

# Searches for new physics in events with leptons at CMS

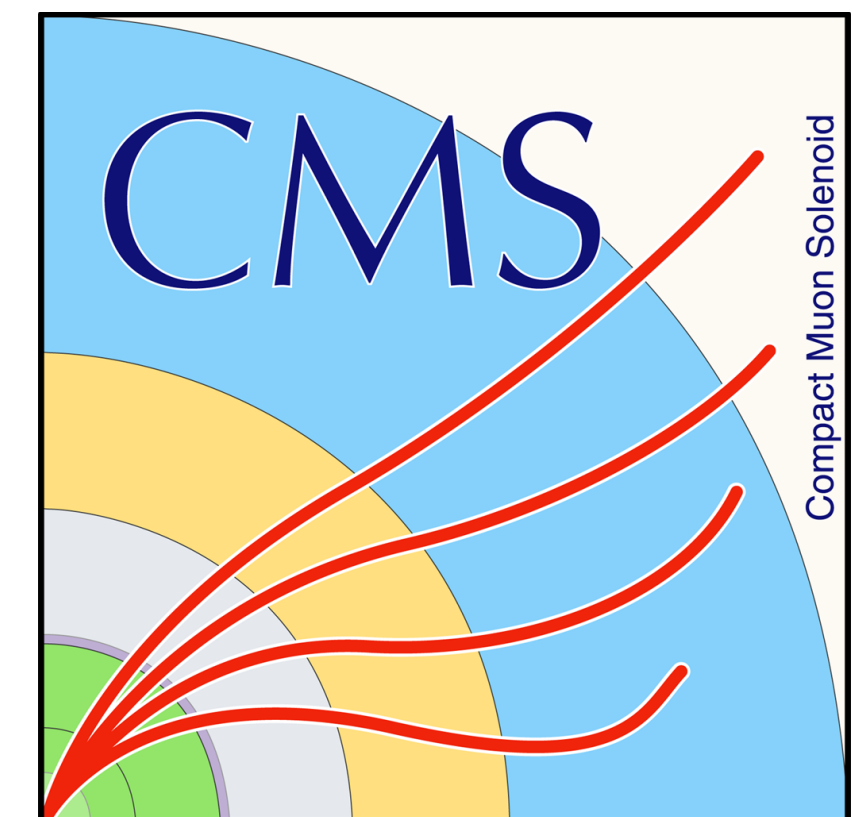
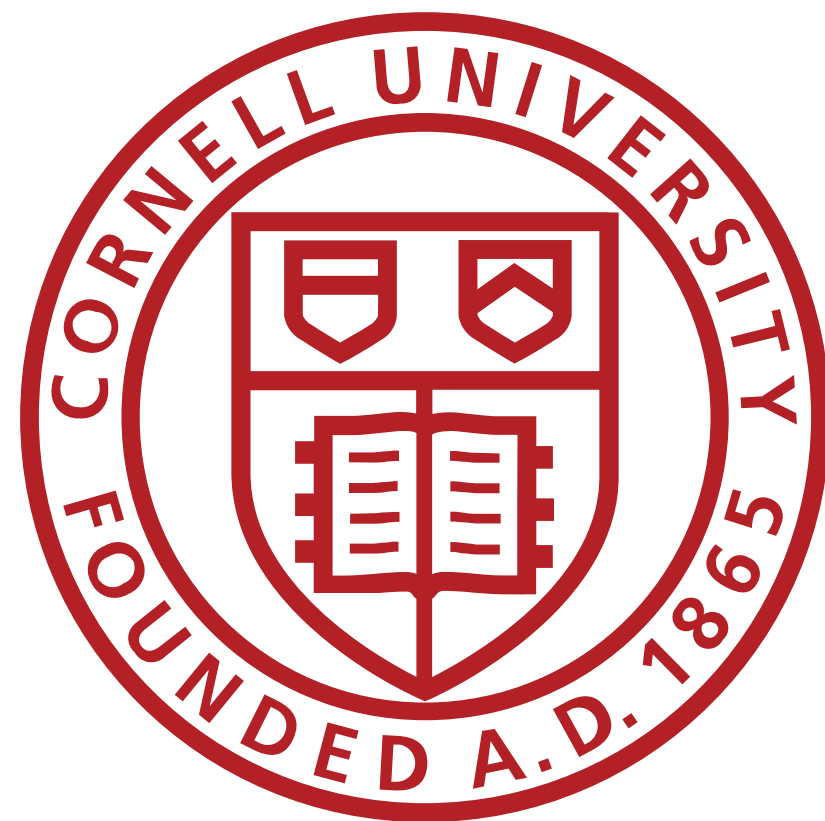
**PHENO** 2022

**Joey Reichert**

on behalf of the  
CMS Collaboration

Cornell University  
joey@cern.ch

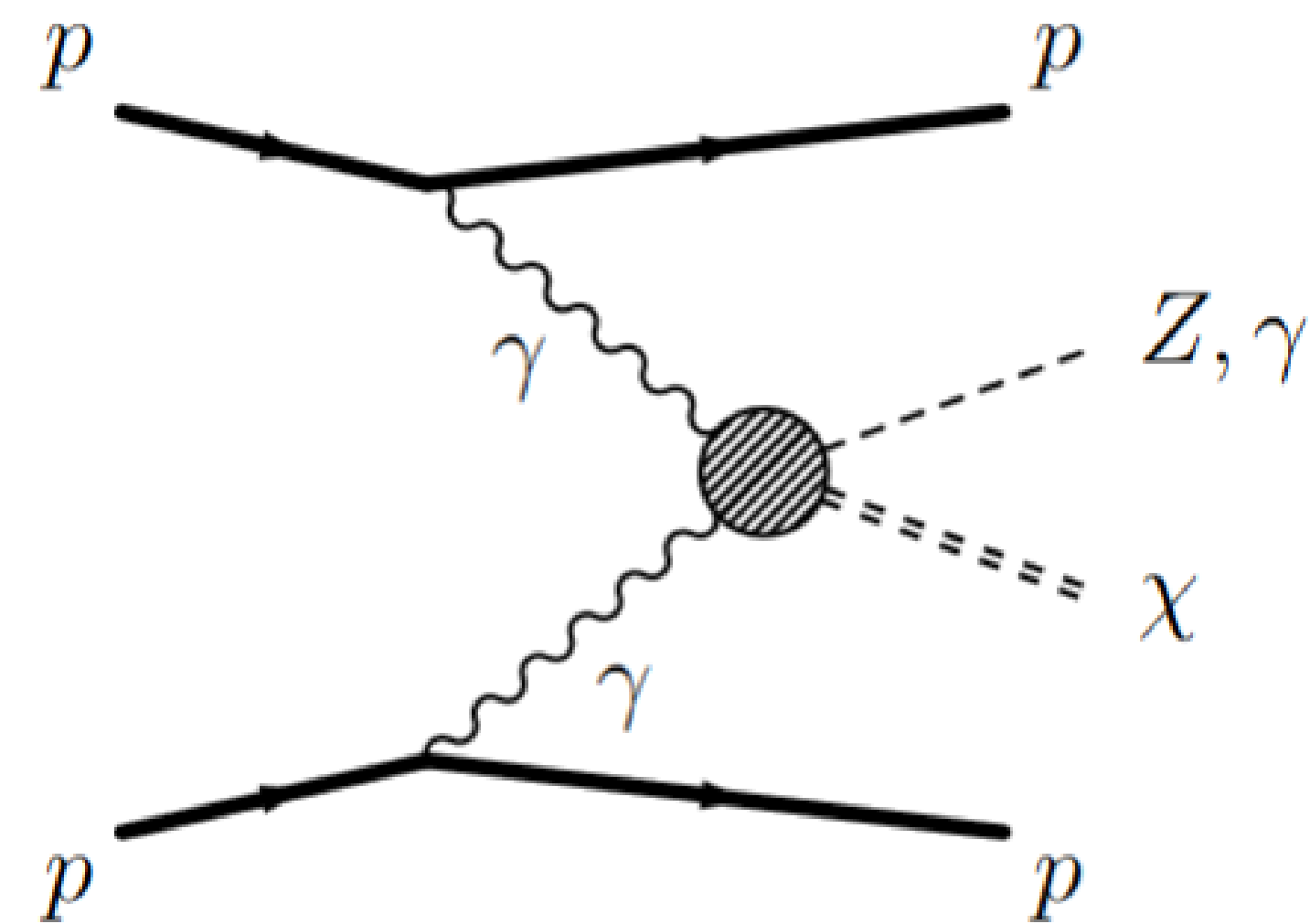
May 9, 2022



# Introduction

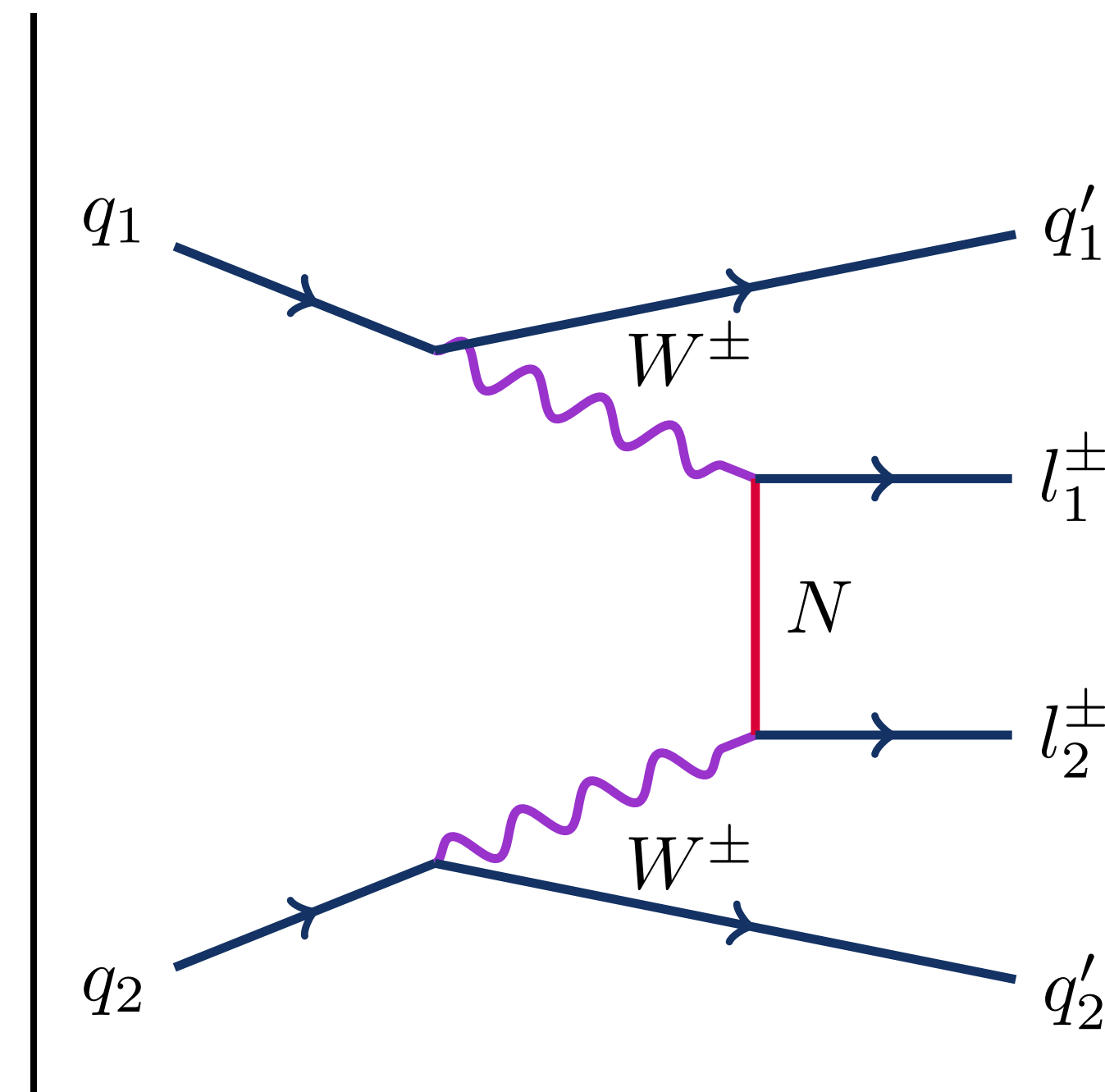
Leptons provide *exceptionally clean* search signatures for new physics at the LHC, and many BSM physics models predict leptonic final states.

*Today, we'll discuss two recent, first-of-their-kind searches w/ leptons at CMS:*



Z or  $\gamma$  + X  
missing mass  
search

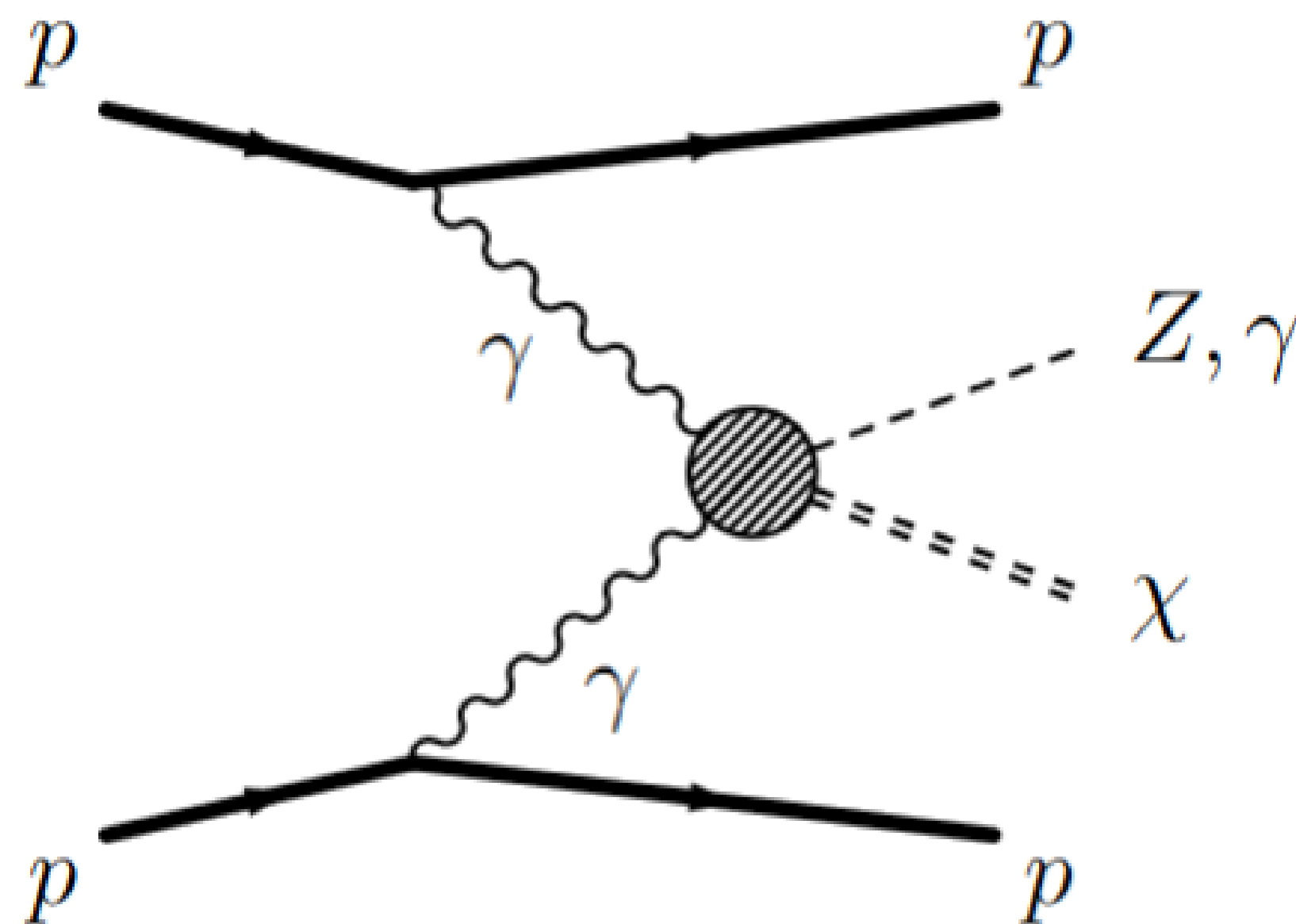
[EXO-19-009](#)



