

N_{eff} constraints on portal interactions with hidden sectors

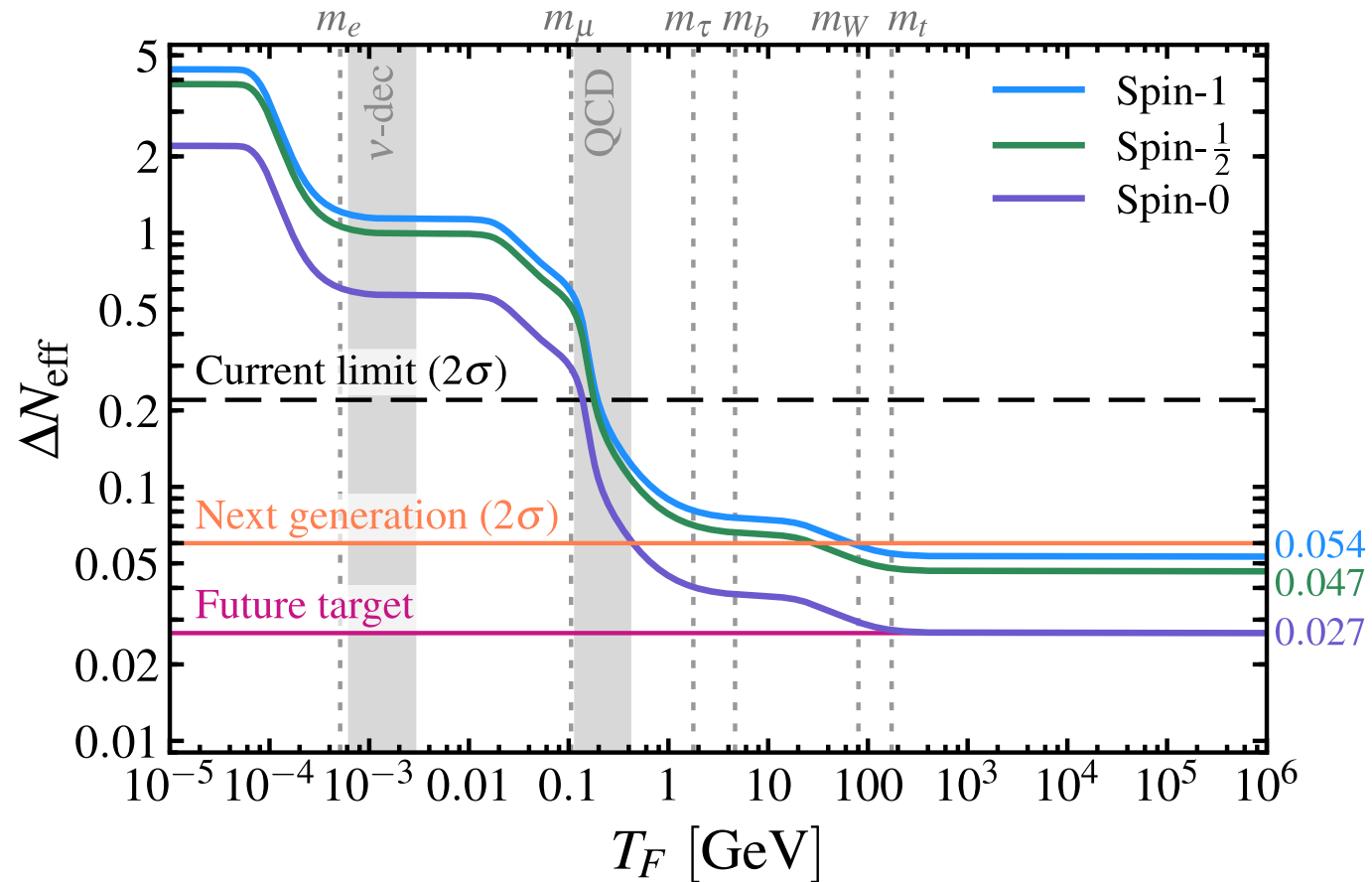
arxiv:2205.xxxxx

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University of Illinois at Urbana-Champaign

Collaborators: Peter Adshead and
Jessie Shelton

ΔN_{eff} : Typically discussed as constraint on decoupling temperature

$$\Delta N_{\text{eff}} = \frac{8}{7} \left(\frac{11}{4} \right)^{4/3} \frac{\rho_{\text{dr}}}{\rho_{\gamma}}$$

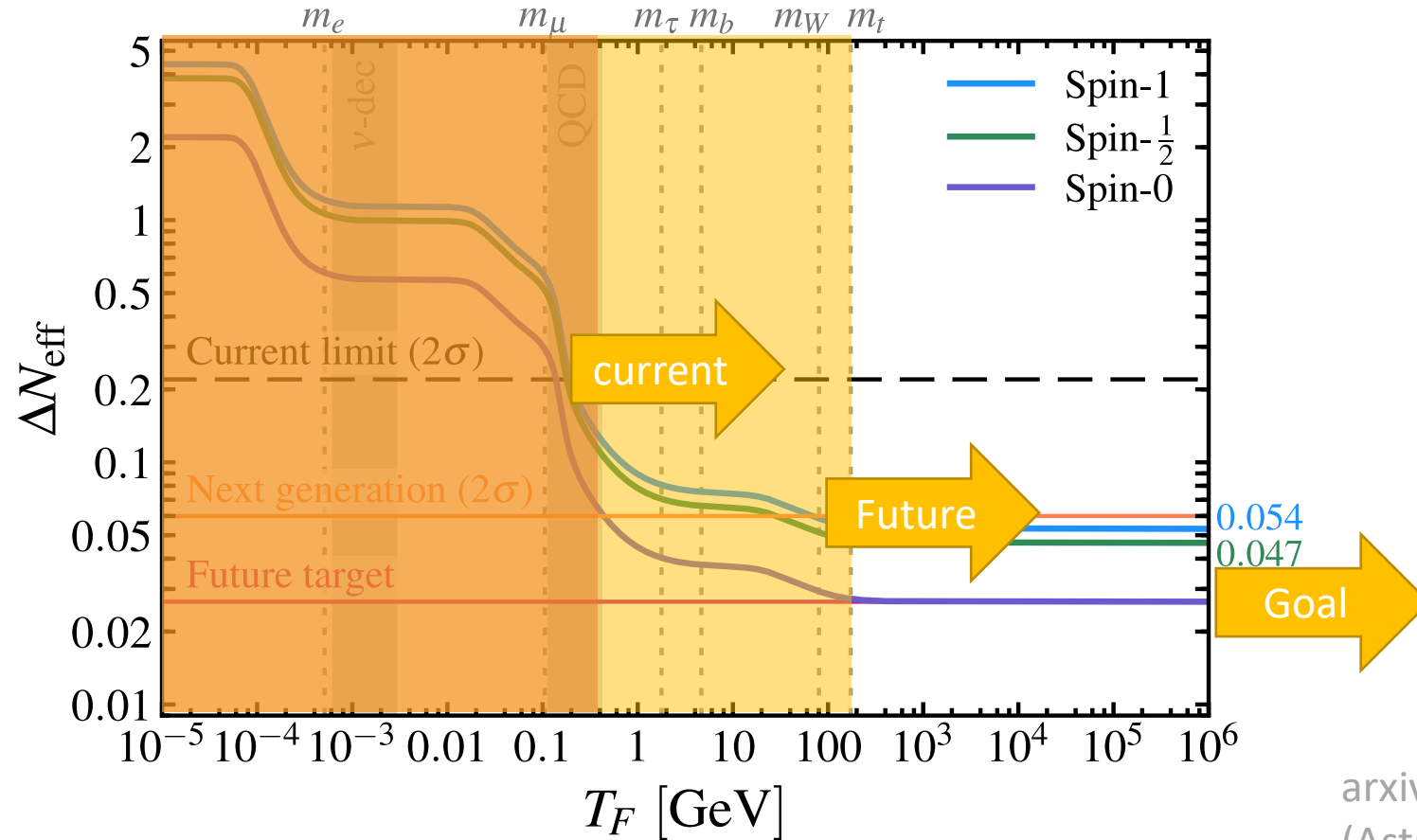


Decoupling temperature of new light particle

arxiv:1903:04763
(Astro 2020 white paper)

Reinterpreting ΔN_{eff} : Constraint on interactions with out-of-equilibrium sectors

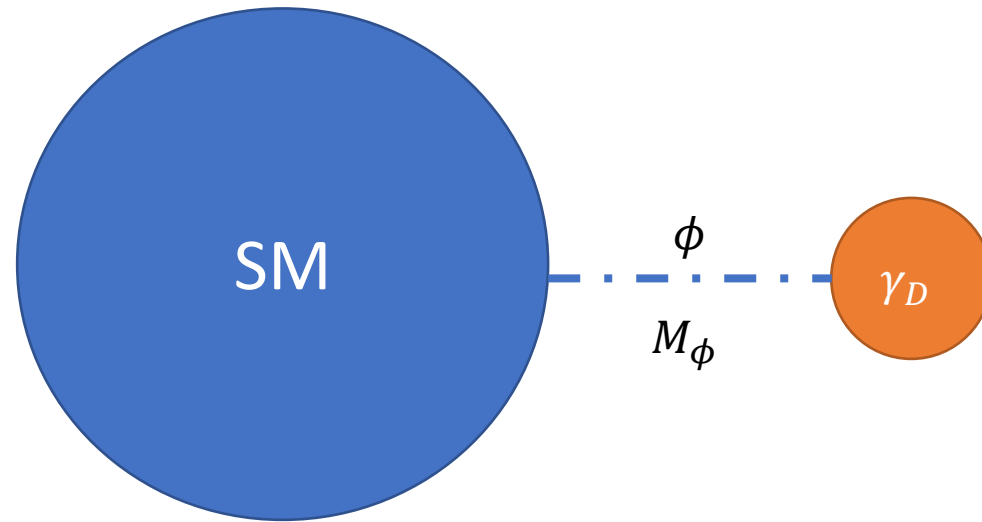
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Highlight implications for model building with
HS with dark radiation

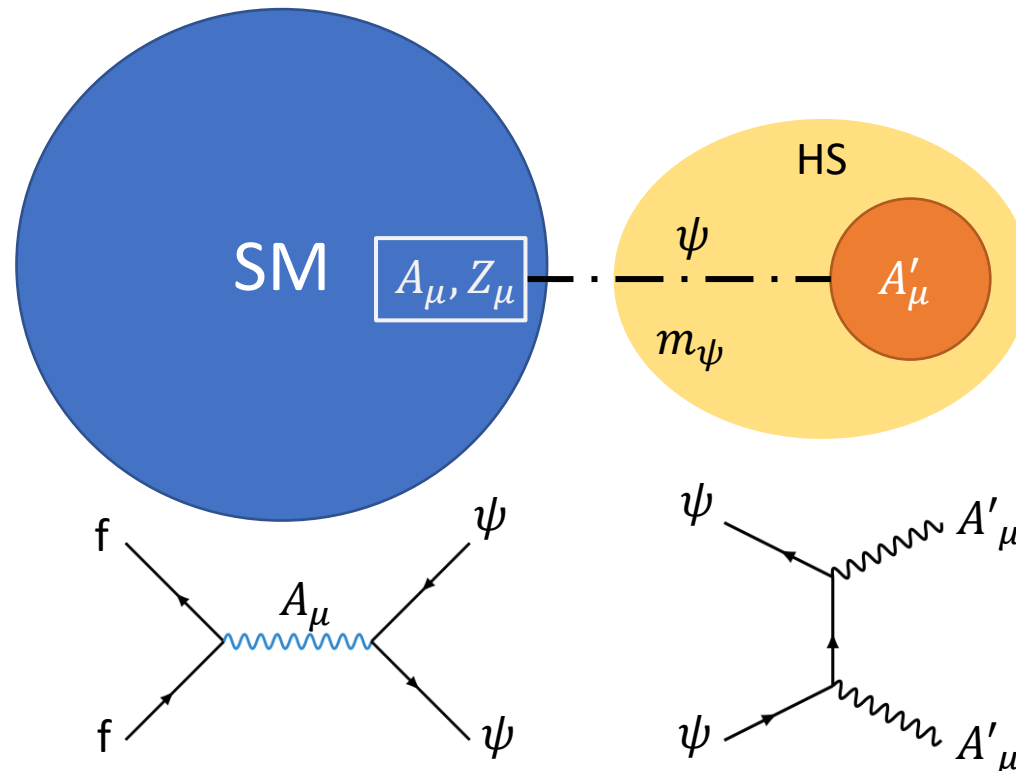
Application of N_{eff} constraint : Relevant types of interaction



