

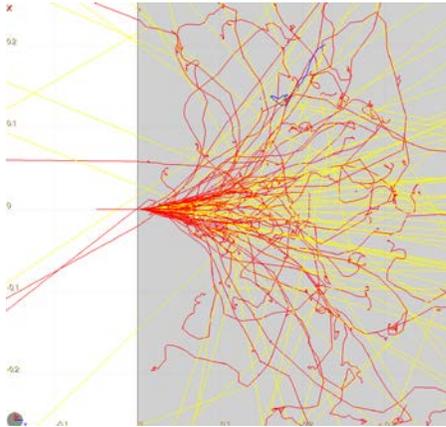
An introduction to the Monte Carlo simulation of radiation transport

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CERN Academic Training
Sep 12, 2022



Monte Carlo simulation tools for radiation transport problems in particle accelerators



Introduction
Mon Sep 12



Tue Sep 13



+ LIVE DEMO!

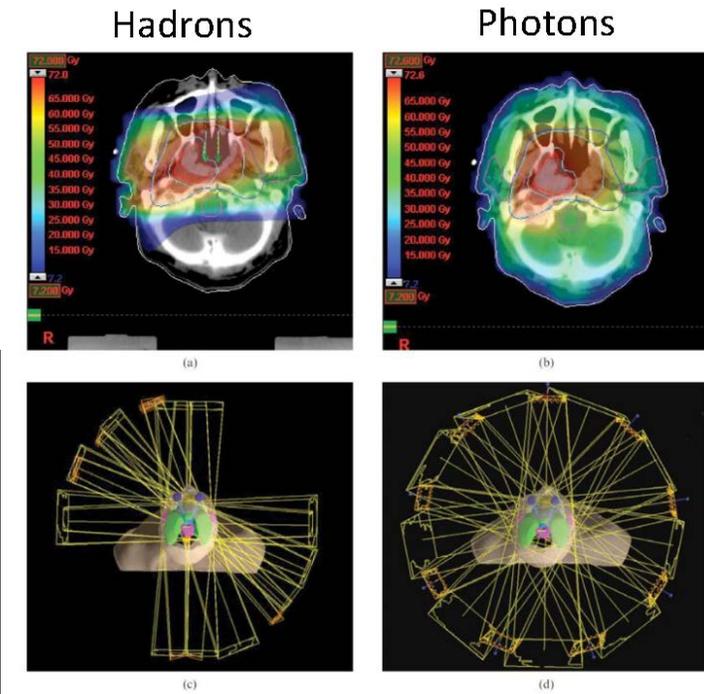
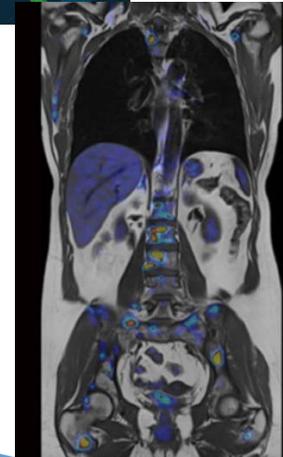
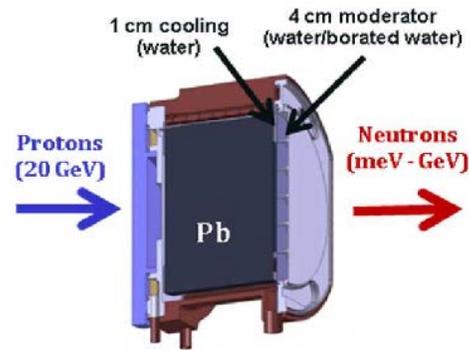
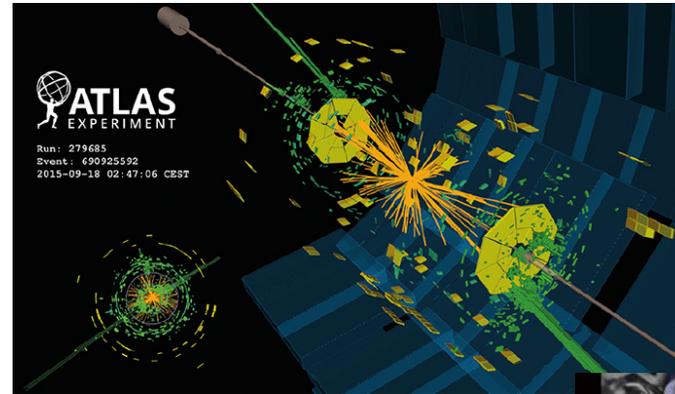
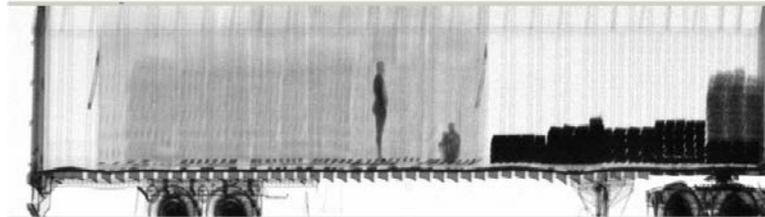
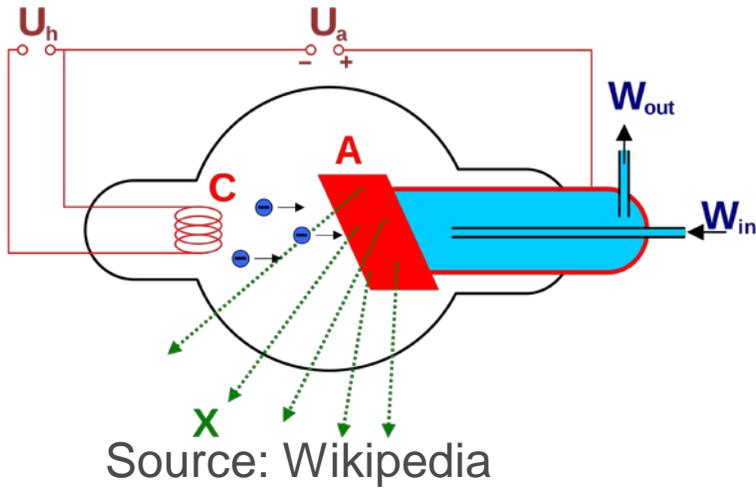
Wed Sep 14

- Other codes exist, e.g. MCNP, PHITS, MARS, PENELOPE,...
- Aim is not to be exhaustive, but to show general capabilities



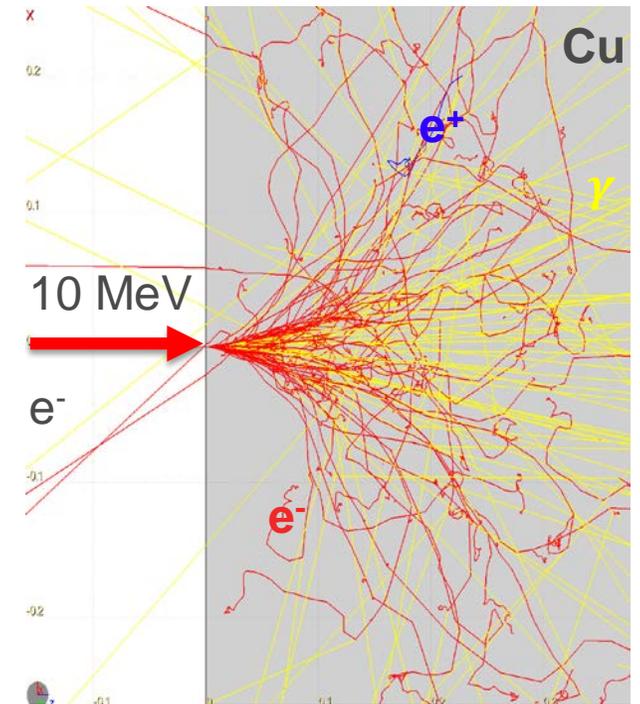
An introduction to the radiation transport problem

Radiation transport



The physical picture

- **Radiation**: ensemble of particles ($\gamma, e^{\pm}, \mu^{\pm}, \nu_i, n, p, \pi^{\pm}, \text{ions}, \dots$) propagating in matter, subject to a series of radiation-matter interaction mechanisms, leading to:
 - **Angular deflections** (target atom displacement)
 - **Energy losses** (energy transfer to material, heating up)
 - Generation of **secondary particles** (formation of a **radiation shower**)
 - Production of **residual nuclei**, possibly radioactive (emission of **prompt/delayed radiation**)

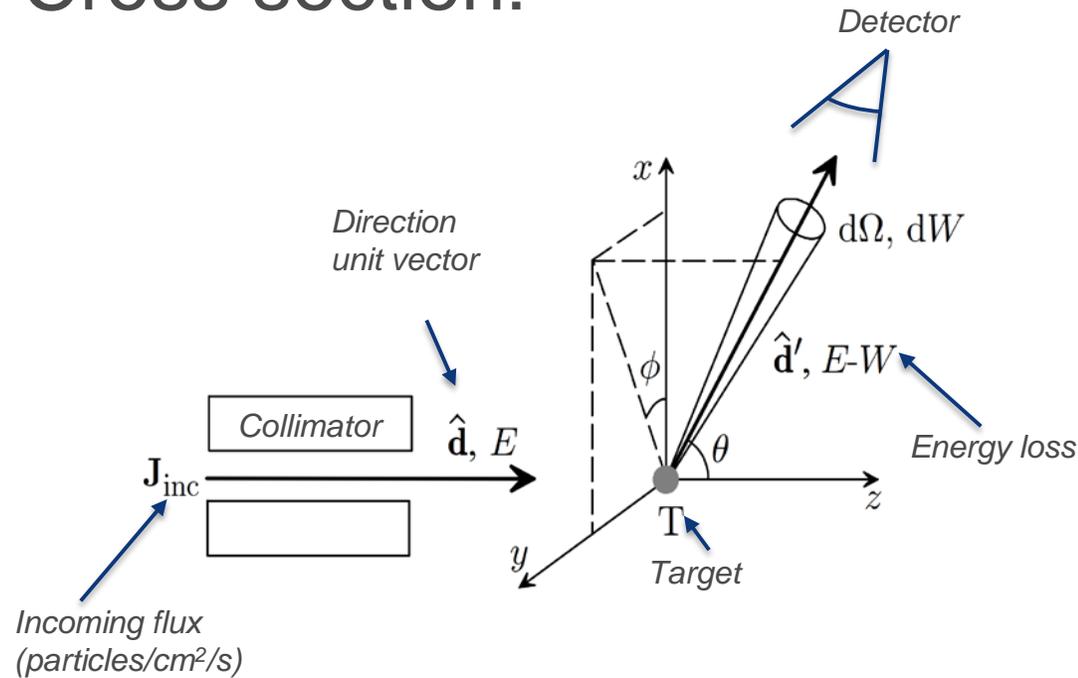


Example: Shower set up by 10-MeV e^- in Cu

Basic quantities

species position energy direction time

- **Particle density** $n_i(\mathbf{r}, E, \Omega, t)$: number of particles of species i per unit volume, unit energy, unit solid angle, at a given time.
- **Cross section:**



Differential cross section

$$\frac{d^2\sigma}{d\Omega dW} \equiv \frac{\dot{N}_{\text{count}}}{|\mathbf{J}_{\text{inc}}| d\Omega dW}$$

Dimensions: Area / Energy / Solid angle

~likelihood of being scattered into a direction Ω with an energy loss W

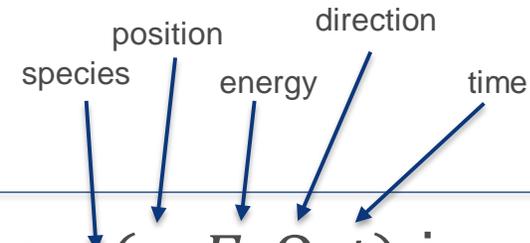
Cross section

$$\sigma \equiv \int \int \frac{d^2\sigma}{d\Omega dW} d\Omega dW$$

Dimensions: Area
Typical unit: 1 barn (=10⁻²⁴ cm²)

~likelihood of being scattered

The transport equation



- Determine time evolution of particle density $n_i(\mathbf{r}, E, \Omega, t)$ in a small volume V
- Transport equation

$$\int_V d\mathbf{r} \frac{\partial n_i(\mathbf{r}, E, \Omega, t)}{\partial t} = - \oint_S dA \mathbf{j}(\mathbf{r}, E, \Omega, t) \cdot \hat{\mathbf{a}}$$

Number of target atoms per unit volume

$$- N \int_V d\mathbf{r} n_i(\mathbf{r}, E, \Omega, t) v(E) \sigma(E)$$

$$+ N \int_V d\mathbf{r} \int dE' \int d\Omega' n_i(\mathbf{r}, E', \Omega', t) v(E') \frac{d\sigma}{d\Omega' dW''}$$

$$+ N \int_V d\mathbf{r} \int dE' \int d\Omega' \sum_j n_j(\mathbf{r}, E', \Omega', t) v(E') \frac{d\sigma_{\text{sec},i}}{d\Omega' dW''}$$

$$+ \int_V d\mathbf{r} Q_{\text{source}}(\mathbf{r}, E, \Omega, t)$$

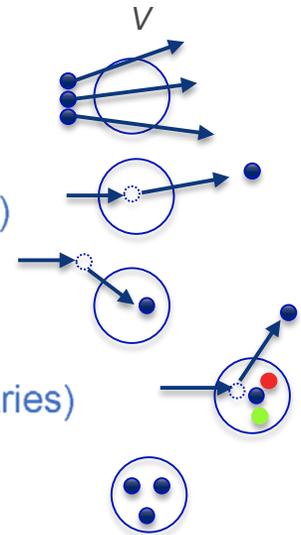
(unscattered particles)

(particles scattered out)

(particles scattered in)

(production of secondaries)

(source)



▪ Integro-differential equation (!)

- Analytical/closed solutions only for simplified scenarios (infinite medium, one/few interaction mechanisms, one/few particle species, etc.)



