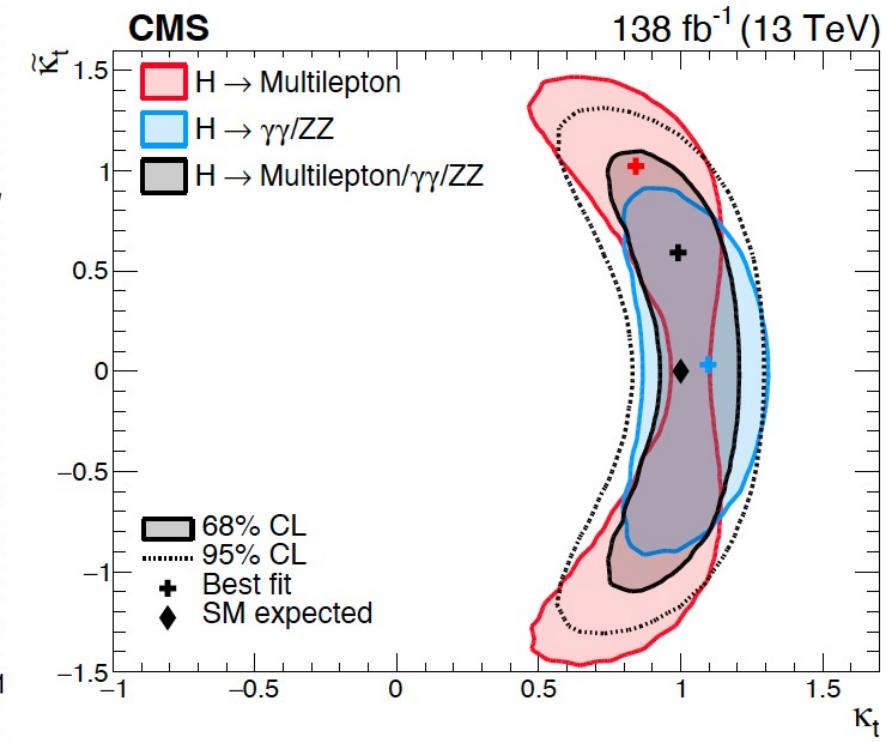
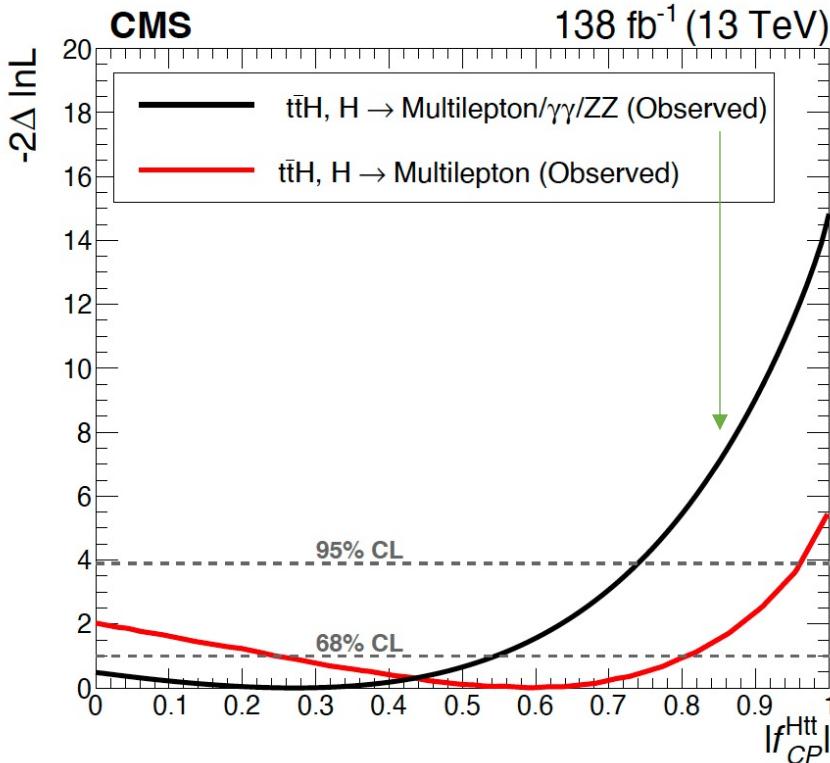
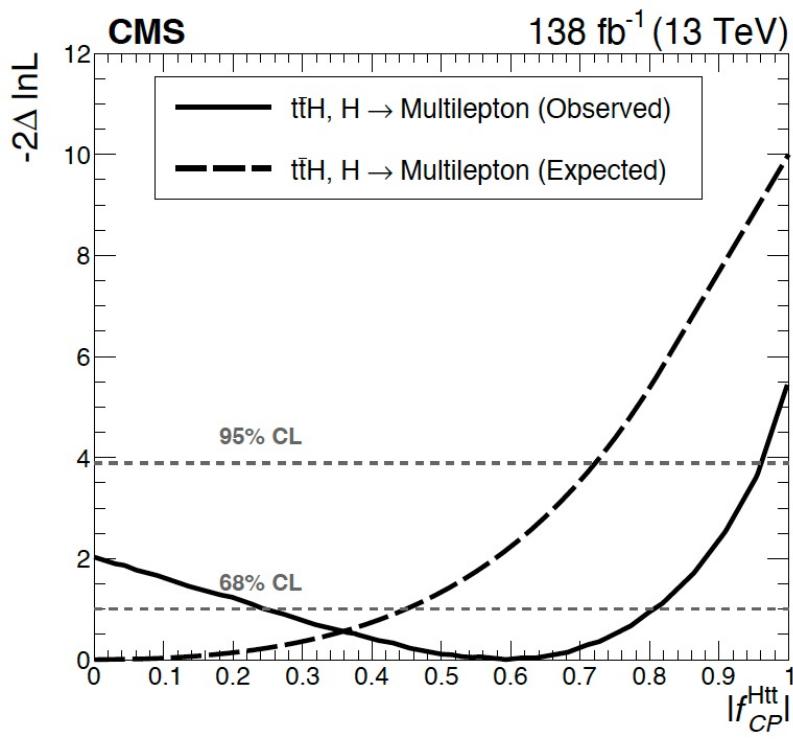
Machine learning techniques:

