

# Dark Solar Wind

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# BSM particles from the Sun



~ keV

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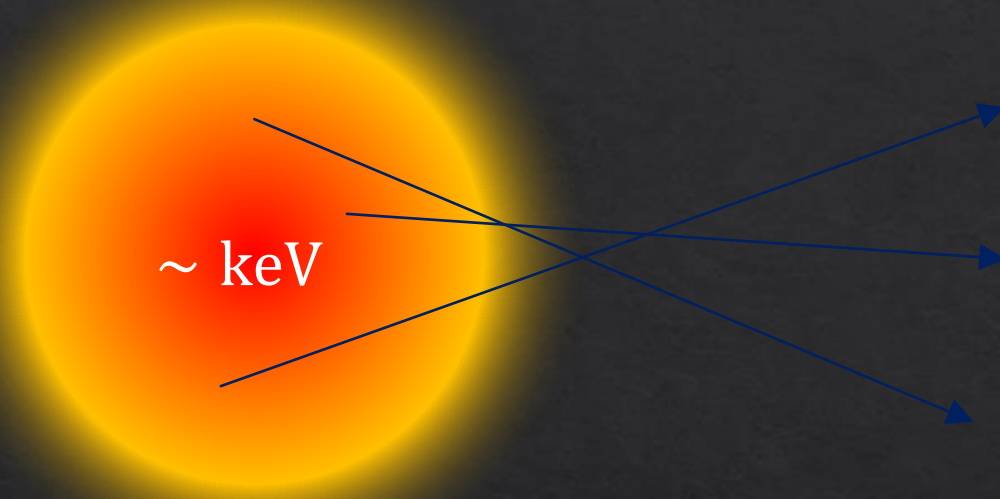
Maximum energy flux on Earth limited by **red giant energy loss**

$$\frac{10^{-2} L_{\odot}}{4\pi \text{ AU}^2} \sim 10^{-4} \frac{\text{GeV}}{\text{cm}^3}$$

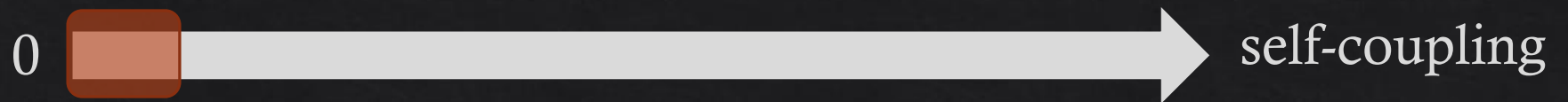
**Comparable** to local DM energy flux  $\rho_{\text{DM}} v_{\text{DM}}$

