

Stasis in an Expanding Universe: A Recipe for Stable Mixed-Component Cosmological Eras

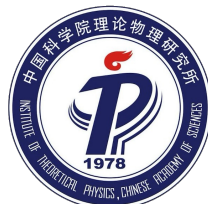
Fei Huang

arXiv: 2111.04753

PRD 105 (2022) 2, 023530

in collaboration with

Keith Dienes, Lucien Heurtier, Doojin Kim, Tim Tait, Brooks Thomas

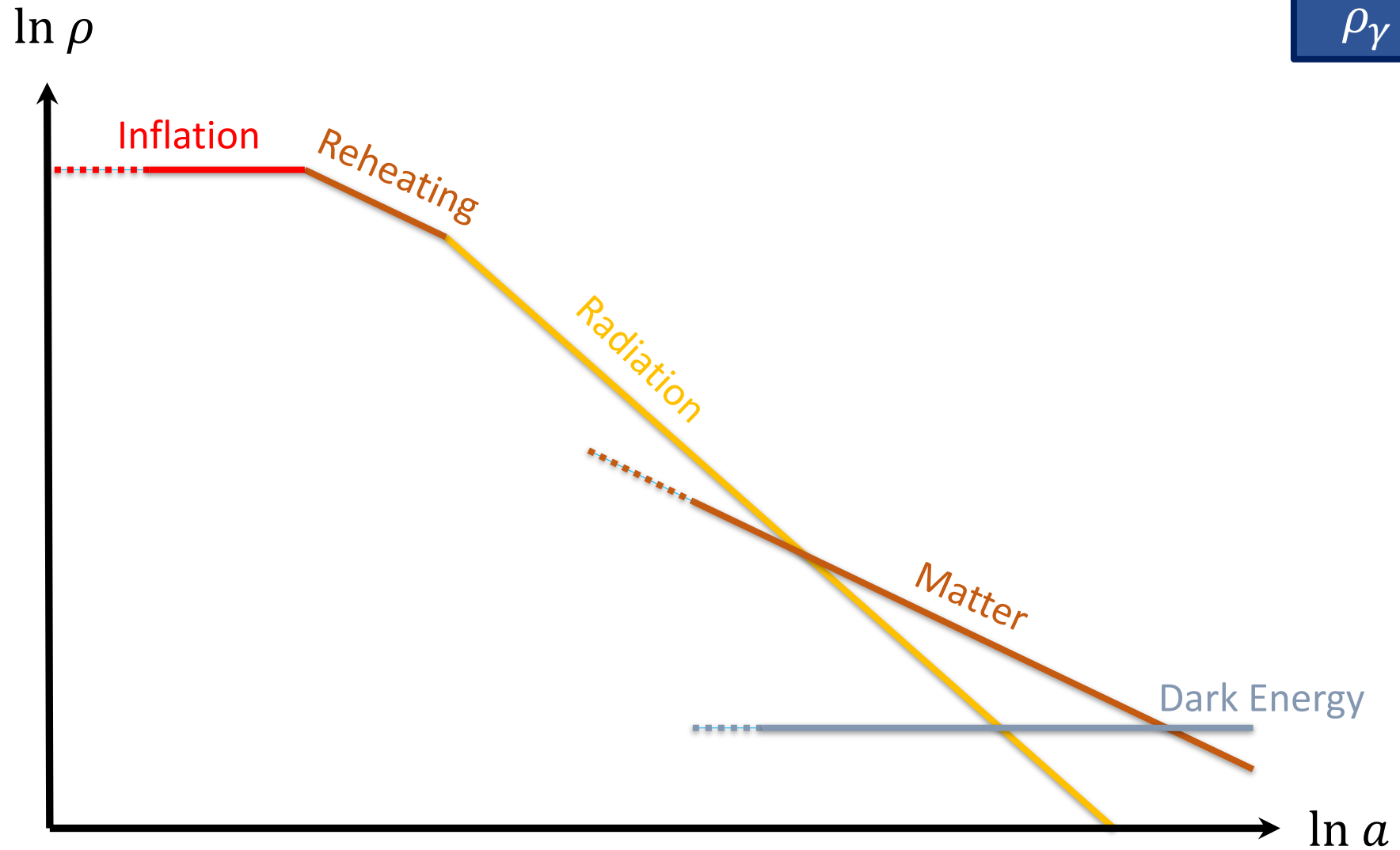


05/10/2022

PHENO
2022

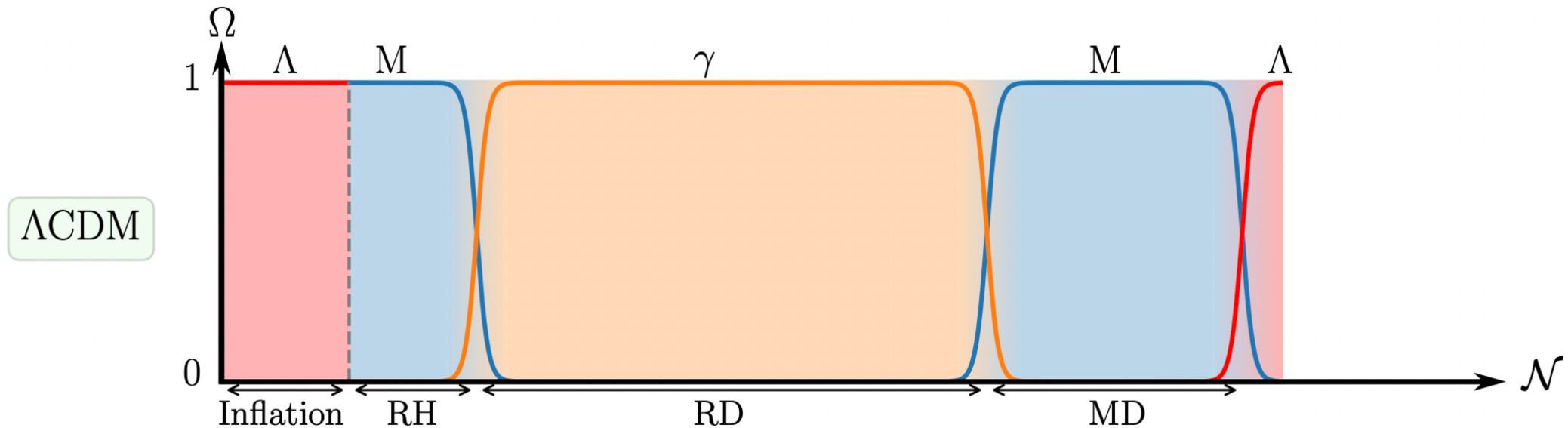
The Standard Lore: Λ CDM

$$\begin{aligned}\rho_\Lambda &\sim a^0 \\ \rho_M &\sim a^{-3} \\ \rho_\gamma &\sim a^{-4}\end{aligned}$$



The Standard Lore: Λ CDM

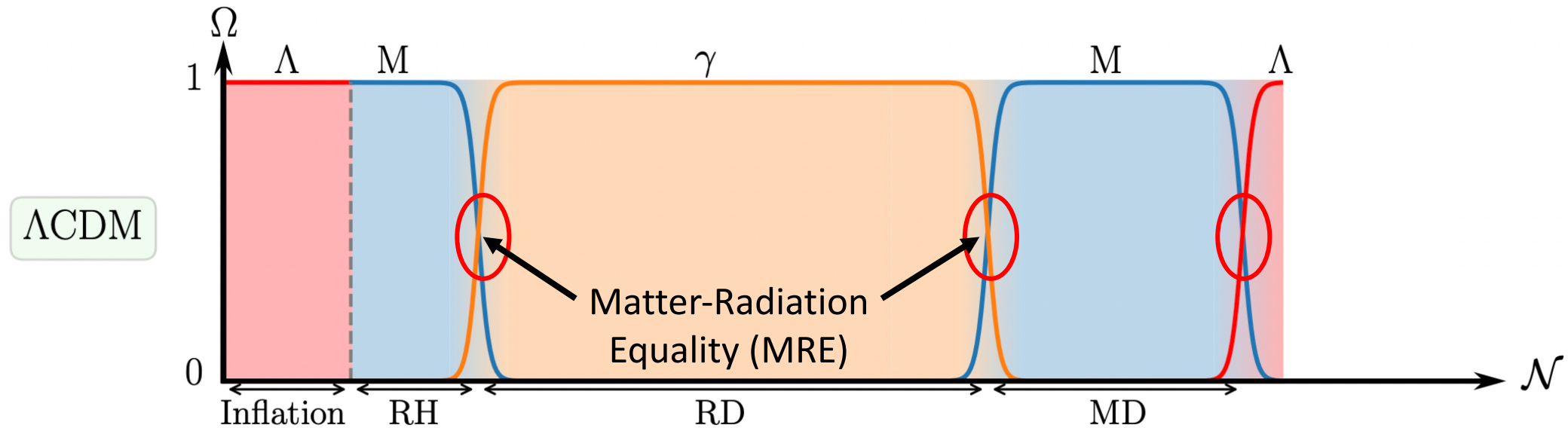
Most parts of the standard cosmological timeline are dominated by just one single energy component.



natural to divide the cosmological history into epochs characterized by the dominant energy component.