- BDT $\rightarrow m_{\text{ttH}}, \Delta R_{jj}, \Delta\eta$



ttH & tH Results | H_tt

$|f_{CP}^{Htt}| = 0$ (SM expectation)



$|f_{CP}^{Htt}| = 0.59$ with an interval of (0.24, 0.81) at 68% CL

Combination: $|f_{CP}^{Htt}| = 0.28$ with an interval of $|f_{CP}^{Htt}| < 0.55$ at 68% CL

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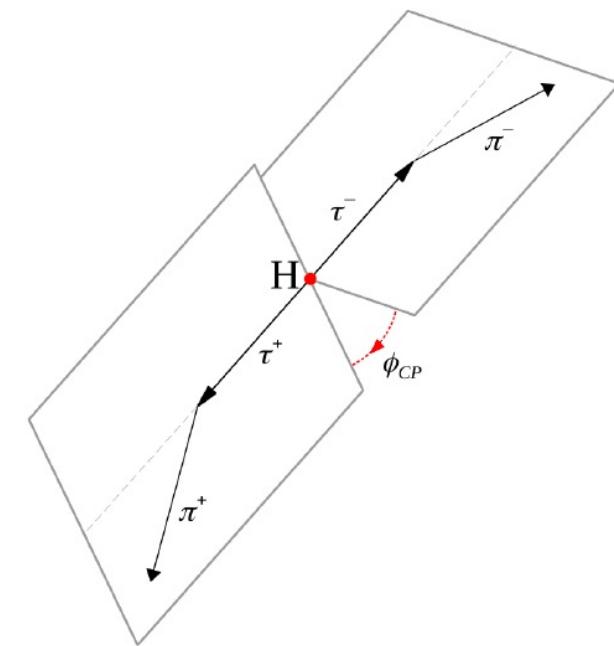
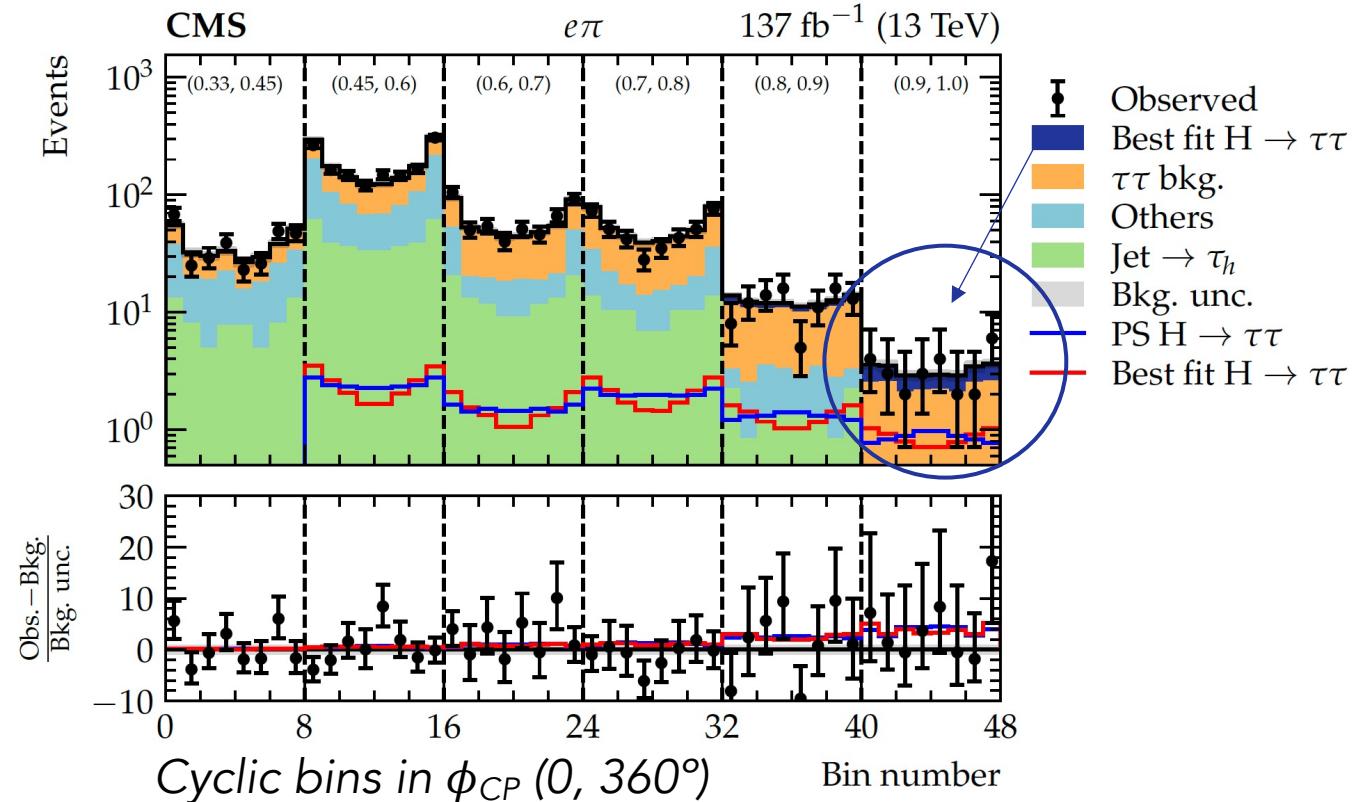
“Analysis of the CP structure of the Yukawa coupling
between the Higgs boson and t leptons in proton-proton
collisions at $\sqrt{s} = 13$ TeV”