Buffon's evaluation
of π using
random numbers
(dating back to the 1700s)

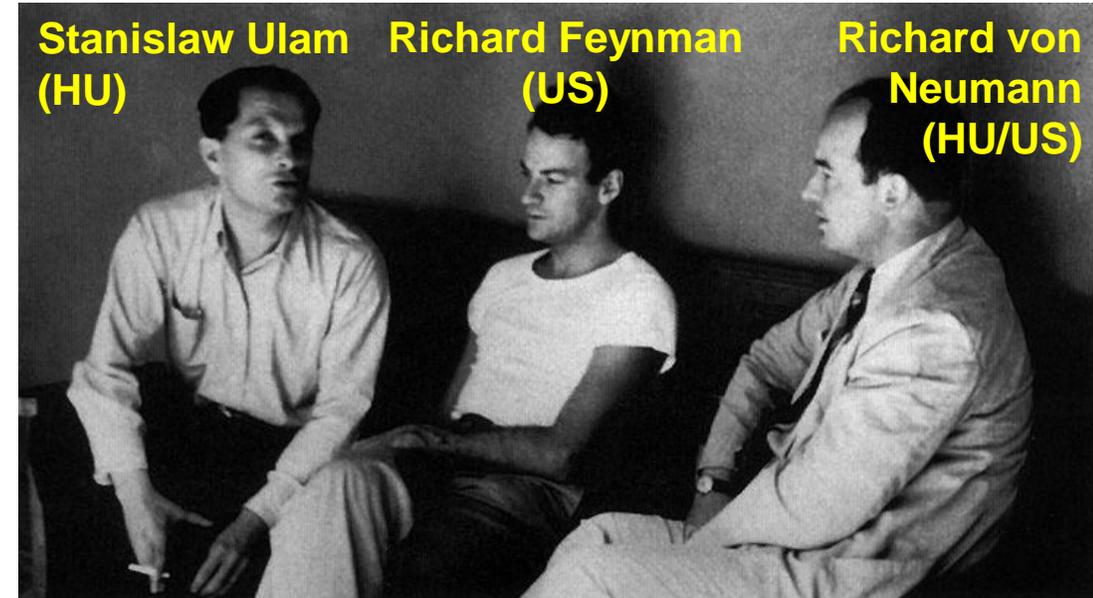
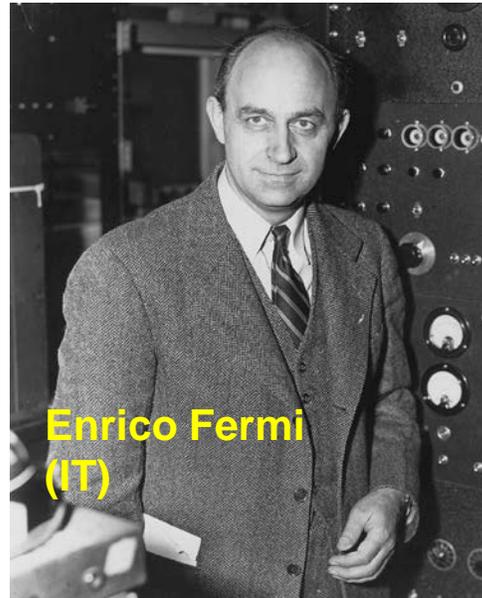


The Monte Carlo approach

Monte Carlo (MC) method for radiation transport – The early days (1930s and 1940s)

- Early ideas by Fermi took off during the Manhattan project effort:

- Nicholas Metropolis
- Stanislaw Ulam
- Richard von Neumann
- Richard Feynman
-



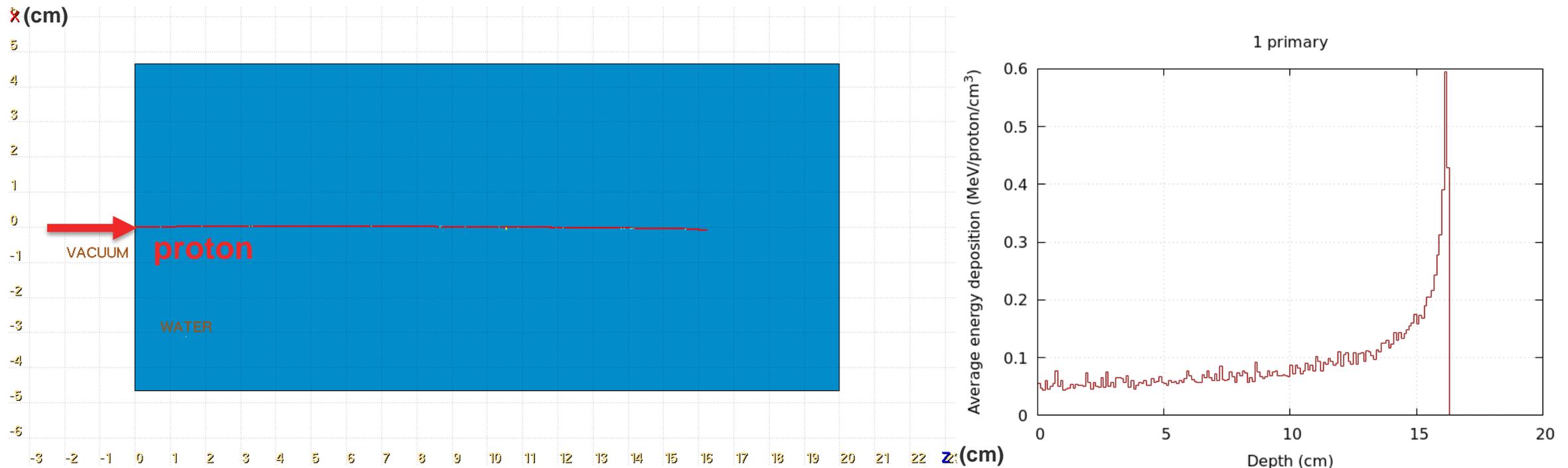
- Parallel to the development of the ENIAC
- *Ref: N. Metropolis “The beginning of the Monte Carlo method” (1987)*

<https://library.lanl.gov/cgi-bin/getfile?00326866.pdf>

See also: <https://ieeexplore.ieee.org/document/6880250>

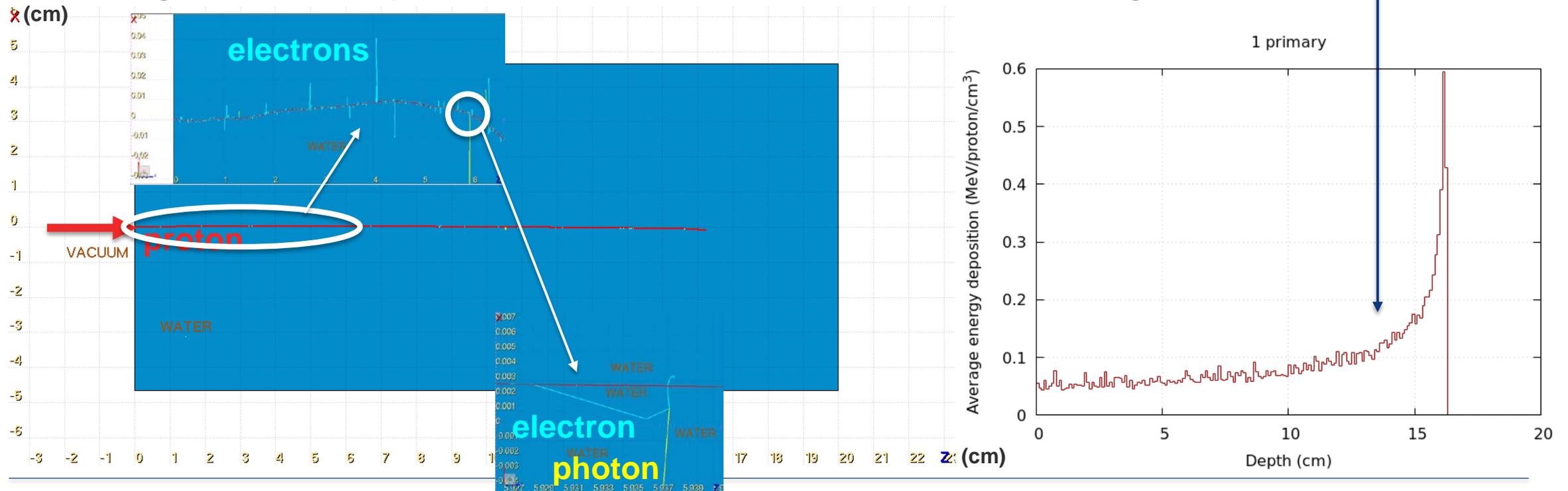
Monte Carlo (MC) approach to radiation transport problems

- **Numerically simulate** a large number of random **particle showers** with given interaction mechanisms
- Accumulate contributions of particle tracks to **statistical estimator** of observable
- Example: 150 MeV protons in water, *energy deposition histogram*



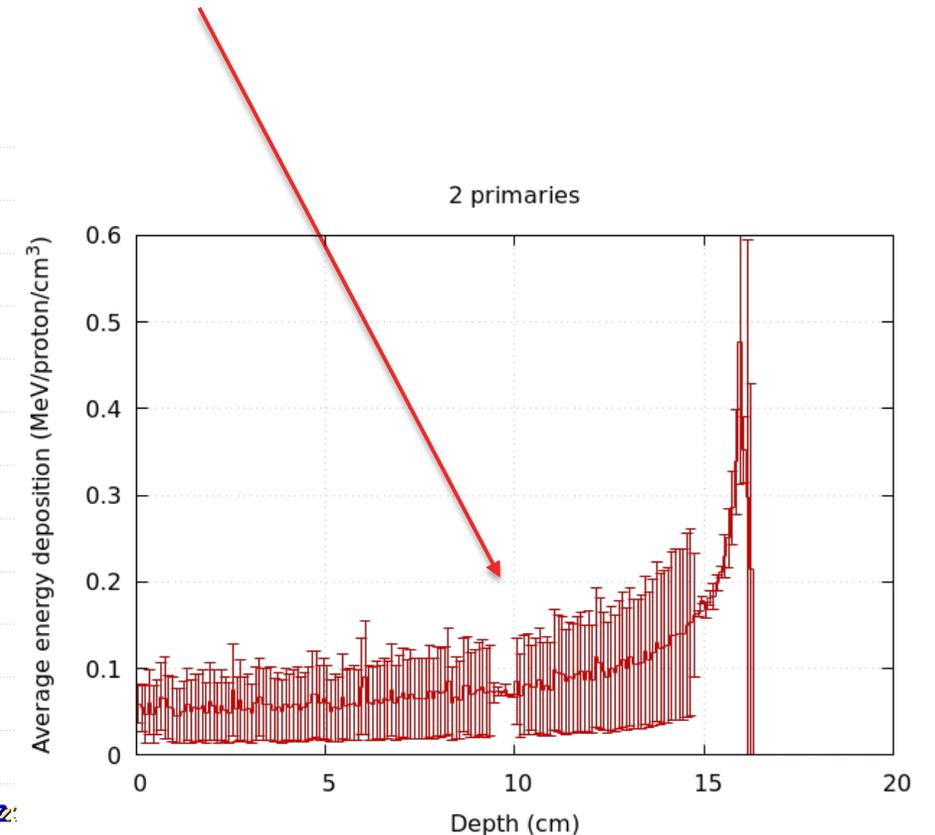
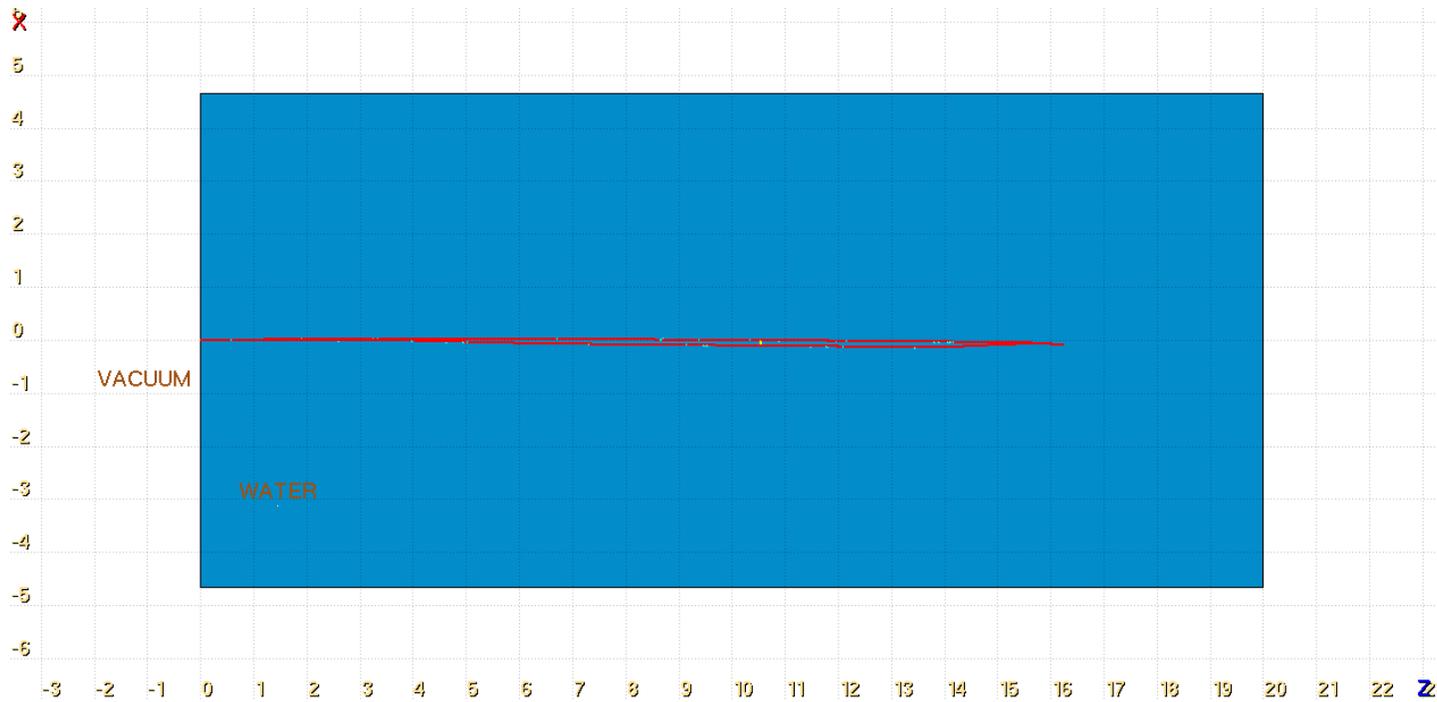
N=1 primary 150-MeV proton in H₂O

- Protons undergo frequent collisions with target e⁻, leading to energy transfer and deposition throughout the depth of the target
- Energy deposition recorded in histogram
- Energetic e⁻ may radiate so-called Bremsstrahlung photons



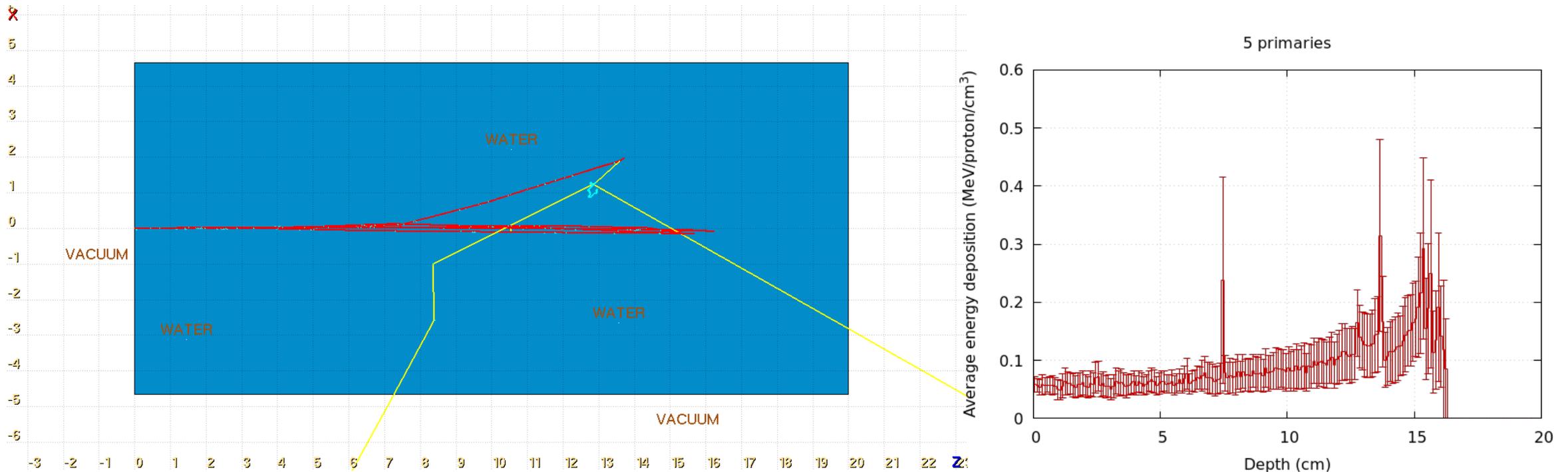
N=2 primary 150-MeV protons in H₂O

- MC simulation results are affected by **statistical uncertainty**:
- All MC results are accompanied by an error bar