The Standard Lore: Λ CDM

Most parts of the standard cosmological timeline are dominated by just one single energy component.



natural to divide the cosmological history into epochs characterized by the dominant energy component.

Goals and Challenges

- In the standard Λ CDM cosmology, critical moments like MRE are fleeting.
- What if the MRE is not just a point, but a sustained era in the cosmological timeline?
- More generally, is it possible to maintain a constant Ω_M and Ω_γ over a sizable period of time?

Goals and Challenges

- In the standard Λ CDM cosmology, critical moments like MRE are fleeting.
- What if the MRE is not just a point, but a sustained era in the cosmological timeline?
- More generally, is it possible to maintain a constant Ω_M and Ω_γ over a sizable period of time?
- The universe is expanding
- $\rho_\gamma \sim a^{-4}$ is diluted faster than $\rho_M \sim a^{-3}$ in an expanding universe

Goals and Challenges

- In the standard Λ CDM cosmology, critical moments like MRE are fleeting.
- What if the MRE is not just a point, but a sustained era in the cosmological timeline?
- More generally, is it possible to maintain a constant Ω_M and Ω_γ over a sizable period of time?
- The universe is expanding
- $\rho_\gamma \sim a^{-4}$ is diluted faster than $\rho_M \sim a^{-3}$ in an expanding universe

Need energy transfer between the matter and radiation sectors

Goals and Challenges

- In the standard Λ CDM cosmology, critical moments like MRE are fleeting.
- What if the MRE is not just a point, but a sustained era in the cosmological timeline?
- More generally, is it possible to maintain a constant Ω_M and Ω_γ over a sizable period of time?

- The universe is expanding
- $\rho_\gamma \sim a^{-4}$ is diluted faster than $\rho_M \sim a^{-3}$ in an expanding universe

Need energy transfer between the matter and radiation sectors



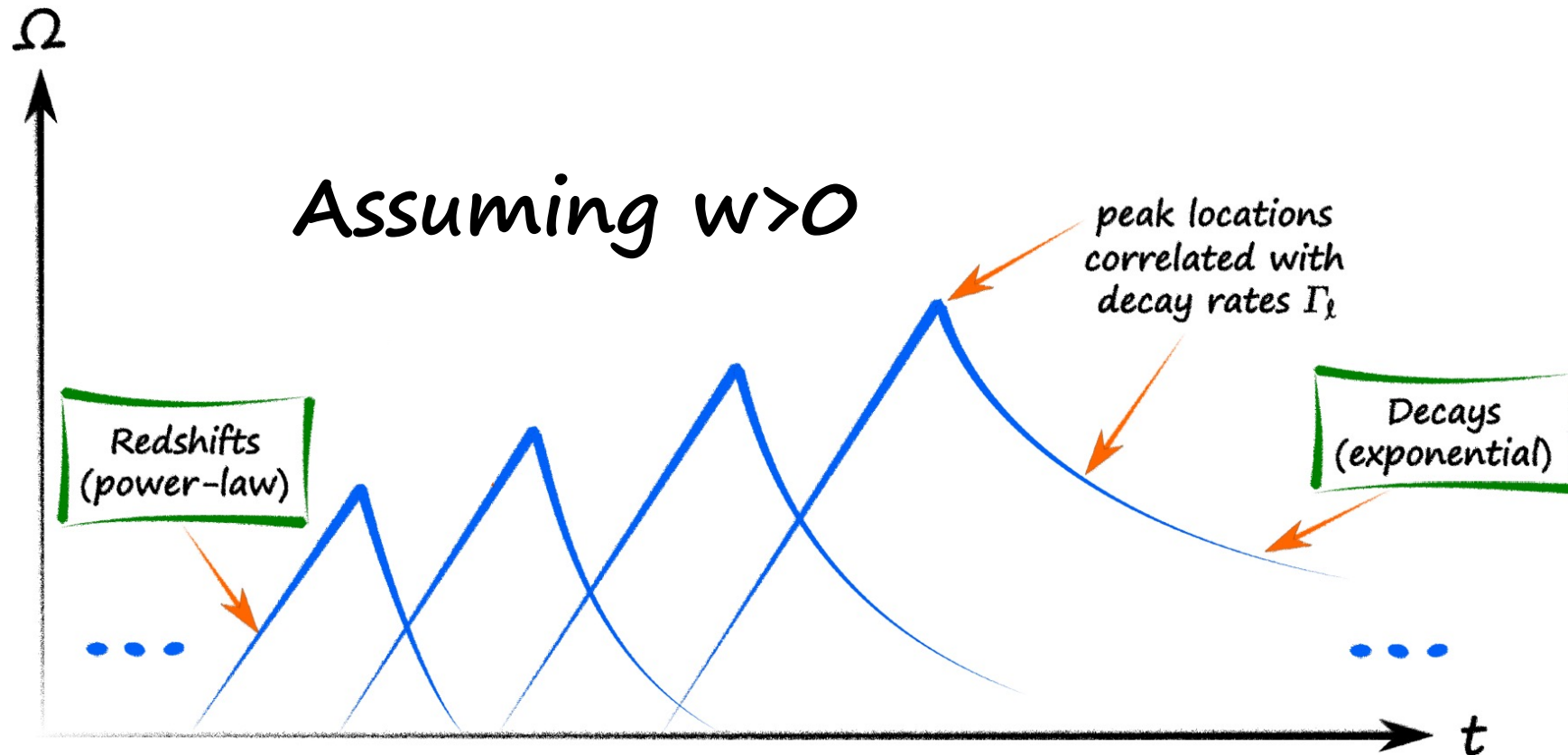
Particle decays!

Goals and Challenges

- However, the decay of a single matter state is not enough since particle decay has an exponential behavior and only occurs during a short window of time