Higgs coupling to tau leptons |

Channels: $\tau_h\tau_h$, $\tau_\mu\tau_h$, $\tau_e\tau_h$
Productions: ggH, VBF, VH

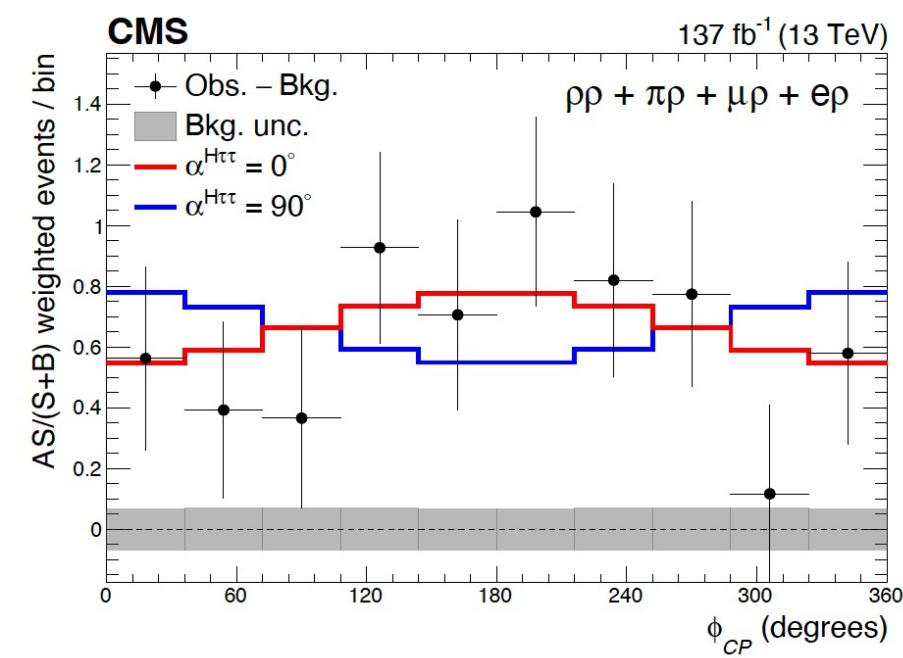
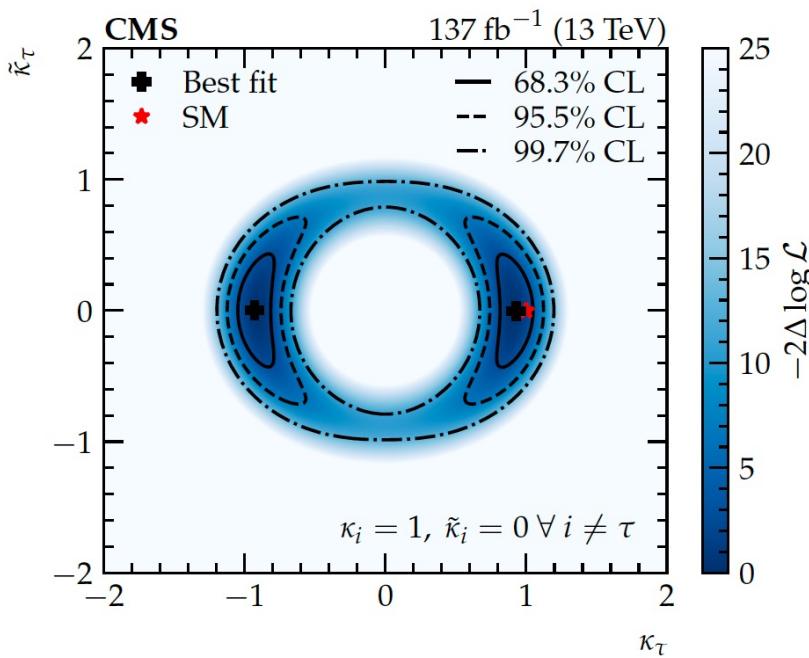
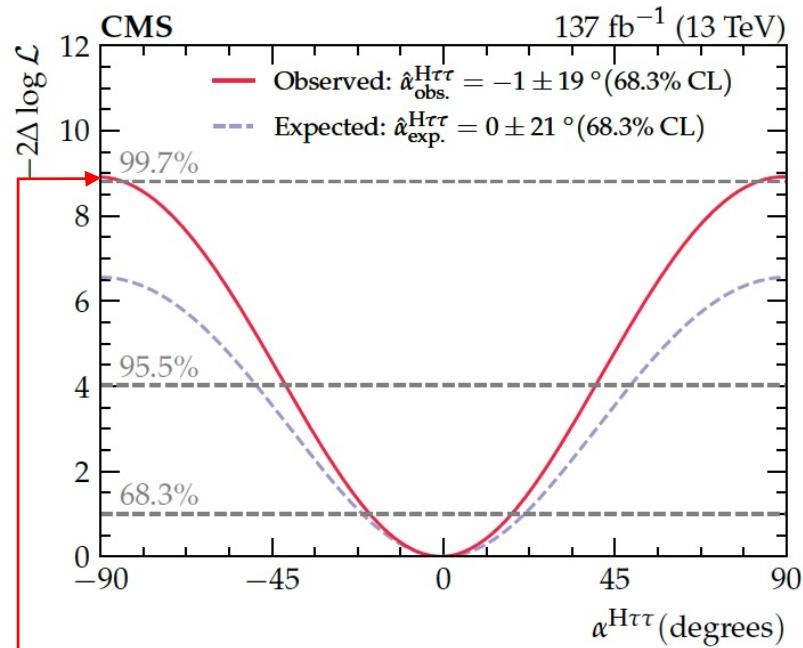
Angular correlation between the decay planes of τ leptons

