

# PROBING B-ANOMALIES VIA DIMUON TAILS AT A FUTURE COLLIDER

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in collaboration with Bradley Garland, Sebastian Jäger and Charanjit K. Khosa  
based on arXiv:2112.05127

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# INTRODUCTION

- B-anomalies:

$$R_{K^{(*)}} = \frac{\text{BR}(B \rightarrow K^{(*)} \bar{\mu} \mu)}{\text{BR}(B \rightarrow K^{(*)} \bar{e} e)}$$

Observable	SM	Measurement	Experiment	
$BR(B_s \rightarrow \mu^+ \mu^-)$	$(3.63 \pm 0.13) \times 10^{-9}$	$(2.8 \pm 0.3) \cdot 10^{-9}$	average [4] of ATLAS [5], CMS [6] and LHCb [7]	$4.2\sigma$
$R_K[1.1, 6]$ $R_K[1, 6]$	$1.0004^{+0.0008}_{-0.0007}$	$0.85 \pm 0.04$ $1.03 \pm 0.28$	LHCb [8] Belle [9]	$3.1\sigma$
$R_{K^*}[0.045, 1.1]$	$0.920^{+0.007}_{-0.006}$	$0.66 \pm 0.11$ $0.52 \pm 0.37$	LHCb [10] Belle [11]	$2.2\sigma$
$R_{K^*}[1.1, 6]$	$0.996 \pm 0.002$	$0.69 \pm 0.12$ $0.96 \pm 0.45$	LHCb [10] Belle [11]	$2.4\sigma$

- Goal: model independent and conservative, focus on inclusive  $pp \rightarrow \mu^+ \mu^-$

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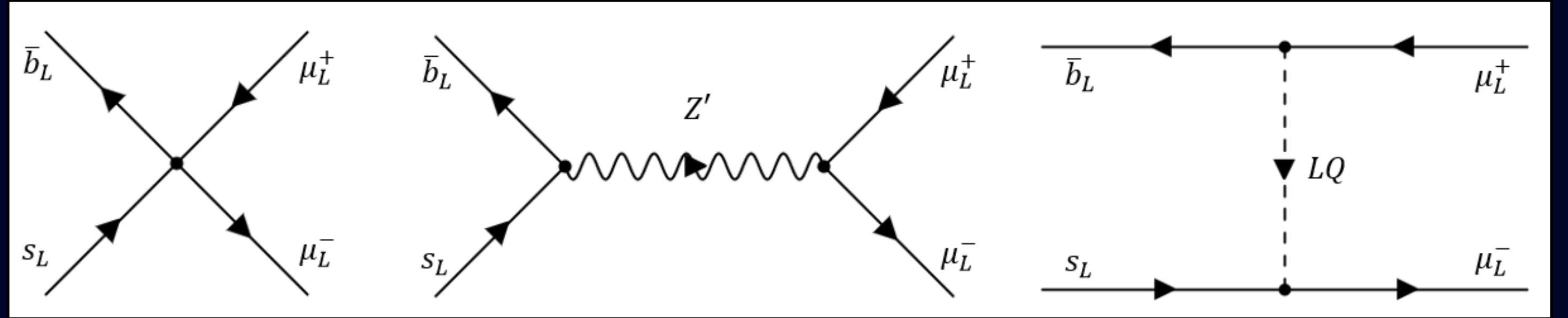
- Goal: model independent and conservative, focus on inclusive  $pp \rightarrow \mu^+ \mu^-$
- Described by BSM four-fermion contact interactions in the low-energy effective weak Hamiltonian involving muons only

$$\begin{aligned} \mathcal{L}_{\text{eff}}^{\text{BSM}} = - \mathcal{H}_{\text{eff}}^{\text{BSM}} &= \frac{4G_F}{\sqrt{2}} V_{ts} V_{tb}^* \frac{e^2}{16\pi^2} \left\{ C_9^{\text{NP}} \left( \bar{b}_L \gamma_\mu s_L \right) \left( \bar{\mu} \gamma^\mu \mu \right) + C_{10}^{\text{NP}} \left( \bar{b}_L \gamma_\mu s_L \right) \left( \bar{\mu} \gamma^\mu \gamma^5 \mu \right) + \dots \right\} \\ &\equiv \frac{1}{\Lambda_{\text{LL}}^2} \left( \bar{b}_L \gamma_\mu s_L \right) \left( \bar{\mu}_L \gamma^\mu \mu_L \right) + \frac{1}{\Lambda_{\text{LR}}^2} \left( \bar{b}_L \gamma_\mu s_L \right) \left( \bar{\mu}_R \gamma^\mu \mu_R \right) + \dots \end{aligned}$$

- Excellent fit of  $b \rightarrow s \ell^+ \ell^-$  data obtained by LH muons only give  $C_9^{\text{NP}} = -C_{10}^{\text{NP}} \equiv C_L \approx -0.40$ , PhysRevD.104.035029, arXiv:2103.13370, arXiv:2104.08921



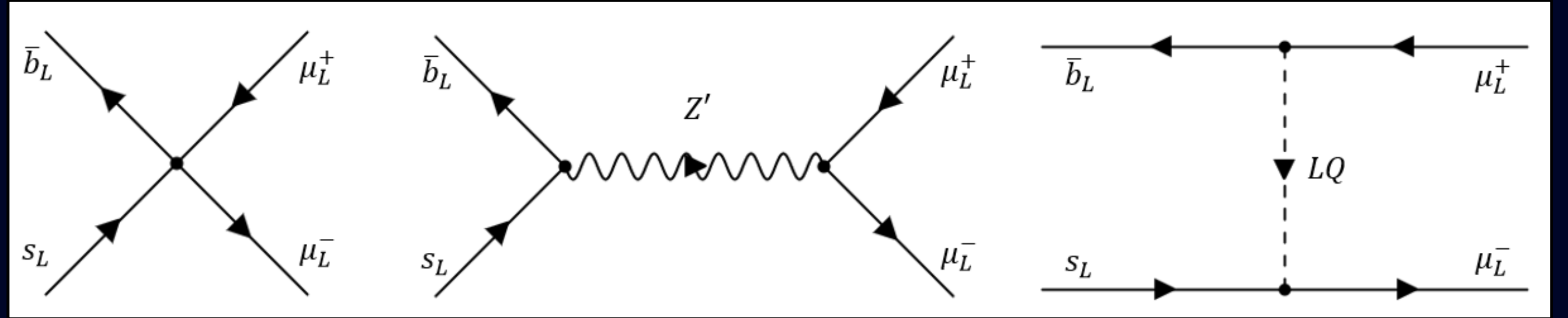
# INTRODUCTION



- Rare  $B$  decay dataset points to the presence of a left-handed contact interaction

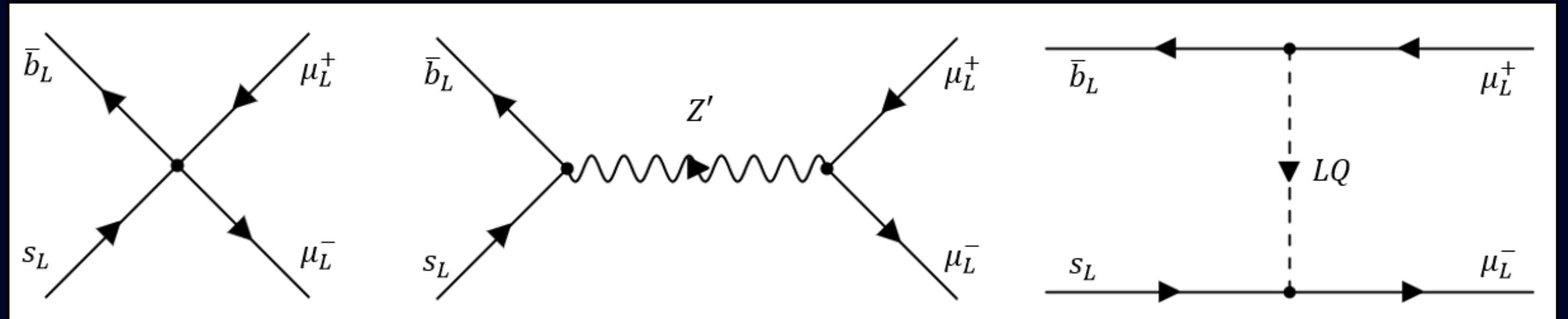
$$\mathcal{L}_{\text{eff}} \supset \frac{1}{\Lambda^2} \mathcal{O}_{LL}, \text{ where } \mathcal{O}_{LL} = \left( \bar{b}_L \gamma_\mu s_L \right) \left( \bar{\mu}_L \gamma^\mu \mu_L \right) \text{ and } \Lambda = (39 \pm 4) \text{ TeV}$$

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- Low-energy effective description of SM extension, eg.  $Z'$  or (spin-1 or spin-0) leptoquark
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- Perturbative unitarity allows, for  $\Lambda = 39 \text{ TeV}$ , mediators as heavy as 105 TeV
- Look for effects of heavy NP in "low-energy" tails of  $m_{\bar{\mu}\mu}$  within EFT
- To what extent does increased  $\sqrt{s}$  improve the sensitivity to  $bs\mu\mu$  contact interactions?