Probes **different theory space** and **more directly**

# BSM particles from the Sun



Free streaming



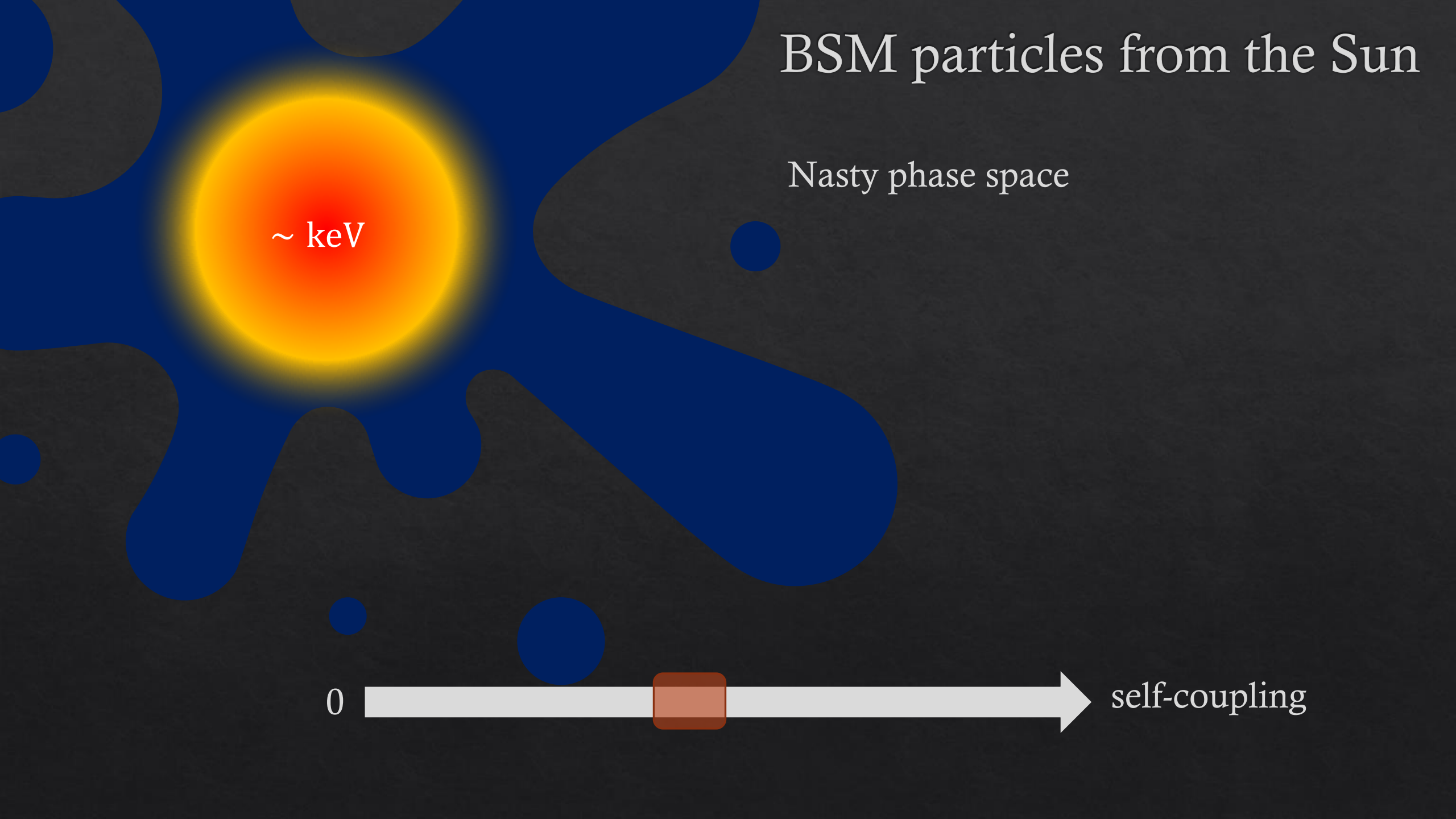
# BSM particles from the Sun

Nasty phase space

$\sim \text{keV}$

0

self-coupling



# BSM particles from the Sun



Self-thermalization

Admits **model-independent** description

**Predictive**

Different and interesting outcomes



# Two outcomes from thermalization

## 1. Denser and Less Energetic

Before thermalization  $n \ll E^3$ . After thermalization  $n \sim T^3 \sim E^3$  (naively)

## 2. Fluid Dynamics

Mean free path  $\sim (\text{coupling})^\# T^{-1} = \text{microscopic}$

Thermal pressure leads to **relativistic bulk velocities**,  $\gamma(1 \text{ AU}) \approx 900$  (for light MCPs)

Boosted thermal distribution:  $n \sim \gamma \tilde{T}^3$ ,  $E \sim \gamma \tilde{T}$

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Overall:  $n$  increased by *at least*  $10^3$   
 $E$  decreased by *at least*  $10^3$  compared to the free-streaming case

# Fluid Dynamics

$$\frac{1}{r^2} \partial_r [r^2 \gamma^2 v (\tilde{\rho} + \tilde{p})] = \dot{Q}$$

$$\frac{1}{r^2} \partial_r [r^2 \gamma^2 v^2 (\tilde{\rho} + \tilde{p})] = -\partial_r \tilde{p}$$

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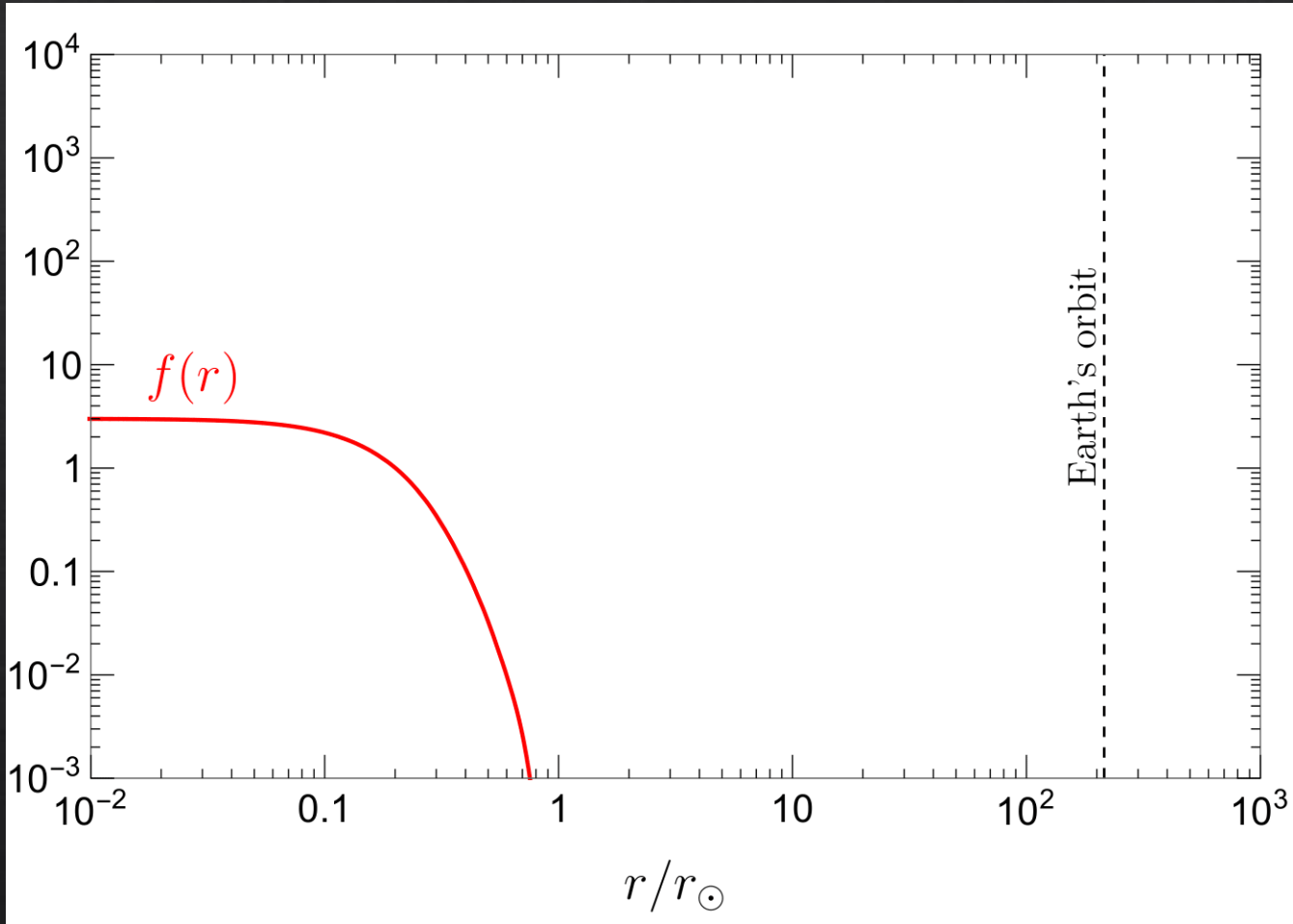
$$p = \rho/3$$



