



AMHERST CENTER FOR FUNDAMENTAL INTERACTIONS

*Physics at the interface: Energy, Intensity, and Cosmic frontiers*

University of Massachusetts Amherst



## From Zero to Hero... to Zero?

$0\nu\beta\beta$ -decay, energy frontier probes, and the origin of matter

**Phenomenology 2022 Symposium**

**May 9 - 11, 2022**

Sebastián Urrutia-Quiroga

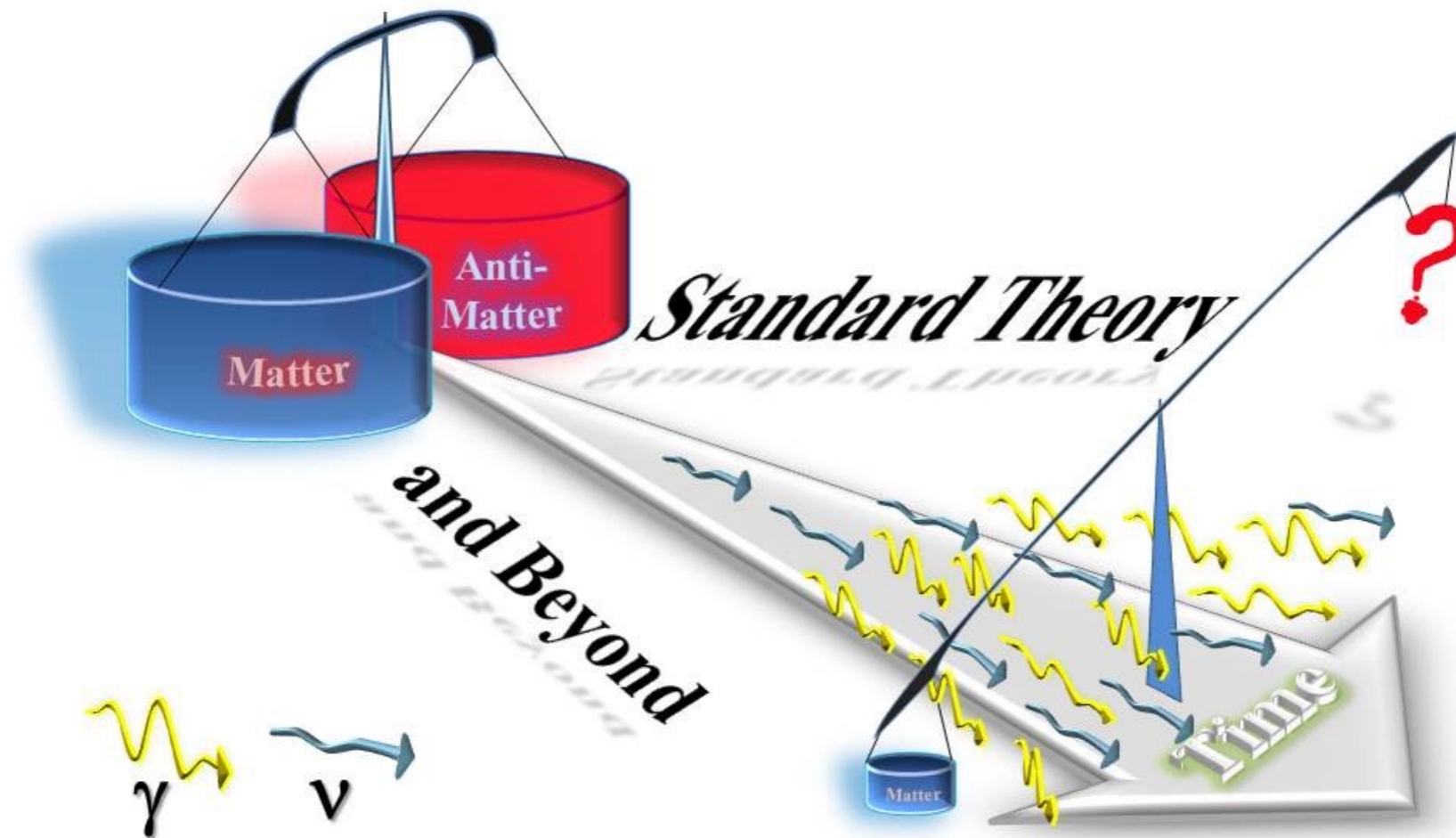
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(in collaboration with Julia Harz, Michael Ramsey-Musolf, and Tianyang Shen)

# I. From Zero to Hero

$$Y_B^{(\text{obs})} = \frac{n_B - n_{\bar{B}}}{s} = (8.66 \pm 0.04) \times 10^{-11}$$

more (**B**aryonic)  
matter than  
anti-matter!



Willmann & Jungmann (2015)

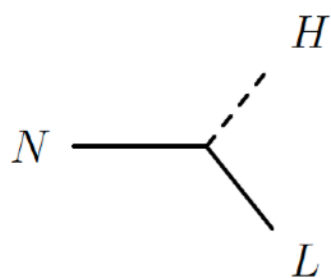
How did we  
go from  
nothing to  
something?

- Simple ingredients: **Sakharov conditions (1967)**

Non-equilibrium dynamics:

Heavy RHNs  $N$  with out-of-equilibrium,  $\Delta\mathcal{L} \neq 0$  decays

$\rightsquigarrow m_{N0}, K$

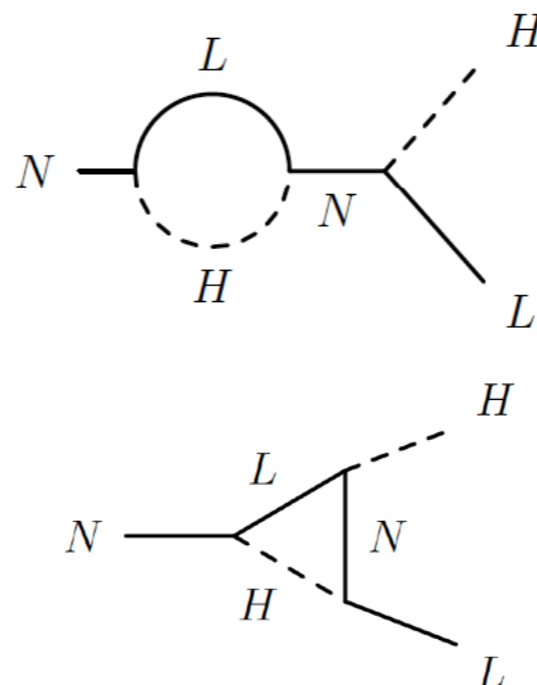


$C$  &  $CP$ -violation:

Quantum corrections,

$$\Gamma(N \rightarrow HL) \neq$$

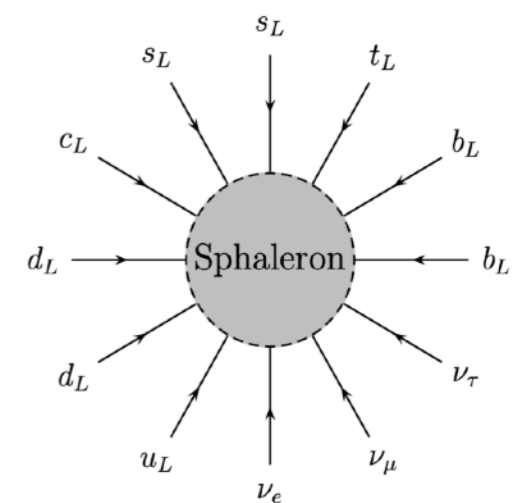
$$\Gamma(N \rightarrow \bar{H}\bar{L}) \rightsquigarrow \epsilon$$



$\mathcal{B}$ -violation:

EW sphaleron processes,

$$\Delta\mathcal{L} \rightarrow \Delta\mathcal{B}$$



Buchmüller (1998)

## II. From Hero... to Zero?

- Thermal LG works, but requires  $m_{N0} \gtrsim 10^9 \text{ GeV}$  Davidson & Ibarra (2002)

- What if we include accessible effects? **Energy scale?**

- LNV at the **TeV-scale**: accessibility and broad impact

Helo *et al.* (2013)

Deppisch *et al.* (2014)

De Gouvêa *et al.* (2019)

**Positive  
experimental  
result**

*hero or villain?*



# Simplified Model

We study a simplified model: **TeV-scale LNV** Peng et al. (2015)



$$\text{SM} + S : (\mathbf{1}, \mathbf{2})_{1/2} \text{ and } F : (\mathbf{1}, \mathbf{1})_0$$
$$m_F, m_S \sim \text{TeV}$$

Minimal interactions (*ingredients*) to address the key questions:

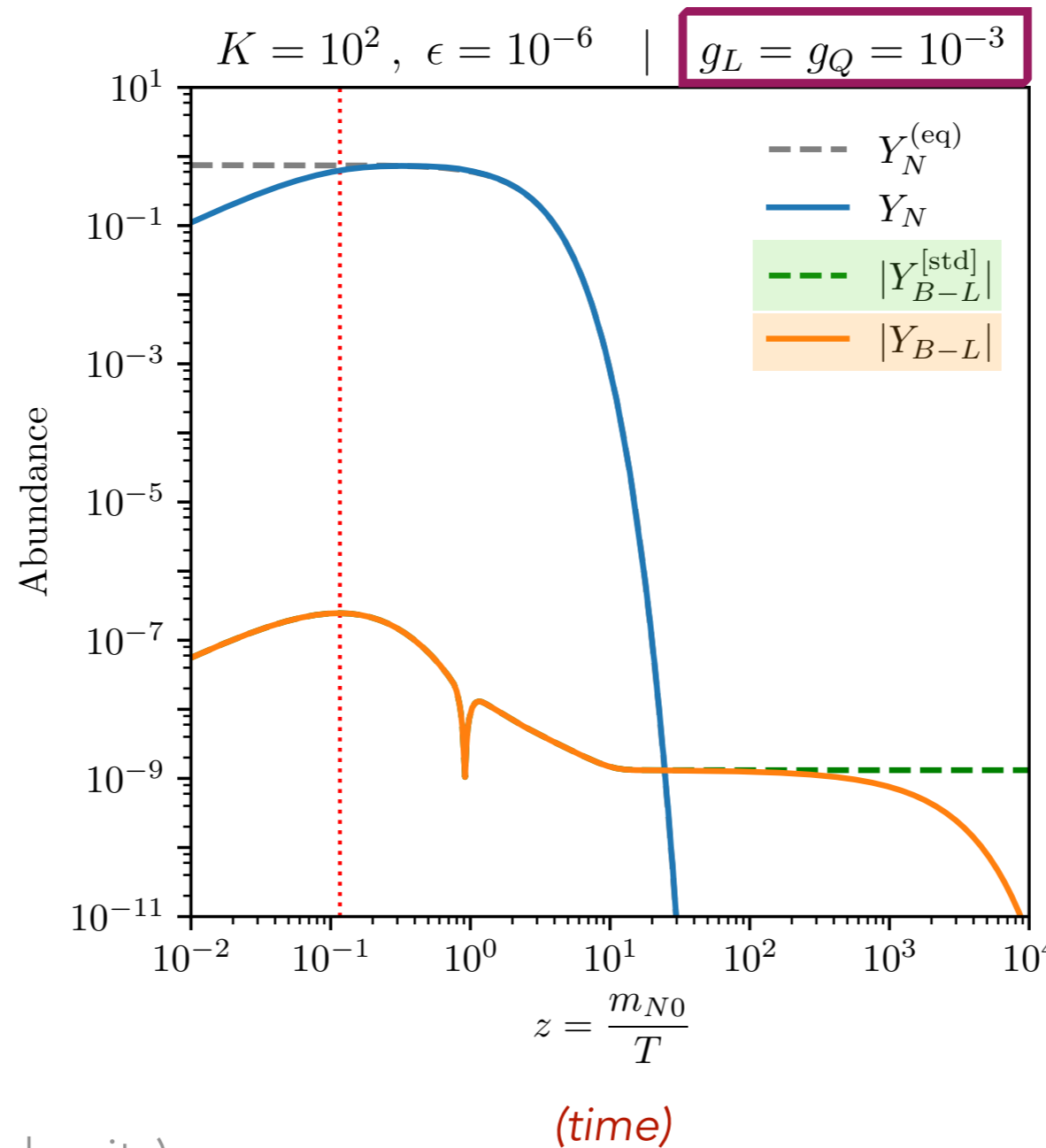
$$\mathcal{L}_{\text{SimpMod}} = g_Q \bar{Q} S d_R + g_L \bar{L} \tilde{S} F - m_S^2 S^\dagger S - \frac{m_F}{2} \bar{F}^c F + \lambda_{HS} (S^\dagger H)^2$$

$$\tilde{X} \equiv \epsilon X^*, \quad \epsilon^{12} = +1 \quad \left( SU(3)_C, SU(2)_L \right)_{U(1)_Y}$$

# TeV-scale LNV effects

- What if we include accessible TeV-scale effects? **Simplified Model**

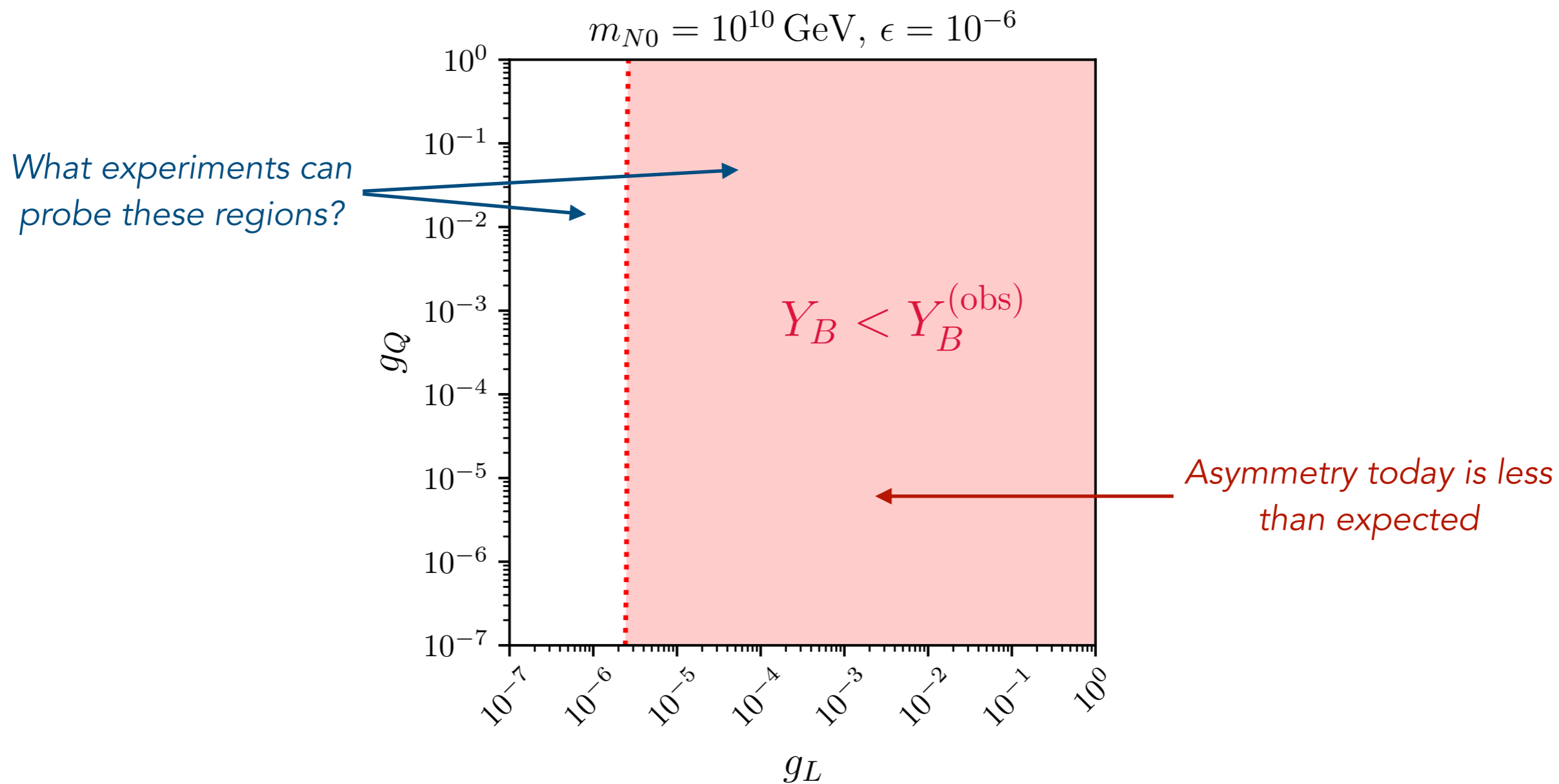
$|Y_{B-L}| \mapsto |Y_B|$   
via Sphaleron



$Y_X = n_X/s$  (normalized number density)