# SMEFT

$$\mathcal{L}^{\text{SMEFT}} = \mathcal{L}^{\text{SM}} + \sum_i c_i^{(6)} \mathcal{O}_i^{(6)} + \sum_j c_j^{(8)} \mathcal{O}_j^{(8)} + \dots$$

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- Choose a basis such that  $b \rightarrow s\mu^+\mu^-$  governed by single Wilson coeff.  $C_{sb}^+$

$$\begin{aligned} \mathcal{L}^{\text{SMEFT}} \supset & C_{ij}^+ (\bar{d}_L^j \gamma_\rho d_L^i) (\bar{\mu}_L \gamma^\rho \mu_L) + C_{ij}^- (\bar{d}_L^j \gamma_\rho d_L^i) (\bar{\nu}_\mu \gamma^\rho \nu_\mu) \\ & + \sum_{k,l} V_{ki}^* C_{ij}^+ V_{lj} (\bar{u}_L^l \gamma_\rho u_L^k) (\bar{\nu}_\mu \gamma^\rho \nu_\mu) + \sum_{k,l} V_{ki}^* C_{ij}^- V_{lj} (\bar{u}_L^l \gamma_\rho u_L^k) (\bar{\mu}_L \gamma^\rho \mu_L) \end{aligned}$$



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$$C^+ = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & C_{sb}^+ \\ 0 & C_{sb}^{+*} & 0 \end{pmatrix}$$

- Consider minimal/conservative scenario

$$\mathcal{L}^{\text{SMEFT}} \supset C_{sb}^+ (\bar{b}_L \gamma_\rho s_L) (\bar{\mu}_L \gamma^\rho \mu_L) + \sum_{k,l} V_{ks}^* C_{sb}^+ V_{lb} (\bar{u}_L^l \gamma_\rho u_L^k) (\bar{\nu}_\mu \gamma^\rho \nu_\mu) + \text{h.c.}$$

$$C^- = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

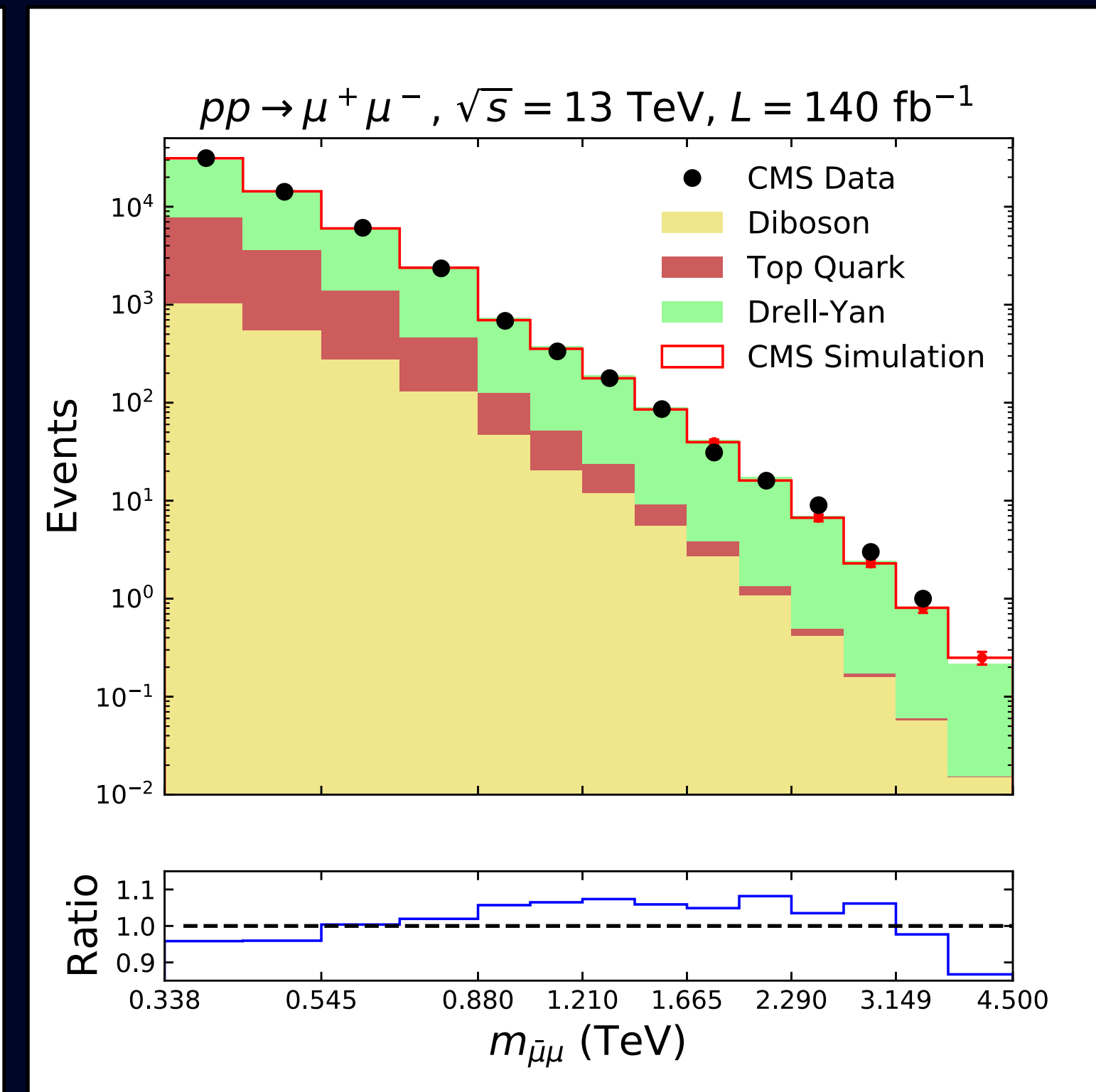
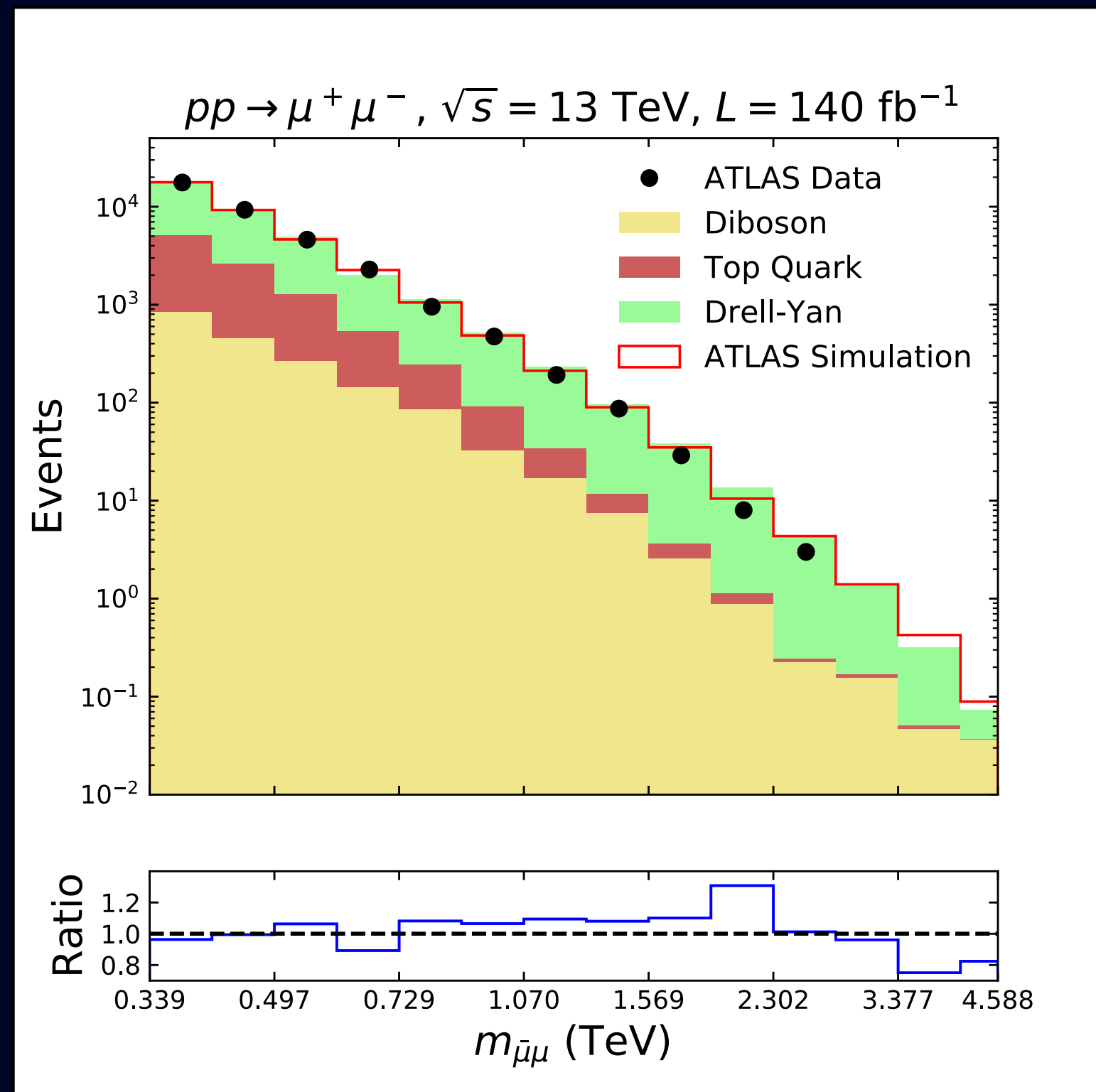
where  $C_{sb}^+ = \frac{1}{\Lambda^2}$

