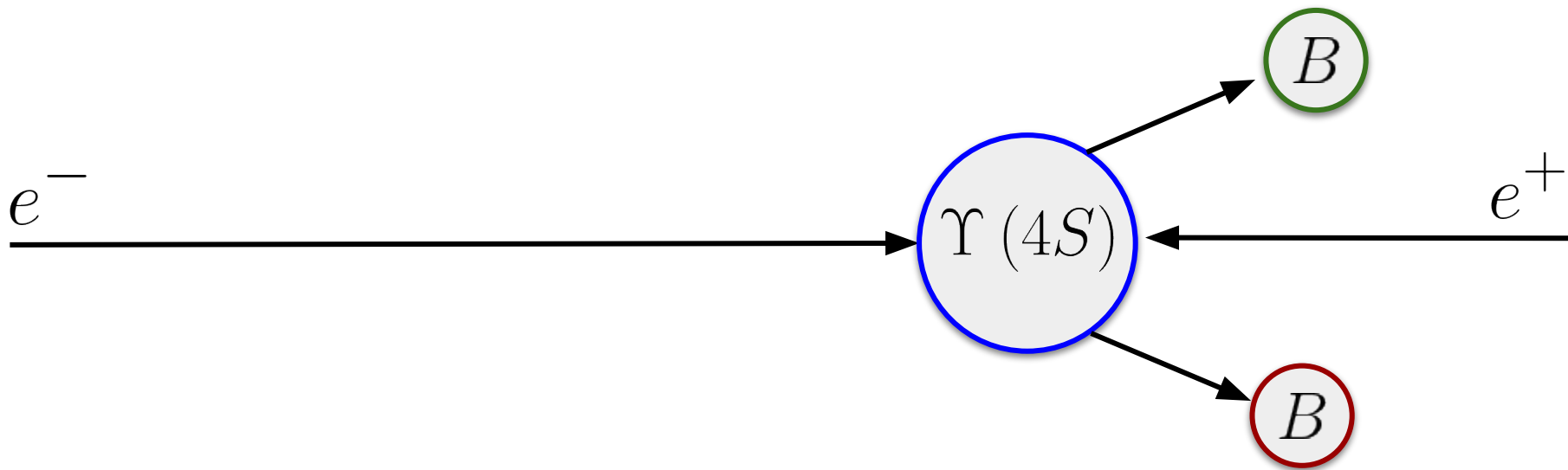
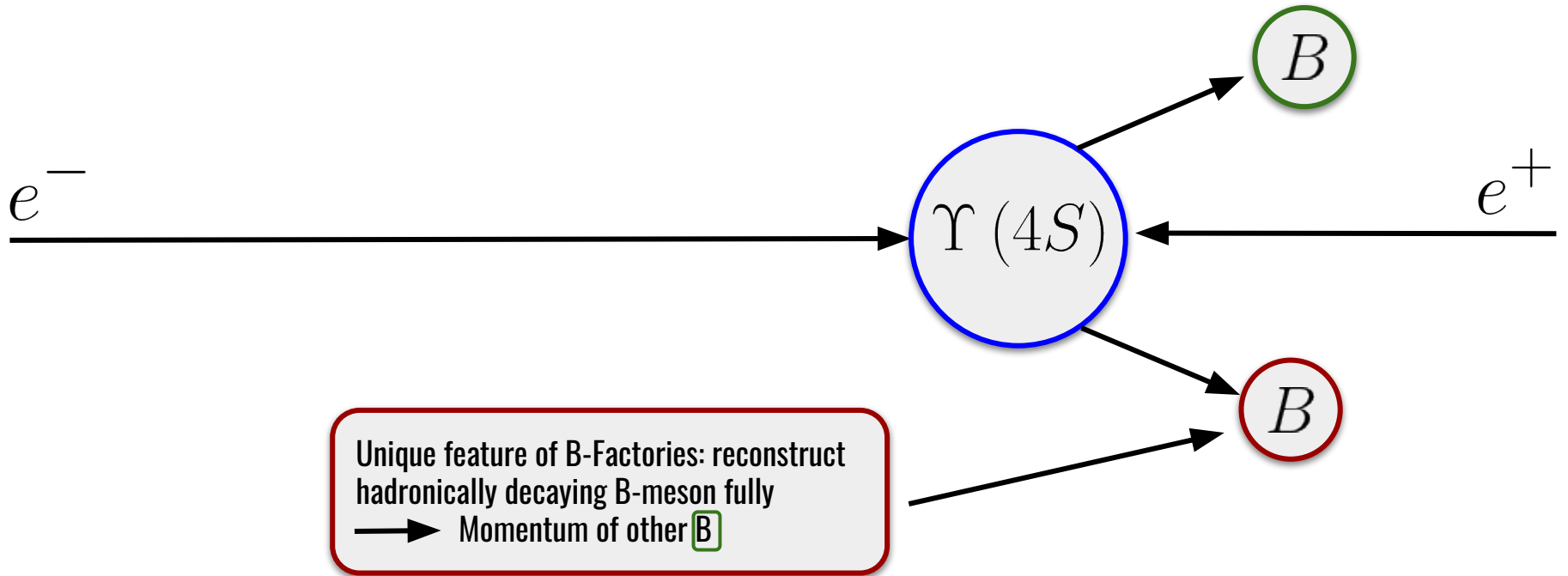


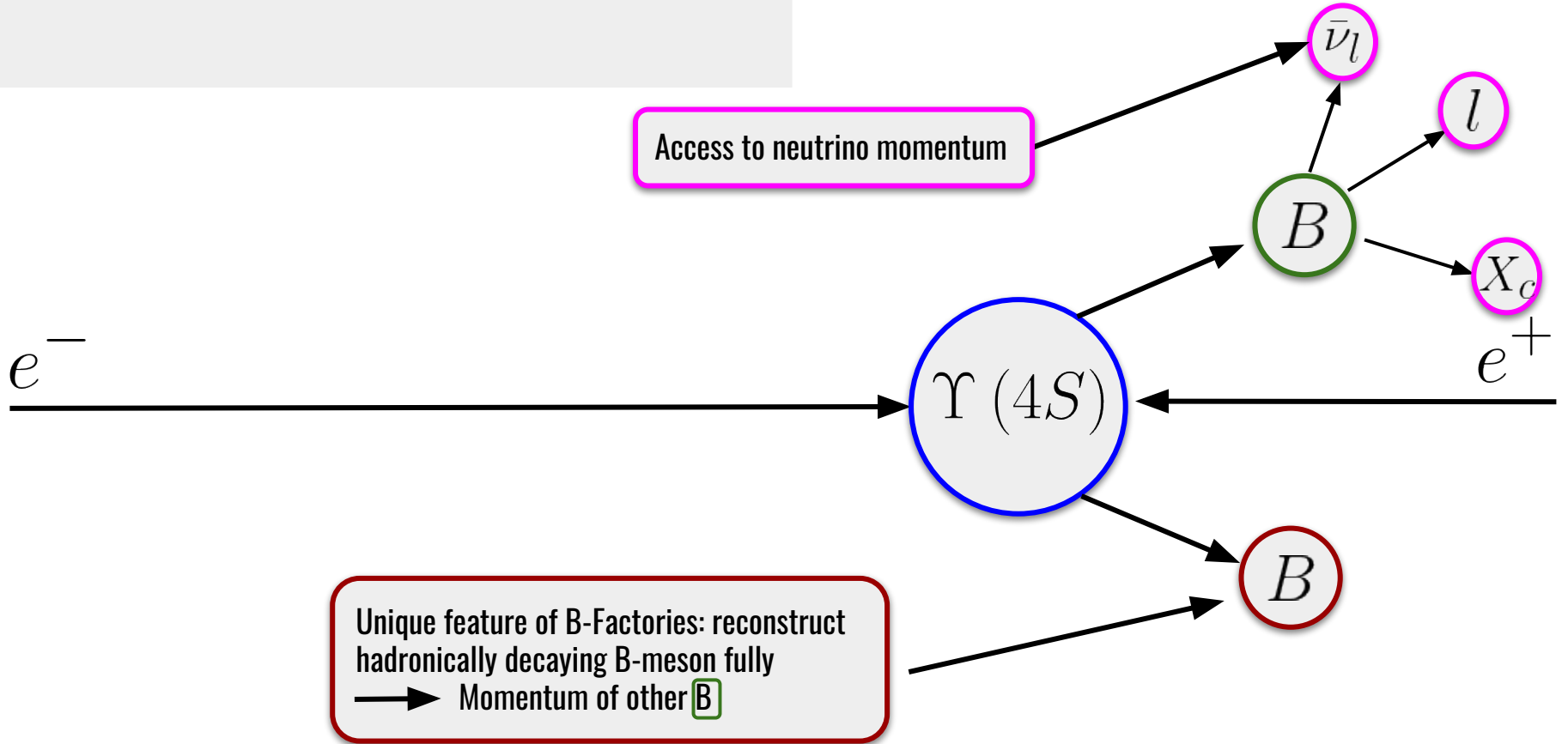
B-Factories



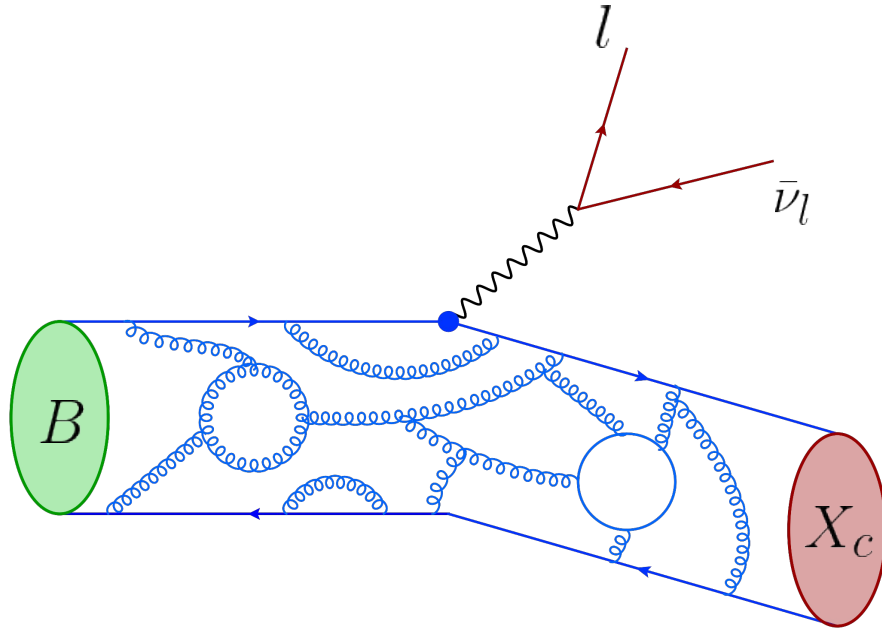
B-Factories



B-Factories

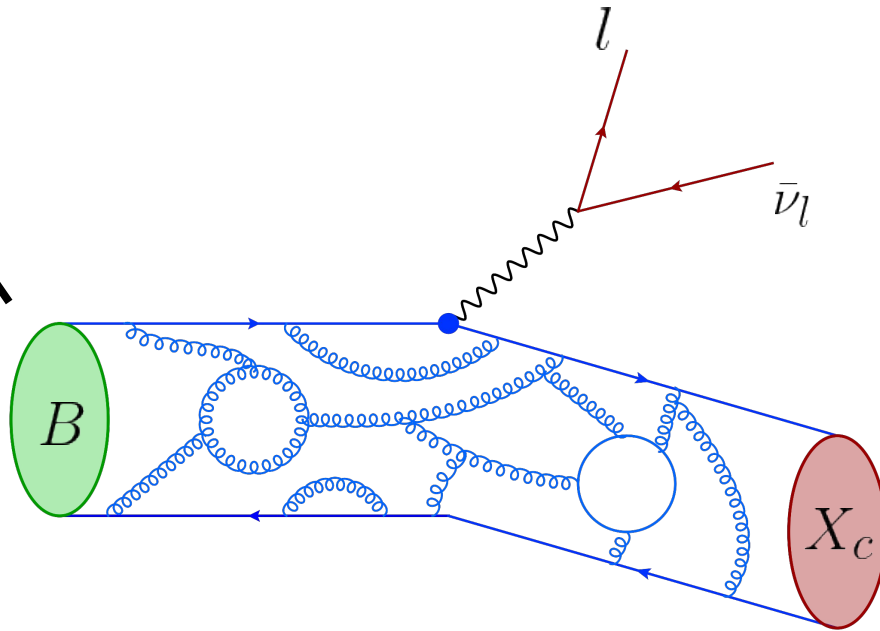
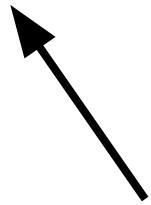