# TeV-scale LNV effects

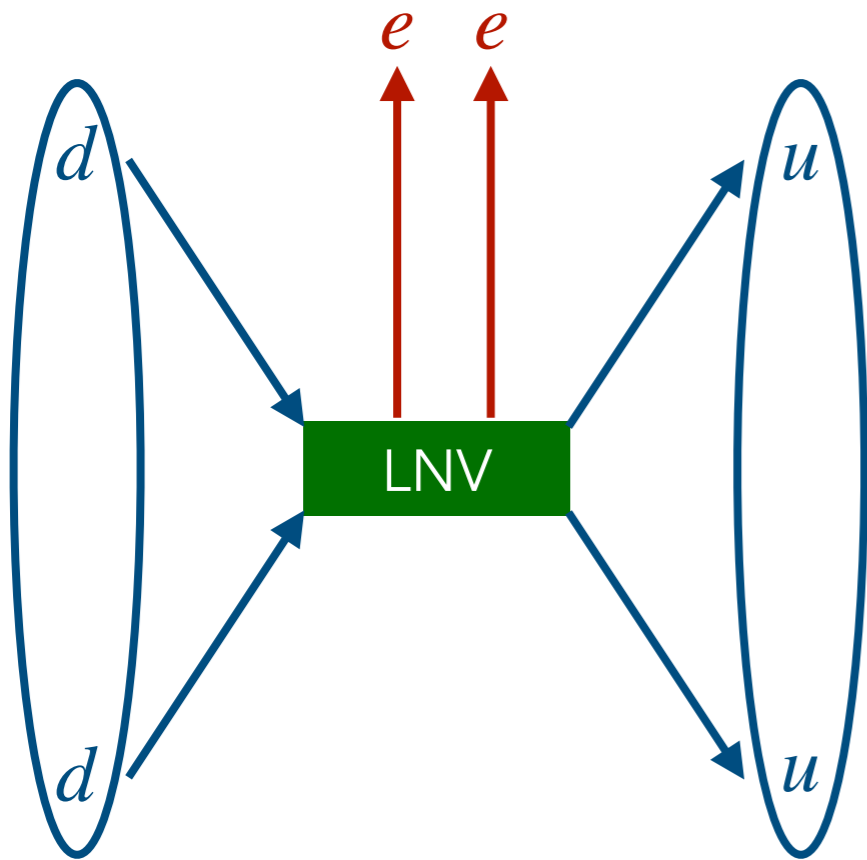
- What if we include accessible TeV-scale effects? **Simplified Model**





