

# Asymmetric Reheating by Primordial Black Holes

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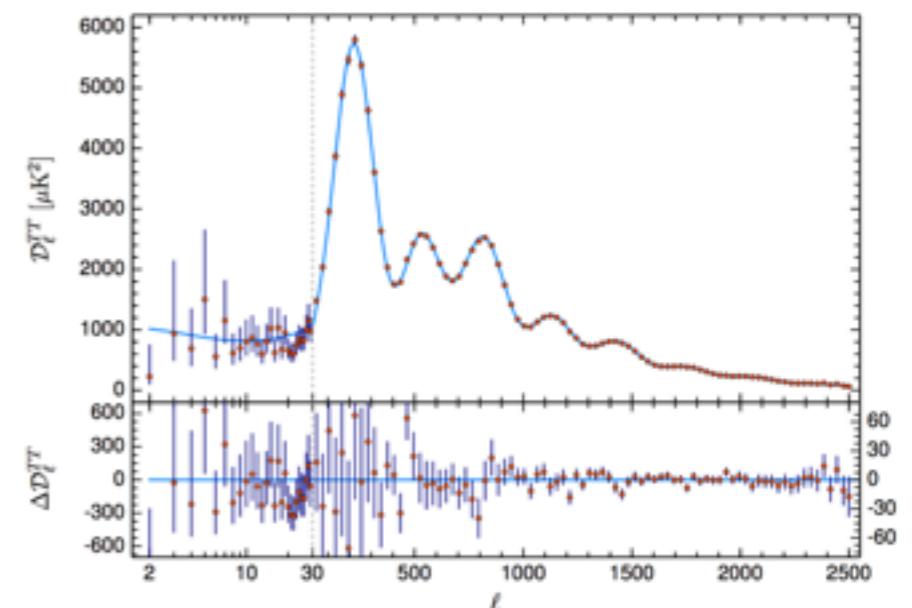
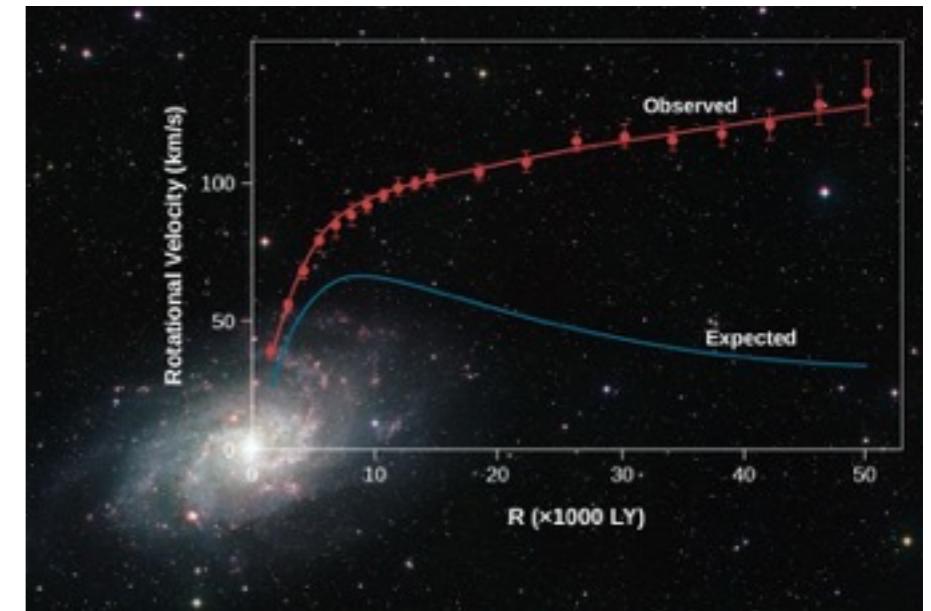
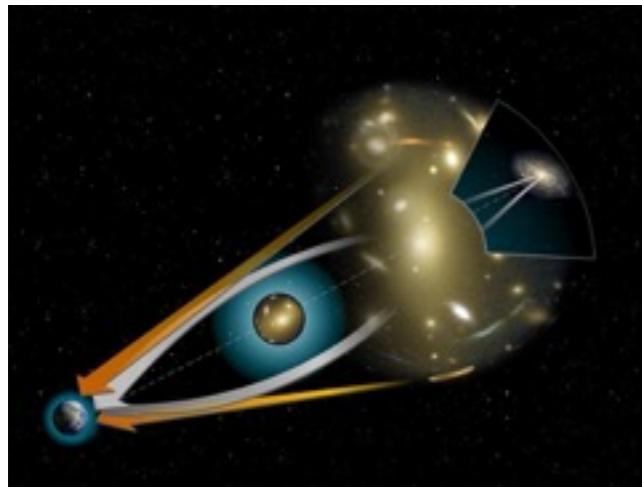
# Evidence for DM: (All from Gravitational Interaction)

Rotation Curves of Galaxies

Cosmic Microwave Background (CMB)