Majorana neutrino  
search in VBF  
 $\mu^\pm\mu^\pm$  events

[EXO-21-003](#)

# Missing mass search w/ the CMS-TOTEM Precision Proton Spectrometer (PPS)

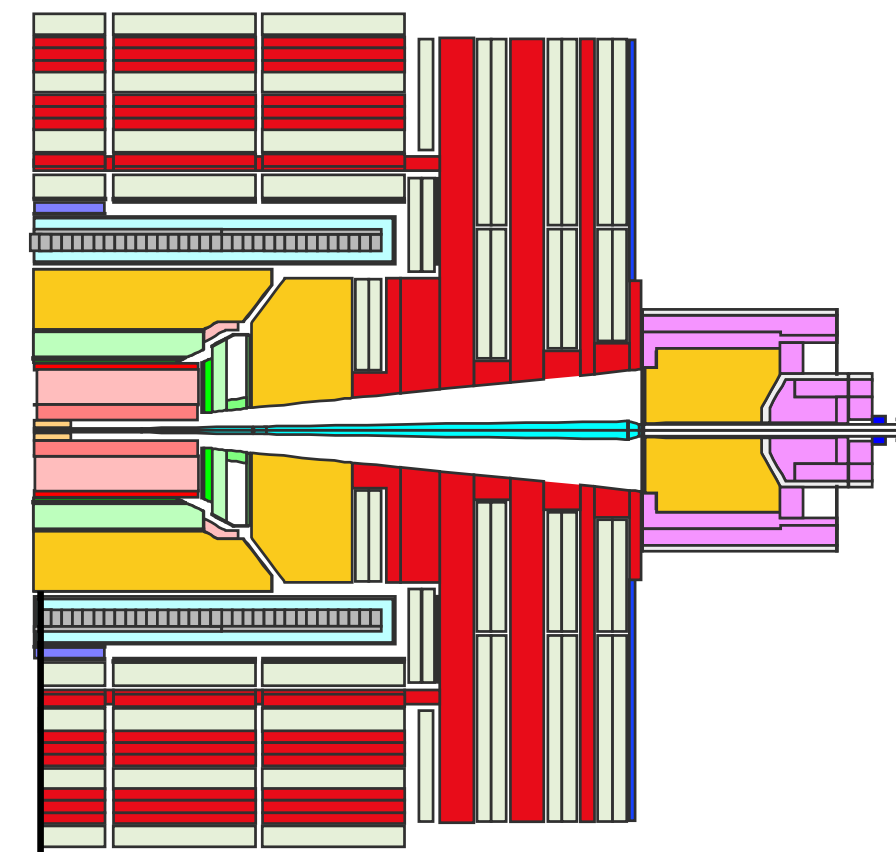


Using up to 37 fb<sup>-1</sup>  
of 2017 data

# Proton Reconstruction w/ PPS

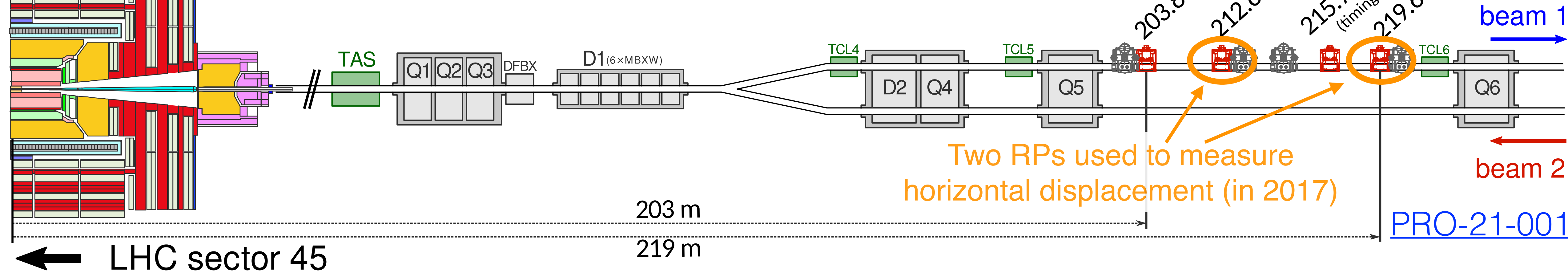
New for Run 2!

CMS central detector



LHC sector 56

Roman Pots

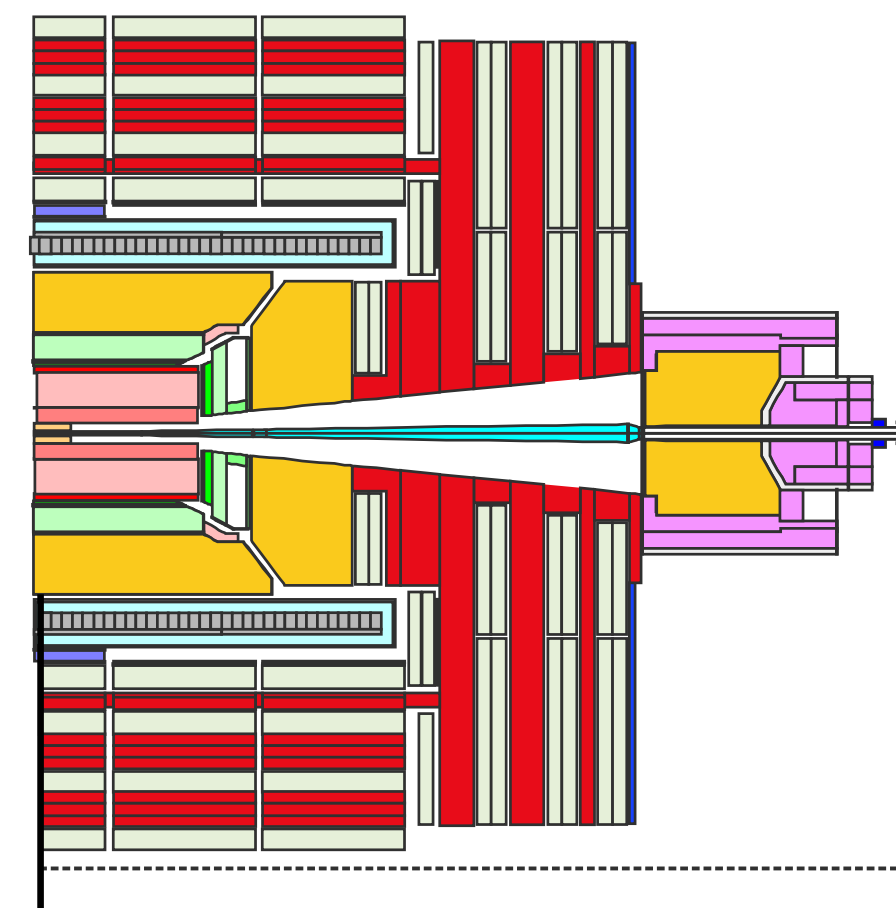


PPS allows CMS to measure **forward protons** using silicon tracking detectors

# Proton Reconstruction w/ PPS

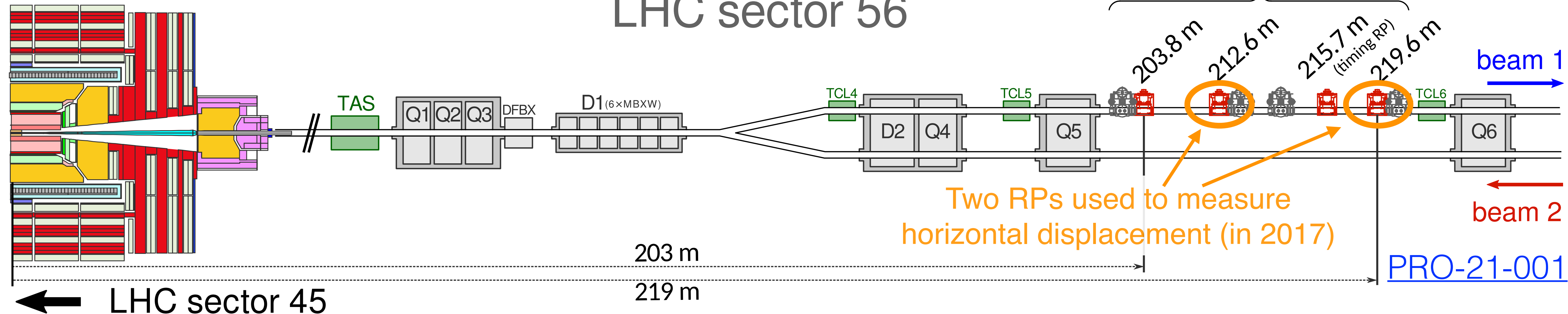
New for  
Run 2!

CMS  
central detector



LHC sector 56

Roman Pots



Two RPs used to measure  
horizontal displacement (in 2017)

[PRO-21-001](#)

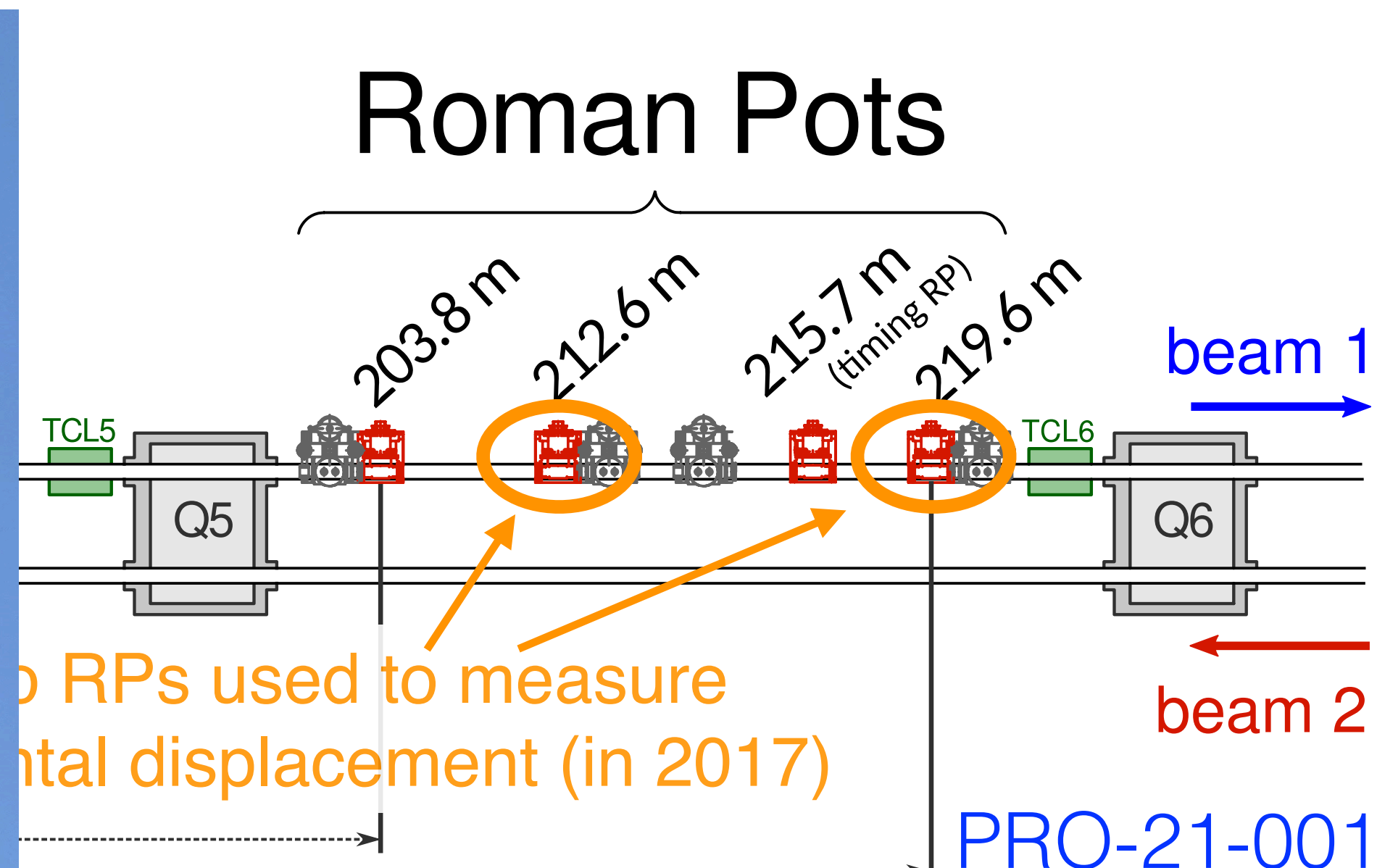
PPS allows CMS to measure **forward protons** using silicon tracking detectors

- Main measurement: proton momentum loss ( $\xi = \Delta p/p$ ) of outgoing protons from pp collisions w/ small scattering angles; sensitive to  $\xi \in [2\%, 20\%]$
- Calibration relies on precise understanding of LHC's  $\vec{B}$  field and beam optics

# Proton Reconstruction w/ PPS

New for  
Run 2!

Cathedral of Learning is “only” 163m tall!



← LHC sector 45

219 m

[PRO-21-001](#)

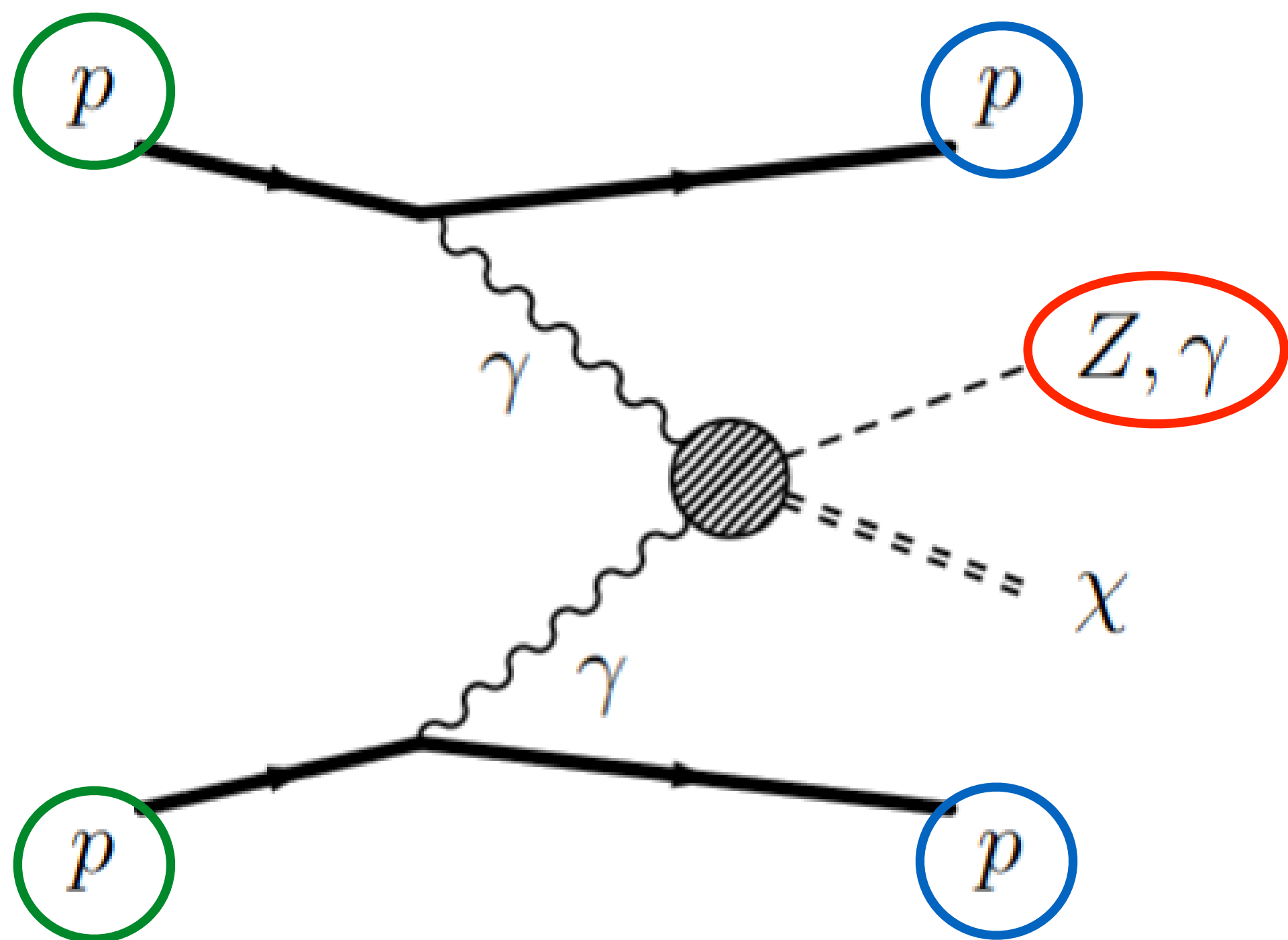
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# Search Strategy