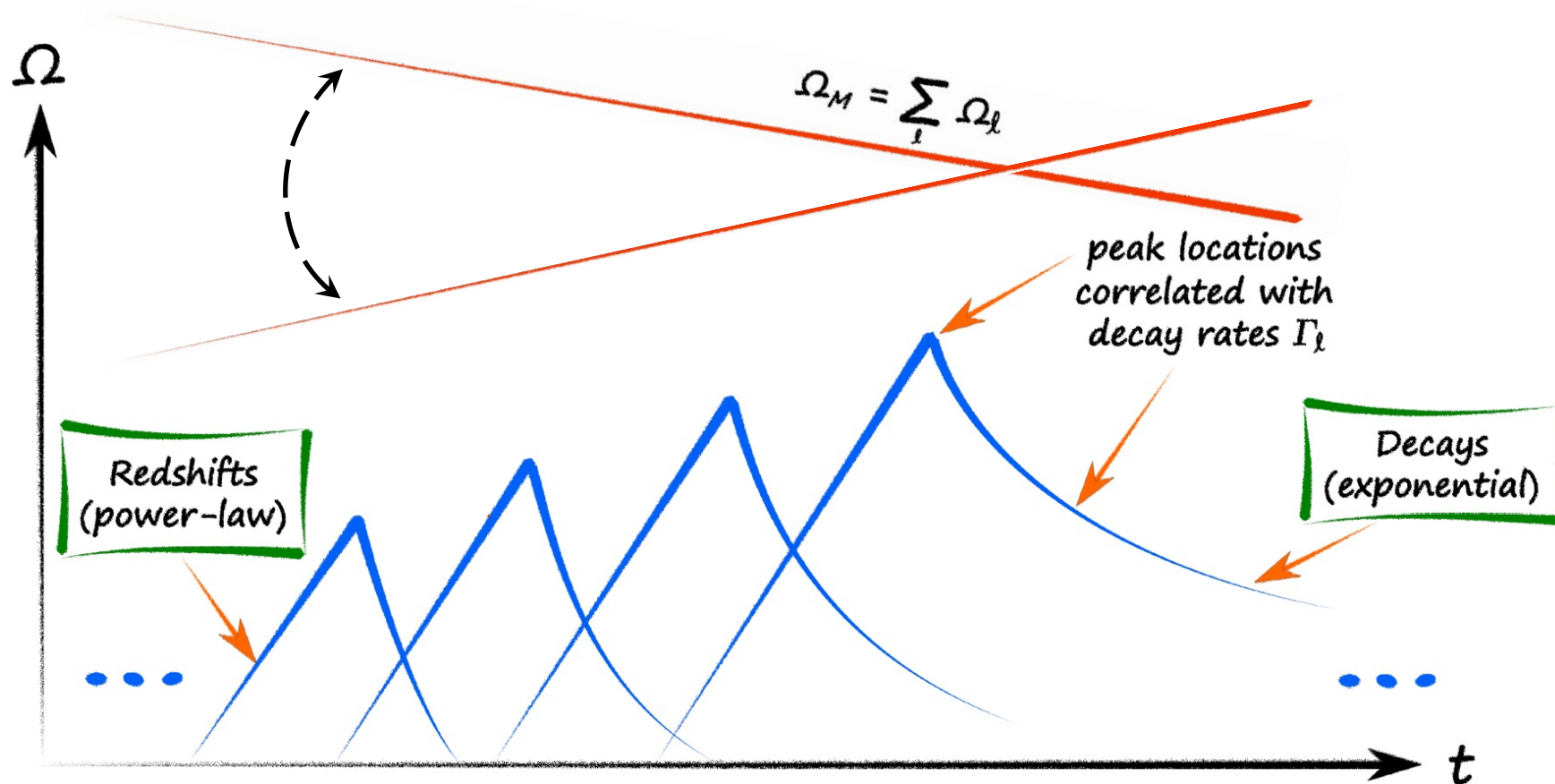
Goals and Challenges

- However, the decay of a single matter state is not enough since particle decay has an exponential behavior and only occurs during a short window of time
- How about a series of decays from many matter states?



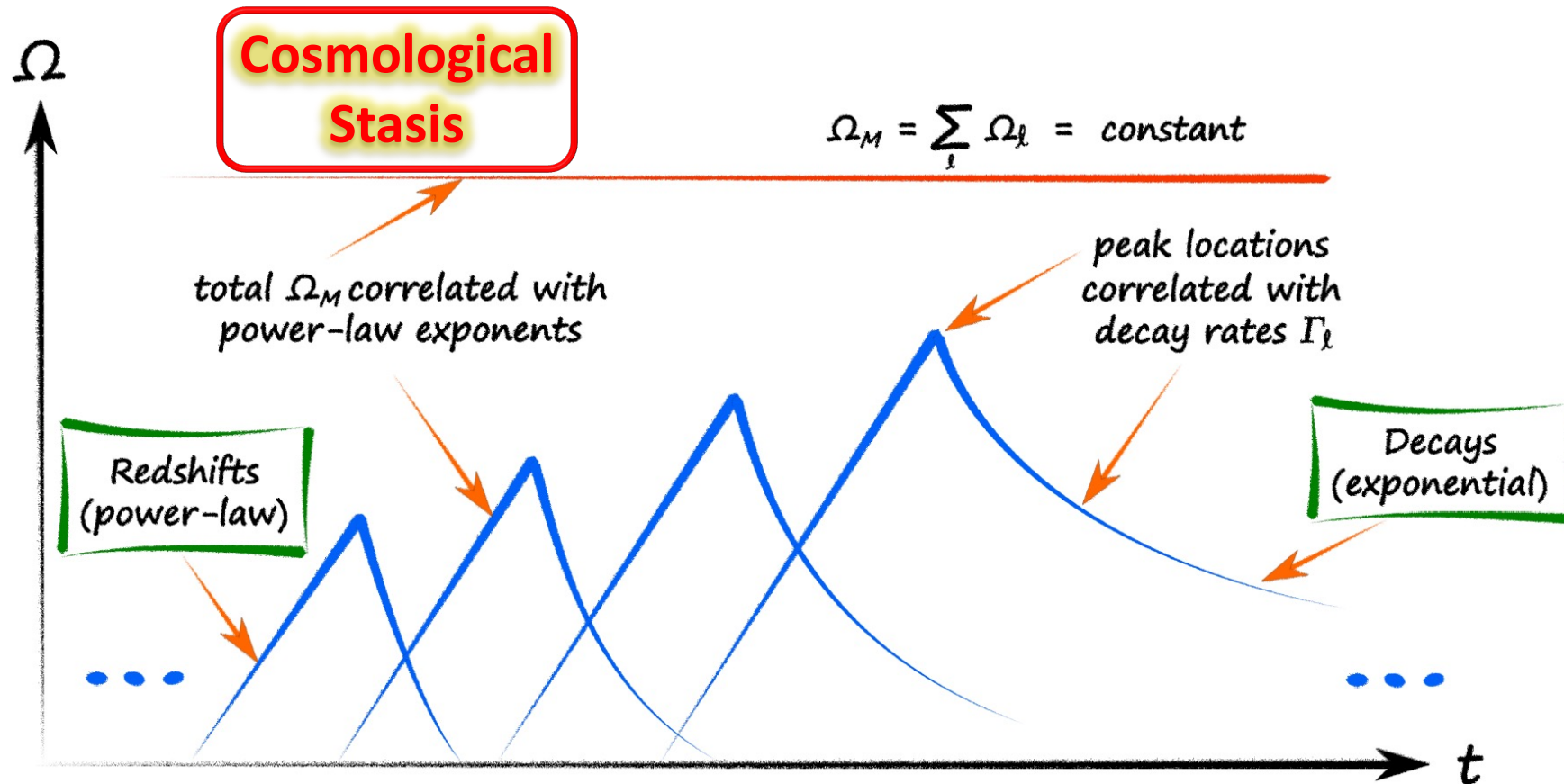
Goals and Challenges

- However, the decay of a single matter state is not enough since particle decay has an exponential behavior and only occurs during a short window of time
- How about a series of decays from many matter states?