# SET-UP

- Main SM contributions to  $pp \rightarrow \mu^+ \mu^-$ :
  1. Drell–Yan via  $Z/\gamma^*$  exchange (~80-90%)
  2. diboson ( $ZZ$ ,  $WZ$  and  $WW$ ) production
  3. top-quark production ( $\bar{t}t$  &  $tW$ )
- Use MadGraph5\_aMC@NLO including fixed order NLO QCD and EW corrections for DY and LO for the rest
- NNPDF31\_nlo\_as\_0118\_luxqed 5 flavour
- Consider  $\sqrt{s} = 13$  TeV LHC,  $\sqrt{s} = 14$  TeV HL-LHC and FCC-hh with  $\sqrt{s} = 100$  TeV
- Signal generated by adding dim 6 ops. to default SM UFO file and including  $R_2$  terms by hand to yield gauge-invariant results

# VALIDATION

- Check that our background simulation is in agreement with CMS and ATLAS
- $p_T > 30$  GeV and  $|\eta| < 2.5$  for ATLAS
- $p_T > 53$  GeV and  $|\eta| < 2.4$  for CMS
- Our background is in excellent agreement with CMS (within  $\sim 10\%$ ), slightly worse for ATLAS due to more complicated selection criteria
- CMS JHEP07 (2021) 208, ATLAS JHEP11 (2020) 006



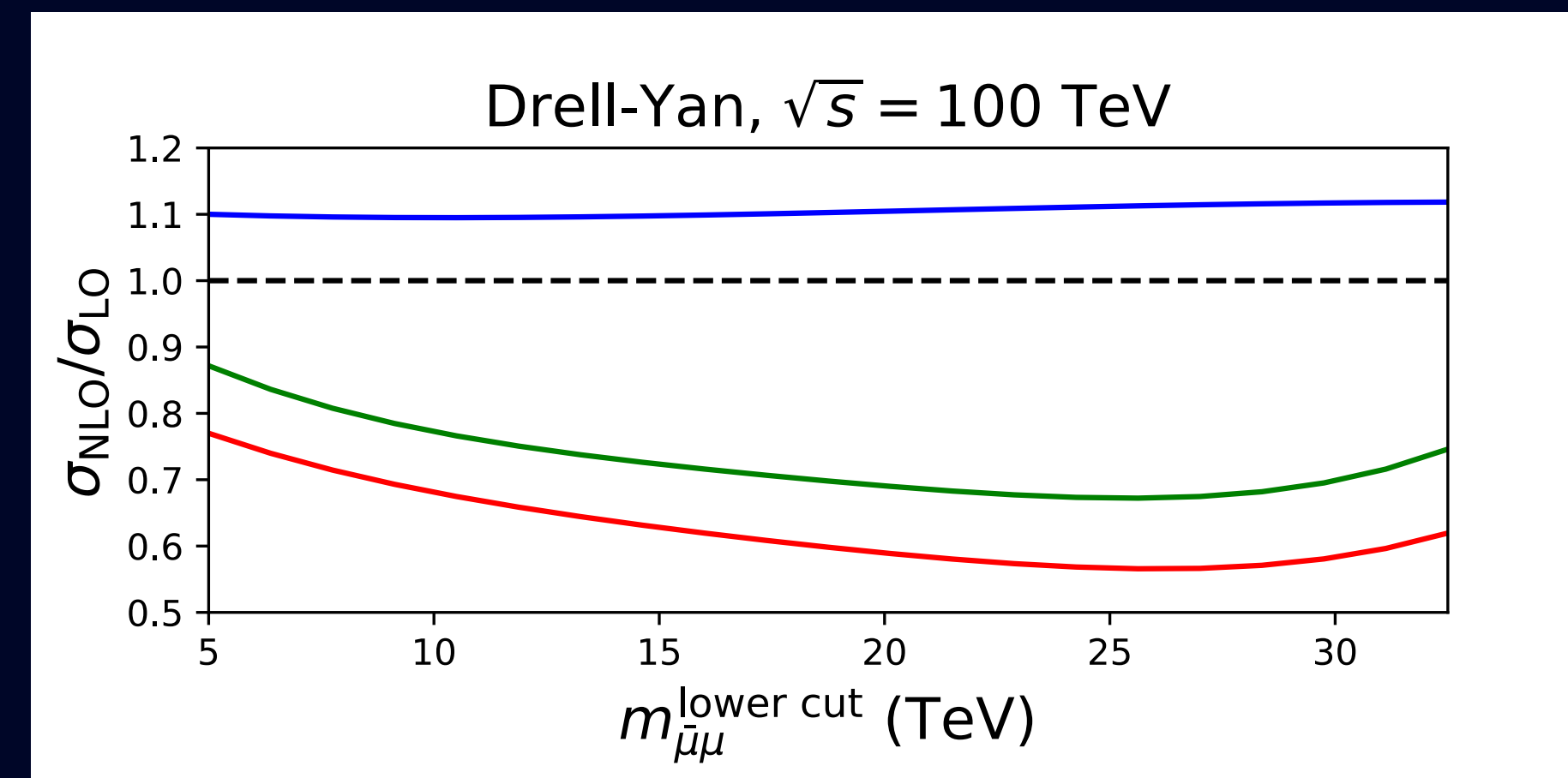
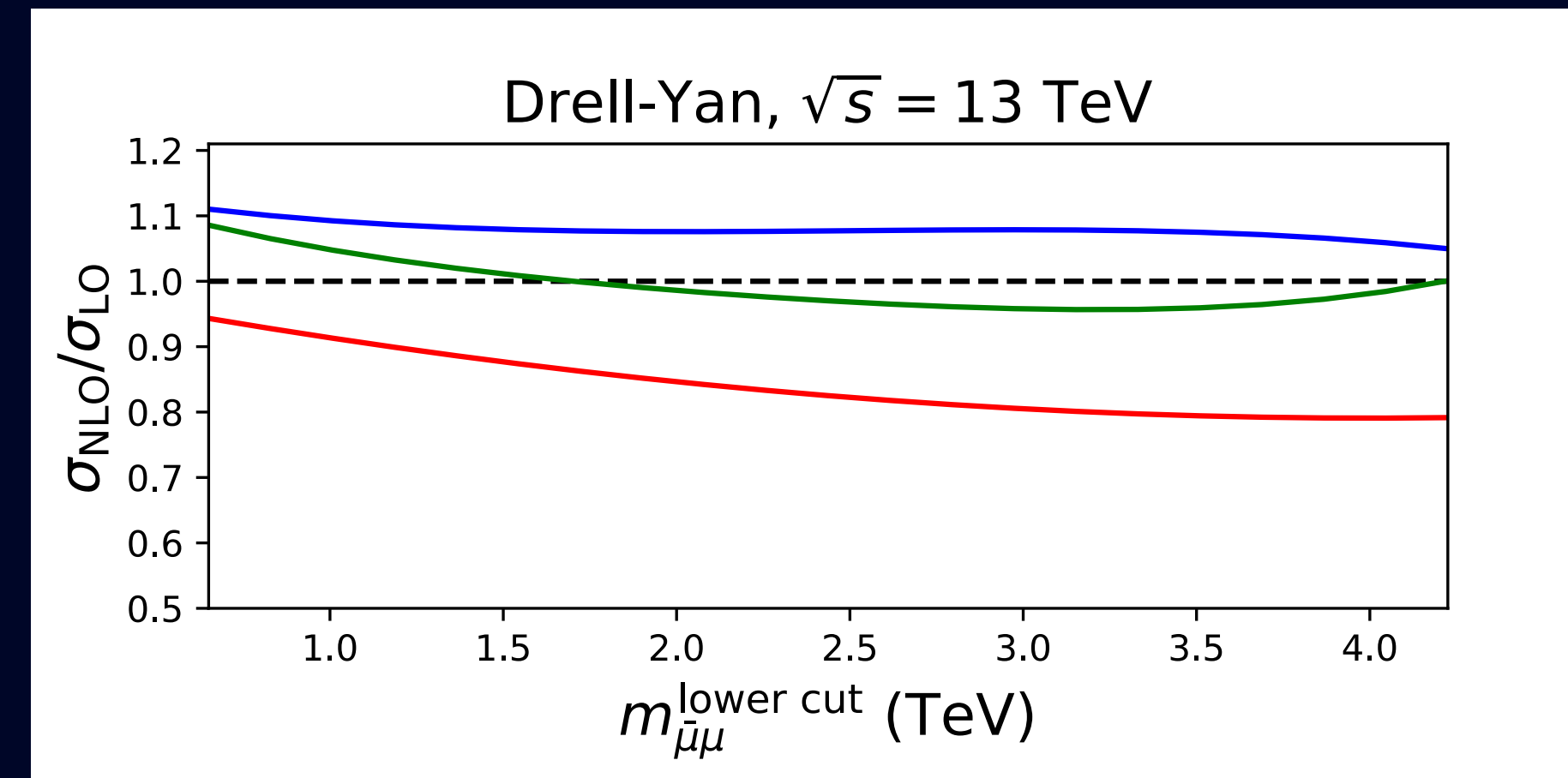
# INCLUSION OF NLO