Inclusive semileptonic decays



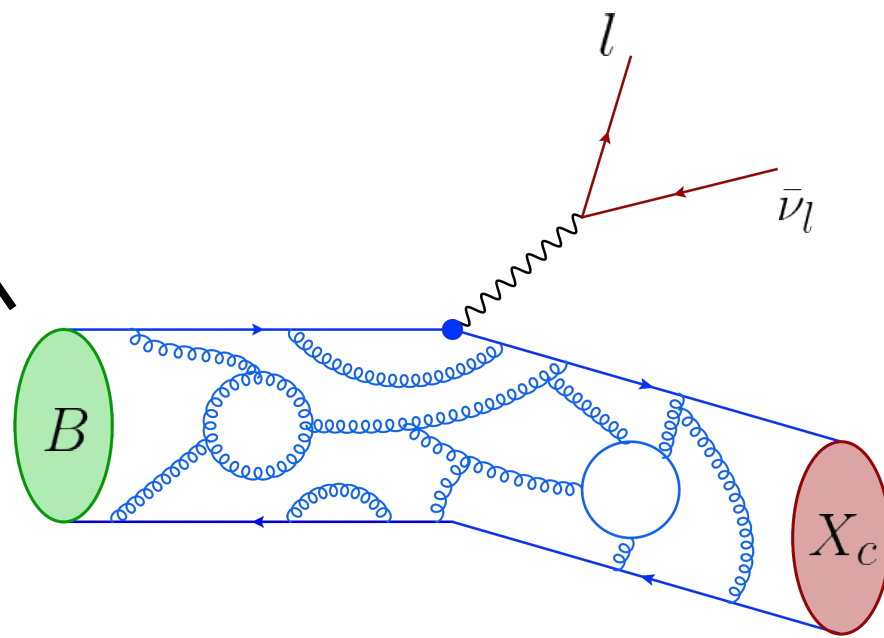
Inclusive semileptonic decays

Γ_{SL}

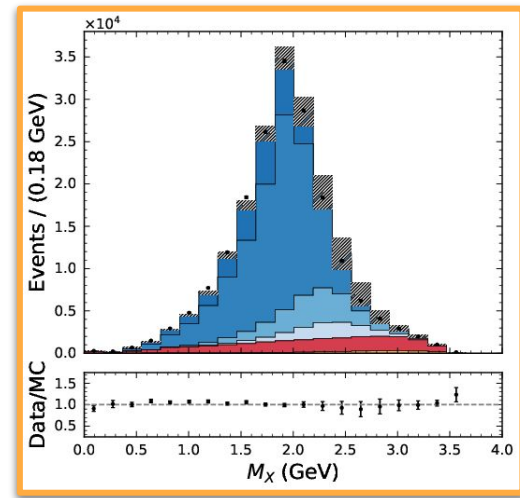


Inclusive semileptonic decays

Γ_{SL}

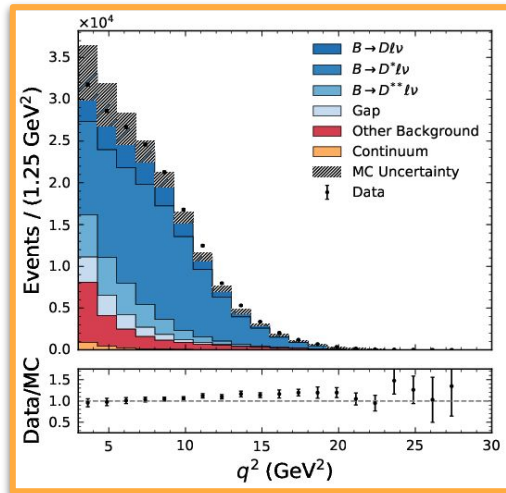
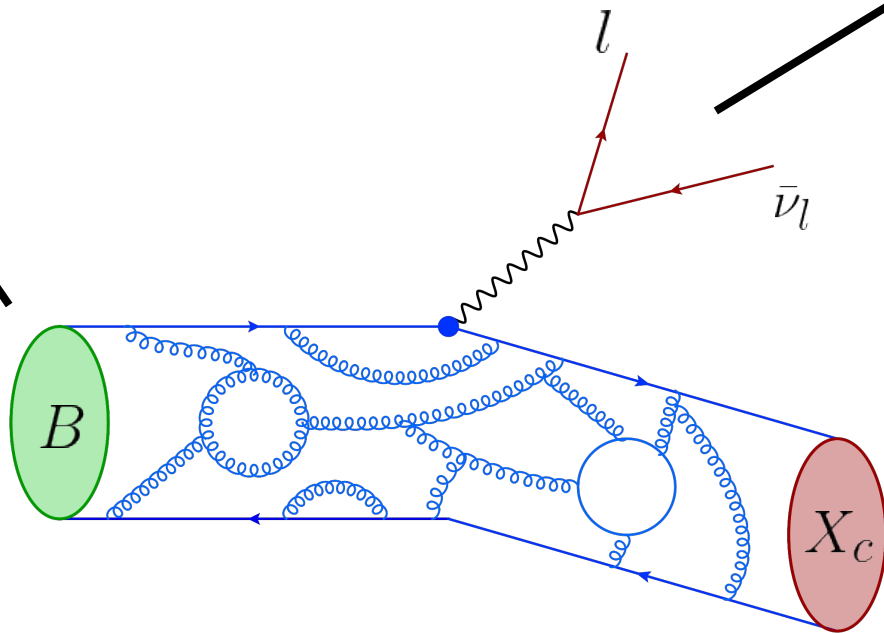


Phys.Rev.D 104 (2021) 11, 112011

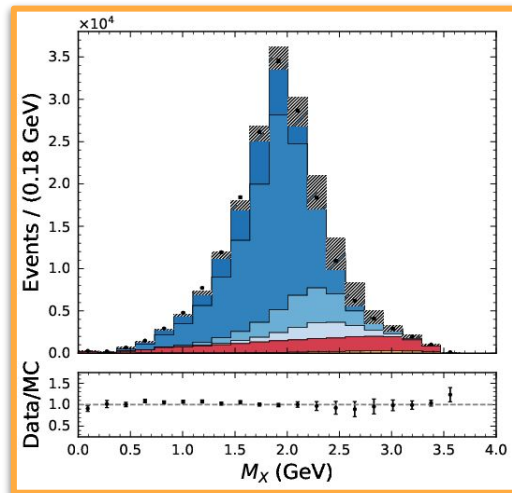


Inclusive semileptonic decays

Γ_{SL}

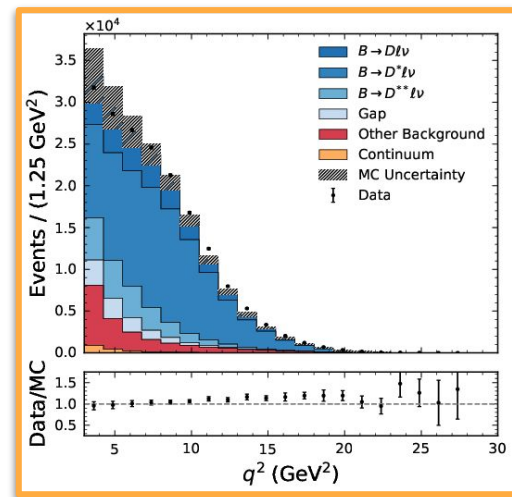
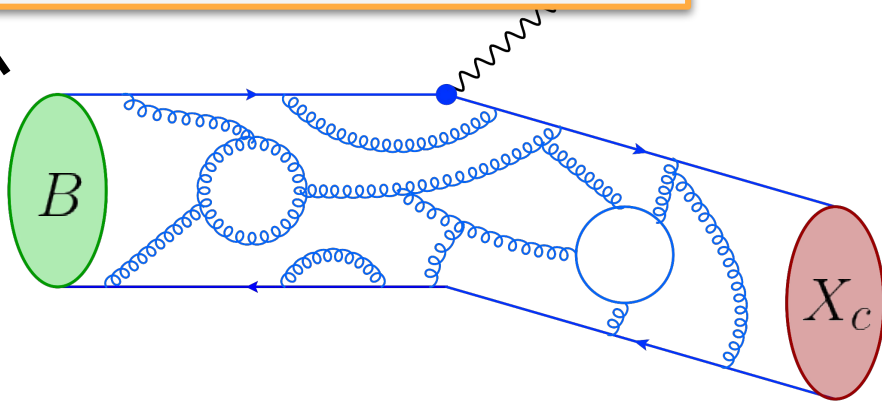


Phys.Rev.D 104 (2021) 11, 112011

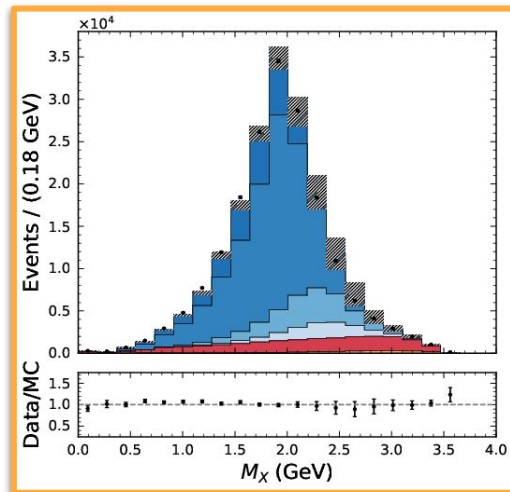


Inclusive semileptonic decays

$$\Gamma_{\text{SL}} = \frac{G_F^2 |V_{cb}|^2}{192\pi^3 m_b^5} \left[\Gamma^{(0)} + \frac{\hat{\mu}_\pi^2}{m_b^2} \Gamma^{(2,\pi)} + \frac{\hat{\mu}_G^2}{m_b^2} \Gamma^{(2,G)} + \frac{\hat{\rho}_D^3}{m_b^3} \Gamma^{(3,D)} + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} \Gamma^{(3,\text{LS})} + \mathcal{O}(m_b^{-4}) \right] \bar{\nu}_l$$