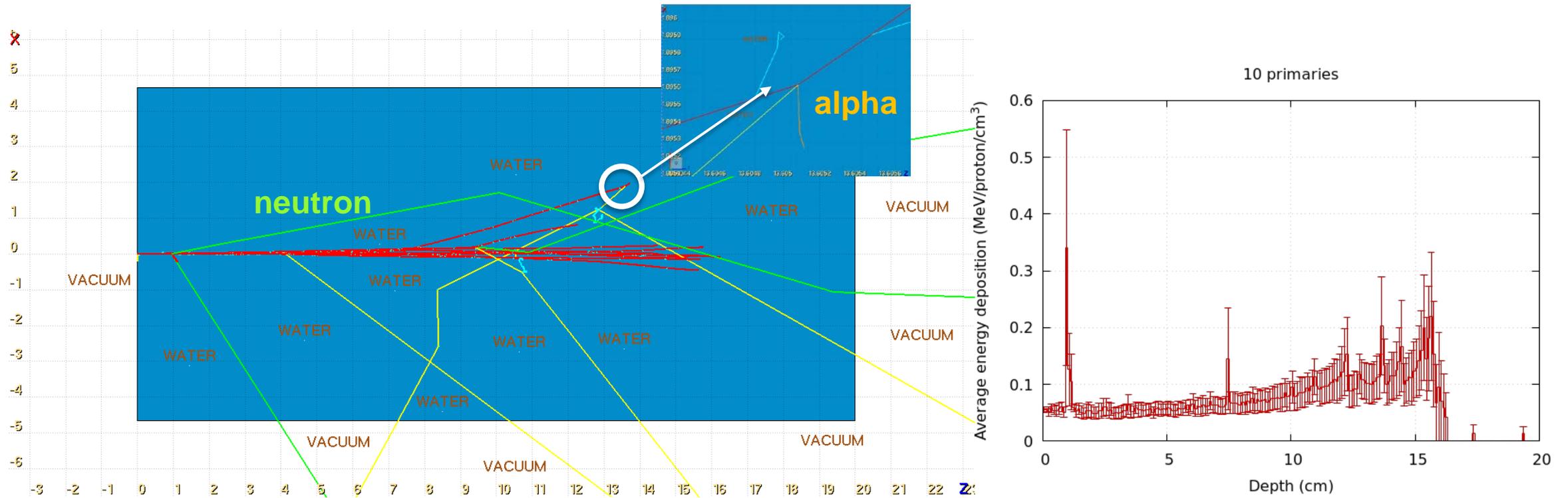
N=5 primary 150-MeV protons in H₂O

- Statistical uncertainty decreases with N
- Still considerable (~80%)



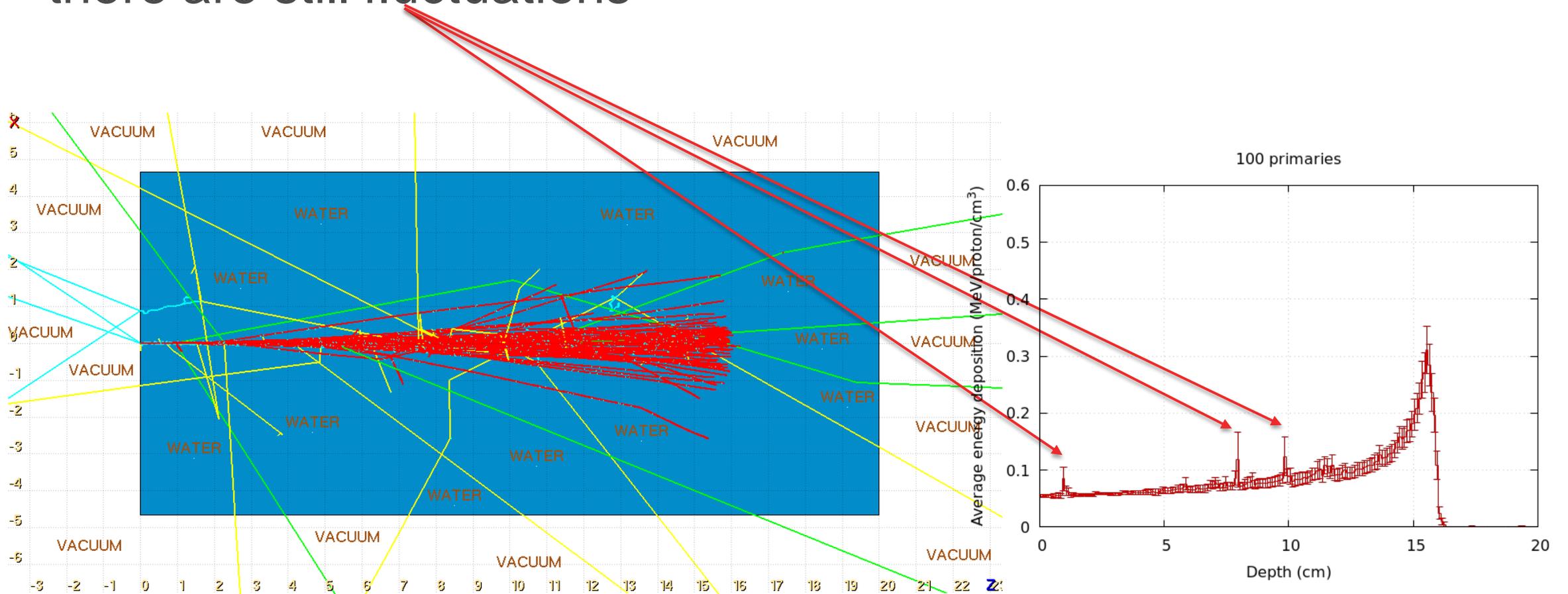
N=10 primary 150-MeV protons in H₂O

- Occasionally, p undergo nuclear inelastic interactions, producing $\gamma, n, p, d, t, {}^3\text{He}, \alpha, \dots$, + (possibly unstable) residual nucleus



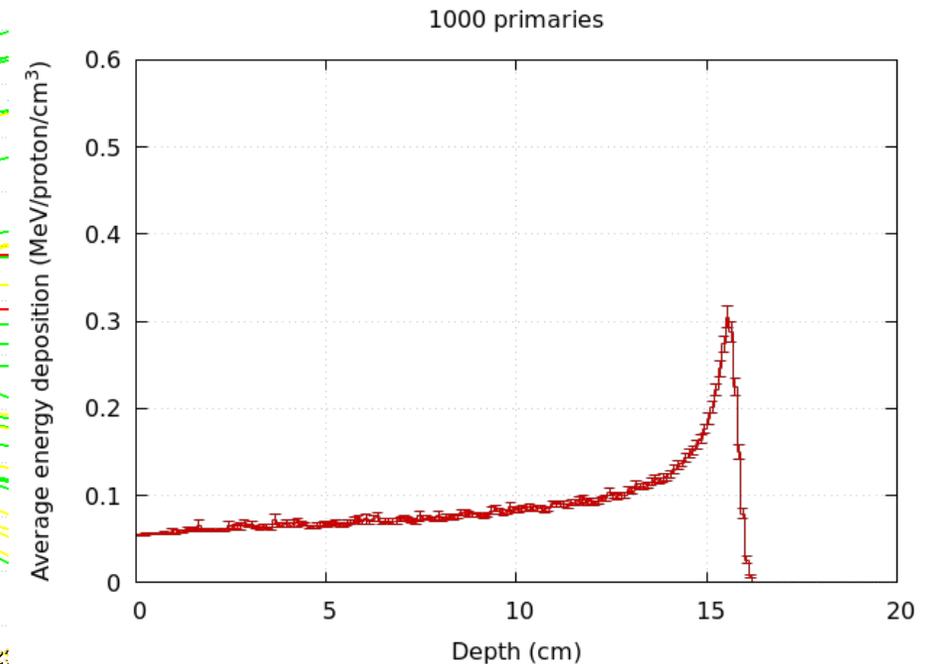
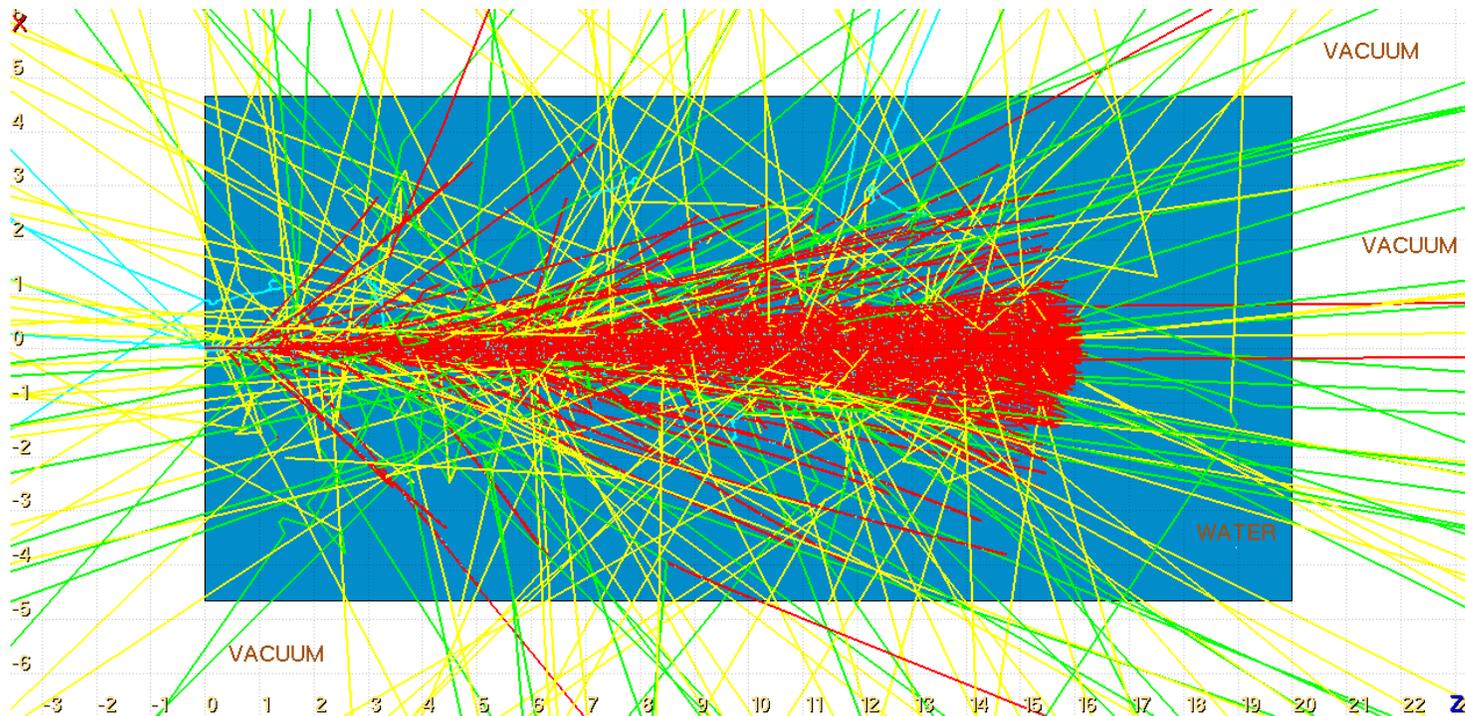
N=100 primary 150-MeV protons in H₂O

- Statistical uncertainty is gradually decreasing (~20-25%), but there are still fluctuations



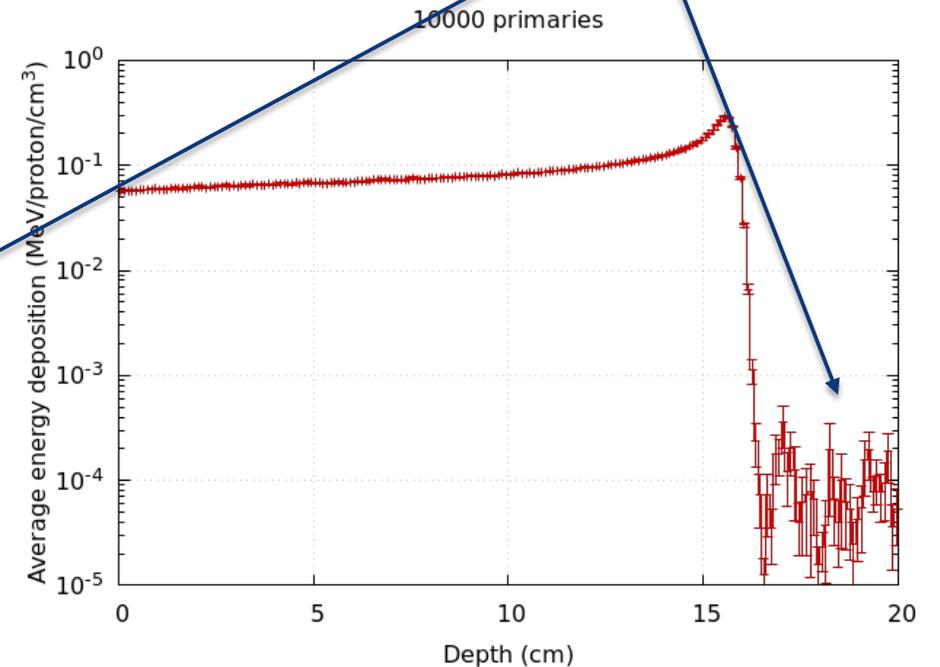
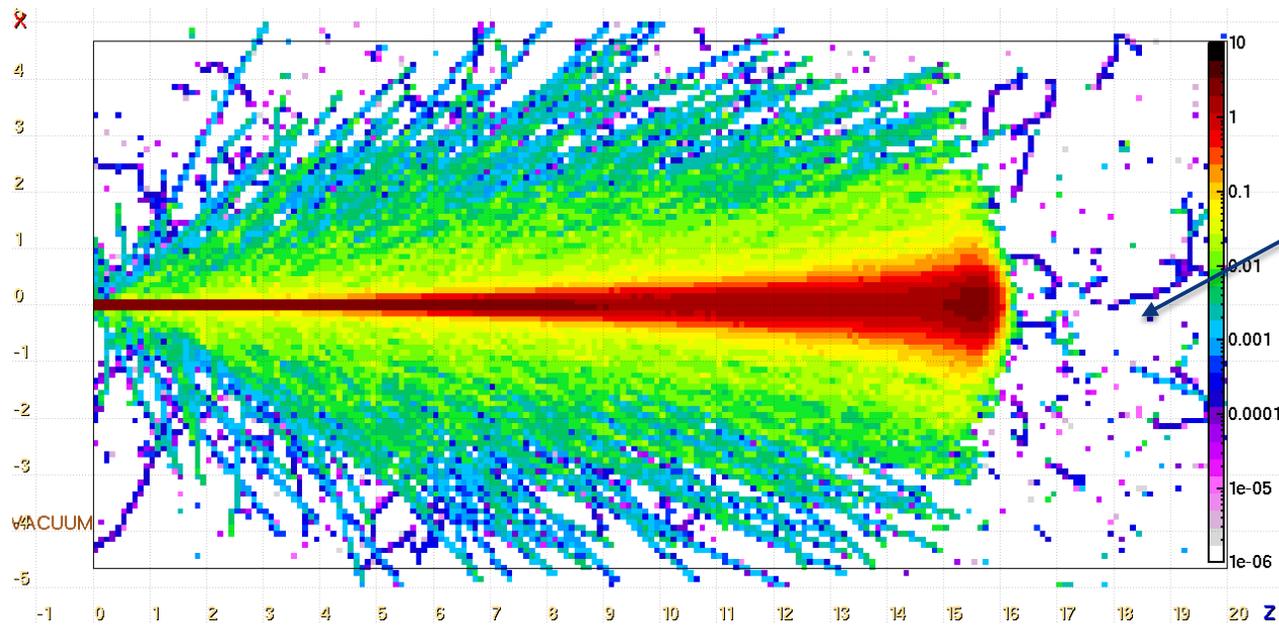
N=1000 primary 150-MeV protons in H₂O