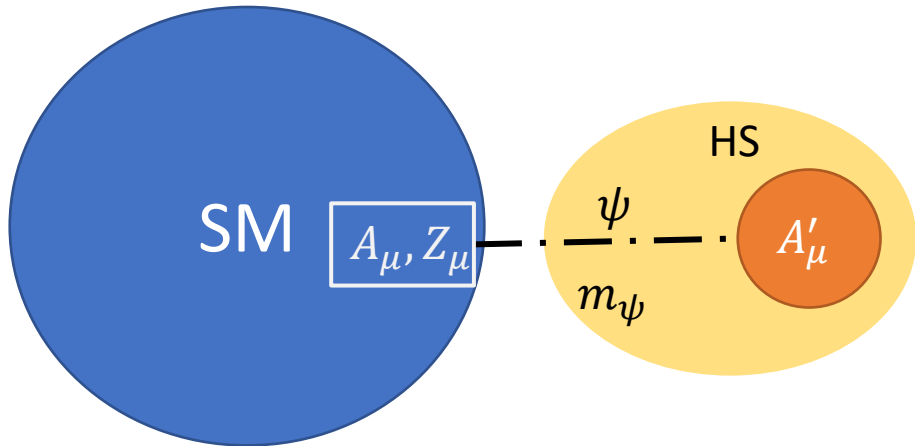
Application of N_{eff} constraint: Millicharged particle example

$$L_{\text{int}} \supset -\frac{\epsilon}{2} B_{\mu\nu} F^{\mu\nu'} + e' A'_\mu \bar{\psi} \gamma^\mu \psi - m_\psi \bar{\psi} \psi$$

$$L_{\text{int}} \supset -eQA_\mu \bar{\psi} \gamma^\mu \psi + e' A'_\mu \bar{\psi} \gamma^\mu \psi + eQZ_\mu \tan \theta_W \bar{\psi} \gamma^\mu \psi - m_\psi \bar{\psi} \psi$$



Physics behind dark radiation production: Boltzmann equations



Boltzmann equations:

$$\frac{d\rho_{SM}}{dt} + 3H(1 + w_{SM})\rho_{SM} = -C$$

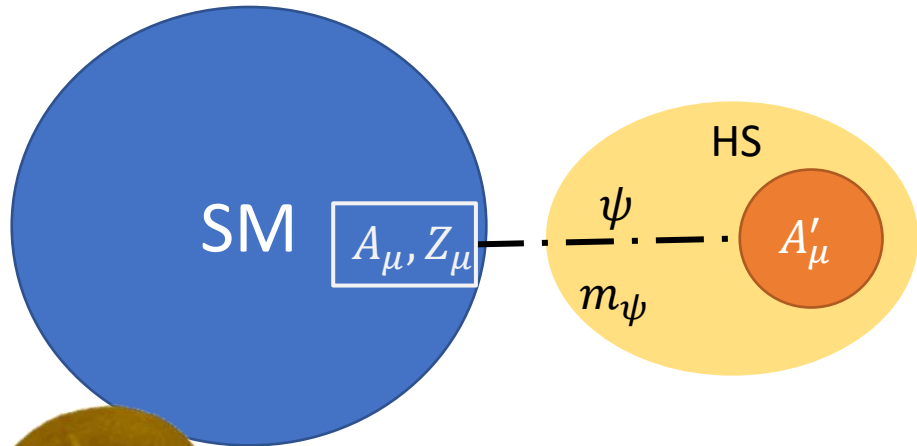
$$\frac{d\rho_{HS}}{dt} + 3H(1 + w_{HS})\rho_{HS} = C$$

$$H = \frac{\sqrt{\rho_{SM} + \rho_{HS}}}{\sqrt{3}M_{\text{Pl}}}$$

$$C = \frac{1}{32\pi^4} \sum_f \int ds (s - 4m_f^2) s \sigma_{ff \rightarrow \psi\psi} [T_{SM} G(\sqrt{s}/T_{SM}) - T_{HS} G(\sqrt{s}/T_{HS})] + \dots$$

↑
Energy transfer collision term

Physics behind dark radiation production: Boltzmann equations



Boltzmann equations:

$$\frac{d\rho_{SM}}{dt} + 3H(1 + w_{SM})\rho_{SM} = -C$$

$$\frac{d\rho_{HS}}{dt} + 3H(1 + w_{HS})\rho_{HS} = C$$

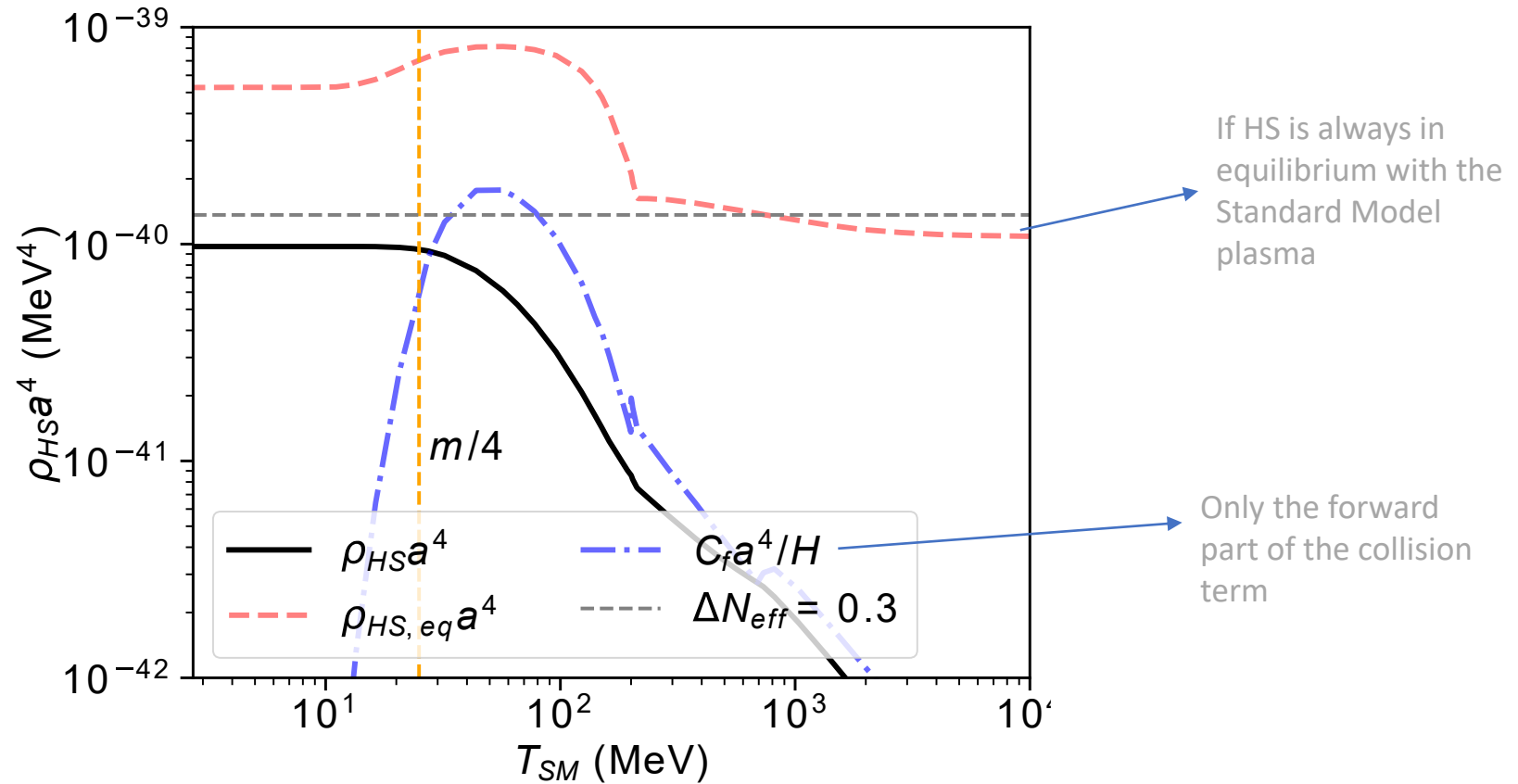
$$H = \frac{\sqrt{\rho_{SM} + \rho_{HS}}}{\sqrt{3}M_{Pl}}$$

$$C = \frac{1}{32\pi^4} \sum_f \int ds (s - 4m_f^2) s \sigma_{ff \rightarrow \psi\psi} [T_{SM} G(\sqrt{s}/T_{SM}) - T_{HS} G(\sqrt{s}/T_{HS})] + \dots$$

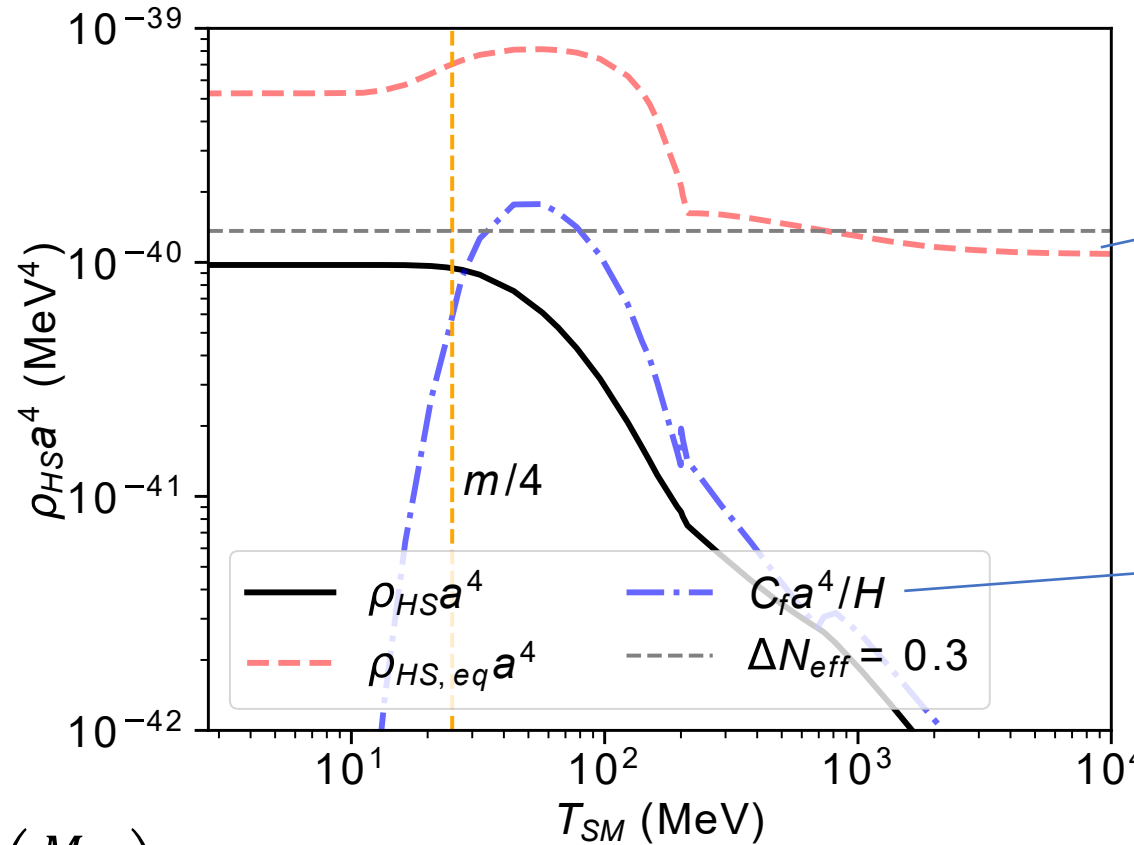
↑
Energy transfer collision term



Physics behind dark radiation production: Plots!