- MVA scores
- ϕ_{CP} distributions



Higgs coupling to tau leptons |

$$\alpha^{Hff} = \tan^{-1} \left(\frac{\tilde{\kappa}_f}{\kappa_f} \right)$$



The data disfavour the pure CP-odd scenario at 3.0σ

Observed (expected) $\alpha^{\text{H}\tau\tau} = -1 \pm 19$ (0 ± 21) at the 68.3% CL

Conclusions |

- JCP studies in Higgs couplings to both fermions and bosons are important to determine its nature.
- The most stringent limits on CP violation and Higgs AC by the CMS experiment have been presented.

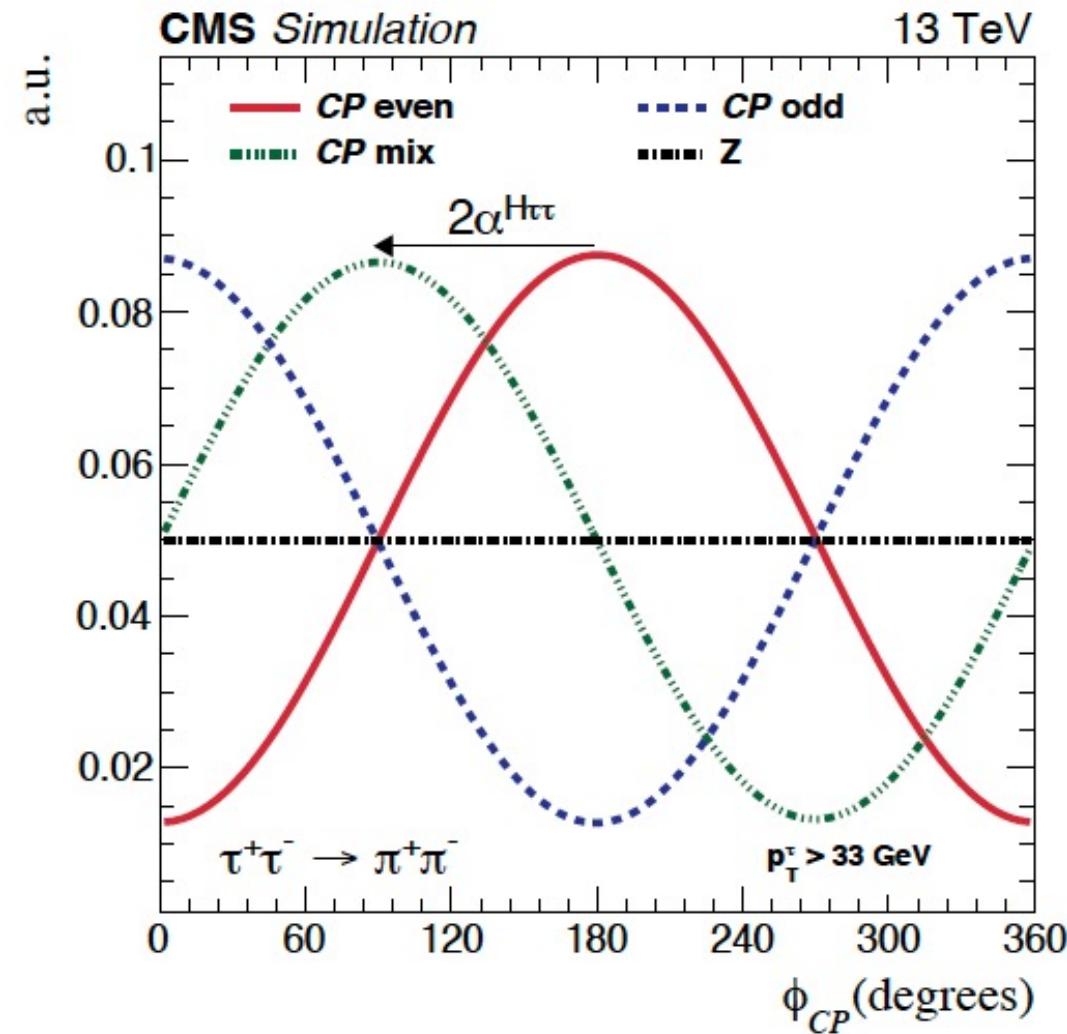
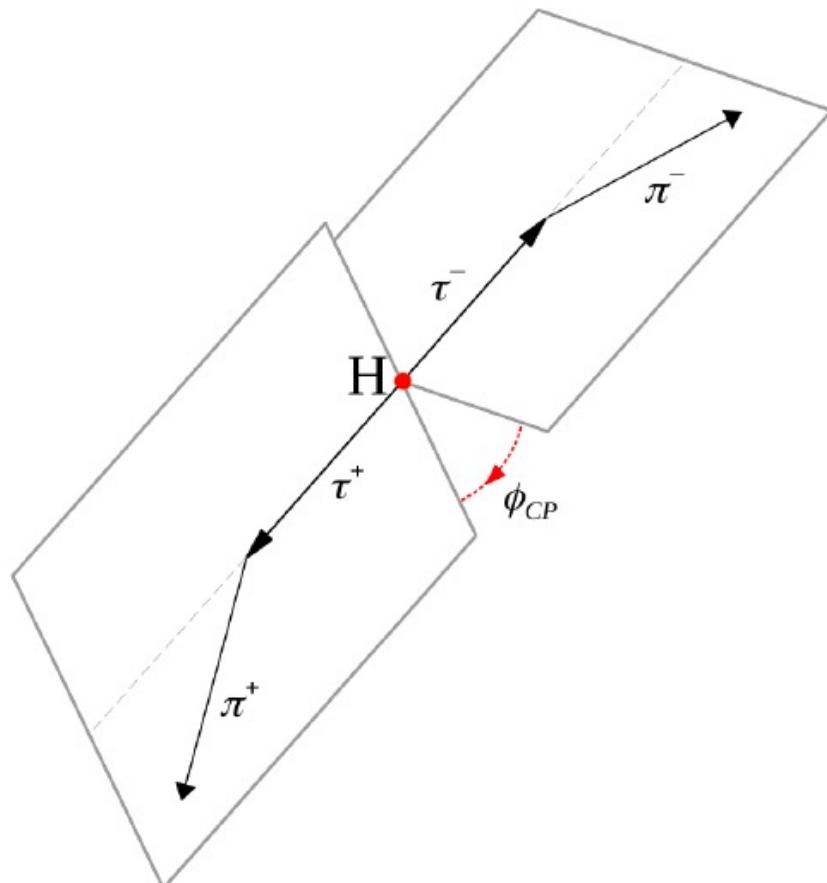
Growing field with many recent updates including new interpretations, and plenty more results to come in the future