# TeV-scale LNV: $0\nu\beta\beta$ -decay & Colliders

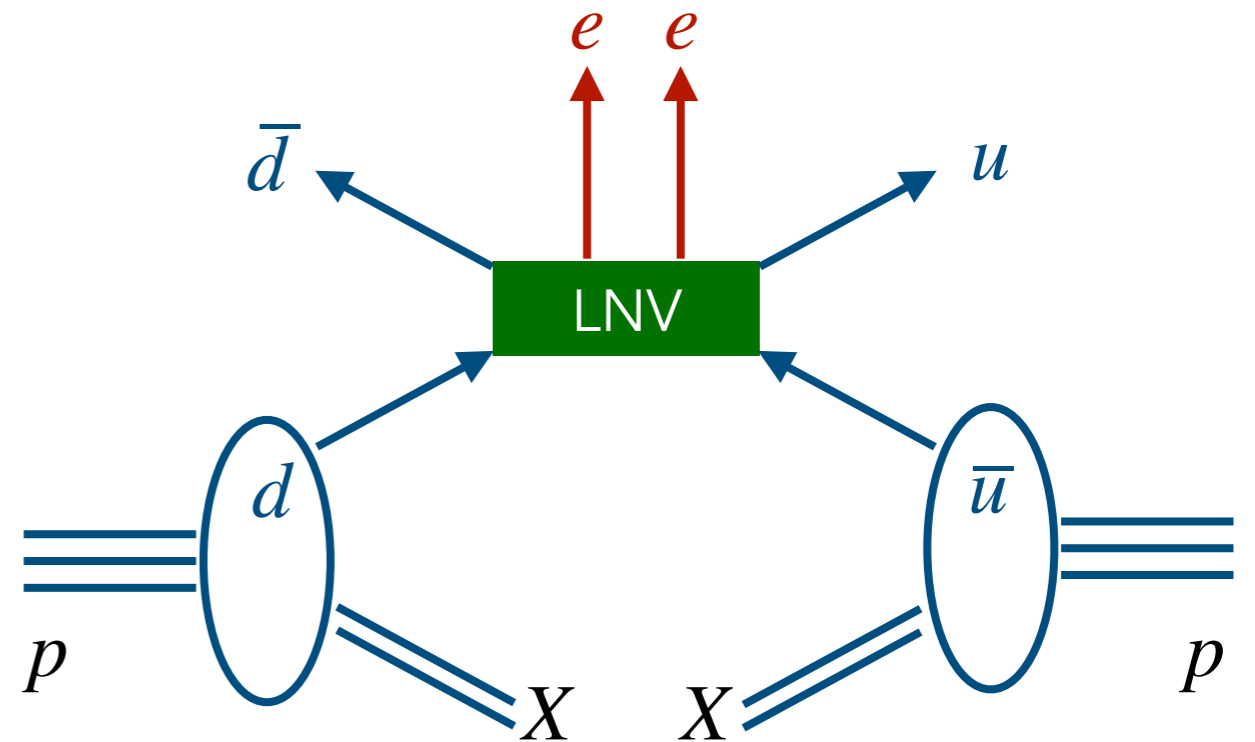
$0\nu\beta\beta$ -decay



$A(Z, N)$

$A(Z + 2, N - 2)$

$pp$  collisions

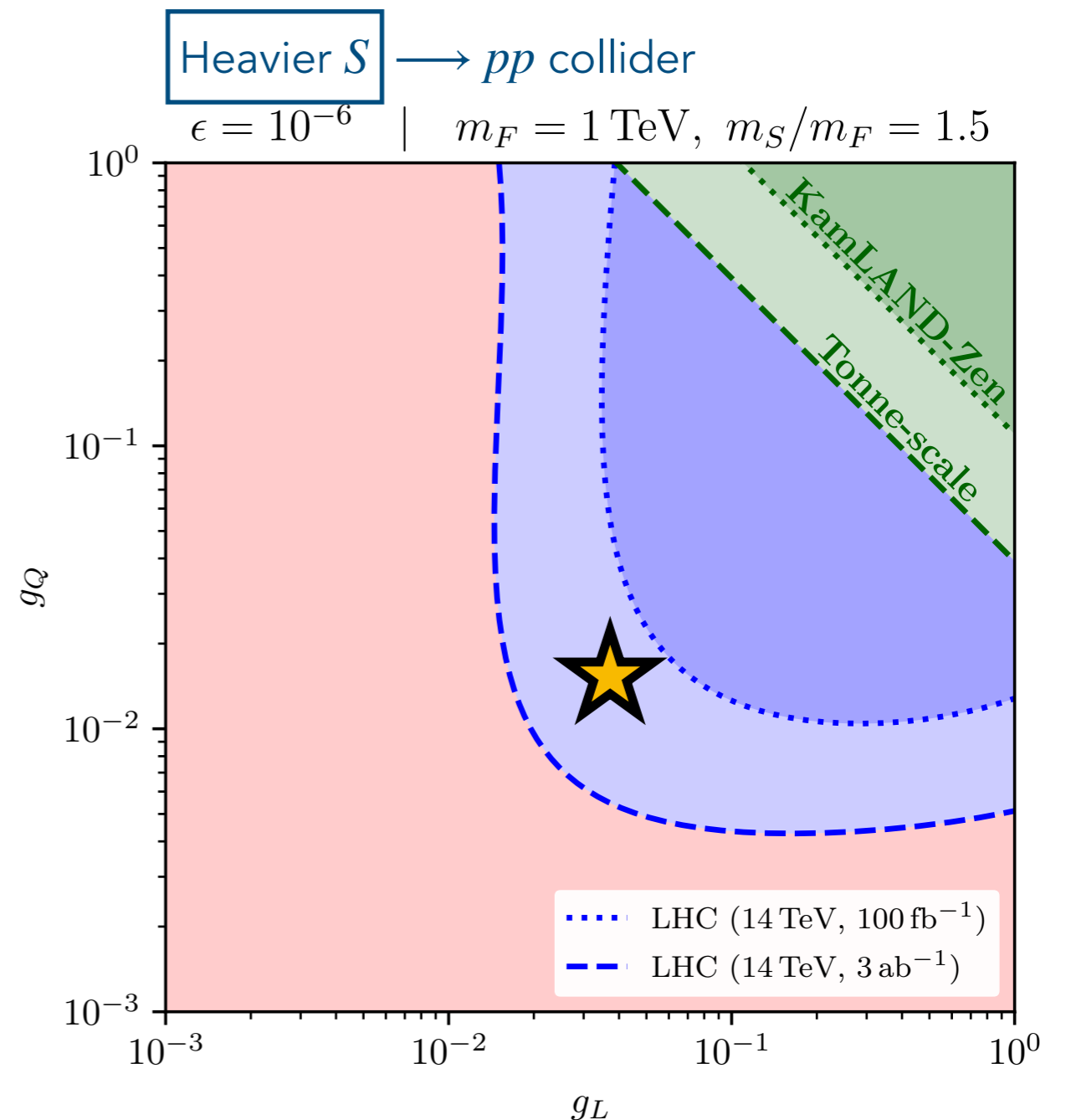
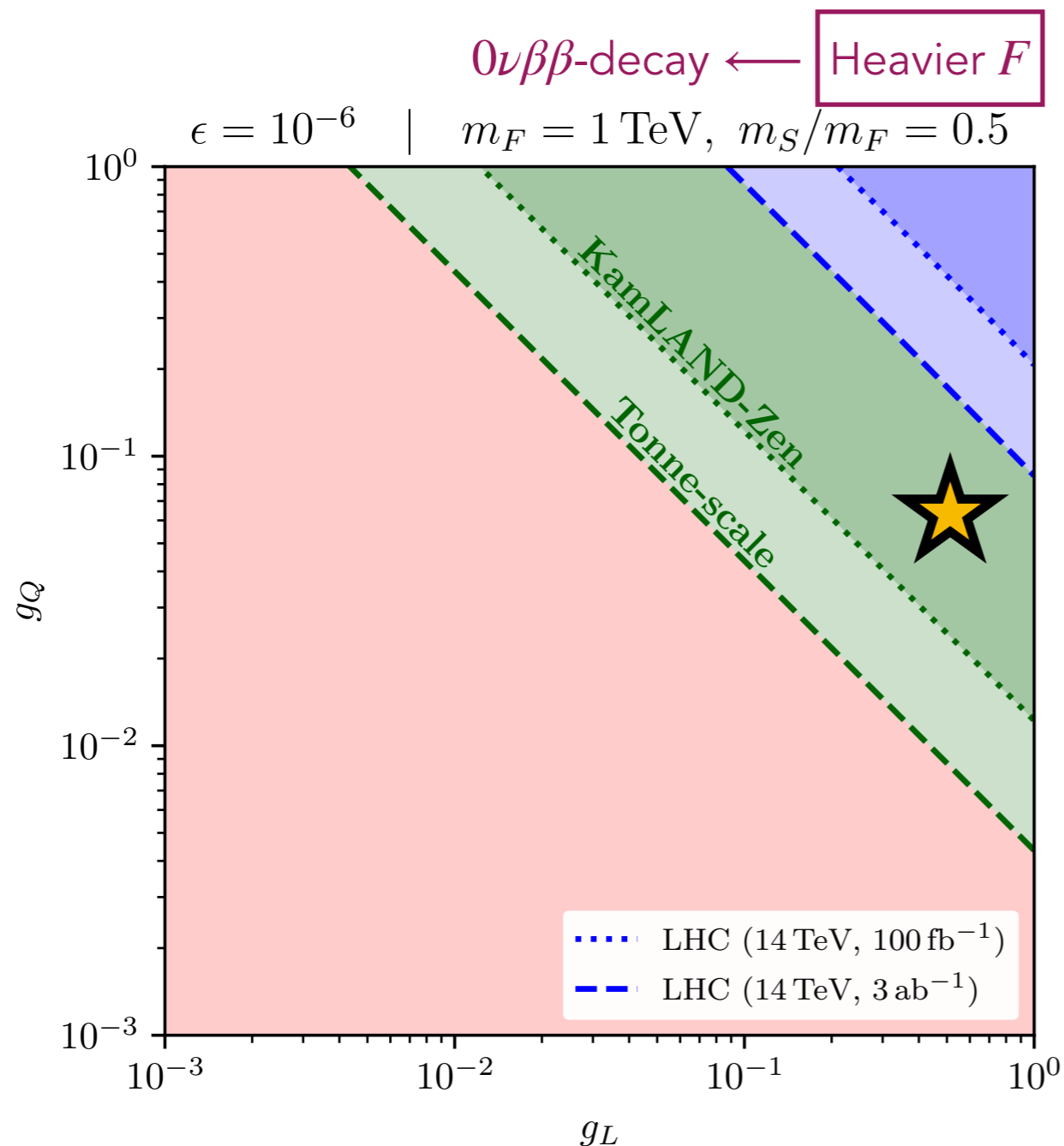


(Thanks! M. Ramsey-Musolf)



# TeV-scale LNV effects: Results

- Impact of  $0\nu\beta\beta$ -decay **and/or** collider searches on leptogenesis
- Both searches provide complementary sensitivity!



# Takeaways

- The observation of an LNV signal at the Large Hadron Collider (LHC) and/or a future 100 TeV  $pp$  collider would preclude the viability of standard thermal leptogenesis
- The observation of  $0\nu\beta\beta$ -decay could also rule out the standard leptogenesis paradigm, assuming the  $0\nu\beta\beta$ -decay amplitude is dominated by the TeV-scale LNV mechanism
- The relative reaches of  $0\nu\beta\beta$ -decay and collider LNV searches depends decisively on the new particle spectrum — a feature not readily seen within the previously used pure EFT approach



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Backup Slides

# Baryon asymmetry: ingredients

## Sakharov conditions (1967)



### 1. $\mathcal{B}$ violation

number of  $B \neq$  number of  $\bar{B}$

SM Sphalerons



### 2. $C$ & $CP$ violation

$\Gamma(A + B \rightarrow C) \neq \Gamma(\bar{A} + \bar{B} \rightarrow \bar{C})$

SM CKM CPV



### 3. Non-equilibrium dynamics\*

Prevent washout by inverse processes

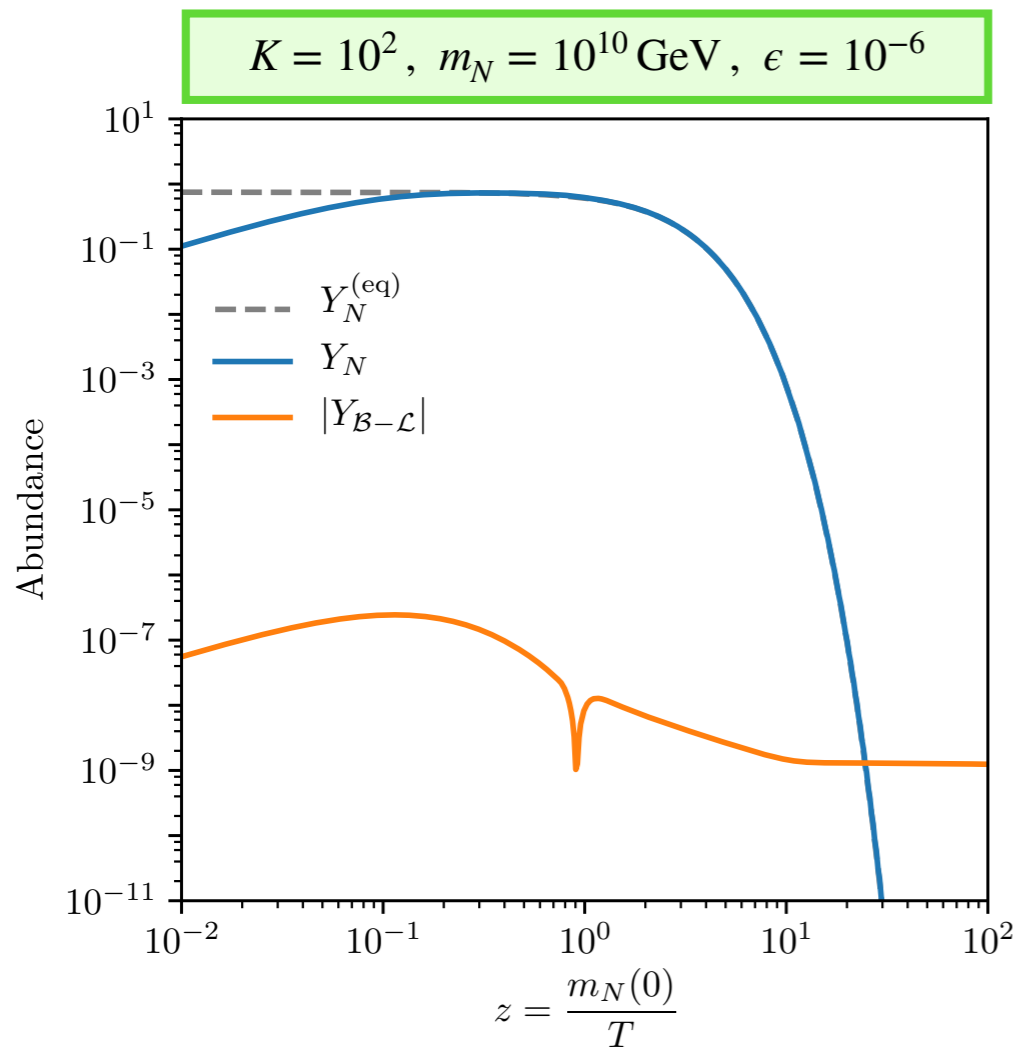
SM EWPT



***New physics!***

# Boltzmann Equations

The dynamics is described by Boltzmann equations:



$$\frac{dY_N}{dz} = - \left( \overset{\text{Decay}}{\downarrow} \boxed{D} + \overset{\text{Scattering}}{\downarrow} \boxed{S} \right) \left( Y_N - Y_N^{(\text{eq})} \right)$$