- Convergence is typically approached when statistical uncertainty is below ~10%



N=10000 primary 150-MeV protons in H₂O

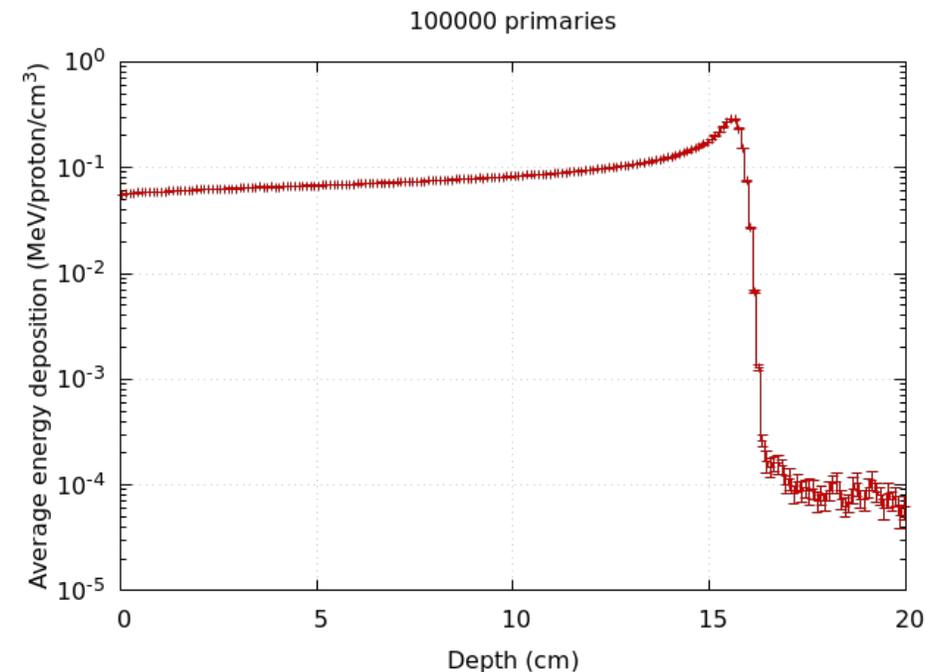
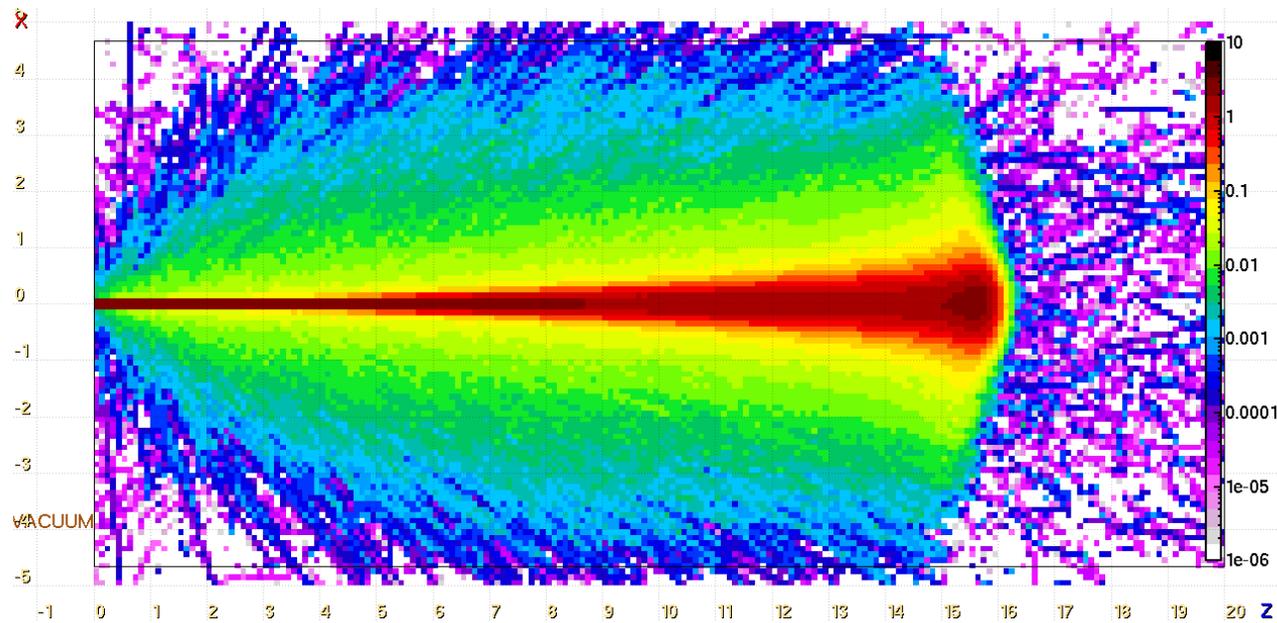
- One does not look at tracks, but rather 2D/1D distributions
- Switching now to log scale to see details
- Neutrons and photons contribute to tails of distribution



N=100000 primary 150-MeV p in H₂O

- Convergence of suppressed features requires further primaries
- There exists a series of tricks to address slow convergence (biasing / variance reduction techniques)

See: <https://www.frontiersin.org/articles/10.3389/fphy.2021.718873/full>

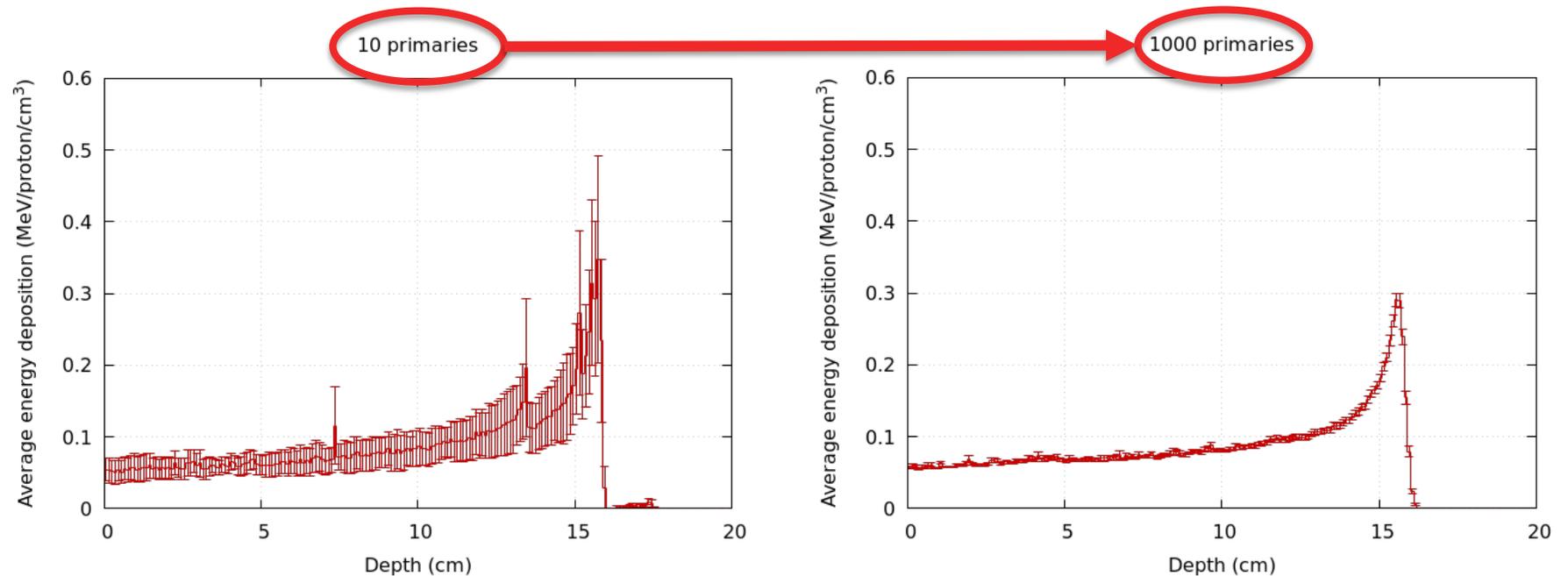


Statistical uncertainty

- Statistical uncertainty in a MC simulation σ drops like

$$\sigma = \frac{1}{\sqrt{N}}$$

i.e. to lower error bars by a factor 10 one needs 100 times more primaries:

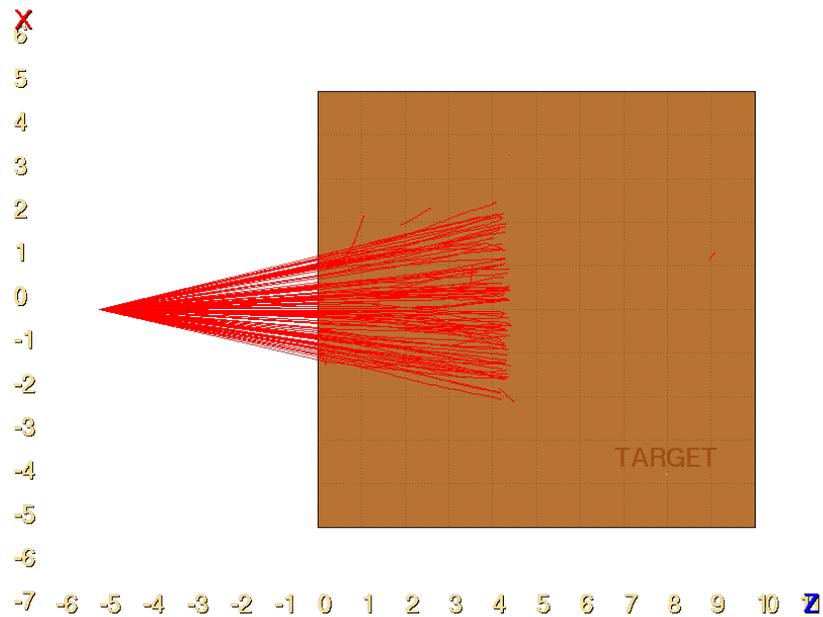




Basic ingredients for a MC code: 1 – Radiation sources

Radiation sources

- Particle beam:
 - Angular divergence
 - Energy spectrum...



- Radioactive isotopes

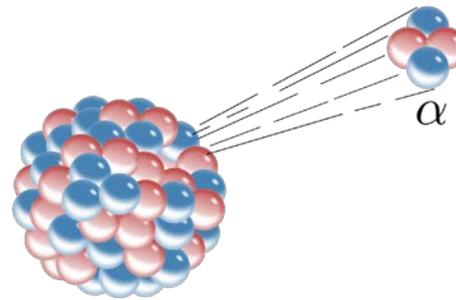
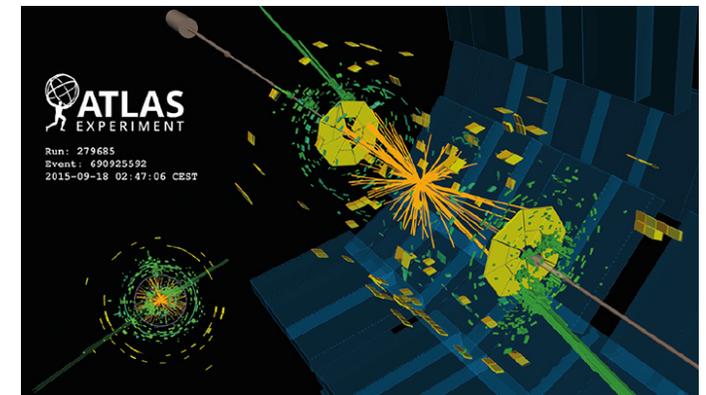
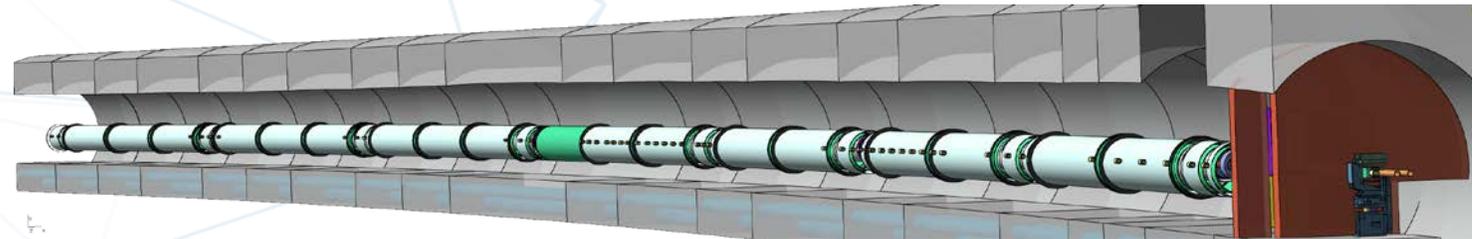
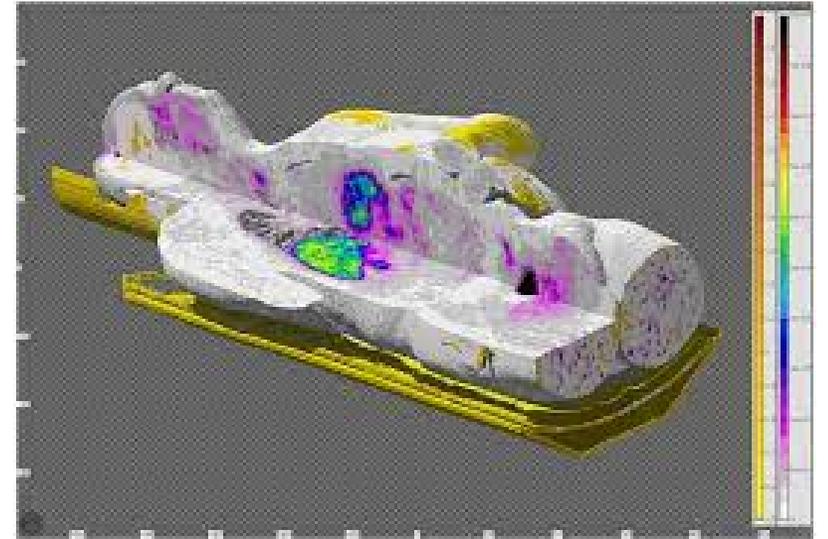
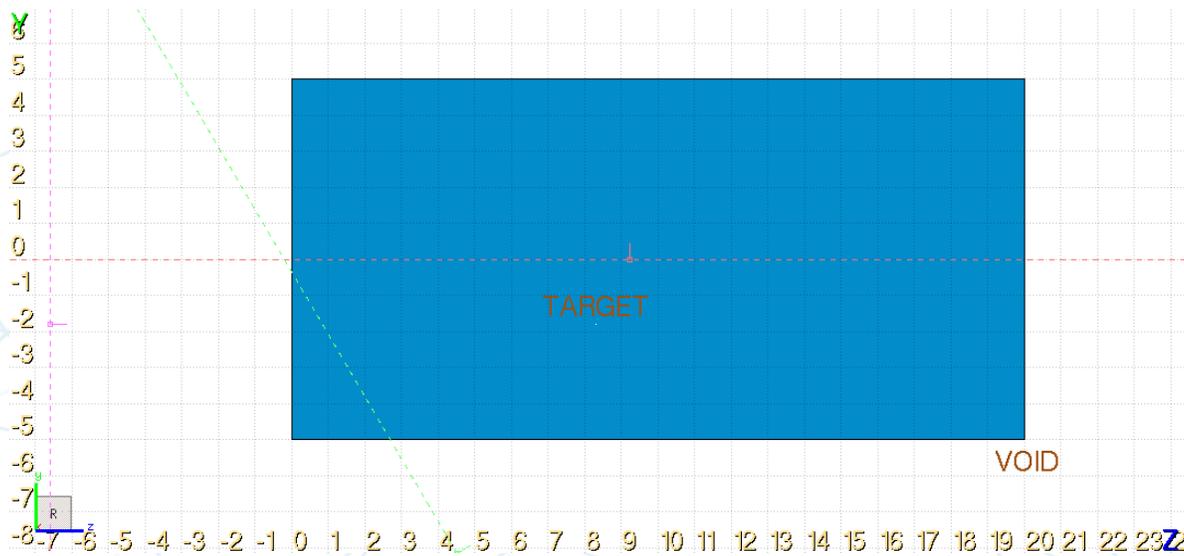