Energy Flux

$$\# \gamma^2 v \tilde{T}^4 = \frac{L_\chi}{4\pi r^2}$$

# Fluid Dynamics



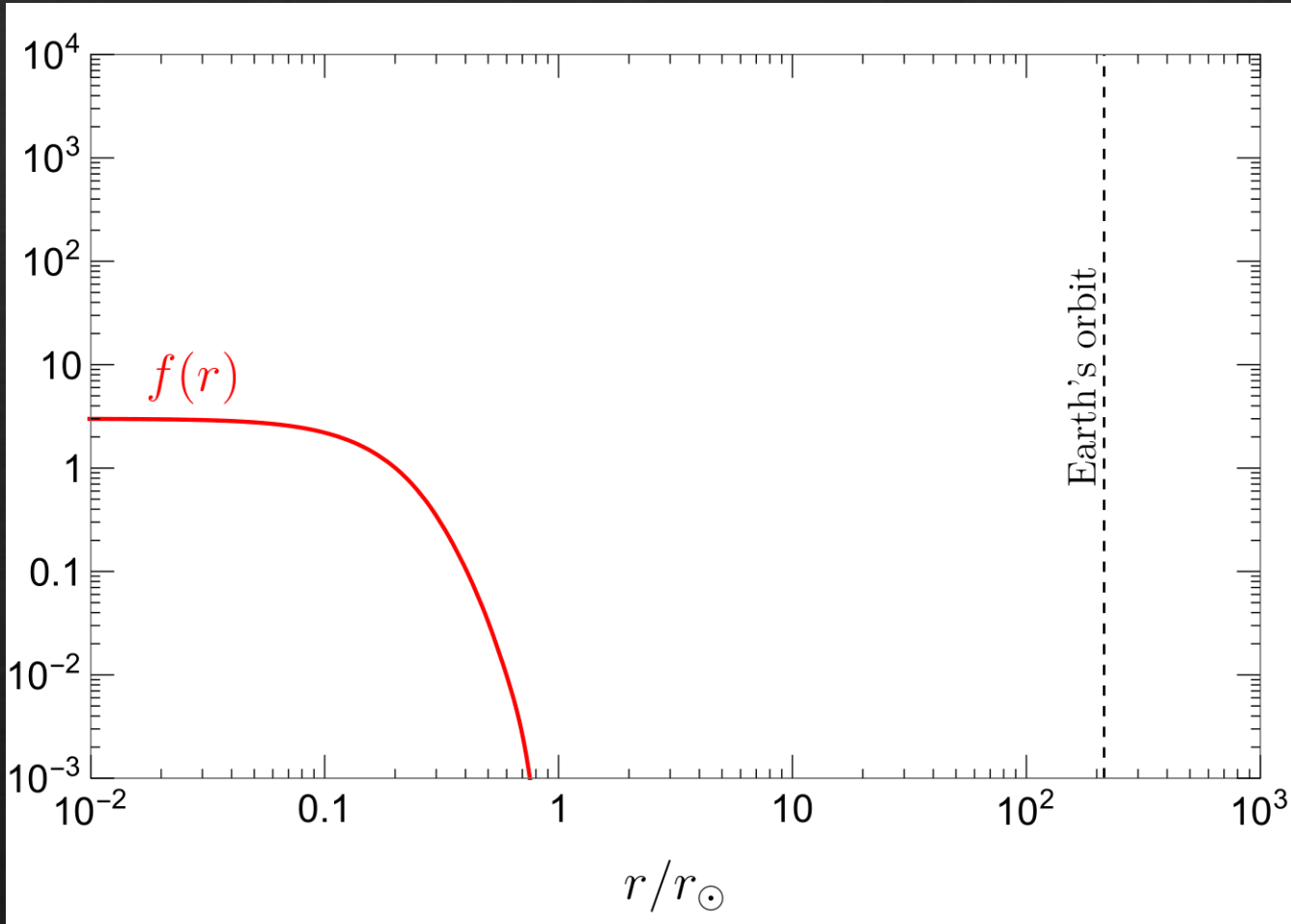
## Energy Flux

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## Bulk Velocity

$$\frac{d \ln v}{d \ln r} = \left( \frac{1/3 + v^2}{1/3 - v^2} \right) \left[ f(r) - \frac{\int_0^r \dot{Q} r'^2 dr'}{1 + 3v^2} \right]$$

# Fluid Dynamics



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$$\frac{d \ln v}{d \ln r} = \left( \frac{1/3 + v^2}{1/3 - v^2} \right) \left[ f(r) - \frac{2(1 - v^2)}{1 + 3v^2} \right]$$