Structure Formation

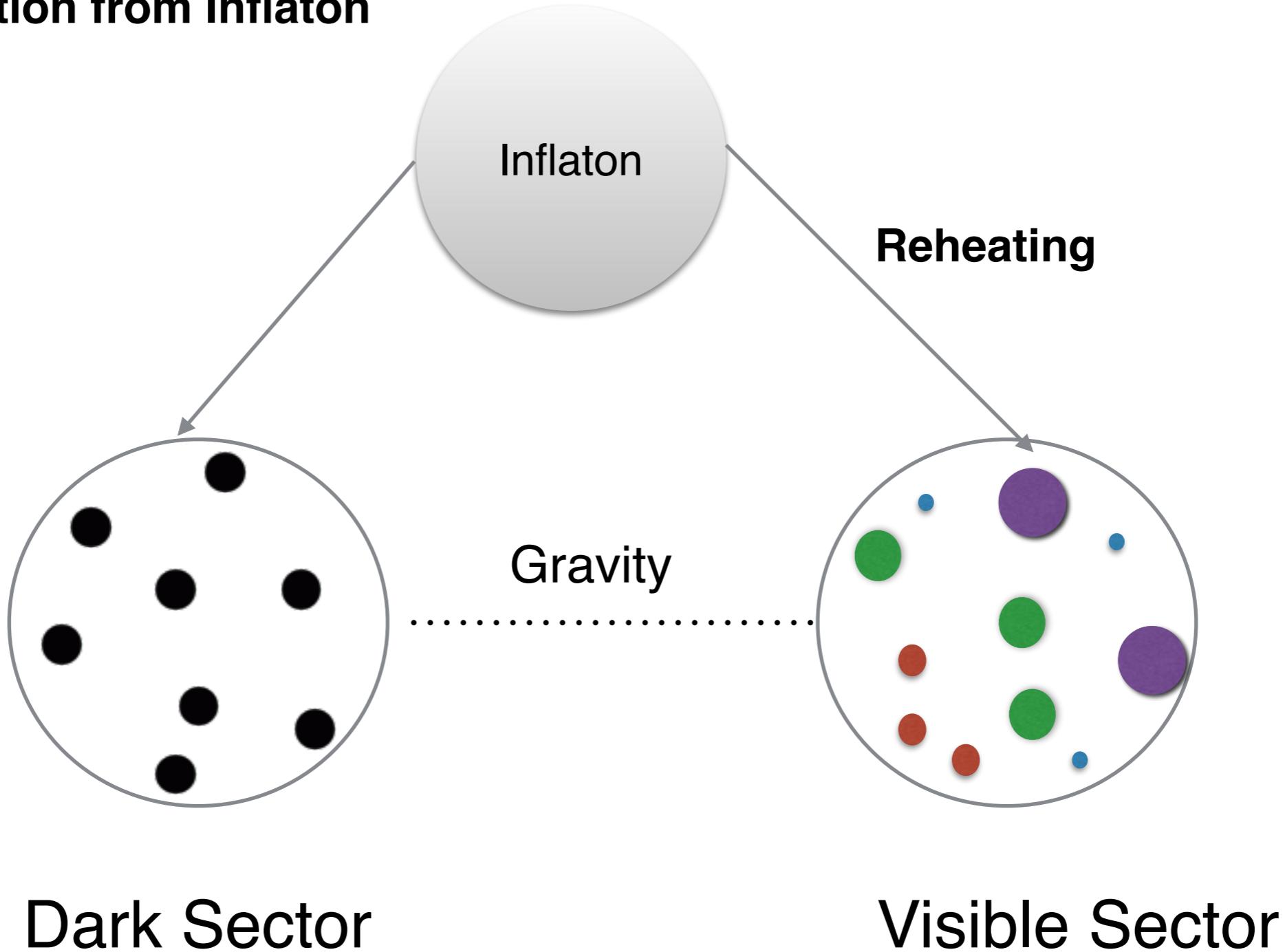
Gravitational Lensing



Null result of search for possible **non-gravitational interactions of DM**: motivates idea of dark sectors coupled only **gravitationally** to visible sector.

**challenge: populating the dark sector**

## Direct Production from Inflaton



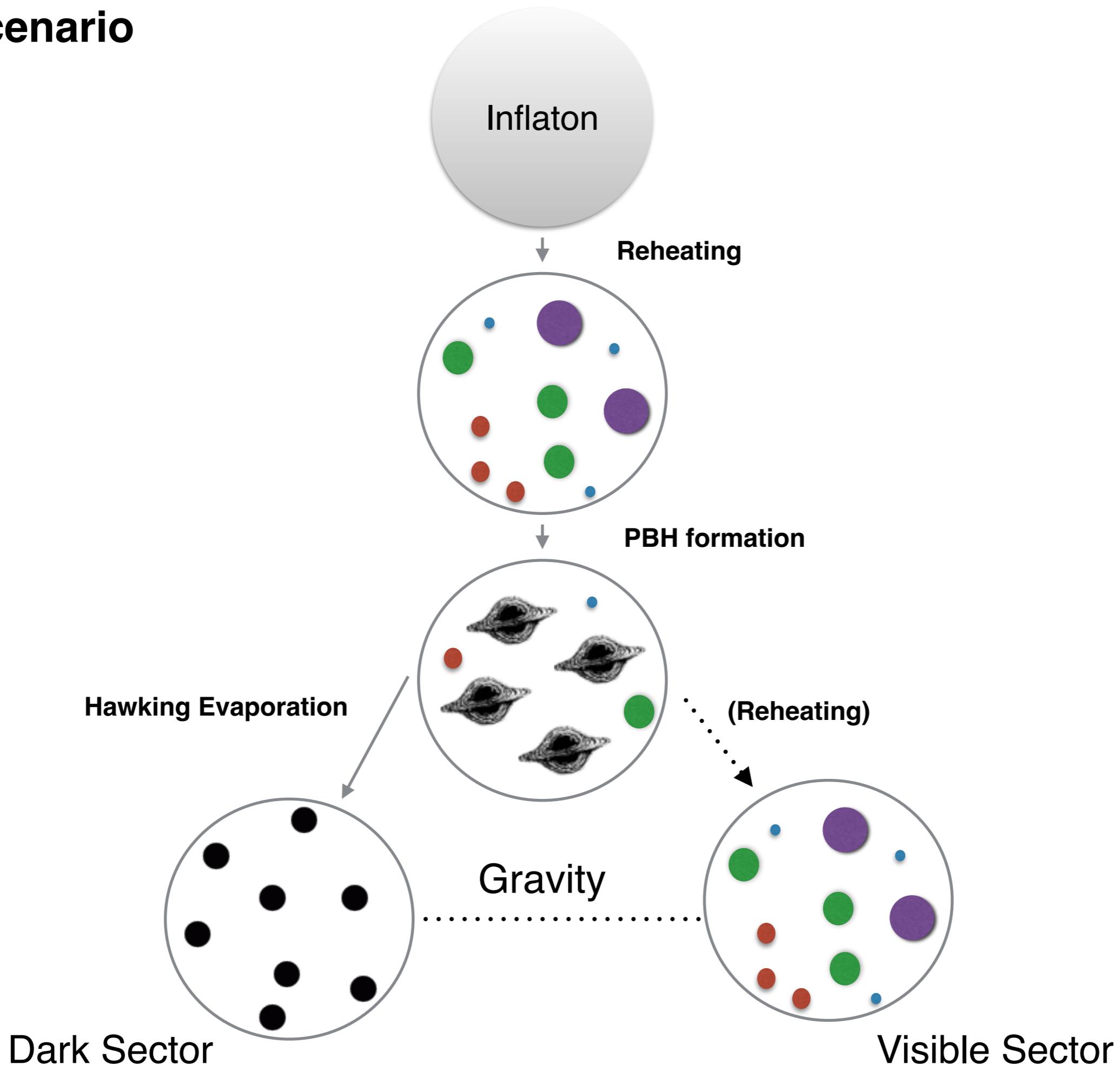
issue: two sectors can exchange inflaton and thermalize

P. Adshead, Y. Cui and J. Shelton, 2016

E. Hardy and J. Unwin, 2017

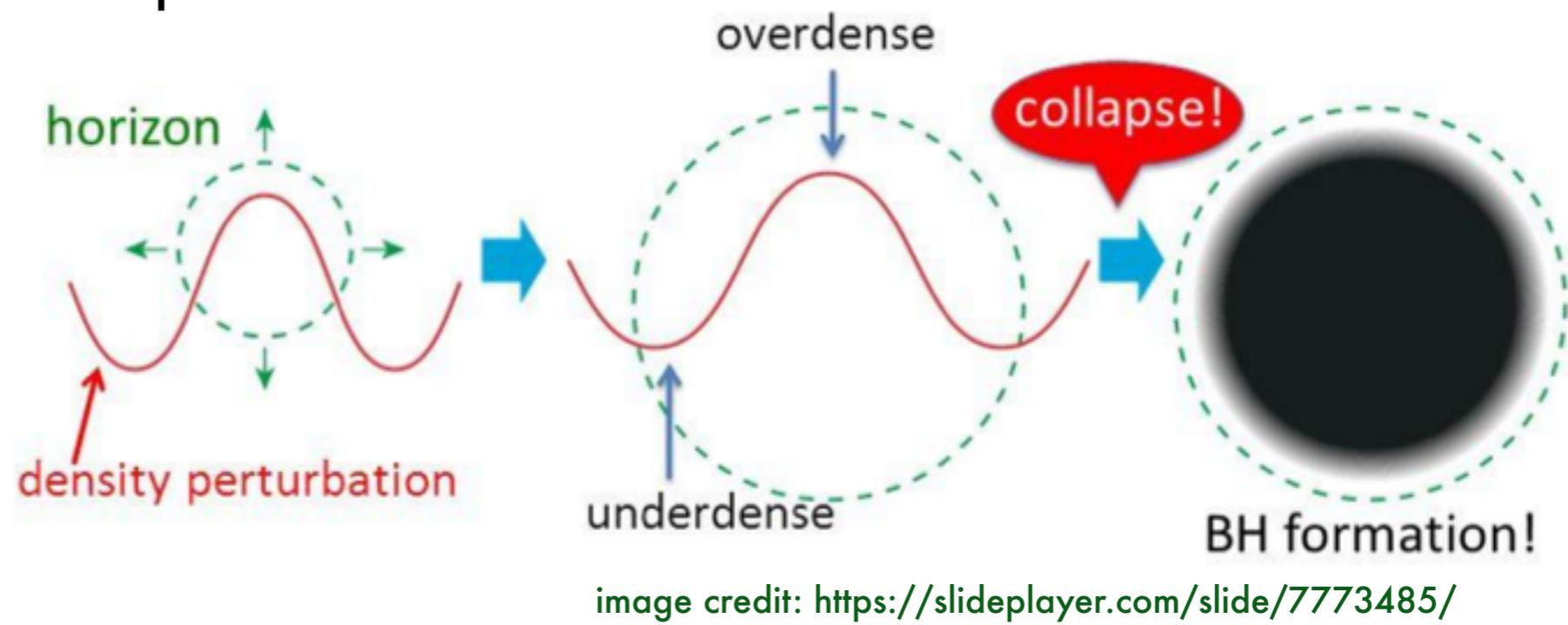
P. Adshead, P. Ralegankar and J. Shelton, 2019 2