Select events with  $Z \rightarrow \ell\ell$  or a photon; use outgoing proton momenta to reconstruct the **missing mass**:

$$m_{\text{miss}}^2 = \left[ \underbrace{(P_{p_1}^{\text{in}} + P_{p_2}^{\text{in}})}_{\text{LHC}} - \left( \underbrace{P_V}_{\text{CMS}} + \underbrace{P_{p_1}^{\text{out}} + P_{p_2}^{\text{out}}}_{\text{PPS}} \right) \right]^2$$



# Search Strategy

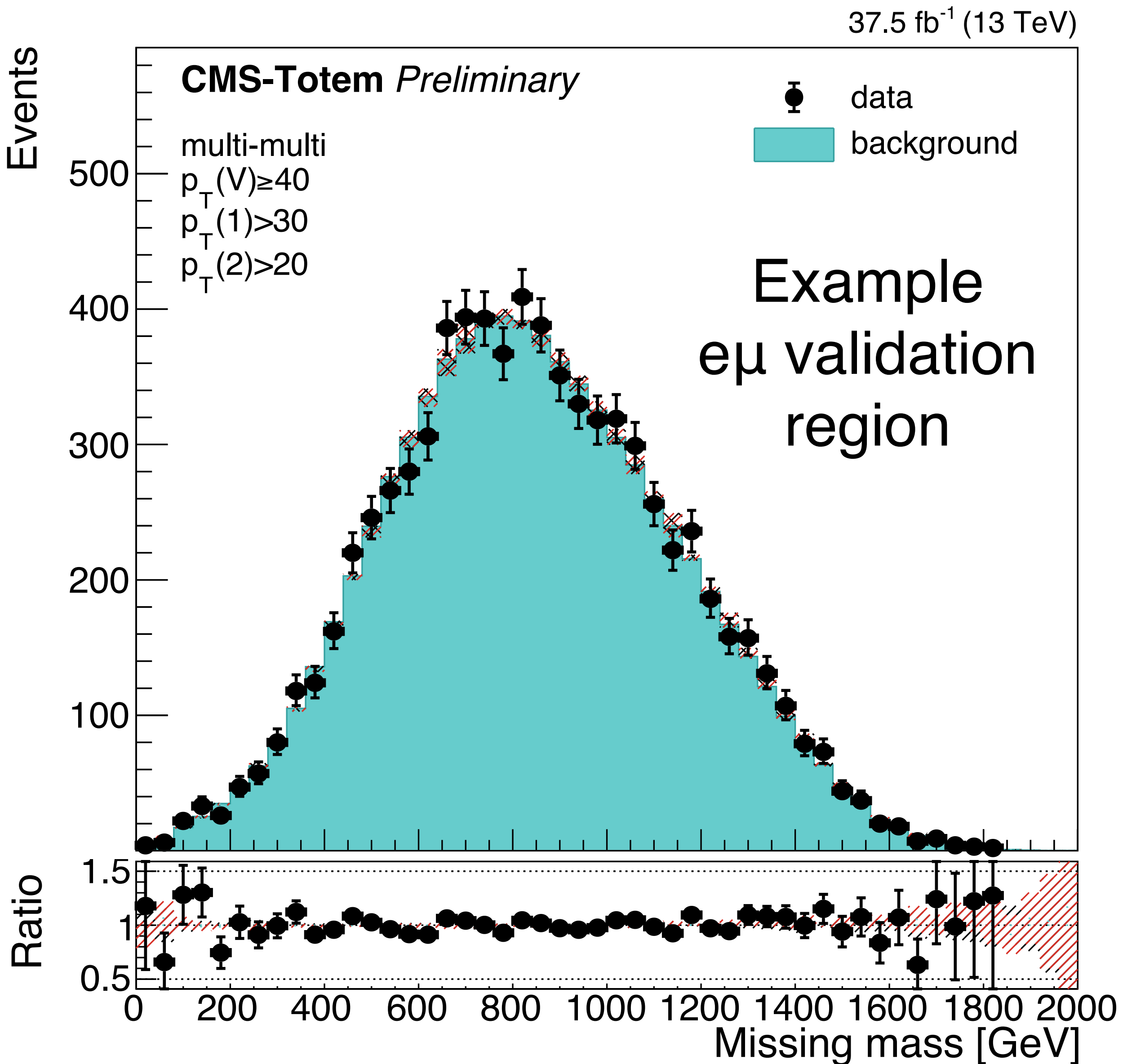
Select events with  $Z \rightarrow \ell\ell$  or a photon; use outgoing proton momenta to reconstruct the **missing mass**:

$$m_{\text{miss}}^2 = \left[ \underbrace{(P_{p_1}^{\text{in}} + P_{p_2}^{\text{in}})}_{\text{LHC}} - \left( \underbrace{P_V}_{\text{CMS}} + \underbrace{P_{p_1}^{\text{out}} + P_{p_2}^{\text{out}}}_{\text{PPS}} \right) \right]^2$$

- Search performed at high mass (600-1600 GeV), where PPS provides mass resolution of  $\sim 2\%$
- **First missing mass-based search at the LHC!**



# Background Estimation



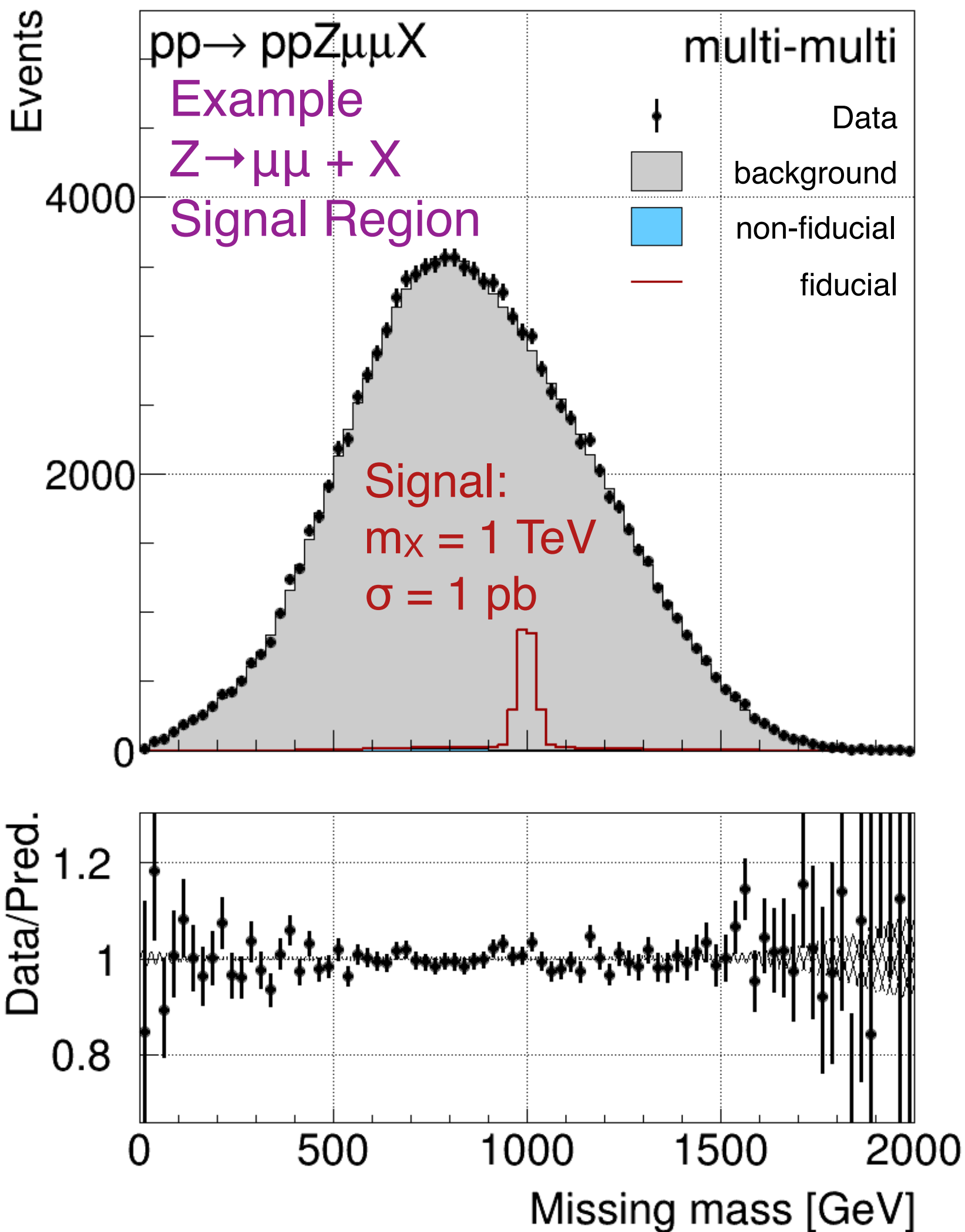
Primary background is combinatorial (coincidence of pileup protons)

Estimated via event mixing:

- **Replace protons** in data events with those from an **orthogonal dataset** of randomly chosen events w/ same beam conditions
- Validate in  $e\mu$  data events, and in simulated Z or  $\gamma$  events

# Results

## CMS-Totem Preliminary

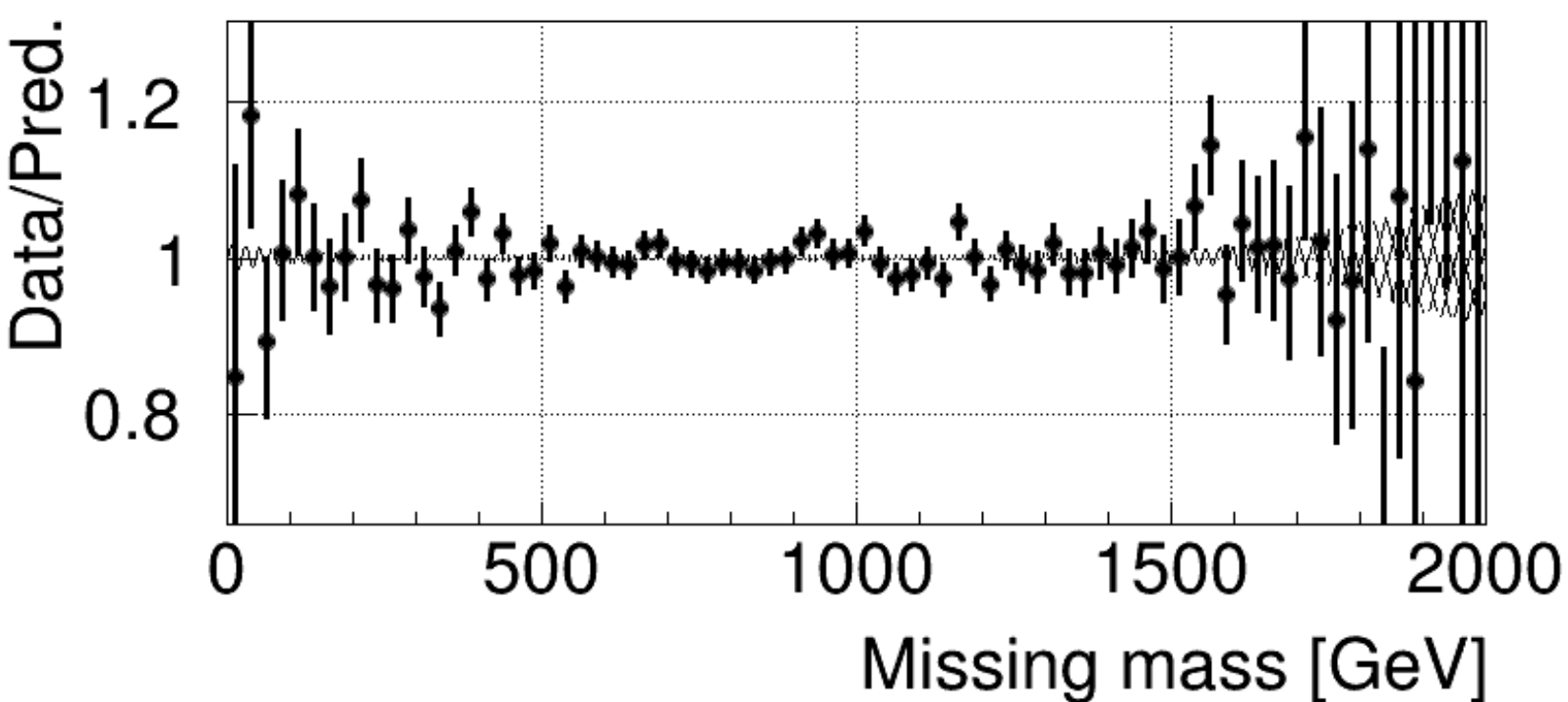
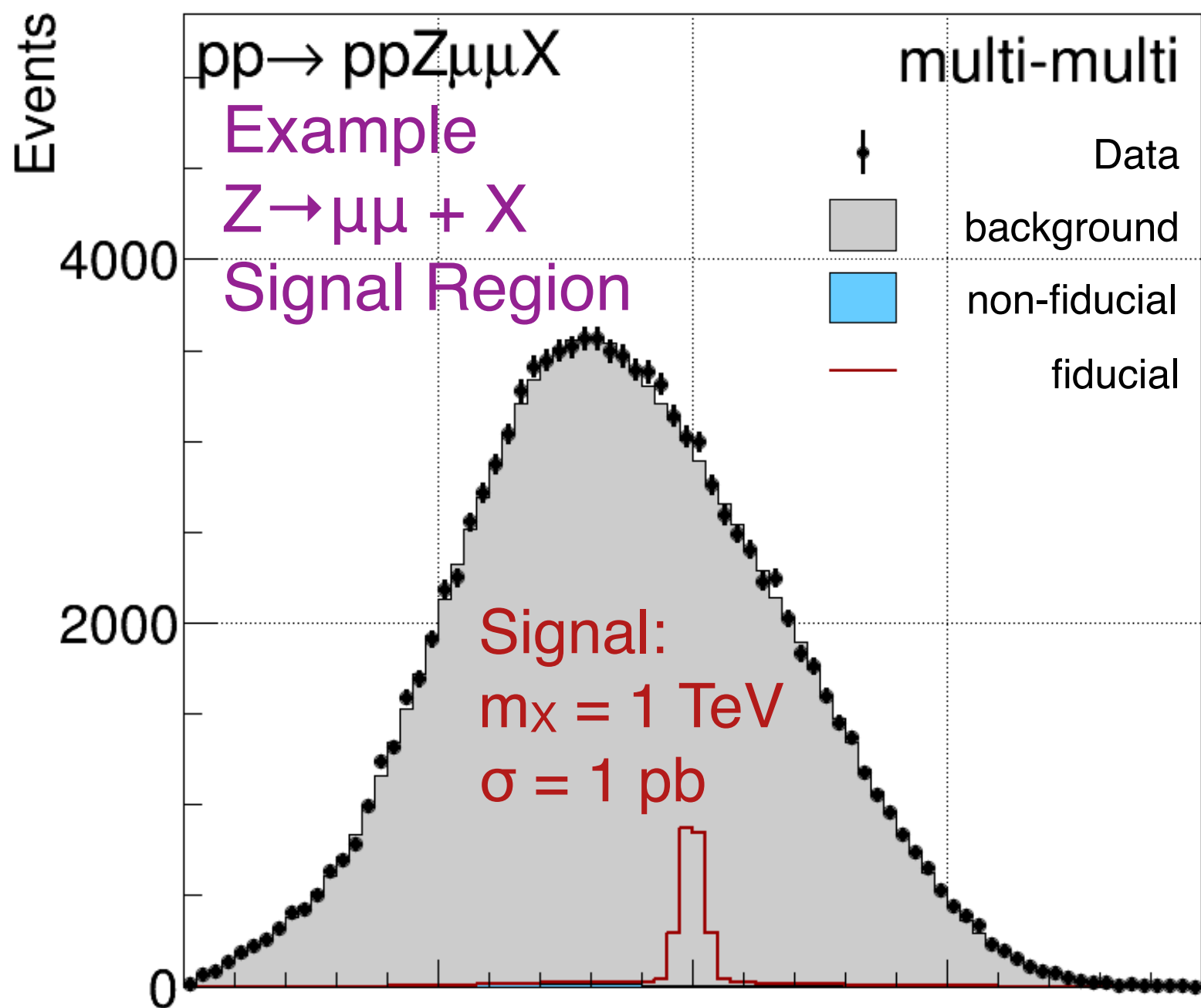


Fiducial selection: good quality, high  $p_T$   $Z \rightarrow \ell\ell$  or  $\gamma$ ; protons within PPS (details in backup).