- DY process receives large EW corrections in the high  $m_{\bar{\mu}\mu}$  region
- NLO-EW corrections contain large negative Sudakov double logarithms that reduce  $d\sigma/dm_{\bar{\mu}\mu}$  of the inclusive DY process at large values of  $m_{\bar{\mu}\mu}$

Blue: NLO-QCD

Red: NLO-EW

Green: NLO QCD&EW



# EFT VALIDITY

- $m_{\bar{\mu}\mu}^{\max}$  cannot be taken to be arbitrarily large
- Tree-level unitarity implies upper bound for  $m_{\bar{\mu}\mu} < \Lambda_*$
- Model dependent bound from UV mediator mass, e.g.  $m_{\bar{\mu}\mu} < m_{Z'}$  or  $m_{\bar{\mu}\mu} < m_{LQ}$



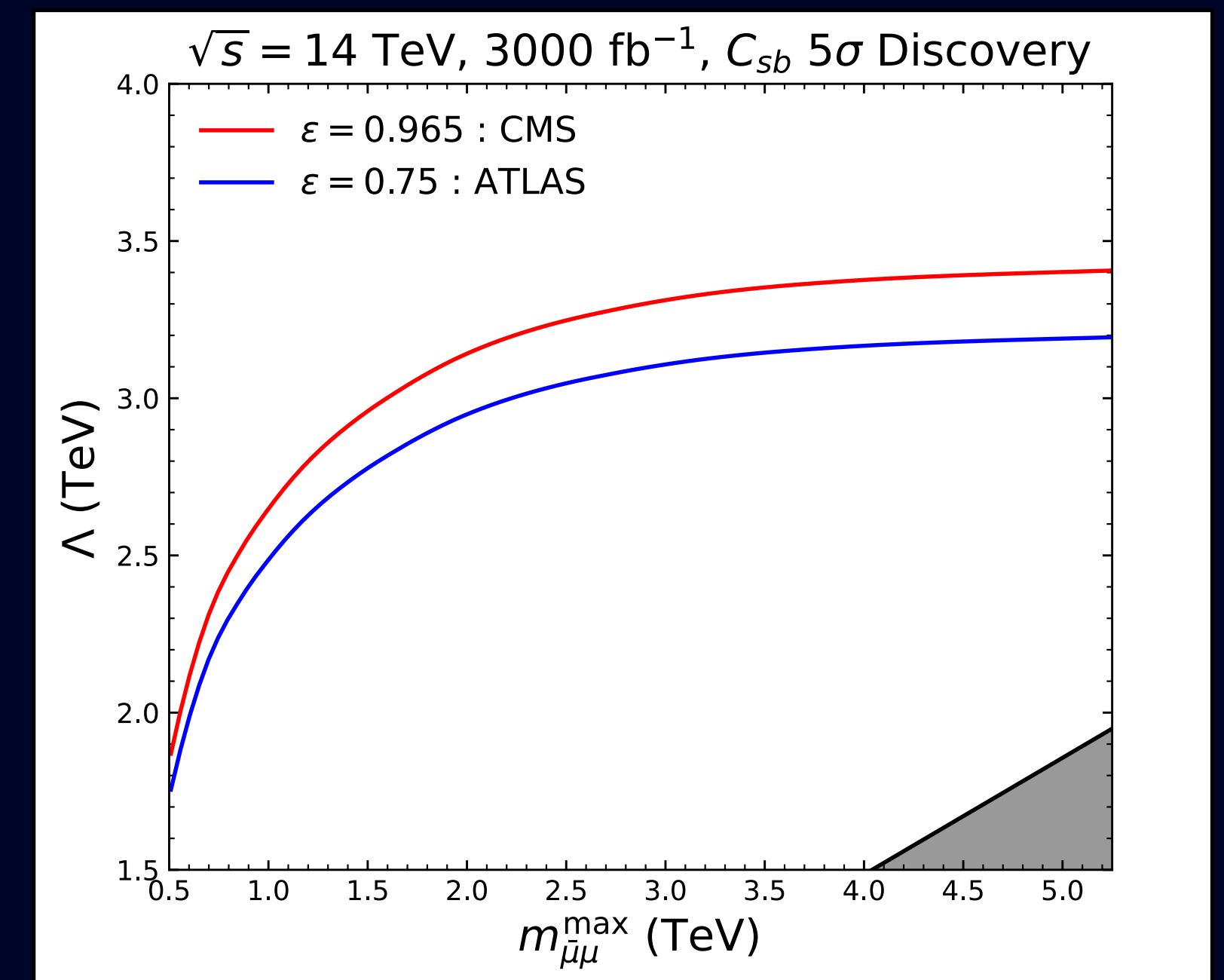
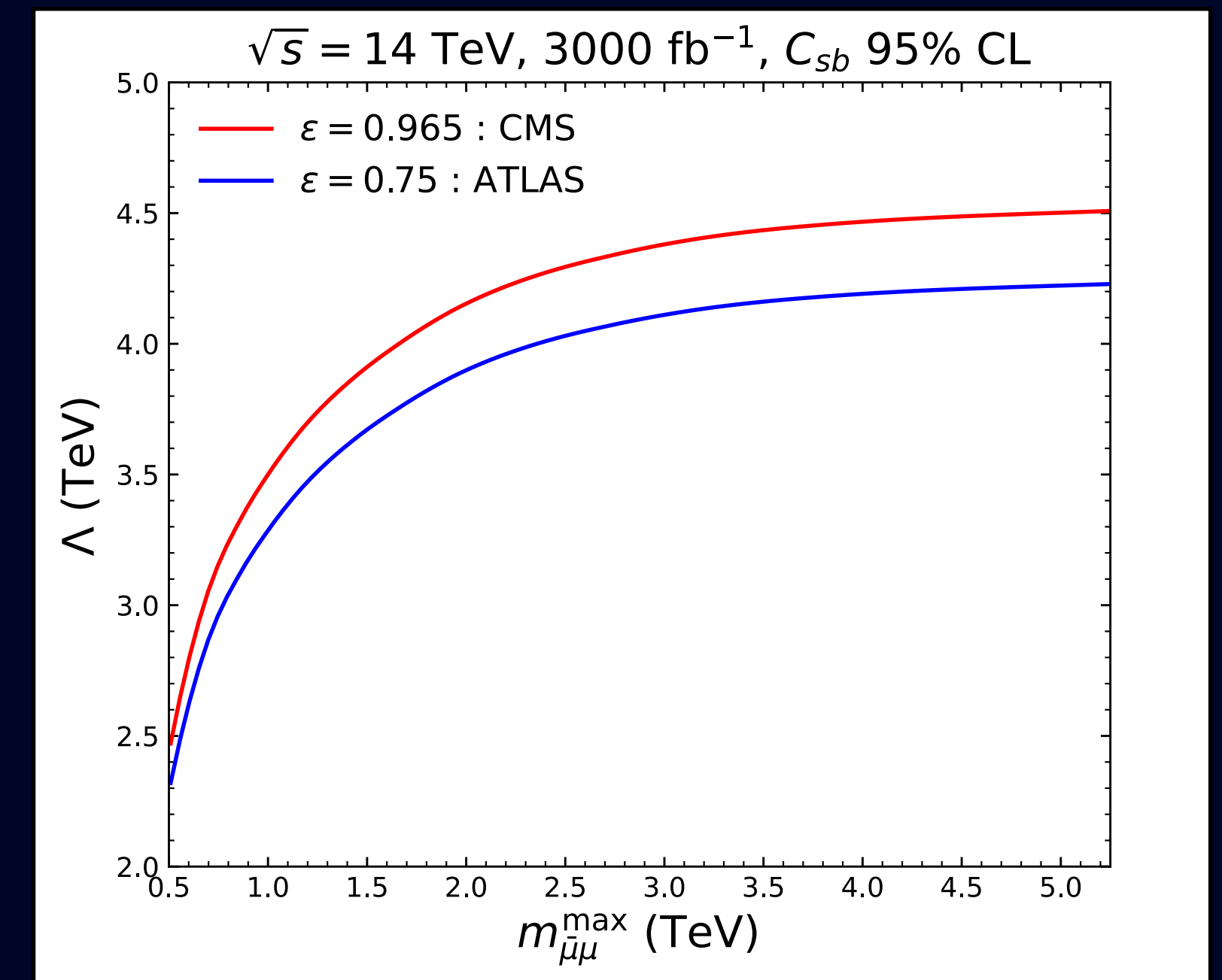
# RESULTS