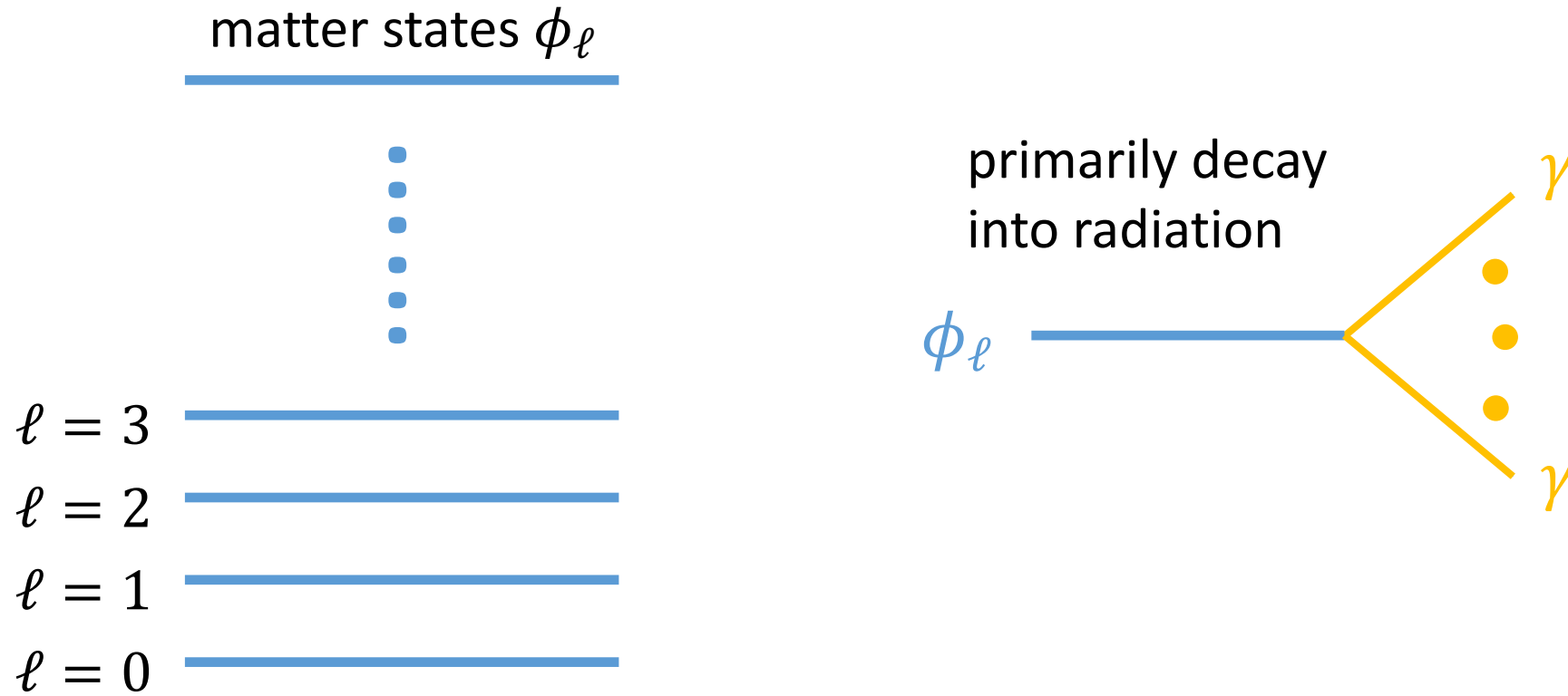


Goals and Challenges

- However, the decay of a single matter state is not enough since particle decay has an exponential behavior and only occurs during a short window of time
- How about a series of decays from many matter states?



Conditions for Stasis



$$\Omega_\ell = \frac{8\pi G}{3H^2} \rho_\ell \quad \Omega_M = \sum_\ell \Omega_\ell \quad \Omega_M + \Omega_\gamma = 1$$

Conditions for Stasis

Boltzmann equations

$$\frac{d\rho_\ell}{dt} = -3H\rho_\ell - \Gamma_\ell\rho_\ell \qquad \frac{d\rho_\gamma}{dt} = -4H\rho_\gamma + \sum_\ell \Gamma_\ell\rho_\ell$$

+

Friedmann equation

$$H^2 = \frac{8\pi G}{3} (\rho_M + \rho_\gamma)$$

Conditions for Stasis

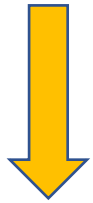
Boltzmann equations

$$\frac{d\rho_\ell}{dt} = -3H\rho_\ell - \Gamma_\ell\rho_\ell \qquad \frac{d\rho_\gamma}{dt} = -4H\rho_\gamma + \sum_\ell \Gamma_\ell\rho_\ell$$

+

Friedmann equation

$$H^2 = \frac{8\pi G}{3} (\rho_M + \rho_\gamma)$$




$$\frac{d\Omega_M}{dt} = - \sum_\ell \Gamma_\ell \Omega_\ell + H(\Omega_M - \Omega_M^2)$$

$$\frac{d\Omega_M}{dt} = 0$$

$$\sum_\ell \Gamma_\ell \Omega_\ell = H(\Omega_M - \Omega_M^2)$$

Conditions for Stasis

$$\frac{d\Omega_M}{dt} = - \sum_{\ell} \Gamma_{\ell} \Omega_{\ell} + H(\Omega_M - \Omega_M^2)$$

$\frac{d\Omega_M}{dt} = 0$


Necessary condition for Stasis,
but **NOT sufficient**

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_M - \Omega_M^2)$$

To have Ω_M remain constant in time, one would require

$$\frac{d^n \Omega_M}{dt^n} = 0$$



an infinite set of constraints

Conditions for Stasis

$$\frac{d\Omega_M}{dt} = - \sum_{\ell} \Gamma_{\ell} \Omega_{\ell} + H(\Omega_M - \Omega_M^2)$$

$\xrightarrow{\frac{d\Omega_M}{dt} = 0}$

Necessary condition for Stasis,
but not sufficient

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_M - \Omega_M^2)$$

To have Ω
constant
one would

Instead, we shall do this
in a quicker way 😎

to set
constraints

Assuming $\frac{d\Omega_M}{dt} = 0$ is satisfied at some **fiducial time** t_* , demand that both sides evolve with time in the same manner.

Conditions for Stasis

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_M - \Omega_M^2)$$

For a universe with only matter and radiation

$$\frac{dH}{dt} = -\frac{1}{2}H^2(4 - \Omega_M)$$

During Stasis $\Omega_M = \bar{\Omega}_M$

$\sim \frac{1}{t}$

$H(t) = \left(\frac{2}{4 - \bar{\Omega}_M} \right) \frac{1}{t}$

Conditions for Stasis

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_M - \Omega_M^2)$$

$$\Omega_{\ell} = \rho_{\ell} / \rho_{tot}$$

$$\rho_{\ell} \sim a^{-3} e^{-\Gamma_{\ell}(t-t_*)}$$

$$\rho_{tot} \sim H^2 \sim t^{-2}$$

$$a \sim t^{2/(4-\bar{\Omega}_M)}$$

$$\sim \frac{1}{t}$$

$$\Omega_{\ell}(t) = \Omega_{\ell}^* \left(\frac{t}{t_*} \right)^{2-6/(4-\bar{\Omega}_M)} e^{-\Gamma_{\ell}(t-t_*)}$$

For a universe with only matter and radiation

$$\frac{dH}{dt} = -\frac{1}{2} H^2 (4 - \Omega_M)$$

During Stasis $\Omega_M = \bar{\Omega}_M$



$$H(t) = \left(\frac{2}{4 - \bar{\Omega}_M} \right) \frac{1}{t}$$

Conditions for Stasis

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell} = H(\Omega_M - \Omega_M^2)$$

For a universe with only matter and radiation

$$\frac{dH}{dt} = -\frac{1}{2} H^2 (4 - \Omega_M)$$

During Stasis $\Omega_M = \bar{\Omega}_M$



$$H(t) = \left(\frac{2}{4 - \bar{\Omega}_M} \right) \frac{1}{t}$$

$$\Omega_{\ell} = \rho_{\ell} / \rho_{tot}$$

$$\rho_{\ell} \sim a^{-3} e^{-\Gamma_{\ell}(t-t_*)}$$

$$\rho_{tot} \sim H^2 \sim t^{-2}$$

$$a \sim t^{2/(4 - \bar{\Omega}_M)}$$

$$\sim \frac{1}{t}$$

$$\Omega_{\ell}(t) = \Omega_{\ell}^* \left(\frac{t}{t_*} \right)^{2 - 6/(4 - \bar{\Omega}_M)} e^{-\Gamma_{\ell}(t-t_*)}$$

**Stasis
(Eternal)**

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t) = \frac{2\bar{\Omega}_M(1 - \bar{\Omega}_M)}{4 - \bar{\Omega}_M} \frac{1}{t}$$

$$\sum_{\ell} \Omega_{\ell}(t) = \bar{\Omega}_M$$

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \bar{\Omega}_M)}{4 - \bar{\Omega}_M} \frac{1}{t}$$

A Model of Stasis

Mass Spectrum

$$m_\ell = m_0 + (\Delta m)\ell^\delta$$

Decay Widths

$$\Gamma_\ell = \Gamma_0 \left(\frac{m_\ell}{m}\right)^\gamma$$

Initial Abundances

$$\Omega_\ell^{(0)} = \Omega_0^{(0)} \left(\frac{m_\ell}{m}\right)^\alpha$$

A Model of Stasis

Mass Spectrum

$$m_\ell = m_0 + (\Delta m)\ell^\delta$$

Decay Widths

$$\Gamma_\ell = \Gamma_0 \left(\frac{m_\ell}{m}\right)^\gamma$$

Initial Abundances

$$\Omega_\ell^{(0)} = \Omega_0^{(0)} \left(\frac{m_\ell}{m}\right)^\alpha$$

Depends on particle physics model

- KK excitations of a 5-d scalar field compactified on a circle of radius R
 - $\delta \sim 1$ for $mR \ll 1$
 - $\delta \sim 2$ for $mR \gg 1$
- Bound states of strongly-coupled gauge theory
 - $\delta \sim 1/2$

Depends on decay mode

- if ϕ_ℓ decays to photons through contact operator $\mathcal{O}_\ell \sim c_\ell \phi_\ell \mathcal{F} / \Lambda^{d-4}$, $\gamma = 2d - 7$, e.g., $\gamma \sim \{3, 5, 7\}$