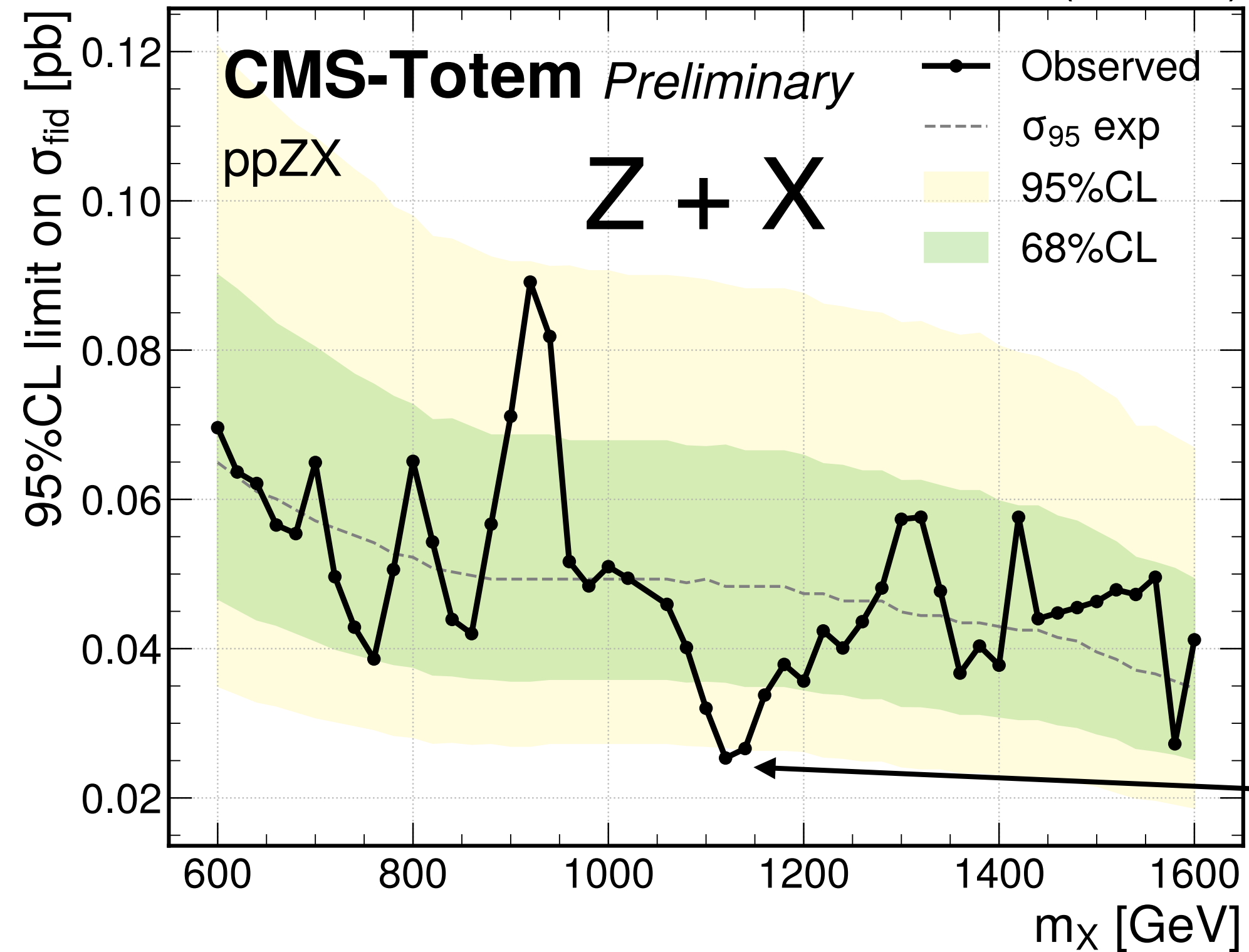
No significant excess in the missing mass distributions!

# Results

**CMS-Totem Preliminary**



37.2 fb<sup>-1</sup> (13 TeV)

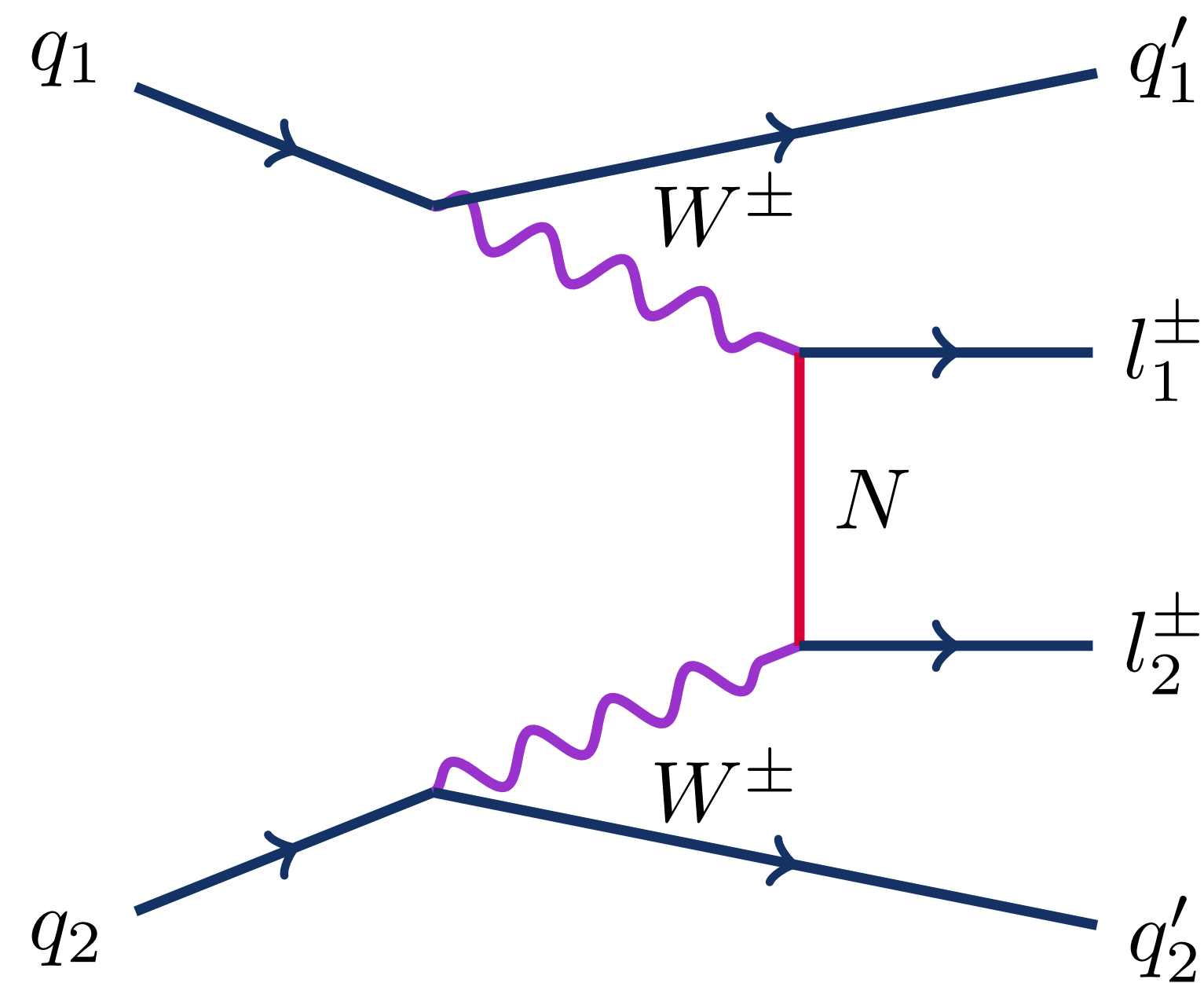


Set limits on  $\sigma_{\text{fid}}$  as low as 25 fb

Fiducial selection: good quality, high  $p_T$   $Z \rightarrow \ell\ell$  or  $\gamma$ ; protons within PPS (details in backup).

No significant excess in the missing mass distributions!

# Search for Majorana neutrinos and the Weinberg operator in VBF $\mu^\pm\mu^\pm$ events

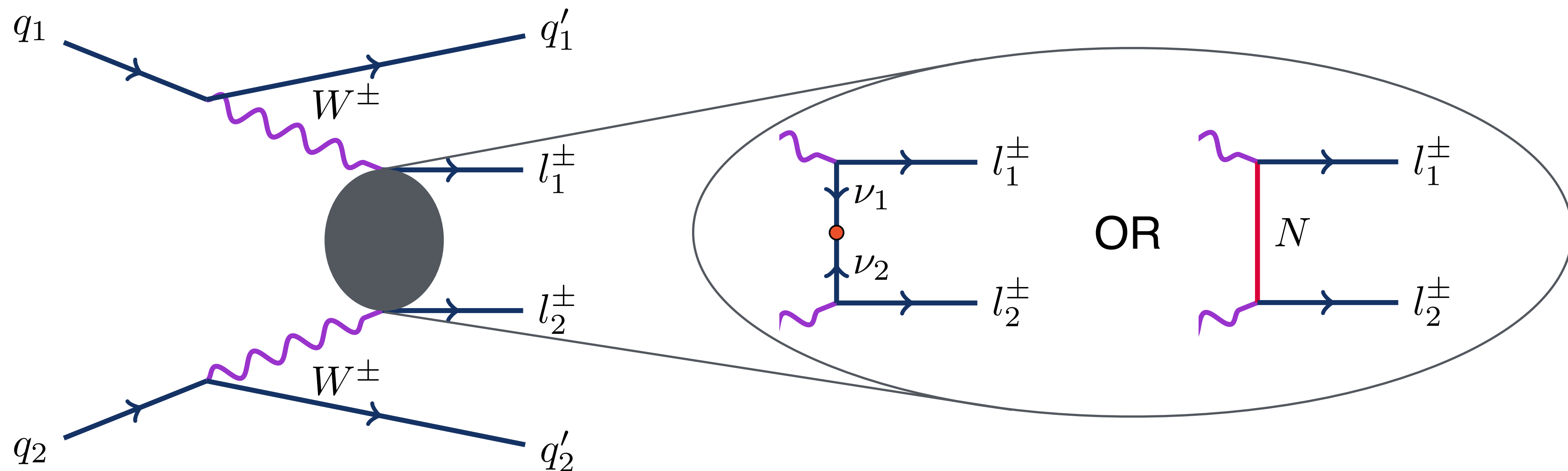


Using 138 fb<sup>-1</sup> of  
Run 2 data

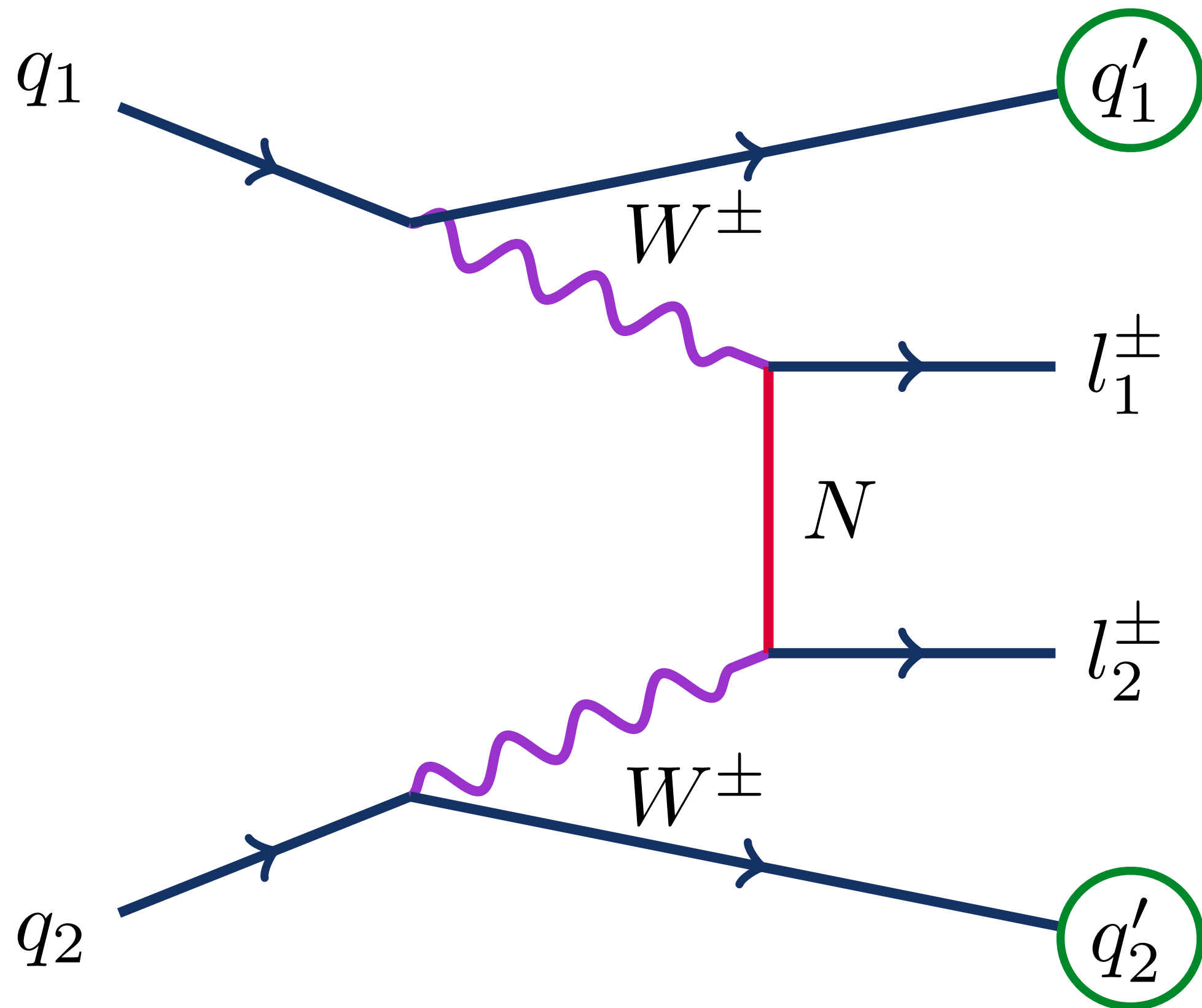
# Motivation

Aims to explain nonzero mass of neutrinos ( $\nu$ ) via either:

- EFT w/ dim-5 [Weinberg operator](#); generates  $m_\nu$  *and* leads to lepton-number violating  $0\nu\beta\beta$  decay
- Heavy Majorana neutrino ( $N$ ) in type-I seesaw model