# LHC & HL-LHC SENSITIVITY

- Collider recast based on JHEP10 (2017) 182 using  $36 \text{ fb}^{-1}$  ATLAS search and in agreement with literature arXiv:1704.09015 and including NLO

$\sqrt{s}$ (TeV)	95% Exclusion				$5\sigma$ Discovery			
	13		14		13		14	
$L$ ( $\text{fb}^{-1}$ )	36	139	3000	3000	36	139	3000	3000
$\Lambda$ (TeV) ( $\epsilon = 0.75$ )	2.3	2.7	4.1	4.2	1.7	2.1	3.1	3.2
$\Lambda$ (TeV) ( $\epsilon = 0.965$ )	2.4	2.9	4.3	4.5	1.8	2.2	3.2	3.4

- $\epsilon$  - ATLAS/CMS muon identification efficiency
- $m_{\bar{\mu}\mu}^{\text{max}} = \{4.5, 5.25\}$  TeV
- Exclude and discover areas of parameter space for a given NP model without the need for a direct search

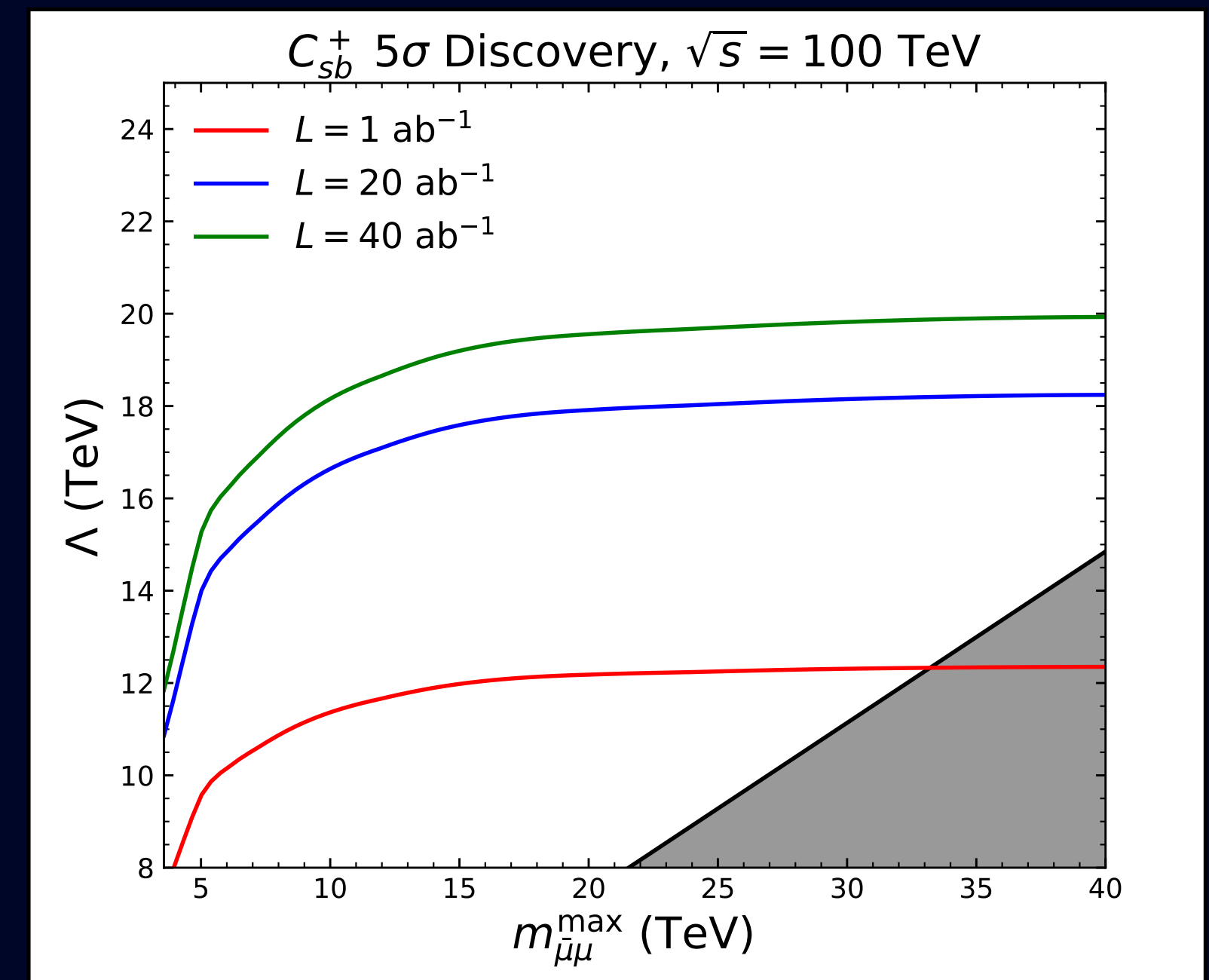
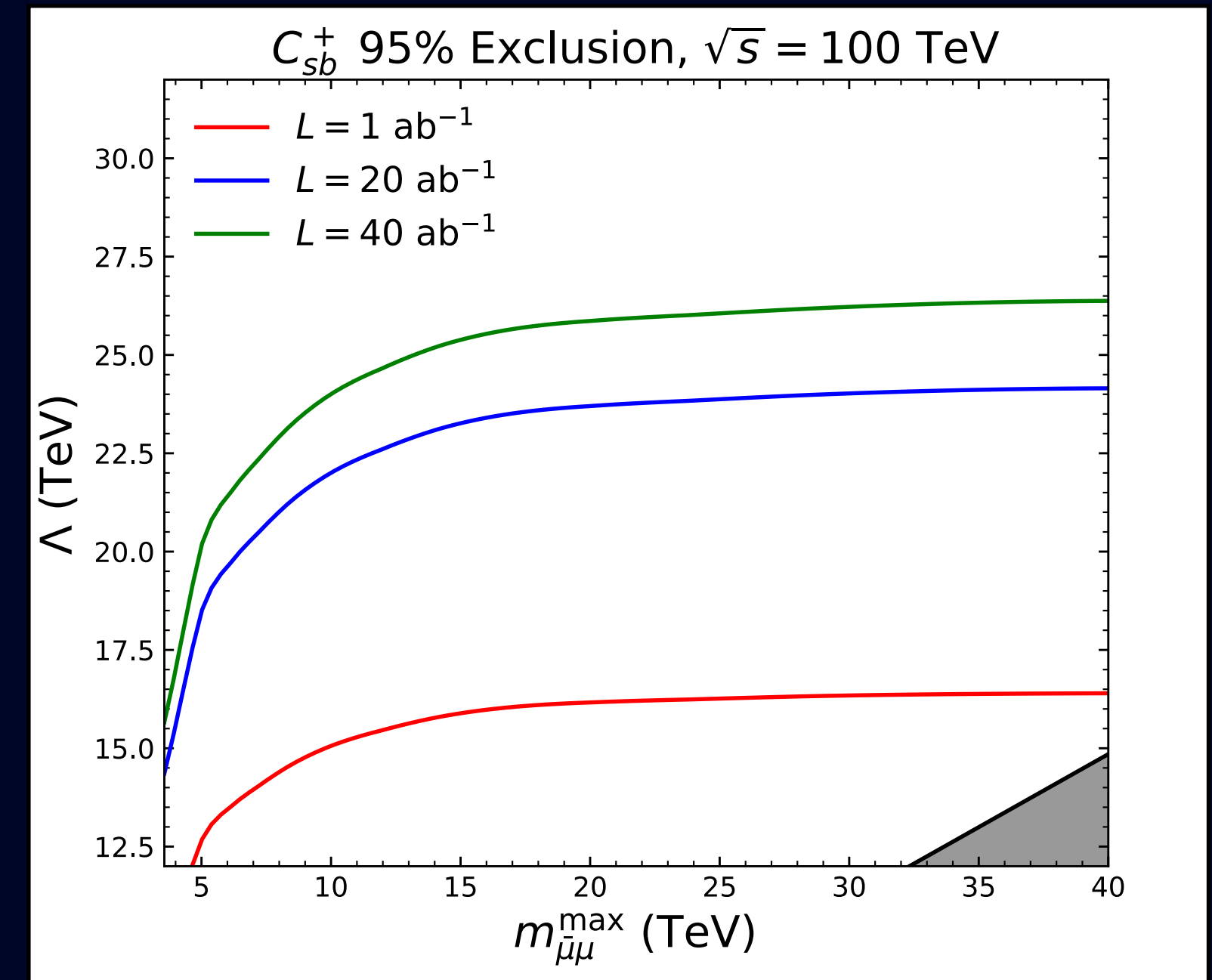