Physics behind dark radiation production: Out-of-equilibrium ρ_{DR} proportional to portal coupling



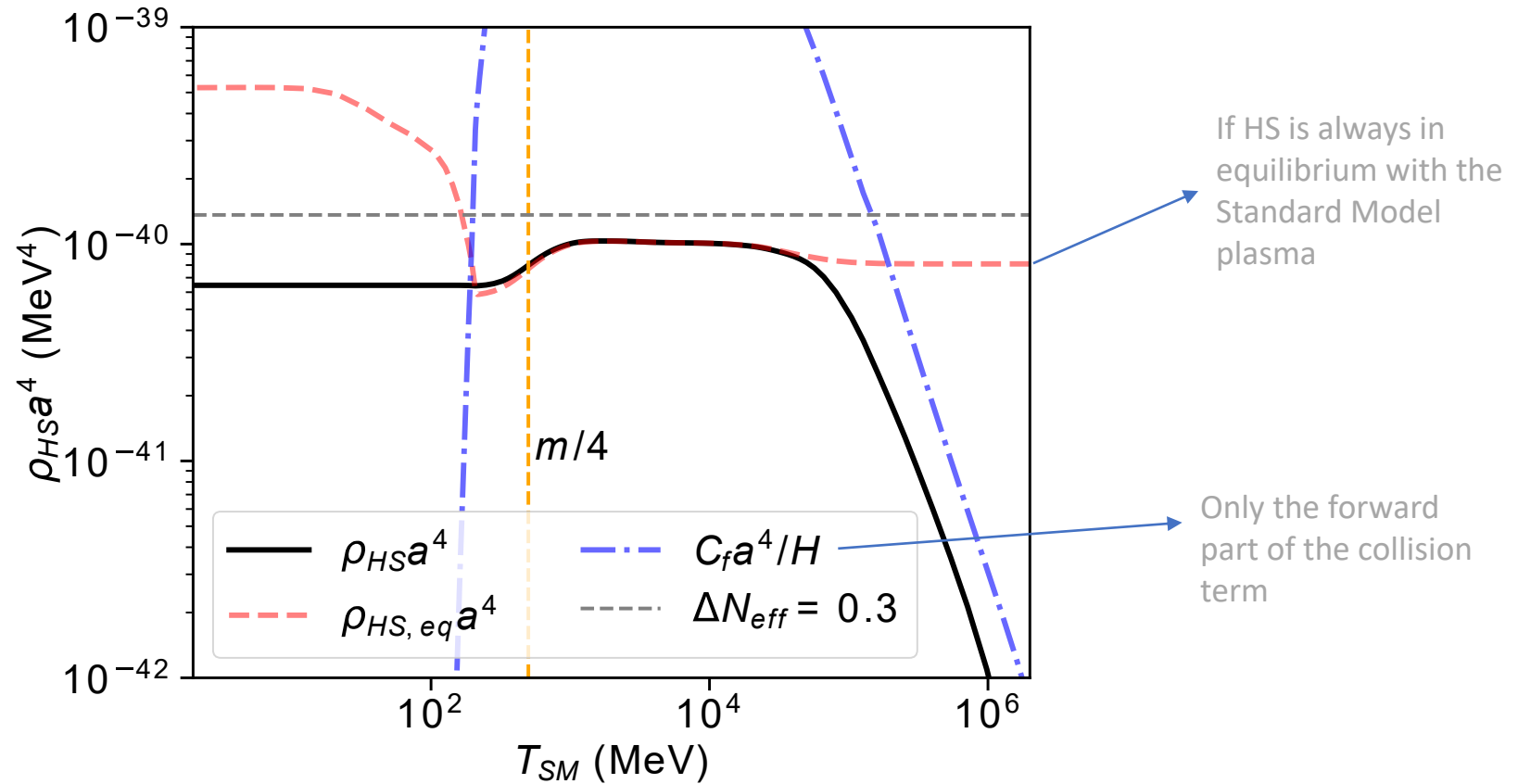
If HS is always in equilibrium with the Standard Model plasma

Only the forward part of the collision term

$$\Gamma_E = \frac{C_f}{\rho_{HS,eq}}$$

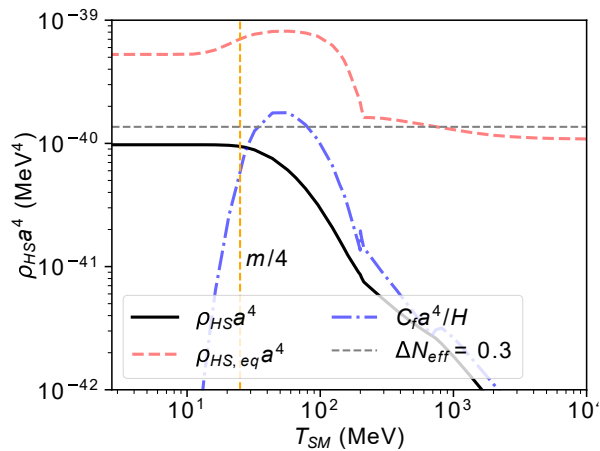
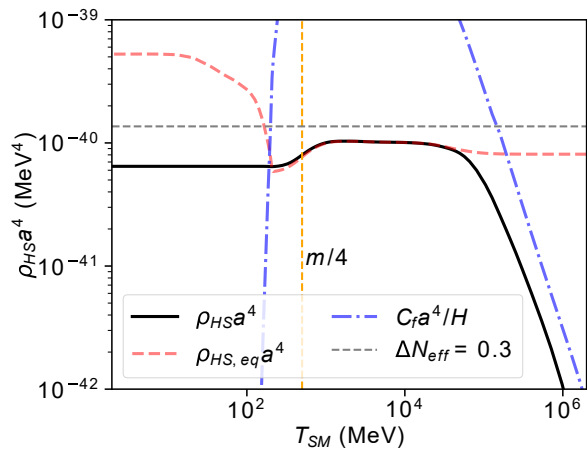
$$\frac{\rho_{HS}}{\rho_{SM}} \sim \left(\frac{\Gamma_E}{H} \right)_{T_{SM} \sim \frac{m}{4}} \propto Q^2 \left(\frac{M_{\text{Pl}}}{m/4} \right)$$

Physics behind dark radiation production: Thermalized ρ_{DR} insensitive to portal coupling

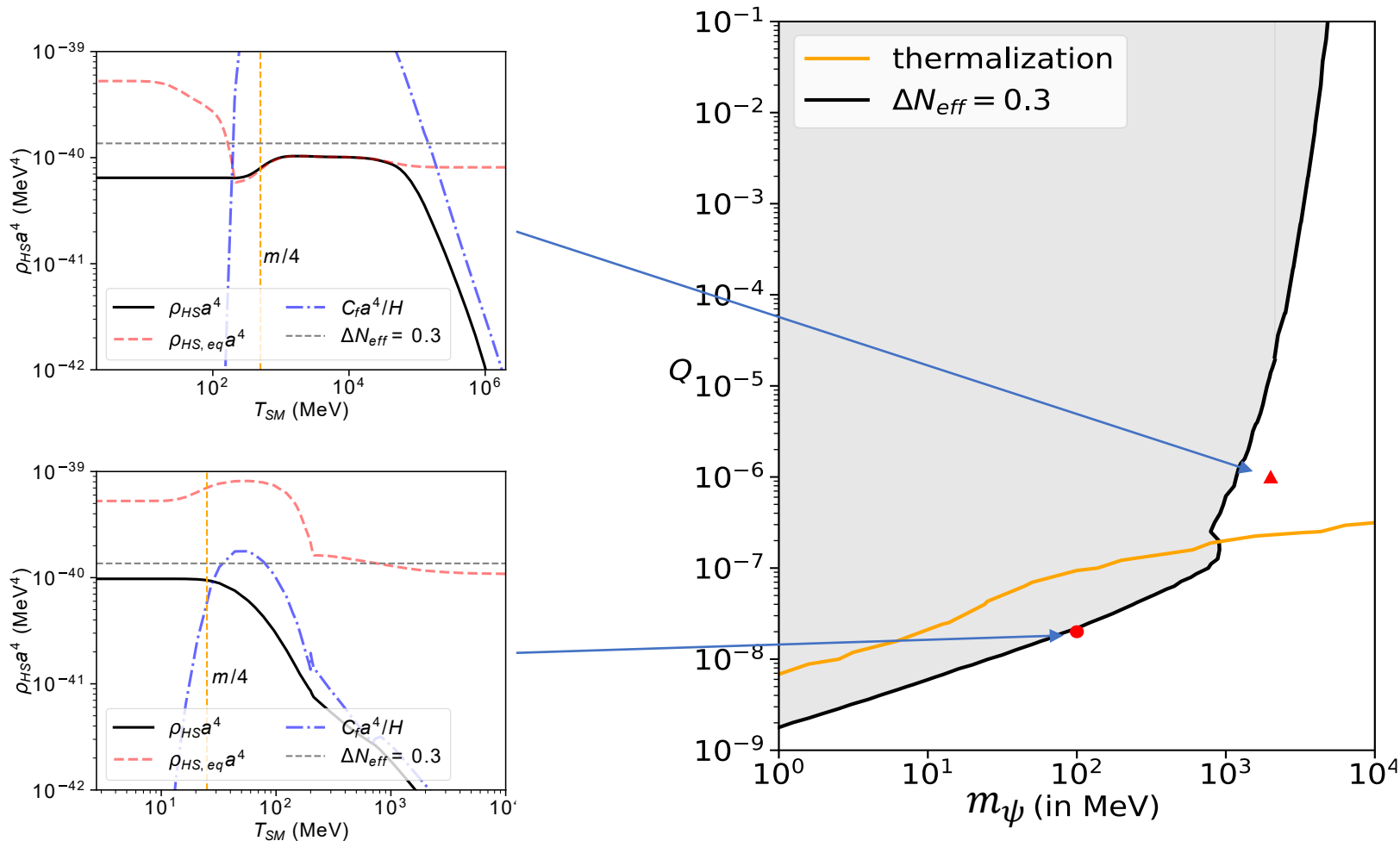


$$\frac{\rho_{HS}}{\rho_{SM}} \propto \frac{g_{HS}}{g_{*SM}(T_d)} \quad \begin{matrix} \rightarrow \\ \rightarrow \end{matrix} \text{Degrees of freedom}$$