# New Scenario



## motivated by many different scenarios

- collapse from inhomogeneities
- sudden reduction in the pressure
- collapse of cosmic loops
- bubble collisions



BHs can be fully characterized by their mass, spin, and charge

mass: of the order of the horizon mass

formed in a radiation dominated era:

spin: negligible

# Hawking Evaporation:

S. Hawking, 1975

$$T_{\text{BH}} = \frac{M_{\text{Pl}}^2}{8\pi M_{\text{BH}}}$$

$$\frac{d^2 u_i(E, t)}{dt dE} = \frac{g_i}{8\pi^2} \frac{E^3}{e^{E/T_{\text{BH}}} \pm 1}$$

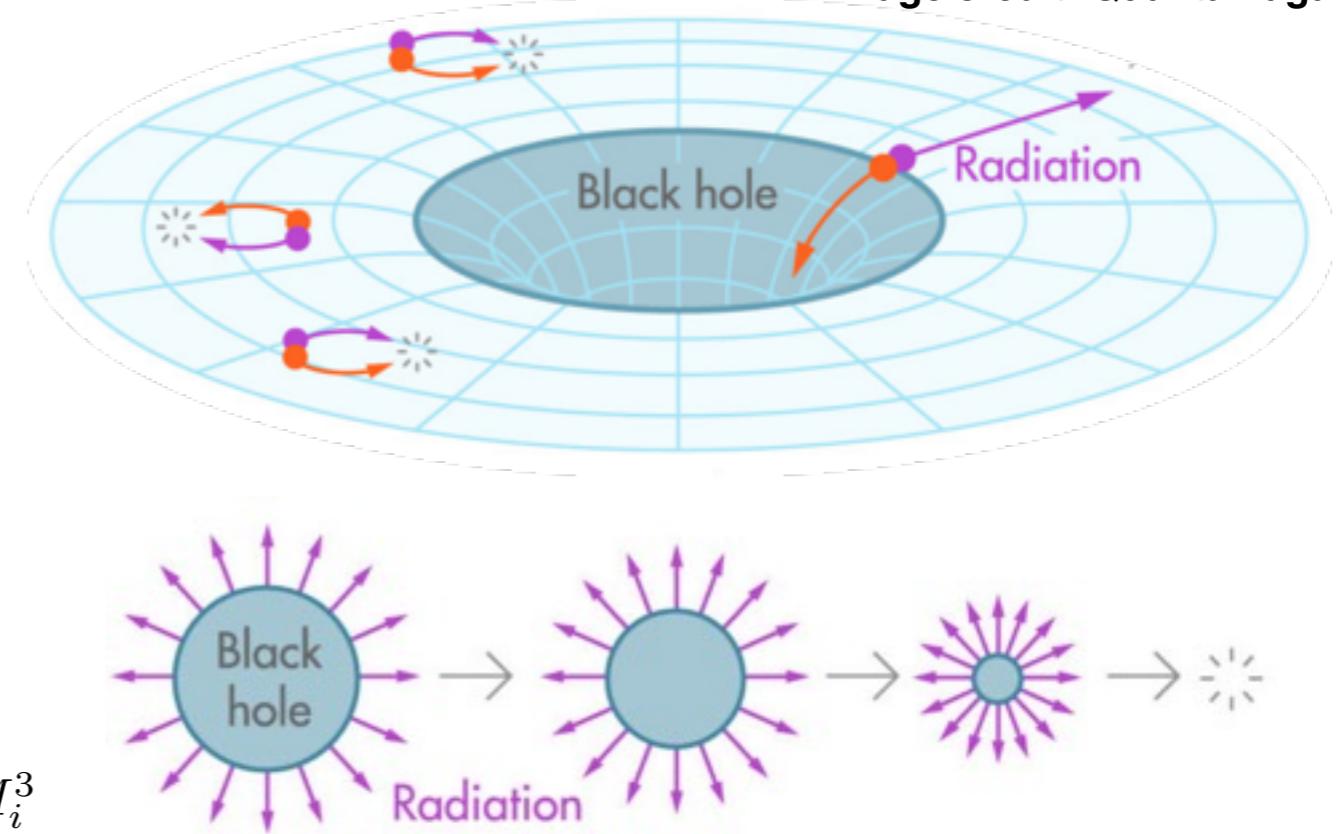
$$M(t) = M_i \left(1 - \frac{(t - t_i)}{\tau}\right)^{1/3}, \quad \tau = \frac{10240\pi}{g_*(T_{\text{BH}})} \frac{M_i^3}{M_{\text{Pl}}^4}$$

$$N_i = \frac{120\zeta(3)}{\pi^3} \frac{g_i}{g_*(T_{\text{BH}})} \frac{M_{\text{BH}}^2}{M_{\text{Pl}}^2}, \quad T_{\text{BH}} > m_i$$

$$N_i = \frac{15\zeta(3)}{8\pi^5} \frac{g_i}{g_*(T_{\text{BH}})} \frac{M_{\text{Pl}}^2}{m_i^2}, \quad T_{\text{BH}} < m_i$$

**CMB:** constraint on the size of Horizon at the end of inflation

$$H_I \lesssim 10^{-5} M_{\text{Pl}} \quad (\text{Planck})$$



$0.1 \text{ g} \lesssim M_{\text{BH}} \lesssim 10^9 \text{ g}$

Not constrained by cosmology

evaporate before BBN

# PBH Energy Content:

initial abundance of PBHs

$$\beta \equiv \frac{\rho_{\text{PBH}}(t_i)}{\rho_{\text{rad}}(t_i)}$$

$$\rho_{\text{PBH}} \propto a^{-3}, \quad \rho_{\text{rad}} \propto a^{-4}$$

$$\rho_{\text{PBH}}(t_{\text{early-eq}})/\rho_{\text{rad}}(t_{\text{early-eq}}) \sim 1 \quad t_{\text{early-eq}} \lesssim t_{\text{eva}}$$