Depends on the production mechanism

- $\alpha < 0$ for misalignment production
- both $\alpha > 0$ or $\alpha < 0$ for thermal freeze-out
- $\alpha = 1$ for universal inflaton decay

A Model of Stasis

Mass Spectrum

$$m_\ell = m_0 + (\Delta m)\ell^\delta$$

Decay Widths

$$\Gamma_\ell = \Gamma_0 \left(\frac{m_\ell}{m}\right)^\gamma$$

Initial Abundances

$$\Omega_\ell^{(0)} = \Omega_0^{(0)} \left(\frac{m_\ell}{m}\right)^\alpha$$

Free parameters

$$\left\{ \alpha, \gamma, \delta, m_0, \Delta m, \Gamma_0, \Omega_0^{(0)}, t^{(0)} \right\}$$

initial
conditions

Depends on particle physics model

- KK excitations of a 5-d scalar field compactified on a circle of radius R
 - $\delta \sim 1$ for $mR \ll 1$
 - $\delta \sim 2$ for $mR \gg 1$
- Bound states of strongly-coupled gauge theory
 - $\delta \sim 1/2$

Depends on decay mode

- if ϕ_ℓ decays to photons through contact operator $\mathcal{O}_\ell \sim c_\ell \phi_\ell \mathcal{F} / \Lambda^{d-4}$, $\gamma = 2d - 7$, e.g., $\gamma \sim \{3, 5, 7\}$

Depends on the production mechanism

- $\alpha < 0$ for misalignment production
- both $\alpha > 0$ or $\alpha < 0$ for thermal freeze-out
- $\alpha = 1$ for universal inflaton decay

A Model of Stasis

Recall the condition for stasis

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \bar{\Omega}_M) 1}{4 - \bar{\Omega}_M} \frac{1}{t}$$

With our particular model,
we can work out the LHS

If we consider the production
of states ϕ_{ℓ} before stasis

$$\Omega_{\ell}(t) = \Omega_{\ell}^{(0)} h(t^{(0)}, t) e^{-\Gamma_{\ell}(t-t^{(0)})}$$

with non-trivial **redshift**
effect encoded in $h(t^{(0)}, t)$

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t) = \Gamma_0 \Omega_0^{(0)} h(t^{(0)}, t) \sum_{\ell} \left(\frac{m_{\ell}}{m_0} \right)^{\alpha+\gamma} e^{-\Gamma_0 \left(\frac{m_{\ell}}{m_0} \right)^{\gamma} (t-t^{(0)})}$$

A Model of Stasis

Recall the condition for stasis

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \bar{\Omega}_M)}{4 - \bar{\Omega}_M} \frac{1}{t}$$

With our particular model,
we can work out the LHS

If we consider the production
of states ϕ_{ℓ} before stasis

$$\Omega_{\ell}(t) = \Omega_{\ell}^{(0)} h(t^{(0)}, t) e^{-\Gamma_{\ell}(t-t^{(0)})}$$

with non-trivial **redshift**
effect encoded in $h(t^{(0)}, t)$

$$\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t) = \Gamma_0 \Omega_0^{(0)} h(t^{(0)}, t) \sum_{\ell} \left(\frac{m_{\ell}}{m_0}\right)^{\alpha+\gamma} e^{-\Gamma_0 \left(\frac{m_{\ell}}{m_0}\right)^{\gamma} (t-t^{(0)})}$$

Continuous limit

$$\begin{aligned} &= \frac{\Gamma_0 \Omega_0^{(0)}}{\delta (\Delta m)^{1/\delta}} h(t^{(0)}, t) \int_0^{\infty} dm m^{1/\delta-1} \left(\frac{m}{m_0}\right)^{\alpha+\gamma} e^{-\Gamma_0 \left(\frac{m}{m_0}\right)^{\gamma} (t-t^{(0)})} \\ &= \frac{\Gamma_0 \Omega_0^{(0)}}{\gamma \delta} \left(\frac{m_0}{\Delta m}\right)^{1/\delta} \Gamma\left(\frac{\alpha + 1/\delta + \gamma}{\gamma}\right) h(t^{(0)}, t) [\Gamma_0 (t - t^{(0)})]^{-(\alpha+1/\delta+\gamma)/\gamma} \end{aligned}$$

A Model of Stasis

Recall the condition for stasis

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \bar{\Omega}_M)}{4 - \bar{\Omega}_M} \frac{1}{t}$$

With our particular model, we can work out the LHS

If we consider the production of states ϕ_{ℓ} before stasis

$$\Omega_{\ell}(t) = \Omega_{\ell}^{(0)} h(t^{(0)}, t) e^{-\Gamma_{\ell}(t-t^{(0)})}$$

with non-trivial **redshift effect** encoded in $h(t^{(0)}, t)$

Therefore

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{\alpha + 1/\delta}{\gamma} \frac{1}{t - t^{(0)}}$$

Stasis condition can be satisfied as long as $t^{(0)} \ll t$

A Model of Stasis

Recall the condition for stasis

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{2(1 - \bar{\Omega}_M)}{4 - \bar{\Omega}_M} \frac{1}{t}$$

With our particular model, we can work out the LHS

If we consider the production of states ϕ_{ℓ} before stasis

$$\Omega_{\ell}(t) = \Omega_{\ell}^{(0)} h(t^{(0)}, t) e^{-\Gamma_{\ell}(t-t^{(0)})}$$

with non-trivial **redshift effect** encoded in $h(t^{(0)}, t)$

Therefore

$$\frac{\sum_{\ell} \Gamma_{\ell} \Omega_{\ell}(t)}{\sum_{\ell} \Omega_{\ell}(t)} = \frac{\alpha + 1/\delta}{\gamma} \frac{1}{t - t^{(0)}}$$

Stasis condition can be satisfied as long as $t^{(0)} \ll t$

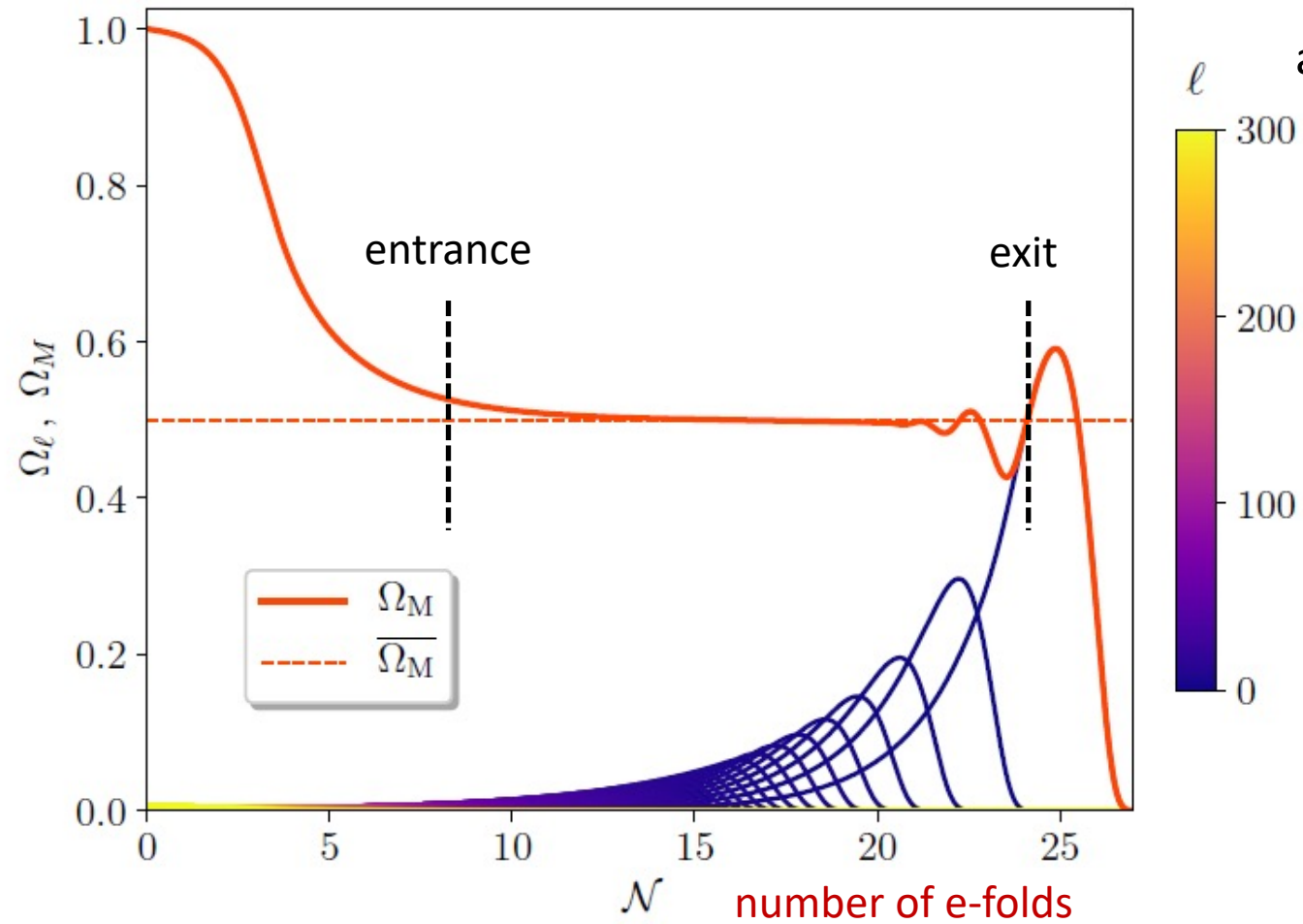
Matter abundance during stasis is determined by model parameters