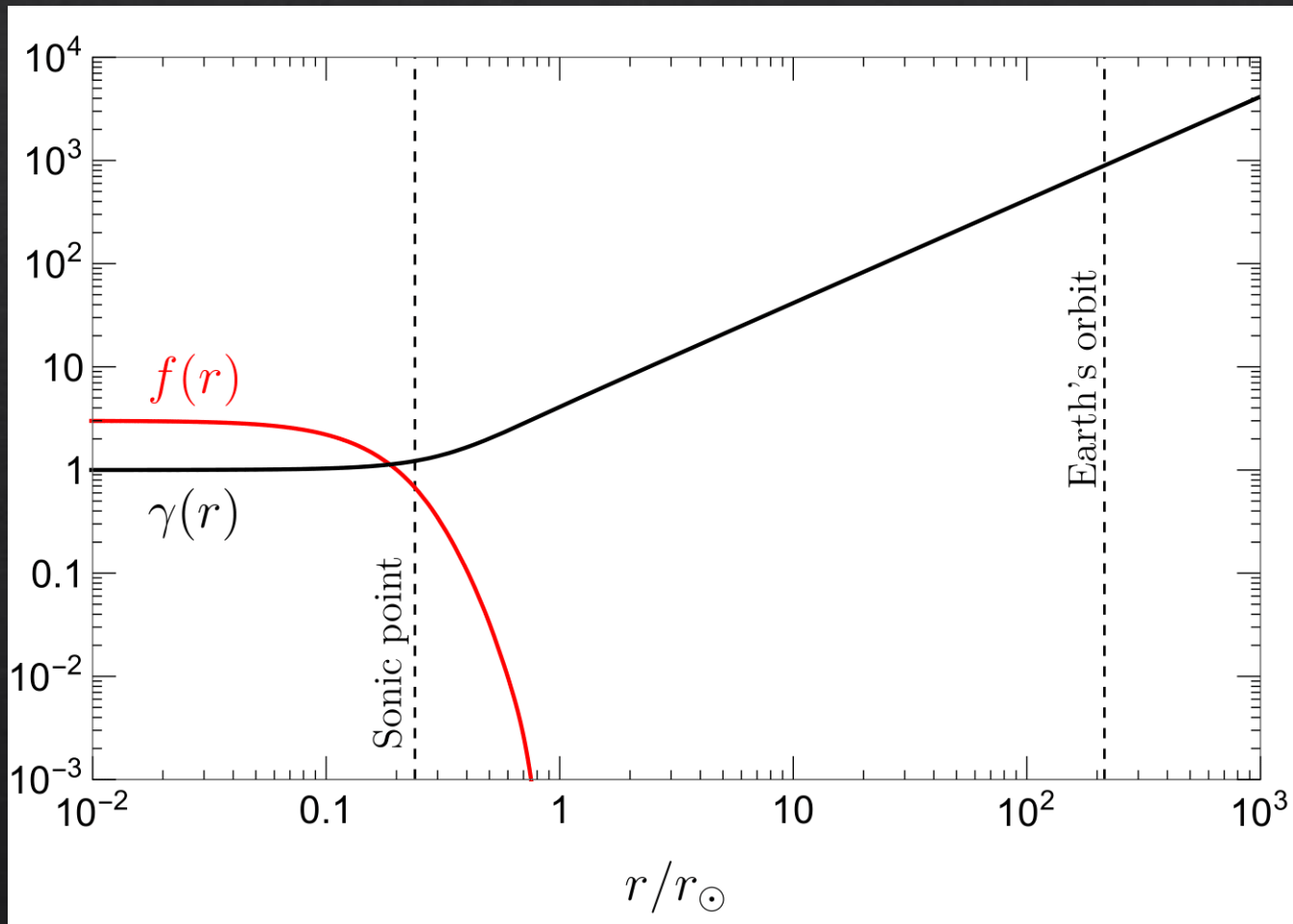
$\int_0^r \dot{Q} r'^2 dr'$

## Boundary Conditions

$$v = 0 \text{ at } r = 0$$

$$\tilde{T} = 0 \text{ at } r \rightarrow \infty$$

# Fluid Dynamics



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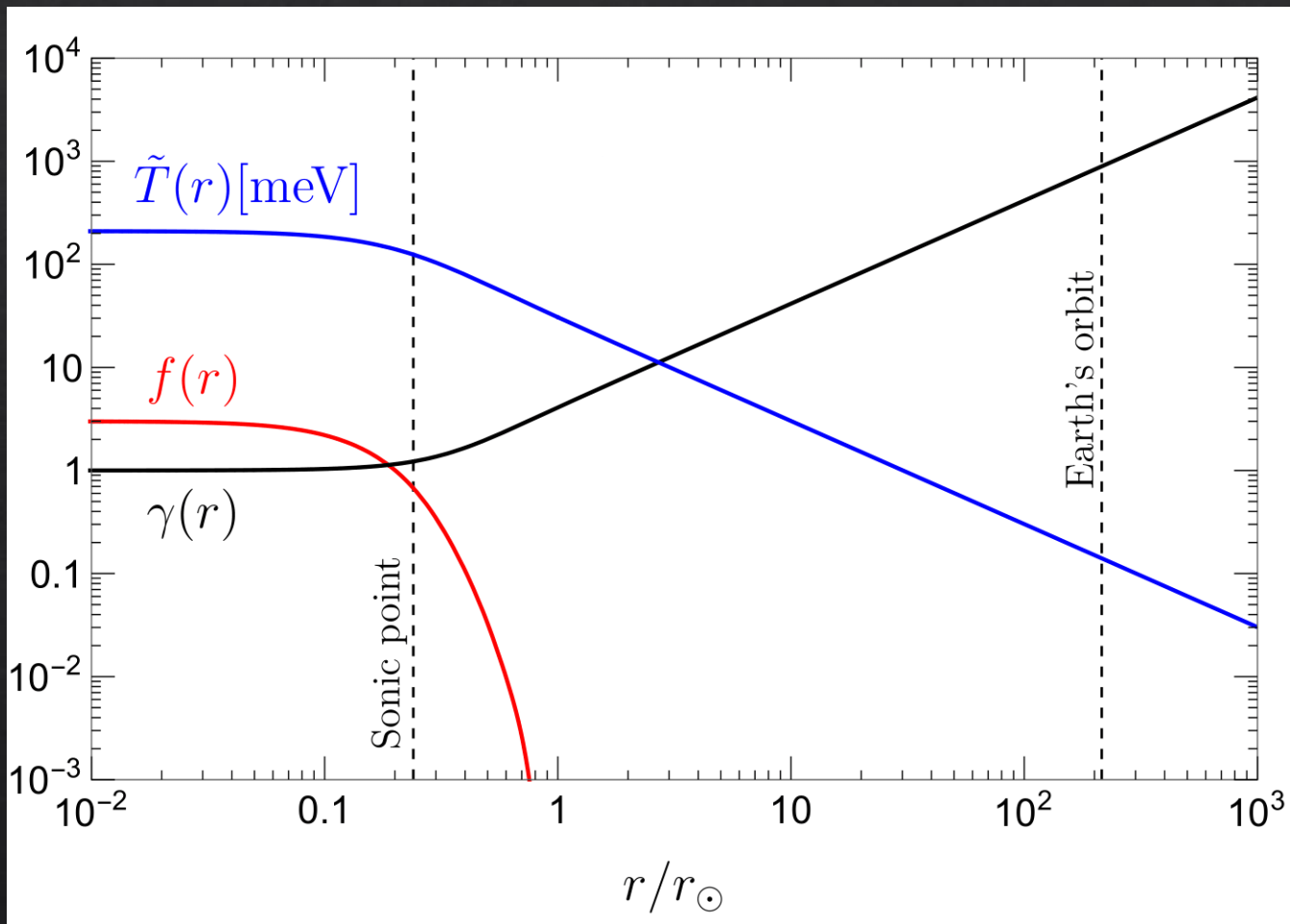
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## Transonic Solution