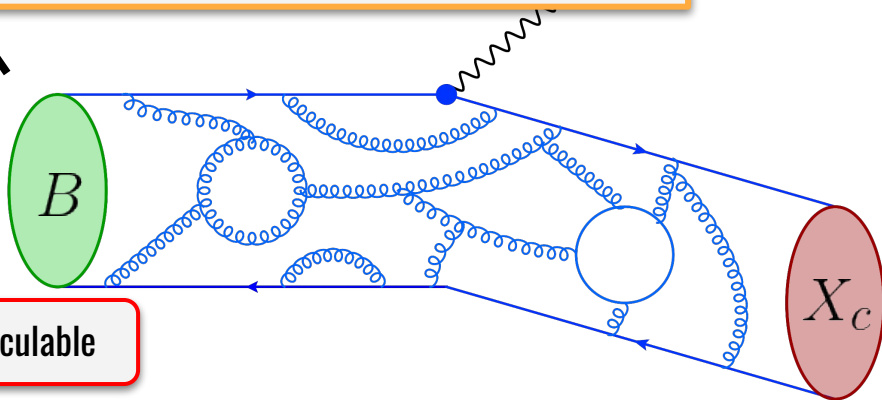


Phys.Rev.D 104 (2021) 11, 112011

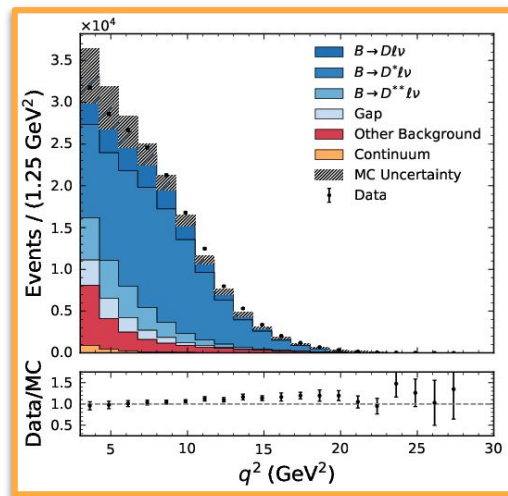


Inclusive semileptonic decays

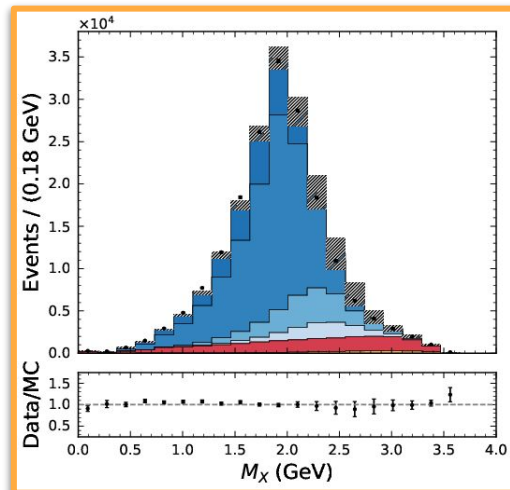
$$\Gamma_{\text{SL}} = \frac{G_F^2 |V_{cb}|^2}{192\pi^3 m_b^5} \left[\Gamma(0) + \frac{\hat{\mu}_\pi^2}{m_b^2} \Gamma(2,\pi) + \frac{\hat{\mu}_G^2}{m_b^2} \Gamma(2,G) + \frac{\hat{\rho}_D^3}{m_b^3} \Gamma(3,D) + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} \Gamma(3,\text{LS}) + \mathcal{O}(m_b^{-4}) \right] \bar{\nu}_l$$



Perturbatively calculable

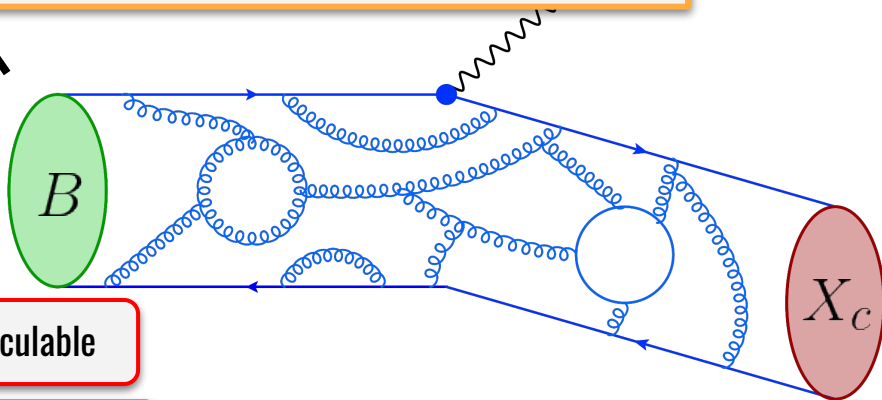


Phys.Rev.D 104 (2021) 11, 112011



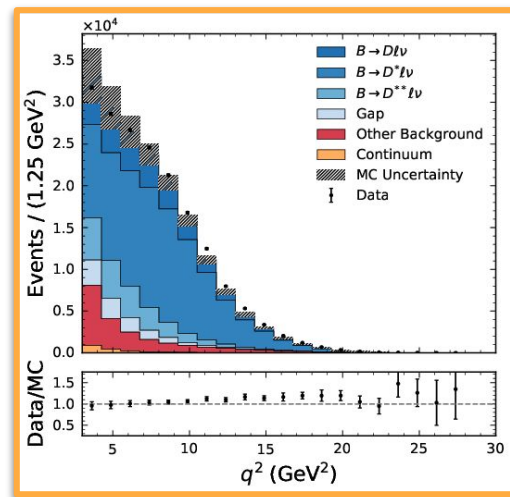
Inclusive semileptonic decays

$$\Gamma_{\text{SL}} = \frac{G_F^2 |V_{cb}|^2}{192\pi^3 m_b^5} \left[\Gamma(0) + \frac{\hat{\mu}_\pi^2}{m_b^2} \Gamma(2,\pi) + \frac{\hat{\mu}_G^2}{m_b^2} \Gamma(2,G) + \frac{\hat{\rho}_D^3}{m_b^3} \Gamma(3,D) + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} \Gamma(3,\text{LS}) + \mathcal{O}(m_b^{-4}) \right] \bar{\nu}_l$$

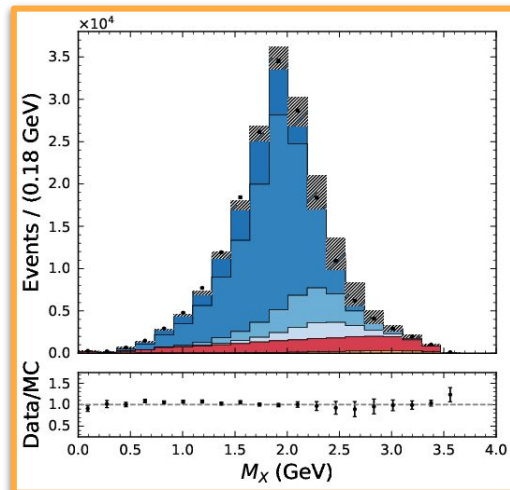


Perturbatively calculable

Nonperturbative parameter

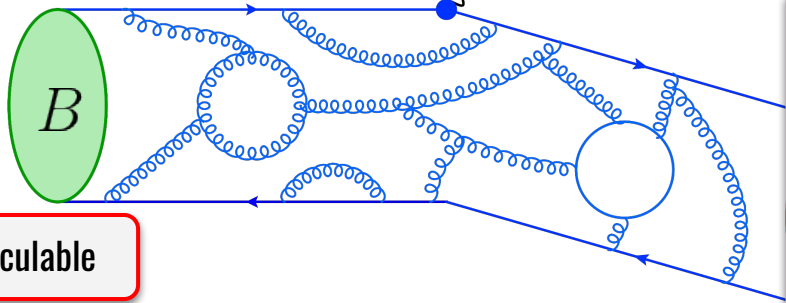


Phys.Rev.D 104 (2021) 11, 112011



Inclusive semileptonic decays

$$\Gamma_{\text{SL}} = \frac{G_F^2 |V_{cb}|^2}{192\pi^3 m_b^5} \left[\Gamma(0) + \frac{\hat{\mu}_\pi^2}{m_b^2} \Gamma(2,\pi) + \frac{\hat{\mu}_G^2}{m_b^2} \Gamma(2,G) + \frac{\hat{\rho}_D^3}{m_b^3} \Gamma(3,D) + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} \Gamma(3,\text{LS}) + \mathcal{O}(m_b^{-4}) \right] \bar{\nu}_l$$

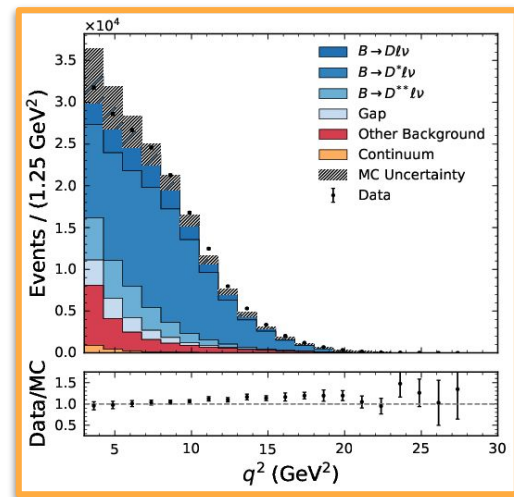


Perturbatively calculable

Nonperturbative parameter

Current uncertainties 10-16%

$$\langle q^2 \rangle = m_b^2 \left[Q(0) + \frac{\hat{\mu}_\pi^2}{m_b^2} Q(2,\pi) + \frac{\hat{\mu}_G^2}{m_b^2} Q(2,G) + \frac{\hat{\rho}_D^3}{m_b^3} Q(3,D) + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} Q(3,\text{LS}) \right]$$



Phys.Rev.D 104 (2021) 11, 112011

Inclusive semileptonic decays

$$\Gamma_{\text{SL}} = \frac{G_F^2 |V_{cb}|^2}{192\pi^3 m_b^5} \left[\Gamma(0) + \frac{\hat{\mu}_\pi^2}{m_b^2} \Gamma(2,\pi) + \frac{\hat{\mu}_G^2}{m_b^2} \Gamma(2,G) + \frac{\hat{\rho}_D^3}{m_b^3} \Gamma(3,D) + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} \Gamma(3,\text{LS}) + \mathcal{O}(m_b^{-4}) \right] \bar{\nu}_l$$

$$|V_{cb}| = 42.16(30)_{\text{th}}(32)_{\text{exp}}(25)_\Gamma 10^{-3}$$

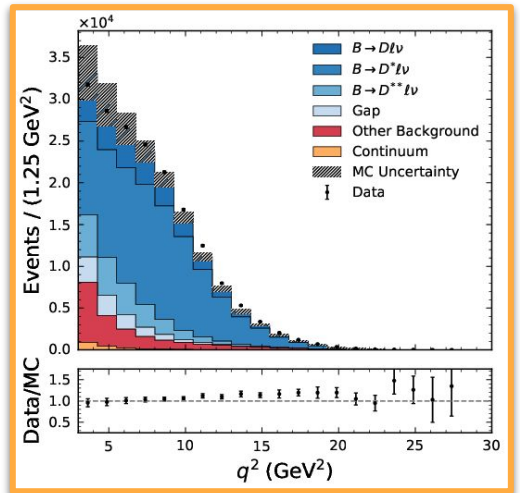
Bordone et al., *Phys.Lett.B* 822 (2021) 136679

Perturbatively calculable

Nonperturbative parameter

Current uncertainties 10-16%

$$\langle q^2 \rangle = m_b^2 \left[Q(0) + \frac{\hat{\mu}_\pi^2}{m_b^2} Q(2,\pi) + \frac{\hat{\mu}_G^2}{m_b^2} Q(2,G) + \frac{\hat{\rho}_D^3}{m_b^3} Q(3,D) + \frac{\hat{\rho}_{\text{LS}}^3}{m_b^3} Q(3,\text{LS}) \right]$$