$$\bar{\Omega}_M = \frac{2\gamma\delta - 4(1 + \alpha\delta)}{2\gamma\delta - (1 + \alpha\delta)}$$

It's time to test it numerically...

Examples

An "Extended MRE"

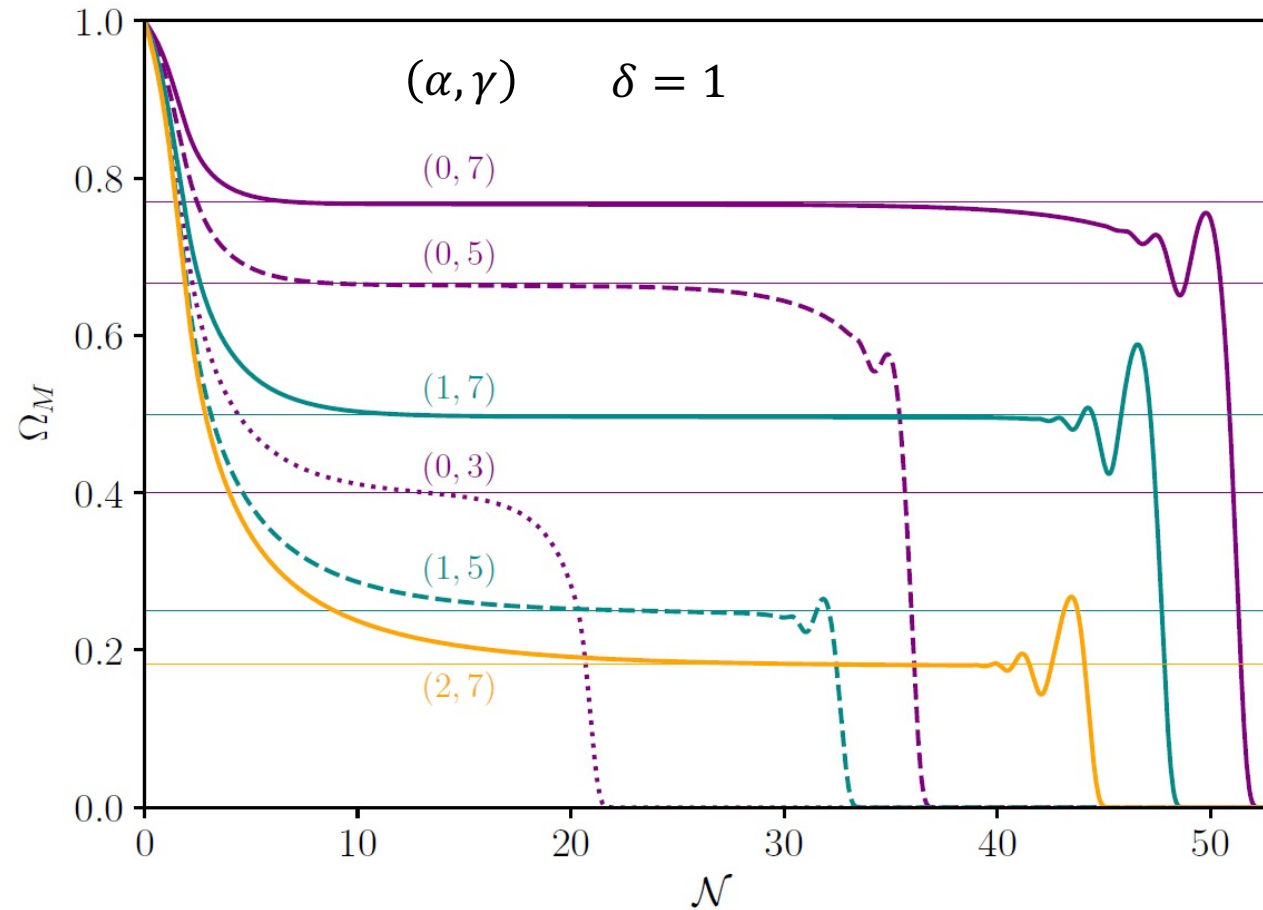


The actual model has an entrance and an exit

$$(\alpha, \gamma, \delta) = (1, 7, 1) \rightarrow \bar{\Omega}_M = 1/2, \Delta m = m_0, \Gamma_{N-1}/H^{(0)} = 0.01, N = 300$$