Image source: Wikipedia

- Particle-particle collision

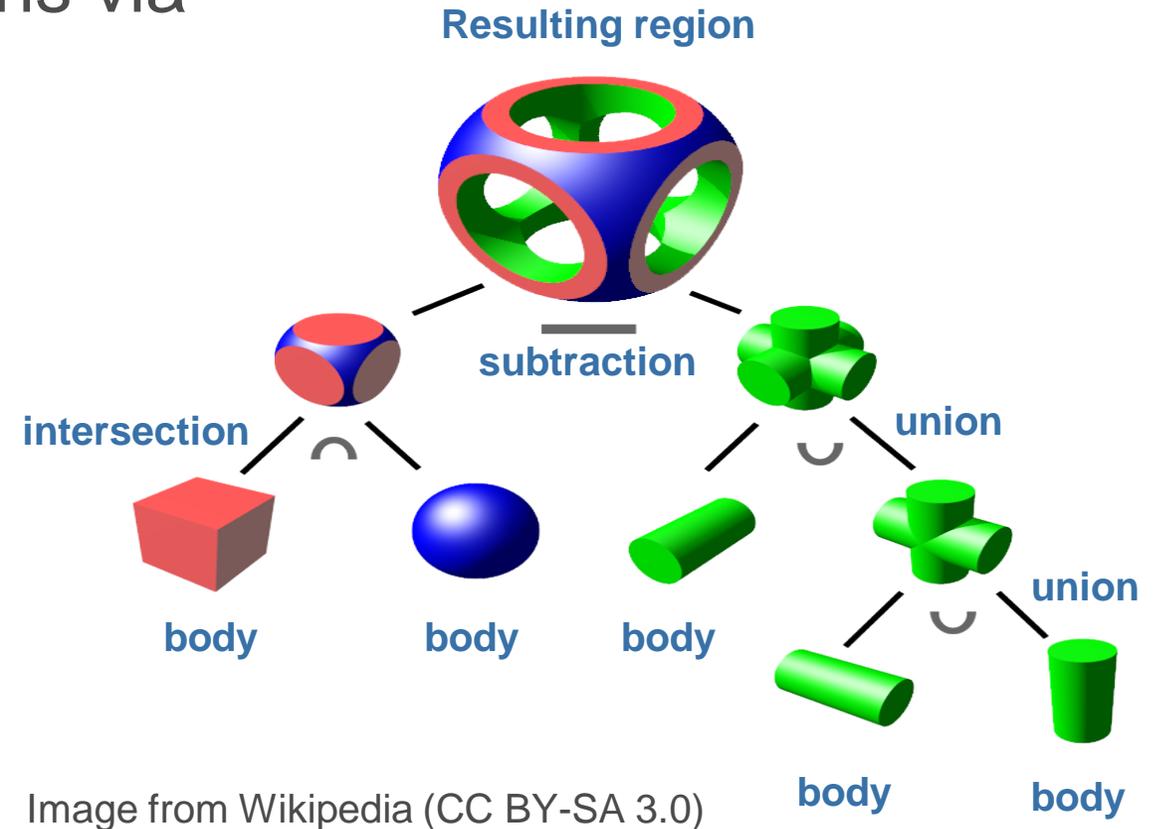




Basic ingredients for a MC code: 2 - Geometry

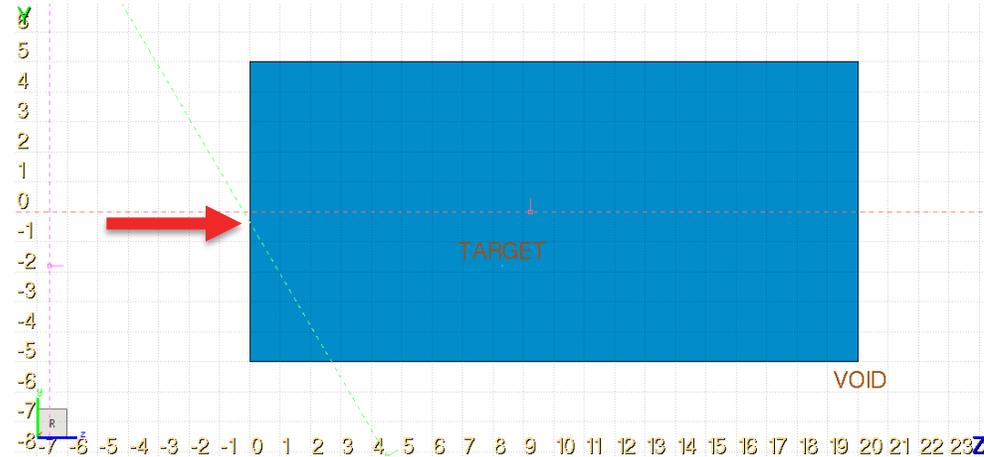
Combinatorial geometry package

- Basic objects (bodies): spheres, cylinders, parallelepipeds, planes...
- Combined to form complex regions via Boolean operations:
 - Union
 - Intersection
 - Subtraction
- Ability to clone/replicate regions
- Build macroscopic geometry to the desired degree of accuracy using these basic objects and operations

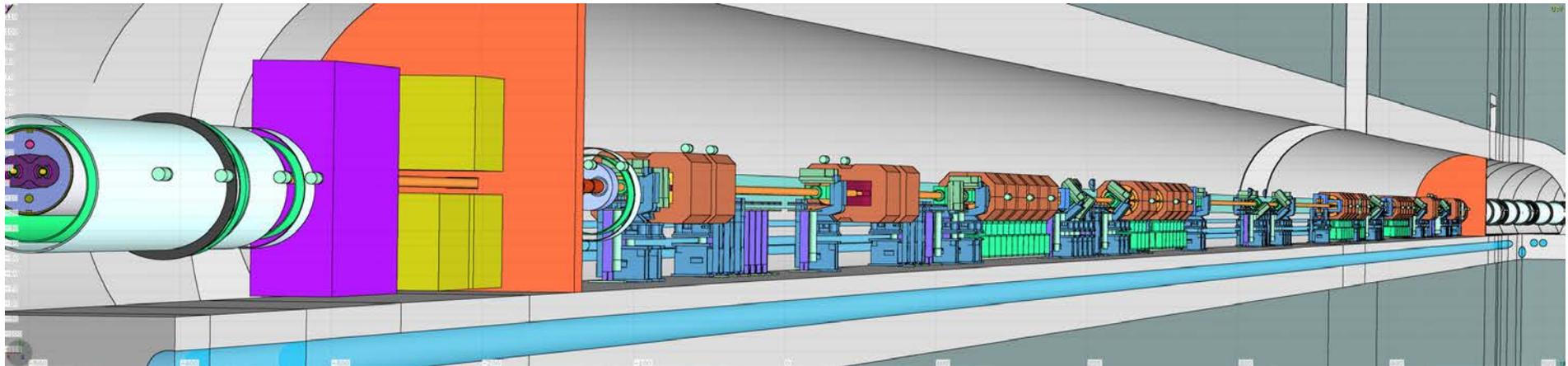


Ability to construct geometries, from trivial to complex

Trivial example: a simple material slab



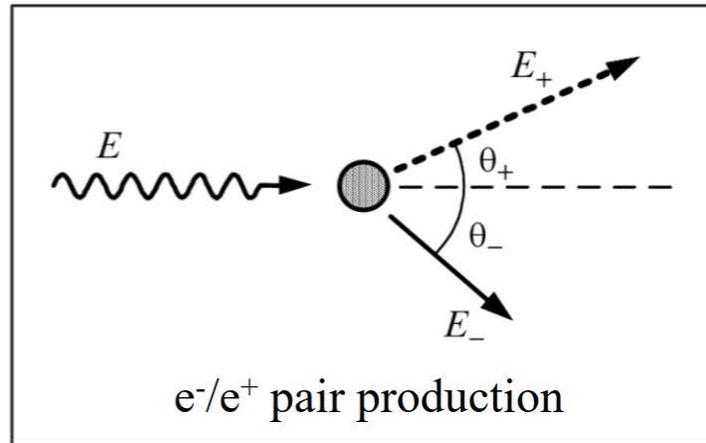
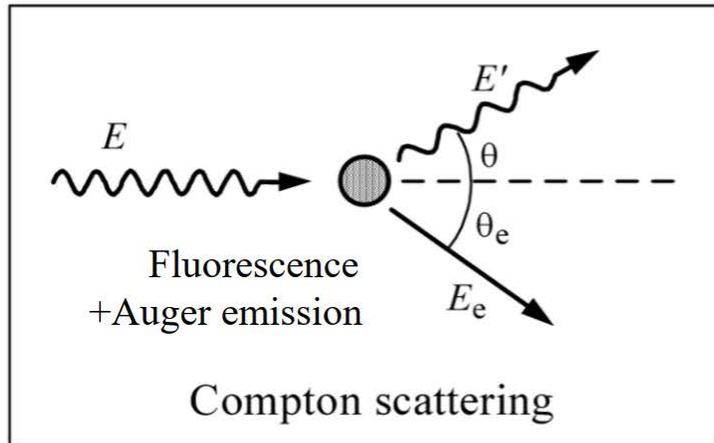
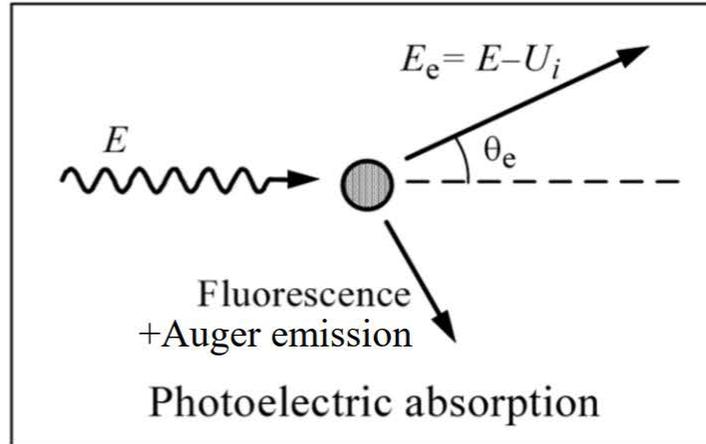
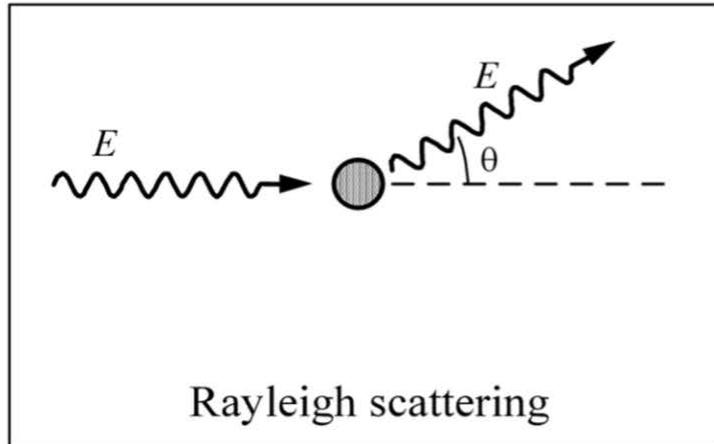
Complex example: LHC collimation insertion region (~100 m) built from the combination of **planes, cylinders, cuboids, ...**:





Basic ingredients for a MC code: 3 – Radiation-matter interactions

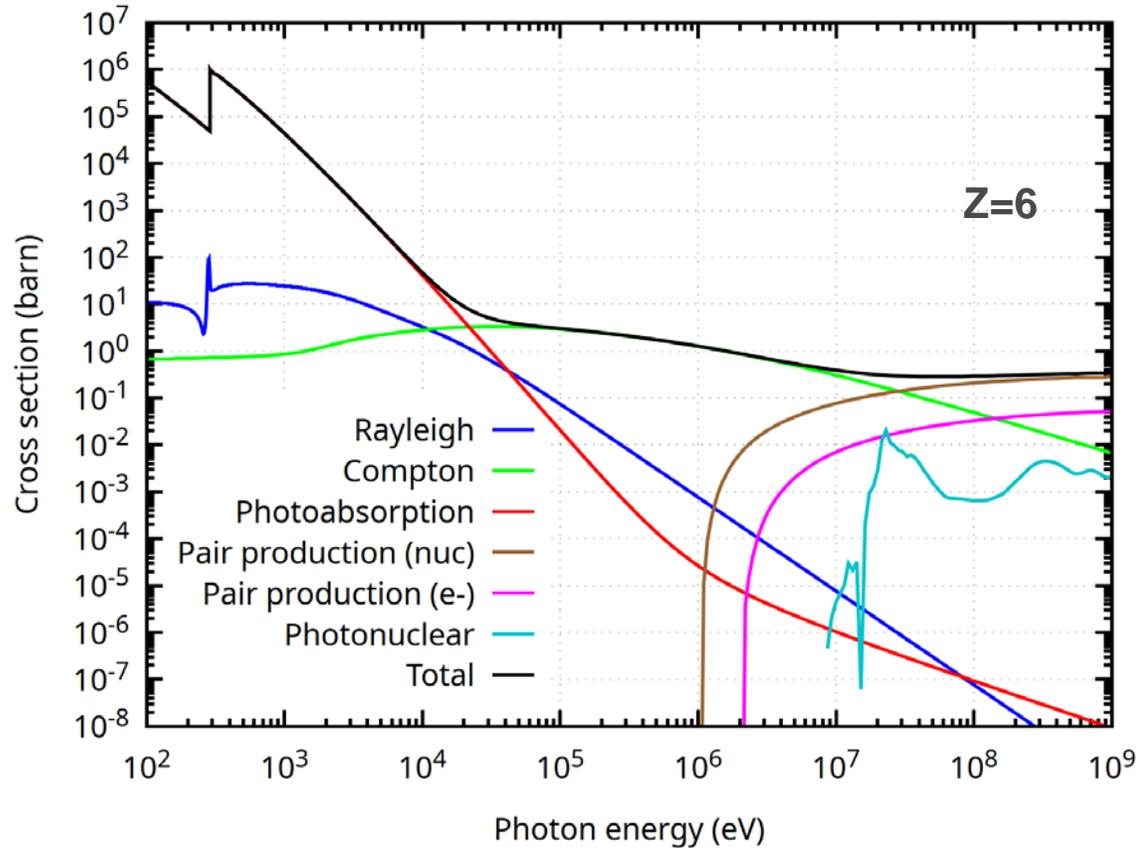
Basic interaction mechanisms of photons



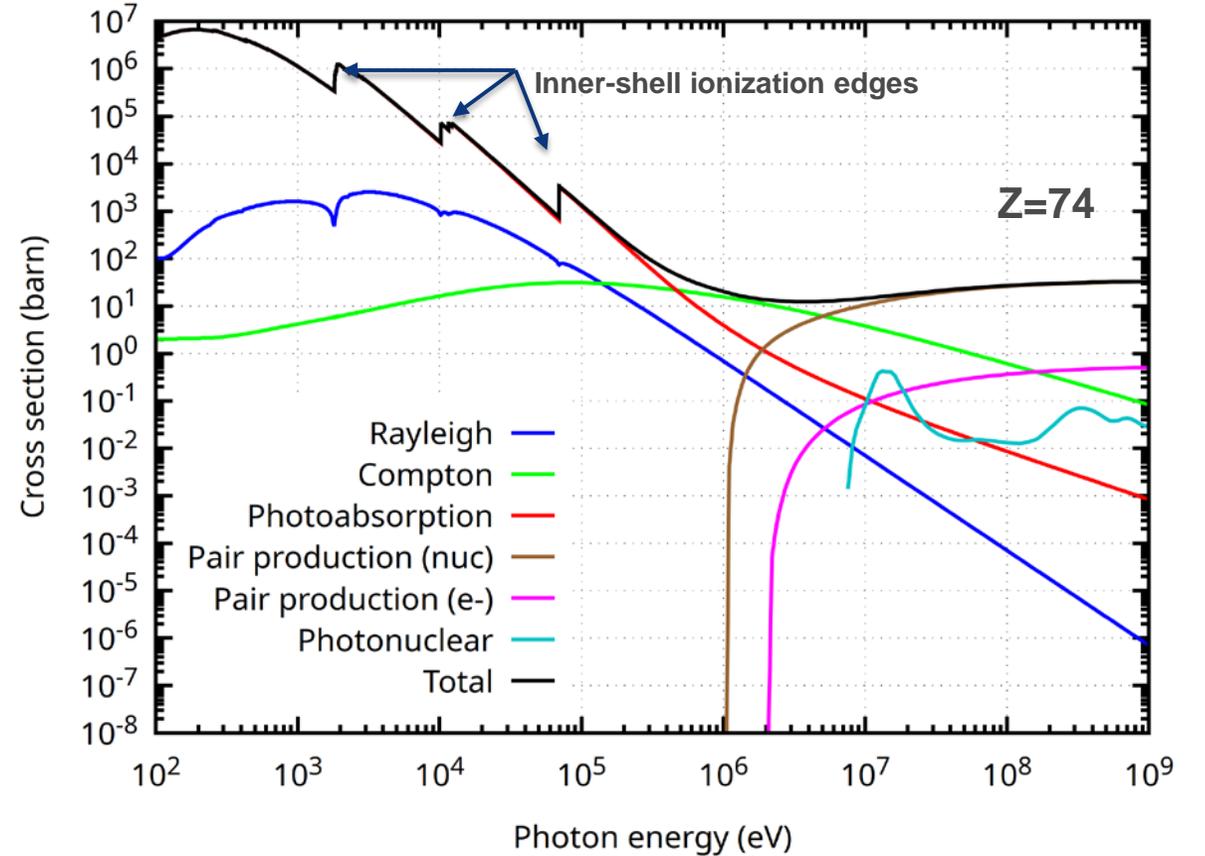
- Photonuclear reactions:
 - (γ, n) , $(\gamma, 2n)$, ...
- μ^\pm pair production
- At least 2 particles in the final state (except for Rayleigh scattering)
- Rapid increase of particles in the shower