analogous to Parker's solar wind, asymptotes to "fireball" solution  $\gamma \propto r/r_{\text{sonic}}$  at  $r \gg r_{\text{sonic}} = 0.2r_{\odot}$

# Fluid Dynamics



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$\frac{\int_0^r \dot{Q} r'^2 dr'}{\dot{Q} r^3}$

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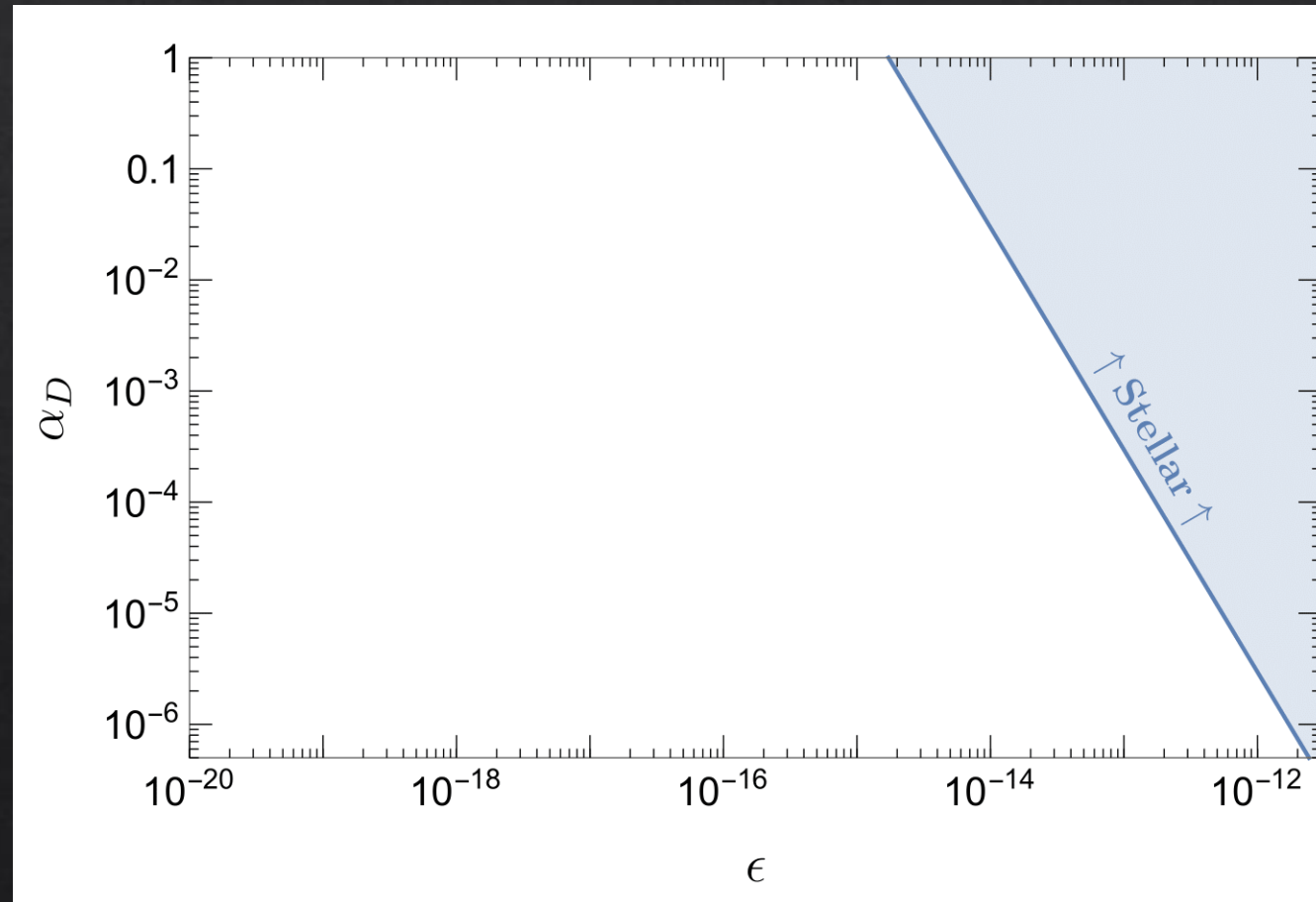
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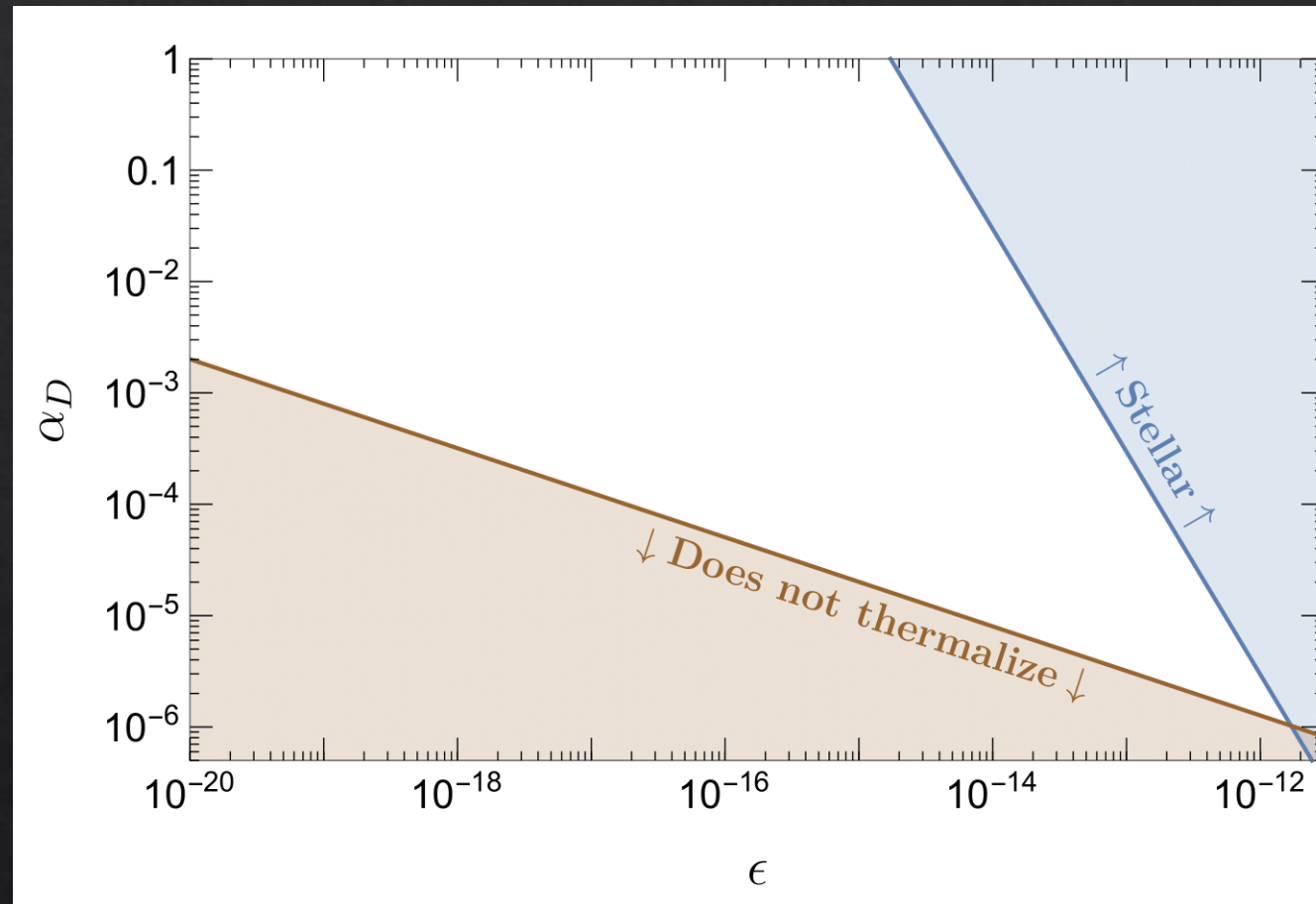
# Millicharged Particles