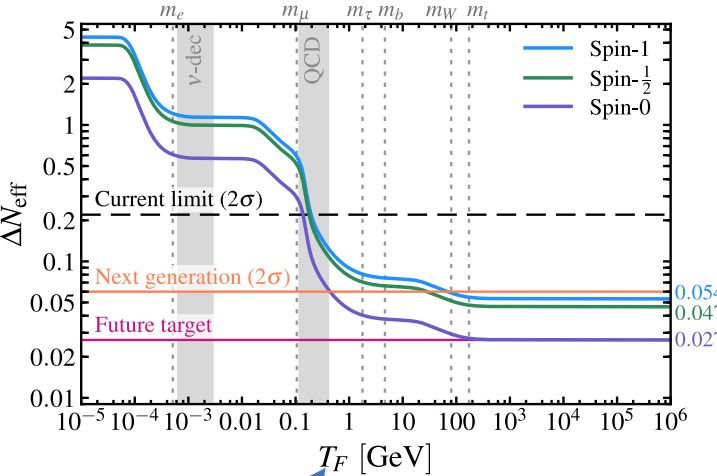
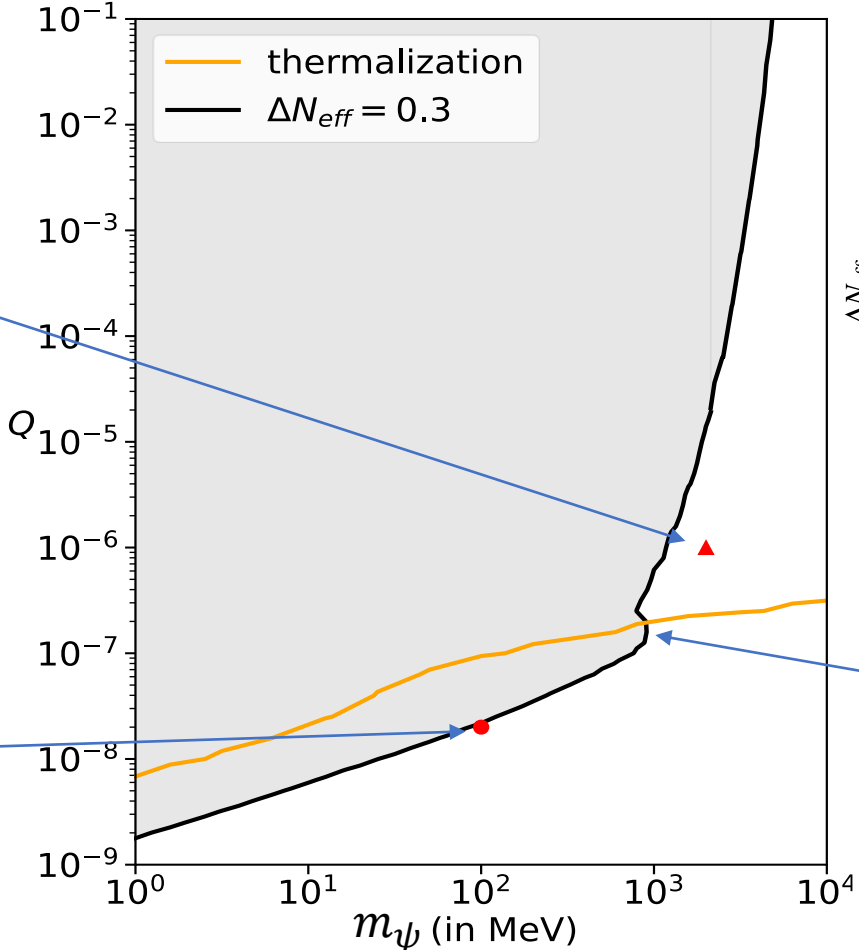
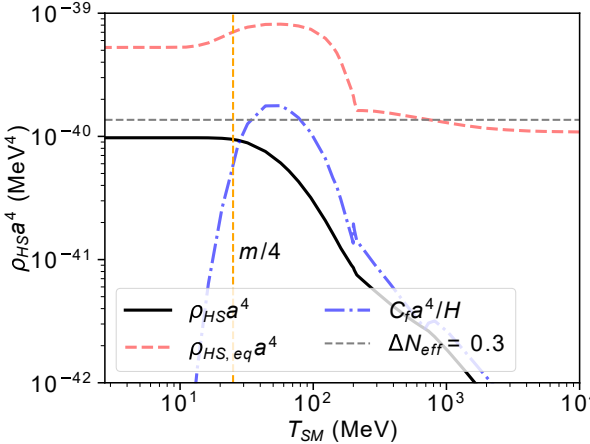
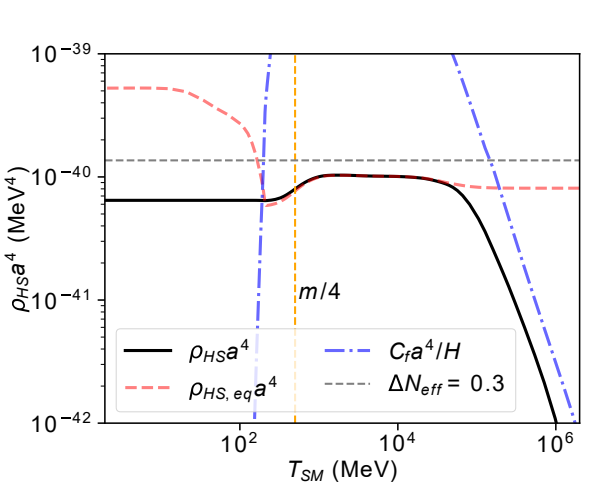
Physics behind dark radiation production: Translating to constraints



Physics behind dark radiation production : Translating to constraints

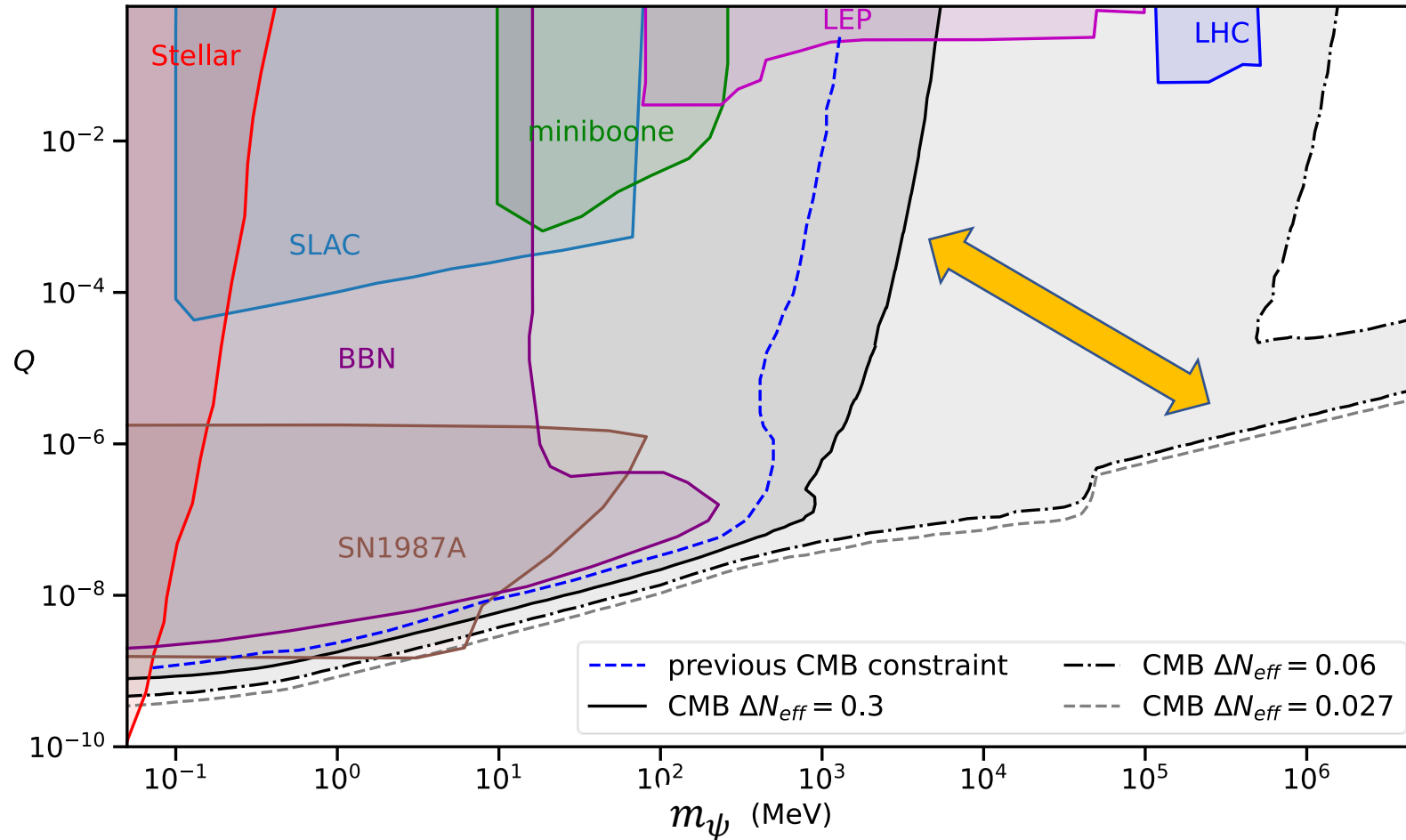


Physics behind dark radiation production : Most relevant when thermally decoupled