- The leading uncertainty is statistical → the precision will increase with the accumulation of more collision data.
 - Additional Run 2 AC analyses are still to be released soon.

Back-up slides

Higgs AC to fermions | Mixing angle

κ_f and $\tilde{\kappa}_f$ are Yukawa coupling strength modifiers related to the mixing angle $\alpha^{Hff} = \tan^{-1} \left(\frac{\tilde{\kappa}_f}{\kappa_f} \right)$



Effective Field Theory |

Direct searches at the LHC find no new particles

Suggests an energy gap between SM and BSM motivating the use of EFT.

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_i \bar{c}_i^{(6)} \mathcal{O}_i^{(6)}$$

- $\mathcal{O}_i^{(5)}$ → violate lepton number
- $\mathcal{O}_i^{(7)}$ → violate lepton or baryon number
- $\mathcal{O}_i^{(8)}$ → Suppressed by $1/\Lambda^4$ → Typically neglected

Particles with $m \ll \Lambda$ → Operators (\mathcal{O}_i) obey SM symmetries.

Effects of **new physics** mapped onto $\mathcal{O}_i^{(6)}$ → contribution scale as $(E/\Lambda)^2$

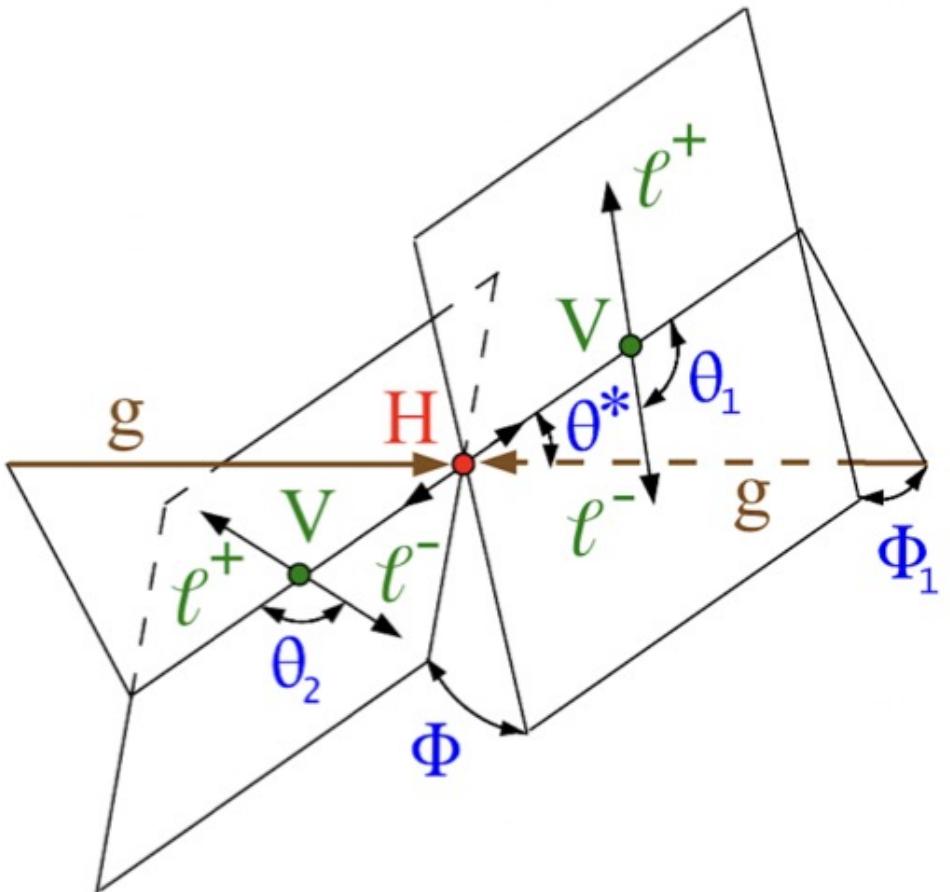
$\bar{c}_i^{(6)}$ (Wilson coefficient) → strength of the new interactions

EFT only valid at
 $E < \Lambda$

Measurement of Higgs couplings |

Dedicated measurements (detector-level)

Dedicated discriminants and full simulation of Signal PDFs.