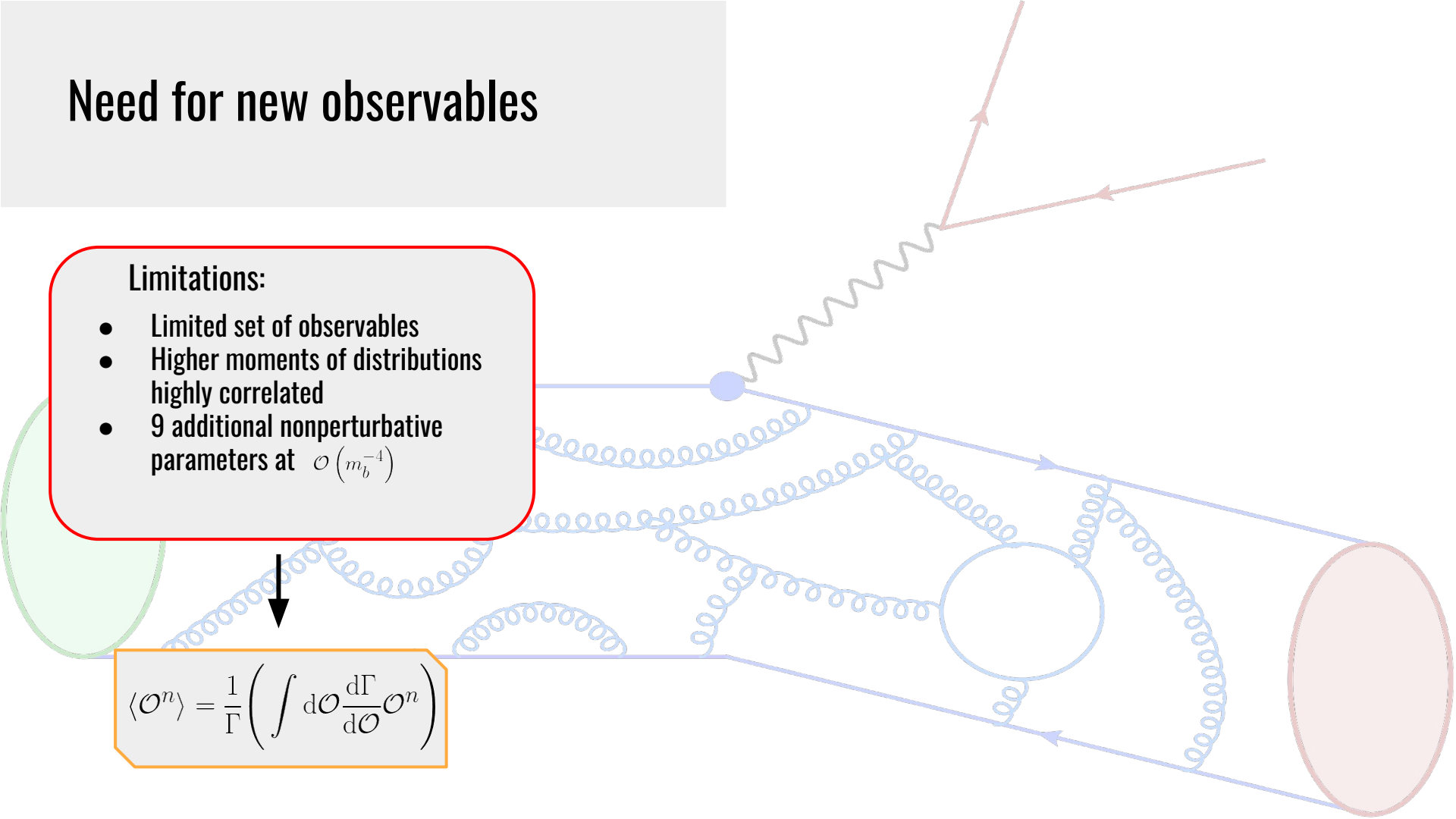
Phys.Rev.D 104 (2021) 11, 112011

Need for new observables

Limitations:

- Limited set of observables
- Higher moments of distributions highly correlated
- 9 additional nonperturbative parameters at $\mathcal{O}(m_b^{-4})$

$$\langle \mathcal{O}^n \rangle = \frac{1}{\Gamma} \left(\int d\mathcal{O} \frac{d\Gamma}{d\mathcal{O}} \mathcal{O}^n \right)$$



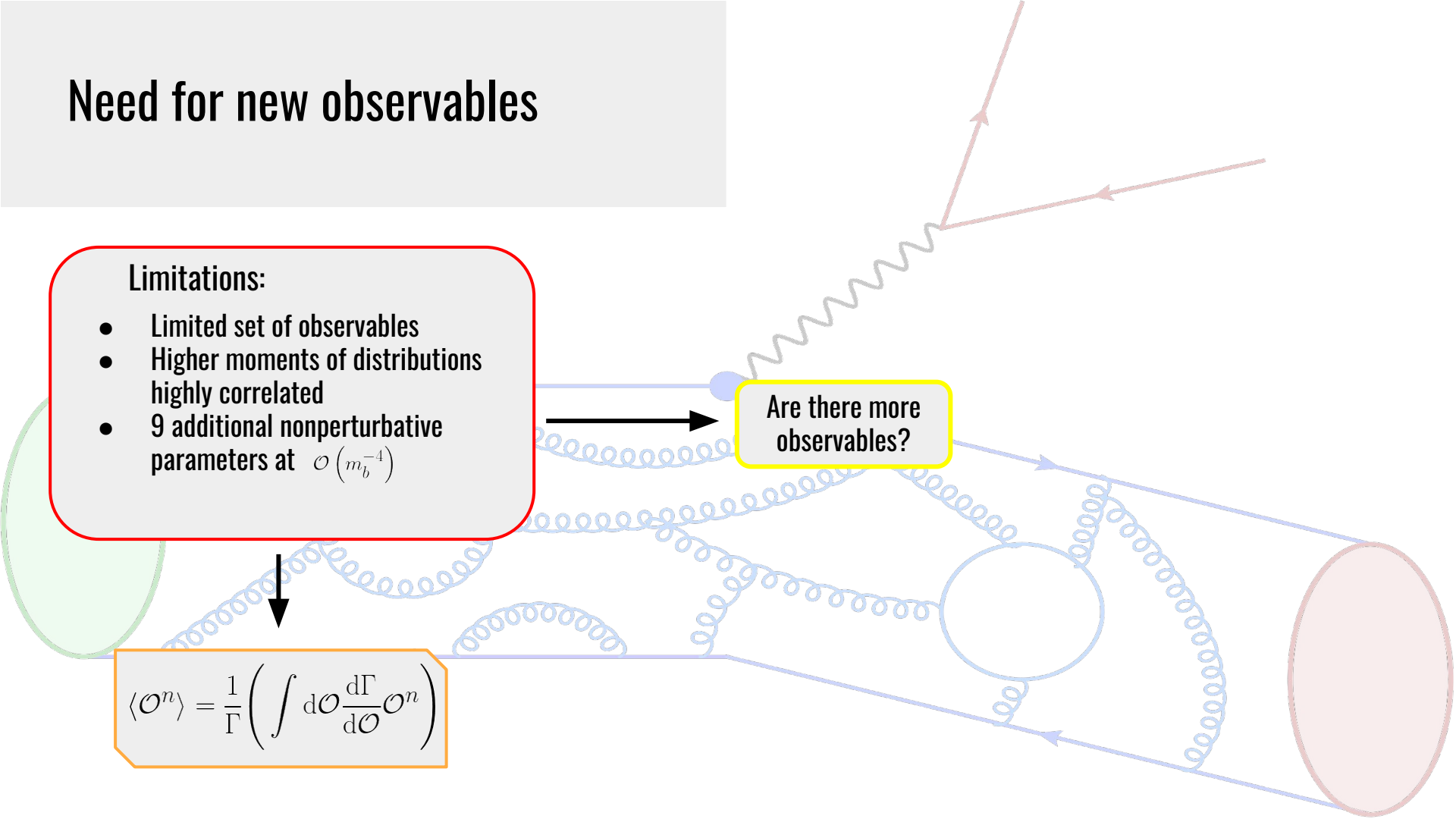
Need for new observables

Limitations:

- Limited set of observables
- Higher moments of distributions highly correlated
- 9 additional nonperturbative parameters at $\mathcal{O}(m_b^{-4})$

Are there more observables?

$$\langle \mathcal{O}^n \rangle = \frac{1}{\Gamma} \left(\int d\mathcal{O} \frac{d\Gamma}{d\mathcal{O}} \mathcal{O}^n \right)$$



Need for new observables

Limitations:

- Limited set of observables
- Higher moments of distributions highly correlated
- 9 additional nonperturbative parameters at $\mathcal{O}(m_b^{-4})$

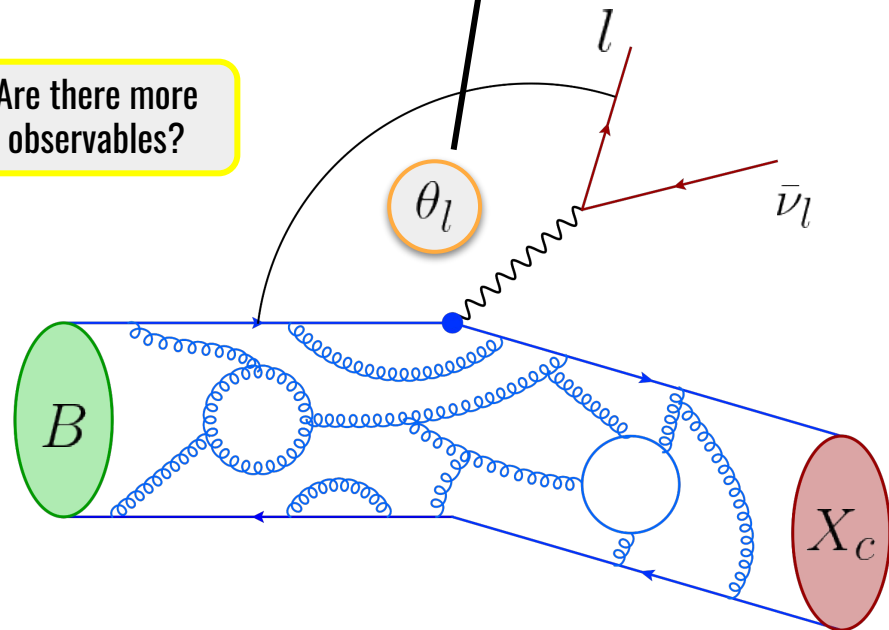
$$\langle \mathcal{O}^n \rangle = \frac{1}{\Gamma} \left(\int d\mathcal{O} \frac{d\Gamma}{d\mathcal{O}} \mathcal{O}^n \right)$$

Yes!
The forward-backward
asymmetry

S. Turczyk, *JHEP* 04 (2016) 131

$$\mathcal{A}_{FB} = \frac{1}{\Gamma} \left(\int_{-1}^0 dz \frac{d\Gamma}{dz} - \int_0^1 dz \frac{d\Gamma}{dz} \right)$$
$$z = \cos \theta_l = \frac{E_\nu - E_l}{\sqrt{(E_\nu + E_l)^2 - q^2}}$$

Are there more
observables?



Need for new observables

Limitations:

- Limited set of observables
- Higher moments of distributions highly correlated
- 9 additional nonperturbative parameters at $\mathcal{O}(m_b^{-4})$

There are even more!

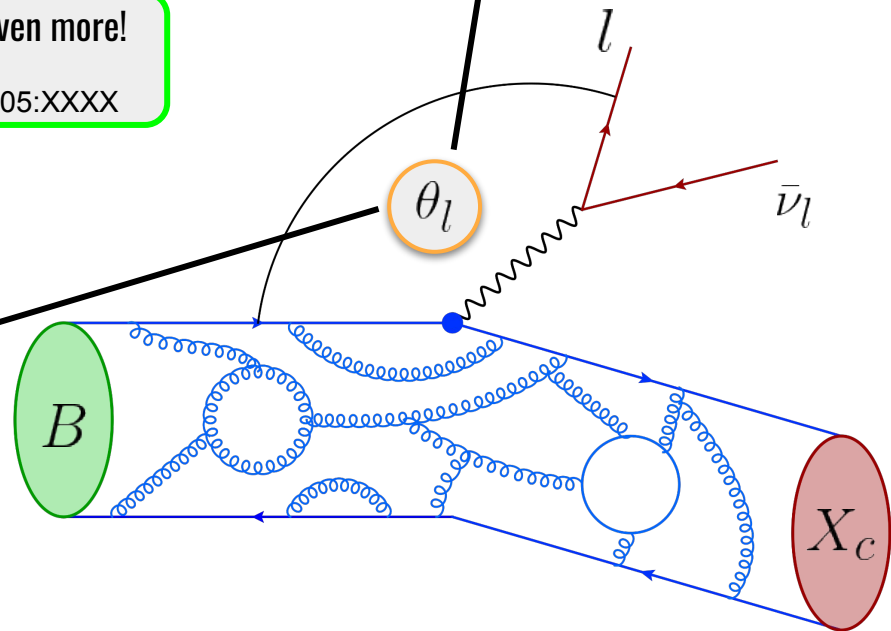
FH, arXiv:2205:XXXX

Yes!
The forward-backward
asymmetry

S. Turczyk, *JHEP* 04 (2016) 131

$$\mathcal{A}_{FB} = \frac{1}{\Gamma} \left(\int_{-1}^0 dz \frac{d\Gamma}{dz} - \int_0^1 dz \frac{d\Gamma}{dz} \right)$$
$$z = \cos \theta_l = \frac{E_\nu - E_l}{\sqrt{(E_\nu + E_l)^2 - q^2}}$$

$$\langle \mathcal{O}^n \rangle_{\pm} = \frac{1}{\Gamma} \left(\int d\mathcal{O} \left[\int_{-1}^0 dz \frac{d\Gamma}{d\mathcal{O}dz} \mathcal{O}^n \pm \int_0^1 dz \frac{d\Gamma}{d\mathcal{O}dz} \mathcal{O}^n \right] \right)$$