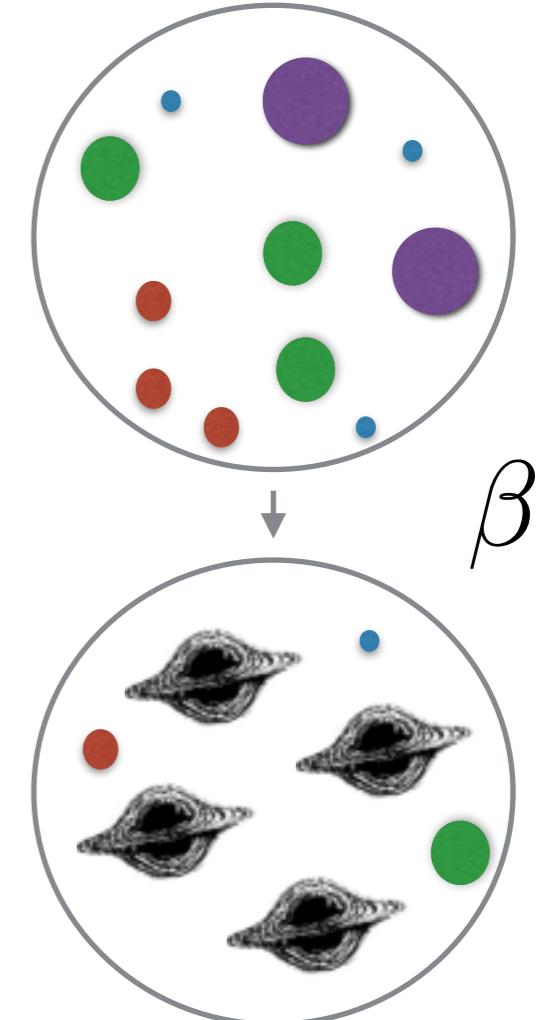
an initially radiation-dominated universe will eventually become matter-dominated if the PBHs are still around.

critical initial abundance

$$\beta_{\text{crit}} \sim \frac{M_{\text{Pl}}}{M_{\text{BH}}}$$

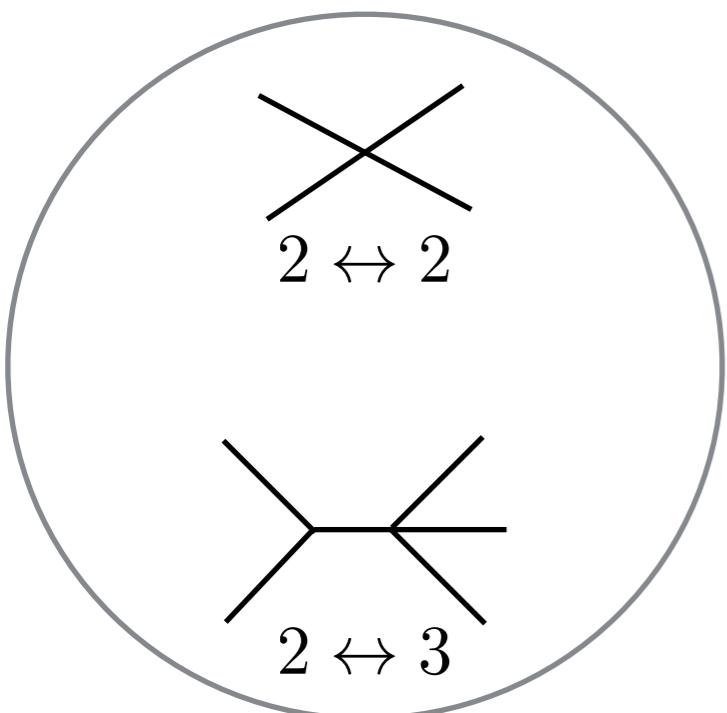
$\beta < \beta_{\text{crit}}$  evaporation happens in a RD Universe

$\beta \geq \beta_{\text{crit}}$  evaporation happens in a MD Universe



# Populating a self-interacting dark sector by relativistic and far from equilibrium particles:

$$\mathcal{L} = \frac{1}{2}\partial_\mu\chi\partial^\mu\chi - \frac{1}{2}m_\chi^2\chi^2 - \frac{m_\chi\lambda}{3!}\chi^3 - \frac{\lambda^2}{4!}\chi^4, \quad 1 \lesssim \lambda \lesssim 4\pi$$



**Dark Sector**

$t = \tau : n_\chi(\tau), \rho_\chi(\tau)$  populating  $(\lambda \simeq 1)$

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$2 \leftrightarrow 2$  kinetic equilibrium:  
leads to a temperature and non-zero chemical potential

$$\frac{\Gamma_{\chi,2\rightarrow 2}(t)}{H(t)} \sim \frac{n_\chi(t)/\bar{E}_\chi(t)^2}{H(t)} \propto a(t), \quad \tau \lesssim t \lesssim t_m$$


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$2 \rightarrow 3$  Chemical equilibrium:  
gives rise to a temperature and zero chemical potential

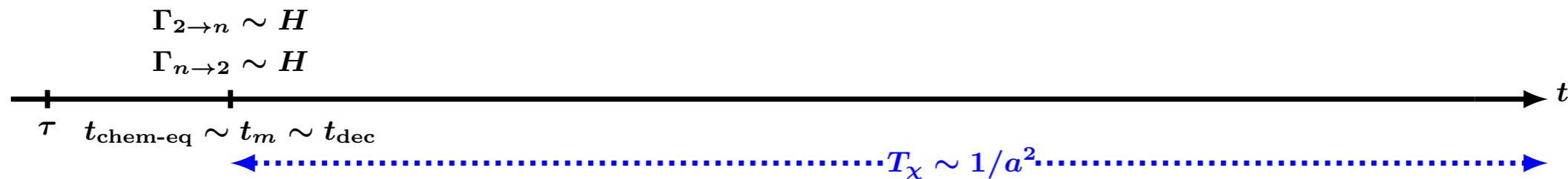
$$\frac{\Gamma_{\chi,2\rightarrow 3}(t)}{H(t)} \sim \frac{n_\chi(t)m_\chi^2/T_\chi^4(t)}{H(t)} \propto a^3(t), \quad t_{\text{kin-eq}} \lesssim t \lesssim t_m$$


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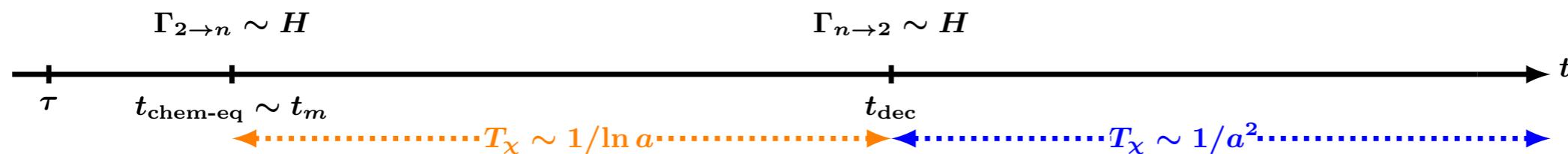
$3 \rightarrow 2$  Cannibalism/decoupling (freeze-out):

$$\frac{\Gamma_{\chi,3\rightarrow 2}(t)}{H(t)} \sim \frac{n_\chi^2(t)/m_\chi^5}{H(t)}, \quad t_m \lesssim t \lesssim t_{\text{dec}}$$

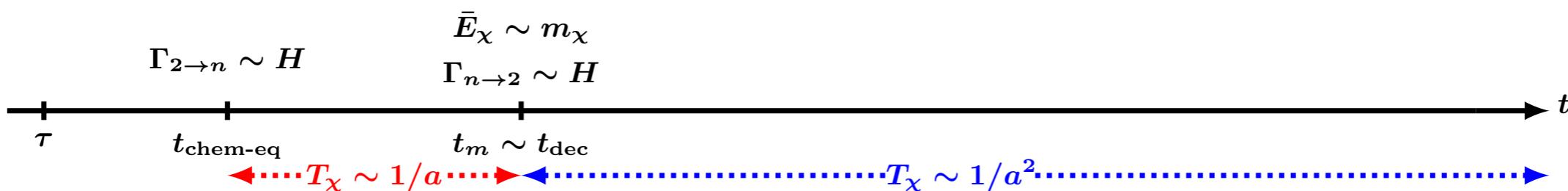
non-relativistic, no cannibalism (NRNC):