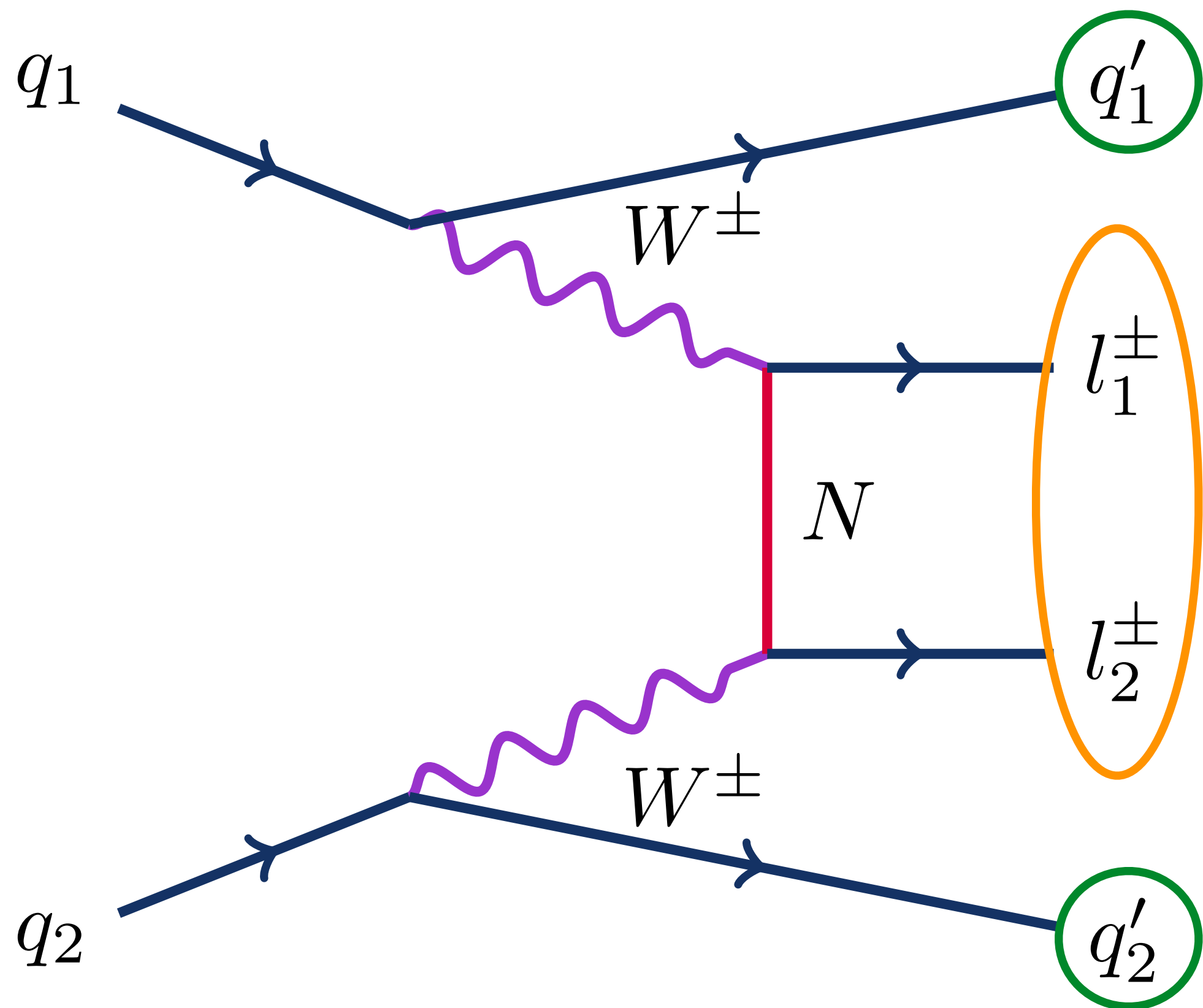
# Signal vs. Background Discrimination



Suppress background by requiring:

- **Forward VBF jet kinematics**  
(large  $m_{jj}$ ,  $\Delta\eta_{jj}$ , etc.)

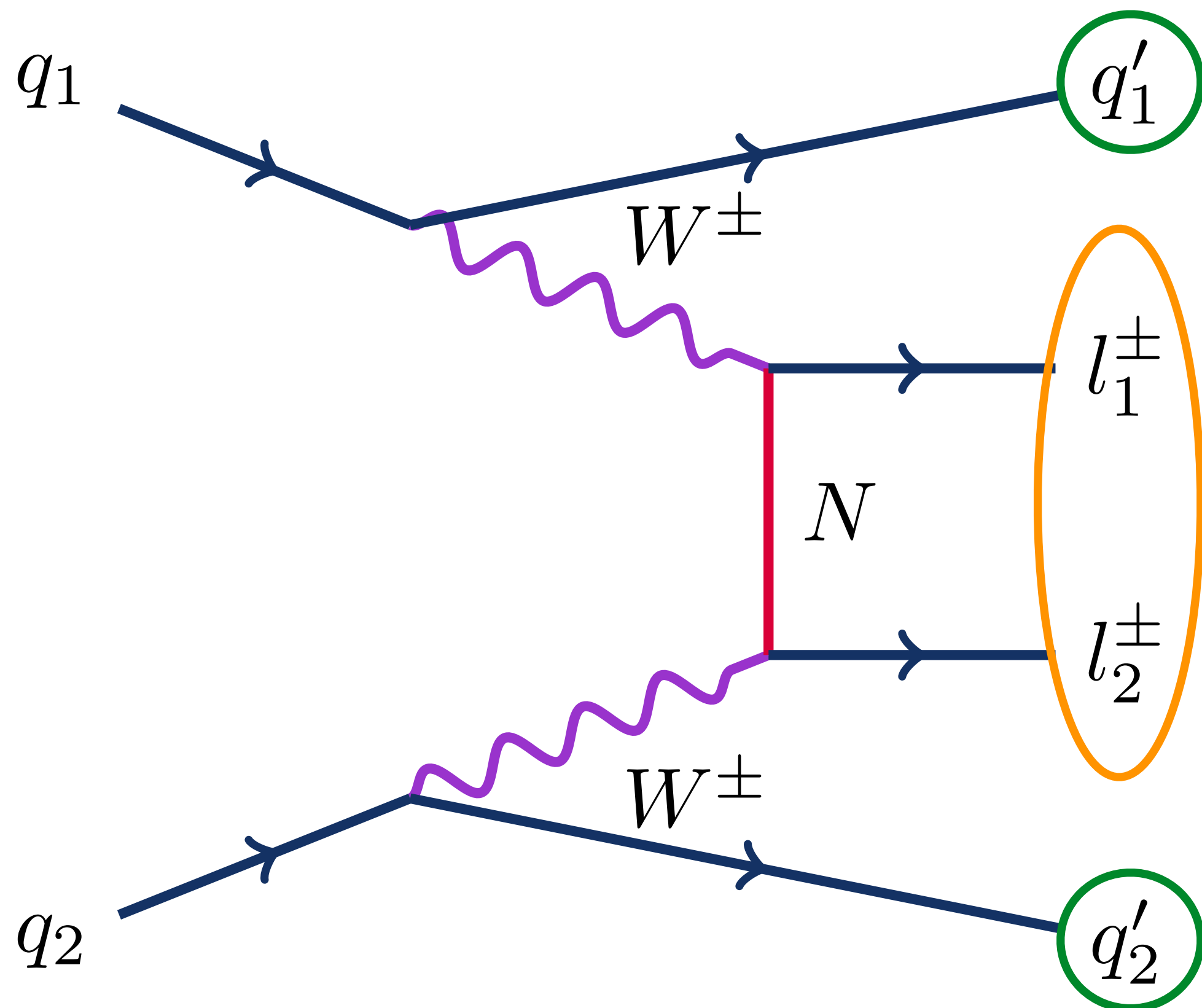
# Signal vs. Background Discrimination



Suppress background by requiring:

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- **Same-sign dimuons** (rare in SM)

# Signal vs. Background Discrimination

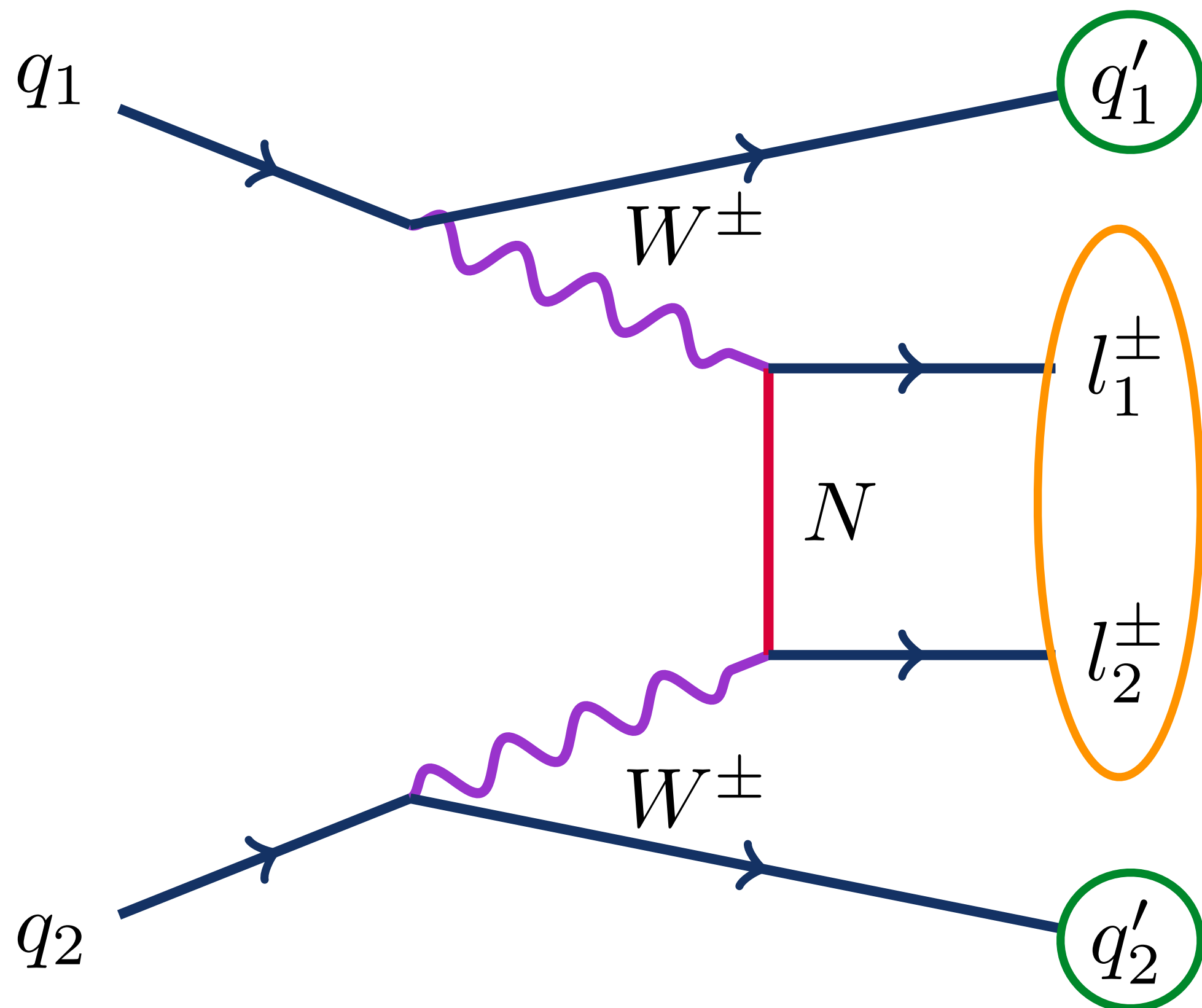


Suppress background by requiring:

- **Forward VBF jet kinematics** (large  $m_{jj}$ ,  $\Delta\eta_{jj}$ , etc.)
- **Same-sign dimuons** (rare in SM)
- Small jet activity ( $H_T = \sum \text{jet } p_T$ ) relative to leptonic activity: search is binned in  $H_T / p_T(\mu_1)$



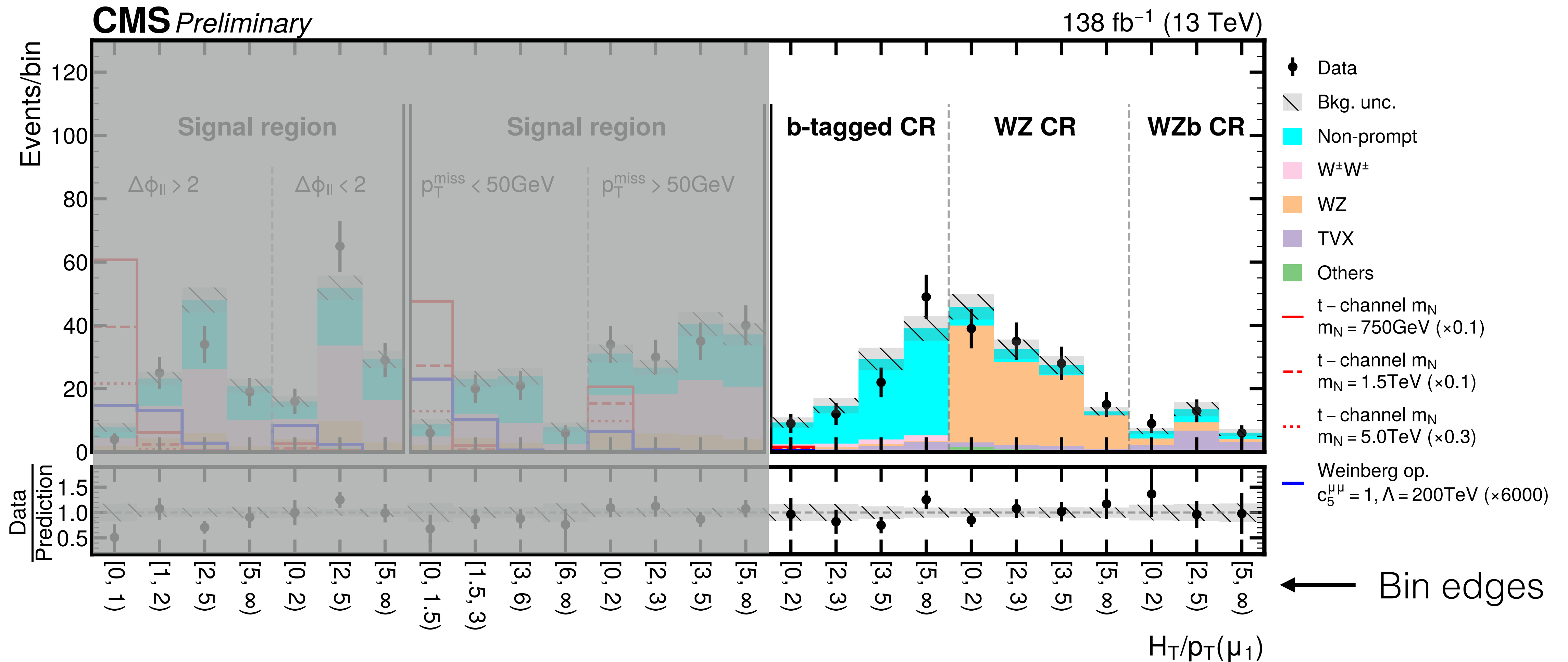
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Suppress background by requiring:

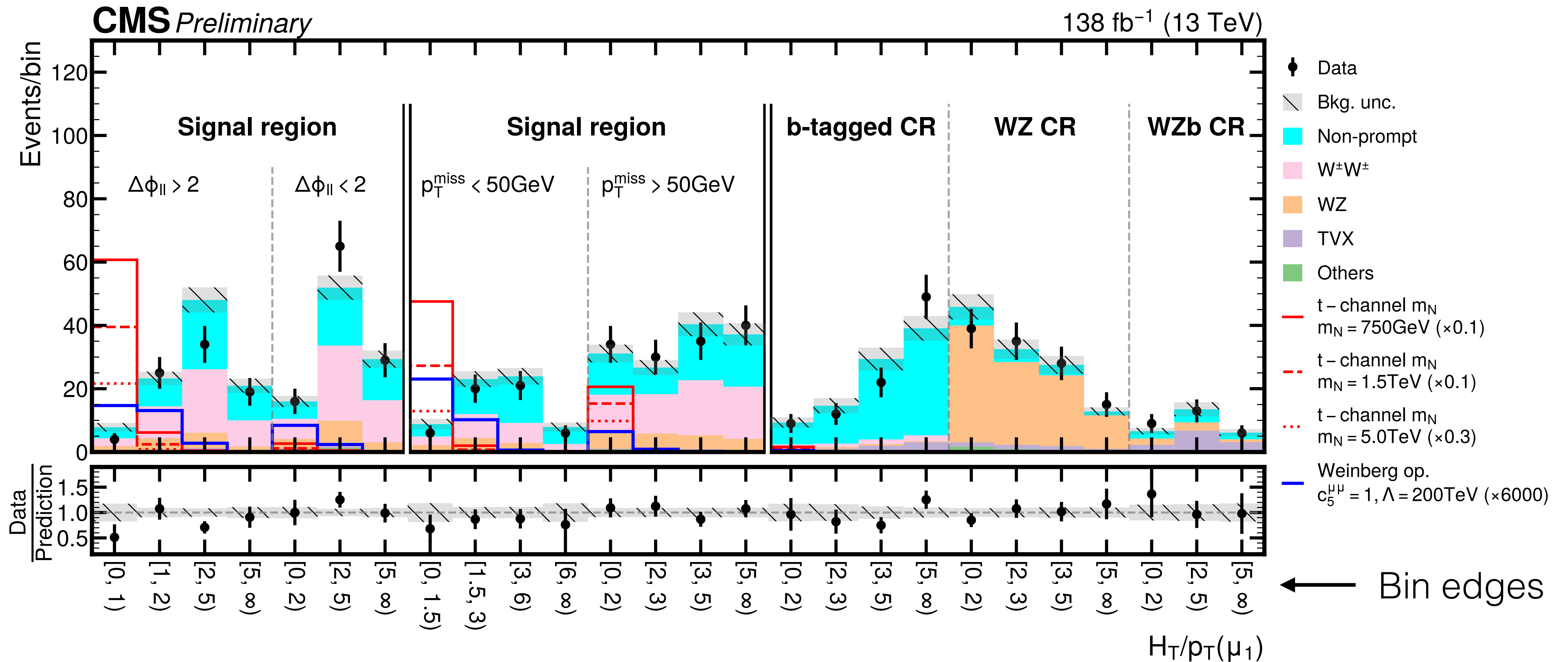
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- Small jet activity ( $H_T = \sum \text{jet } p_T$ ) relative to leptonic activity: search is binned in  $H_T / p_T(\mu_1)$
- Large  $\Delta\phi_{\mu\mu}$
- Small  $p_T^{\text{miss}}$  (no final state neutrinos!)