The angular spectrum

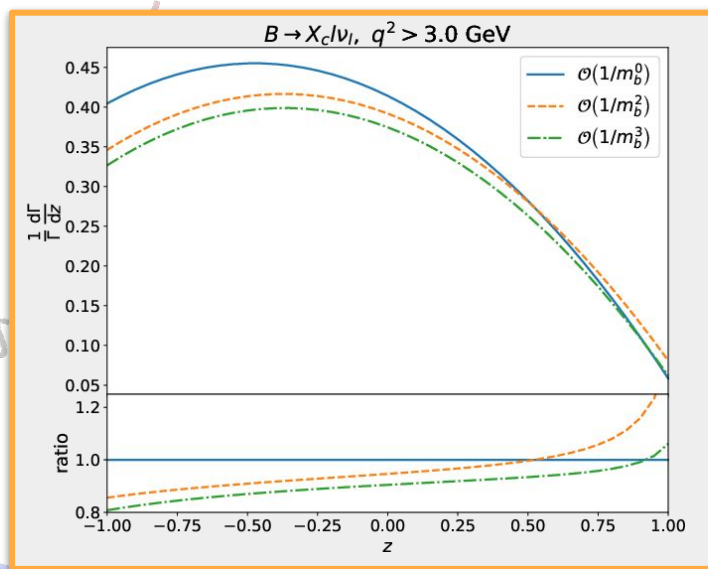
$$\Gamma \approx \frac{G_F^2 |V_{cb}|^2 m_b^5}{192\pi^3} 0.706 \left(1 - 0.011|\hat{\rho}_\pi^2 - 0.029|\hat{\rho}_G^2 - 0.031|\hat{\rho}_D^3 - 0.003|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

$$A_{FB} \approx 0.192 \left(1 - 0.051|\hat{\rho}_\pi^2 - 0.112|\hat{\rho}_G^2 + 0.024|\hat{\rho}_D^3 + 0.009|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

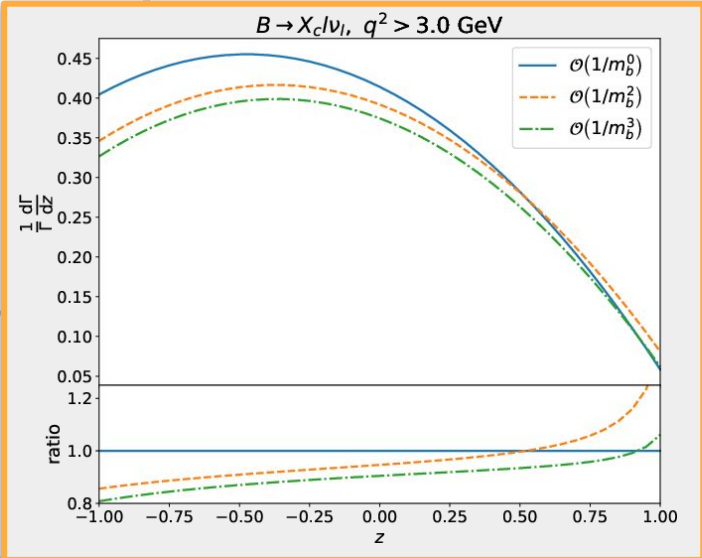
$$\langle M_X^2 \rangle_+ \approx 0.202 m_b^2 \left(1 - 0.076|\hat{\rho}_\pi^2 + 0.035|\hat{\rho}_G^2 + 0.040|\hat{\rho}_D^3 + 0.004|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

$$\langle E_l \rangle_+ \approx 0.316 m_b \left(1 + 0.011|\hat{\rho}_\pi^2 - 0.019|\hat{\rho}_G^2 - 0.011|\hat{\rho}_D^3 - 0.0004|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

$$\langle q^2 \rangle_- \approx 0.056 m_b^2 \left(1 - 0.081|\hat{\rho}_\pi^2 - 0.173|\hat{\rho}_G^2 + 0.014|\hat{\rho}_D^3 + 0.017|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$



The angular spectrum



$$\Gamma \sim \frac{G_F^2 |V_{cb}|^2 m_b^5}{192\pi^3} 0.706 \left(1 - 0.011|\hat{\mu}_\pi^2| - 0.029|\hat{\mu}_G^2| - 0.031|\hat{\rho}_D^3| - 0.003|\hat{\rho}_{LS}^3| + \mathcal{O}(1/m_b^4) \right)$$

$$A_{FB} \approx 0.192 \left(1 - 0.051|\hat{\mu}_\pi^2| - 0.112|\hat{\mu}_G^2| + 0.024|\hat{\rho}_D^3| + 0.009|\hat{\rho}_{LS}^3| + \mathcal{O}(1/m_b^4) \right)$$

$$\langle M_X^2 \rangle_+ \approx 0.202 m_b^2 \left(1 - 0.076|\hat{\mu}_\pi^2| + 0.035|\hat{\mu}_G^2| + 0.040|\hat{\rho}_D^3| + 0.004|\hat{\rho}_{LS}^3| + \mathcal{O}(1/m_b^4) \right)$$

$$\langle E_l \rangle_+ \approx 0.316 m_b \left(1 + 0.011|\hat{\mu}_\pi^2| - 0.019|\hat{\mu}_G^2| - 0.011|\hat{\rho}_D^3| - 0.0004|\hat{\rho}_{LS}^3| + \mathcal{O}(1/m_b^4) \right)$$

Sum of $\hat{\mu}_\pi^2$ and $\hat{\mu}_G^2$

Difference of $\hat{\mu}_\pi^2$ and $\hat{\mu}_G^2$

$$\langle q^2 \rangle_- \approx 0.056 m_b^2 \left(1 - 0.081|\hat{\mu}_\pi^2| - 0.173|\hat{\mu}_G^2| + 0.014|\hat{\rho}_D^3| + 0.017|\hat{\rho}_{LS}^3| + \mathcal{O}(1/m_b^4) \right)$$

The angular spectrum

$$\Gamma \approx \frac{G_F^2 |V_{cb}|^2 m_b^5}{192\pi^3} 0.706 \left(1 - 0.011|\hat{\mu}_\pi^2 - 0.029|\hat{\mu}_G^2 - 0.031|\hat{\rho}_D^3 - 0.003|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

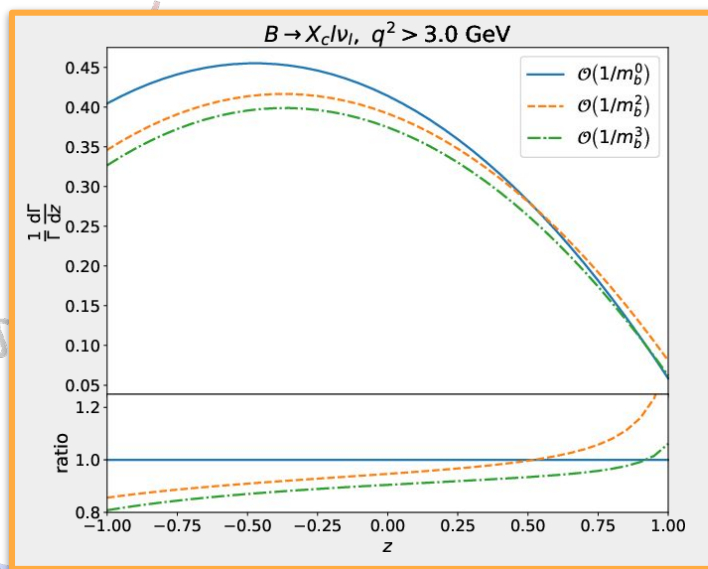
$$A_{FB} \approx 0.192 \left(1 - 0.051|\hat{\mu}_\pi^2 - 0.112|\hat{\mu}_G^2 + 0.024|\hat{\rho}_D^3 + 0.009|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

$$\langle M_X^2 \rangle_+ \approx 0.202 m_b^2 \left(1 - 0.076|\hat{\mu}_\pi^2 - 0.035|\hat{\mu}_G^2 + 0.040|\hat{\rho}_D^3 + 0.004|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

$$\langle E_l \rangle_+ \approx 0.316 m_b \left(1 + 0.011|\hat{\mu}_\pi^2 - 0.019|\hat{\mu}_G^2 - 0.011|\hat{\rho}_D^3 - 0.0004|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

Strongest dependence on $\hat{\rho}_D^3$

$$\langle q^2 \rangle_- \approx 0.056 m_b^2 \left(1 - 0.081|\hat{\mu}_\pi^2 - 0.173|\hat{\mu}_G^2 + 0.014|\hat{\rho}_D^3 + 0.017|\hat{\rho}_{LS}^3 + \mathcal{O}(1/m_b^4) \right)$$

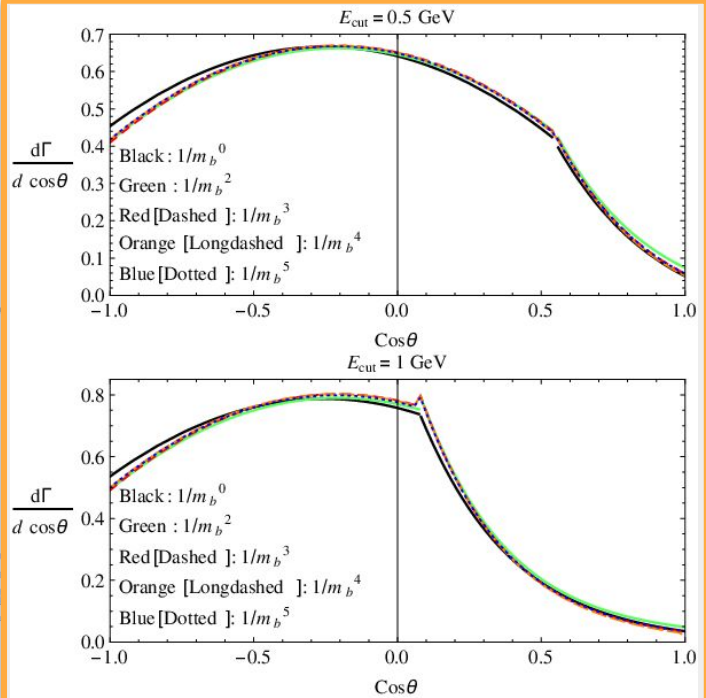


Possibility to determine $\hat{\rho}_{LS}^3$ more precisely?

