Experimentally distinguishable origin for electroweak symmetry breaking

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Coleman-Weinberg Mechanism Radiative Symmetry breaking as origin of SM Higgs potential

$$V = \left[-m^2(\Phi^{\dagger}\Phi)\right] + \lambda_{\Phi}(\Phi^{\dagger}\Phi)^2, \quad m^2 > 0$$
 (1)

- Effective negative mass-squared term can be induced via the Coleman-Weinberg mechanism (Coleman & Weinberg, 1973).
 - ► Large top Yukawa coupling in SM prevents straightforward application.
- Extend SM minimally with a new hidden U(1) gauge group, dubbed U(1)_H, containing a Higgs scalar Φ.
- Implement CW mechanism for SM-extended Φ sector by imposing classical conformality (Iso, Okada, & Orikasa, 2009).

Coleman-Weinberg Mechanism

Radiative Symmetry breaking as origin of SM Higgs potential

• Hidden $U(1)_H$ sector scalar effective potential of the form

$$V_{\phi} = \lambda_{\phi} \left(\Phi^{\dagger} \Phi \right)^{2} + V_{1-loop}$$

= $\frac{1}{4} \lambda_{\phi} \phi^{4} + \frac{\beta_{\phi}}{8} \phi^{4} \left(ln \left[\frac{\phi^{2}}{v_{\phi}^{2}} \right] - \frac{25}{6} \right)$, where $\phi = \sqrt{2} \text{Re} \left[\Phi \right]$ (2)

 \bullet Combined SM Higgs and Φ scalar potential is

$$V = \lambda_h \left(H^{\dagger} H \right)^2 - \left[\lambda_{mix} \left(H^{\dagger} H \right) \left(\Phi^{\dagger} \Phi \right) \right] + V_{\phi}$$
(3)

- Radiative symmetry breaking in $U(1)_H$ sector at $\langle \phi
 angle = {\it v}_\phi$
- With $\lambda_{mix} > 0$, $\langle \phi \rangle = v_{\phi}$ generates negative SM Higgs mass squared term, driving EW symmetry breaking.

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Coupling Analysis

Conventional:
$$V = \frac{\lambda_h}{4}(h^2 - v_h^2)^2 + \frac{\lambda_\phi}{4}(\phi^2 - v_\phi^2)^2 + \frac{\lambda_{mix}}{4}(h^2 - v_h^2)(\phi^2 - v_\phi^2)$$

Conformal: $V = \frac{\lambda_h}{4}h^4 + \frac{\lambda_\phi}{4}\phi^4 + \frac{\beta_\phi}{8}\phi^4 \left(\ln\left[\frac{\phi^2}{v_\phi^2}\right] - \frac{25}{6}\right) - \frac{\lambda_{mix}h^2\phi^2}{4}$

• Mass-squared matrices defined as

$$M_{sq} = \begin{pmatrix} \partial_h^2 V & \partial_h \partial_\phi V \\ \partial_\phi \partial_h V & \partial_\phi^2 V \end{pmatrix} \Big|_{h=v_h, \phi=v_\phi} = \begin{pmatrix} m_h^2 & M^2 \\ M^2 & m_\phi^2 \end{pmatrix}$$

• Diagonalize M_{sq} to find mixing of eigenstates:

$$h = h_1 \cos(\theta) + h_2 \sin(\theta)$$

$$\phi = -h_1 \sin(\theta) + h_2 \cos(\theta)$$

In our analysis, we set $M_{h_1} > 2M_{h_2}, \theta \ll 1 \Rightarrow \begin{bmatrix} h_1 \sim h, h_2 \sim \phi \end{bmatrix}$

Coupling Analysis

• Express potentials in terms of observables and extract couplings.