# SENSITIVITY AT FCC-HH

- $m_{\bar{\mu}\mu}^{\max} = \{15, 40, 30\}$  TeV

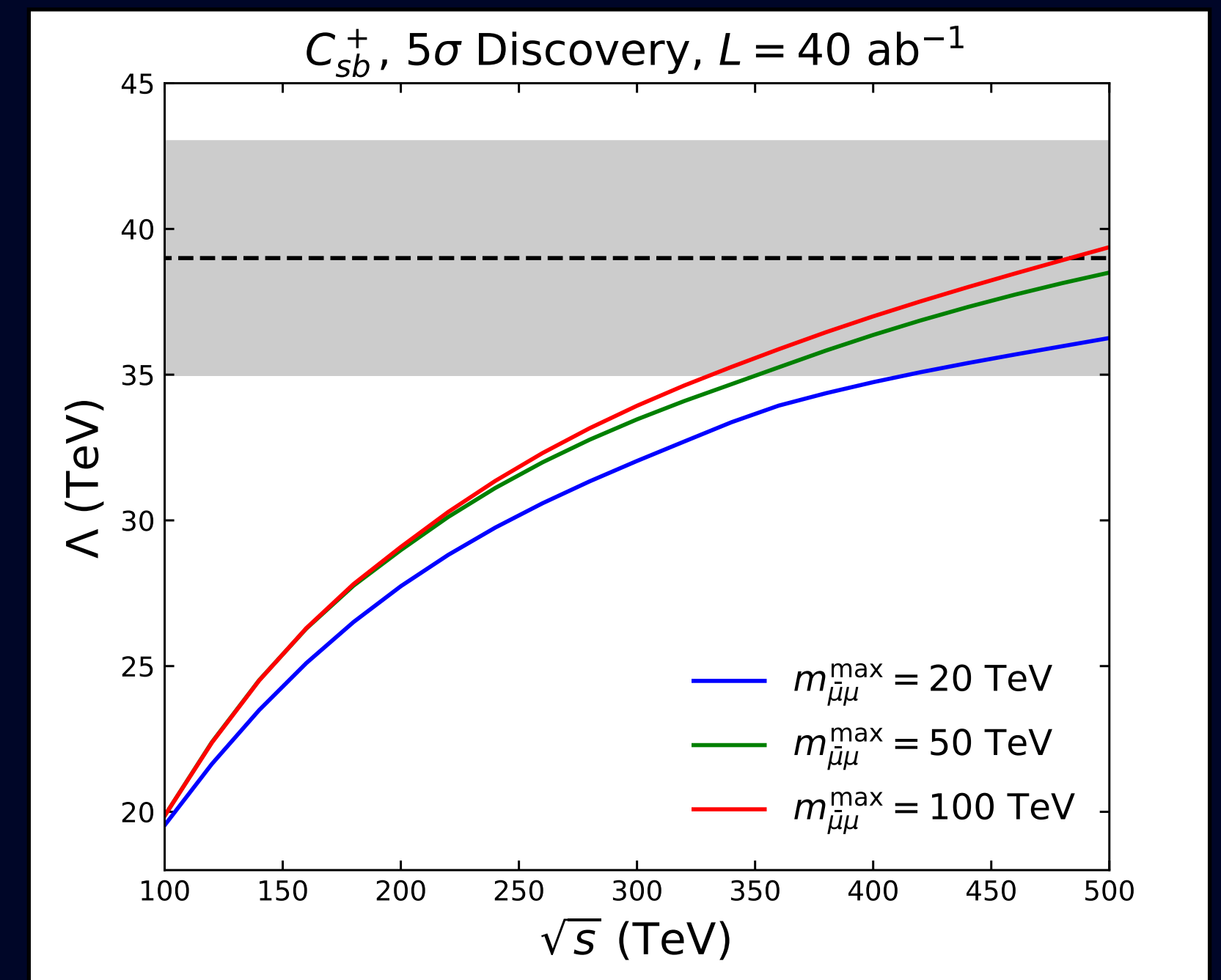
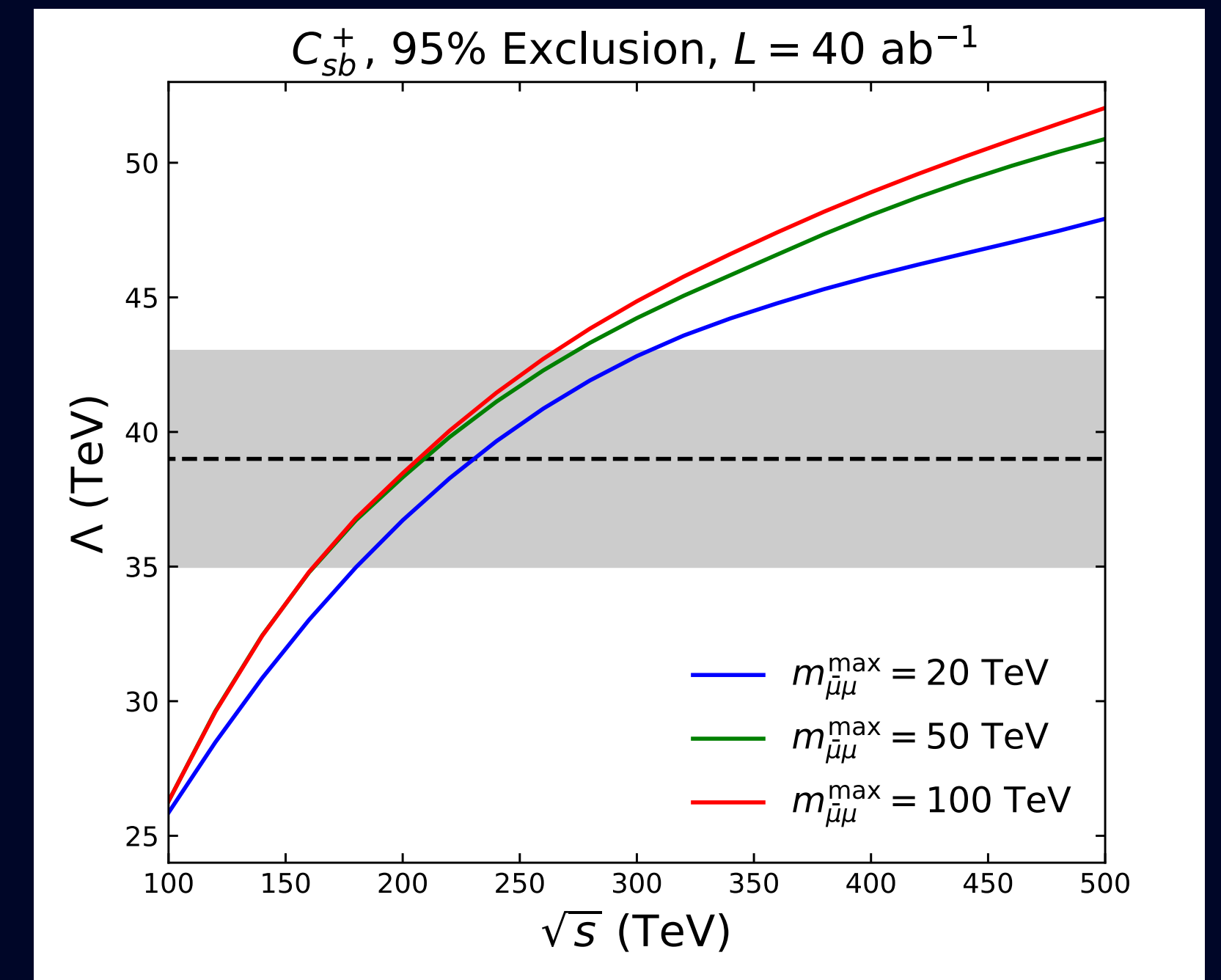
	95% Exclusion			5 $\sigma$ Discovery			
	$L$ (ab $^{-1}$ )	1	20	40	1	20	40
$\Lambda$ (TeV)		15.8	24.1	26.4	12.0	18.1	19.8