No measurement without cuts

Original idea: cut on lepton energy

S. Turczyk, *JHEP* 04 (2016) 131

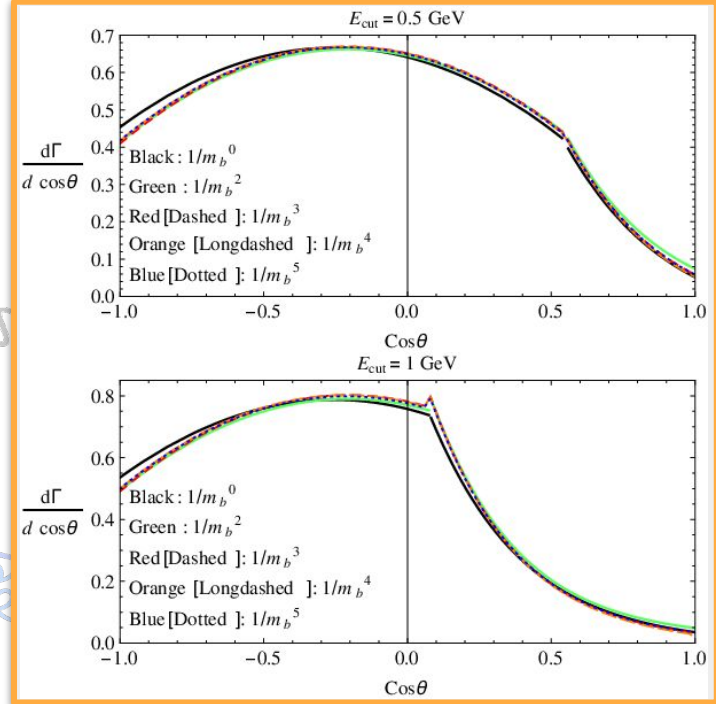


No measurement without cuts

Original idea: cut on lepton energy

S. Turczyk, *JHEP* 04 (2016) 131

- Discontinuity in spectrum
- Cut of 0.5 GeV unrealistic for muons

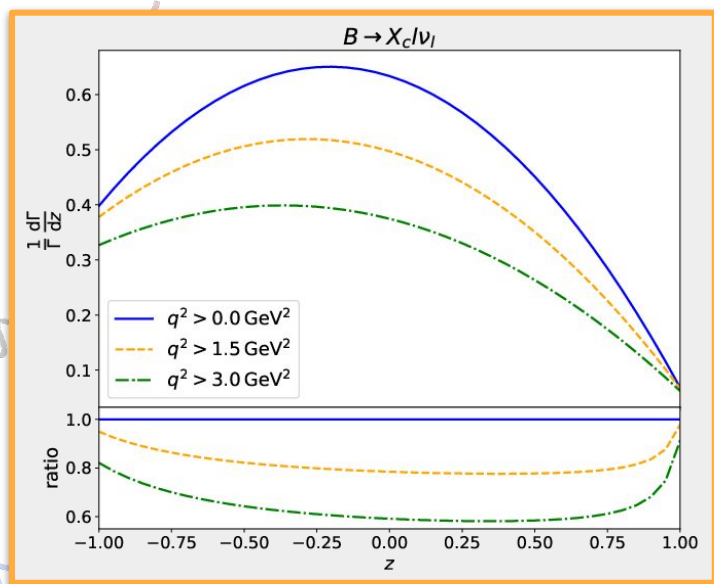


No measurement without cuts

Better idea: cut on q^2

FH, arXiv:2205:XXXX

- Discontinuity in spectrum
- Cut of 0.5 GeV unrealistic for muons

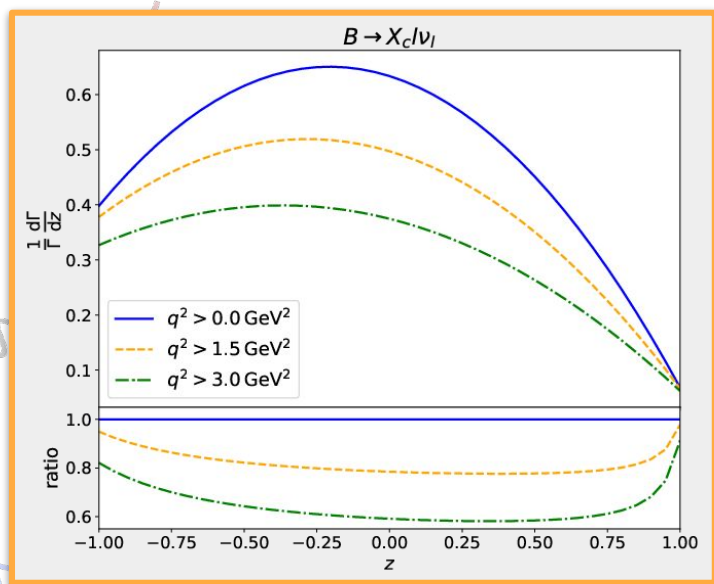


No measurement without cuts

Better idea: cut on q^2

FH, arXiv:2205:XXXX

- Remains polynomial
- Cut can be varied

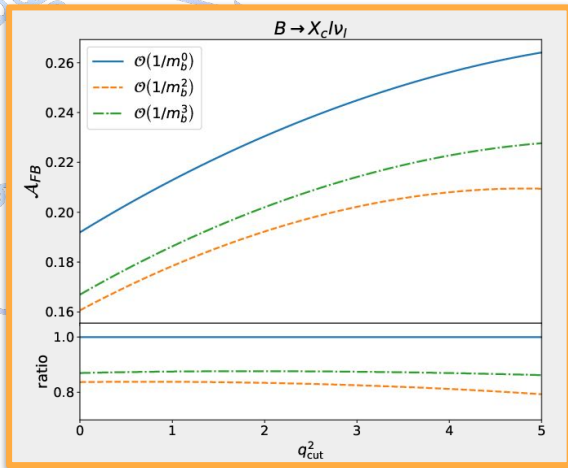
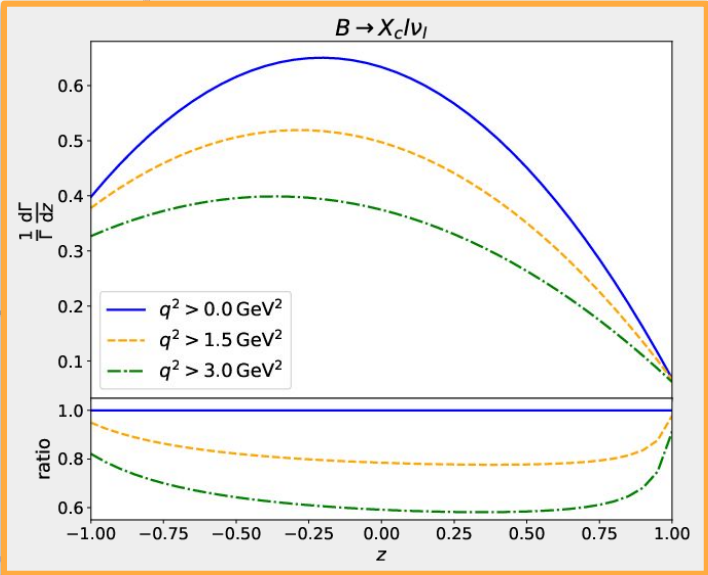


No measurement without cuts

Better idea: cut on q^2

FH, arXiv:2205:XXXX

- Remains polynomial
- Cut can be varied



No measurement without cuts

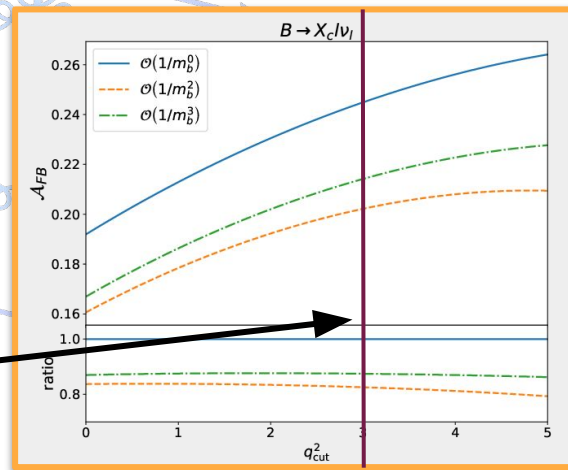
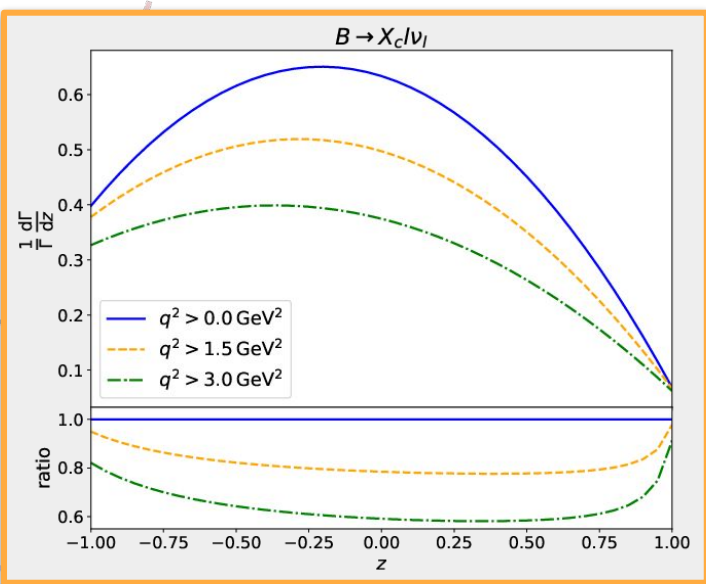
Better idea: cut on q^2

FH, arXiv:2205:XXXX

- Remains polynomial
- Cut can be varied

A cut of 3 GeV^2 realistic

Phys.Rev.D 104 (2021) 11, 112011



No measurement without cuts

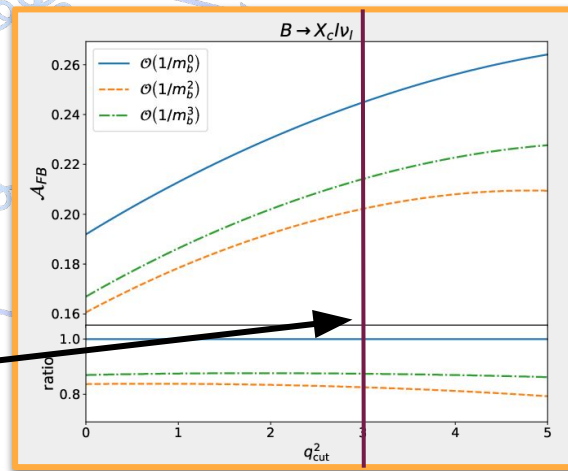
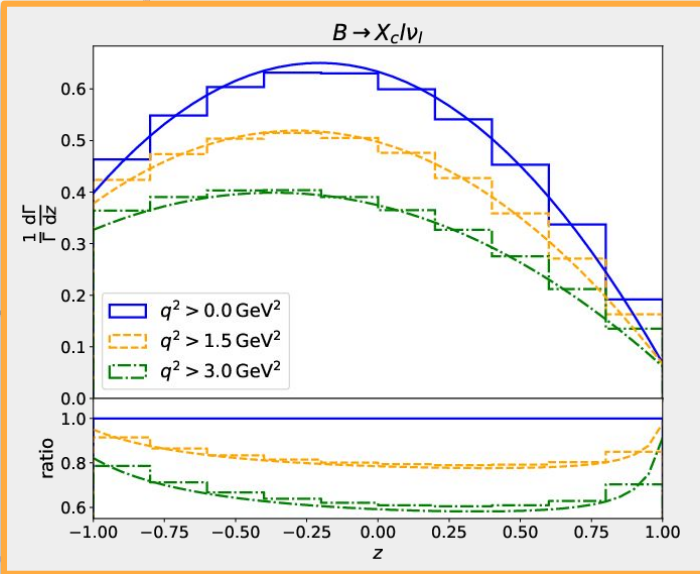
Compare to Sherpa