# Background Estimation



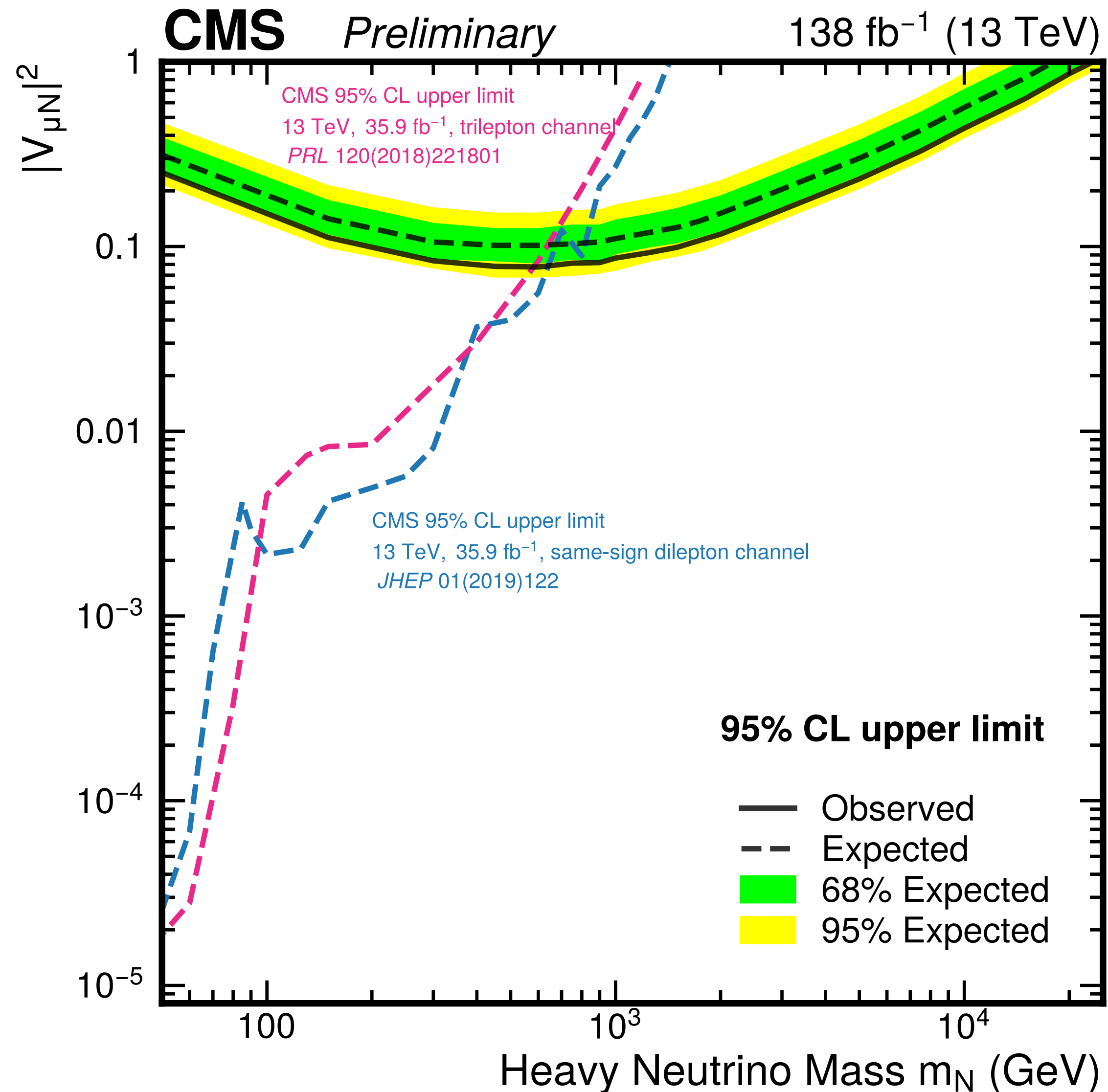
Use data control regions to estimate backgrounds in VBF-enriched regions, primarily due to fake/non-prompt leptons and dibosons.

# Results



Signal tends to have small  $H_T/p_T(\mu_1)$ , small MET, and large  $\Delta\phi_{\mu\mu}$  relative to backgrounds.  
 Data agrees well with SM prediction  $\Rightarrow$  set limits on these models.

# Results



Outperforms previous heavy neutrino searches for  $m_N > 650$  GeV, and excludes masses up to 23 TeV for  $|V_{\mu N}|^2 = 1$ .

For the Weinberg operator EFT, the search sets limits on the effective Majorana mass up to  $|m_{\mu\mu}| = 10.84$  GeV.

⇒ first constraints on this process!

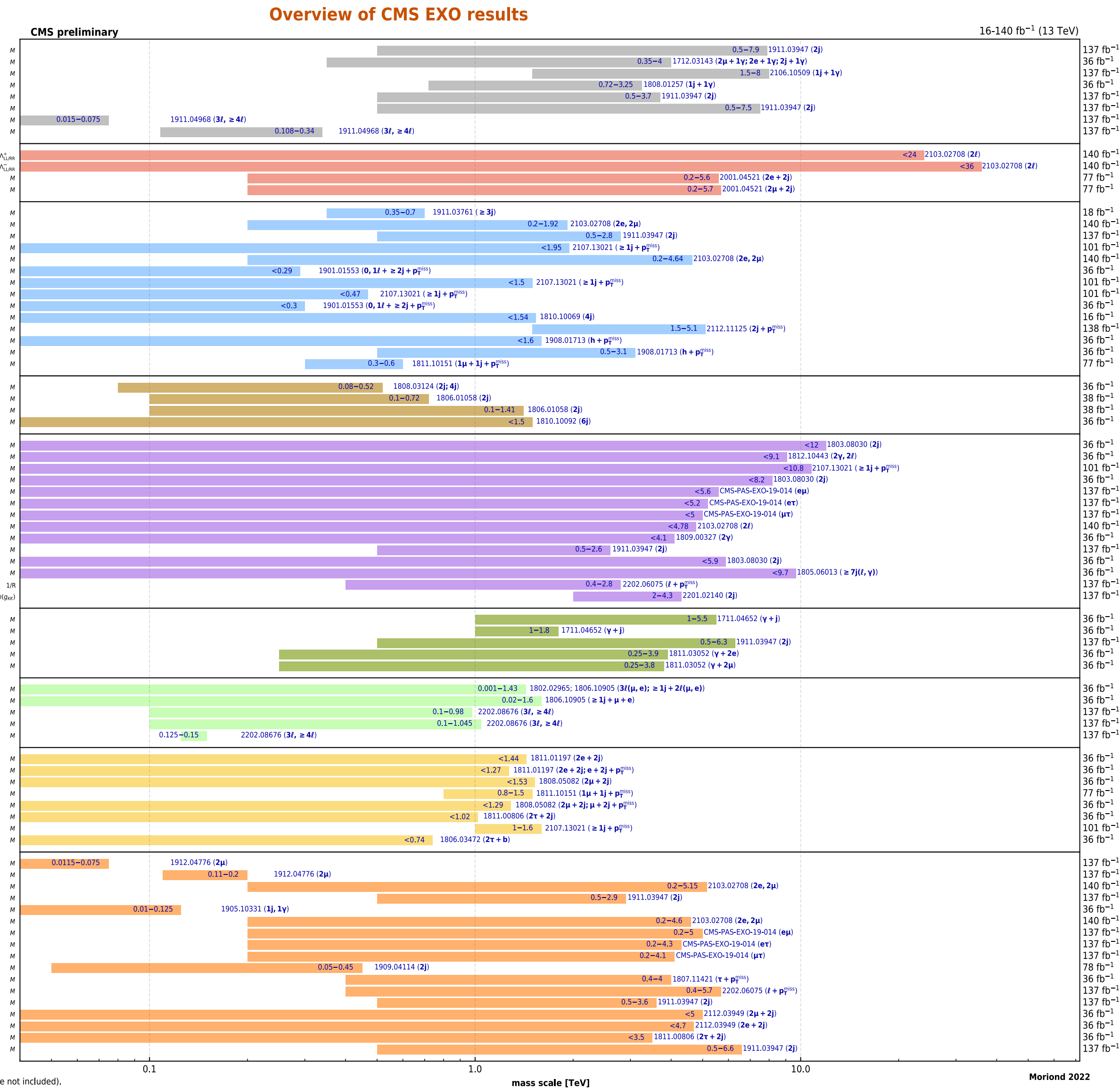
# Summary

No evidence for new physics in these leptonic / photonic searches yet, but CMS has set strong constraints across many final states.

However, many new ideas are still being explored!

- Maybe new physics is hiding in some exotic area of phase space that has not yet been probed
- Or perhaps we haven't collected enough data yet! Run 3 and the HL-LHC data will change that
  - Run 2 was *less than 5%* of the total expected LHC dataset!

# CMS EXO Summary Plot ([link](#))



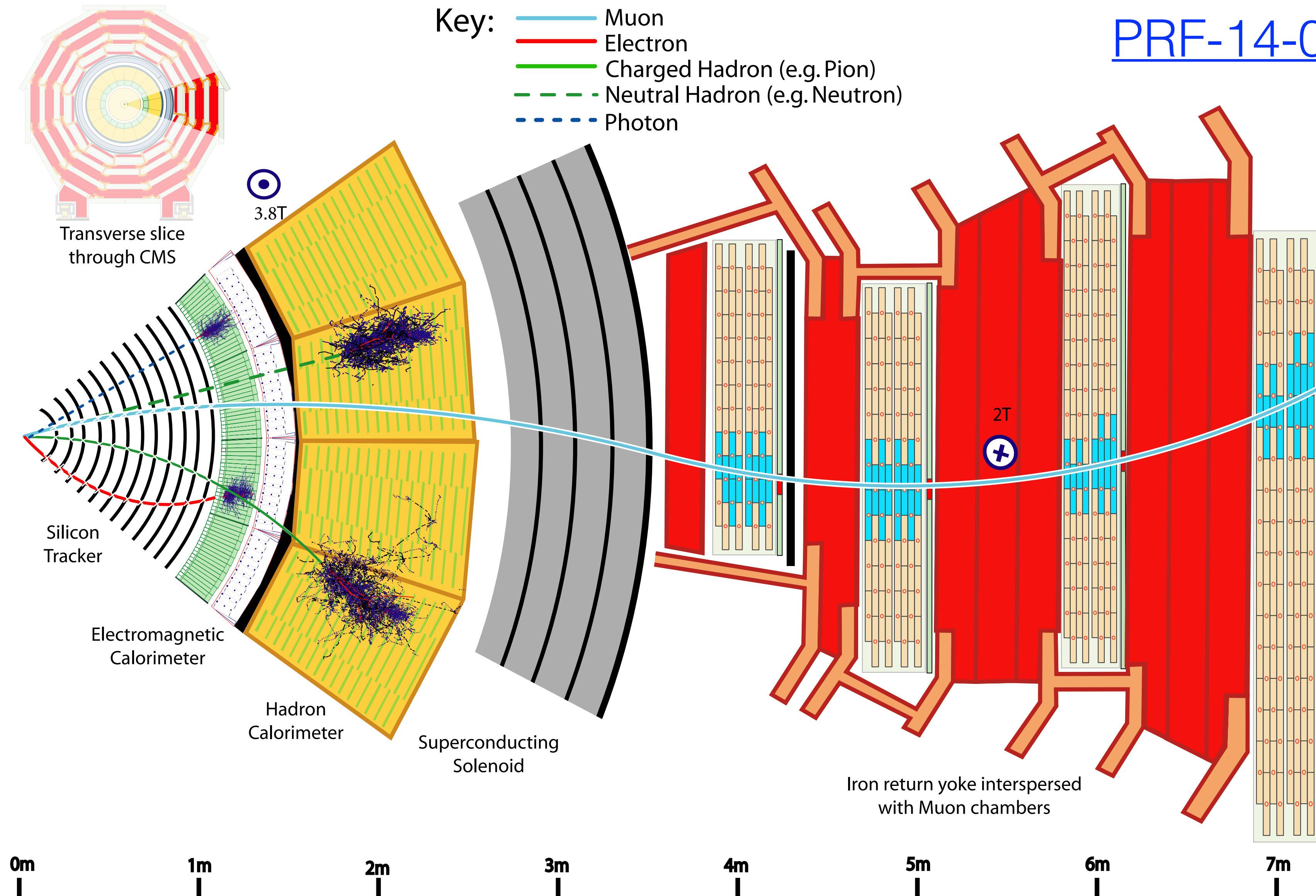
Other new lepton- or photon-based CMS searches since Pheno 2021:

- LFV resonance search in  $e\mu$ ,  $e\tau$ , and  $\mu\tau$  events ([EXO-19-014](#))
- Inclusive nonresonant multilepton search ([EXO-21-002](#))

# Backup

# The CMS Detector

[PRF-14-001](#)