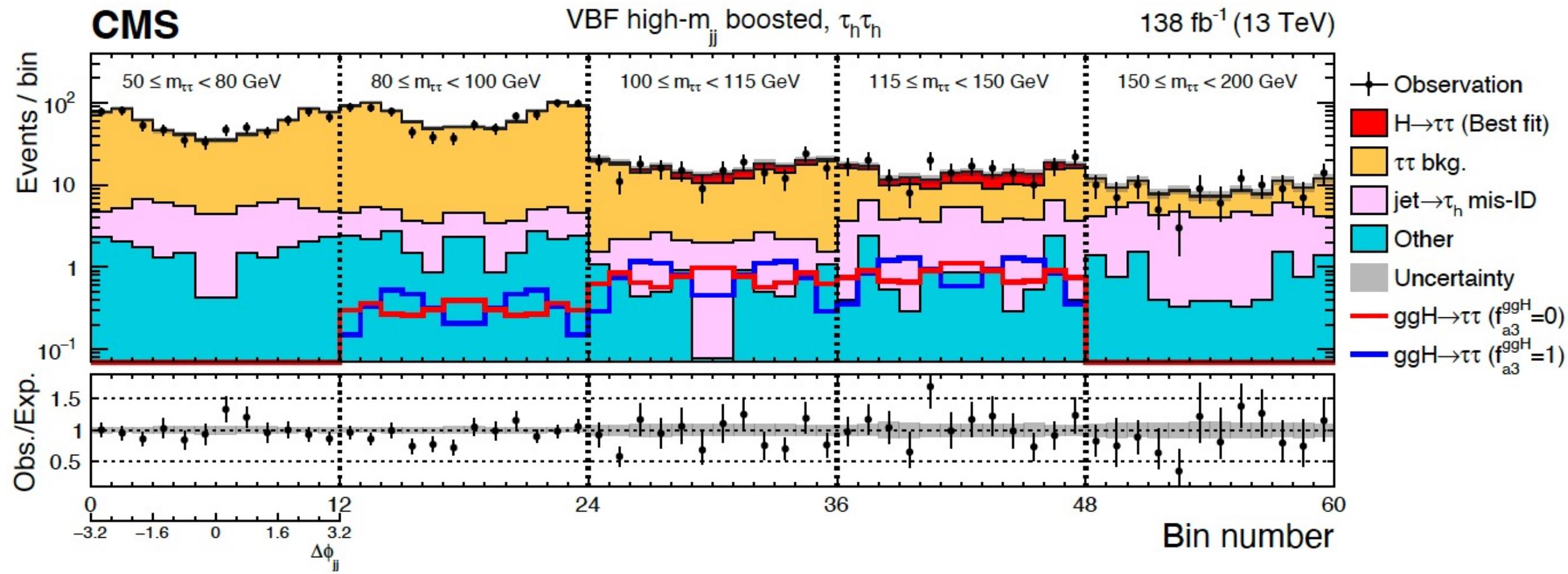


Exploit full production and decay kinematic information
with ME or ML-based discriminants.

Sensitive to higher dimensional operators in the EFT

- ✓ Can target production mode, Higgs coupling and interference.

$H \rightarrow \tau\tau$ analysis | Discriminants



Higgs AC to bosons | Lagrangian

$$A(HV_1V_2) = \frac{1}{v} \left\{ M_{V_1}^2 \left(g_1^{VV} + \frac{\kappa_1^{VV} q_1^2 + \kappa_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} + \frac{\kappa_3^{VV} (q_1 + q_2)^2}{(\Lambda_Q^{VV})^2} + \frac{2q_1 \cdot q_2}{M_{V_1}^2} g_2^{VV} \right) (\epsilon_1 \cdot \epsilon_2) \right. \\ \left. - 2g_2^{VV} (\epsilon_1 \cdot q_2) (\epsilon_2 \cdot q_1) - 2g_4^{VV} \epsilon_{\epsilon_1 \epsilon_2 q_1 q_2} \right\},$$

"The parametrization of the amplitude can be related to a fundamental Lagrange density function using the Higgs basis which is based on an effective field theory expansion up to dimension six. The relevant $SU(3) \times SU(2) \times U(1)$ invariant Lagrangian for H boson interactions with gauge bosons"

$$\mathcal{L}_{hvv} = \frac{h}{v} \left[(1 + \delta c_z) \frac{(g^2 + g'^2)v^2}{4} Z_\mu Z_\mu + c_{zz} \frac{g^2 + g'^2}{4} Z_{\mu\nu} Z_{\mu\nu} + c_{z\square} g^2 Z_\mu \partial_\nu Z_{\mu\nu} + \tilde{c}_{zz} \frac{g^2 + g'^2}{4} Z_{\mu\nu} \tilde{Z}_{\mu\nu} \right. \\ + (1 + \delta c_w) \frac{g^2 v^2}{2} W_\mu^+ W_\mu^- + c_{ww} \frac{g^2}{2} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{w\square} g^2 (W_\mu^- \partial_\nu W_{\mu\nu}^+ + \text{H.c.}) + \tilde{c}_{ww} \frac{g^2}{2} W_{\mu\nu}^+ \tilde{W}_{\mu\nu}^- \\ + c_{z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} A_{\mu\nu} + \tilde{c}_{z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} \tilde{A}_{\mu\nu} + c_{\gamma\square} gg' Z_\mu \partial_\nu A_{\mu\nu} \\ \left. + c_{rr} \frac{e^2}{4} A_{\mu\nu} A_{\mu\nu} + \tilde{c}_{rr} \frac{e^2}{4} A_{\mu\nu} \tilde{A}_{\mu\nu} + c_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a G_{\mu\nu}^a + \tilde{c}_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a \right],$$