FH, arXiv:2205:XXXX

- Remains polynomial
- Cut can be varied

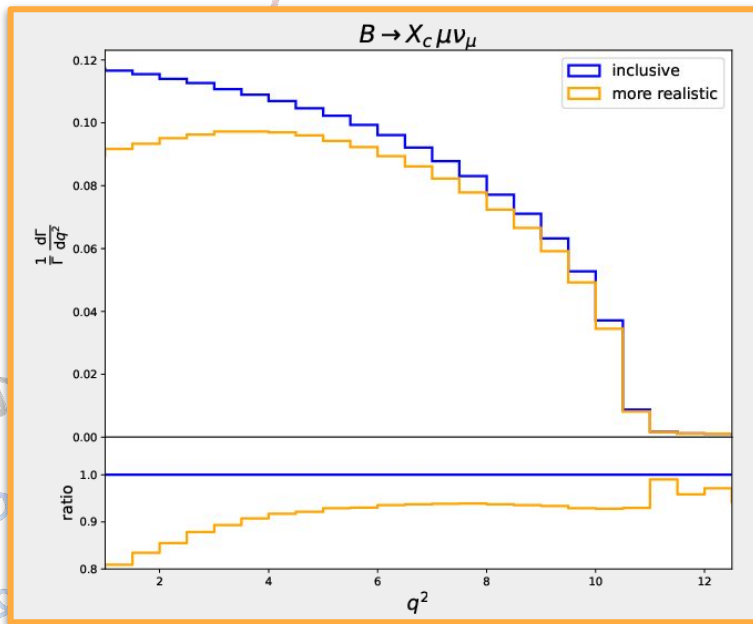
A cut of 3 GeV^2 realistic

Phys.Rev.D 104 (2021) 11, 112011



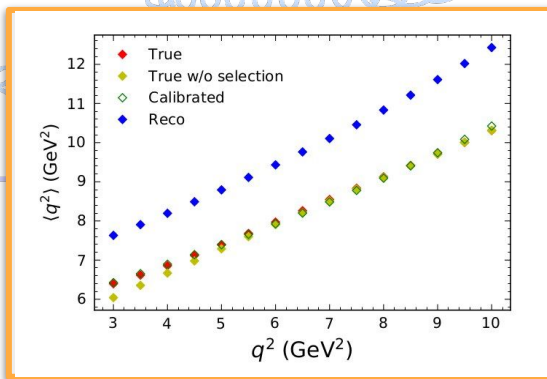
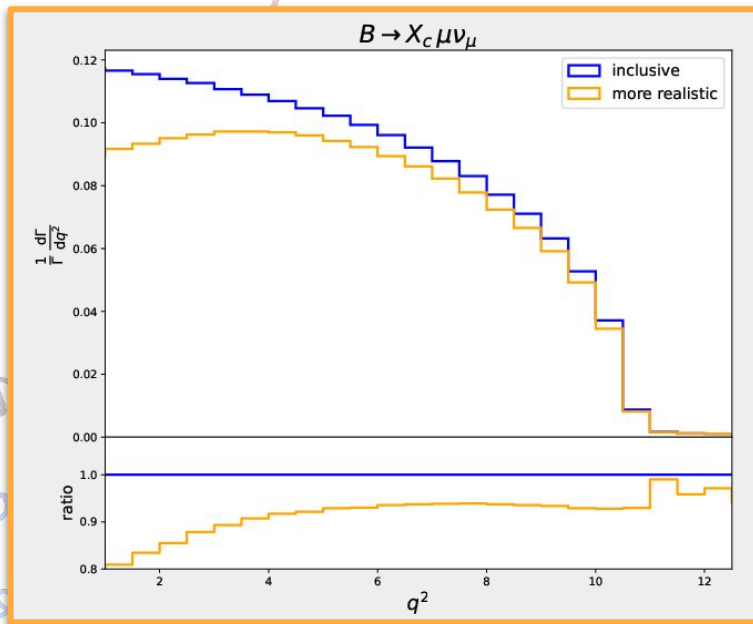
Getting closer to the experiment

- We are interested in fully inclusive observables
- Experiment has to deal with final state radiation and additional cuts



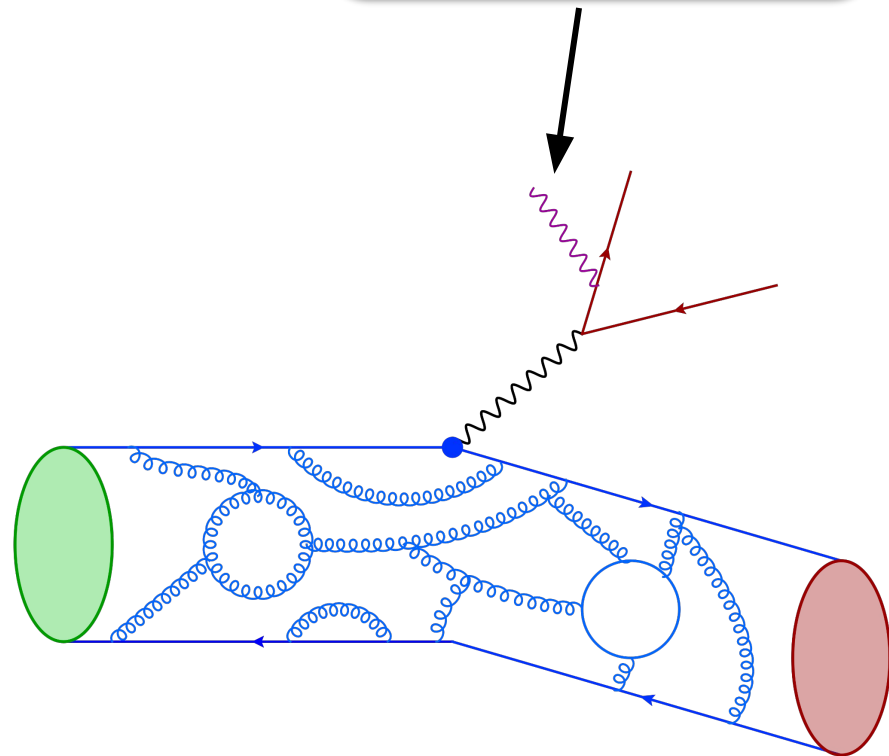
Getting closer to the experiment

- We are interested in fully inclusive observables
- Experiment has to deal with final state radiation and additional cuts



Finally: Some Radiation (FSR)

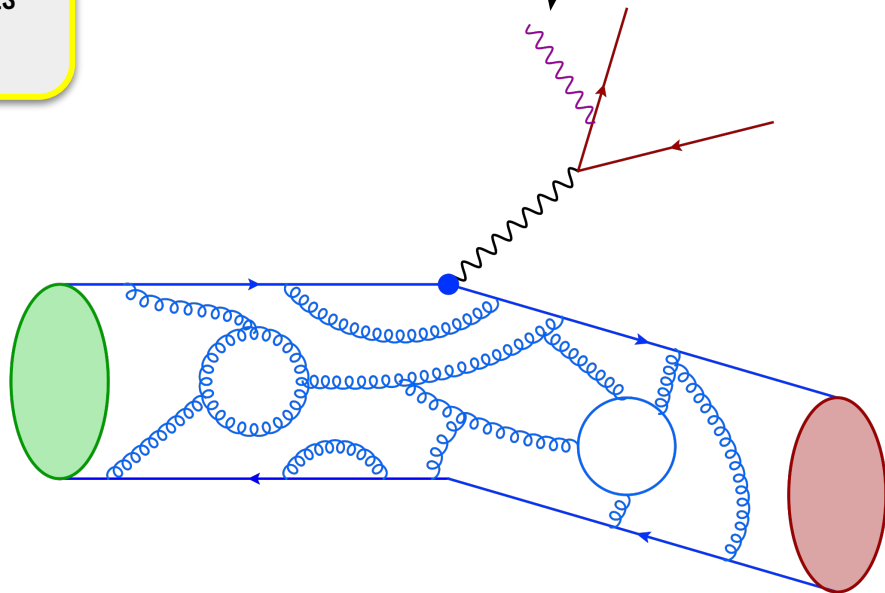
- FSR affects electrons more than muons
- Possibly larger impact on angular observables



Finally: Some Radiation (FSR)

- FSR affects electrons more than muons
- Possibly larger impact on angular observables

Check if simulations and corrections by experiments capture FSR correctly

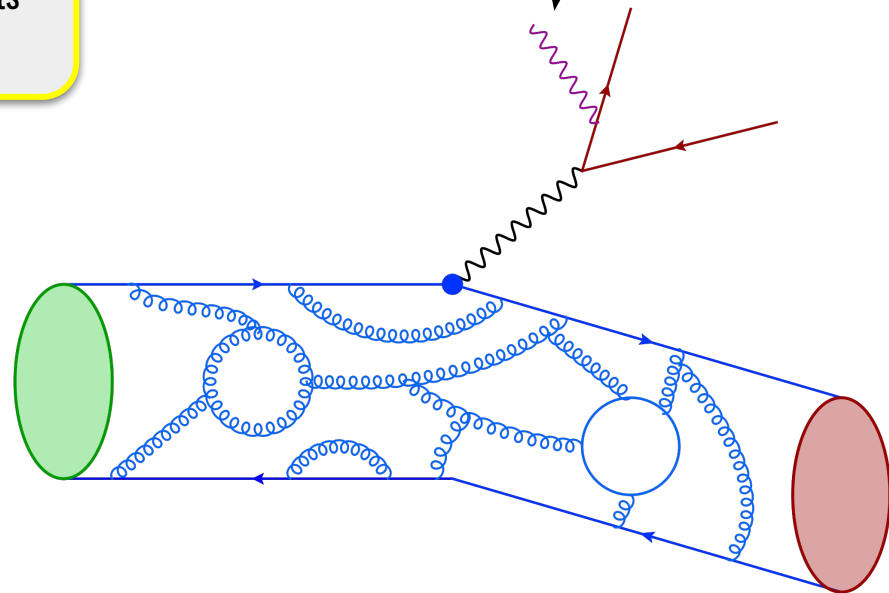


Finally: Some Radiation (FSR)

- FSR affects electrons more than muons
- Possibly larger impact on angular observables

Check if simulations and corrections by experiments capture FSR correctly

- Simulation with PHOTOS
- Merging of electrons with hard-collinear photons



Finally: Some Radiation (FSR)

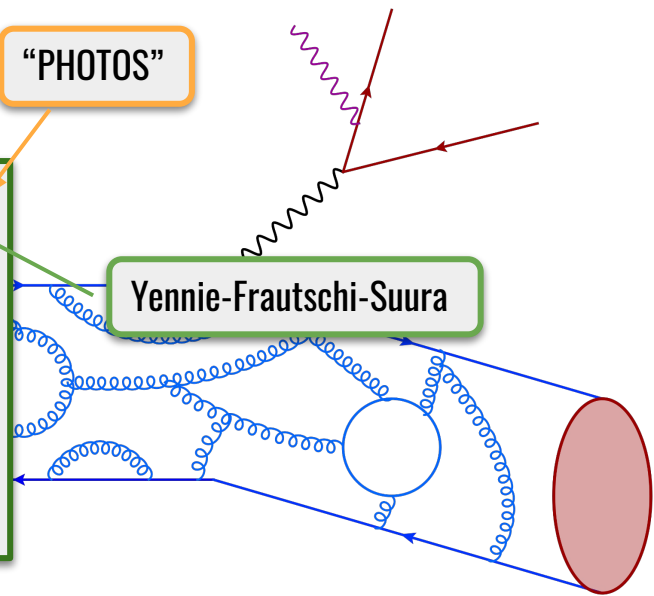
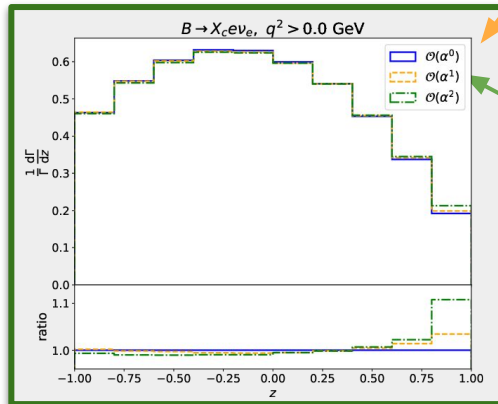
- FSR affects electrons more than muons
- Possibly larger impact on angular observables

Check if simulations and corrections by experiments capture FSR correctly

- Simulation with PHOTOS
- Merging of electrons with hard-collinear photons

“PHOTOS”

Yennie-Frautschi-Suura



Finally: Some Radiation (FSR)

- FSR affects electrons more than muons
- Possibly larger impact on angular observables

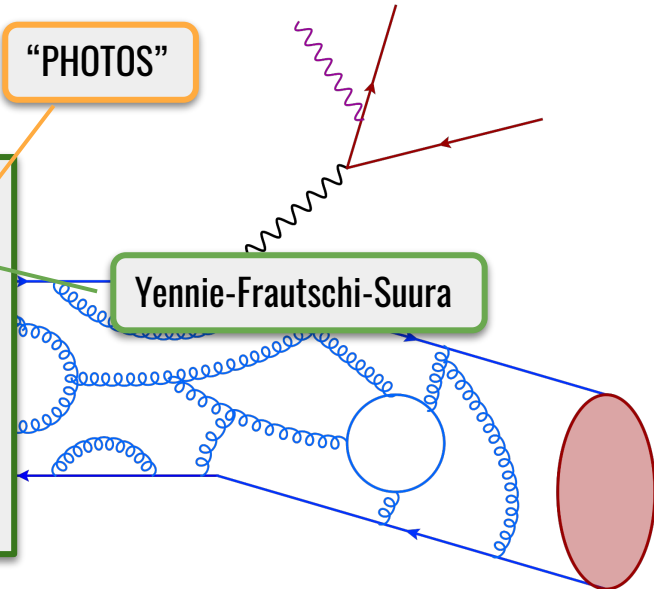
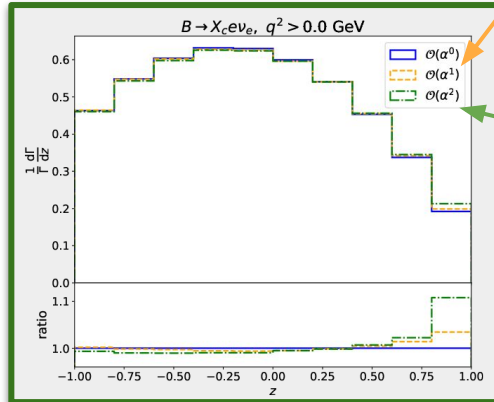
Check if simulations and corrections by experiments capture FSR correctly