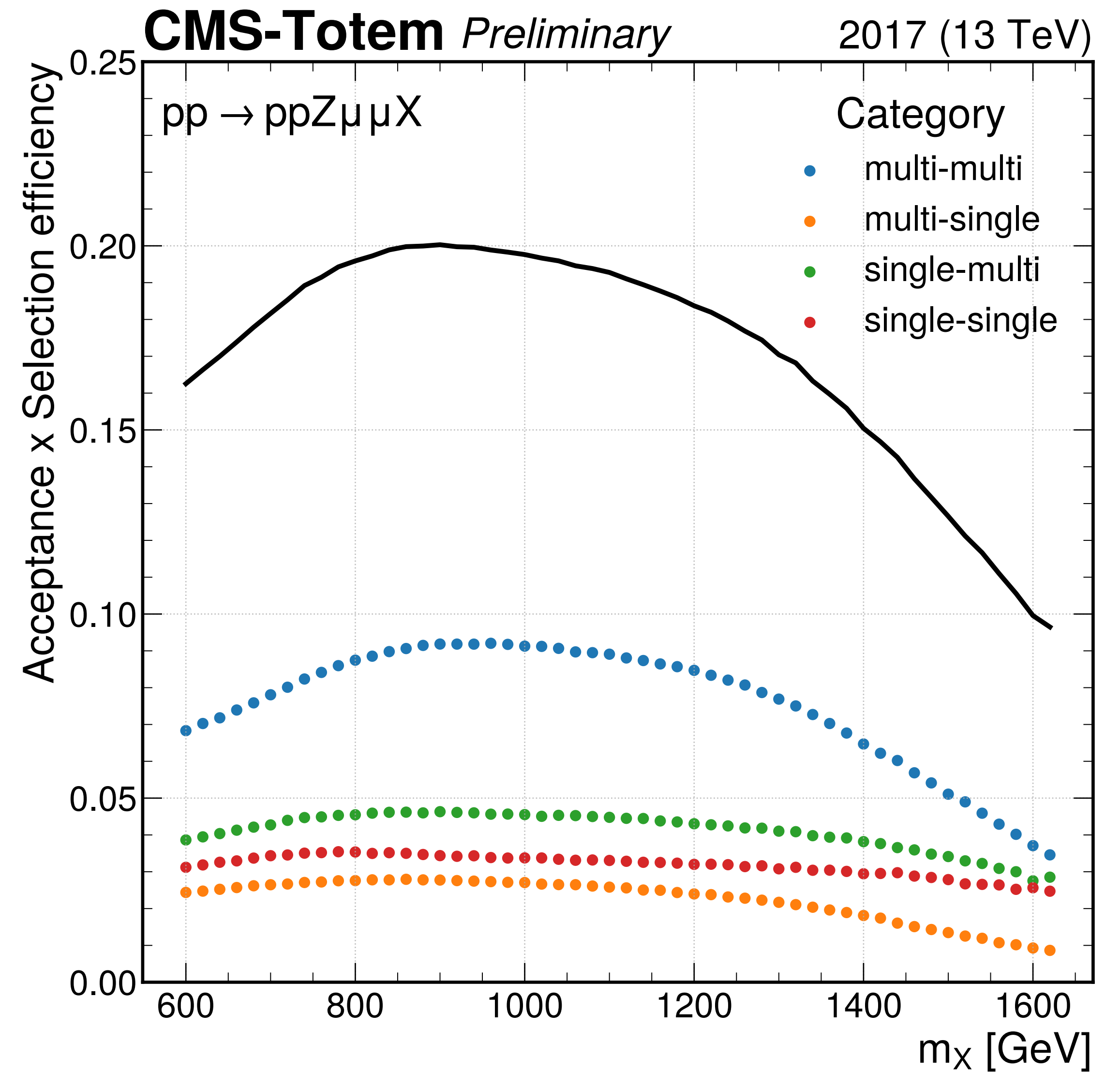
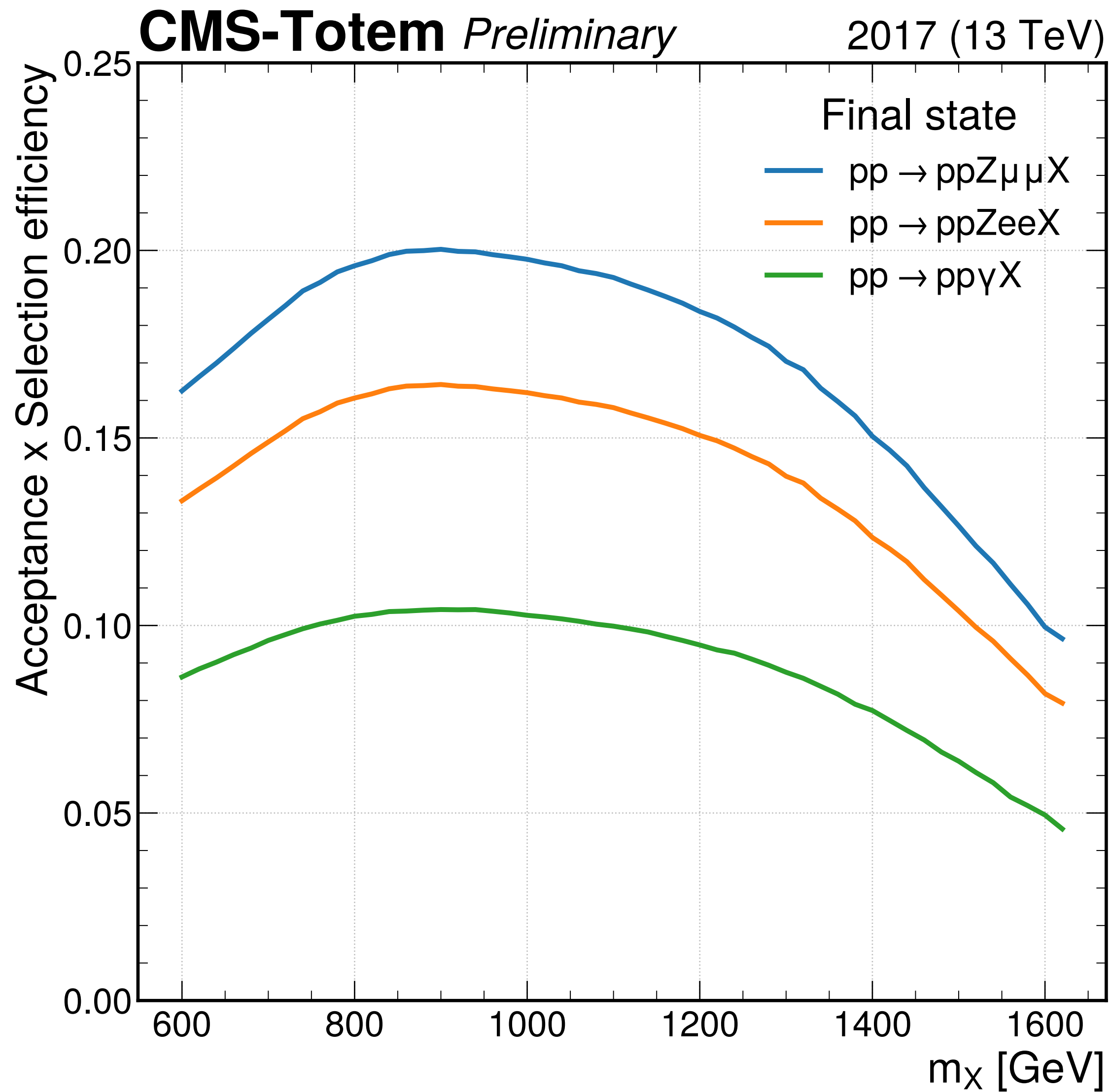


# EXO-19-009

Fiducial requirements:

Selection / Analysis	$Z \rightarrow e^+e^- / Z \rightarrow \mu^+\mu^-$	$\gamma$
Leptons/Photons	$\geq 2$ same flavour leptons (e or $\mu$ ) opposite electric charge $p_T(\ell_1, \ell_2) > 30, 20$ GeV $ \eta(\ell)  < 2.4$ $ m(\ell_1\ell_2) - m_Z  < 10$ GeV	=1 $\gamma$ within $ \eta(\gamma)  < 1.4442$
Boson $p_T$	$p_T(Z) > 40$ GeV	$p_T(\gamma) > 95$ GeV
Protons	$0.02 < \zeta_+^{gen} < 0.16 \quad \wedge \quad 0.03 < \zeta_-^{gen} < 0.18$	

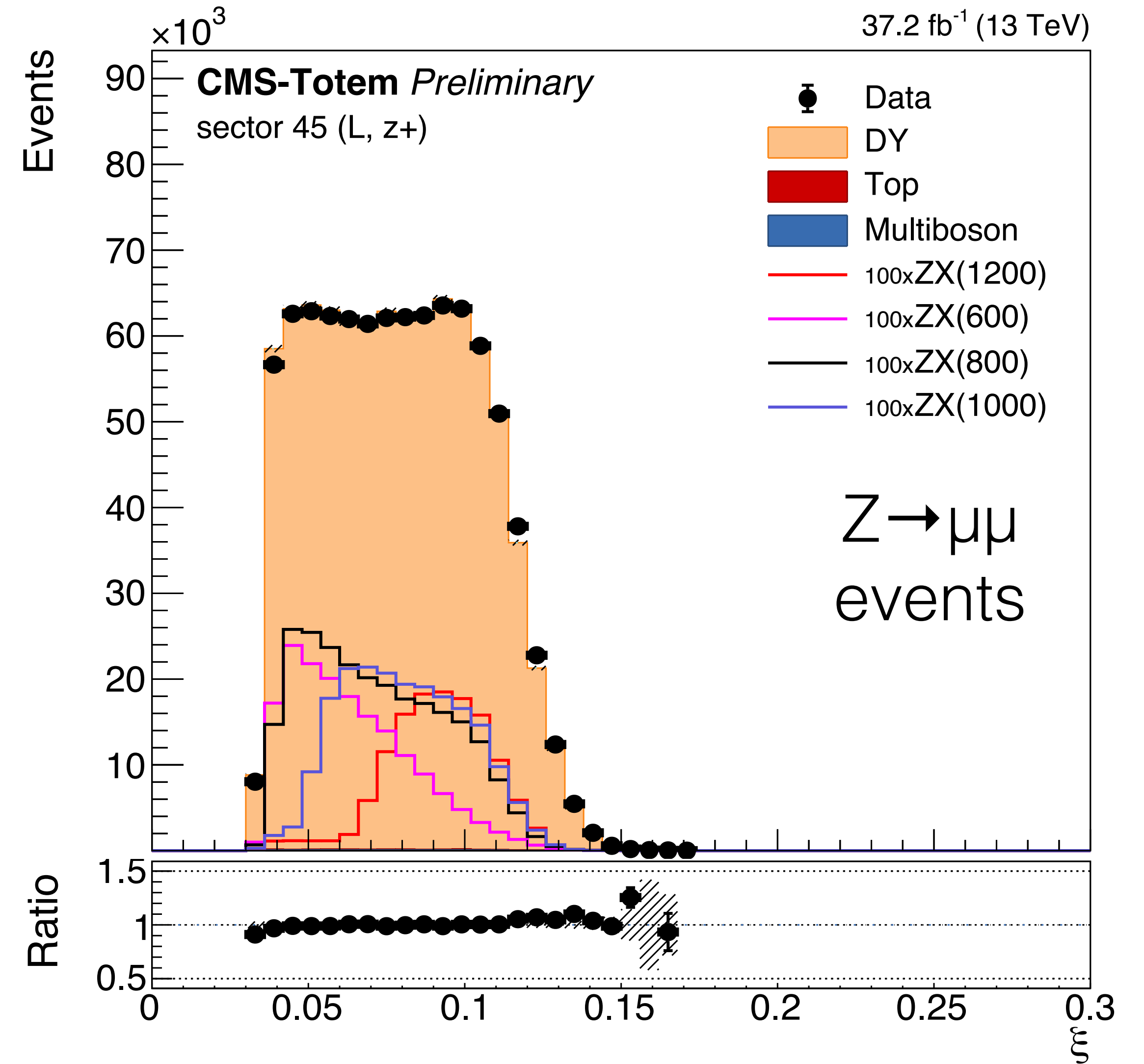
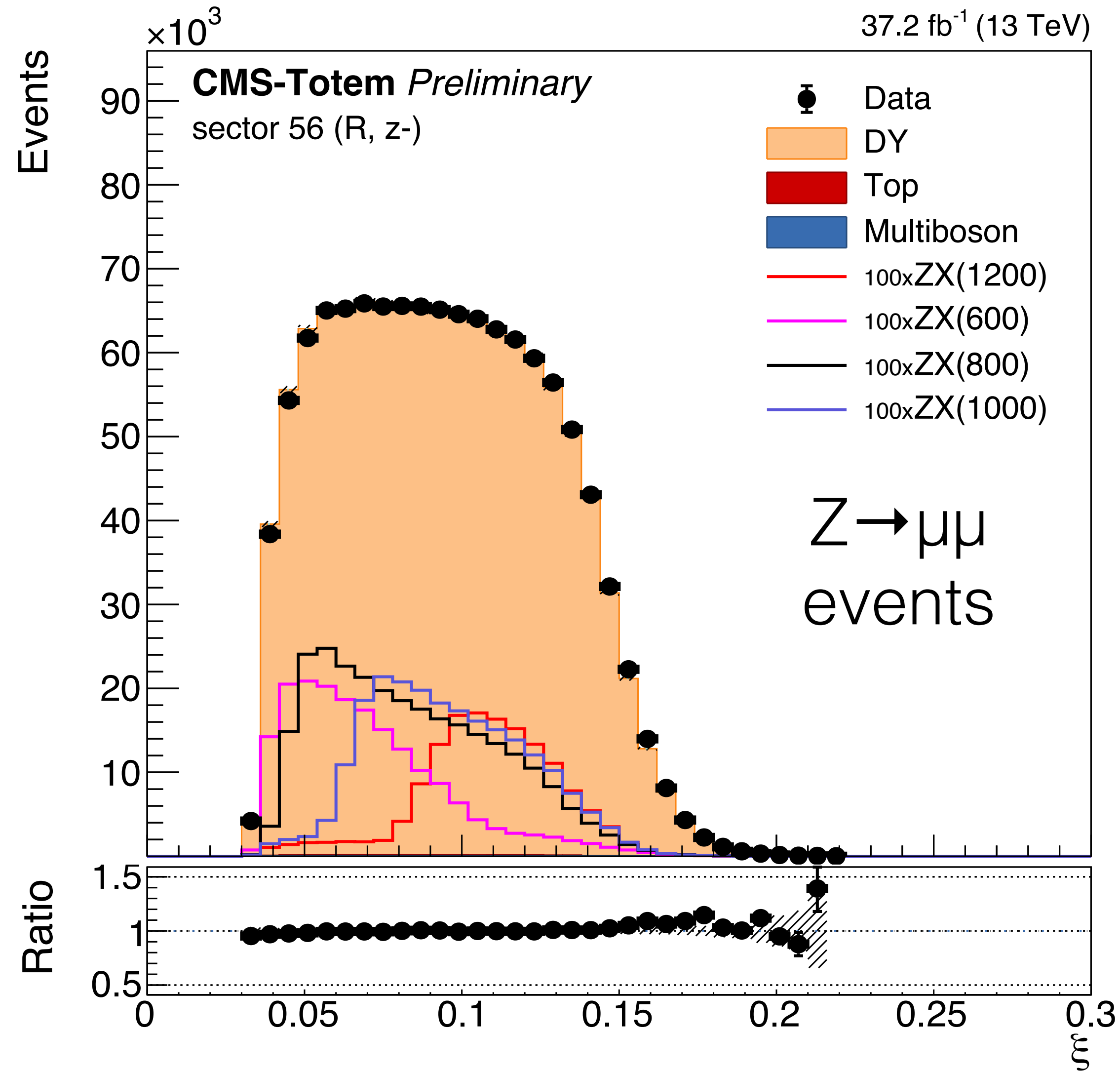
# EXO-19-009