Higgs AC to bosons | Lagrangian

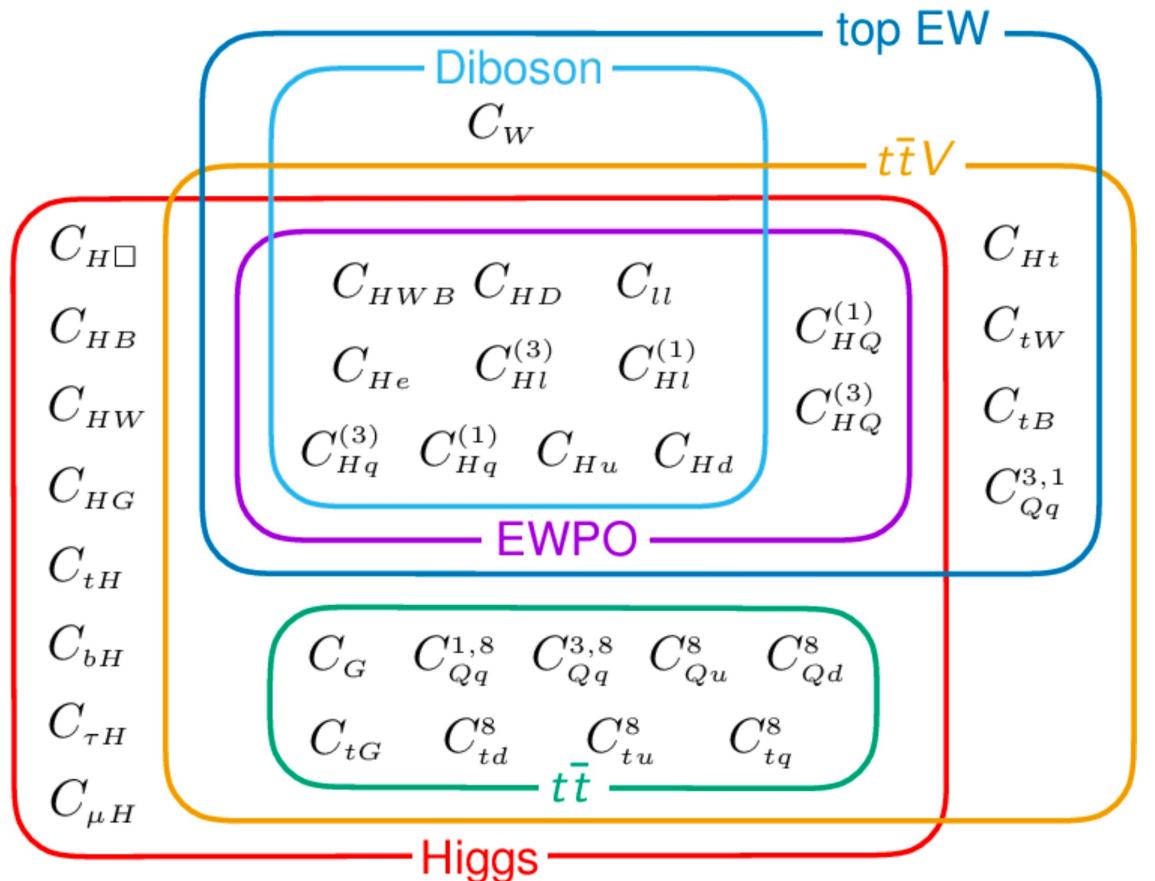
$$A(HV_1V_2) = \frac{1}{v} \left\{ M_{V_1}^2 \left(g_1^{VV} + \frac{\kappa_1^{VV} q_1^2 + \kappa_2^{VV} q_2^2}{(\Lambda_1^{VV})^2} + \frac{\kappa_3^{VV} (q_1 + q_2)^2}{(\Lambda_Q^{VV})^2} + \frac{2q_1 \cdot q_2}{M_{V_1}^2} g_2^{VV} \right) (\epsilon_1 \cdot \epsilon_2) \right. \\ \left. - 2g_2^{VV} (\epsilon_1 \cdot q_2) (\epsilon_2 \cdot q_1) - 2g_4^{VV} \epsilon_{\epsilon_1 \epsilon_2 q_1 q_2} \right\},$$

"The generality of the amplitude parametrization allows to uniquely represent each EFT coefficient in the Lagrangian by an anomalous coupling limiting our couplings to real-valued numbers"

$$\begin{aligned} \delta c_z &= \frac{1}{2} g_1^{ZZ} - 1, & c_{zz} &= -\frac{2s_w^2 c_w^2}{e^2} g_2^{ZZ}, & c_{z\square} &= \frac{M_Z^2 s_w^2}{e^2} \frac{\kappa_1^{ZZ}}{(\Lambda_1^{ZZ})^2}, & \tilde{c}_{zz} &= -\frac{2s_w^2 c_w^2}{e^2} g_4^{ZZ}, \\ \delta c_w &= \frac{1}{2} g_1^{WW} - 1, & c_{ww} &= -\frac{2s_w^2}{e^2} g_2^{WW}, & c_{w\square} &= \frac{M_W^2 s_w^2}{e^2} \frac{\kappa_1^{WW}}{(\Lambda_1^{WW})^2}, & \tilde{c}_{ww} &= -\frac{2s_w^2}{e^2} g_4^{WW}, \\ c_{z\gamma} &= -\frac{2s_w c_w}{e^2} g_2^{Z\gamma}, & \tilde{c}_{z\gamma} &= -\frac{2s_w c_w}{e^2} g_4^{Z\gamma}, & c_{\gamma\square} &= \frac{s_w c_w}{e^2} \frac{M_Z^2}{(\Lambda_1^{Z\gamma})^2} \kappa_2^{Z\gamma}, \\ c_{\gamma\gamma} &= -\frac{2}{e^2} g_2^{\gamma\gamma}, & \tilde{c}_{\gamma\gamma} &= -\frac{2}{e^2} g_4^{\gamma\gamma}, & c_{gg} &= -\frac{2}{g_s^2} g_2^{gg}, & \tilde{c}_{gg} &= -\frac{2}{g_s^2} g_4^{gg}. \end{aligned}$$