dark fine-structure constant  $\alpha_D$ , kinetic mixing  $\epsilon$ , negligible masses



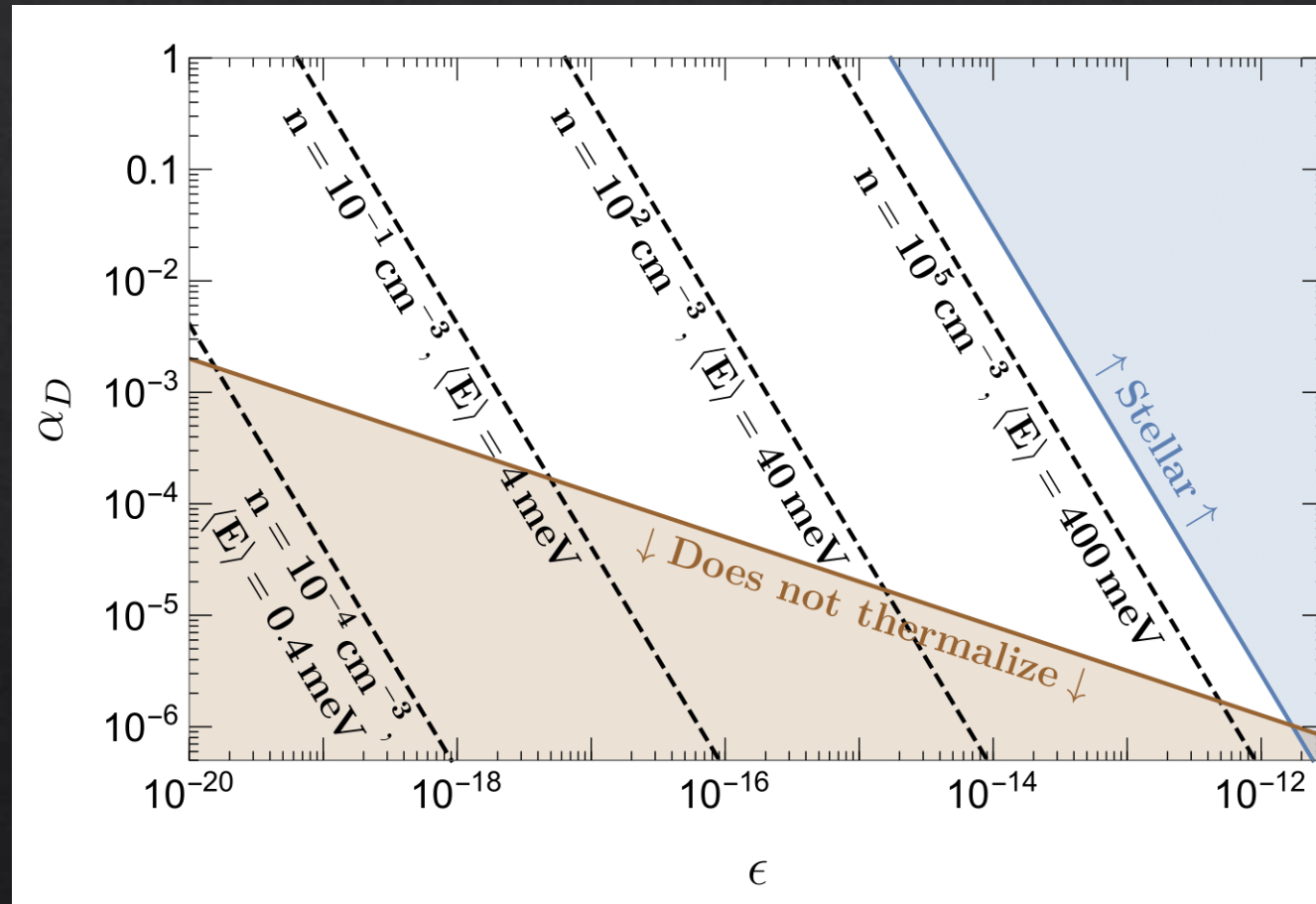
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# Summary

