Higgs bosons with large couplings to up type quarks

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Flavor symmetries of the quark sector

$$\diamondsuit$$
 $\begin{pmatrix} u \\ d \end{pmatrix}$, $\begin{pmatrix} c \\ s \end{pmatrix}$, $\begin{pmatrix} t \\ b \end{pmatrix}$ Quarks are grouped in 3 Families

- \diamondsuit Flavor Symmetries of quark sector: $U(3)^3 \equiv U(3)_Q \times U(3)_{\bar{u}} \times U(3)_{\bar{d}}$
 - → The Gauge sector respects all the flavor symmetries
 - ightarrow Yukawa interactions ("spurions") break the flavor symmetries i.e. $Y_u\sim (3,\bar{3},1)$ and $Y_d\sim (3,1,\bar{3})$
- Choosing a flavor basis we can write:

$$\lambda^u \equiv V^T Y^u$$
 and $\lambda^d \equiv Y^d$

 $V \rightarrow \mathsf{CKM}$ matrix



Extended Flavor Sectors

- ♦ Flavorful New Physics generically leads to Flavor Changing Neutral Currents
 - FCNCs are suppressed in the SM
 - Assumptions on the structure of NP are needed
 - Most common assumption: Minimal Flavor Violation (MFV)
- ♦ Minimal Flavor Violation:
 - Any NP spurions are tensor products of y^u and y^d
 - New Physics scale is allowed to be closer to the energy reach of current experiments
 - FCNCs follow SM suppression

Example: $\frac{\left(\bar{d}_L V_{tb} y_t^2 V_{td}^* b_L\right)^2}{\Lambda_{NP}^2} \to \Lambda_{NP} \sim \textit{TeV Vs} \ \Lambda_{NP} \sim 10^3 - 10^4 \textit{TeV} \ \text{for generic } \mathcal{O}(1) \ \text{coupling}$



Minimal Flavor Violation

What we know:

→ FCNCs bounds force flavor alignment of any New Physics couplings

Using MFV:

- → Constrained to SM yukawa hierarchy
- → Limited collider phenomenology

Spontaneous Flavor Violation PhysRevLett.123.031802

- ♦ Key idea: The only source of flavor changing processes is the CKM matrix
- ♦ What we gain: New spurions that are allowed to deviate from SM hierarchy

Spontaneous Flavor Violation

♦ Flavor-changing processes & CP breaking come from the wave function renormalization of right handed SM quarks

$$\mathcal{L} \supset i Z^u_{ij} \, ar{u}^\dagger_i ar{\sigma}^\mu D_\mu ar{u}_j + ar{d}^\dagger_i ar{\sigma}^\mu D_\mu ar{d}_i - \left[\lambda^u_{ij} Q_i H ar{u}_j - \lambda^d_{ij} Q_i H^c ar{d}_j + h.c.
ight]$$

♦ Up-Type:

$$Z^u_{ij} \bar{u}^\dagger_i \bar{\sigma}^\mu D_\mu \bar{u}_j$$

$$\bar{u} o \left(\sqrt{Z^u}\right)^{-1} \bar{u}$$

$$\lambda^u \to \left(\sqrt{Z^u}\right)^{-1} \lambda^u = V^T Y^u$$

♦ Down-Type:

$$Z^d_{ij} \bar d^\dagger_i \bar \sigma^\mu D_\mu \bar d_j$$

$$ar{d}
ightarrow \left(\sqrt{Z^d}
ight)^{-1} ar{d}$$

$$\lambda^d \to \left(\sqrt{Z^d}\right)^{-1} \lambda^d = V^T Y^d$$



Spontaneous Flavor Violating Two Higgs Doublet Model

 \Diamond Two complex scalar fields: H_{α} , $\alpha = 1, 2$

$$D_{\mu}H_{\alpha}^{\dagger}D^{\mu}H_{\alpha}-V\left(H_{1},H_{2}\right)-\left[\lambda_{\alpha ij}^{u}Q_{i}H_{\alpha}\bar{u}_{j}-\lambda_{\alpha ij}^{d\dagger}Q_{i}H_{\alpha}^{c}\bar{d}_{j}-\lambda_{\alpha ij}^{\ell\dagger}L_{i}H_{\alpha}^{c}\bar{\ell}_{j}+h.c.\right] \tag{1}$$