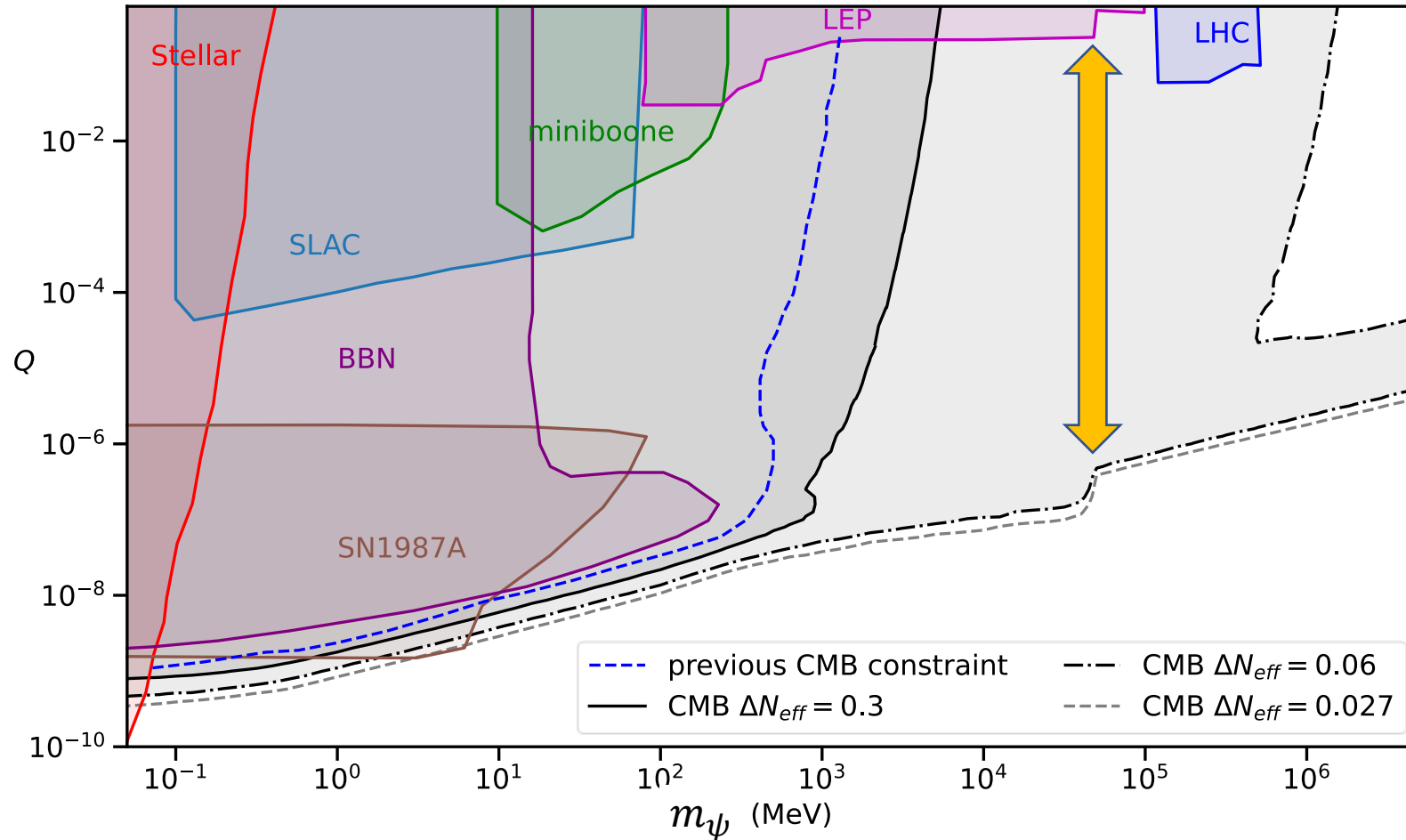


$m_{th} = 4T_{d,s}$

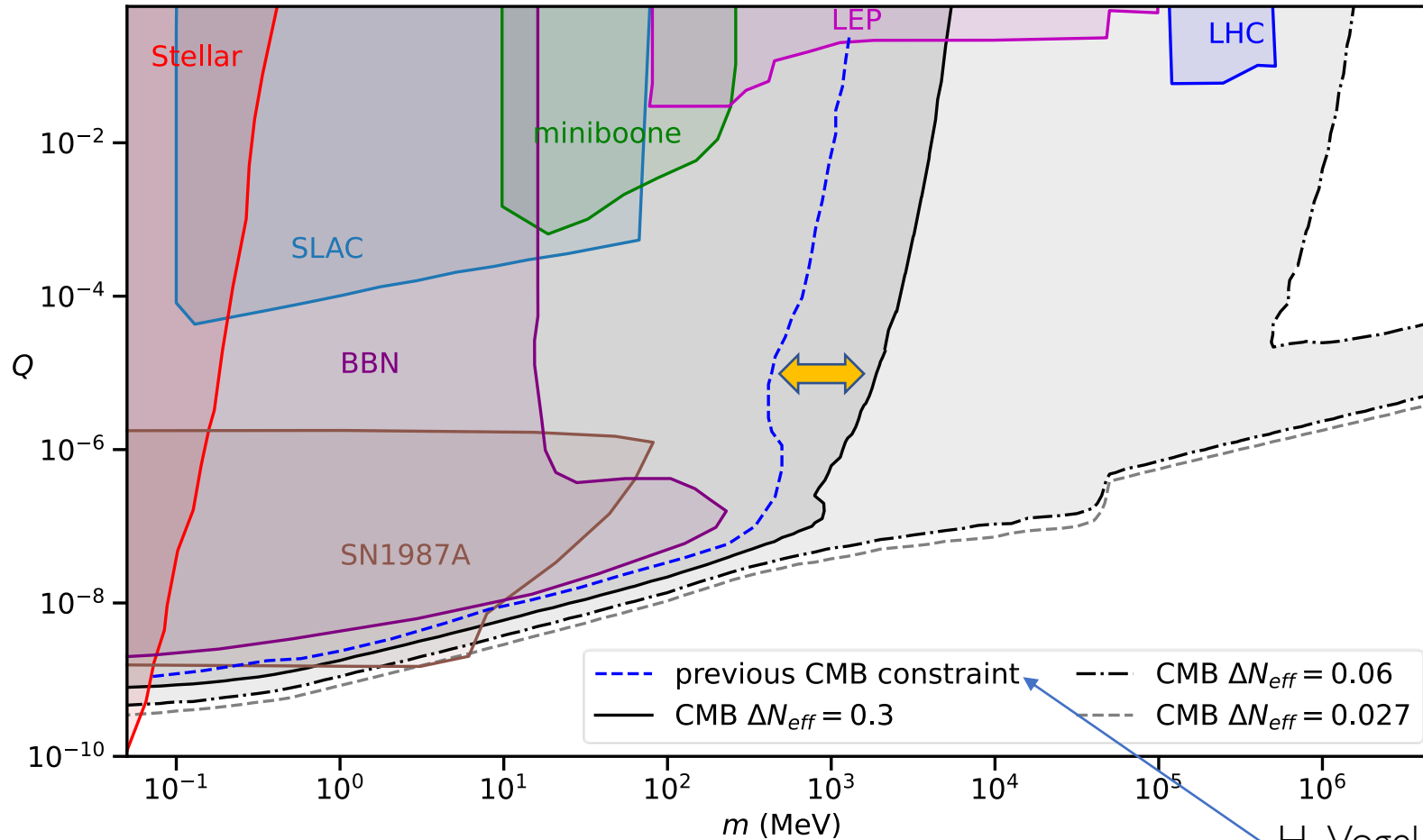
Comparing N_{eff} constraints: Future constraint will extend to much larger parameter space



Comparing Neff constraints: Dominant for $M_\psi > 0.1$ MeV

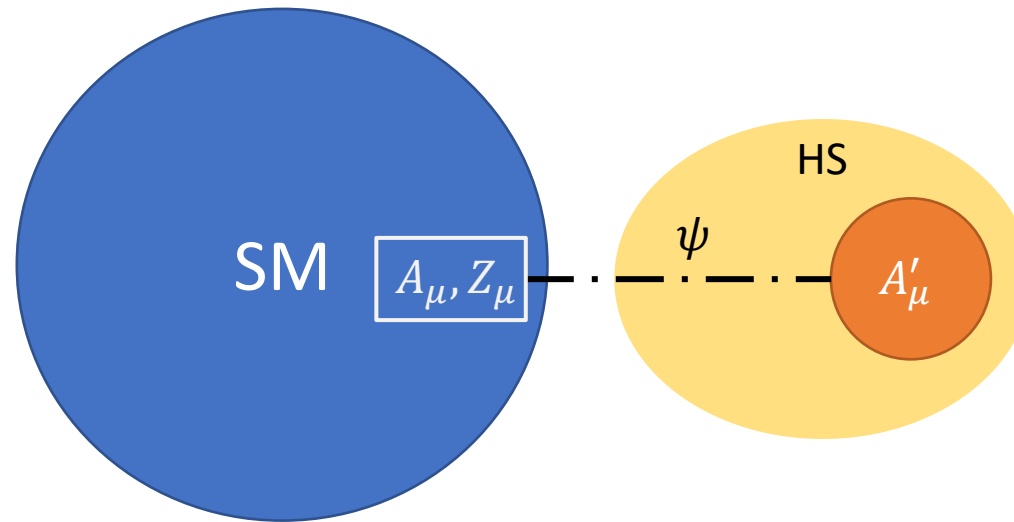


Comparing Neff constraints: Updating previous constraint

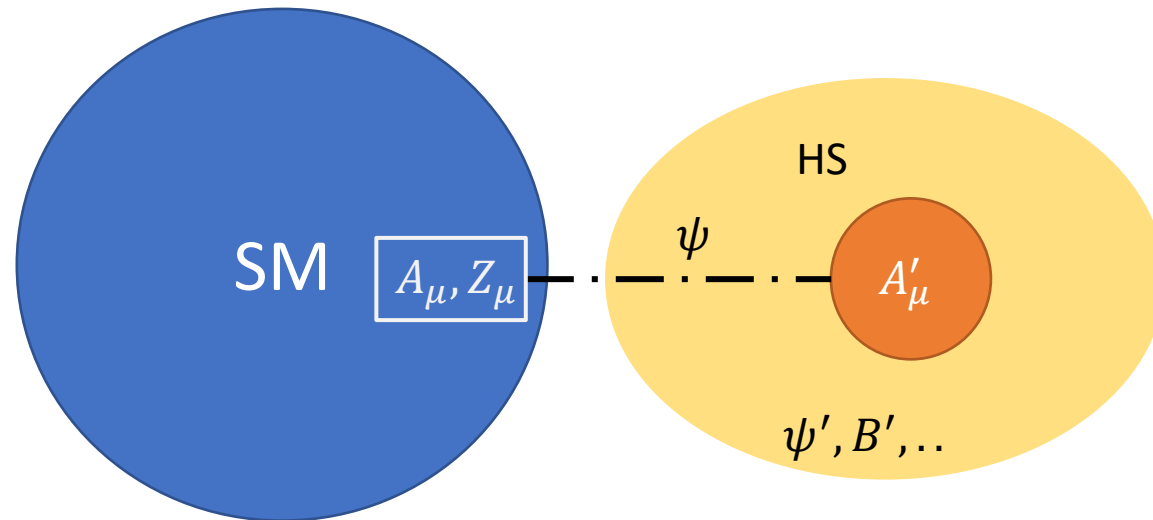


H. Vogel and J. Redondo, JCAP 02 (2014) 029.

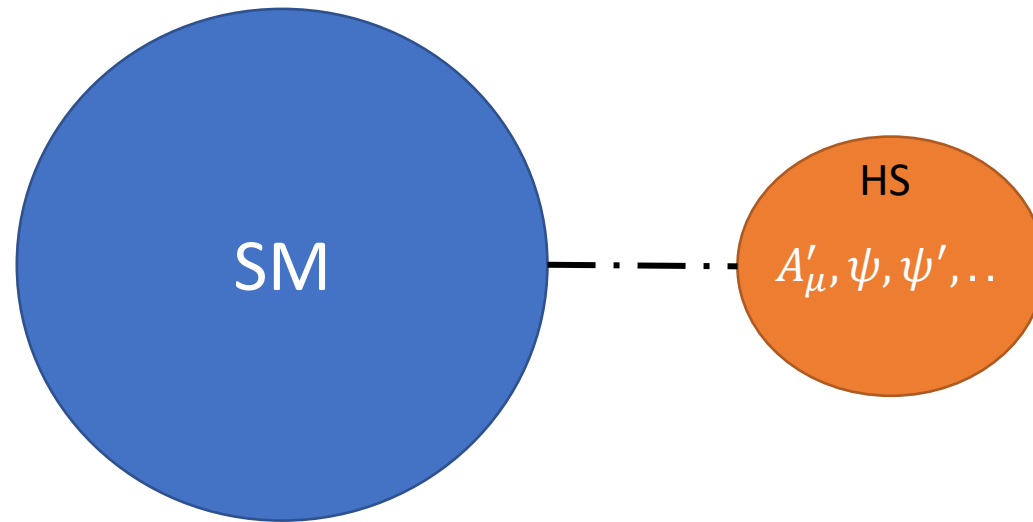
Extending to general hidden sector



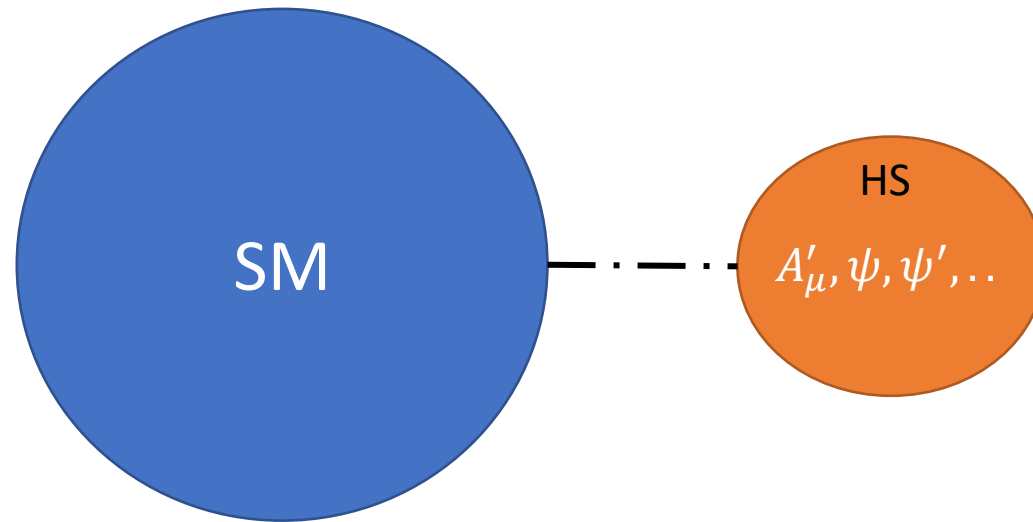
Extending to general hidden sector



Extending to general hidden sector

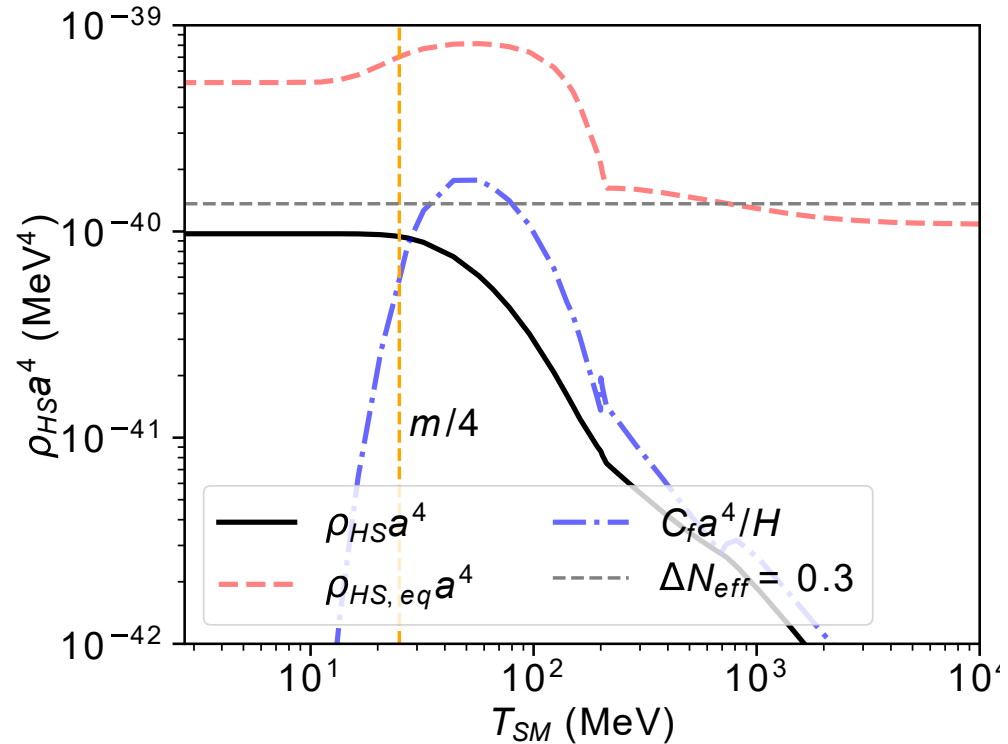
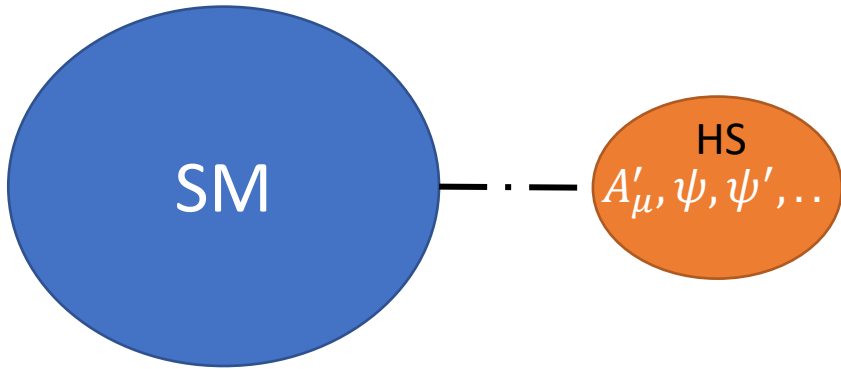


