


$$g \log(2) = \lambda_1 \log(2) + \nu_2(2i\pi)$$

Observation of Odderon Effects at LHC energies

A Real Extended Bialas-Bzdak Model Study

based on [T. Csörgő, I. Szanyi, Eur. Phys. J. C **81**, 611 \(2021\)](#) and other recent results

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Day of Femtoscopy 2021

October 28, 2021, Gyöngyös, Hungary

Unitarity and the elastic amplitude

- unitarity of the scattering matrix S :

$$SS^\dagger = I$$

$$S = I + iT$$

$$T - T^\dagger = iTT^\dagger$$

- unitarity equation in impact parameter \vec{b} representation:

$$2 \operatorname{Im} t_{\text{el}}(s, \vec{b}) = |t_{\text{el}}(s, \vec{b})|^2 + \tilde{\sigma}_{in}(s, \vec{b})$$

(s is the squared CM energy)

- the elastic amplitude can be written as a solution of the unitarity equation in terms of $\tilde{\sigma}_{in}$
- $0 \leq \tilde{\sigma}_{in}(s, \vec{b}) \leq 1$ and $\tilde{\sigma}_{in}$ can be calculated using the rules of the probability calculus based on Glauber's multiple scattering theory

Bialas-Bzdak p=(q,d) model

$$\tilde{\sigma}_{in}(\vec{b}) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} d^2\vec{s}_q d^2\vec{s}'_q d^2\vec{s}_d d^2\vec{s}'_d D(\vec{s}_q, \vec{s}_d) D(\vec{s}'_q, \vec{s}'_d) \sigma(\vec{s}_q, \vec{s}_d; \vec{s}'_q, \vec{s}'_d; \vec{b})$$

- quark-diquark distribution inside the proton:

$$D(\vec{s}_q, \vec{s}_d) = \frac{1 + \lambda^2}{R_{qd}^2 \pi} e^{-\frac{s_q^2 + s_d^2}{R_{qd}^2}} \delta^2(\vec{s}_q + \lambda \vec{s}_d)$$

$$\begin{aligned} \vec{s}_d &= -\lambda \vec{s}_q \\ \lambda &= \frac{m_q}{m_d} \\ \vec{s}'_d &= -\lambda \vec{s}'_q \end{aligned}$$

A. Bialas, A. Bzdak Acta
Phys.Polon. B 38, 159-168 (2007)

- interaction probability of the constituents:

$$\sigma(\vec{s}_q, \vec{s}_d; \vec{s}'_q, \vec{s}'_d; \vec{b}) = 1 - \prod_a \prod_b [1 - \sigma_{ab}(\vec{b} + \vec{s}'_a - \vec{s}_b)]$$

$$\sigma_{ab}(\vec{s}) = A_{ab} e^{-|\vec{s}|^2/S_{ab}^2} \quad S_{ab}^2 = R_a^2 + R_b^2 \quad a, b \in \{q, d\}$$

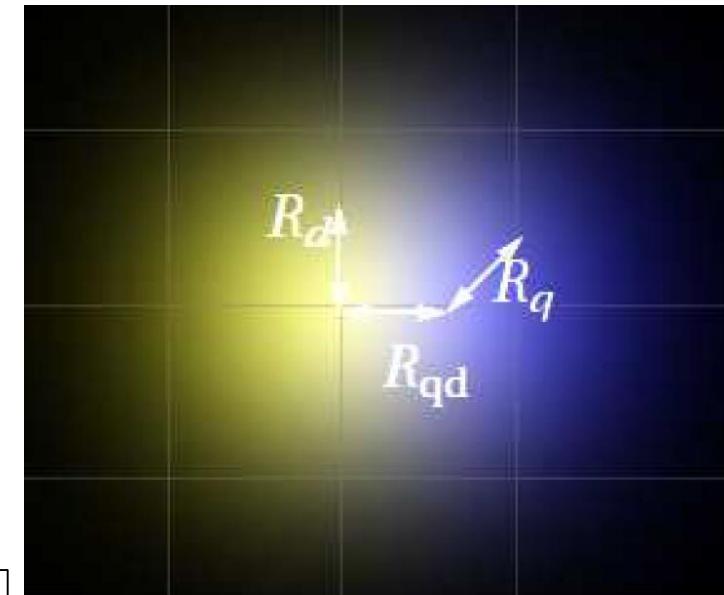
- inelastic cross-sections of quark, diquark scatterings :

$$\sigma_{ab,in} = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \sigma_{ab}(\vec{s}) d^2\vec{s}$$

$$\sigma_{qq,in} : \sigma_{qd,in} : \sigma_{dd,in} = 1 : 2 : 4$$

- free parameters:

$$A_{qq}, \lambda, R_q, R_d, R_{qd}, (A_{qq} = 1 \text{ and } \lambda = 0.5 \text{ can be fixed})$$



Proton-(anti)proton scattering in
the quark-diquark model.

Unitarily Real Extended Bialas-Bzdak (ReBB) model

- elastic scattering amplitude in the impact parameter space:

$$t_{el}(s, \vec{b}) = i [1 - e^{-\Omega(s, \vec{b})}]$$

arXiv:1505.01415

F. Nemes, T. Csörgő, M. Csanád, Int. J. Mod. Phys. A Vol. 30 (2015) 1550076

- the opacity function:

$$\Omega(s, \vec{b}) = Re\Omega(s, \vec{b}) + i Im\Omega(s, \vec{b})$$

$Im\Omega \neq 0$ as the real part of the amplitude is not negligibly small

$$Re\Omega(s, \vec{b}) = -\frac{1}{2} \ln[1 - \tilde{\sigma}_{in}(s, \vec{b})]$$

$$Im\Omega(s, \vec{b}) = -\alpha \tilde{\sigma}_{in}(s, \vec{b})$$



NEW FREE PARAMETER

- elastic scattering amplitude in momentum space:

$$T(s, t) = 2\pi \int_0^\infty t_{el}(s, |\vec{b}|) J_0(|\vec{\Delta}| |\vec{b}|) |\vec{b}| d|\vec{b}|$$

$$|\vec{\Delta}| \equiv \sqrt{-t} \text{ as } \sqrt{s} \rightarrow \infty$$

(t is the squared momentum transfer)

Measurable quantities

- differential cross section:

$$\frac{d\sigma}{dt}(s, t) = \frac{1}{4\pi} |T(s, t)|^2$$

- total, elastic and inelastic cross sections:

$$\sigma_{tot}(s) = 2Im T(s, t = 0)$$

$$\sigma_{el}(s) = \int_{-\infty}^0 \frac{d\sigma(s, t)}{dt} dt$$

$$\sigma_{in}(s) = \sigma_{tot}(s) - \sigma_{el}(s)$$

- ratio ρ_0 :

$$\rho_0(s) = \lim_{t \rightarrow 0} \rho(s, t) \equiv \frac{Re T(s, t \rightarrow 0)}{Im T(s, t \rightarrow 0)}$$

- slope of $d\sigma/dt$:

$$B_0(s) = \lim_{t \rightarrow 0} B(s, t) \equiv \left. \frac{d}{dt} \left(\ln \frac{d\sigma}{dt}(s, t) \right) \right|_{t \rightarrow 0}$$

The $p p$ and $p \bar{p}$ elastic scattering amplitude

- according to the Regge formalism the strong scattering amplitude for $p p$ and $p \bar{p}$ scattering can be written in terms of $C = +1$ and $C = -1$ exchange components

$$T^{pp}(s, t) = T^+(s, t) - T^-(s, t)$$

$$T^{p\bar{p}}(s, t) = T^+(s, t) + T^-(s, t)$$

- for $\sqrt{s} \gtrsim 1$ TeV the mesonic reggeon exchanges are negligible only the gluonic Pomeron and Odderon exchanges are present implying that

$$T^+(s, t) \equiv T^P(s, t)$$

$$T^-(s, t) \equiv T^O(s, t)$$



$$T^P(s, t) = \frac{1}{2}(T^{pp}(s, t) + T^{p\bar{p}}(s, t))$$

$$T^O(s, t) = \frac{1}{2}(T^{p\bar{p}}(s, t) - T^{pp}(s, t))$$

- a simple and model independent consequence:

if $\frac{d\sigma^{pp}}{dt}(s, t) \neq \frac{d\sigma^{p\bar{p}}}{dt}(s, t)$ for $\sqrt{s} \gtrsim 1$ TeV then $T^O(s, t) \neq 0$

Fit method

- least squares fitting with:

$$\chi^2 = \left(\sum_{j=1}^M \left(\sum_{i=1}^{n_j} \frac{(d_{ij} + \epsilon_b j \tilde{\sigma}_{bij} + \epsilon_c j d_{ij} \sigma_{cj} - t h_{ij})^2}{\tilde{\sigma}_{ij}^2} \right) + \epsilon_b^2 + \epsilon_c^2 \right) + \left(\frac{d_{\sigma_{tot}} - t h_{\sigma_{tot}}}{\delta \sigma_{tot}} \right)^2 + \left(\frac{d_{\rho_0} - t h_{\rho_0}}{\delta \rho_0} \right)^2$$