Photon cross sections

Interaction cross sections for photons on C

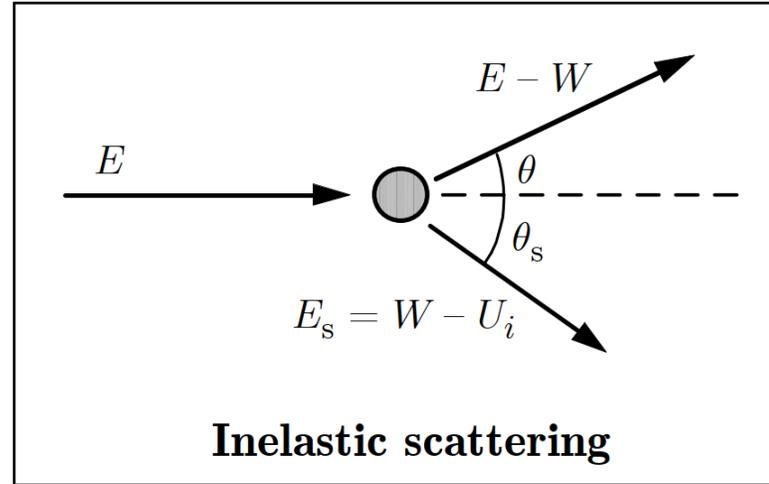
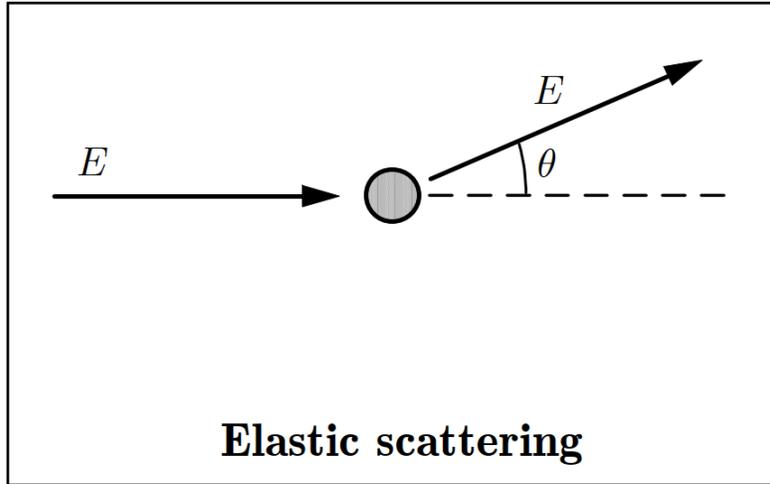


Interaction cross sections for photons on W

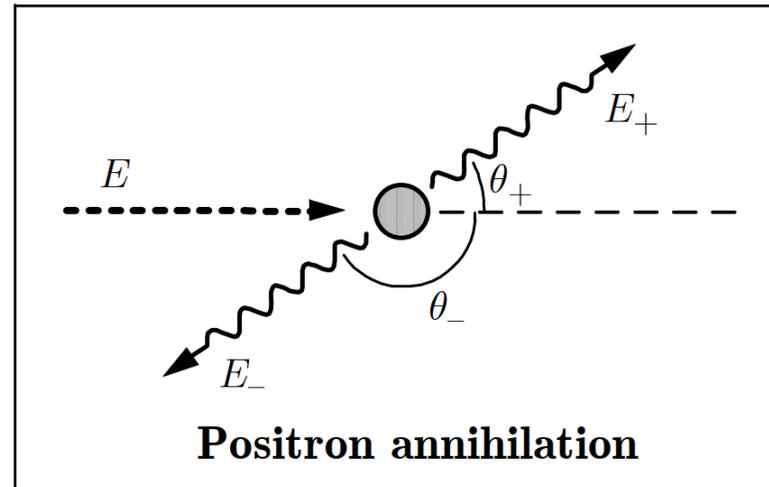
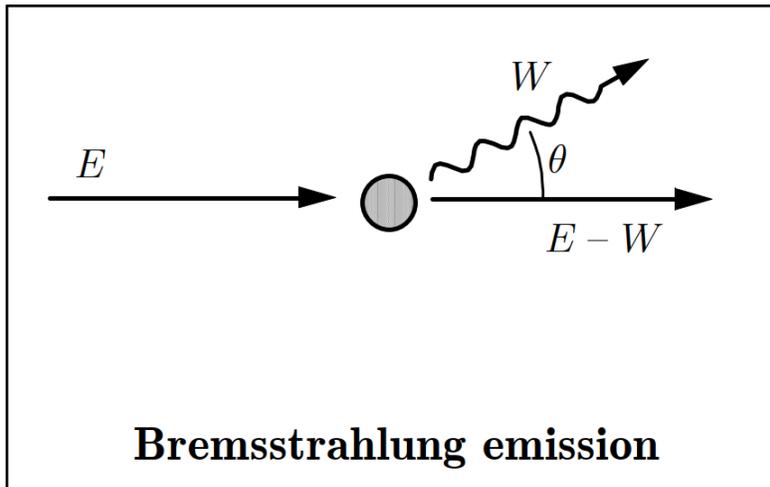


Basic interaction mechanisms of e^\pm

(Multiple scattering algorithms)

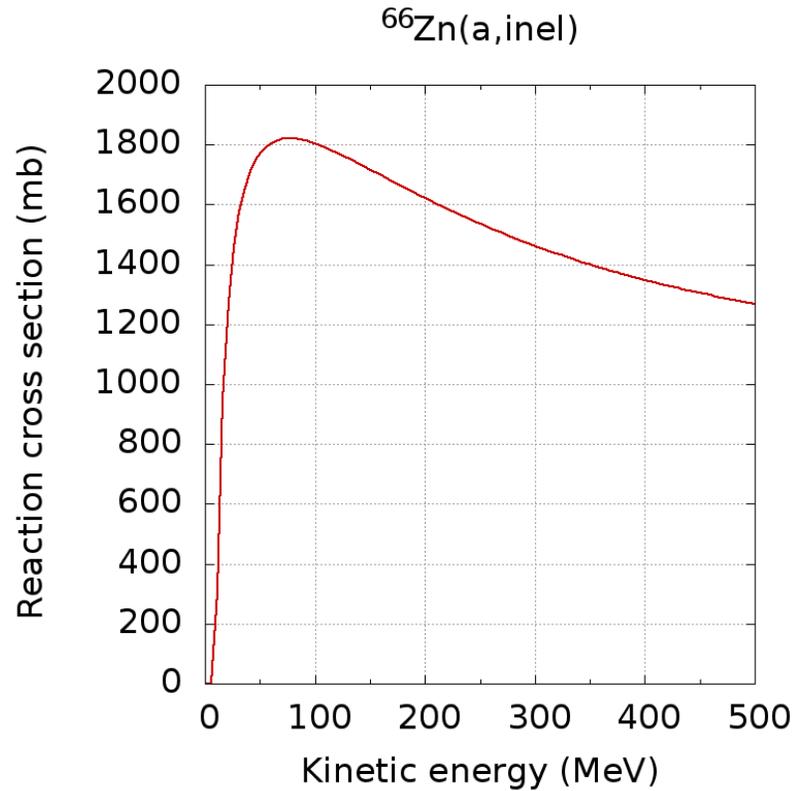


(Aggregate energy loss along particle step via stopping power. See backup slide)



Hadronic interactions

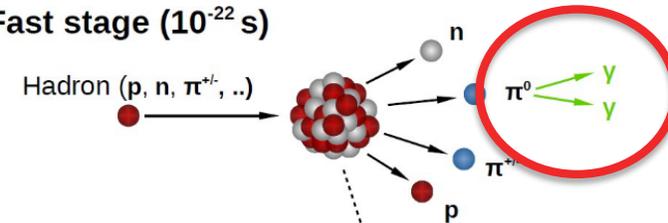
Reaction cross section (typically parametrized)



Interaction model

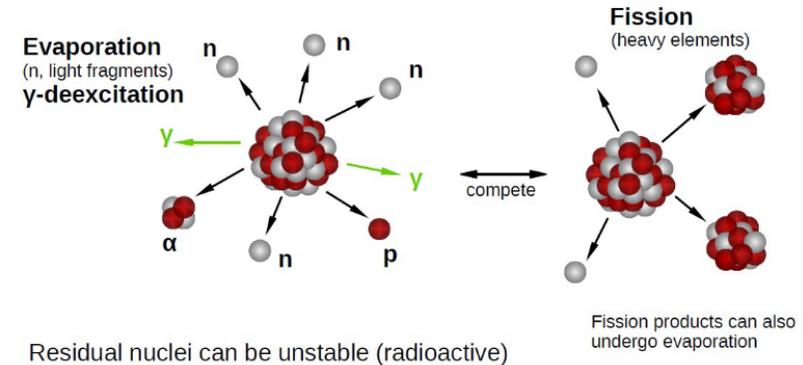
High-energy hadron on nucleus:

Fast stage (10^{-22} s)



Photons from π^0 decay couple hadronic shower with EM shower

Slow stage (10^{-16} s)

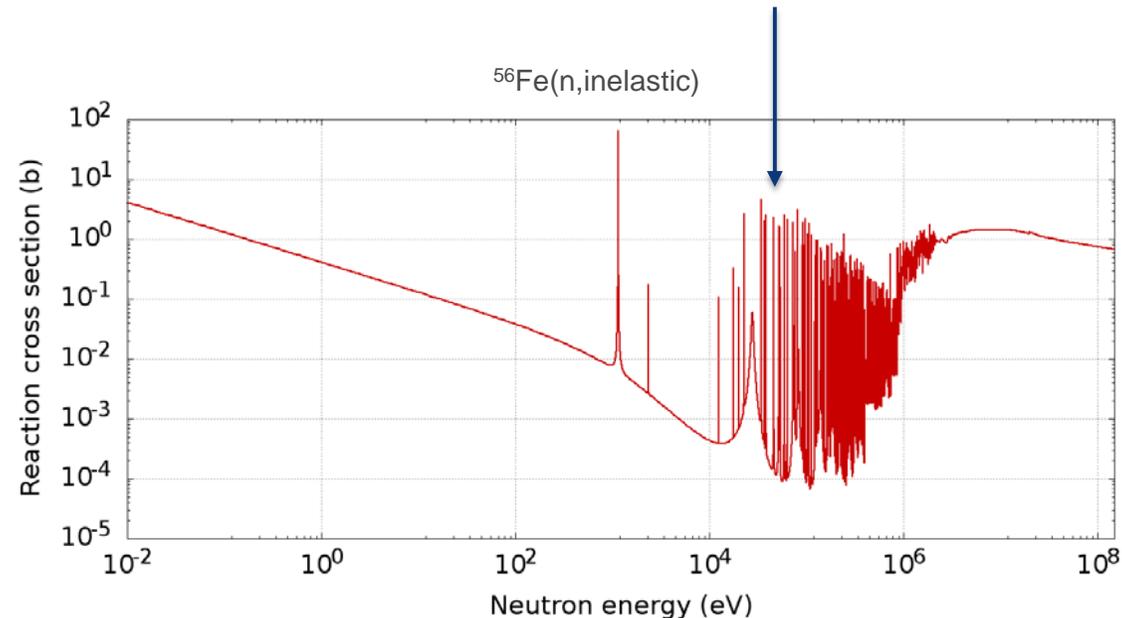


Residual nuclei can be unstable (radioactive)

Fission products can also undergo evaporation

Low-energy neutron interactions

- Neutron energy degrades quickly (elastic scattering or lower energy n secondaries)
- Low energies: **neutron capture** has high cross section and leads to MeV photons + production of residual nuclei. Relevant for radiation protection and shielding aspects
- Neutron cross sections have **rich resonant structure**:

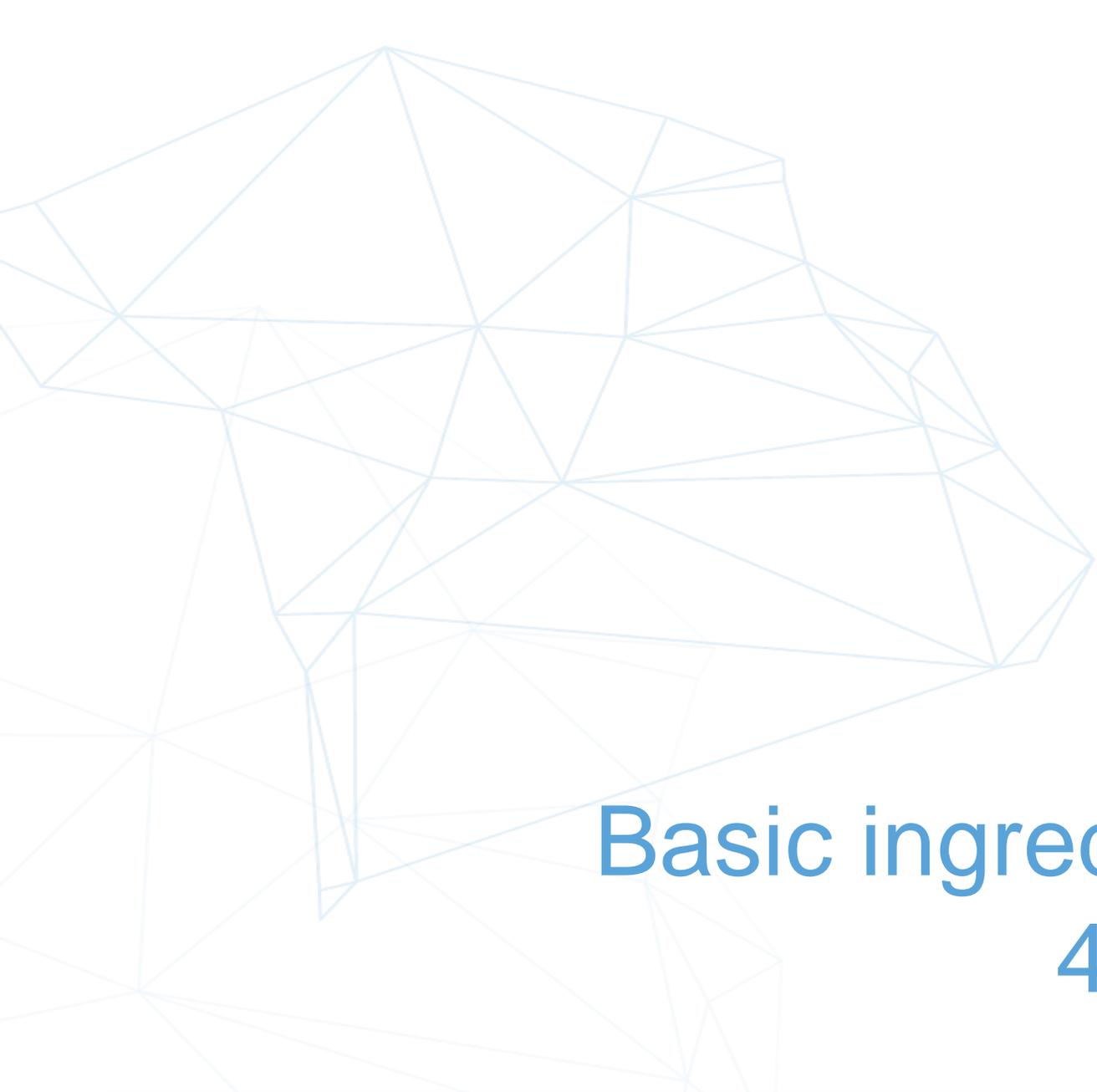


- No model available to effectively reproduce all resonant details

Low-energy neutron interactions - Approaches

- Either relying on a coarse-grained version of neutron cross sections (**group-wise** neutron interactions) or retaining as many details as possible (**point-wise** neutron interactions).
- Both rely on evaluated nuclear data libraries, e.g.:
 - Joint Evaluated Fission and Fusion (JEFF)
https://www.oecd-nea.org/jcms/pl_27365/the-jeff-nuclear-data-library
 - Evaluated Nuclear Data Files (ENDF)
<https://www-nds.iaea.org/exfor/endl.htm>
 - Japanese Evaluated Nuclear Data Library (JENDL)
<https://wwwndc.jaea.go.jp/jendl/jendl.html>
 - etc
- Covering: elastic scattering, capture, fission, and explicit inelastic channels e.g. (n,p), (n,a), etc.