- Unitary bound is more relevant at low luminosities
- Improved sensitivity at FCC-hh by  $\sim 5$  times
- Need  $\sim 2$  times more luminosity to exclude  $\Lambda \approx 30$  TeV
- Need  $\sim 20$  times more luminosity to exclude  $\Lambda \approx 40$  TeV

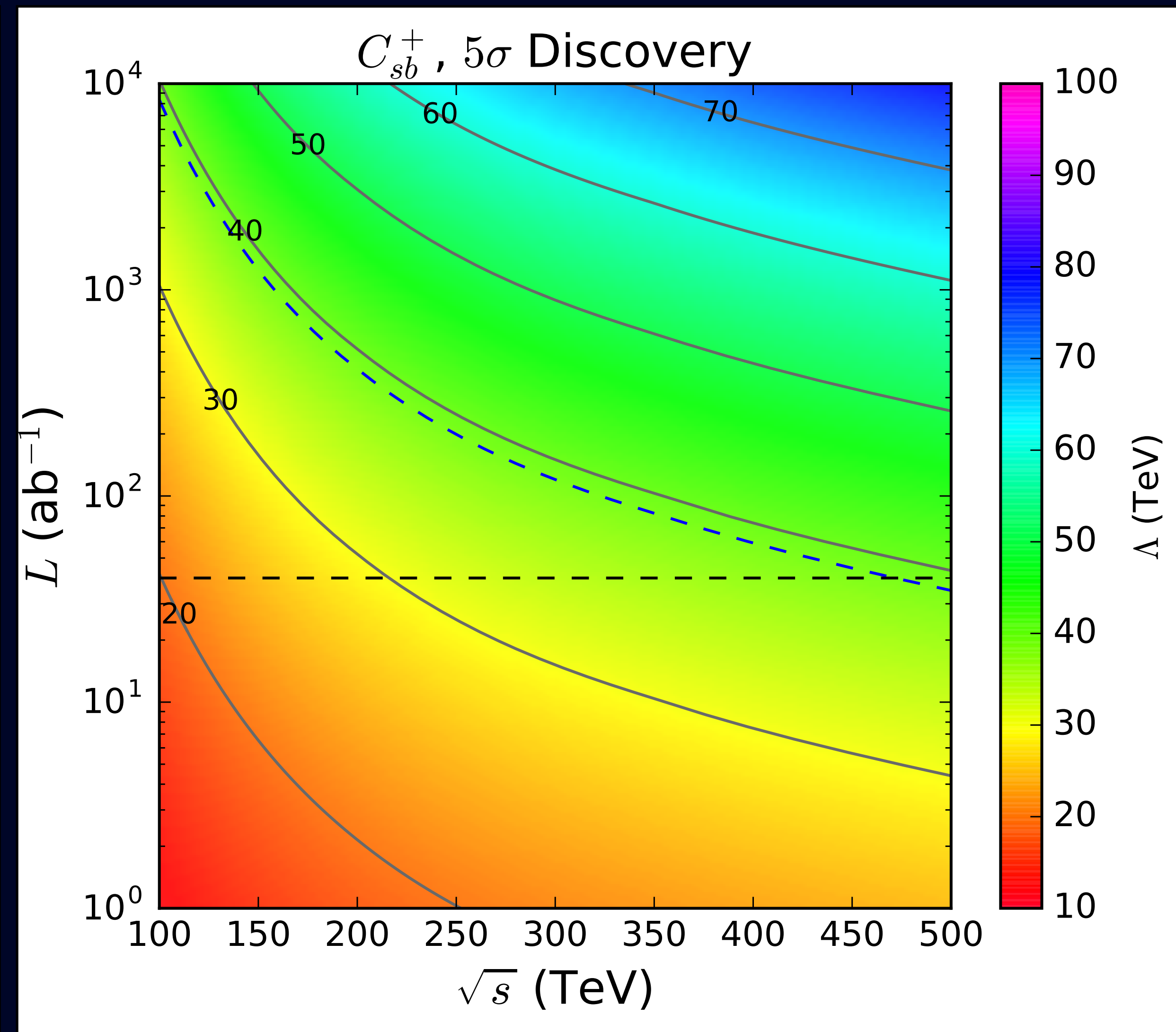
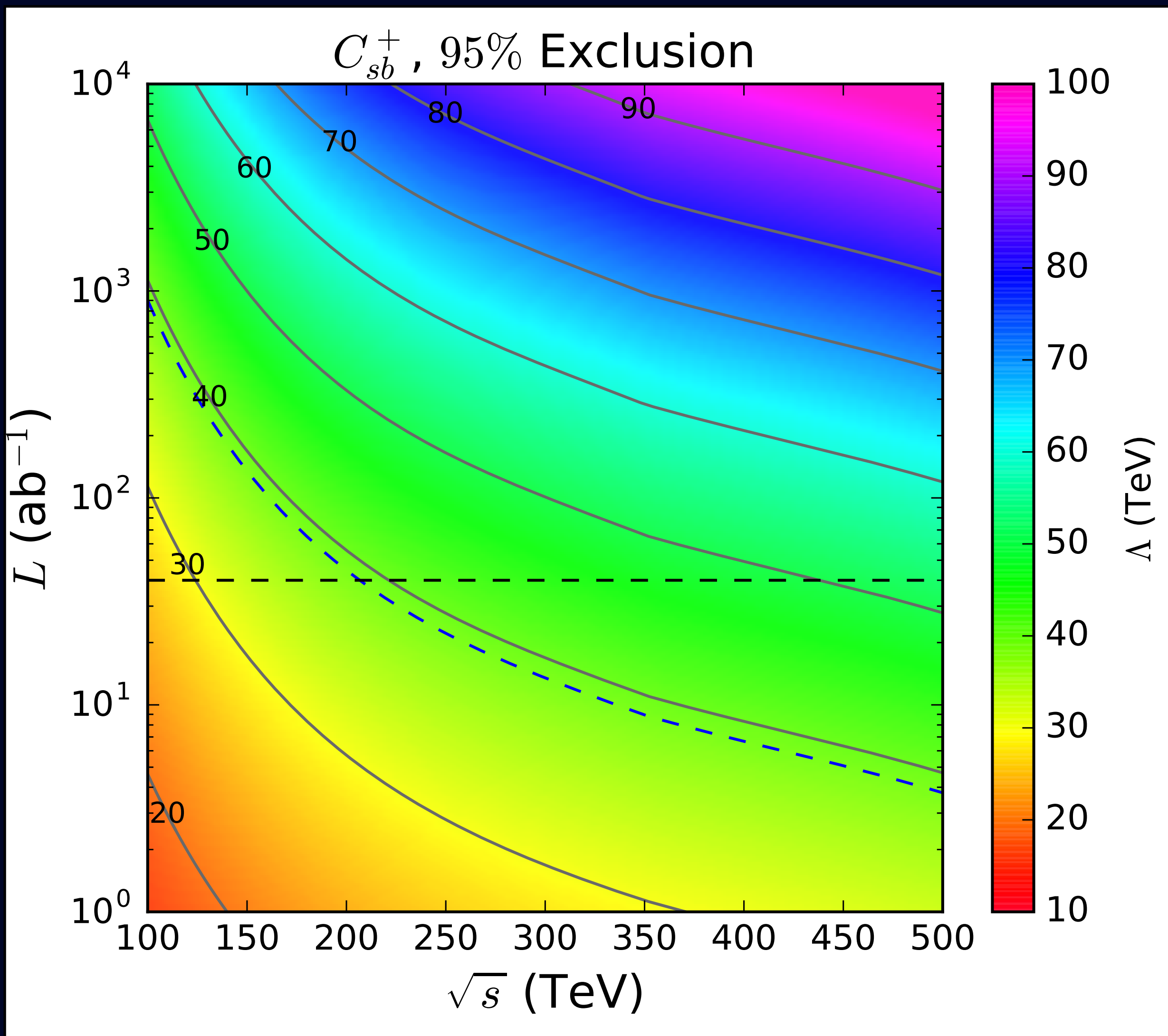