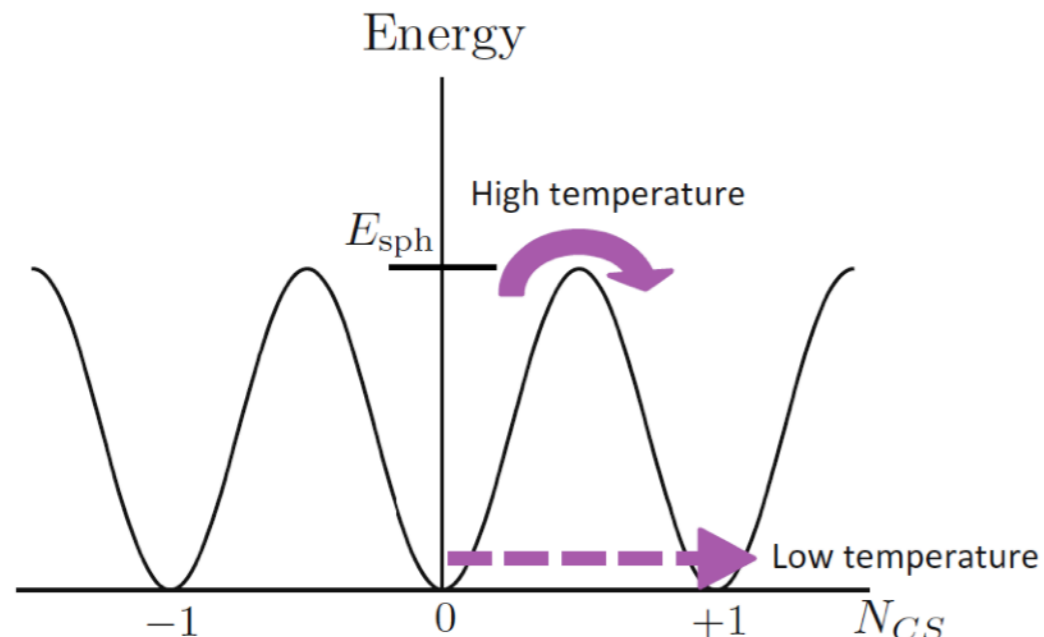
$$\frac{dY_{\mathcal{B}-\mathcal{L}}}{dz} = - \underset{\substack{\uparrow \\ \text{CPV source} \\ \text{(decay)}}}{\boxed{\epsilon D}} \left( Y_N - Y_N^{(\text{eq})} \right) - \underset{\substack{\uparrow \\ \text{Washout} \\ \text{(inverse decay, } \Delta\mathcal{L} = 1, 2)}}{\boxed{W}} Y_{\mathcal{B}-\mathcal{L}}$$



$Y_X = n_X/s$  (normalized number density)

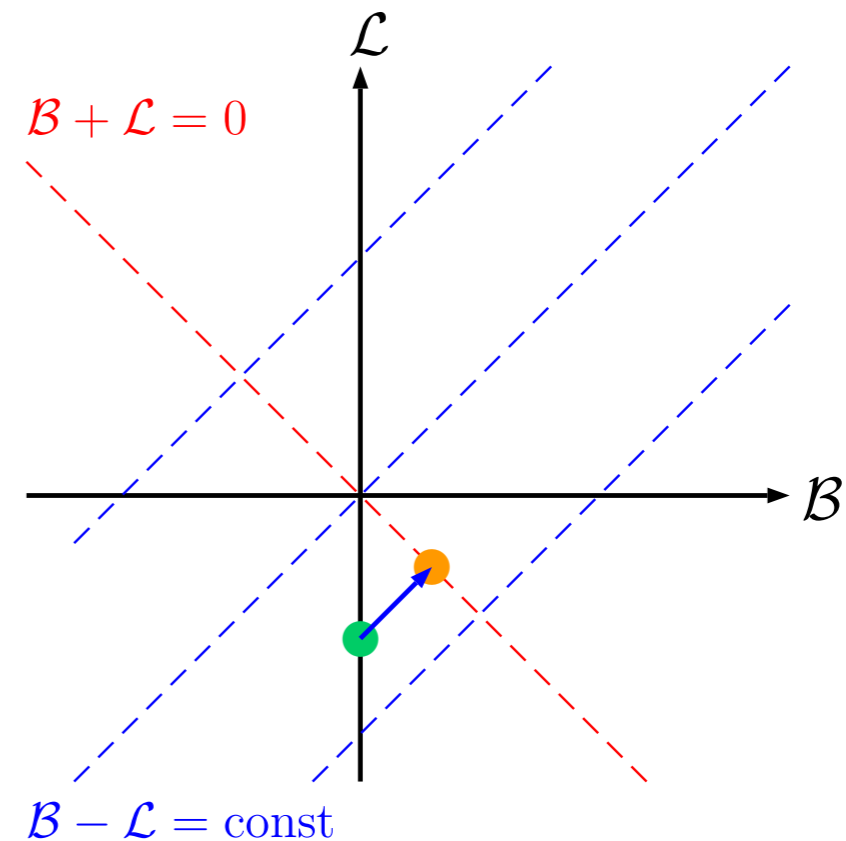
# Baryon asymmetry: EW Sphalerons

- In the SM, **non-perturbative effects** can violate  $\mathcal{B}$  and  $\mathcal{L}$ . This violation is such that  $\mathcal{B} - \mathcal{L}$  is conserved (while  $\mathcal{B} + \mathcal{L}$  is violated)



Fuyuto (Springer Theses, 2018)

$$\Delta\mathcal{B} = \Delta\mathcal{L} = 3N_{CS}$$



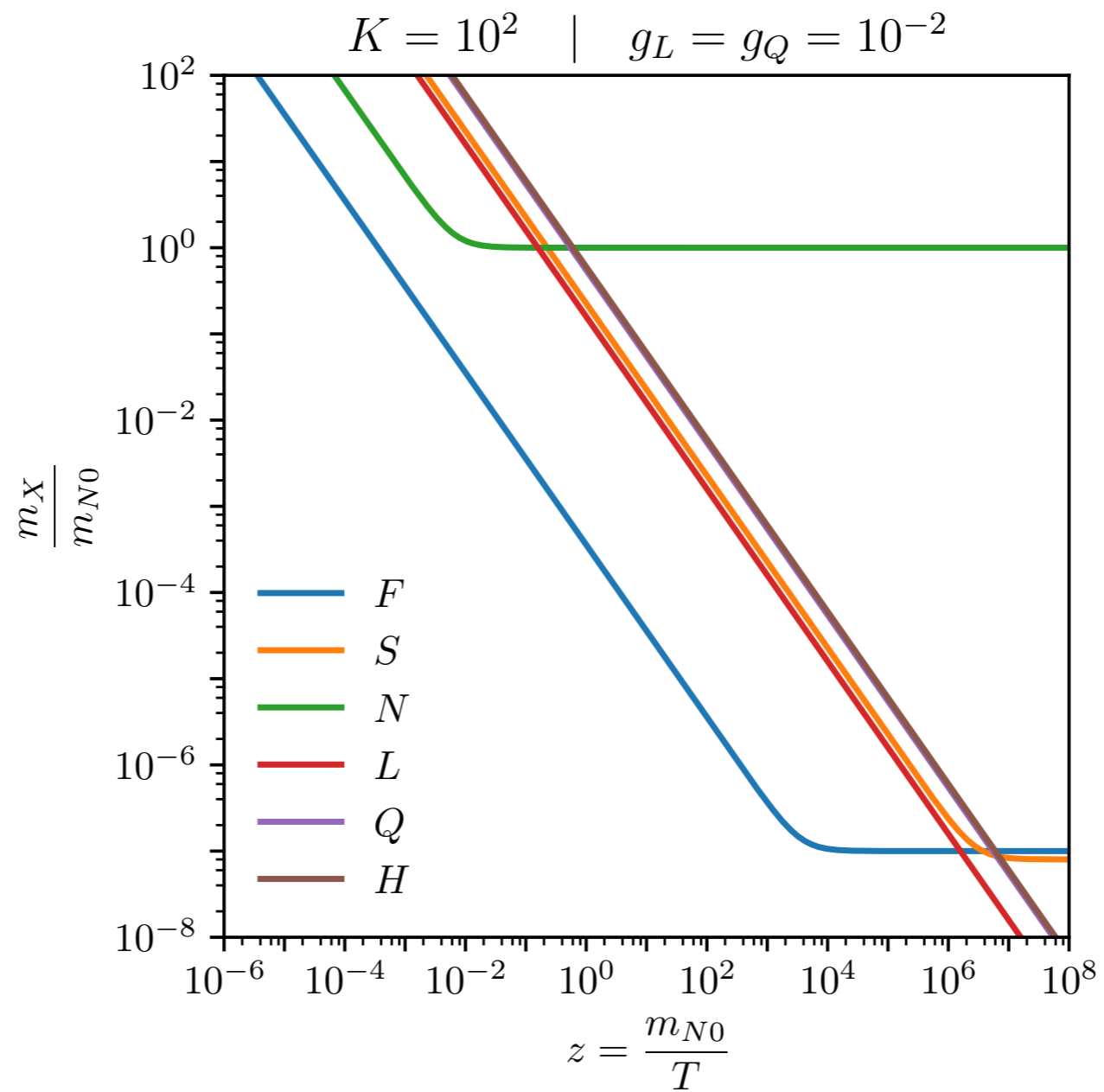
Sphaleron transitions:

- Conserve  $\mathcal{B} - \mathcal{L}$
- Wash  $\mathcal{B} + \mathcal{L}$  out

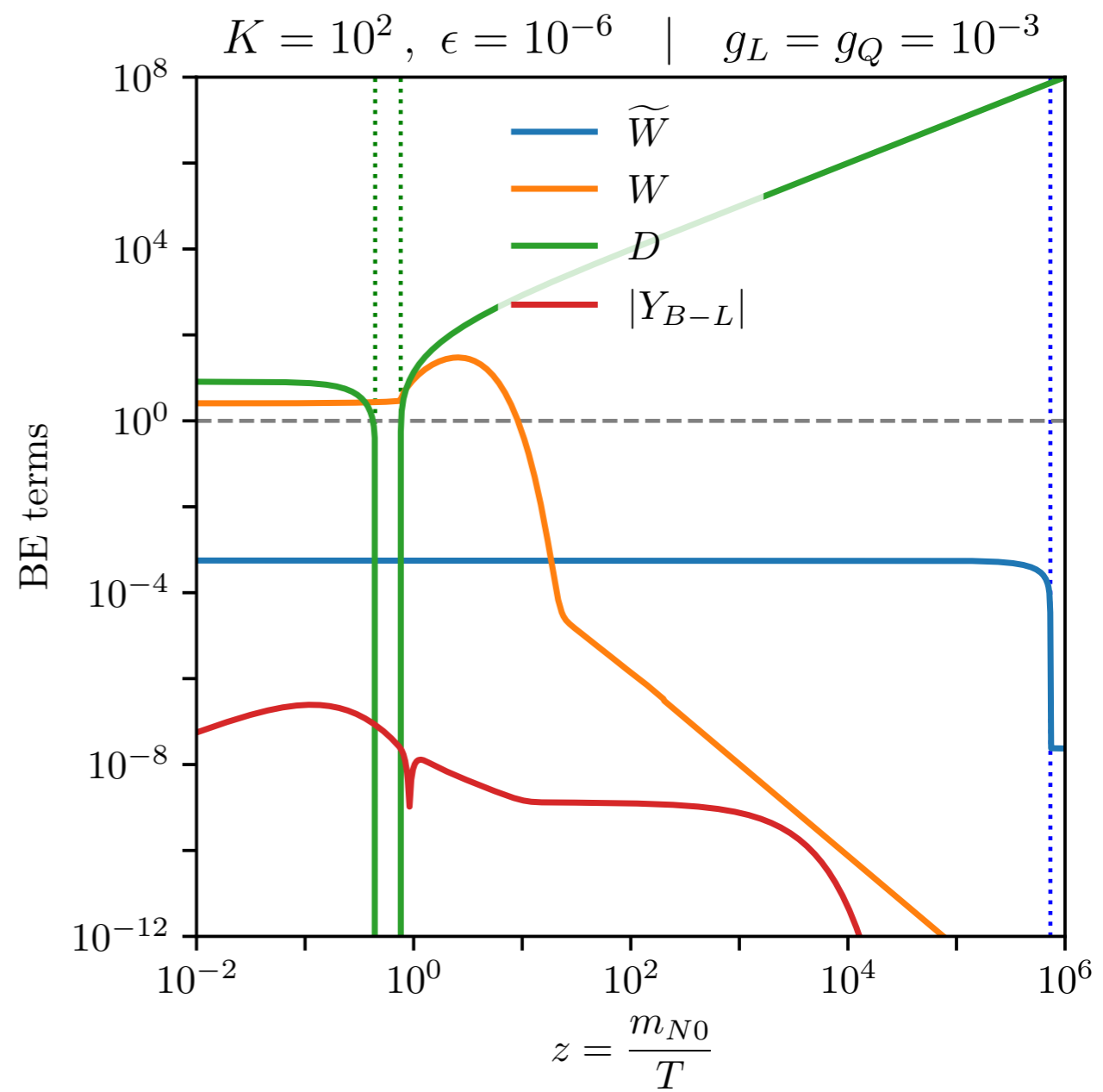
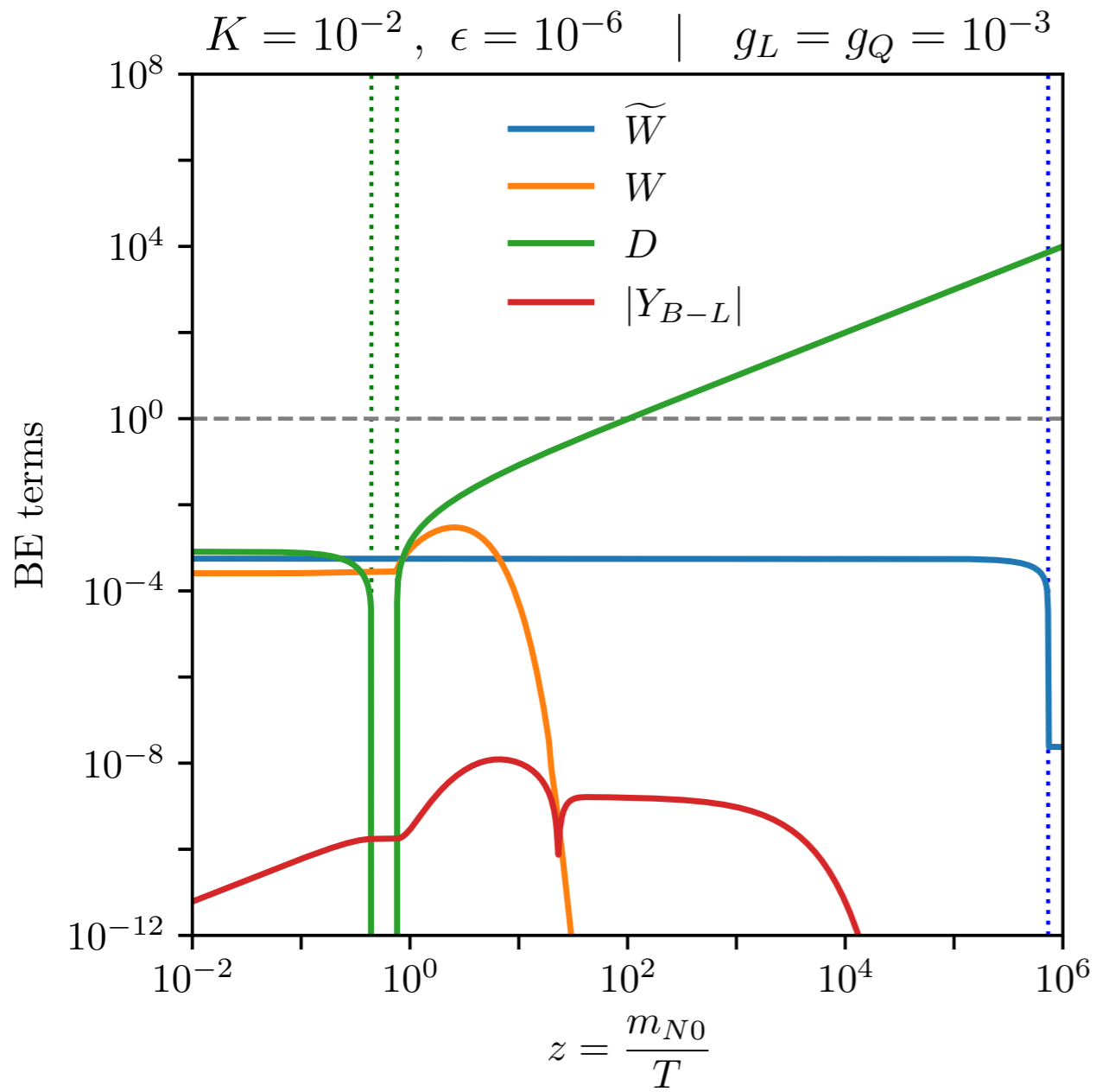
$\mathcal{L}$  asymmetry is reprocessed into  $\mathcal{B}$  asymmetry