Extending to general hidden sector

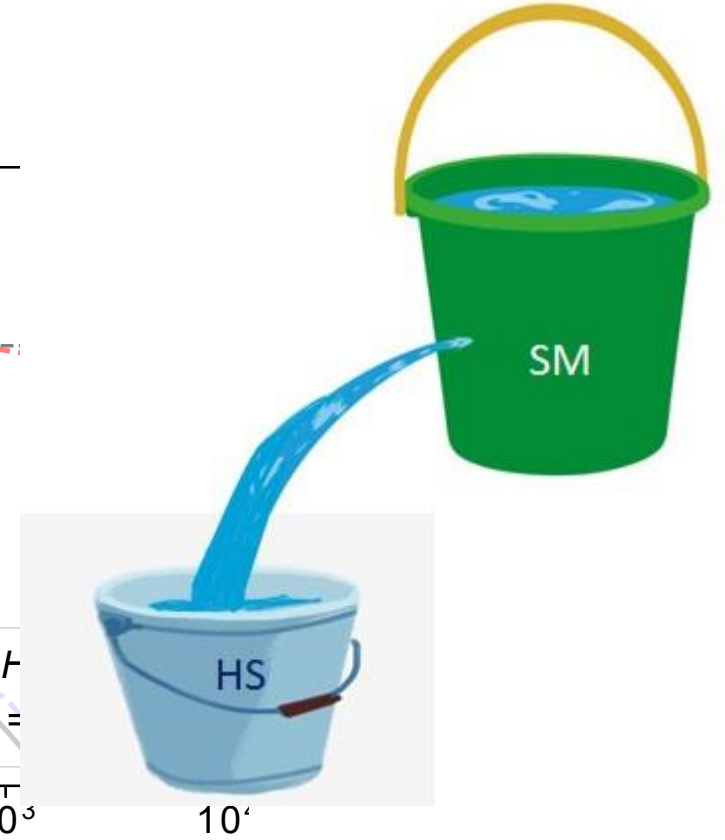
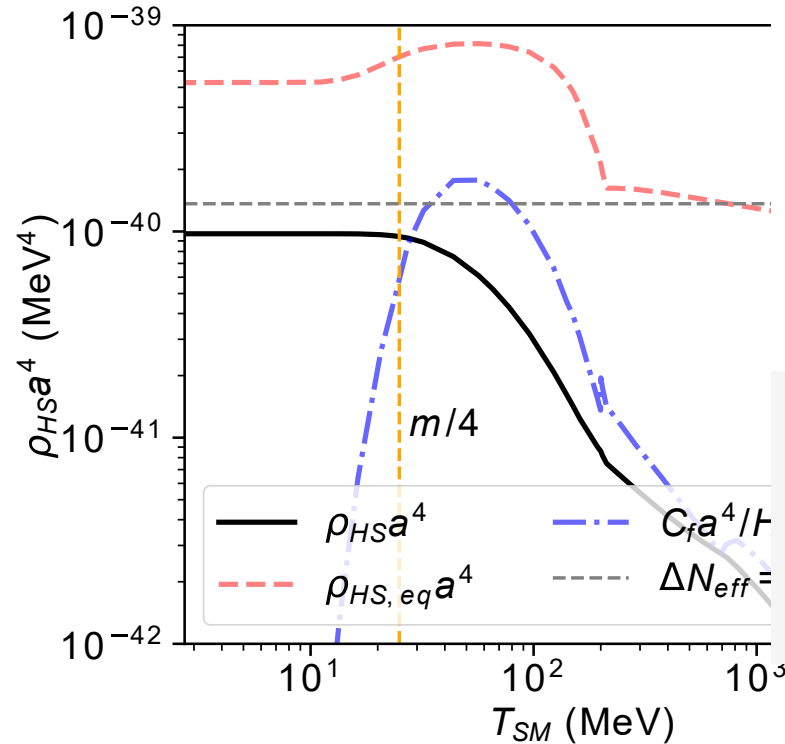
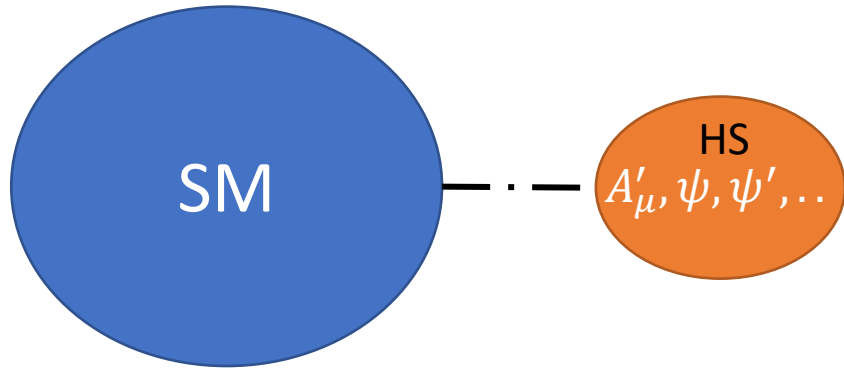


One can calculate a conservative N_{eff} constraint on the millicharge interaction that is independent of details of hidden sector.

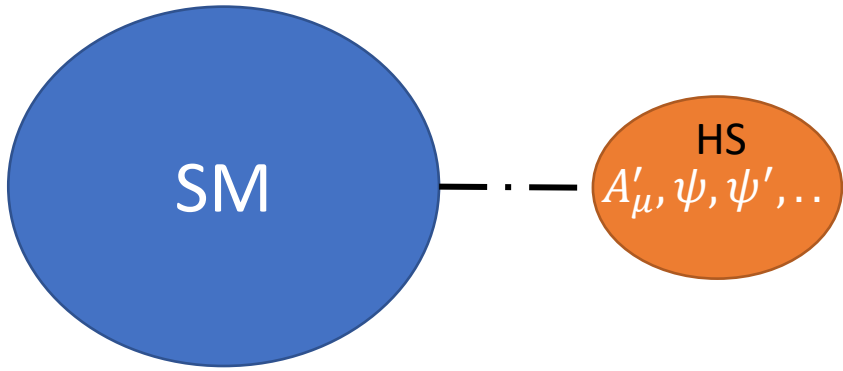
Minimum leaked energy independent of details within HS



Minimum leaked energy independent of details within HS

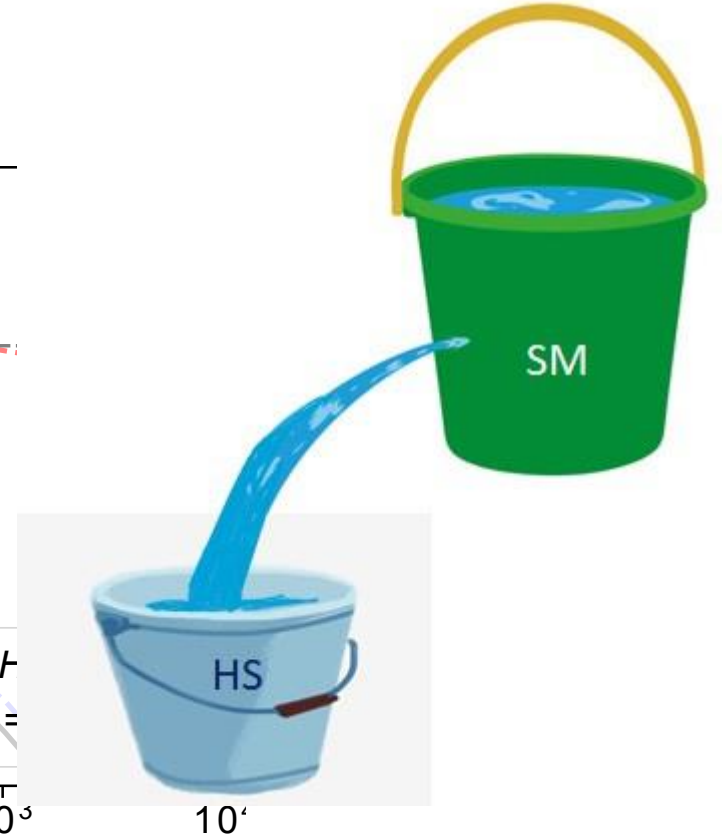
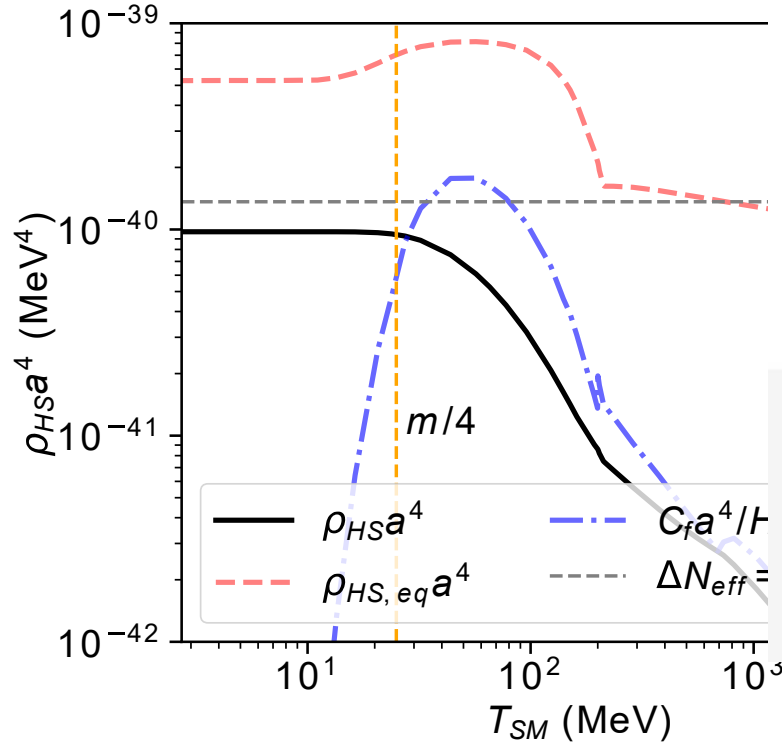