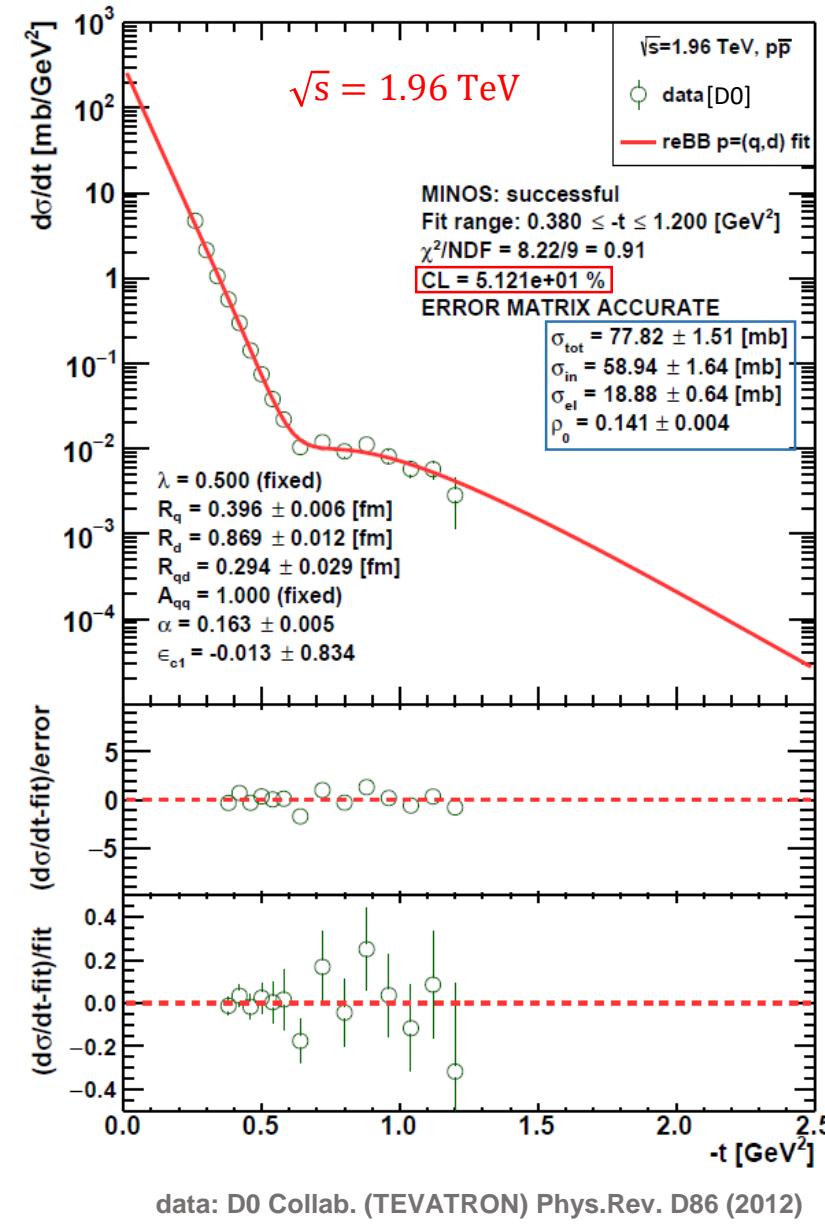
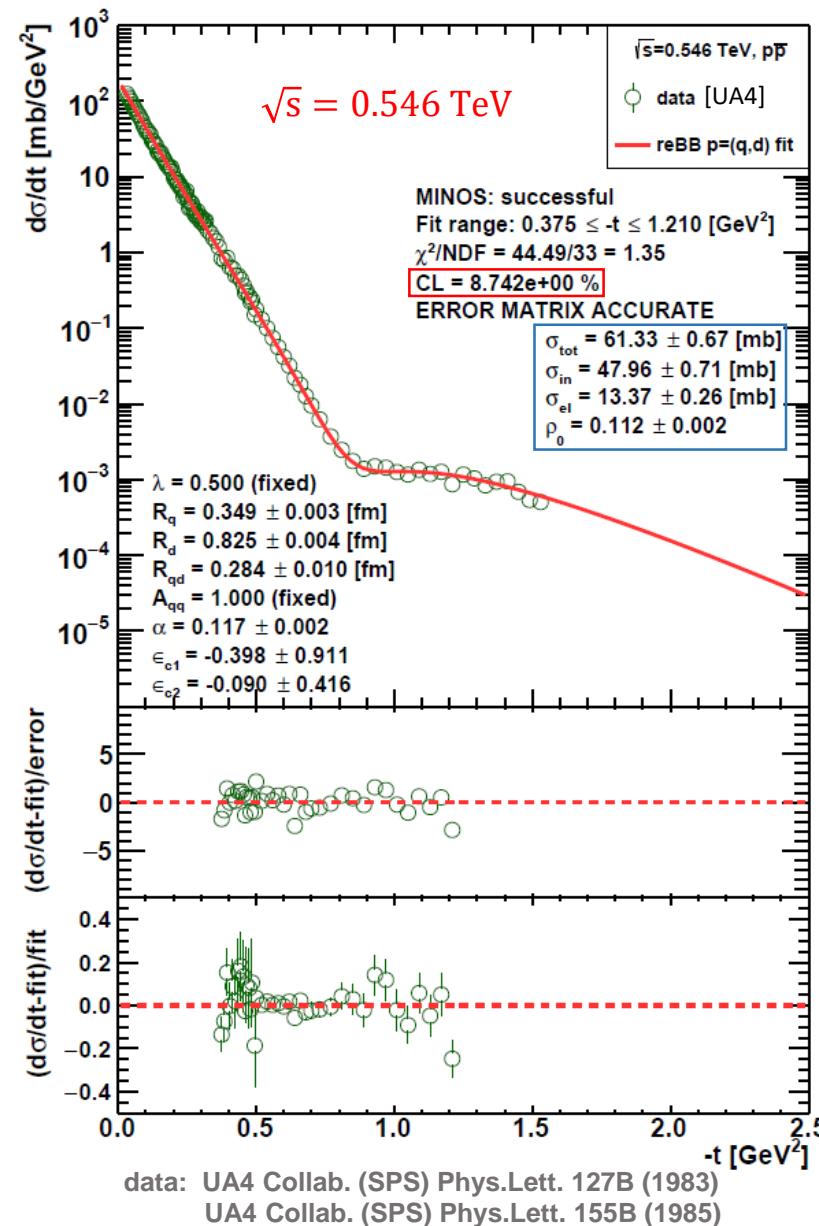
$$\tilde{\sigma}_{ij}^2 = \tilde{\sigma}_{aij} \left(\frac{d_{ij} + \epsilon_b j \tilde{\sigma}_{bij} + \epsilon_c j d_{ij} \sigma_{cj}}{d_{ij}} \right)$$

$$\tilde{\sigma}_{kij} = \sqrt{\sigma_k^2 + (d'_{ij} \delta_k t_{ij})^2}, \quad k \in \{a, b\}$$

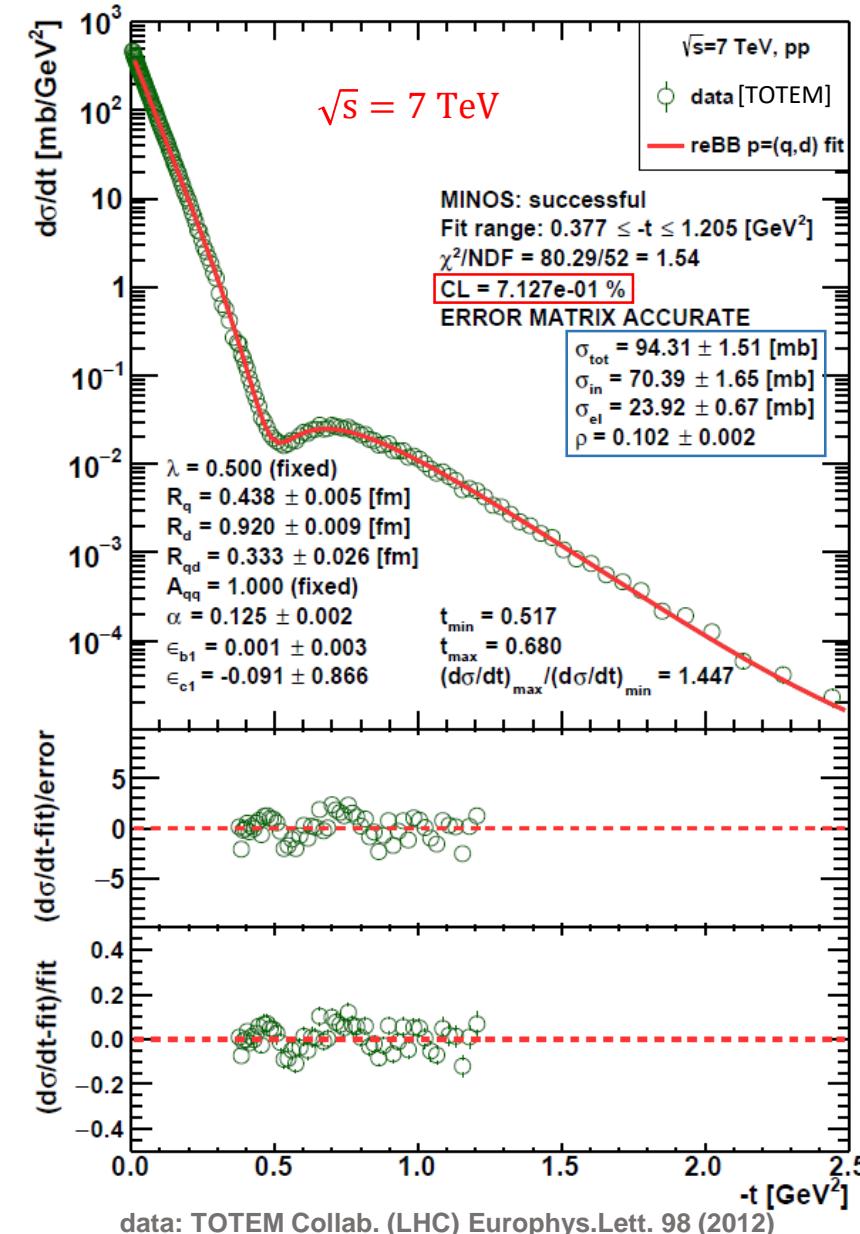
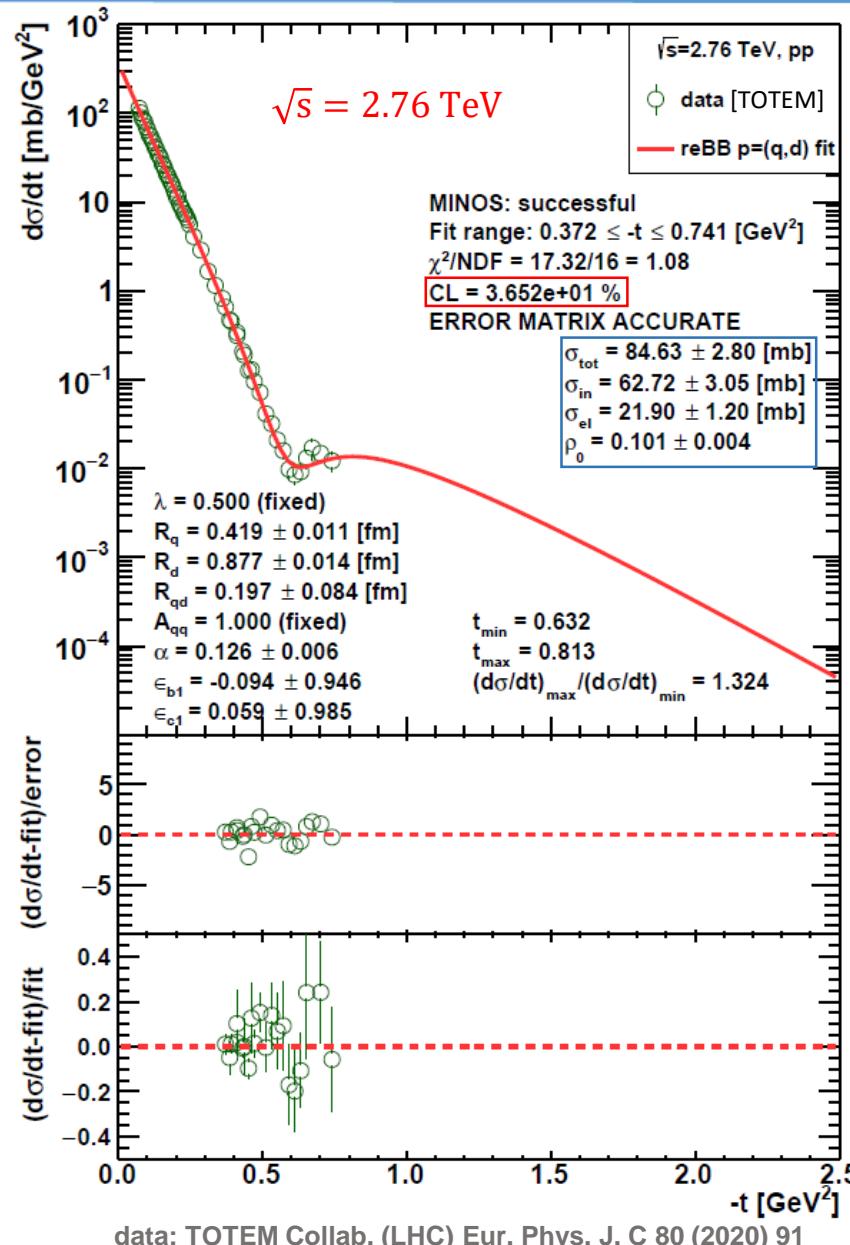
A. Adare *et al.* (PHENIX Collab.)
Phys. Rev. C 77, 064907

- it takes into account (in M separately measured t ranges):
 - the t -dependent statistical (a) and systematic (b) errors (both vertical σ_k and horizontal $\delta_k t$) $\rightarrow \epsilon_b$ parameters;
 - the t -independent σ_c normalization uncertainties $\rightarrow \epsilon_c$ parameters;
 - the measured total cross-section $d_{\sigma_{tot}}$ and ratio d_{ρ_0} and their total uncertainties $\delta \sigma_{tot}$ and $\delta \rho_0$.
- minimization with **CERN Root MINUIT**, parameter error estimation by **MINOS**.