- Simulation with PHOTOS
- Merging of electrons with hard-collinear photons

“PHOTOS”

Yennie-Frautschi-Suura

- Sizeable differences
- FSR affects spectrum asymmetrically



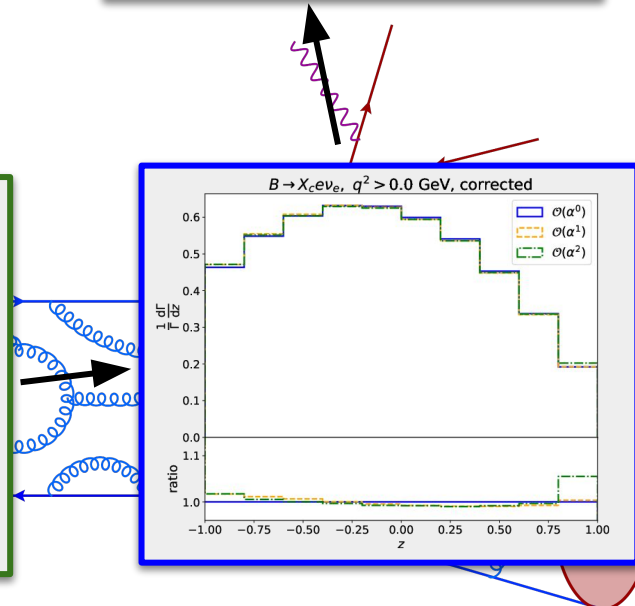
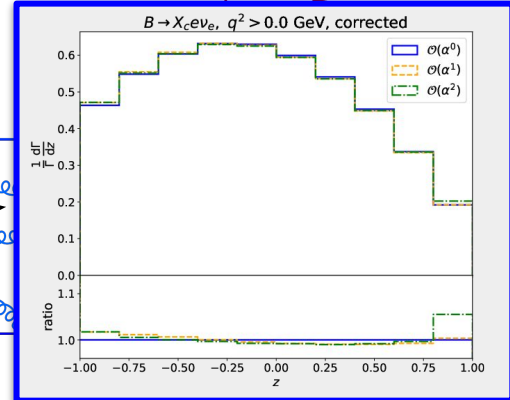
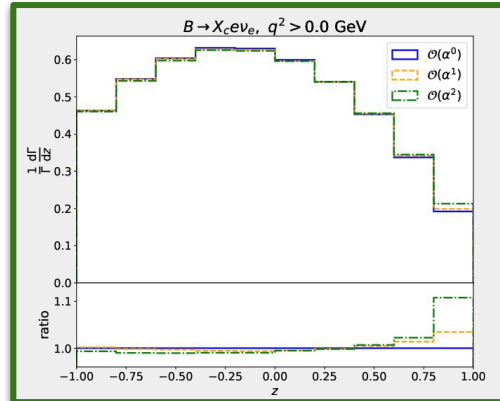
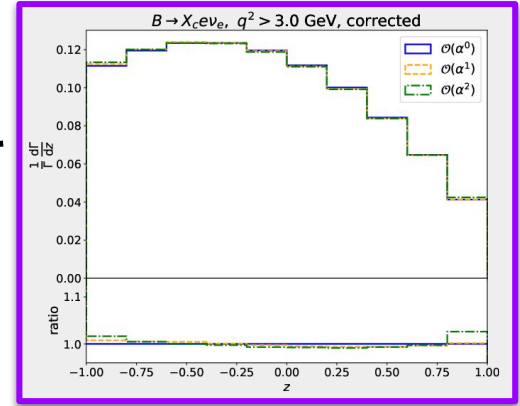
Finally: Some Radiation (FSR)

Cutting on q^2
improves situation

Check if simulations and
corrections by experiments
capture FSR correctly ✓

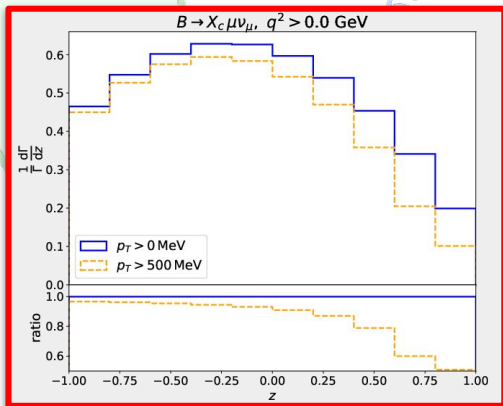
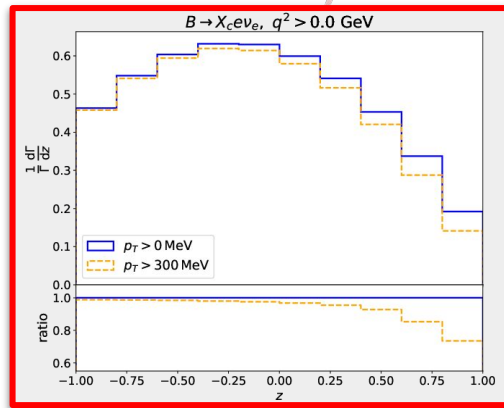
- Simulation with PHOTOS
- Merging of electrons with hard-collinear photons

- Sizeable differences
- FSR affects spectrum asymmetrically



Lepton identification

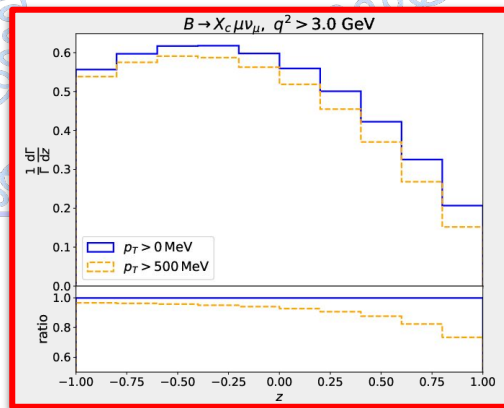
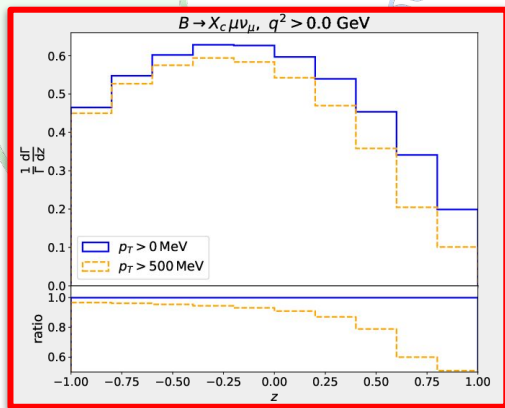
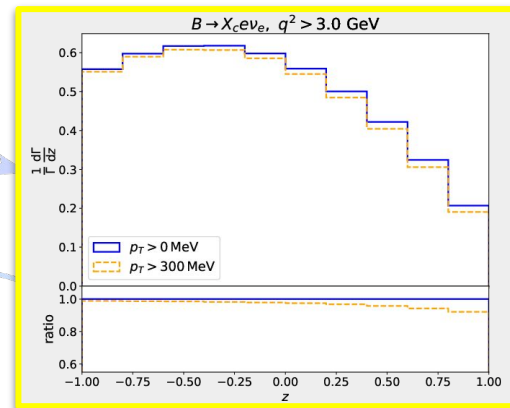
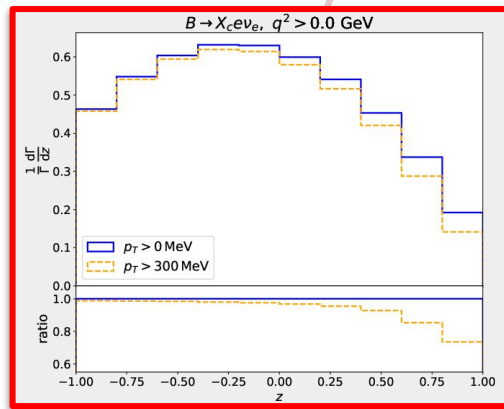
- Leptons need to reach respective subsystems
- Transv. momentum of 300 MeV and 500 MeV



Lepton identification

- Leptons need to reach respective subsystems
- Transv. momentum of 300 MeV and 500 MeV

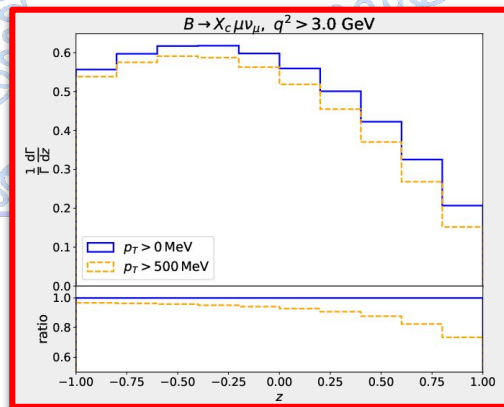
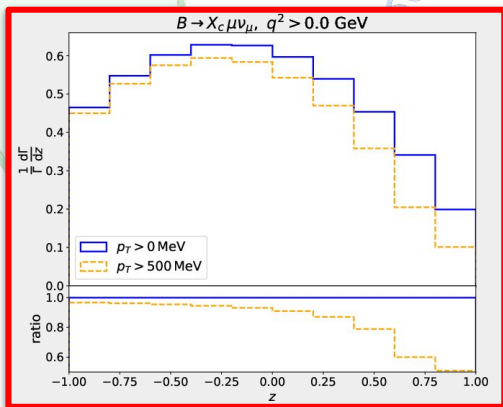
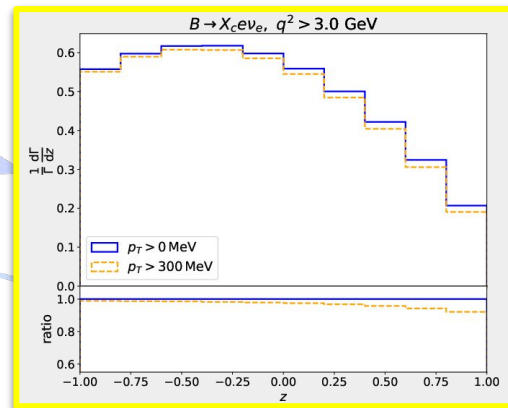
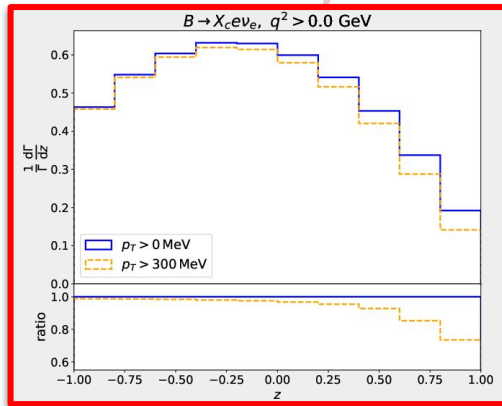
Cut on q^2



Lepton identification

- Leptons need to reach respective subsystems
- Transv. momentum of 300 MeV and 500 MeV

Cut on q^2



- Shifts \mathcal{A}_{FB} for electrons by 6%
 - Shifts \mathcal{A}_{FB} for muons by 20%
- ➔ Unfolding the distribution?

Conclusions & Outlook

Paper out later today!

Conclusions:

- Belle II has the unique ability to study angular observables in inclusive decays
- Angular observables provide additional information on nonperturbative parameters
- A cut on the invariant mass of the lepton-neutrino system is advantageous
- Final state radiation is under control
- Experimental analyses need to carefully correct for lepton identification requirements

Outlook:

- Inclusive decays can cross-check tensions with the SM in exclusive decays
- Measurement of the forward-backward asymmetry could be done already
- Need higher order radiative corrections to take full advantage of a measurement
- Can we do this for the tau?
Need mass effects in the calculation
- Are there additional observables that could benefit CKM element determinations?