- ◇ The Sun is a good source of light BSM particles
- ◇ Self-thermalization of these particles leads to very different outcomes compared to the ordinary free-streaming scenario:
  - ◇  $\gtrsim 10^3$  denser,
  - ◇  $\gtrsim 10^3$  less energetic,
  - ◇ relativistic fluid outflow

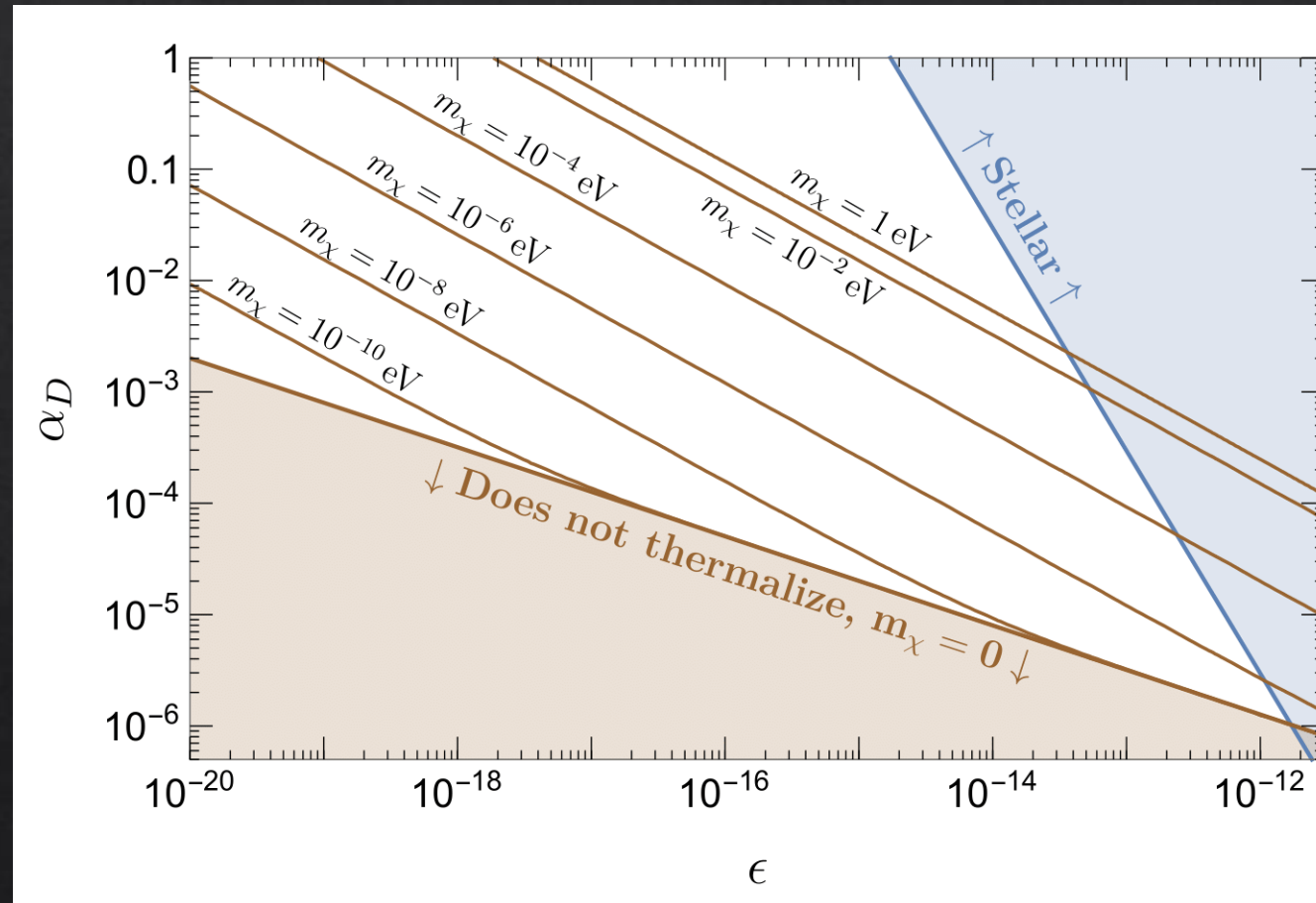
Dark Solar Wind!