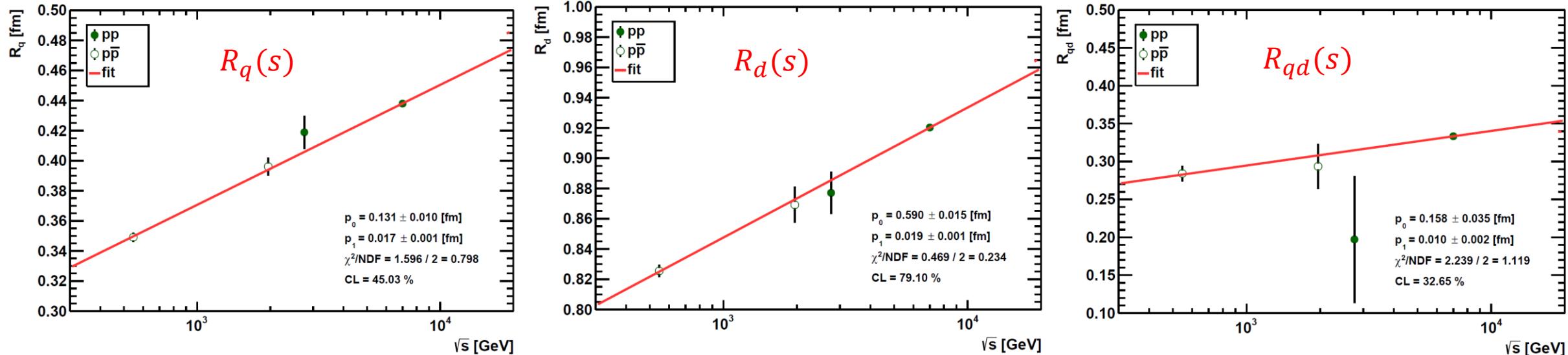
Satisfactory ReBB model fits for $p\bar{p}$ $d\sigma/dt$ data



Satisfactory ReBB model fits for pp $d\sigma/dt$ data



Energy dependences of the scale parameters



The energy dependences of the scale parameters, R_q , R_d and R_{qd} , are the same for pp and p \bar{p} processes!

$$P(s) = p_0 + p_1 \ln(s/s_0)$$

$$P \in \{R_q, R_d, R_{qd}, \alpha\}$$

$$s_0 = 1 \text{ GeV}^2$$

Parameter	R_q [fm]	R_d [fm]	R_{qd} [fm]
χ^2/NDF	1.596/2	0.469/2	2.239/2
CL [%]	45.03	79.10	32.65
p_0	0.131 ± 0.010	0.590 ± 0.015	0.158 ± 0.035
p_1	0.017 ± 0.001	0.019 ± 0.001	0.010 ± 0.002

Parameters which define the energy dependence
of the ReBB model scale parameters

New: proportionality between $\rho_0(s)$ and $\alpha(s)$

$$t_{el}(s, b) = i \left(1 - e^{i \alpha \tilde{\sigma}_{in}(s, b)} \sqrt{1 - \tilde{\sigma}_{in}(s, b)} \right)$$

$$\alpha \tilde{\sigma}_{in} \ll 1$$

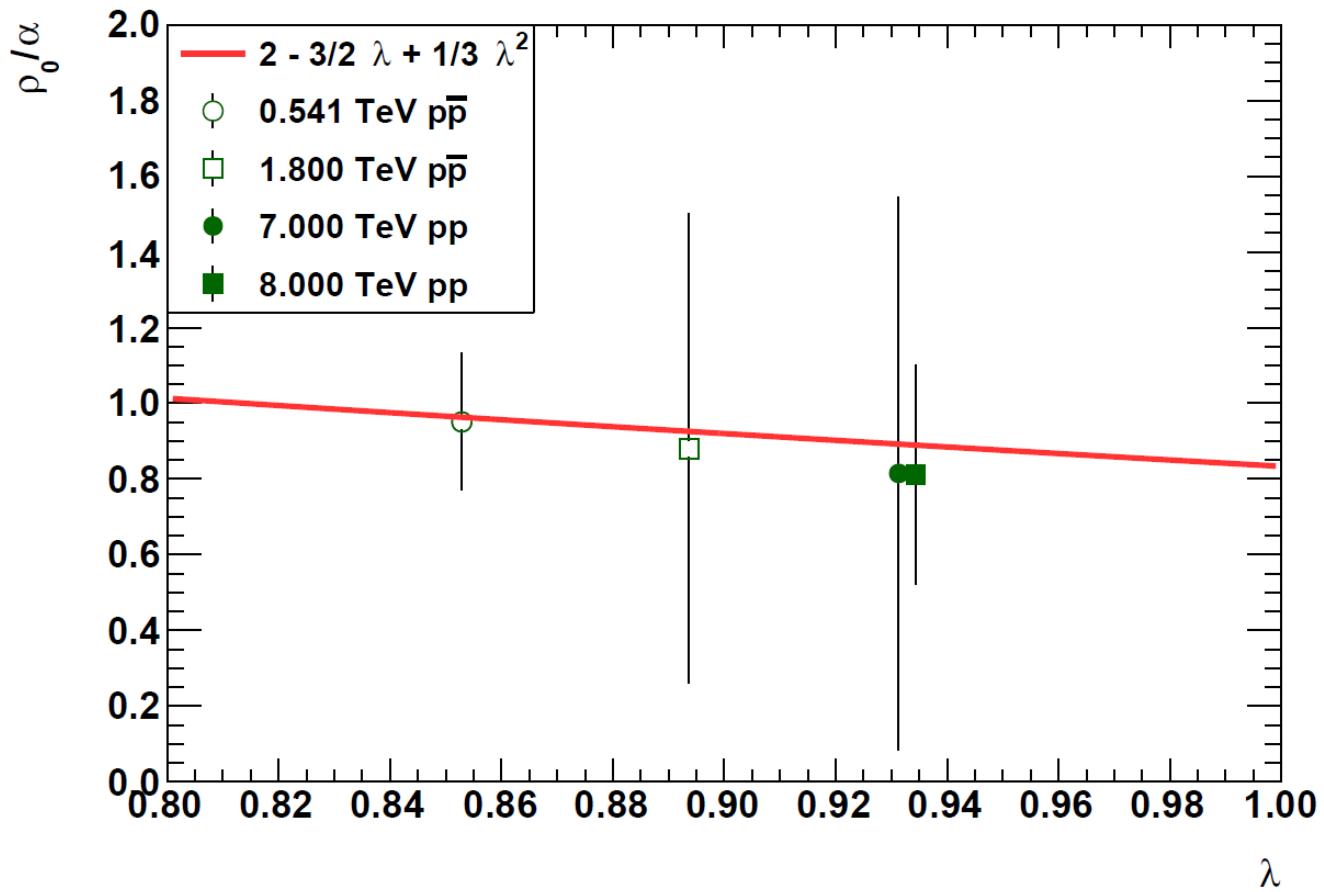
$$\text{Im } t_{el}(s, b) \simeq \lambda(s) \exp \left(-\frac{b^2}{2R^2(s)} \right)$$



$$\rho_0(s) = \alpha(s) \left(2 - \frac{3}{2} \lambda(s) + \frac{1}{3} \lambda^2(s) \right)$$

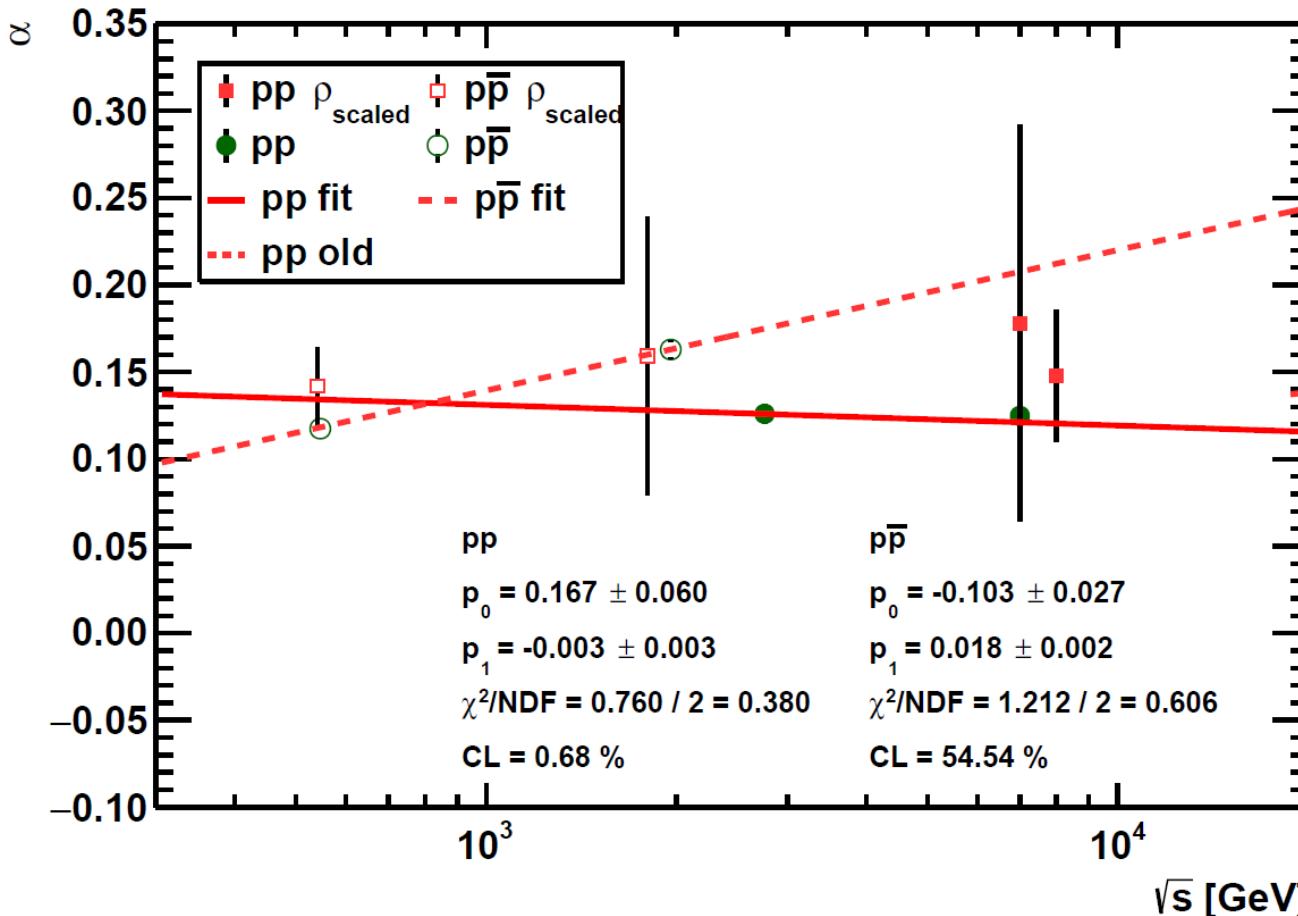
$$\lambda(s) = \text{Im } t_{el}(s, b = 0)$$

→ by rescaling one can get additional α parameter values at energies where ρ_0 is measured



The dependence of ρ_0/α on $\lambda = \text{Im } t_{el}(s, b = 0)$ in the TeV energy range. The data points are generated numerically by using the trends of the ReBB model scale parameters and the experimentally measured ρ -parameter values.