# BEYOND FCC-HH

- Consider FCC-hh with increased c.o.m energy, i. e.  $\sqrt{s} > 100$  TeV and  $L = 40 \text{ ab}^{-1}$
- As c.o.m energy increases the DY process and the EFT signal is suppressed by increasingly larger negative Sudakov Double Logarithms
- Diboson process dominates the background around  $\sqrt{s} = 300$  TeV
- $m_{\bar{\mu}\mu}^{\min} = 3$  TeV gives the optimised sensitivity in the region of  $\sqrt{s} = [100, 500]$  TeV



# BEYOND FCC-HH



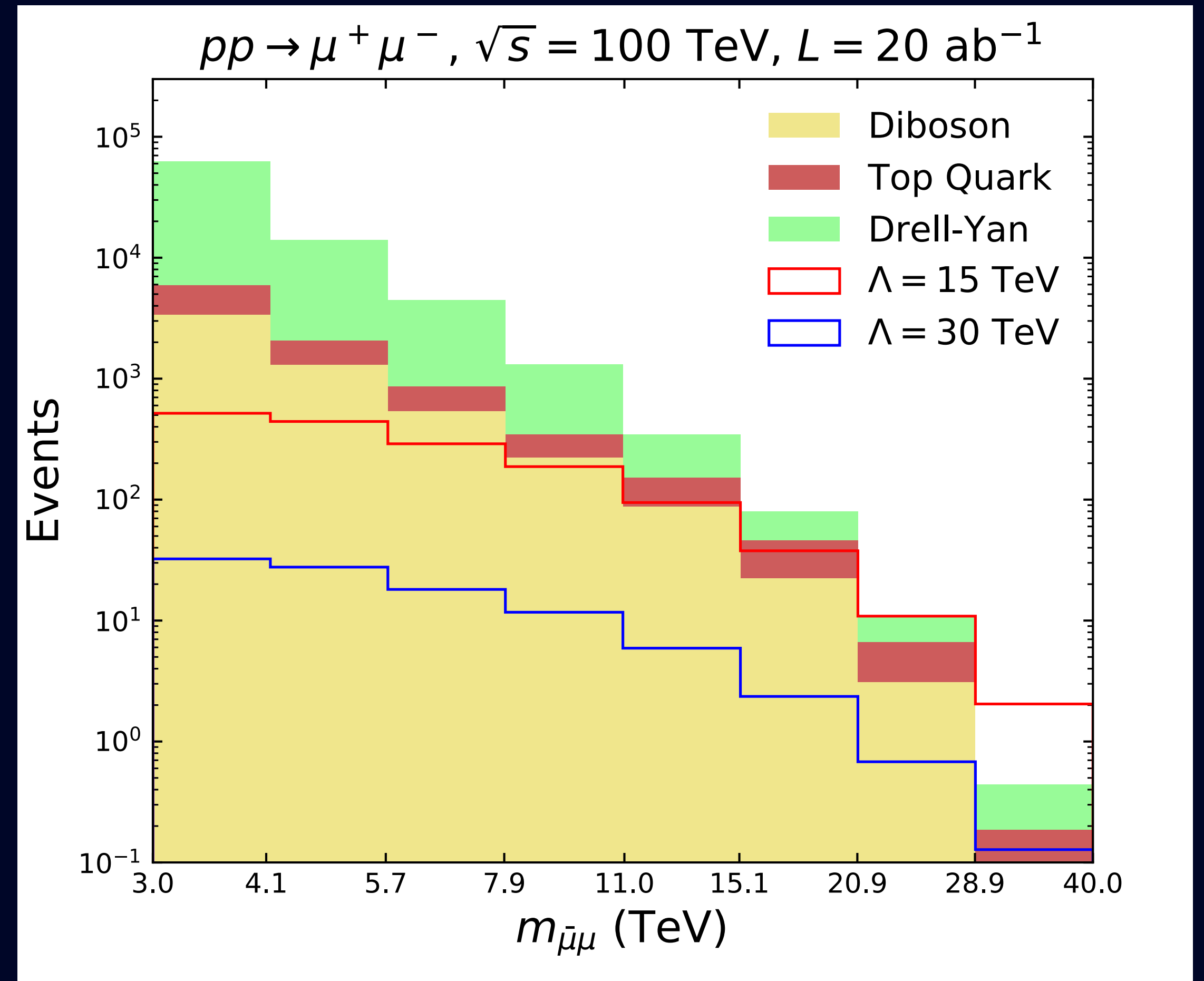
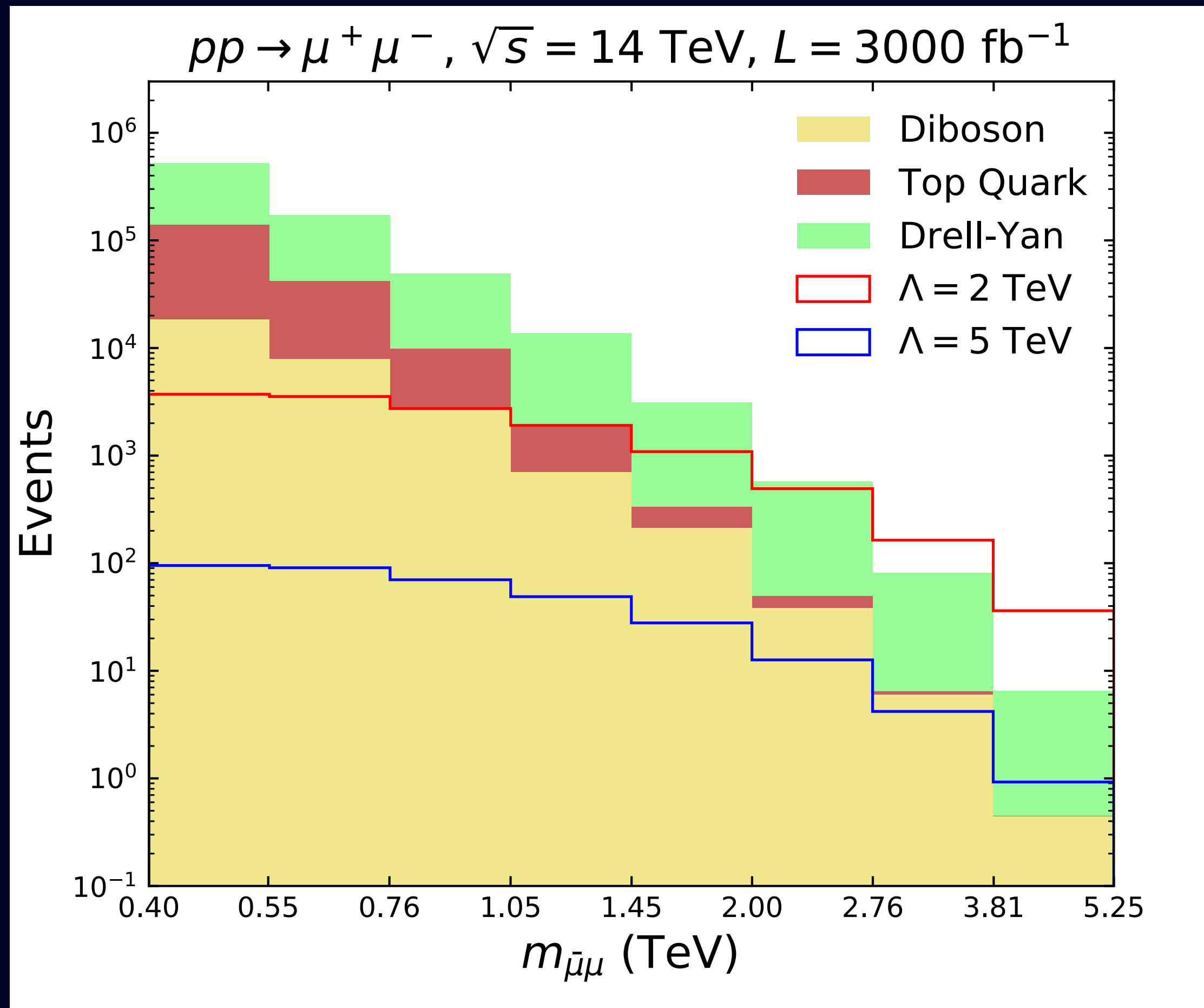


# CONCLUSION

- Model-independent and most conservative analysis
- Probing a value of  $\Lambda \approx 39$  TeV as suggested by the B-anomalies would be possible with a machine with higher energy and/or luminosity than the FCC-hh
- A comprehensive case would combine contact interaction searches (for heavy mediators) with direct searches (for lighter mediators)
- Could improve sensitivity by considering a more complex final state with a reduced SM background, eg. exclusive  $\mu^+\mu^-b$  search JHEP08 (2018) 056
- Muon collider offers encouraging results PhysRevD.105.015013

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