# Leptogenesis with Simplified Model - I

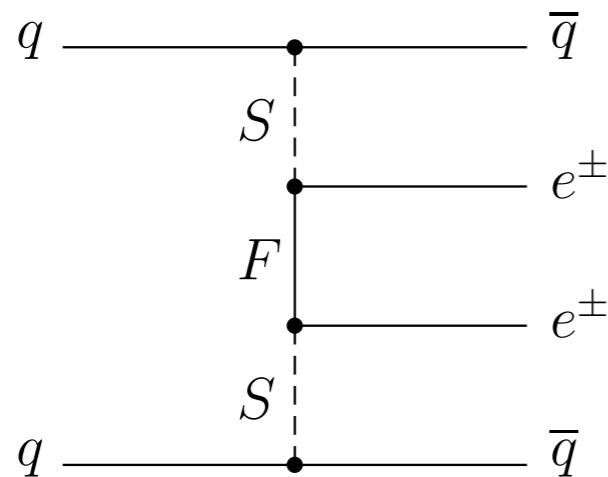


# Leptogenesis with Simplified Model - II

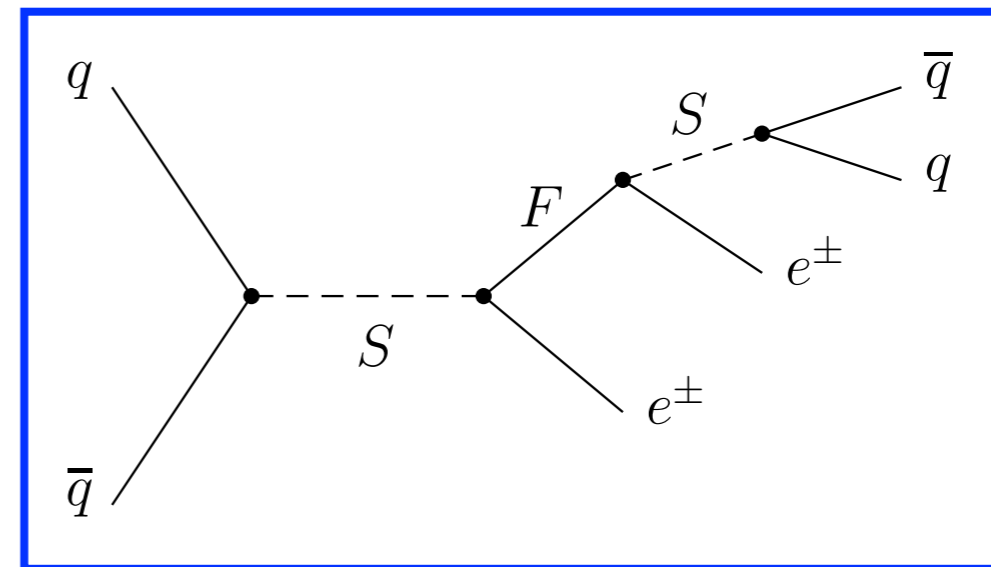


# LNV searches at colliders

- **Signal:** two same-sign electrons and two jets,  $pp \rightarrow jje^\pm e^\pm$



(a)  $t$ -channel



Dominant contribution

(b)  $s$ -channel

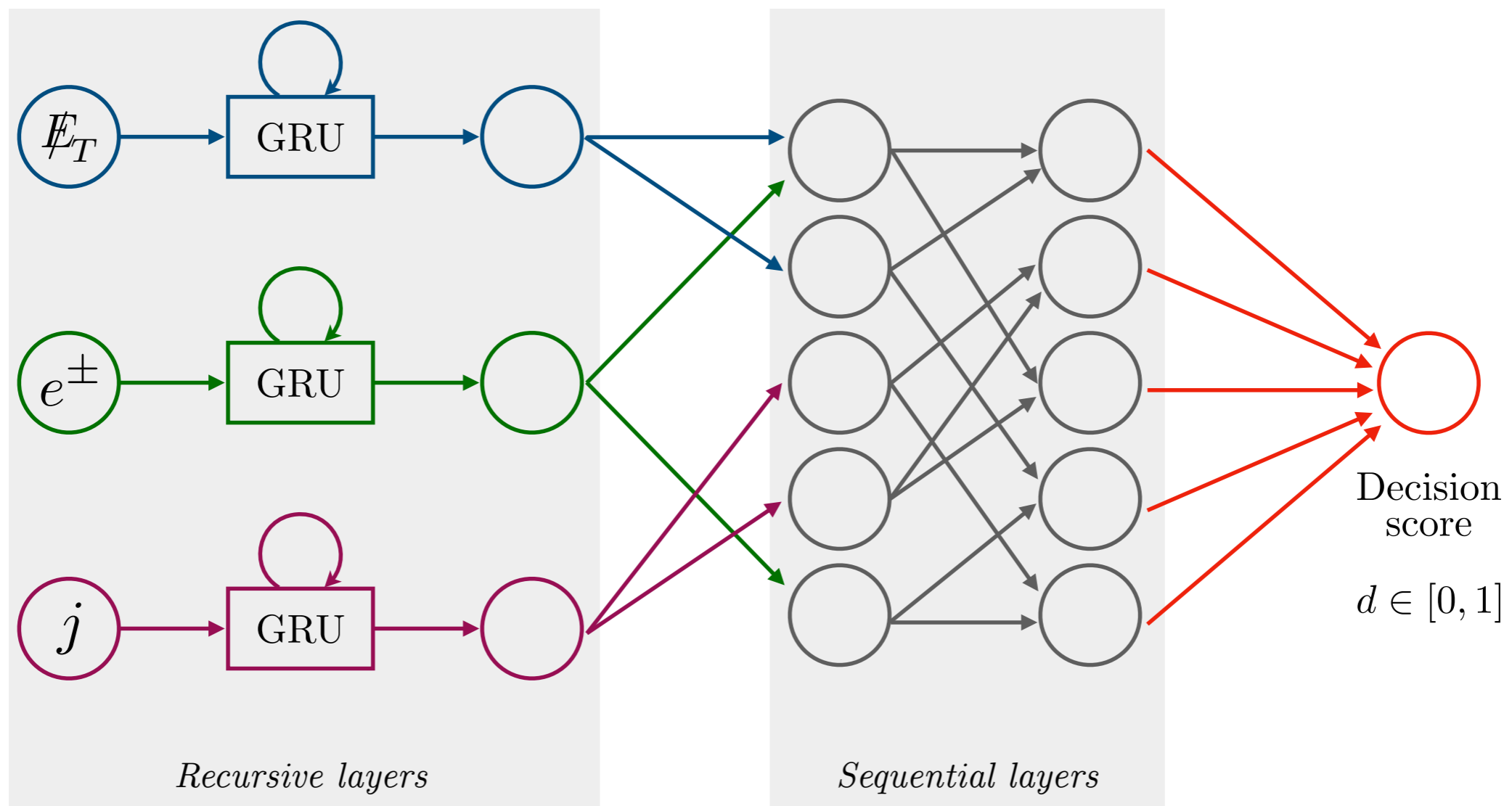
- **Background:**

- SM processes with 2SS leptons (e.g.,  $jjWW$ ,  $t\bar{t}W$ ,  $t\bar{t}Z$ )
- Charge misidentification (e.g.,  $t\bar{t}$ ,  $Z/\gamma^*$ )
- Jet-fakes (e.g.,  $t\bar{t}$ ,  $W + 3j$ )

Dominant contributions

# RNN for signal/background discrimination

- Madgraph + Pythia 8 + Delphes for MC simulation
- Neural Networks for signal/background discrimination



# Effective operators - I

Prézeau, Ramsey-Musolf, Vogel (2003)