Energy dependence of the α parameter



$$P(s) = p_0 + p_1 \ln(s/s_0)$$

$$P \in \{R_q, R_d, R_{qd}, \alpha\}$$

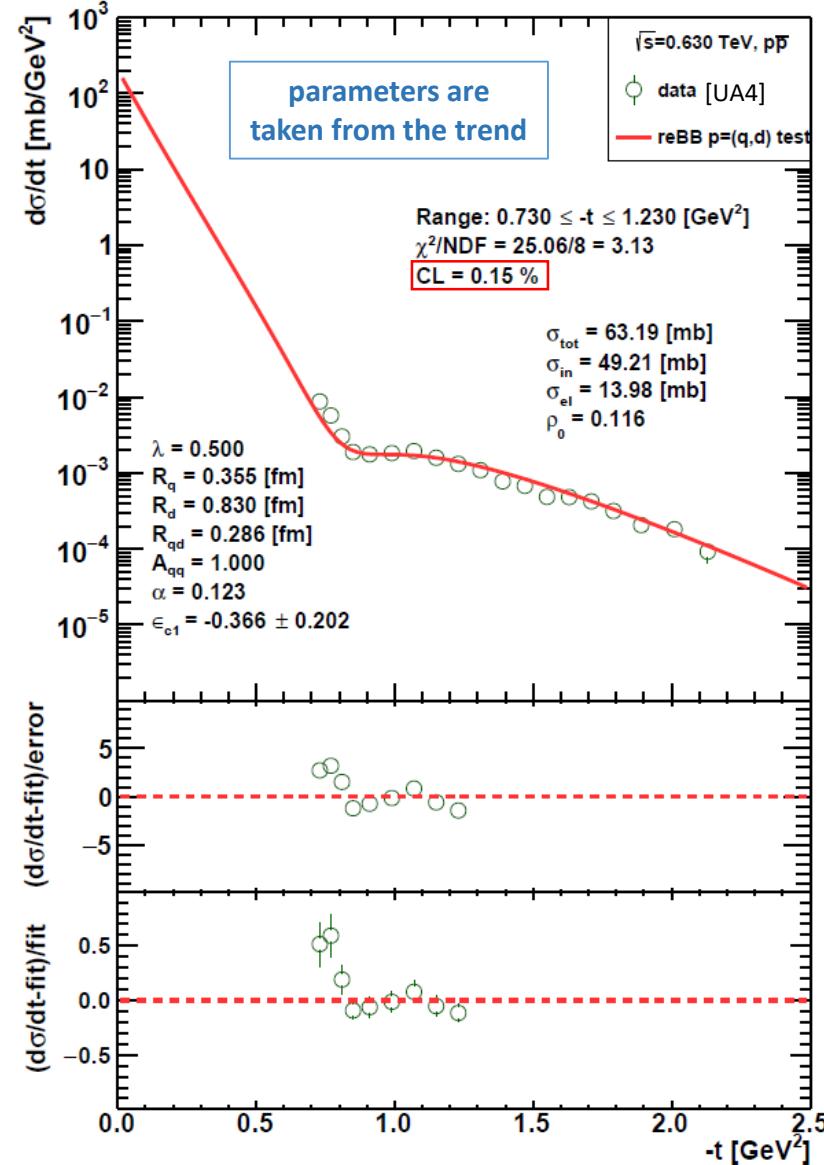
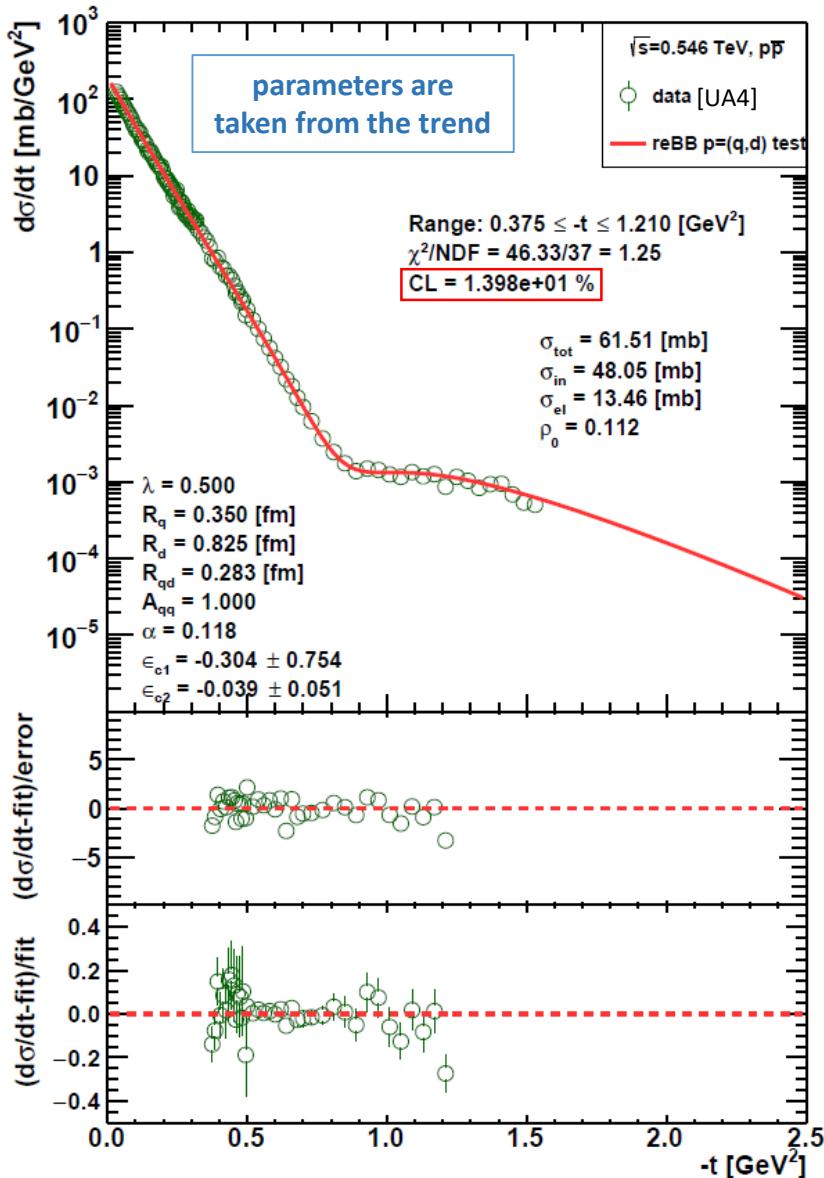
$$s_0 = 1 \text{ GeV}^2$$

Parameter	$\alpha (pp)$	$\alpha (p\bar{p})$
χ^2/NDF	$0.760/2$	$1.212/2$
CL [%]	0.68	54.54
p_0	0.167 ± 0.060	-0.103 ± 0.027
p_1	-0.003 ± 0.003	0.018 ± 0.002

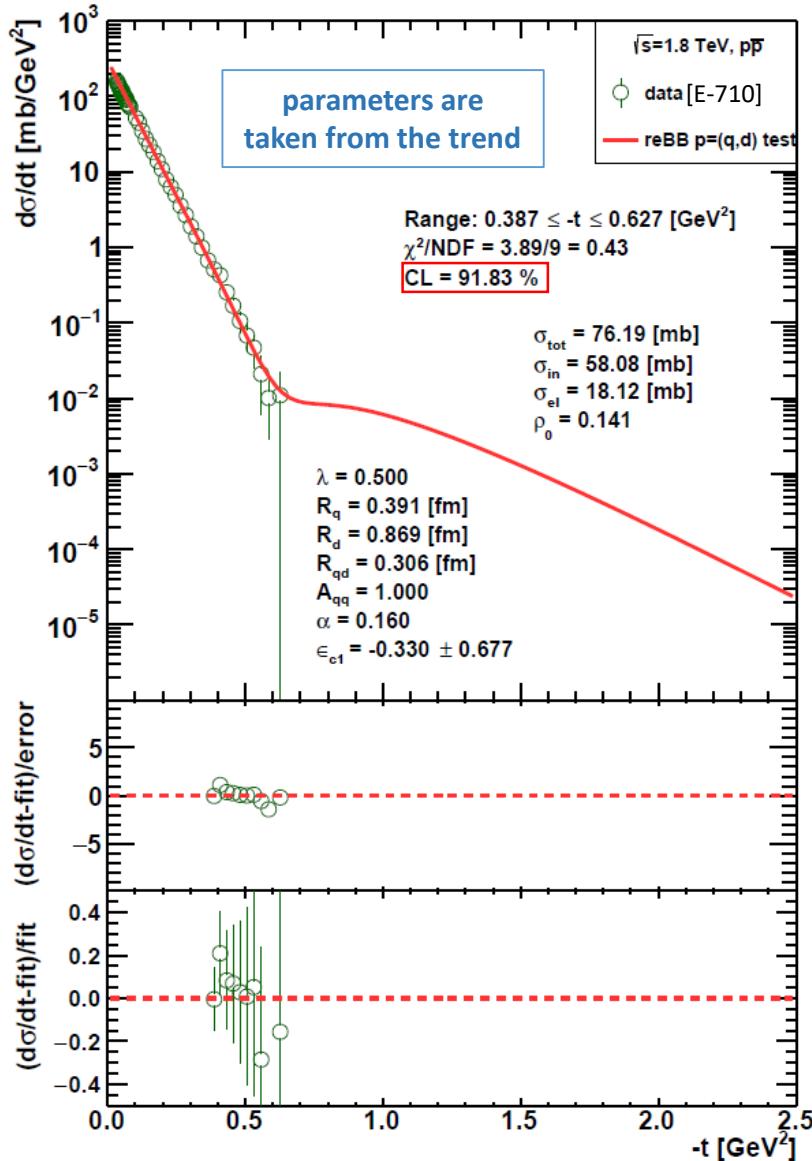
Parameters which define the energy dependence of the ReBB model α parameters for pp and $p\bar{p}$

The energy dependence of the α parameter is not the same for pp and $p\bar{p}$ processes
 → the Odderon is characterized by a single parameter, α !