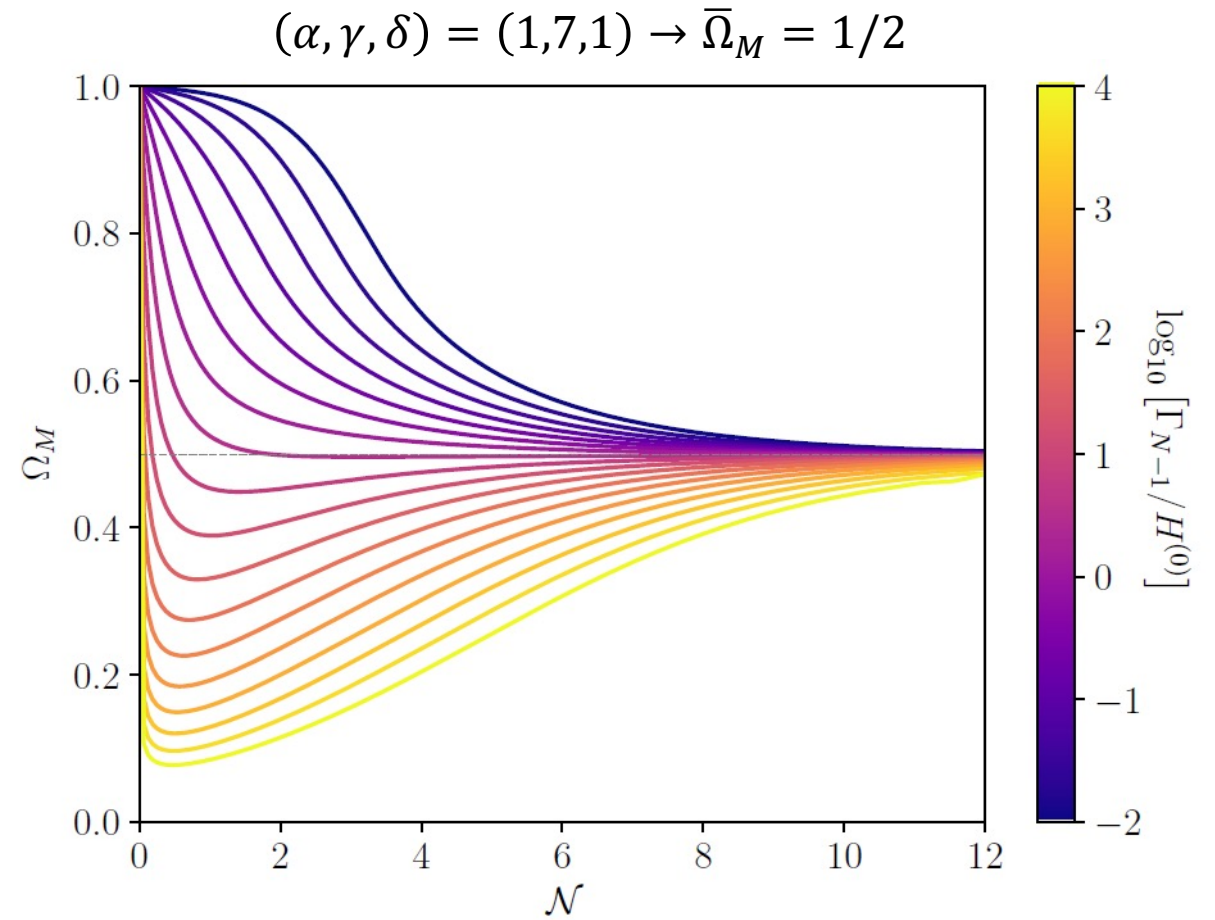
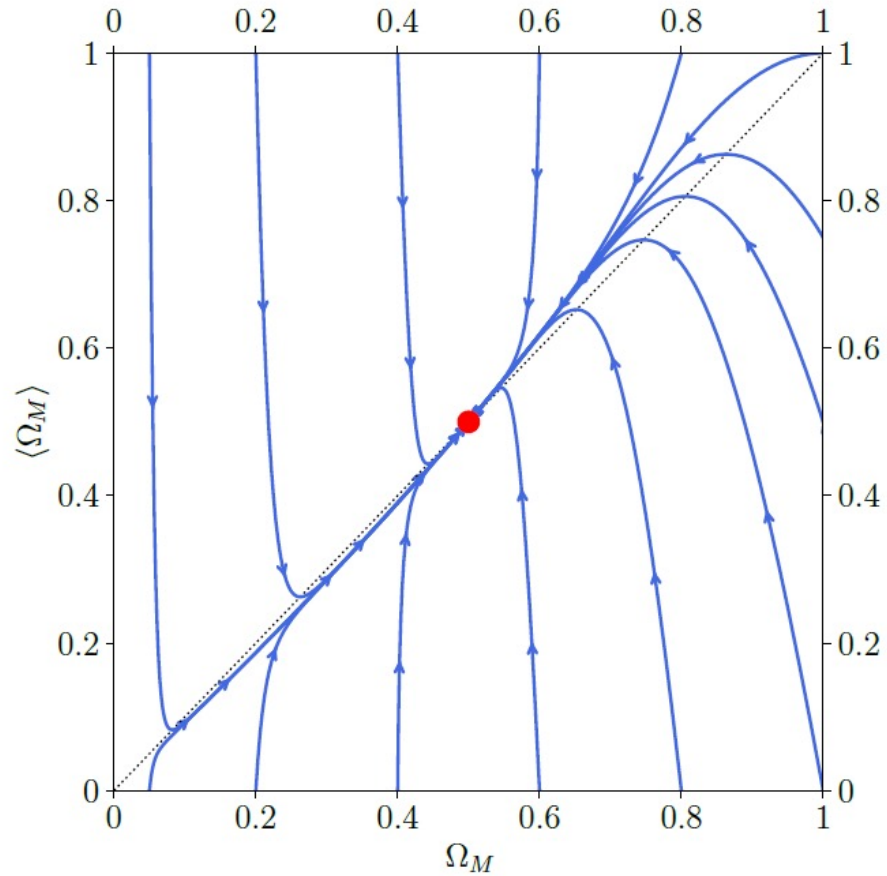
Examples

Fine tuned?



Prolonged epochs in which $\Omega_M \sim \bar{\Omega}_M$ can be achieved in multiple cases!

Stasis as a Global Attractor



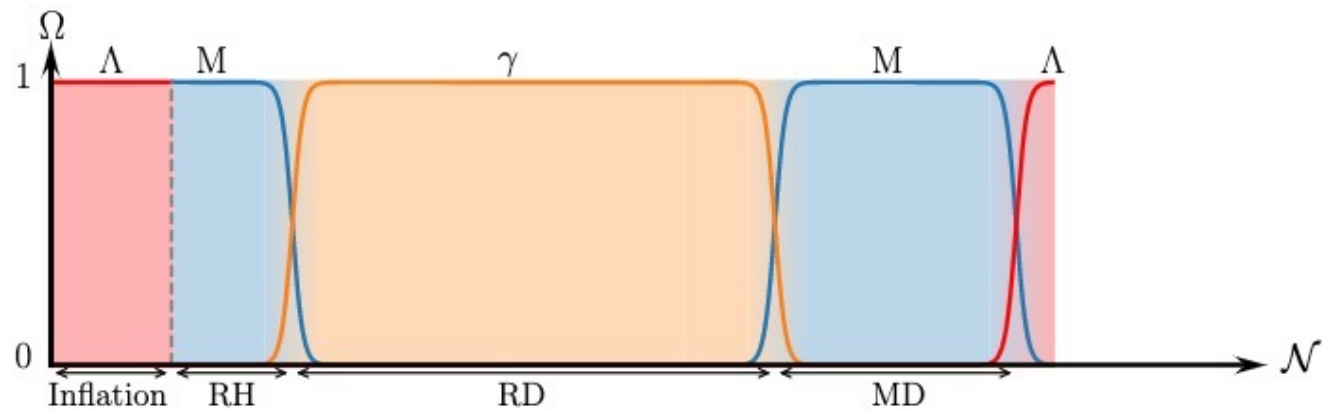
The attractor is GLOBAL!

Splicing Stasis into the Cosmological Timeline

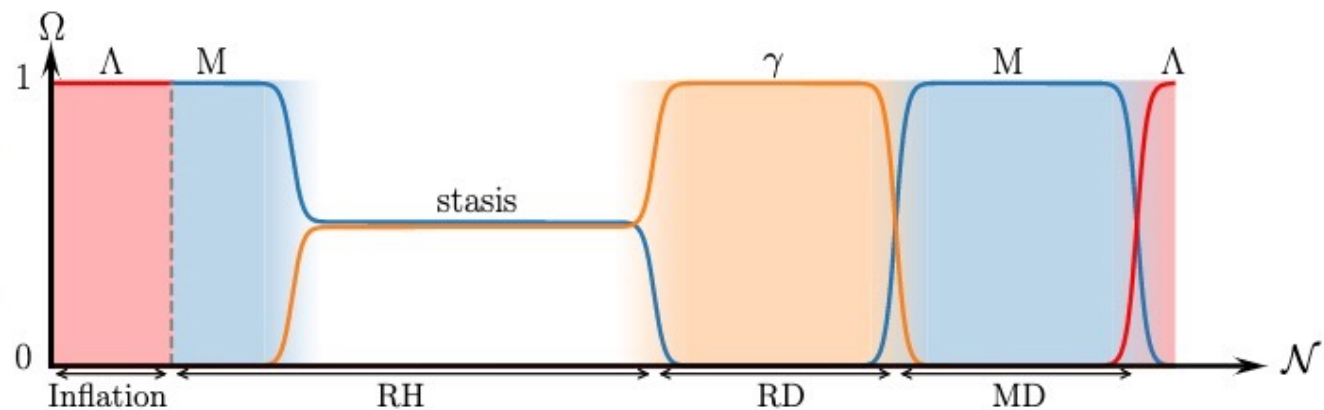
Reheating occurs during the stasis epoch and results from the decays of ϕ_ℓ

The presence of multiple matter fields first leads to an early matter-dominated era (EMDE), then stasis occurs when decays start