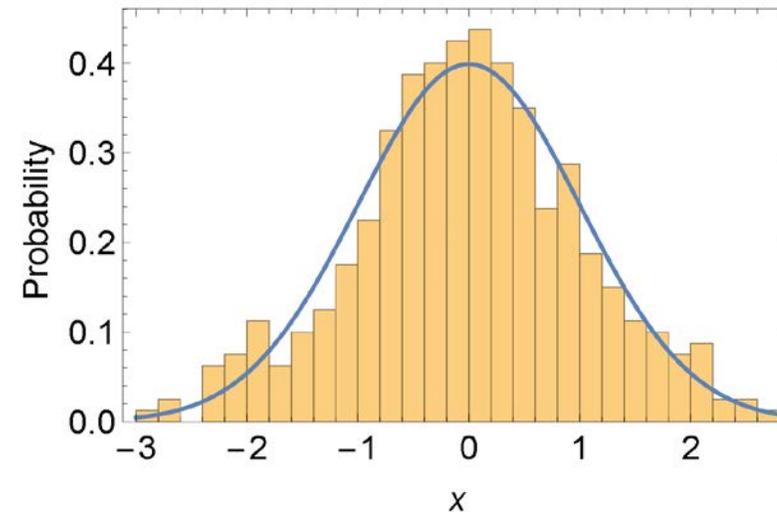
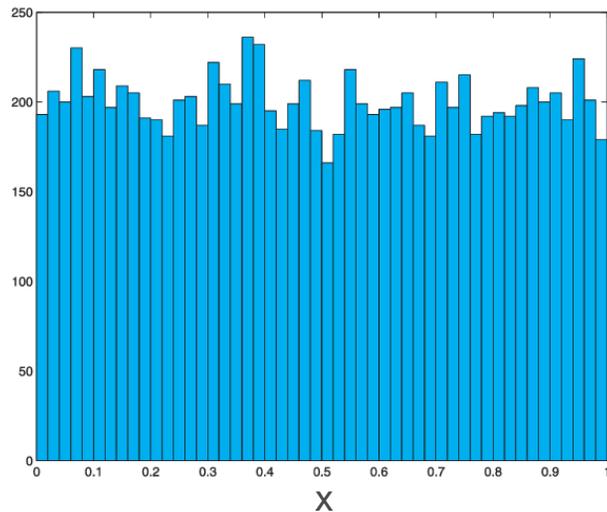




Basic ingredients for a MC code: 4 – Random sampling

(Pseudo-)Random number generator (PRNG)

- Function returning homogeneously sampled values from 0 to 1
- **Sampling techniques:** exploit PRNG values to sample from arbitrary distributions (e.g. angular distribution of secondaries)



<https://www.researchgate.net/publication/230860021> An introduction to the numerical analysis of the Boltzmann equation

<https://www.researchgate.net/publication/360098338> Optimisation of Thimble Simulations and Quantum Dynamics of Multiple Fields in Real Time



Basic flow diagram of a Monte Carlo simulation of radiation transport

Mean free path

Cross section

$$\sigma \equiv \int \int \frac{d^2\sigma}{d\Omega dW} d\Omega dW$$

~likelihood of being scattered

Dimensions: Area

Typical unit: 1 barn (=10⁻²⁴ cm²)

- Cross section (slide 6):

- Consider a medium with \mathcal{N} atoms per unit volume
- The mean free path

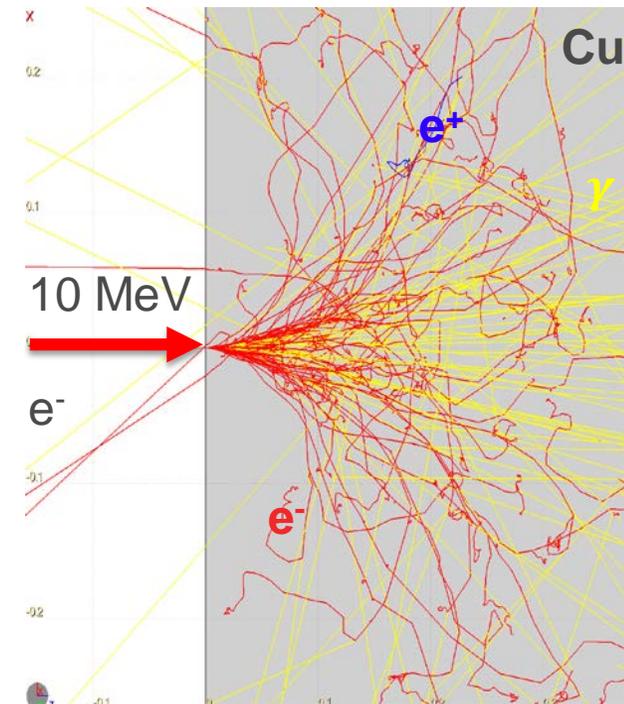
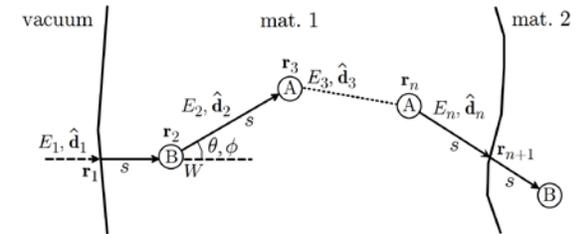
$$\lambda = \frac{1}{\mathcal{N}\sigma}$$

is a quantity with units of length that gives the average distance to the next interaction

- One can see (backup slides) that it follows an exponential distribution

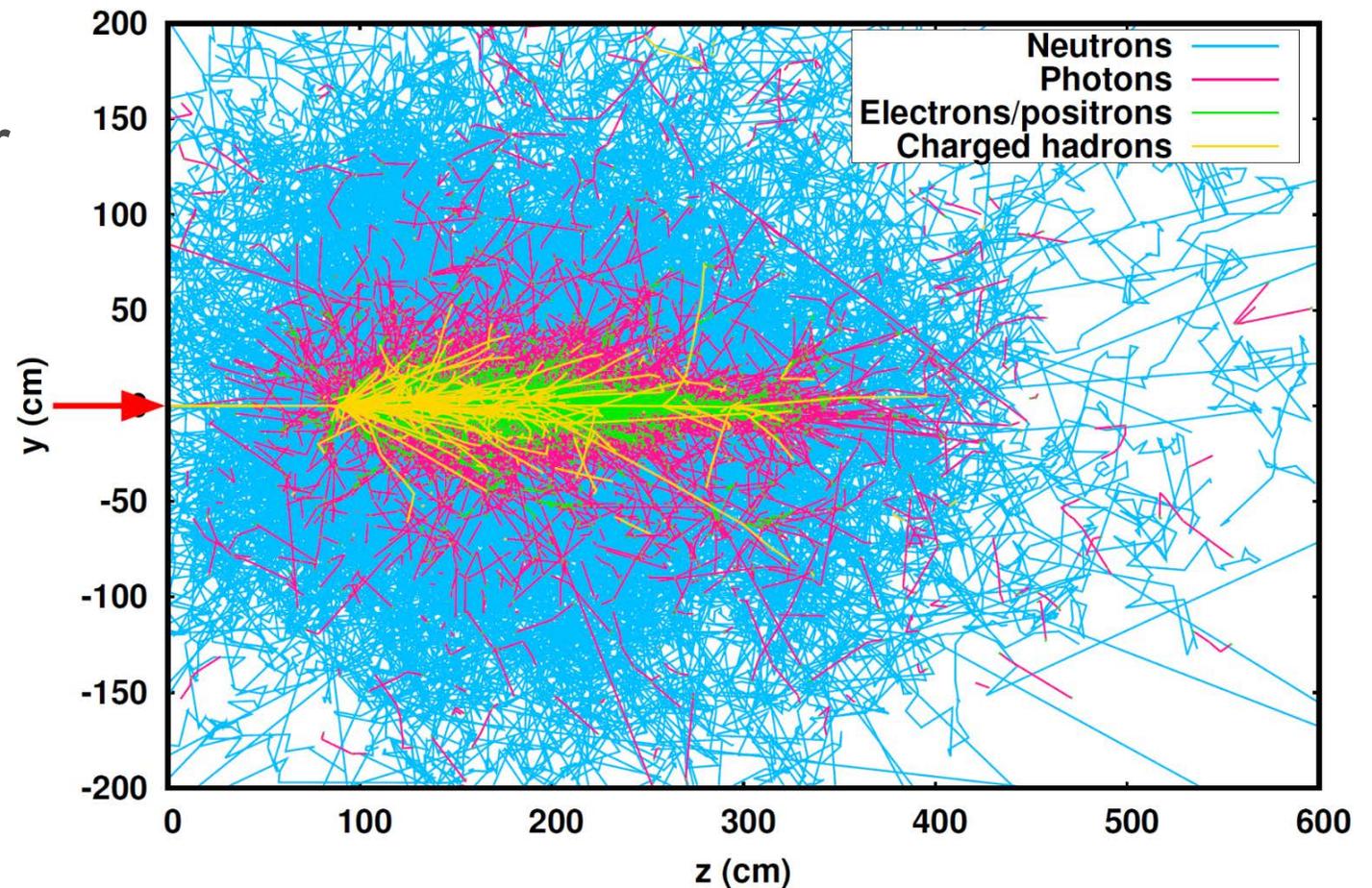
Simplified flowchart of a MC simulation of particle transport

- Define radiation source and material geometry
- Do for a sufficiently large number N of primary particles:
 1. New primary event: add 1 (or more!) particles to stack
 2. Pick up particle from stack, advance to next material medium (score contribution of this step to observables)
 3. Evaluate mean free path $\lambda(E)$
 4. Sample random step length to next interaction
 5. Accumulate contribution of step to statistical estimators
 6. Decide kind of interaction that takes place
 7. Sample final state (add possible secondaries to particle stack)
 8. Accumulate contribution of interaction to statistical estimators
 9. If primary particle survives, update energy/direction, go to 3. Otherwise, go to 2 (or go to 1 if no more particles in stack).



MC simulation of hadronic and electromagnetic radiation showers

- Coupled hadronic and electromagnetic shower set up by a **single** 450 GeV proton in Al





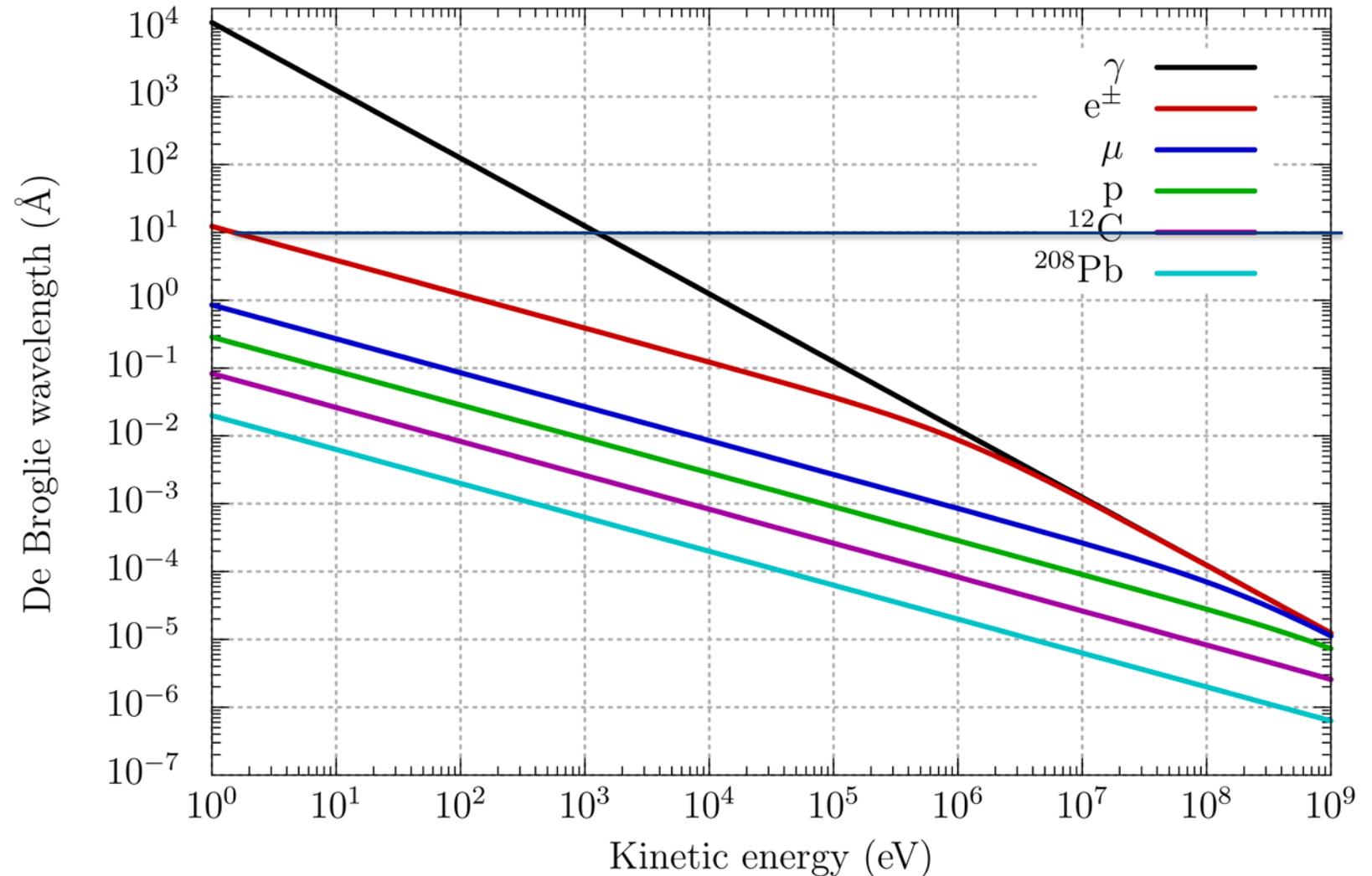
Limitations

Assumptions underlying MC simulation of radiation transport

- Medium is **amorphous** and **isotropic**:
 - No diffraction effects out of the box
 - Small subset: e.g. channeled charged particles in bent crystals (FLUKA, Geant4)
- **Material is not affected** by previous particles
- Shower particles do not interact with each other
- Transport is assumed to be Markovian: fate of a particle does not depend on its past history
- Particles interact on an individual atom/nucleus