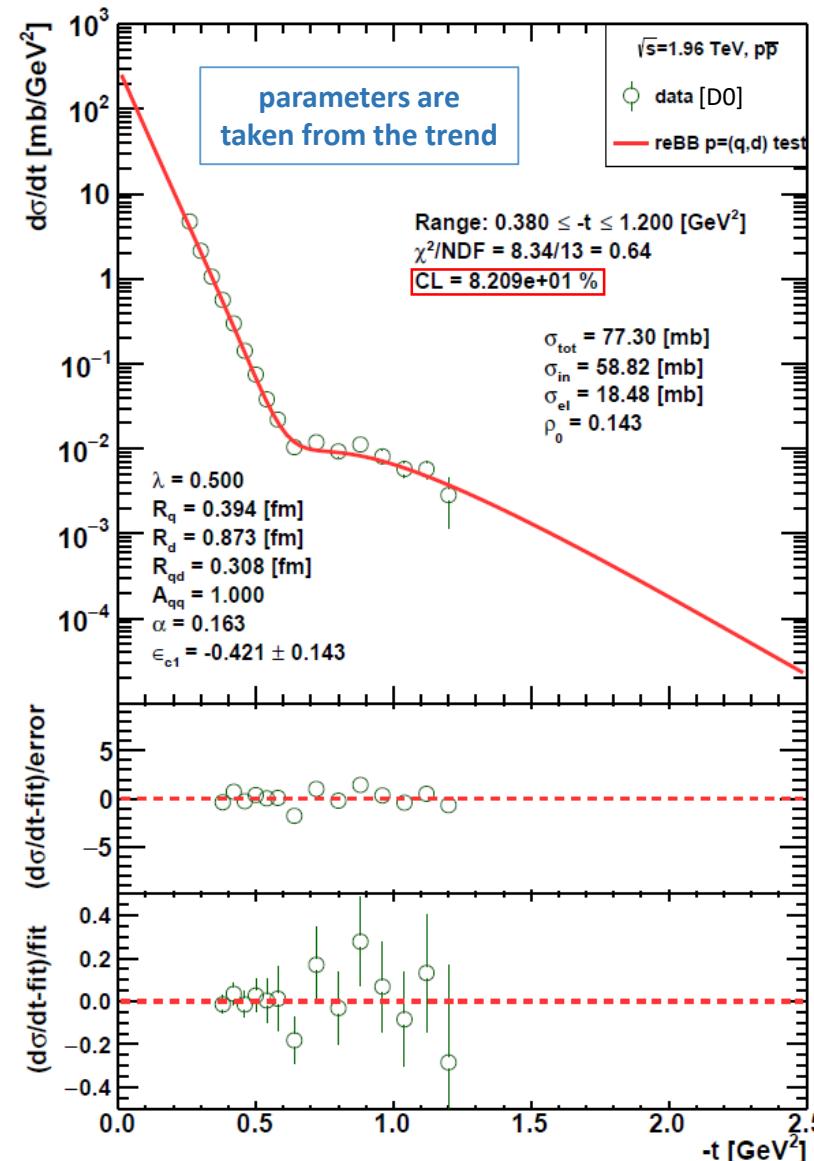
Tests @ 0.546 & 0.630 TeV ✓



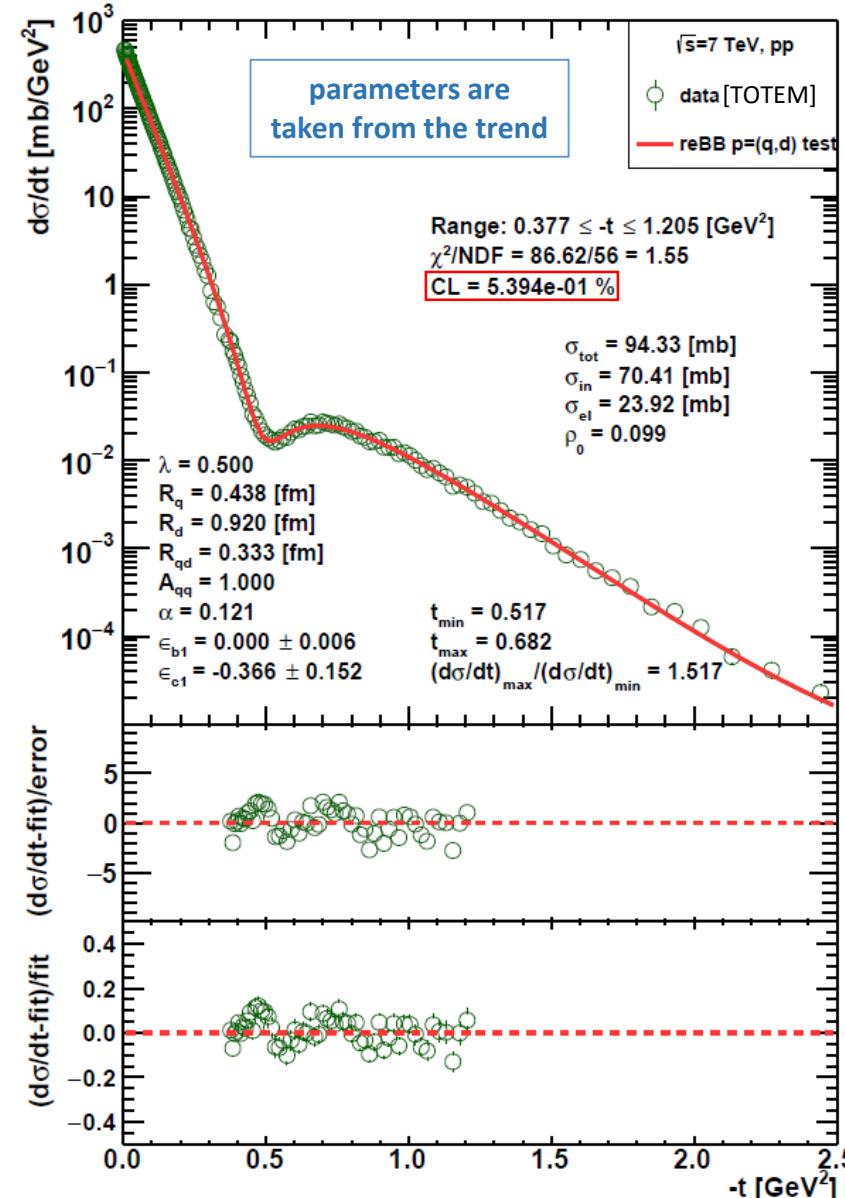
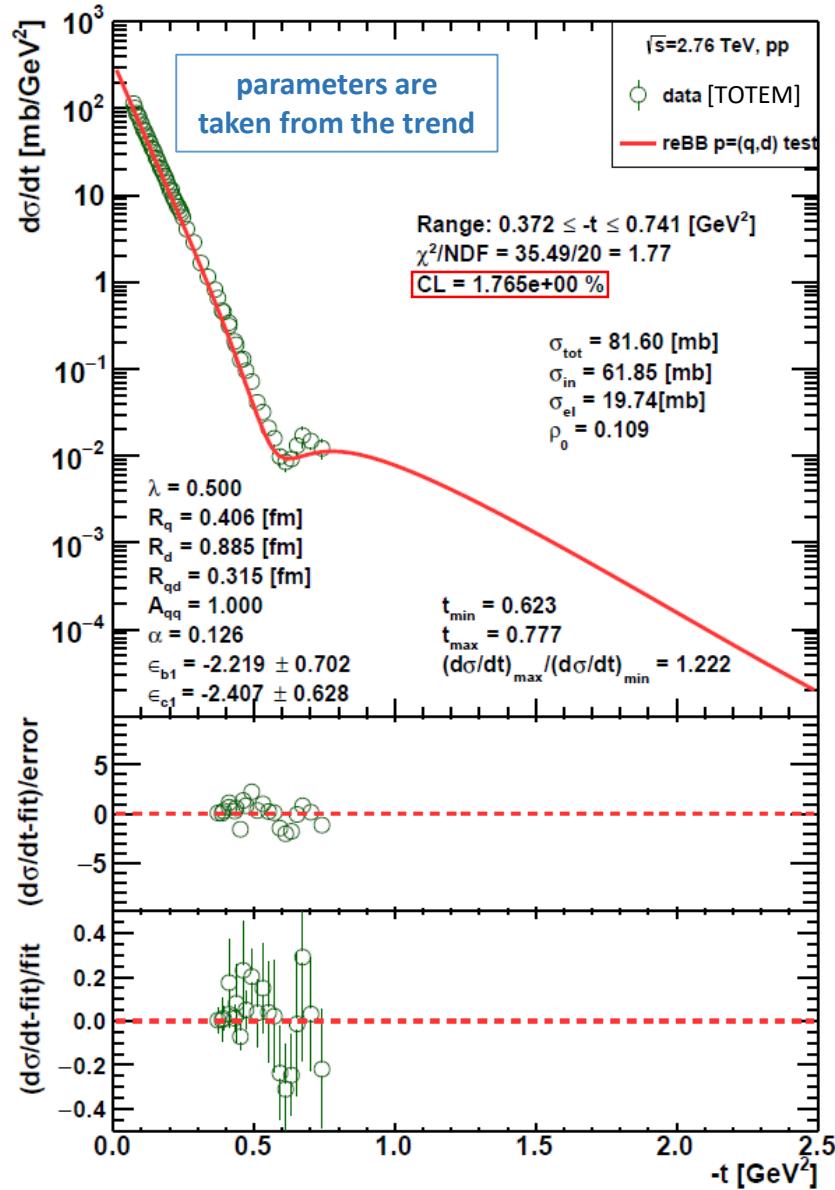
Tests @ 1.8 & 1.96 TeV



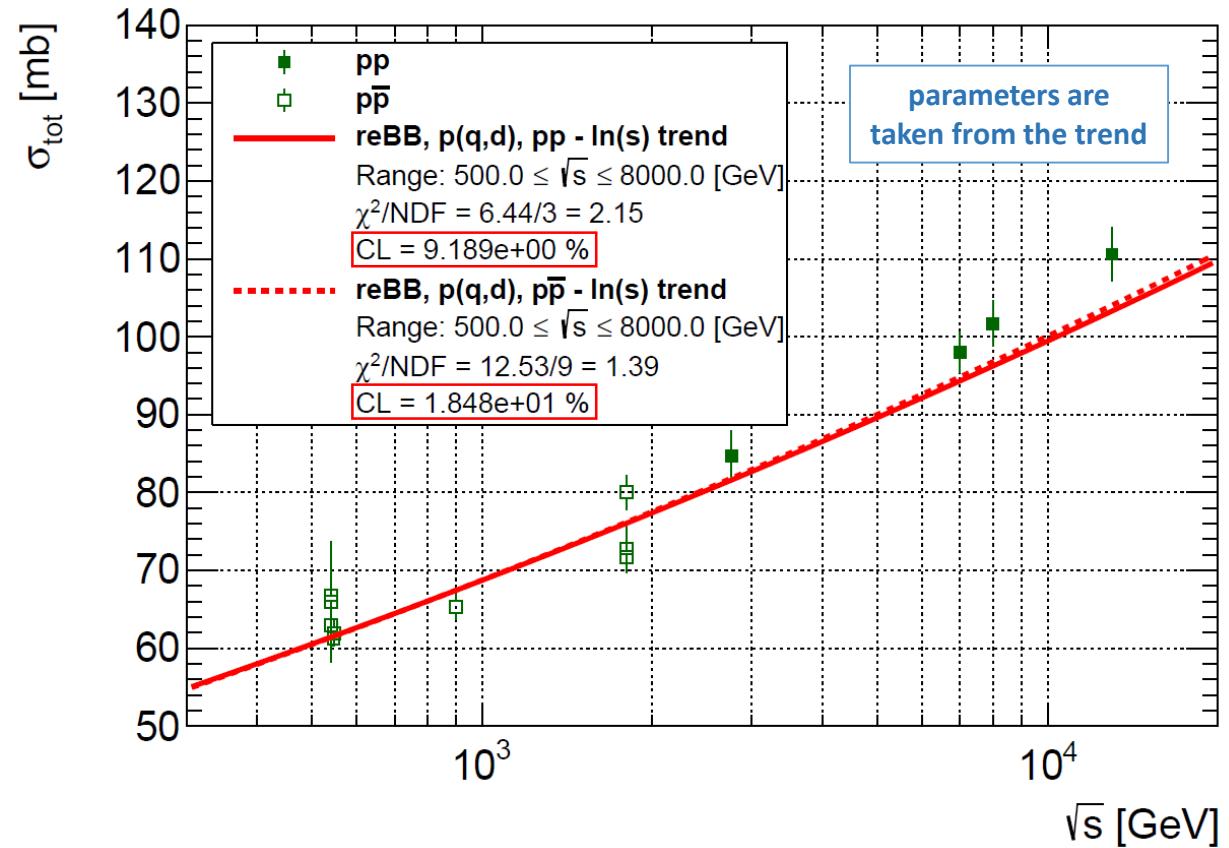
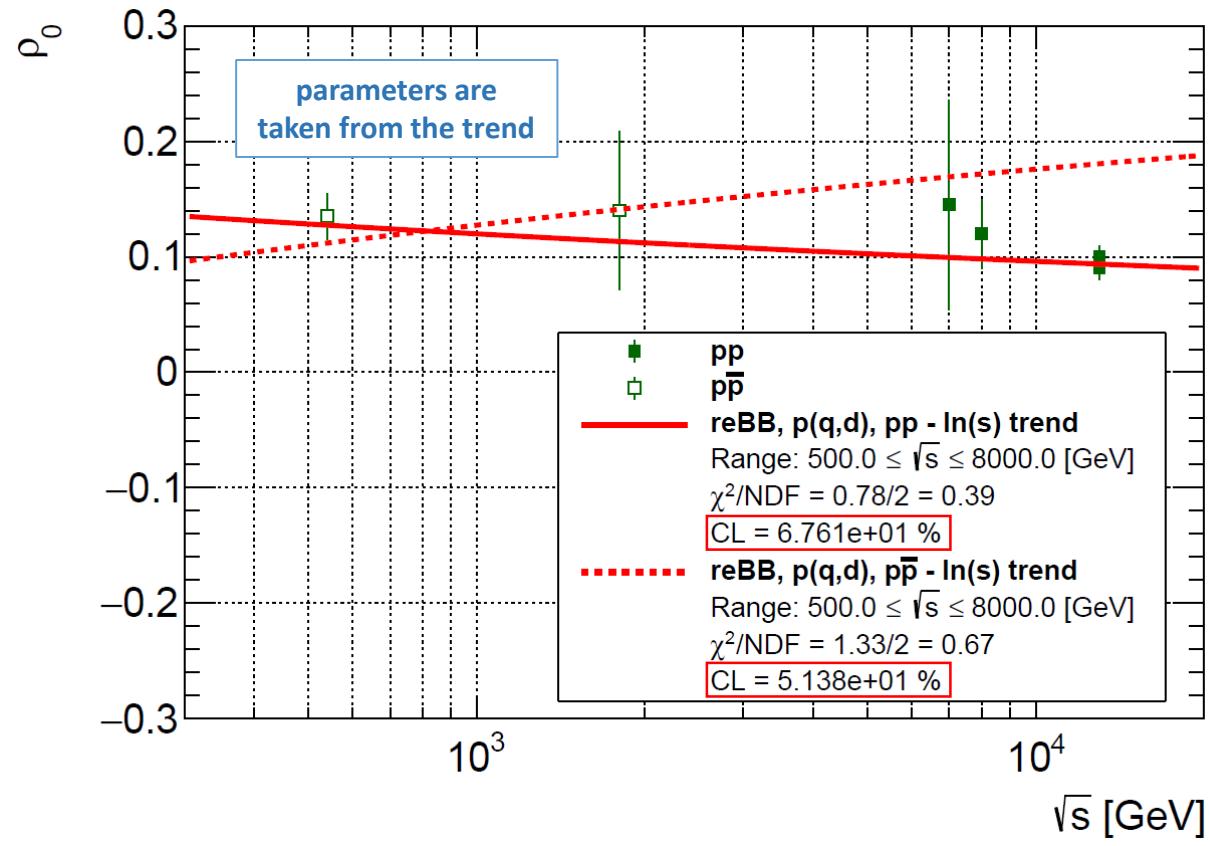
data: E-710 Collab. (TEVATRON) Phys.Lett. B247 (1990)