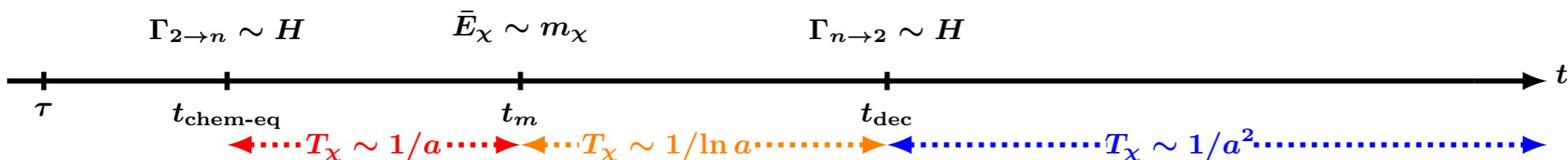
non-relativistic, with cannibalism (NRC):



relativistic, no cannibalism (RNC):



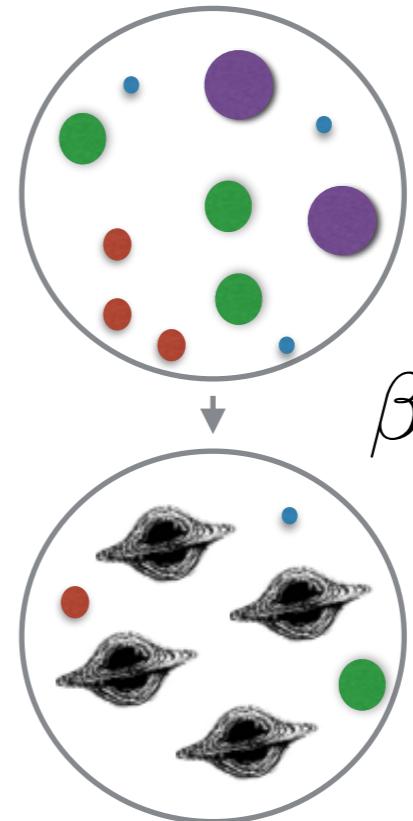
relativistic, with cannibalism (RC):



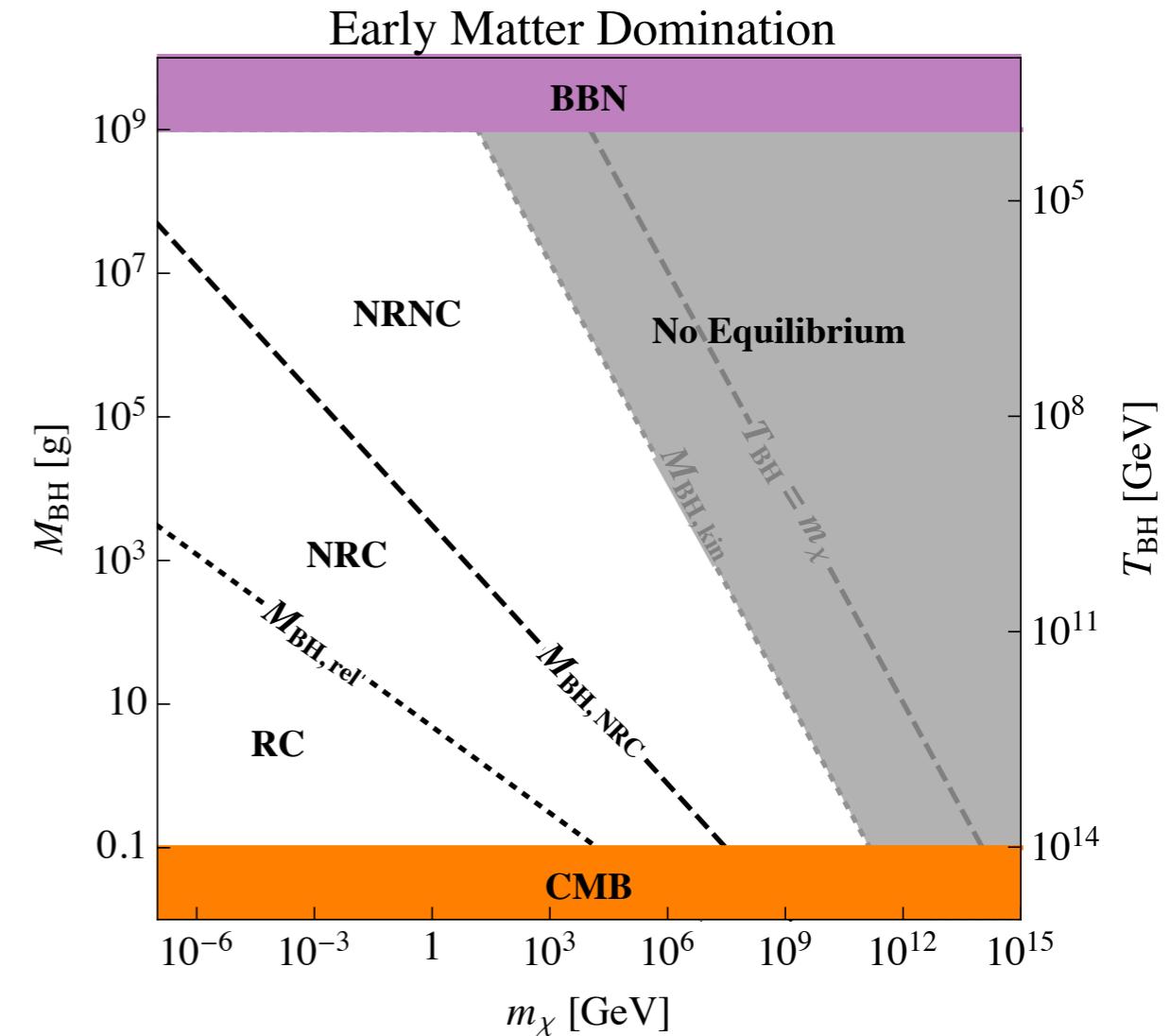
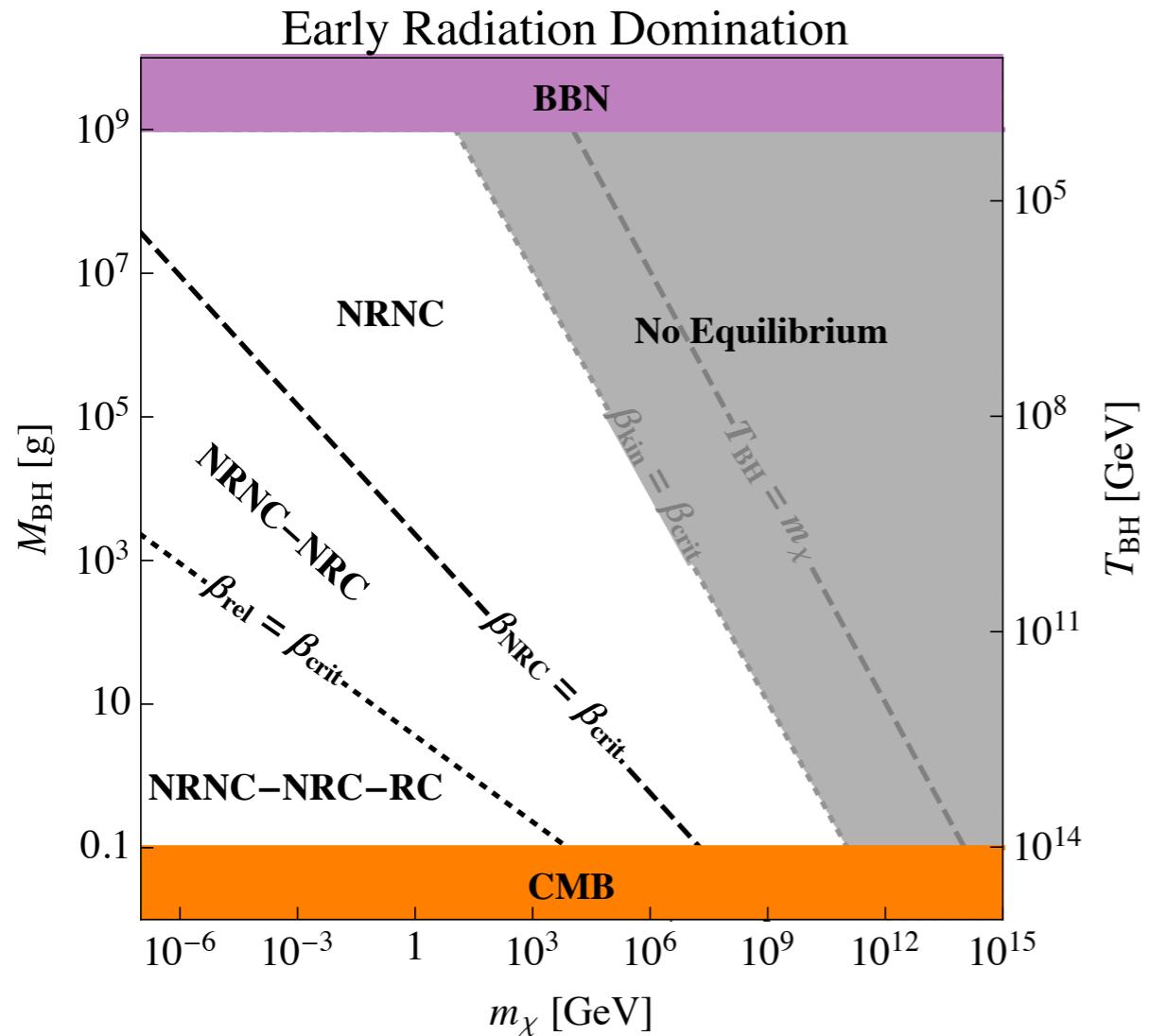
# Populating a self-interacting dark sector by Hawking evaporation of PBHs:

$$\beta \equiv \frac{\rho_{\text{PBH}}(t_i)}{\rho_{\text{rad}}(t_i)}$$

$$\beta_{\text{crit}} \sim \frac{M_{\text{Pl}}}{M_{\text{BH}}}$$



Thermal History	Early Radiation Domination	Early Matter Domination
NRNC (non-relativistic, no cannibalism)	$\beta_{\text{kin}} \lesssim \beta \lesssim \beta_{\text{NRC}}$	$M_{\text{BH, NRC}} \lesssim M_{\text{BH}} \lesssim M_{\text{BH, kin}}$
NRC (non-relativistic, cannibalism)	$\beta_{\text{NRC}} \lesssim \beta \lesssim \beta_{\text{rel}}$	$M_{\text{BH, rel}} \lesssim M_{\text{BH}} \lesssim M_{\text{BH, NRC}}$
RNC (relativistic, no cannibalism)	$\beta_{\text{rel}} \lesssim \beta \lesssim \beta_{\text{RC}}$	$M_{\text{BH, RC}} \lesssim M_{\text{BH}} \lesssim M_{\text{BH, rel}}$
RC (relativistic, cannibalism)	$\beta_{\text{RC}} \lesssim \beta \lesssim \beta_{\text{crit}}$	$M_{\text{BH}} \lesssim M_{\text{BH, RC}}$