Minimum leaked energy independent of details within HS

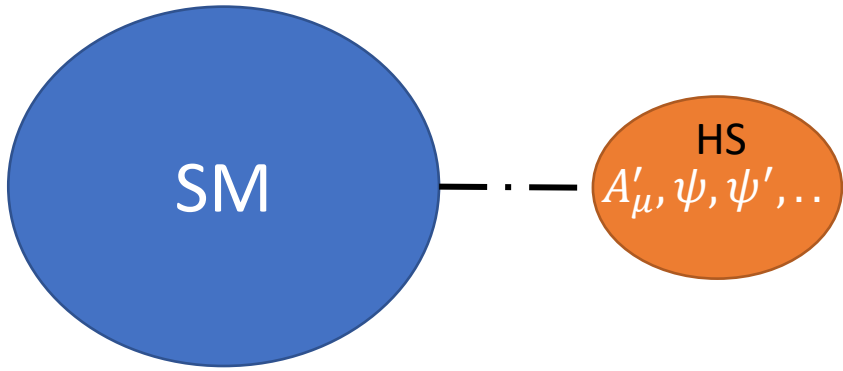


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

Leak factor $L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$



Minimum leaked energy independent of details within HS

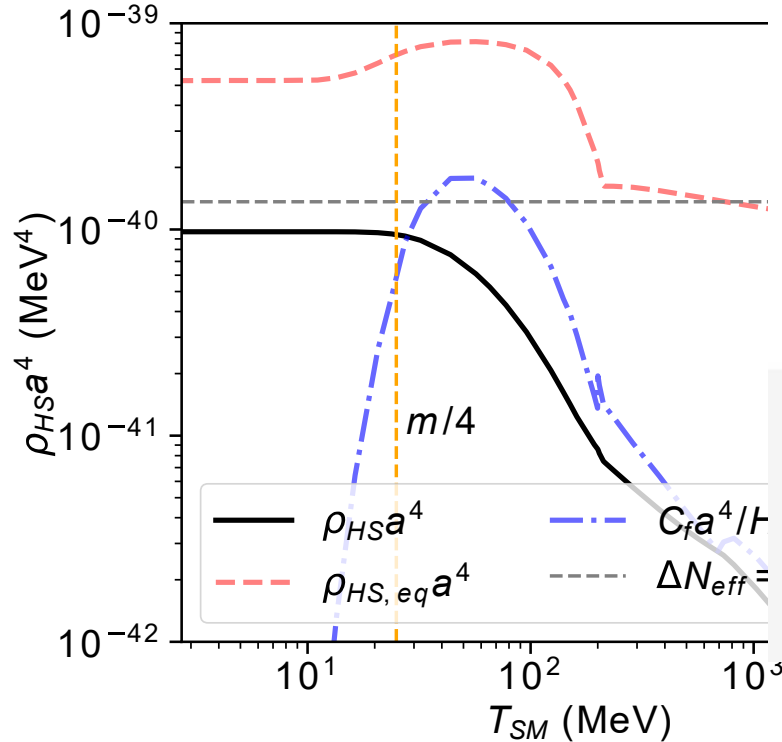


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

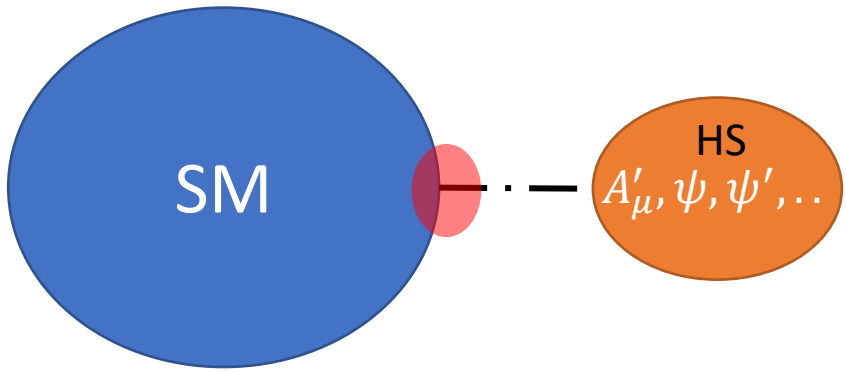
Leak factor

$$L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$$

Independent of g_{HS} or T_{HS} or value of dark charge.



Minimum leaked energy independent of details within HS: Depends only on one BSM coupling

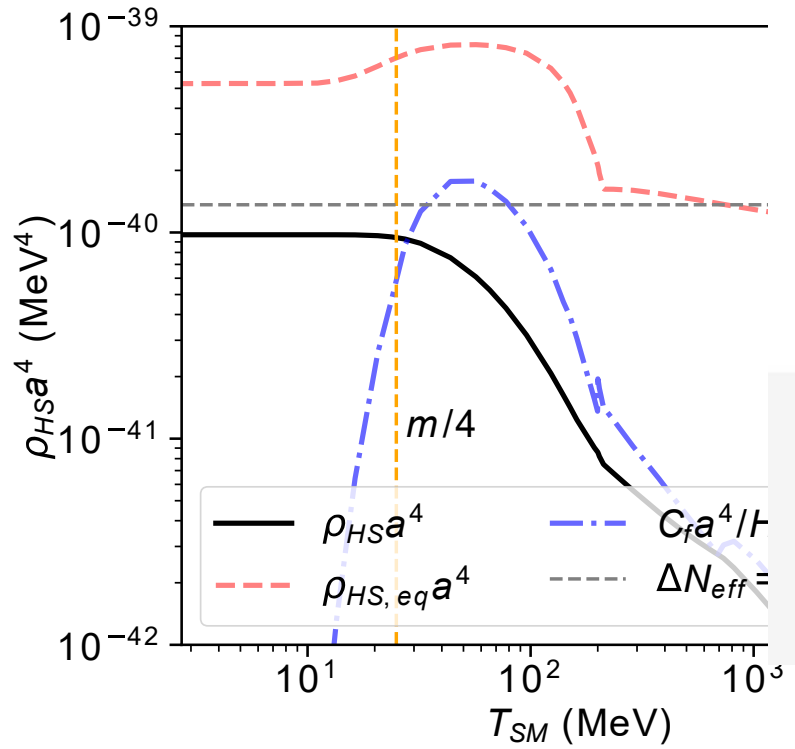


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

Leak factor $L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$

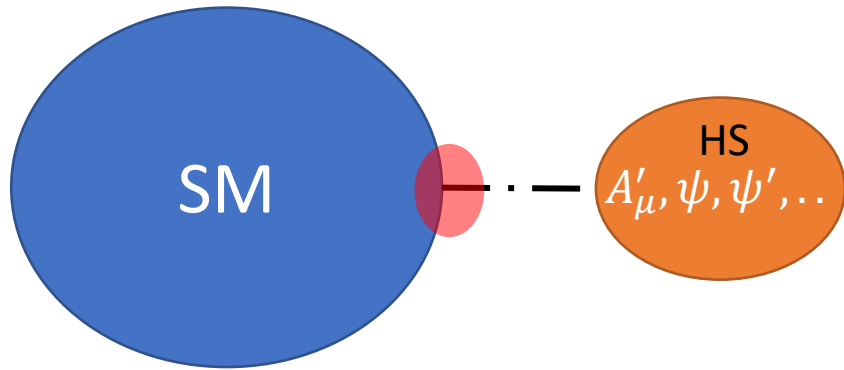
Independent of g_{HS} or T_{HS} or value of dark charge.

Beyond Standard Model



Minimum leaked energy independent of details within HS: Conservative constraint on BSM coupling

Beyond Standard Model

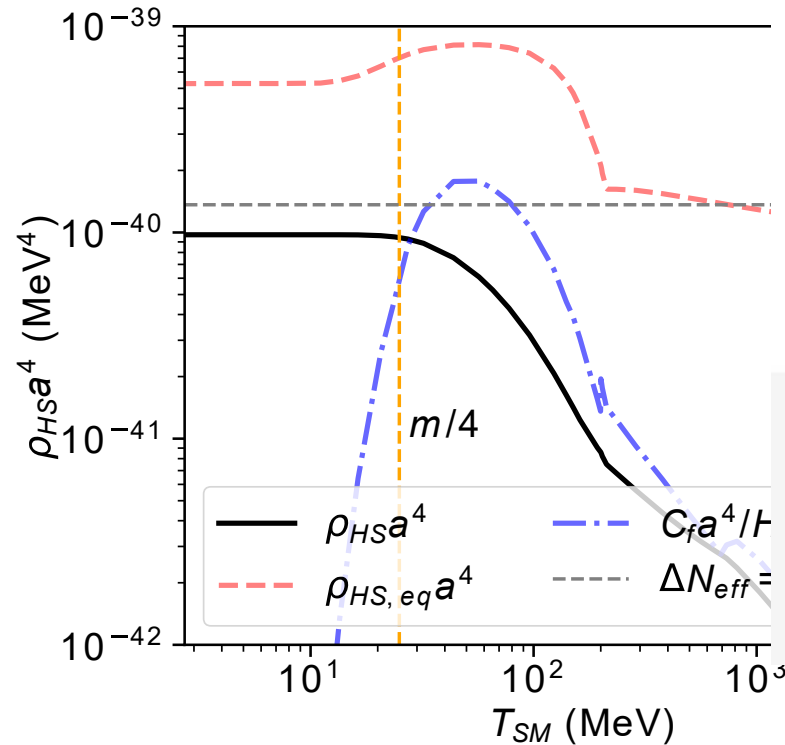


$$\left(\frac{\rho_{HS}}{\rho_{SM}}\right)_{\min} \propto \frac{M_{\text{Pl}}}{m_\psi} L$$

$$L = m_\psi \int ds \frac{s - 4m_f^2}{s\sqrt{s}} \sigma_{ff \rightarrow \psi\psi} \propto Q^2$$

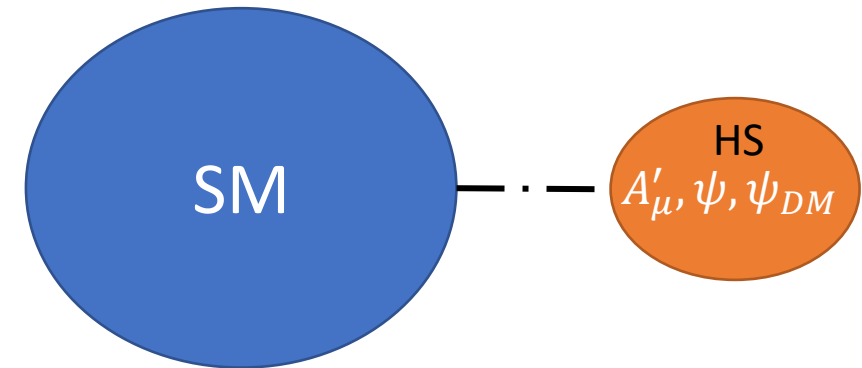
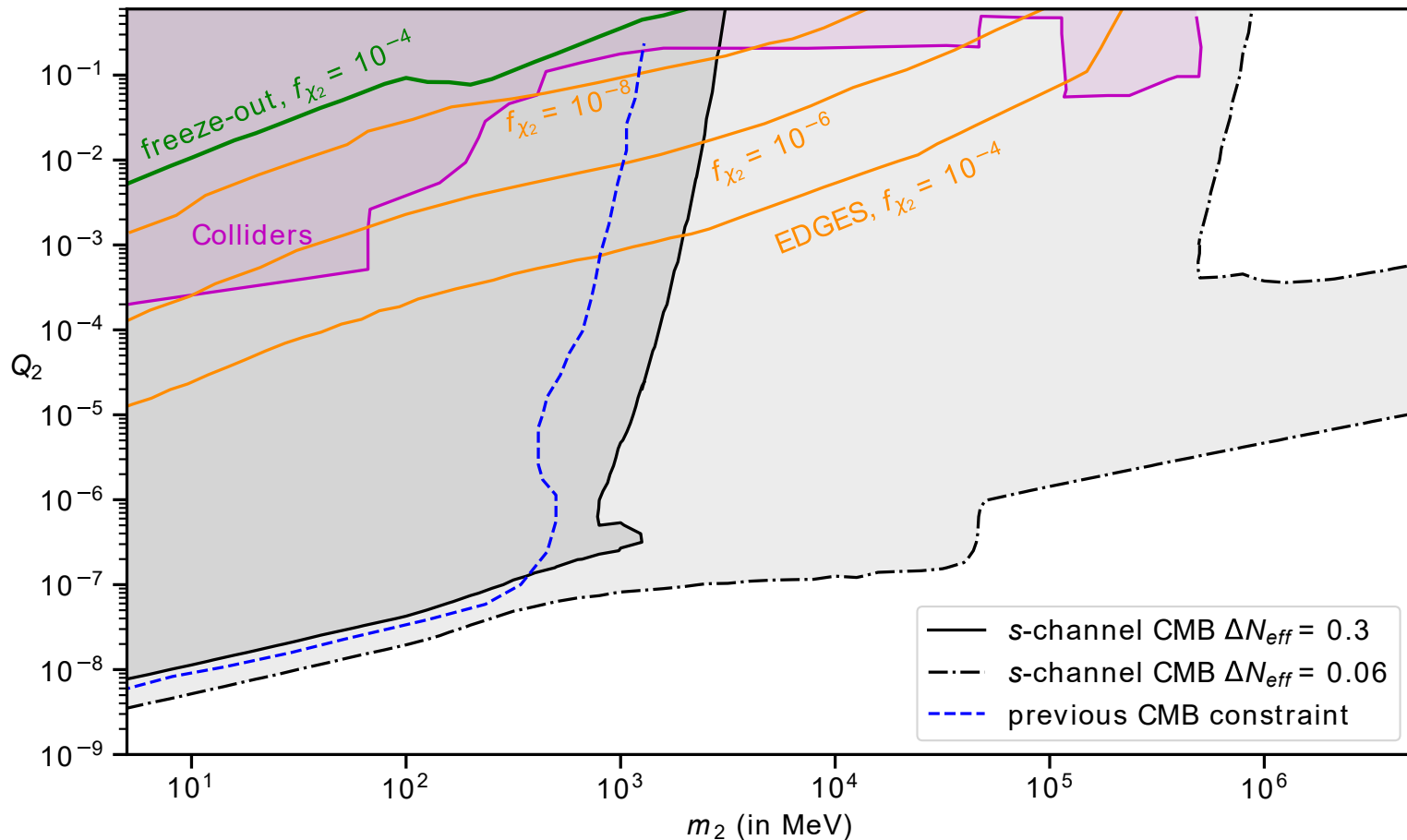
Leak factor