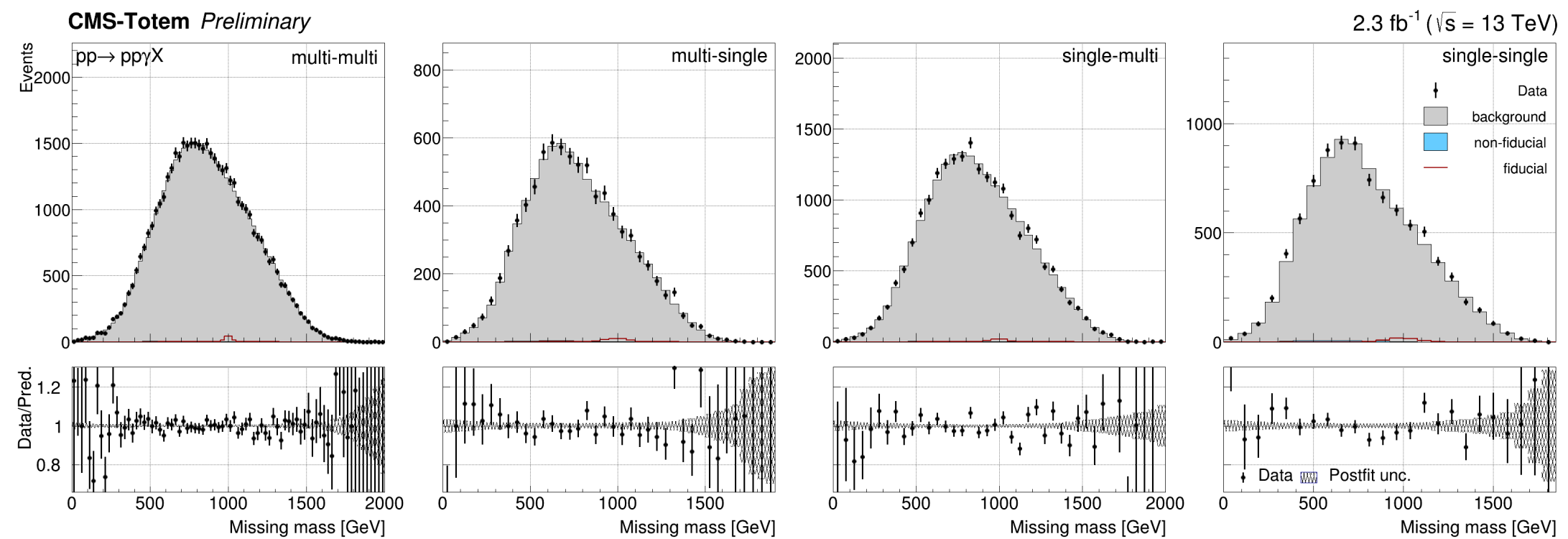
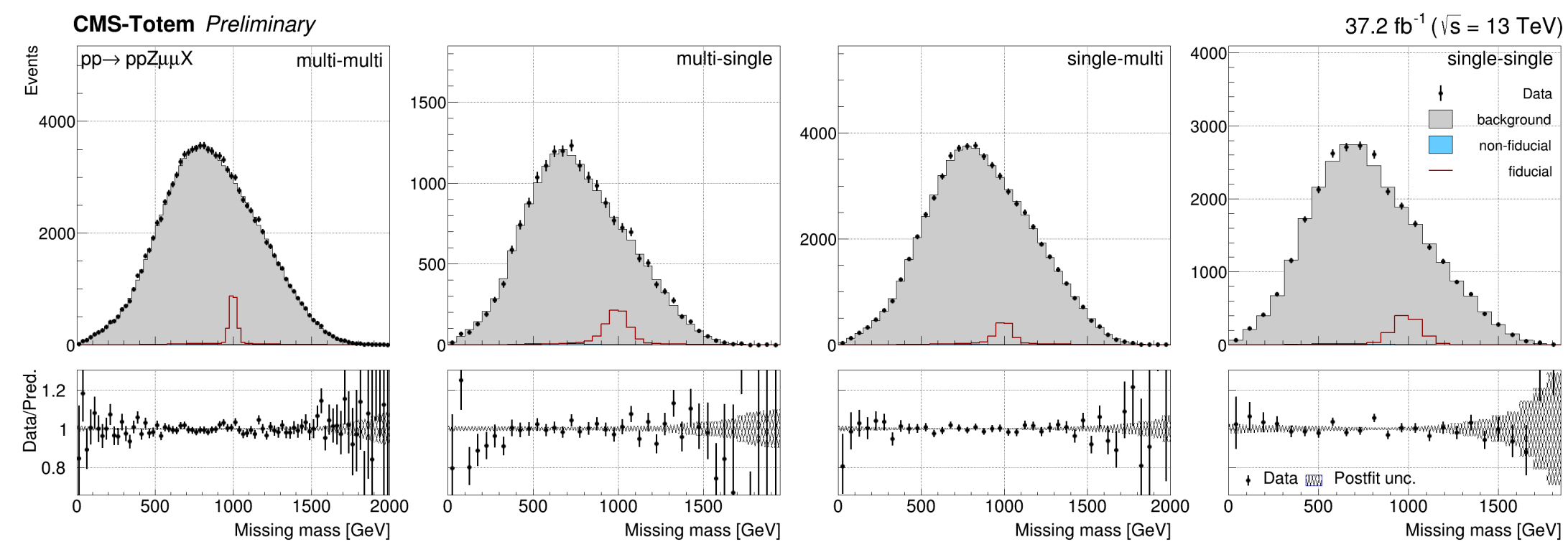
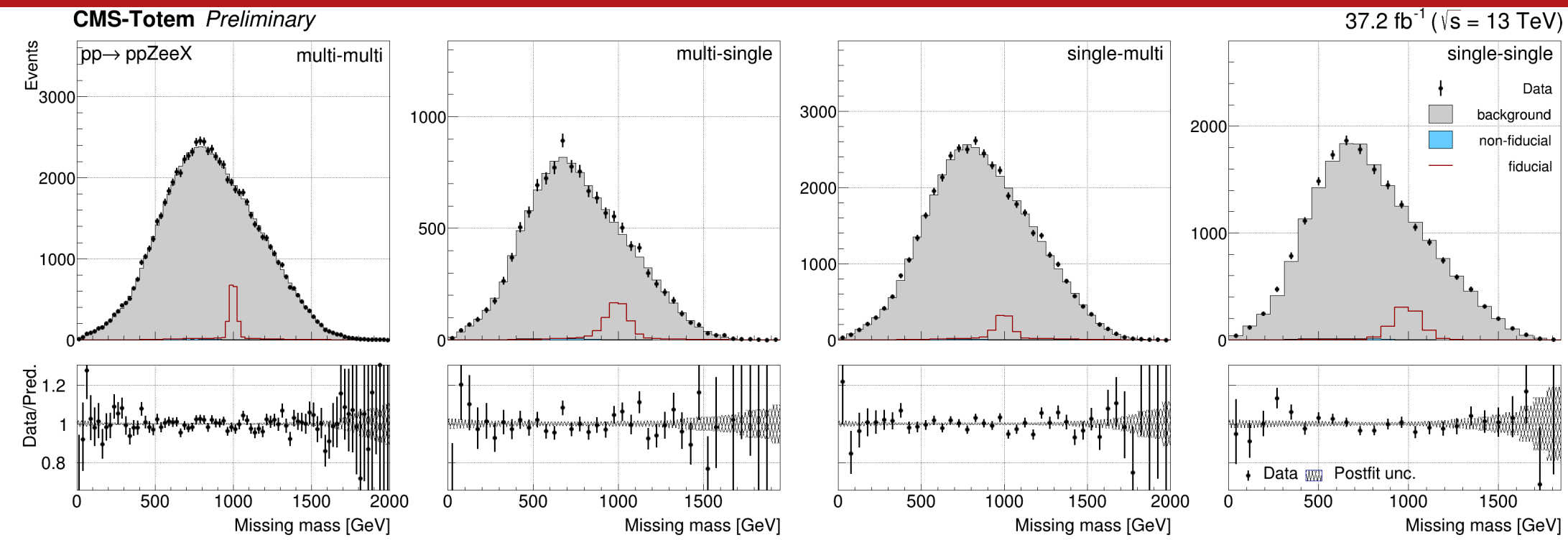
# EXO-19-009

The tracking stations provide a measurement of the proton trajectories with respect to the beam position. Knowledge of the magnetic fields traversed by the proton from the IP to the RPs allows for the reconstruction of its fractional momentum loss  $\tilde{\zeta} = \Delta p / p \sim D_x^{-1} \cdot x$  (with respect to the momentum of incident proton), where  $x$  represents the horizontal displacement of the scattered proton at the RP location, and  $D_x$  the horizontal dispersion, a property of the accelerator optics. The techniques used for the alignment and calibration of the apparatus and proton reconstruction are detailed in Refs. [13–16].

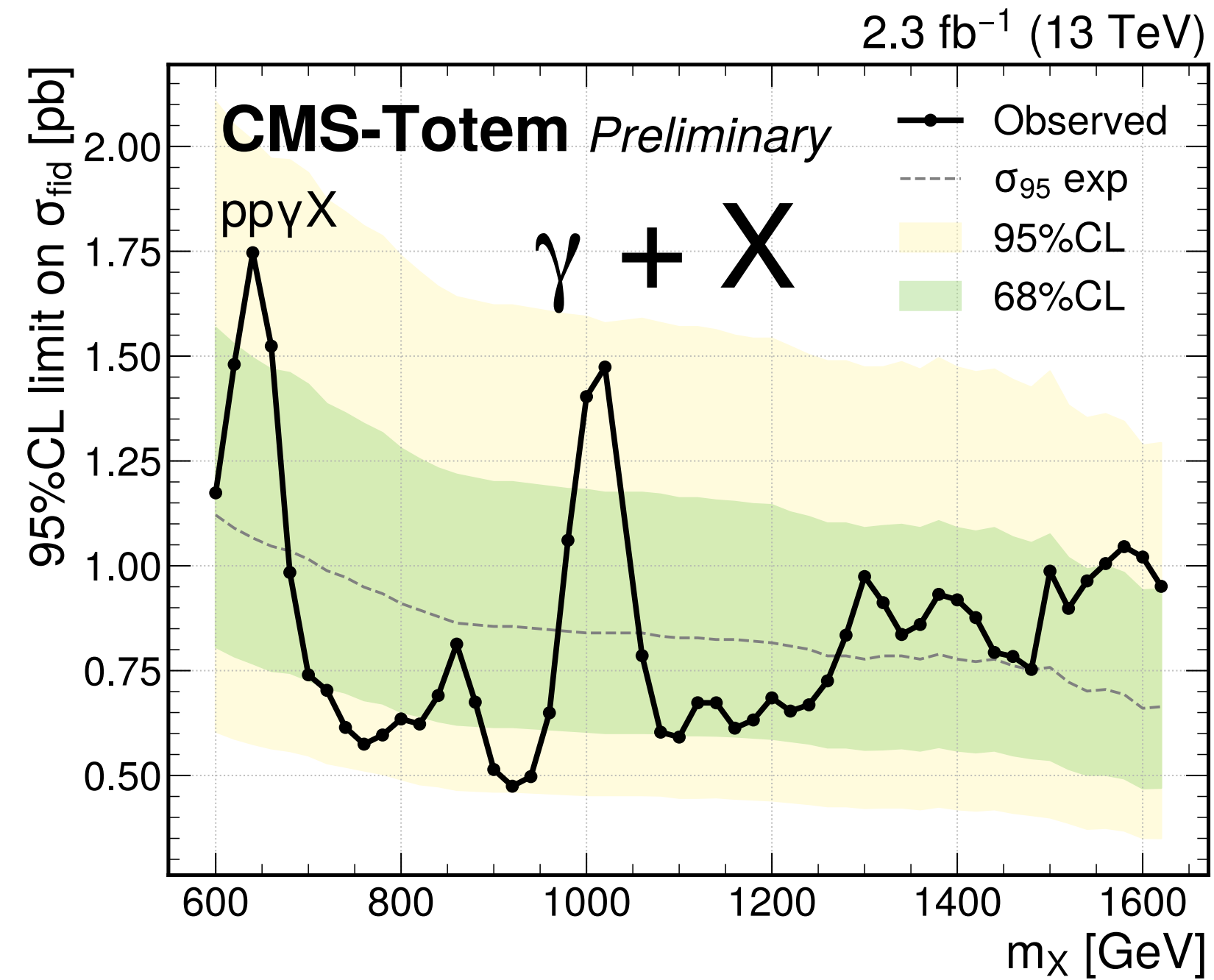
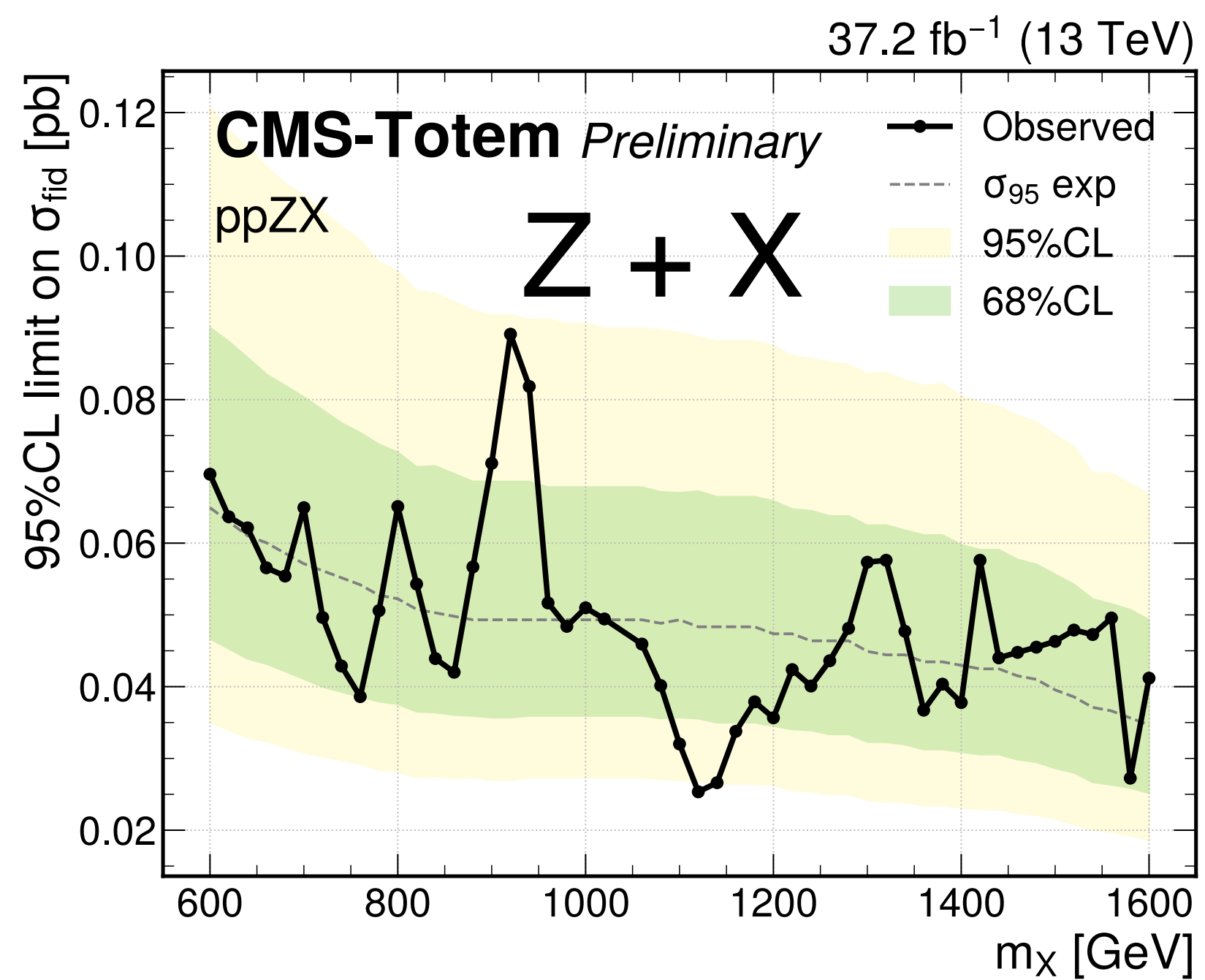
# EXO-19-009



# EXO-19-009



# EXO-19-009



# EXO-21-003

SR and CR require two jets w/  $p_T > 30$  GeV,  $|\Delta\eta_{jj}| > 2.5$ ,  $m_{jj} > 750$  GeV

Variable	SR	WZ CR
lepton $p_T$	$> 30/30$ GeV	$> 10/25/25$ GeV
dilepton mass $m_{\ell\ell}$	$> 20$ GeV	—
$ m_{\ell\ell} - m_Z $	—	$< 15$ GeV (for one $\mu^+\mu^-$ pair)
trilepton mass $m_{\ell\ell\ell}$	—	$> 100$ GeV
$\max \mathcal{Z}(\ell)$	$< 0.75$	$< 1.0$
$p_T^{\text{miss}}$	—	$> 30$ GeV

with Zeppenfeld variable: 
$$\mathcal{Z}(\ell) = \left| \eta_\ell - \frac{\eta_{j1} + \eta_{j2}}{2} \right| / \left| \Delta\eta_{jj} \right|$$

Events w/ additional e,  $\mu$ , or  $\tau$  leptons, or b-tagged jets, are vetoed.

# EXO-21-003

The simplest formalism in which neutrino masses can arise is through a dimension-5 operator as shown by Weinberg [2]:  $\mathcal{L}_5 = C_5^{\ell\ell'} / \Lambda [\Phi \cdot \bar{L}_\ell^c] [L_{\ell'} \cdot \Phi]$ , where  $\Lambda$  is the scale at which the particles responsible for neutrino masses become relevant degrees of freedom;  $C_5^{\ell\ell'}$  is a flavor-dependent Wilson coefficient;  $L_\ell^T = (\nu_\ell, \ell)$  is the left-handed lepton doublet; and  $\Phi$  is the SM Higgs doublet with a vacuum expectation value  $v = \sqrt{2} \langle \Phi \rangle \approx 246$  GeV. The Weinberg operator generates the Majorana neutrino masses as  $m_{\ell\ell'} = C_5^{\ell\ell'} v^2 / \Lambda$ . It leads to an experimental signature involving lepton number violation (LNV), as the Majorana neutrino is its own antiparticle. The LNV process has been tested extensively in searches for neutrinoless double beta decay in the decays of heavy nuclei. These experiments search for a signature consisting of two same-charge electrons and the absence of neutrinos in the final state. These searches set stringent upper limits, in the range 79–180 meV, on  $|m_{ee}|$  at the 90% confidence level (CL) [3]. Since muons and tau leptons are much heavier than electrons, same-charge lepton final states involving muons and tau leptons are kinematically forbidden in low-energy nuclear experiments and have to be studied at colliders [4].