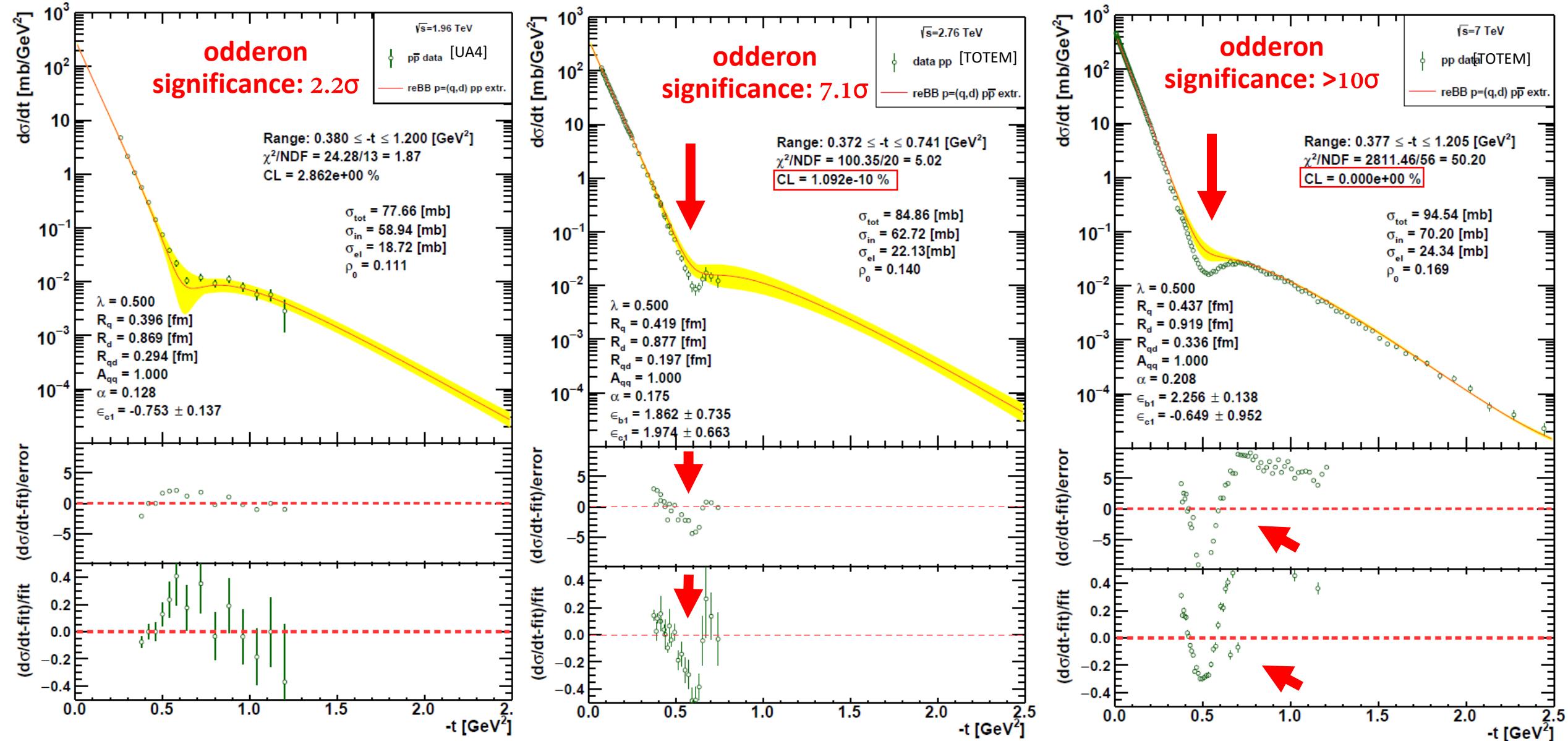
Tests @ 2.76 & 7.0 TeV ✓



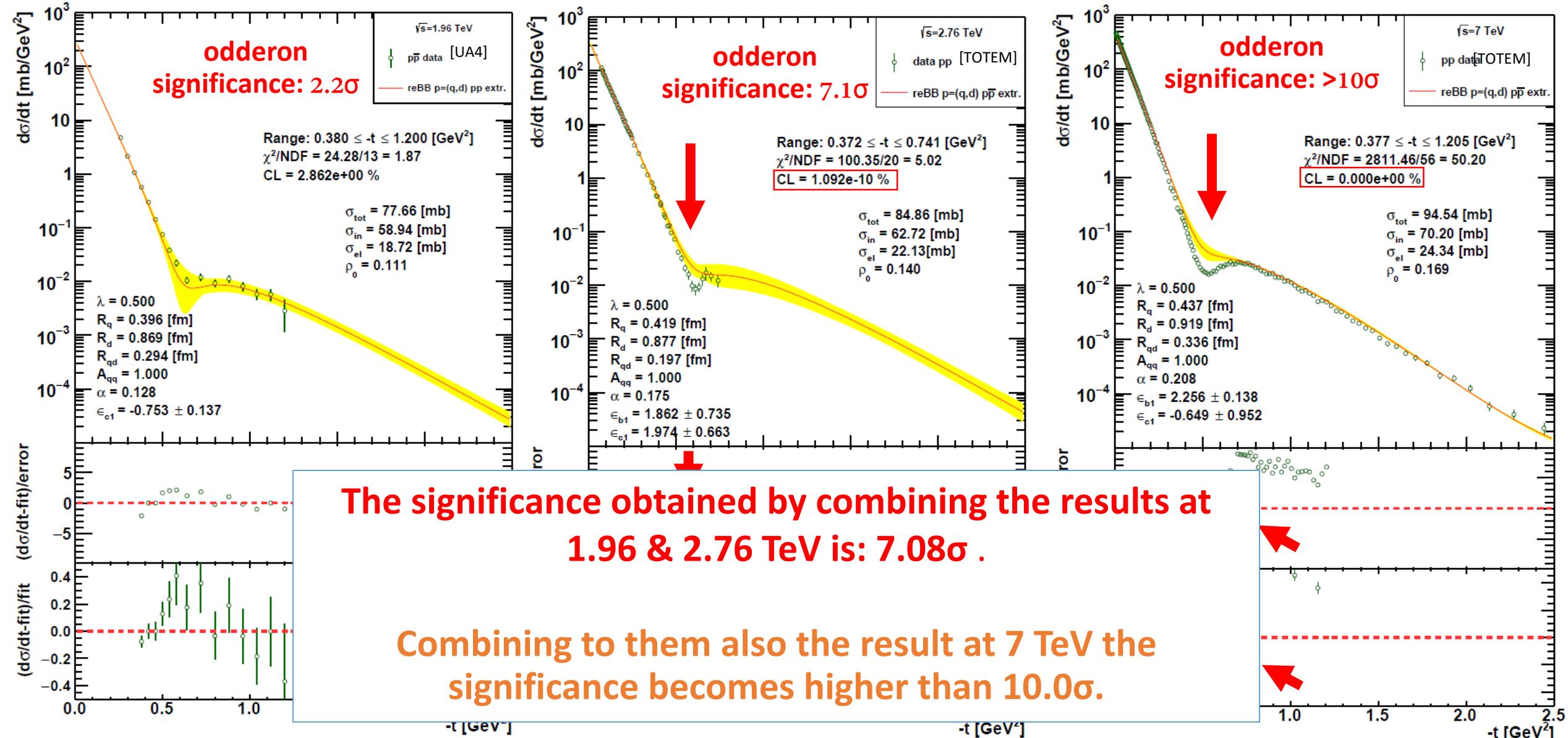
Tests for σ_{tot} and ρ_0 ✓



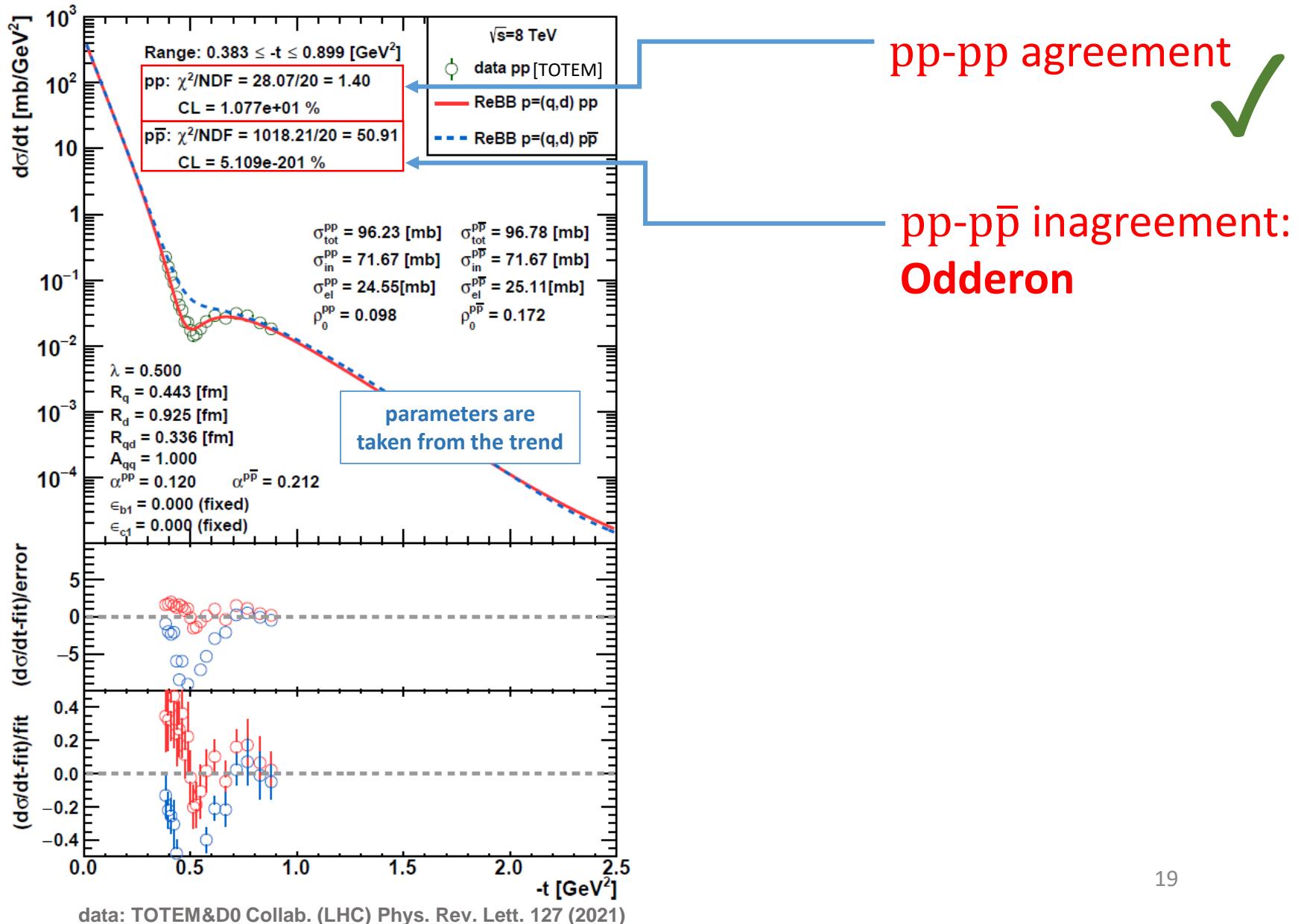
Extrapolations → ODDERON



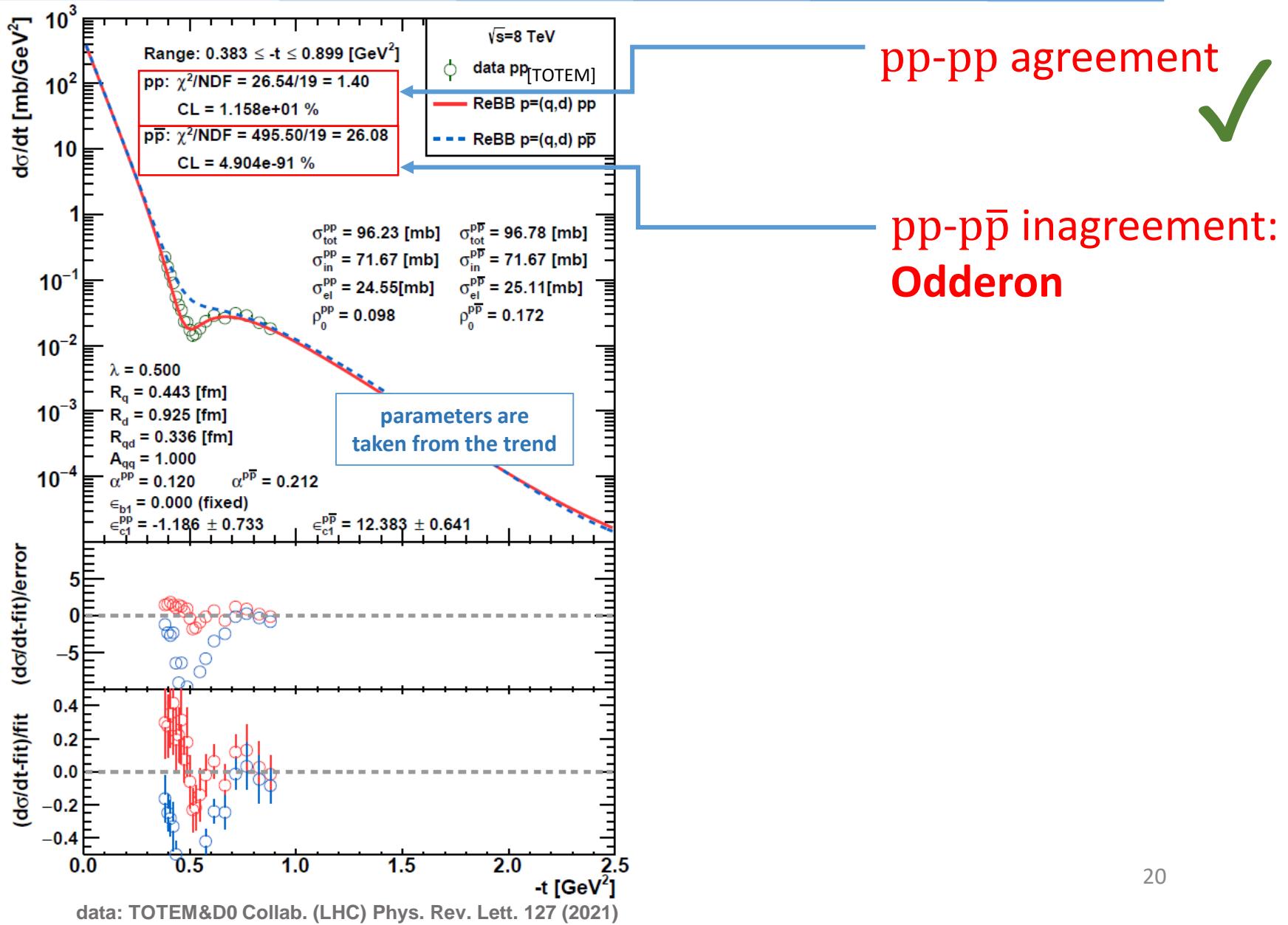
Extrapolations → ODDERON



Confirmation @ 8.0 TeV



Confirmation @ 8.0 TeV (with 3% C-type error)



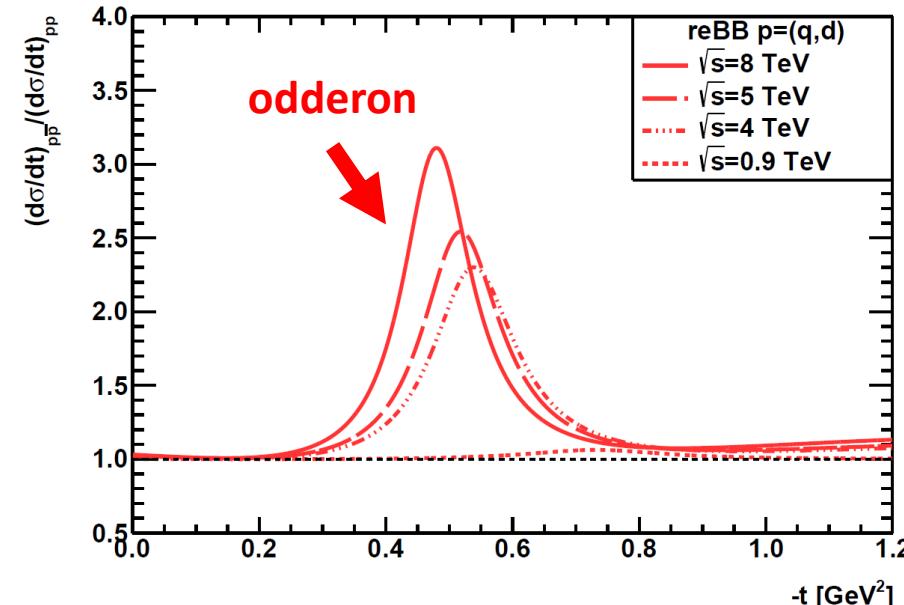
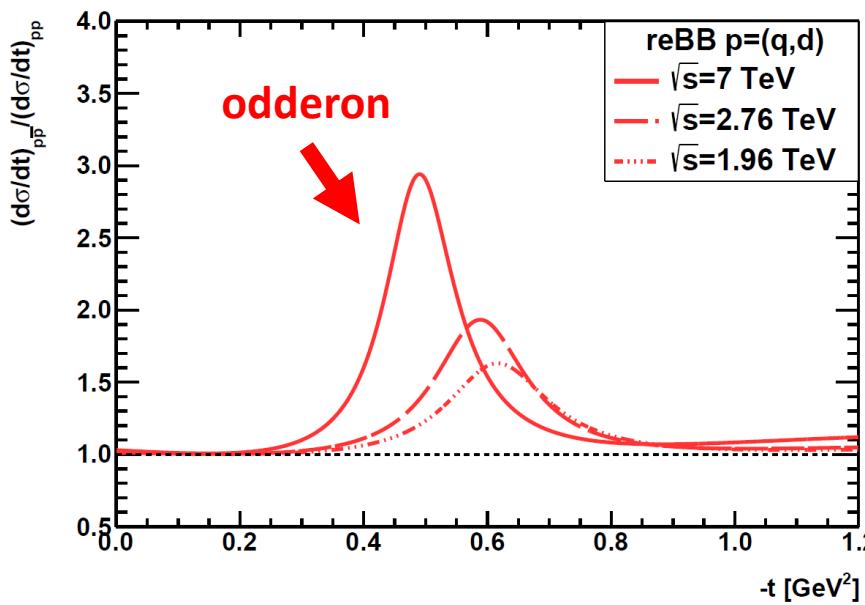
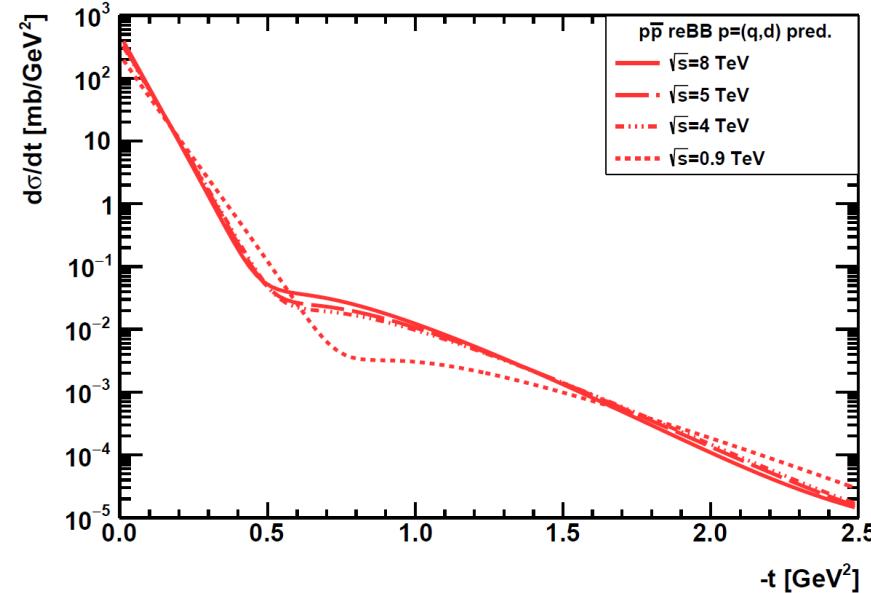
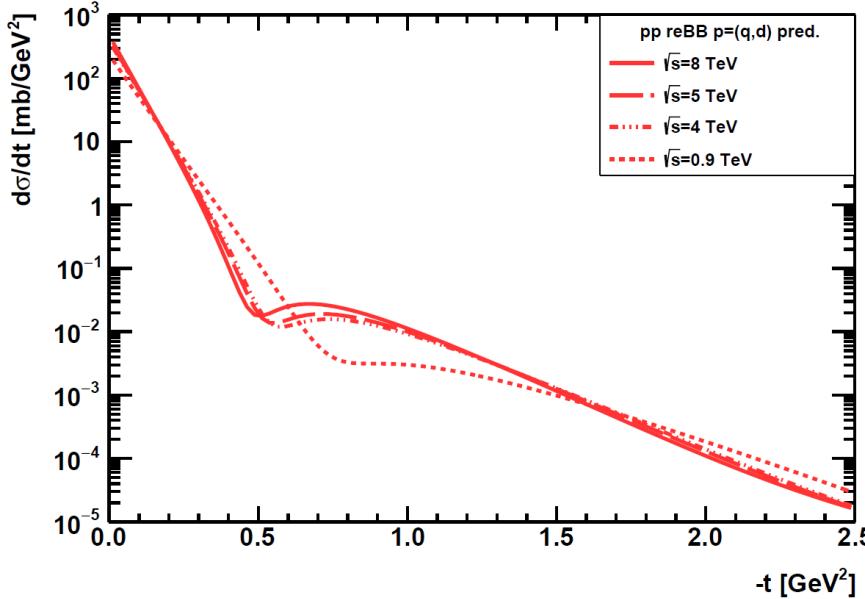
Summary

- **ReBB model fits to pp and p \bar{p} d σ /dt data**
 - satisfactory description in the energy range of $0.546 \leq \sqrt{s} \leq 7(8)$ TeV and squared momentum transfer range of $0.37 \leq -t \leq 1.2$ GeV 2
- **determination of the energy dependence of the parameters**
 - $R_q(s)$, $R_d(s)$ and $R_{qd}(s)$ are the same for pp and p \bar{p} processes, $\alpha(s)$ is not
 - **Odderon effect**
 - lack of pp and p \bar{p} d σ /dt measured data at the same energies at the TeV energy region
 - **extrapolations to accomplish comparative study and obtain a significance for the Odderon effect**
 - **difference between pp and p \bar{p} d σ /dt in the dip region**
 - **model-dependent evidence for Odderon exchange in t-channel at $\sqrt{s} \gtrsim 1$ TeV with a significance of at least 7.08σ**