Independent of g_{HS} or T_{HS} or value of dark charge.



$$L < g_*^{3/4}(m) g_*^{1/3} \left(\frac{m}{4}\right) \frac{m/4}{M_{\text{Pl}}} (\Delta N_{\text{eff}})_{\text{max}}$$

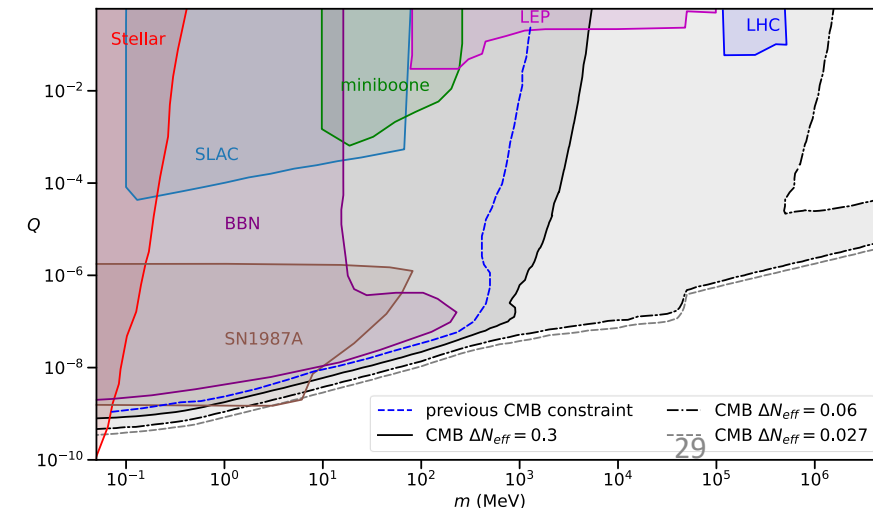
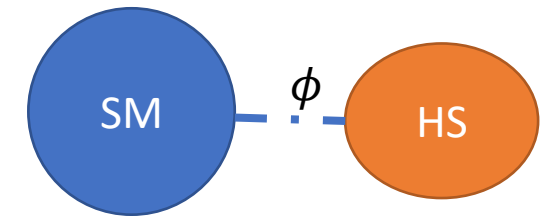
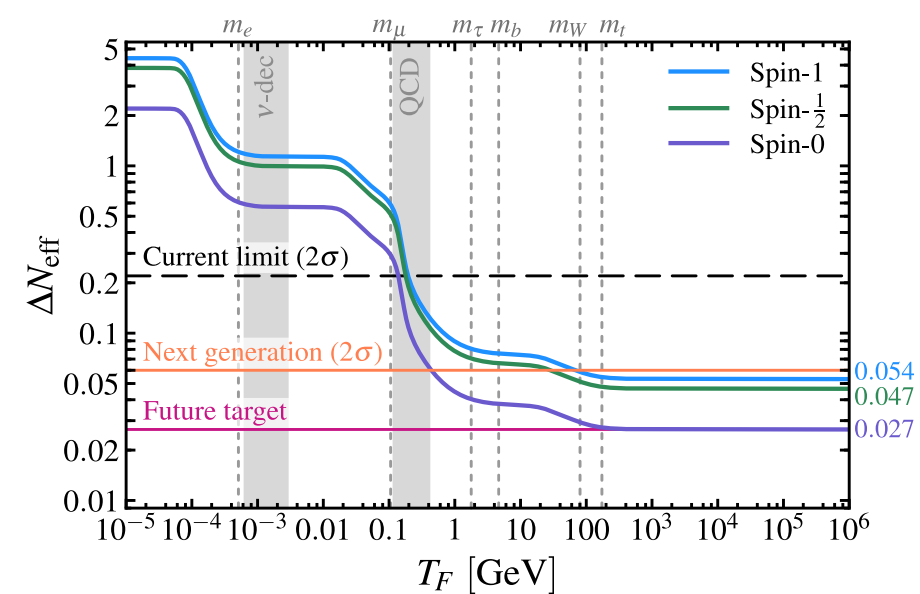
Neff constraints applicable for wide class of hidden sectors: Application to EDGES



H. Liu, N. J. Outmezguine, D. Redigolo, and T. Volansky, Phys. Rev. D 100 no. 12, (2019) 123011.

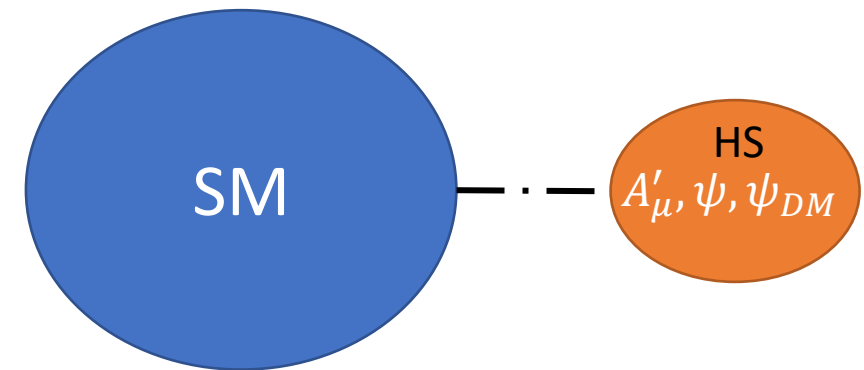
Summary

- With improving N_{eff} measurements, we should interpret them as constraints on portal interactions with out-of-equilibrium sectors
- N_{eff} constraints on out-of-equilibrium particles are:
 - Most relevant for portal interactions mediated by a particle heavier than 0.1 MeV
 - Orders of magnitude stronger than collider experiments
 - Constraints largely independent of internal hidden sector model
- Simple way to calculate: $L < g_*^{3/4} (4\Lambda) g_*^{1/3} (\Lambda) \frac{\Lambda}{M_{\text{Pl}}} (\Delta N_{\text{eff}})_{\text{max}}$



Backup slides

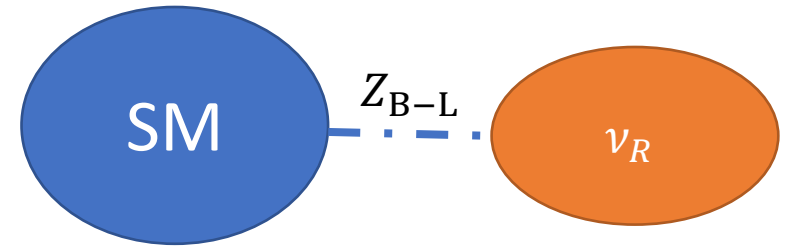
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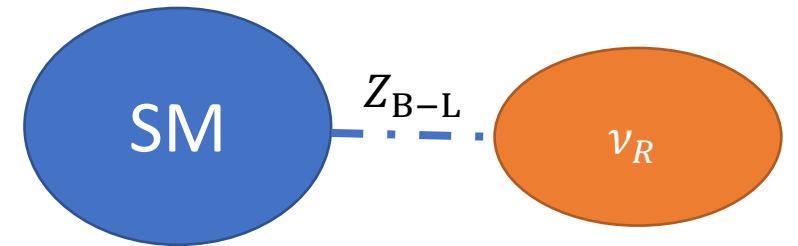
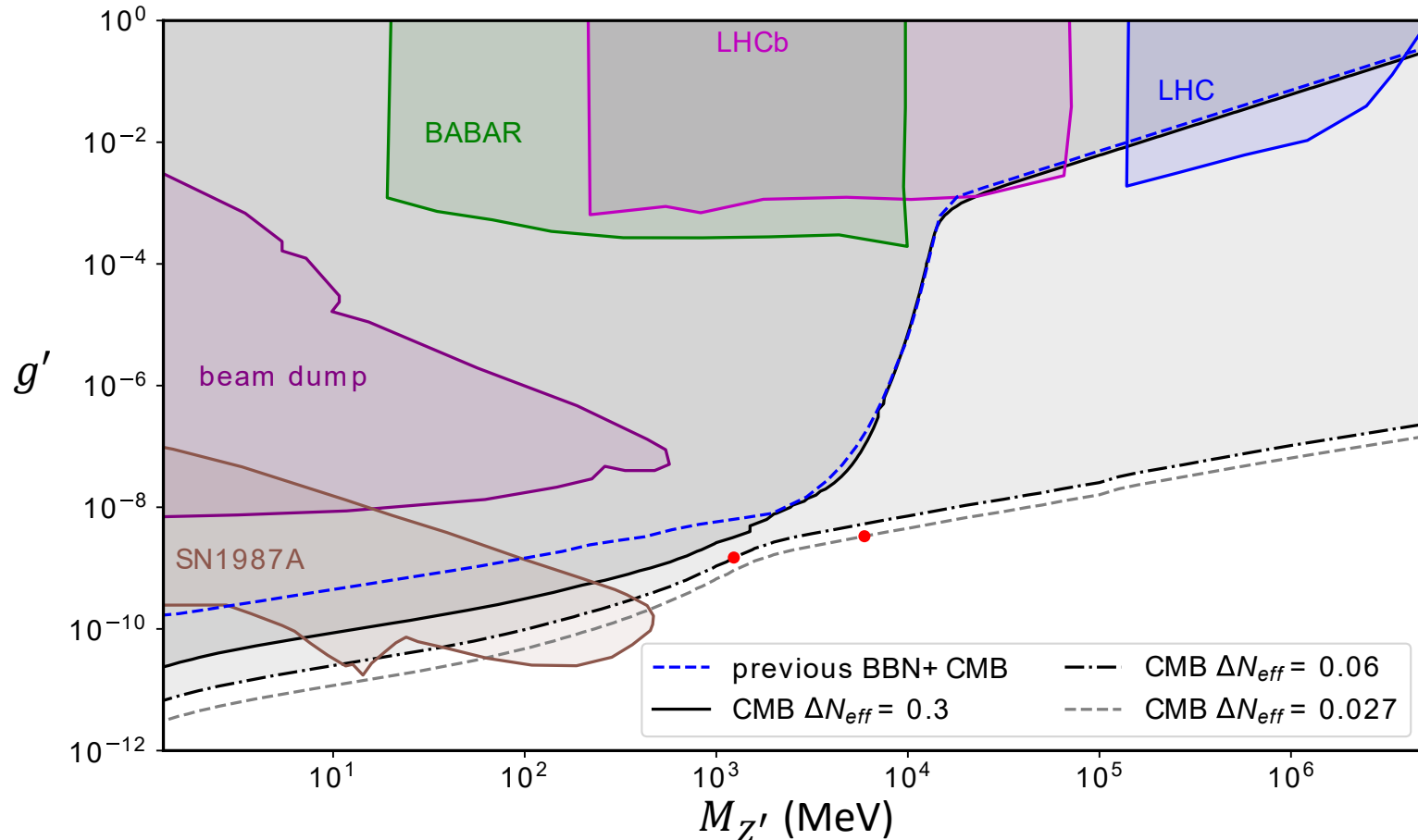
Neff constraints applicable for wide class of hidden sectors: B-L model

$$L_{int} \supset -\frac{1}{4}F'_{\mu\nu}F^{\mu\nu'} + g'Z'_\mu J_{B-L,SM}^\mu - g'Z'_\mu \sum_i \bar{\nu}_{R,i}\gamma^\mu\nu_{R,i} + \frac{1}{2}M_{Z'}^2 Z'^\mu Z'_\mu$$



Neff constraints applicable for wide class of hidden sectors: B-L model

$$L_{int} \supset -\frac{1}{4}F'_{\mu\nu}F^{\mu\nu'} + g'Z'_\mu J_{B-L,SM}^\mu - g'Z'_\mu \sum_i \bar{\nu}_{R,i}\gamma^\mu\nu_{R,i} + \frac{1}{2}M_{Z'}^2 Z'^\mu Z'_\mu$$



Millicharged particles must dominantly annihilate into dark photons