♦ Up-Type:

$$\lambda_2^u = \xi V^T Y^u$$

$$\lambda_2^d = K^d \equiv diag(\kappa_d, \kappa_s, \kappa_b)$$

♦ Down-Type:

$$\lambda_2^u = V^T K^u \equiv V^T diag(\kappa_u, \kappa_c, \kappa_t)$$

 $\lambda_2^d = \xi Y^d$

 \checkmark The yukawa couplings of H_2 are free parameters, allowed to deviate from SM hierarchy

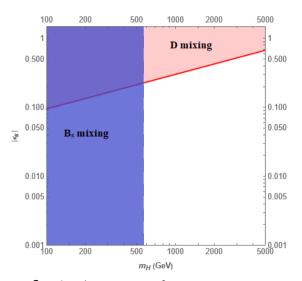
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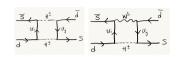
Flavor Phenomenology of neutral meson mixing

- \Diamond Higgs mass eigenstates: H^{\pm} , H, h, & A
- $\Diamond \ \mathcal{H}_{eff}^{\Delta F=2} = \sum_{i=1} C_i \mathcal{O}_i$
 - $\bullet \ \mathcal{O}_1 = \left(\bar{d}_J \gamma_\mu P_L d_I\right) \left(\bar{d}_J \gamma^\mu P_L d_I\right)$
 - $\mathcal{O}_2 = \left(\bar{d}_J P_L d_I\right) \left(\bar{d}_J P_R d_I\right)$
 - And a few more
- \Diamond Use bounds on the C_i of these operators [UTfit Collaboration, JHEP 03 (2008) 049]

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Bounds from neutral meson mixing (up-type)



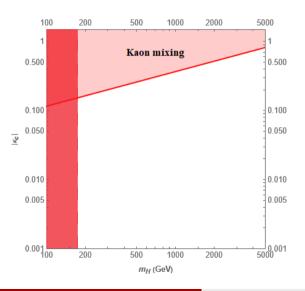


- D meson mixing \rightarrow constraints $(\lambda_2^u = \xi V^T Y^u, \lambda_2^d = K^d)$
- ullet Couplings of ~ 0.1 are allowed for masses $500\,GeV$
- Depending on value of ξ the B_s bounds become irrelevant

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Bounds from neutral meson mixing (down-type)



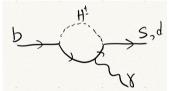
- Kaon mixing is important $(\lambda_2^u = V^T K^u, \lambda_2^d = \xi Y^d)$
- ullet Couplings of ~ 0.1 are allowed for masses $\lesssim 200 GeV$
- Relevant to recent ATLAS result for charm yukawa
 CERN-EP-2021-251

Set:
$$\xi = 1$$
, $\kappa_d = \kappa_b = 0$

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More bounds?

- \Diamond Previously \to bounds from neutral meson mixing
- \diamondsuit Can use $B o X_{s,d} \gamma$ transitions $\left[\mathcal{O}_7 \sim rac{e}{16\pi^2} m_b \left(Q_2 \sigma_{\mu\nu} ar{d}_3
 ight)^\dagger F^{\mu
 u}
 ight]$



Work in progress

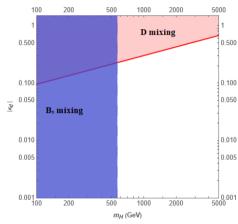
Conclusions

- ♦ Extended Higgs sectors do not need to follow SM coupling hierarchy
- ♦ Viable parameter space for flavorful couplings of significant size
- Motivation for studying theories with large couplings to light quarks

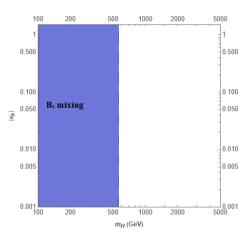
Thank you!

Extra Slides

Bounds from neutral meson mixing (up-type)



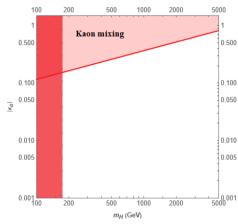
Set:
$$\xi = 1$$
, $\kappa_s = \kappa_b = 0$



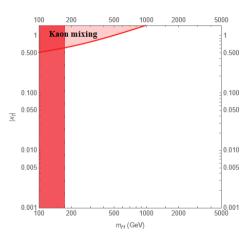
Set:
$$\xi = 1$$
, $\kappa_d = \kappa_s = 0$



Bounds from neutral meson mixing (down-type)



Set: $\xi = 1$, $\kappa_t = \kappa_c = 0$



Set:
$$\xi = 1$$
, $\kappa_{\mu} = \kappa_{c} = 0$