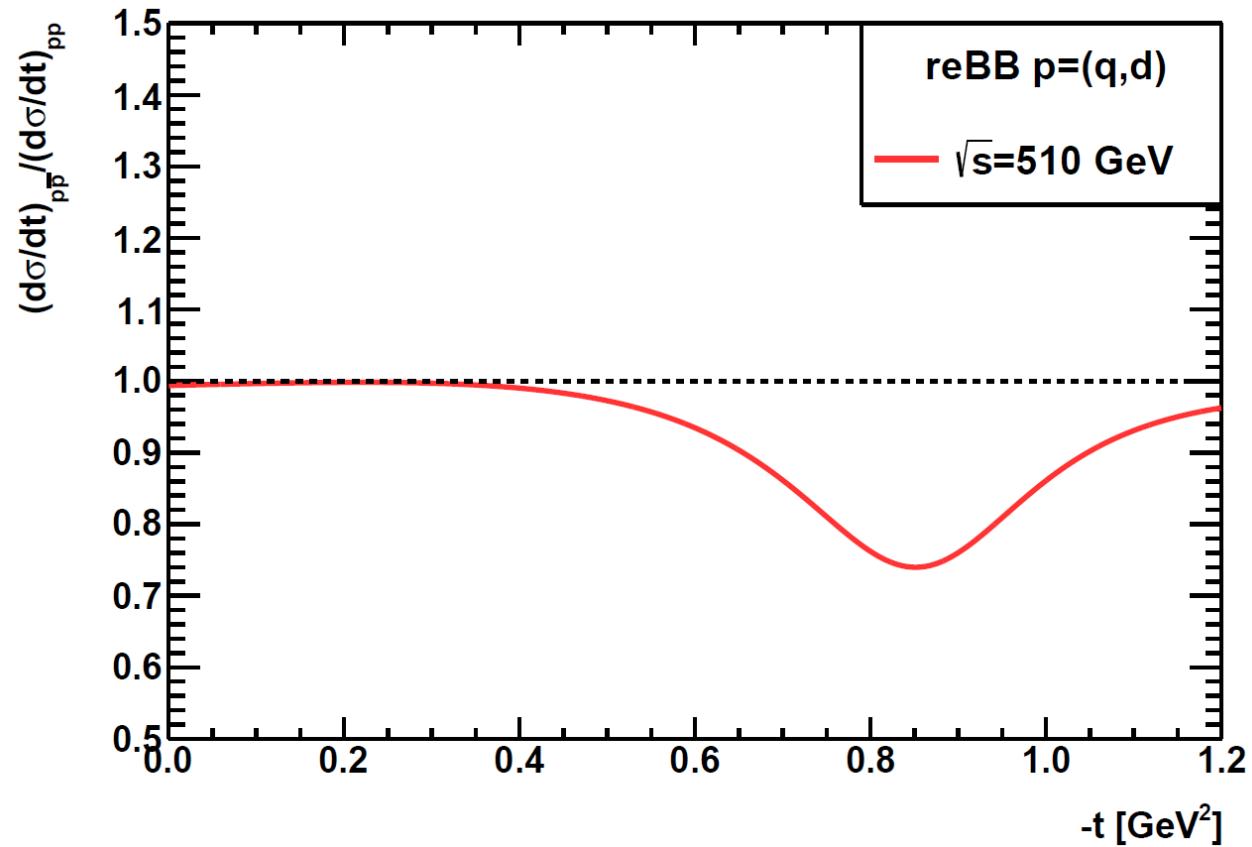
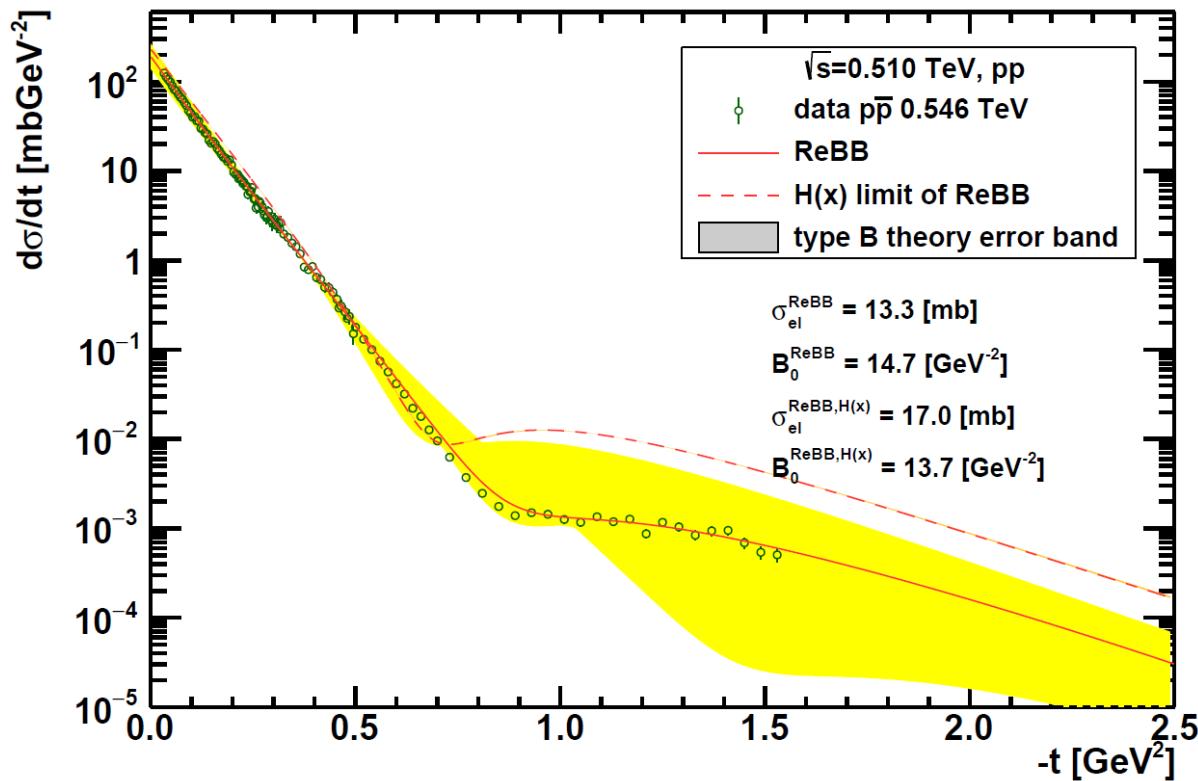
Thank you for your attention!

Backup slides

Predictions for pp and p \bar{p} $d\sigma/dt$ and their ratios

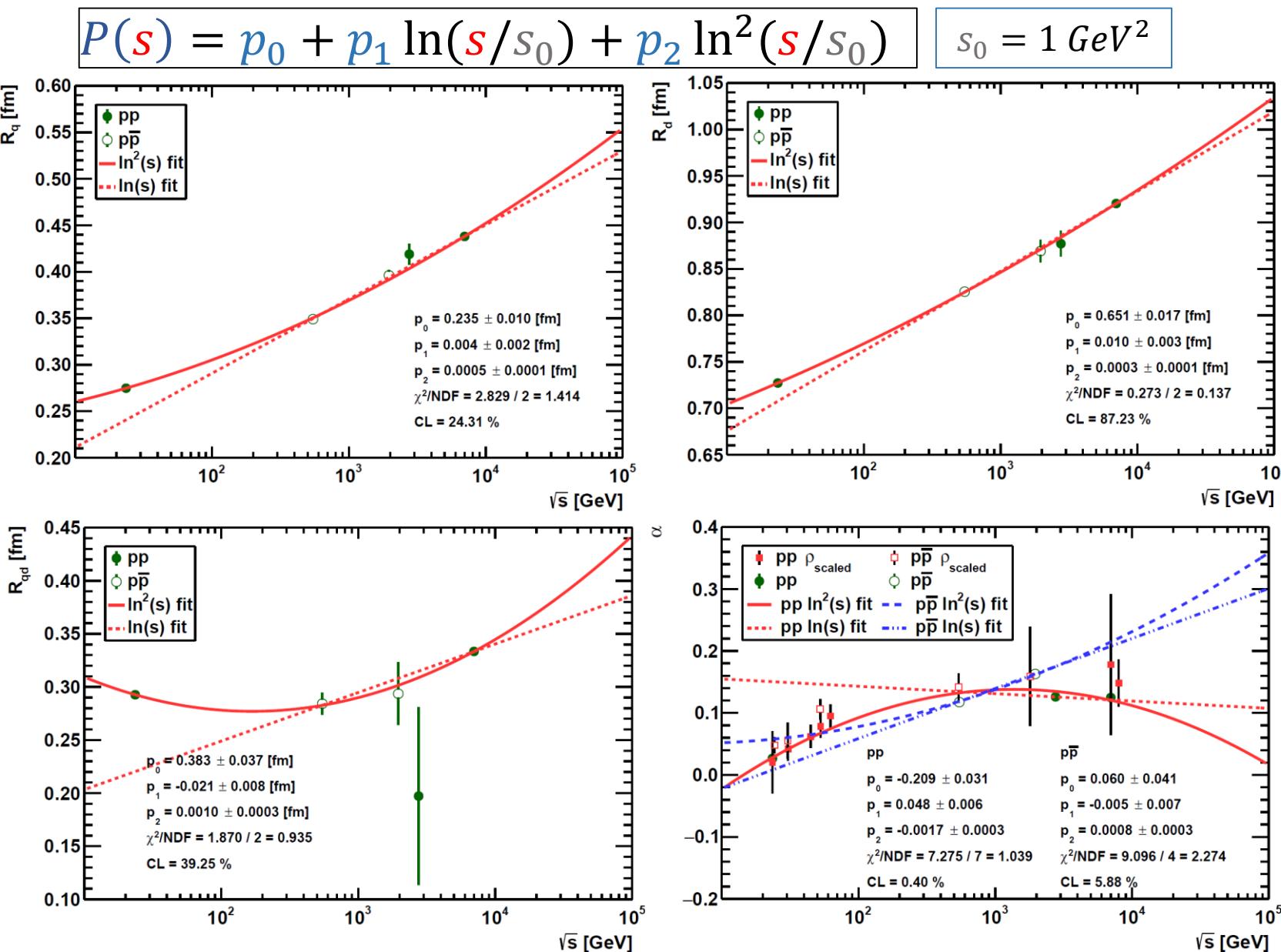
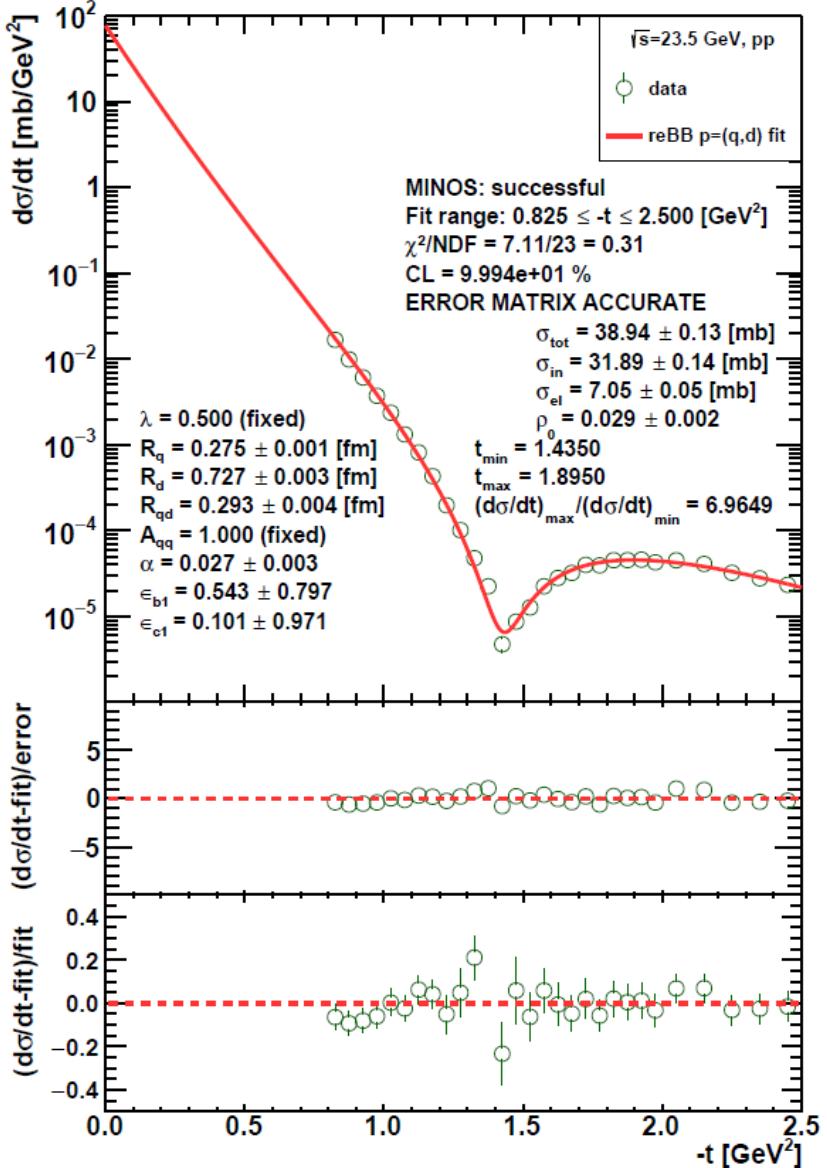


ReBB prediction for pp $d\sigma/dt$ @ 0.51 TeV



- pp data measured by the STAR Collab. is expected at 0.51 TeV

Fit at pp 23.5 GeV & $\ln^2(s)$ energy dependence



Various significance combinations