For $M_{h_1} > 2M_{h_2}, \theta \ll 1$, conventional system coupling goes as

$$g_{h_{1}h_{2}h_{2}} \simeq -\frac{M_{h_{1}}^{2}}{2v_{\phi}} \left(1 + 2\frac{M_{h_{2}}^{2}}{M_{h_{1}}^{2}}\right) \theta \quad \text{for} \quad \theta \ll \frac{v_{h}}{v_{\phi}}, \qquad (4)$$

$$g_{h_{1}h_{2}h_{2}} \simeq \frac{M_{h_{1}}^{2}}{2v_{h}} \left(1 + 2\frac{M_{h_{2}}^{2}}{M_{h_{1}}^{2}}\right) \theta^{2} \quad \text{for} \quad \frac{v_{h}}{v_{\phi}} \lesssim \theta, \qquad (5)$$

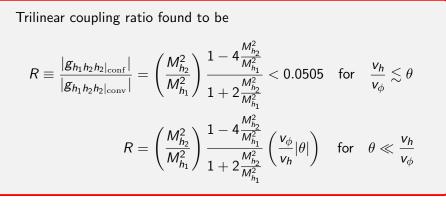
while conformal system coupling goes as

$$g_{h_1h_2h_2} \simeq -\frac{M_{h_2}^2}{2v_h} \left(1 - 4\frac{M_{h_2}^2}{M_{h_1}^2}\right) \theta^2.$$
 (6)

 Combination of cancellation of lower order θ terms and unique structure leads to coupling suppression in the conformal system.

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Coupling Analysis



• Using $M_{h_1} = 125$ GeV, $M_{h_2} = 25$ GeV, $v_h = 246$ GeV, $v_{\phi} = 10^4$ GeV, and $|\theta| = 0.1$:

Conventional system: $g_{h_1h_2h_2} = 0.424$ CW system: $g_{h_1h_2h_2} = -0.0107$ Numerical Analysis: Br $(h_1 \rightarrow h_2 h_2)$

- Gray regions excluded by LHC (ATLAS, 2020) and LEP-II (for $M_{h_2} = 25$ GeV) (LEP-II, 2003)
- Prospective ILC search reach indicated by blue region for anomalous Higgs decay (Liu, Wang, Zhang, 2017) and red region for anomalous coupling (Barklow et. al., 2018).

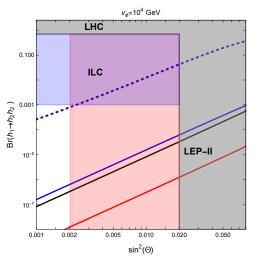
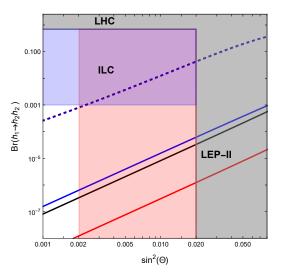


Figure: Conventional (dashed) and Conformal (solid) branching ratios. $M_{h_2} = 10$ (red), 25 (black), and 50 (blue) GeV.

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Numerical Analysis: Br $(h_1 \rightarrow h_2 h_2)$

- Conformal and conventional scenarios are distinguishable!
- Anomalous Higgs decay signal is the key
 - Simultaneous measurement of $h_1 \rightarrow h_2 h_2$ and θ points to conventional model
 - Lack of $h_1 \rightarrow h_2 h_2$ measurement with θ measurement points to conformal model



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Summary

- Classical conformal structure & Coleman-Weinberg mechanism as origin of EW Symmetry breaking.
 - ► Radiative symmetry breaking in U(1)_H sector induces negative SM Higgs mass term, driving EW symmetry breaking
- Unique structure of conformal potential greatly affects Higgs physics/phenomenology
 - ► Most notably, trilinear coupling g_{h1h2h2} suppression in CW model vs. conventional model
 - ▶ Models distinguishable by precision measurement of anomalous Higgs coupling alongside (non-)observation of anomalous Higgs decay $h_1 \rightarrow h_2 h_2 \rightarrow b \bar{b} b \bar{b}$ at future e^+e^- colliders (ILC)
- Can consider U(1)_H as hidden sector, home to DM candidate (subsequent work)
- Extend to hidden non-Abelian gauge groups.