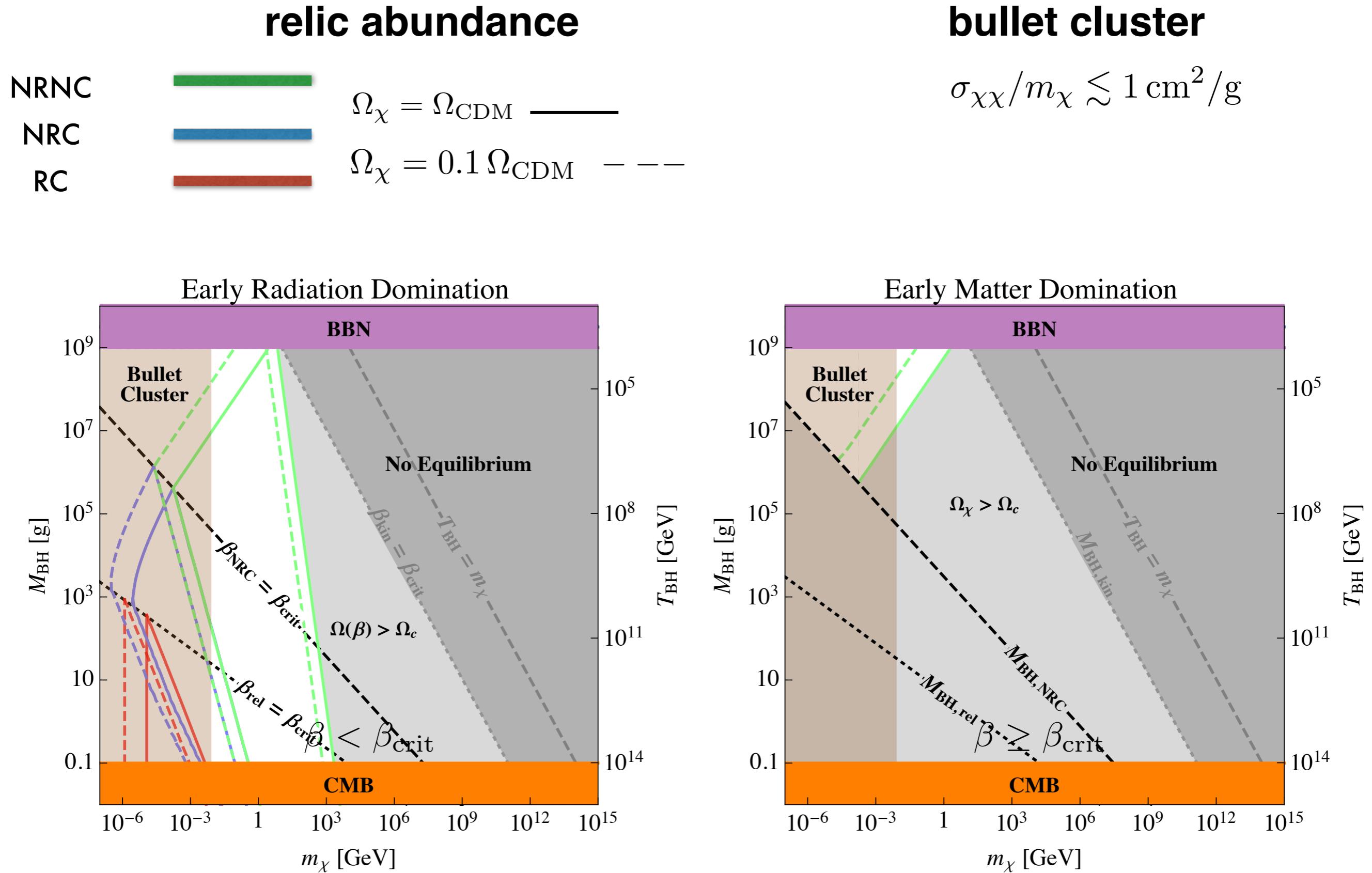


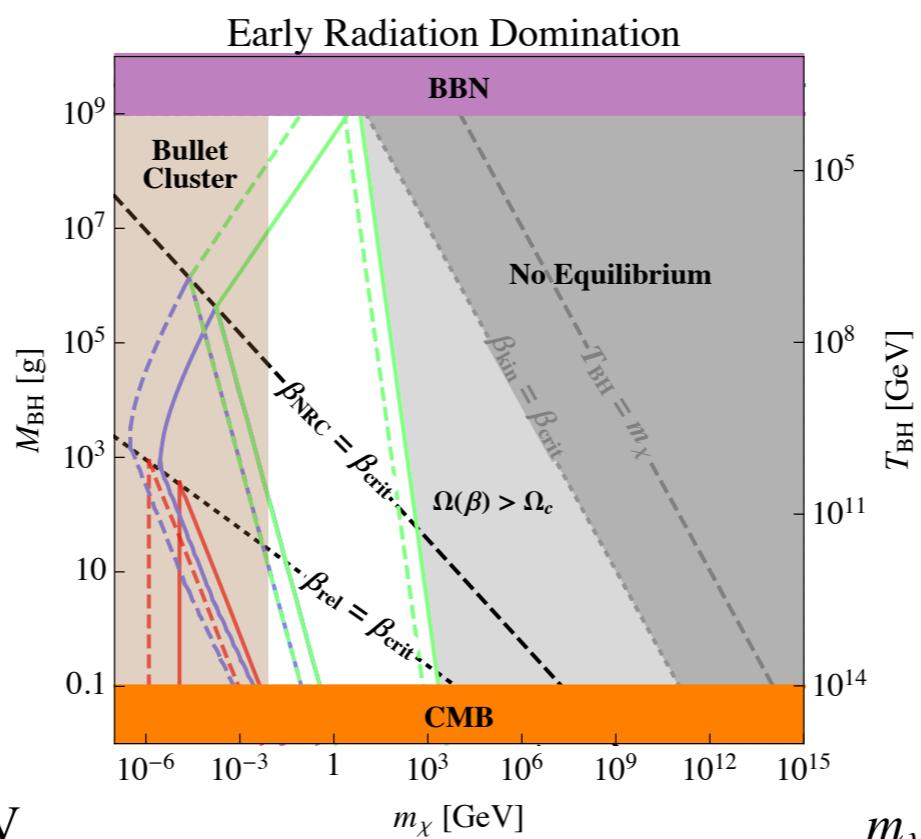
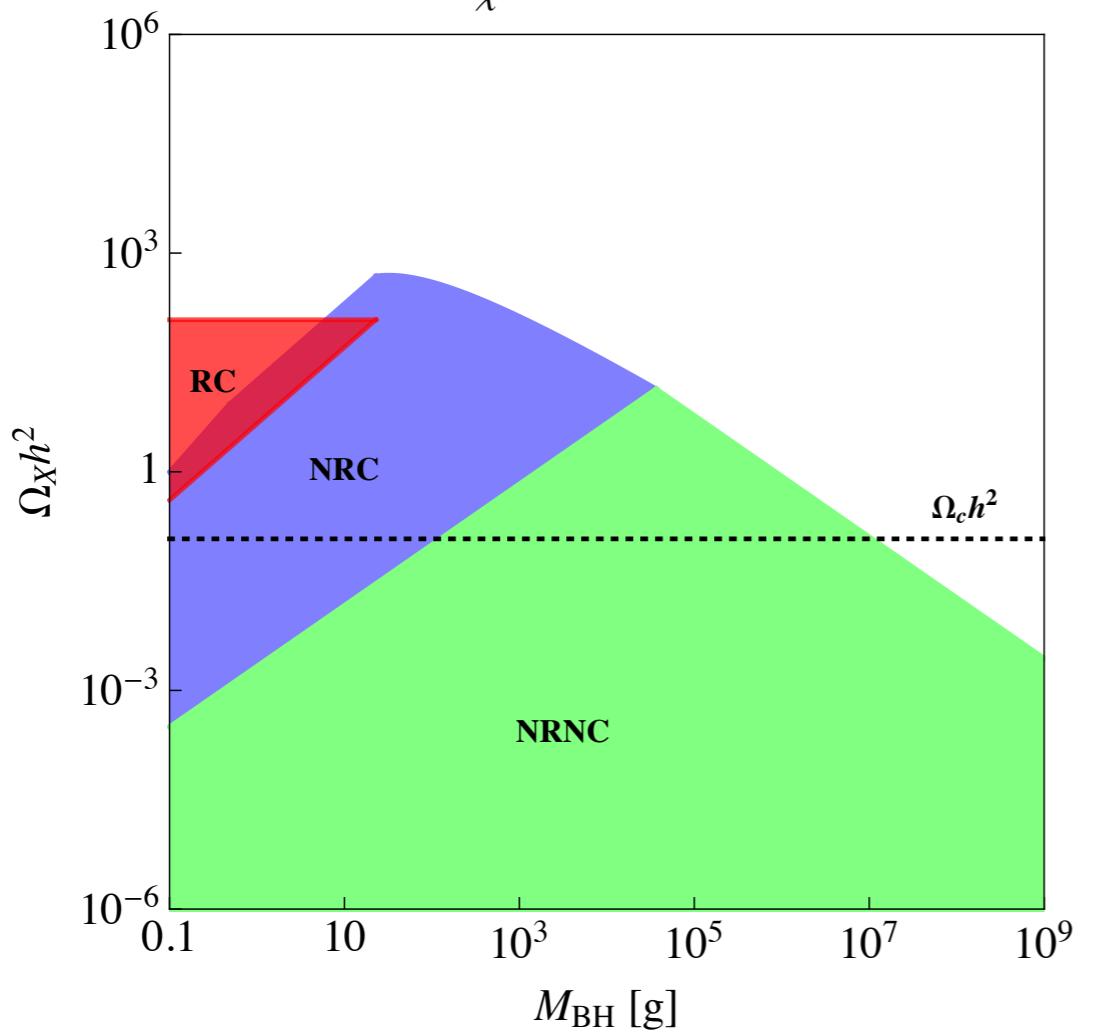
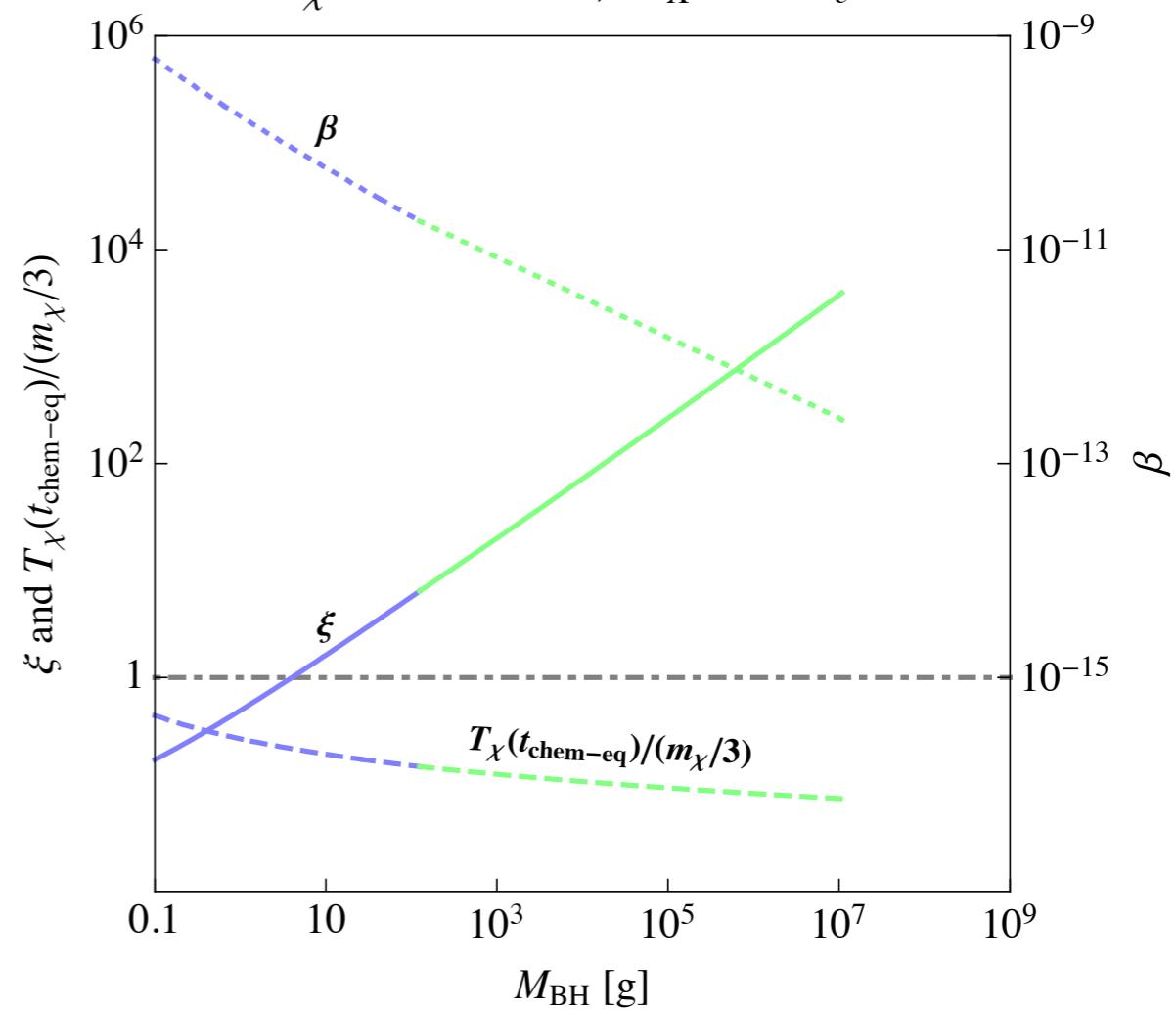
$$\beta < \beta_{\text{crit}}$$