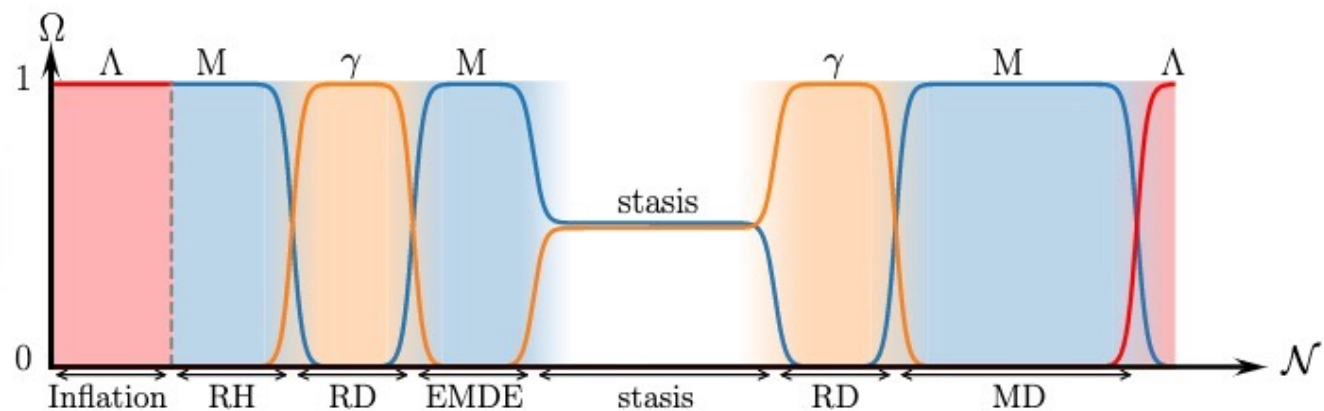
Λ CDM



Stasis spliced into RH



Stasis spliced into RD

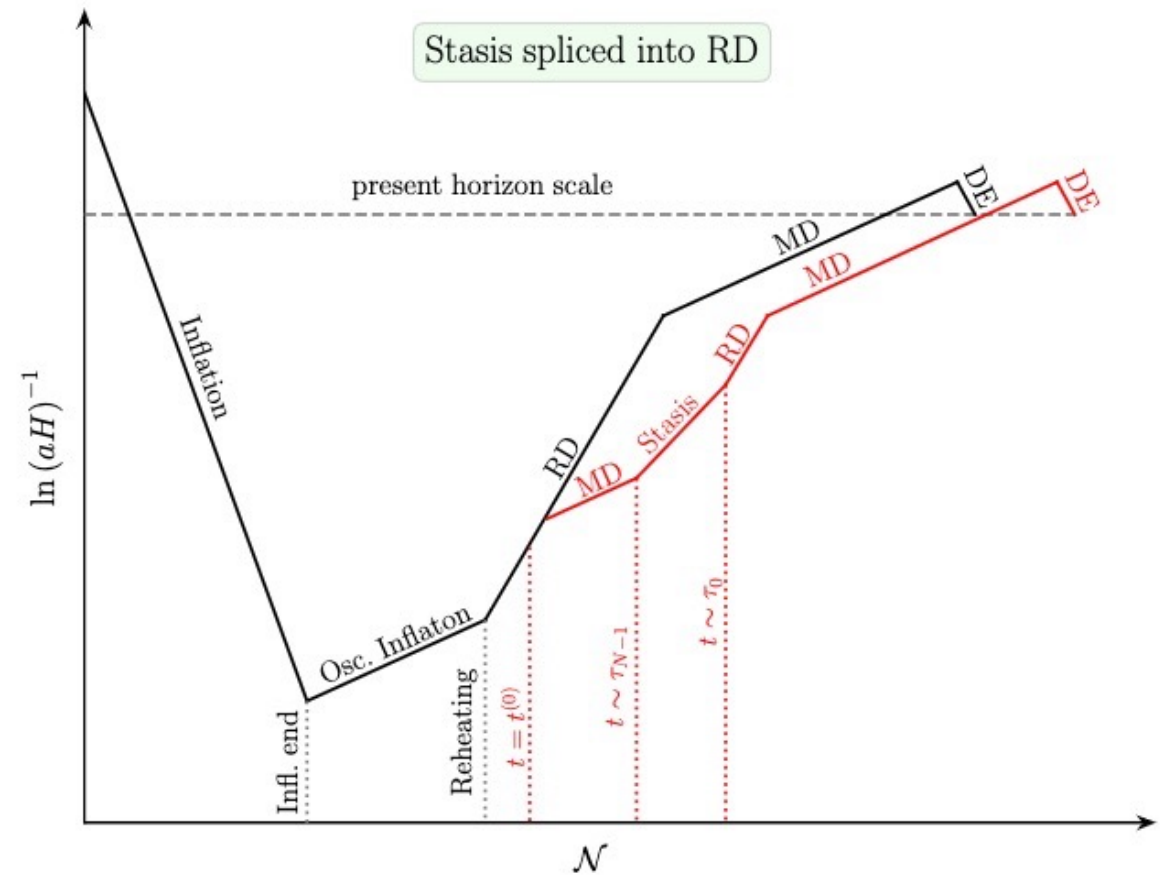
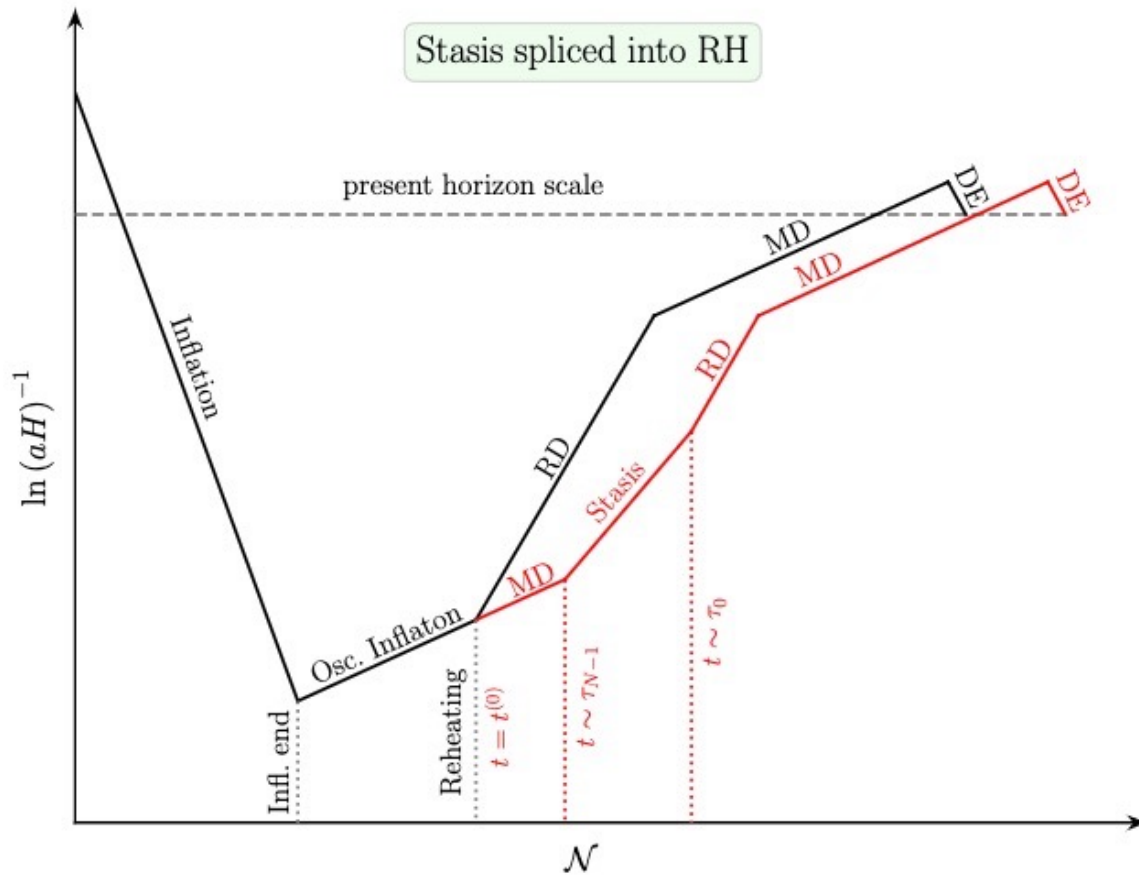


Summary and Future Directions

- Decays of a tower of states could result in an epoch in which the matter and radiation abundances are constant and can take any value.
- Such a scenario could occur naturally in models motivated by particle physics, like KK theory and strongly-coupled gauge theories.
- Stability analysis shows the stasis solution is a global attractor, No fine-tuning needed.
- Since the comoving Hubble radius grows more slowly, perturbation modes re-enter the horizon at a later time, can affect predictions for ***inflationary observables***
- ***Density perturbations*** grow faster during stasis than in Λ CDM, can result in formation of compact objects such as PBH or compact minihalos
- ***DM relic abundance*** would be affected if produced prior to or during Stasis due to different expansion history and injection of entropy. DM can also be the decay products of ϕ_ℓ or the lightest ϕ_ℓ with smaller decay width and parametrically smaller initial abundance.
- ***Generation of lepton or baryon asymmetry*** would also be affected if it occurs prior to or during Stasis, or if the asymmetry is produced through ϕ_ℓ decays
- Numerous directions to be explored...

Cosmological Implications

The insertion of a stasis period delays the subsequent timeline relative to Λ CDM expectations



Cosmological Implications

Stasis and EMDE can place the universe on the same cosmological timeline