# Future Directions

- ◇ Detection strategies
- ◇ Neighboring parameter space
- ◇ Other models
- ◇ Other astrophysical objects
- ◇ Astrophysical consequences

Thank You

# Massive Millicharged Particles

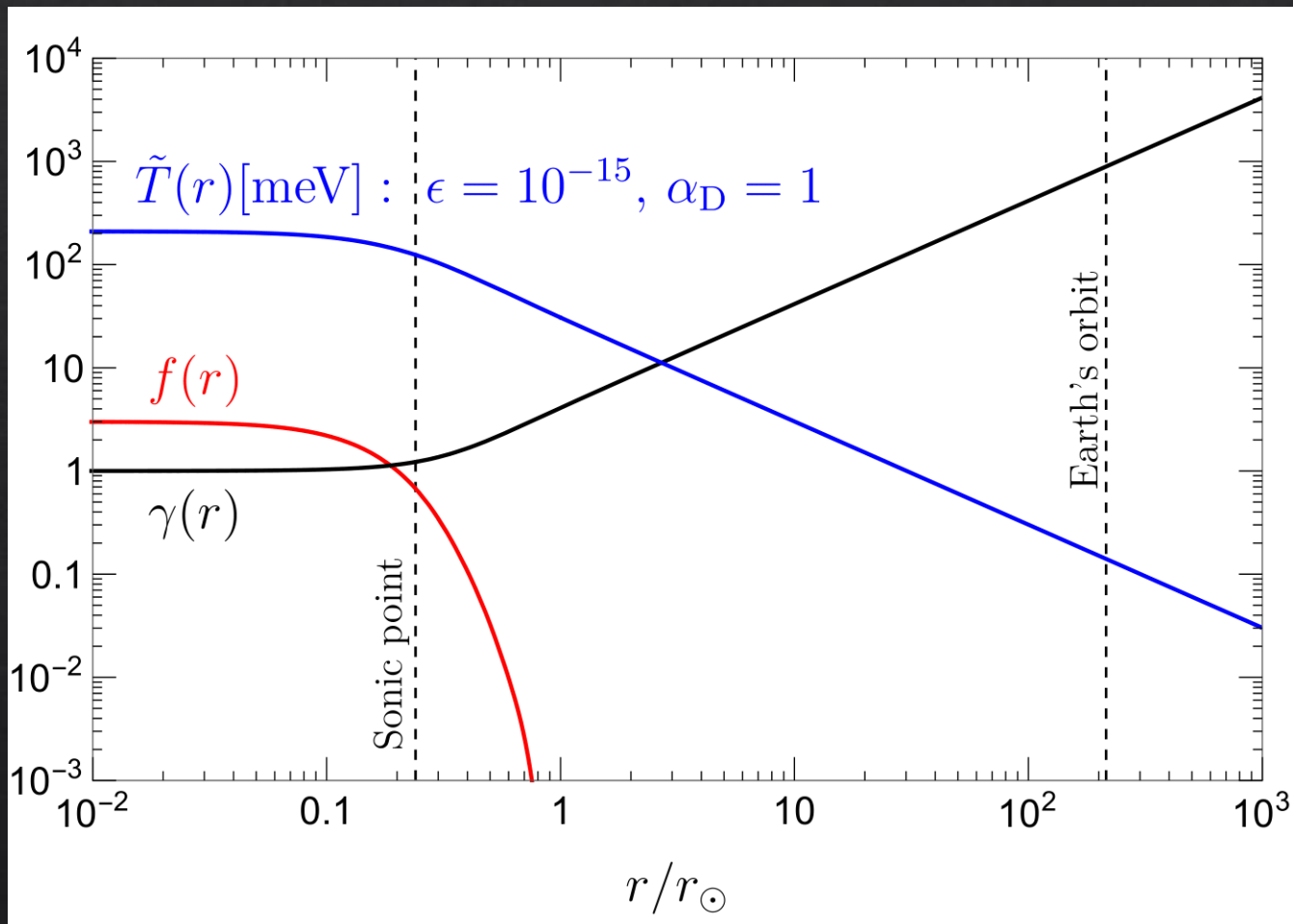


# Millicharged Particle Predictions

- ◆ Bulk velocity:  $\gamma(r) \approx 893 \left( \frac{r}{1 \text{ AU}} \right)$
- ◆ Comoving temperature:  $\tilde{T} \approx 0.1 \text{ meV} \left( \frac{\epsilon}{10^{-15}} \right)^{\frac{1}{2}} \left( \frac{\alpha_D}{1} \right)^{\frac{1}{4}} \left( \frac{r}{1 \text{ AU}} \right)^{-1}$
- ◆ Average energy:  $\langle E \rangle \approx 4\gamma\tilde{T} \approx 0.5 \text{ eV} \left( \frac{\epsilon}{10^{-15}} \right)^{\frac{1}{2}} \left( \frac{\alpha_D}{1} \right)^{\frac{1}{4}}$
- ◆ Number density:  $n = \frac{5\zeta(3)}{\pi^2} \gamma\tilde{T}^3 \approx 2 \times 10^5 \text{ cm}^{-3} \left( \frac{\epsilon}{10^{-15}} \right)^{\frac{3}{2}} \left( \frac{\alpha_D}{1} \right)^{\frac{3}{4}} \left( \frac{r}{1 \text{ AU}} \right)^{-2}$



# Sonic Point



$$\# \gamma^2 v \tilde{T}^4 = \frac{L_\chi}{4\pi r^2} \frac{\dot{Q} r^3}{\int_0^r \dot{Q} r'^2 dr'}$$

$$\frac{d \ln v}{d \ln r} = \left( \frac{1/3 + v^2}{1/3 - v^2} \right) \left[ f(r) - \frac{2(1 - v^2)}{1 + 3v^2} \right]$$

Two types of solutions:

- Subsonic:  $\tilde{p} = \text{finite}$  at  $r \rightarrow \infty$
- Transonic:  $\tilde{p} = 0$  at  $r \rightarrow \infty$

Crossing  $c_s = 1/\sqrt{3}$  smoothly requires:  
 $[\dots]_{v=c_s} = 0 \rightarrow r_{\text{sonic}} \approx 0.2 r_\odot$