De Broglie wavelength vs typical interatomic distance

- At low energies, wavelength extends to several interatomic distance
- Picture of a particle interacting on a single atom breaks down



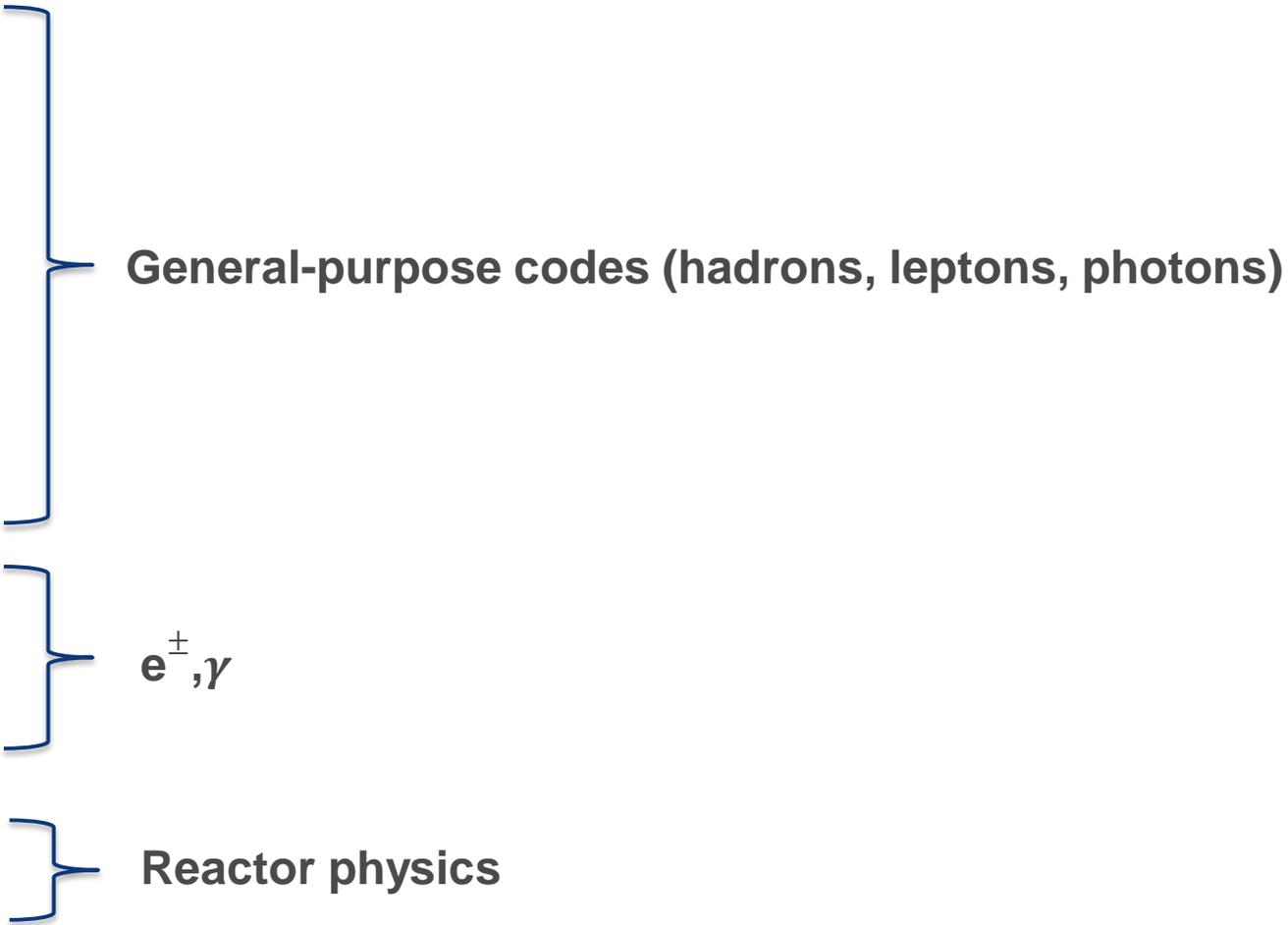


General-purpose MC codes for particle interaction simulations

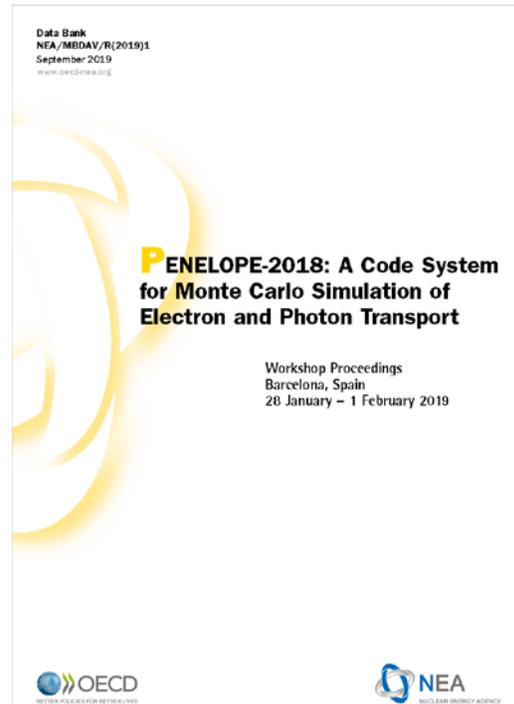
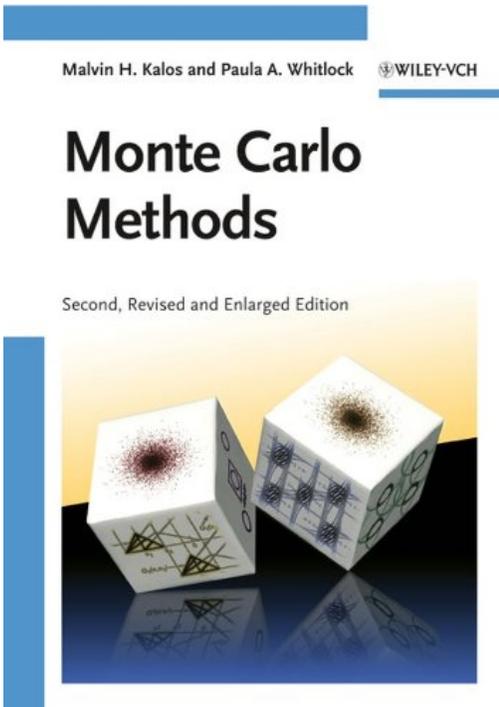
Capabilities of general-purpose MC codes

- Transport e^- , e^+ , γ as well as muons, hadrons, and ions
- Allow for particle beams, particle-particle collision, or radioactive isotope as radiation source
- Cover a broad energy range:
 - Up to TeV for accelerator physics
 - Down to 1 keV for e^- , e^+ , gamma
 - Down to meV for neutrons
- Offer up-to-date and thoroughly benchmarked physics
- Offer a flexible and efficient geometry package
- Capability to score physical observables of interest: energy deposition, particle spectra, radioactive inventories, etc.
- Capability to run in parallel / distributed

General-purpose codes for MC simulation of radiation transport

- Geant4
 - FLUKA
 - PHITS
 - MARS
 - MCNP
 - PENELOPE
 - EGS
 - Tripoli
 - ...
- General-purpose codes (hadrons, leptons, photons)
- e^{\pm}, γ
- Reactor physics
- 
- A diagram showing a list of Monte Carlo codes on the left, grouped into three categories on the right. The first group, 'General-purpose codes (hadrons, leptons, photons)', includes Geant4, FLUKA, PHITS, MARS, and MCNP. The second group, 'e±, γ', includes PENELOPE and EGS. The third group, 'Reactor physics', includes Tripoli and an ellipsis '...'. Each group is indicated by a blue bracket on the right side of the list.

References



CERN

5 Sep 2022



Geant4 Collaboration

Rev5.0: December 4th, 2020

Thank you very much
for your attention!



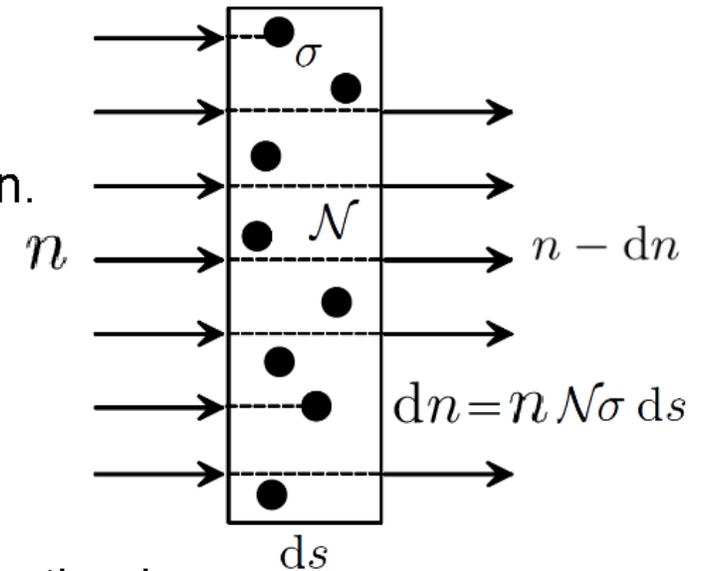
The exponential distribution of path lengths

- Let n particles per unit time and surface impinge normally on a thin material slab of width ds with a density of N scattering centers per unit volume, each having a cross sectional area σ .
- Number of particles that interacted: $dn = n N \sigma ds$.
- The interaction probability in ds : $dn/n = N \sigma ds$
- Let $p(s)$ be the distribution of path lengths to the next interaction.
- The probability that the next interaction is within ds of s is

$$p(s) = [1 - \int_0^s ds' p(s')] (N \sigma) = \int_s^{inf} p(s') (N \sigma) ds'$$
- The solution of this diff eq is

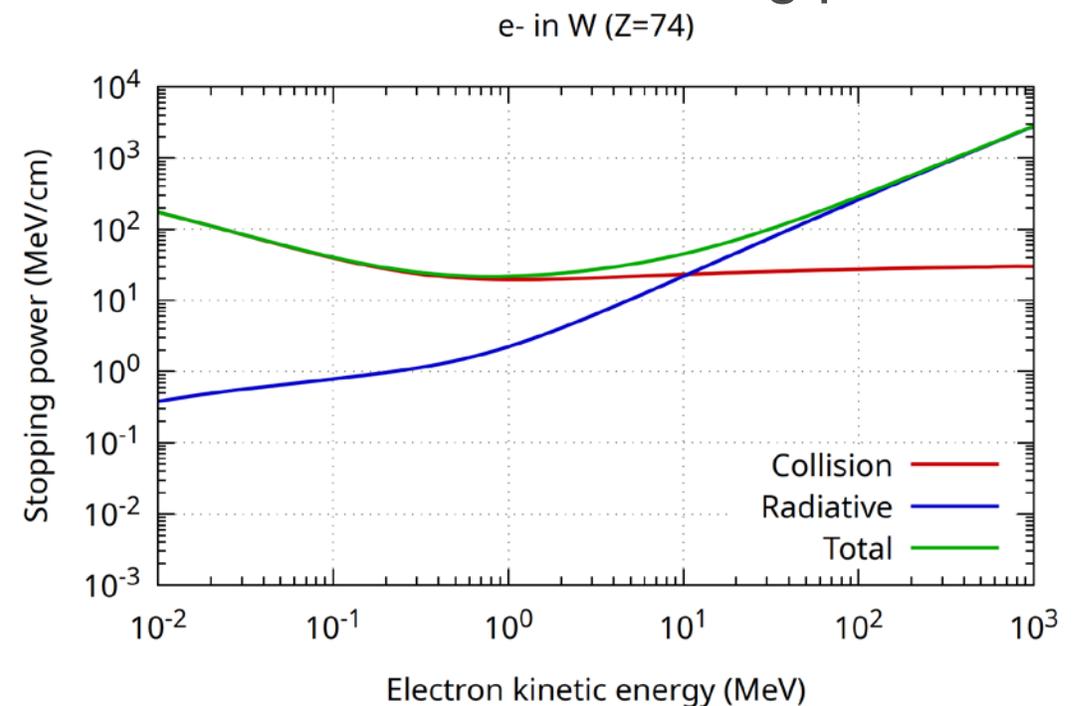
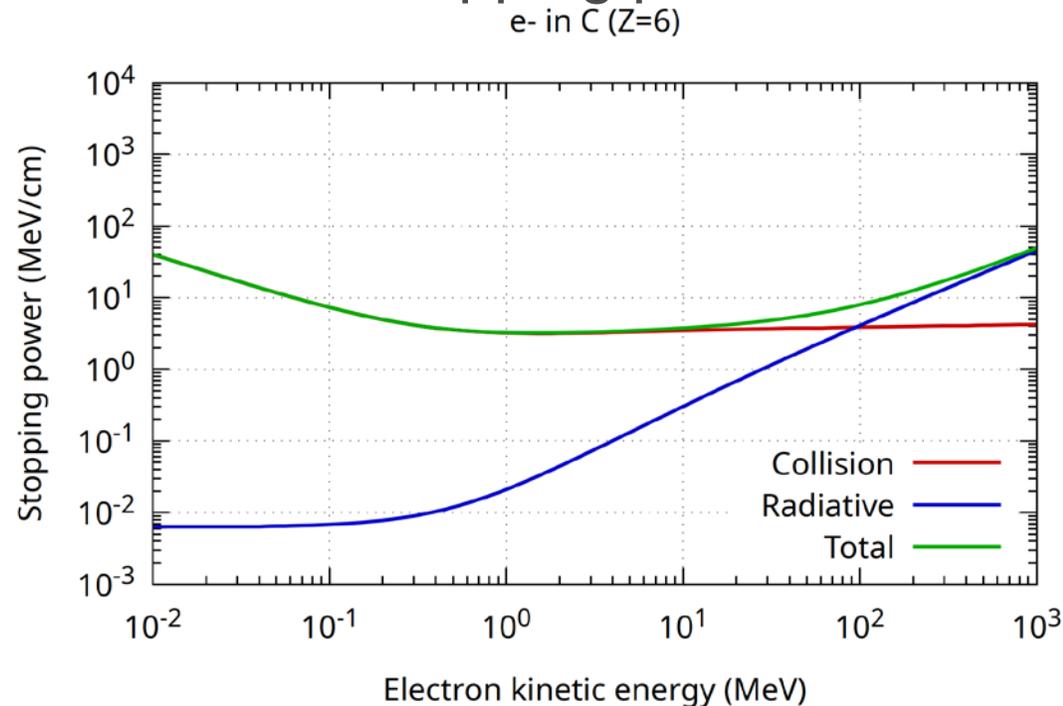
$$p(s) = (N \sigma) e^{-s(N \sigma)}$$
- Thus the path length to the next interaction follows an exponential distribution. The average distance to the next interaction is:

$$\langle s \rangle = 1/(N \sigma) = \lambda,$$
 i.e., we recover the expression of the mean free path given above.



Stopping power of charged particles in matter

- **Stopping power:** energy loss per unit path length
 - **Collision** stopping power: due to collisions with target e^-
 - **Radiative** stopping power: due to the emission of Bremsstrahlung photons



See: <https://physics.nist.gov/PhysRefData/Star/Text/intro.html>