Basis of dimension 6 operators |

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_i \bar{c}_i^{(6)} \mathcal{O}_i^{(6)}$$



Warsaw basis

Popular SMEFT basis in the theoretical community

Higgs basis

Experimentally convenient to express EFT in terms of mass eigenstates after EWSB

Related to Warsaw basis through linear transformations

H \rightarrow ZZ Results | HVV (SMEFT approach - Warsaw basis)

HVV parametrization can be related to the EFT couplings using the Warsaw basis.

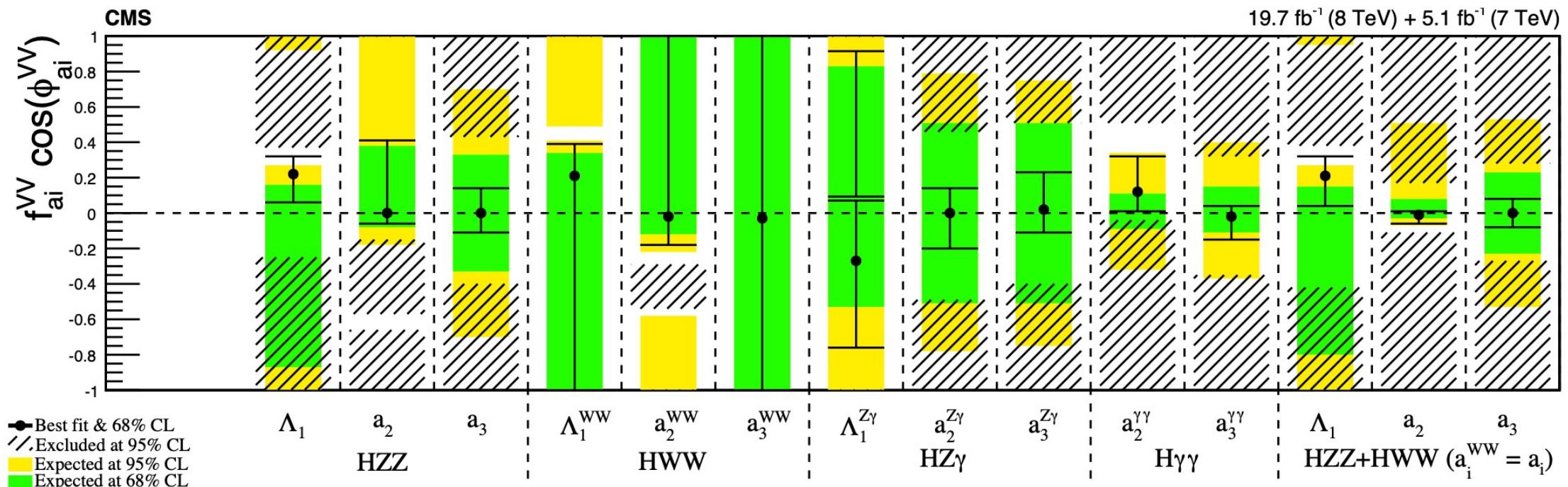
TABLE IX. Summary of constraints on the HVV coupling parameters in the Warsaw basis of SMEFT. For each coupling constraint reported, three other independent operators are left unconstrained, where only one of the three operators c_{HW} , $c_{H\tilde{W}}$, and $c_{H\tilde{B}}$ is independent, and only one of $c_{H\tilde{W}}$, $c_{H\tilde{W}B}$, and $c_{H\tilde{B}}$ is independent.

Channels	Coupling	Observed	Expected
VBF & VH & $H \rightarrow 4\ell$	$c_{H\square}$	$0.04^{+0.43}_{-0.45}$	$0.00^{+0.75}_{-0.93}$
	c_{HD}	$-0.73^{+0.97}_{-4.21}$	$0.00^{+1.06}_{-4.60}$
	c_{HW}	$0.01^{+0.18}_{-0.17}$	$0.00^{+0.39}_{-0.28}$
	$c_{H\tilde{W}B}$	$0.01^{+0.20}_{-0.18}$	$0.00^{+0.42}_{-0.31}$
	c_{HB}	$0.00^{+0.05}_{-0.05}$	$0.00^{+0.03}_{-0.08}$
	$c_{H\tilde{W}}$	$-0.23^{+0.51}_{-0.52}$	$0.00^{+1.11}_{-1.11}$
	$c_{H\tilde{W}B}$	$-0.25^{+0.56}_{-0.57}$	$0.00^{+1.21}_{-1.21}$
	$c_{H\tilde{B}}$	$-0.06^{+0.15}_{-0.16}$	$0.00^{+0.33}_{-0.33}$

CMS analysis from Run 1

LHC Run 1 → Couplings to EW bosons

"Constraints on the spin-parity and anomalous HVV couplings of the Higgs boson in proton collisions at 7 and 8 TeV"



$$f_{ai} = \frac{|a_i|^2 \sigma_i}{\sum_{j=1,2,3...} |a_j|^2 \sigma_j} \text{ sign} \left(\frac{a_i}{a_1} \right)$$