$$\beta \geq \beta_{\text{crit}}$$

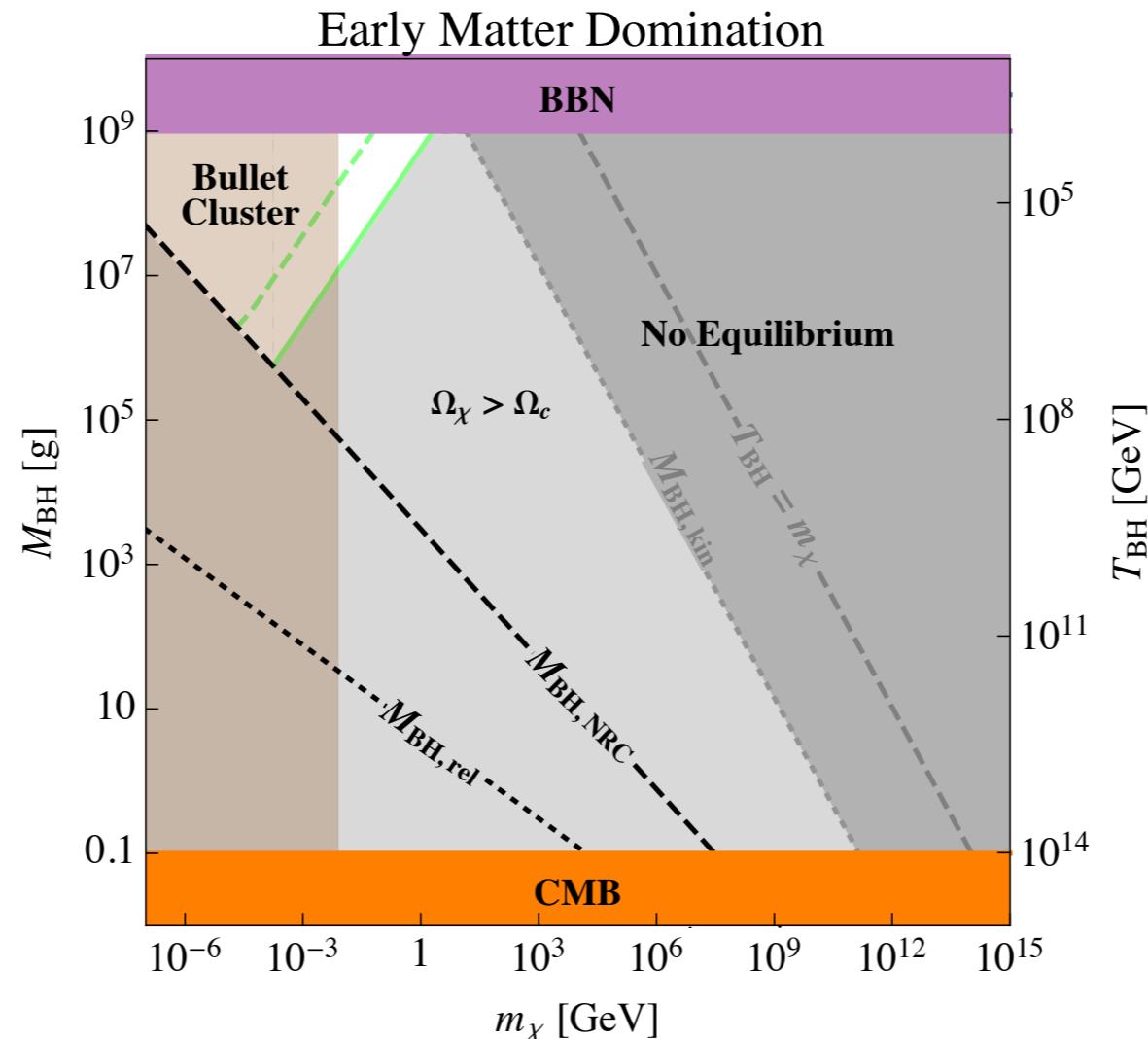
$T_{\text{BH}} < m_\chi$       no equilibrium  
**relativistic: cannibalism is inevitable**

## Constraints:




 $m_\chi = 10^{-2} \text{ GeV}$ 

 $m_\chi [ \text{GeV} ]$ 
 $m_\chi = 10^{-2} \text{ GeV}, \Omega_X h^2 = \Omega_c h^2$ 


$$\xi = \frac{T_\chi(t_{\text{chem-eq}})}{T_V(t_{\text{chem-eq}})}$$



$$(m_\chi = 8 \text{ MeV}, M_{\text{BH}} = 1.2 \times 10^7 \text{ g}, \xi = 3.3 \times 10^3)$$

$$(m_\chi = 2.1 \text{ GeV}, M_{\text{BH}} = 10^9 \text{ g}, \xi = 5.4 \times 10^5)$$

$$\xi = \frac{T_\chi(t_{\text{chem-eq}})}{T_V(t_{\text{chem-eq}})}$$

## **Early radiation-dominated Universe:**

the NRNC and NRC thermal histories can accommodate  $\xi$  less than, equal to, or greater than one.

For RC thermal histories, on the other hand, only  $\xi < 1$  is allowed.

## **Early matter (PBH)-dominated Universe:**

an NRNC or NRC thermal history forces  $\xi > 1$ , while an RC thermal history forces  $\xi = 1$ .

$$\xi = \frac{T_\chi(t_{\text{chem-eq}})}{T_V(t_{\text{chem-eq}})}$$

# Conclusions

- PBHs are motivated by many scenarios.
- Hawking evaporation of PBHs can populate a self-interacting dark sector which is hotter or colder than the visible sector.
- The asymmetry in temperatures of the two sectors is persistent.
- DM with right relic abundance points to MeV-TeV mass range.