$d = 9$  effective operators ( $\mathcal{O}_i \propto \bar{u}\bar{u}dd\bar{e}e^c$ ) contributing to  $0\nu\beta\beta$ -decay:

$$\mathcal{L}_{0\nu\beta\beta} = \frac{1}{\Lambda^5} \left\{ \begin{aligned} & \left( o_1 \mathcal{O}_{1+}^{++} + o_2 \mathcal{O}_{2+}^{++} + o_3 \mathcal{O}_{2-}^{++} + o_4 \mathcal{O}_{3+}^{++} + o_5 \mathcal{O}_{3-}^{++} \right) \bar{e}e^c + \\ & \left( o_6 \mathcal{O}_{1+}^{++} + o_7 \mathcal{O}_{2+}^{++} + o_8 \mathcal{O}_{2-}^{++} + o_9 \mathcal{O}_{3+}^{++} + o_{10} \mathcal{O}_{3-}^{++} \right) \bar{e}\gamma_5 e^c + \\ & \left( o_{11} \mathcal{O}_{4+}^{++,\mu} + o_{12} \mathcal{O}_{4-}^{++,\mu} + o_{13} \mathcal{O}_{5+}^{++,\mu} + o_{14} \mathcal{O}_{5-}^{++,\mu} \right) \bar{e}\gamma_\mu \gamma_5 e^c + \text{h.c.} \end{aligned} \right\}$$

$$\mathcal{O}_{1+}^{ab} = (\bar{q}_L \tau^a \gamma^\mu q_L) (\bar{q}_R \tau^b \gamma_\mu q_R) ,$$

$$\mathcal{O}_{2\pm}^{ab} = (\bar{q}_R \tau^a q_L) (\bar{q}_R \tau^b q_L) \pm (\bar{q}_L \tau^a q_R) (\bar{q}_L \tau^b q_R) ,$$

$$\mathcal{O}_{3\pm}^{ab} = (\bar{q}_L \tau^a \gamma^\mu q_L) (\bar{q}_L \tau^b \gamma_\mu q_L) \pm (\bar{q}_R \tau^a \gamma^\mu q_R) (\bar{q}_R \tau^b \gamma_\mu q_R) ,$$

$$\mathcal{O}_{4\pm}^{ab,\mu} = (\bar{q}_L \tau^a \gamma^\mu q_L \mp \bar{q}_R \tau^a \gamma^\mu q_R) (\bar{q}_L \tau^b q_R - \bar{q}_R \tau^b q_L) ,$$

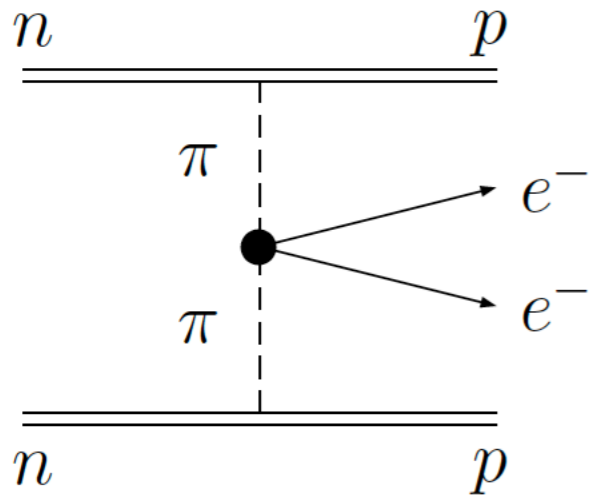
$$\mathcal{O}_{5\pm}^{ab,\mu} = (\bar{q}_L \tau^a \gamma^\mu q_L \pm \bar{q}_R \tau^a \gamma^\mu q_R) (\bar{q}_L \tau^b q_R + \bar{q}_R \tau^b q_L) .$$

Different combinations of  
4 quarks and 2 electrons

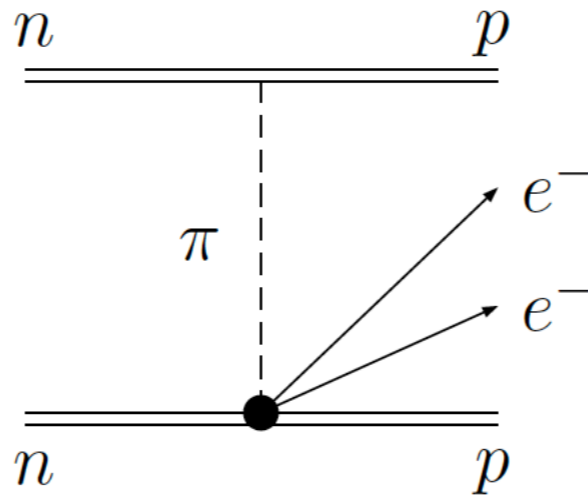
Operators  $\longleftrightarrow$  Models

# Effective operators - II

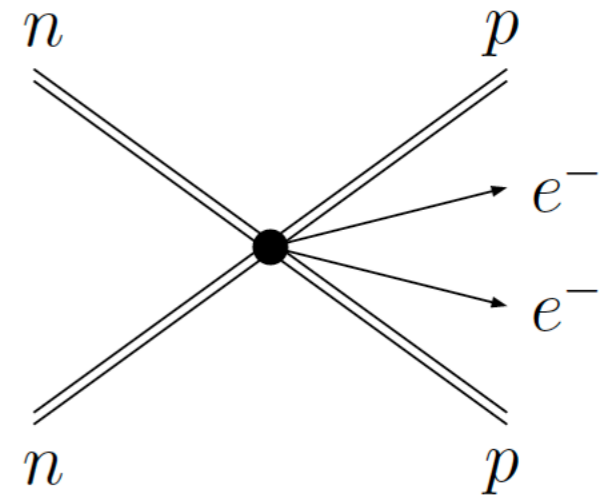
Prézeau, Ramsey-Musolf, Vogel (2003)



$\pi\pi ee$



$NN\pi ee$



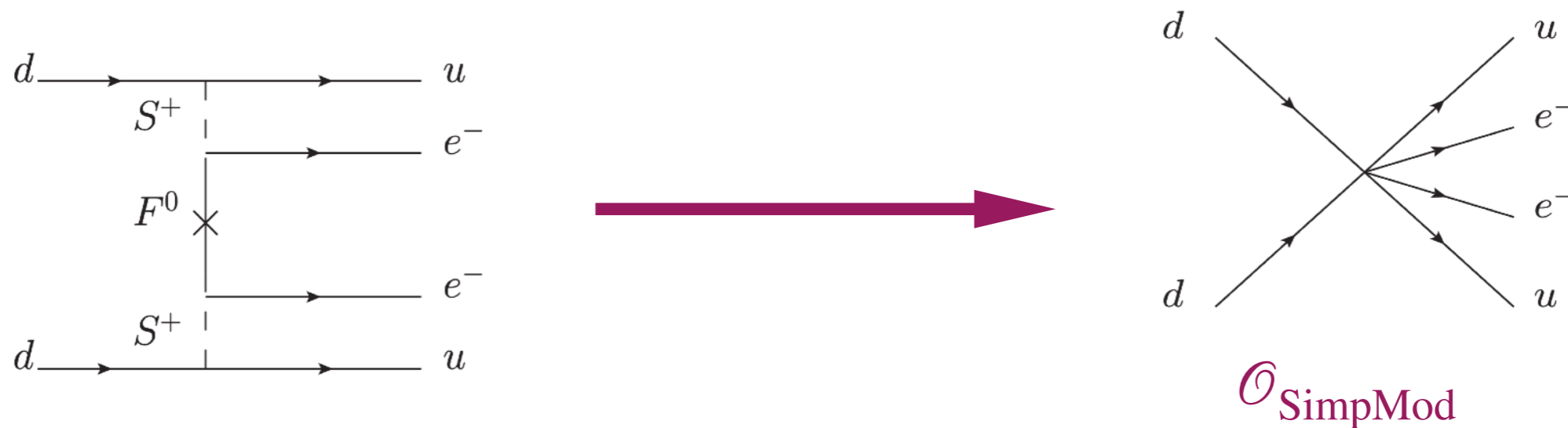
$NNNN ee$

Operator	$\mathcal{O}_{1+}^{\pm\pm}$	$\mathcal{O}_{2+}^{\pm\pm}$	$\mathcal{O}_{2-}^{\pm\pm}$	$\mathcal{O}_{3+}^{\pm\pm}$	$\mathcal{O}_{3-}^{\pm\pm}$	$\mathcal{O}_{4+}^{\pm\pm,\mu}$	$\mathcal{O}_{4-}^{\pm\pm,\mu}$	$\mathcal{O}_{5+}^{\pm\pm,\mu}$	$\mathcal{O}_{5-}^{\pm\pm,\mu}$
$\pi\pi ee$ LO	✓	✓	✗	✗	✗	✗	✗	✗	✗
$\pi\pi ee$ NNLO	✓	✓	✗	✓	✗	✗	✗	✗	✗
$NN\pi ee$ LO	✗	✗	✓	✗	✗	✓	✓	✓	✓
$NN\pi ee$ NLO	✗	✓	✗	✓	✗	✓	✓	✓	✓
$NNNN ee$ LO	✓	✓	✗	✓	✗	✓	✓	✓	✓

# Simplified models $\mapsto$ Effective operators - I

1. Obtain the  $0\nu\beta\beta$ -decay operator after integrating out heavy d.o.f:

$$\mathcal{O} \propto \bar{u}d\bar{u}d\bar{e}e^c$$



$$\mathcal{L}(q, e) = \frac{G_F^2}{\Lambda_{\beta\beta}} \sum_i C_i \mathcal{O}_i \bar{e}\Gamma_i e^c + \text{h.c.}$$

$\mathcal{O}_i$   $\rightarrow$   $dim - 9$  LNV effective operators

$C_i$   $\rightarrow$  Wilson coefficients

Prezeau et al. (2003)

Graesser (2016)

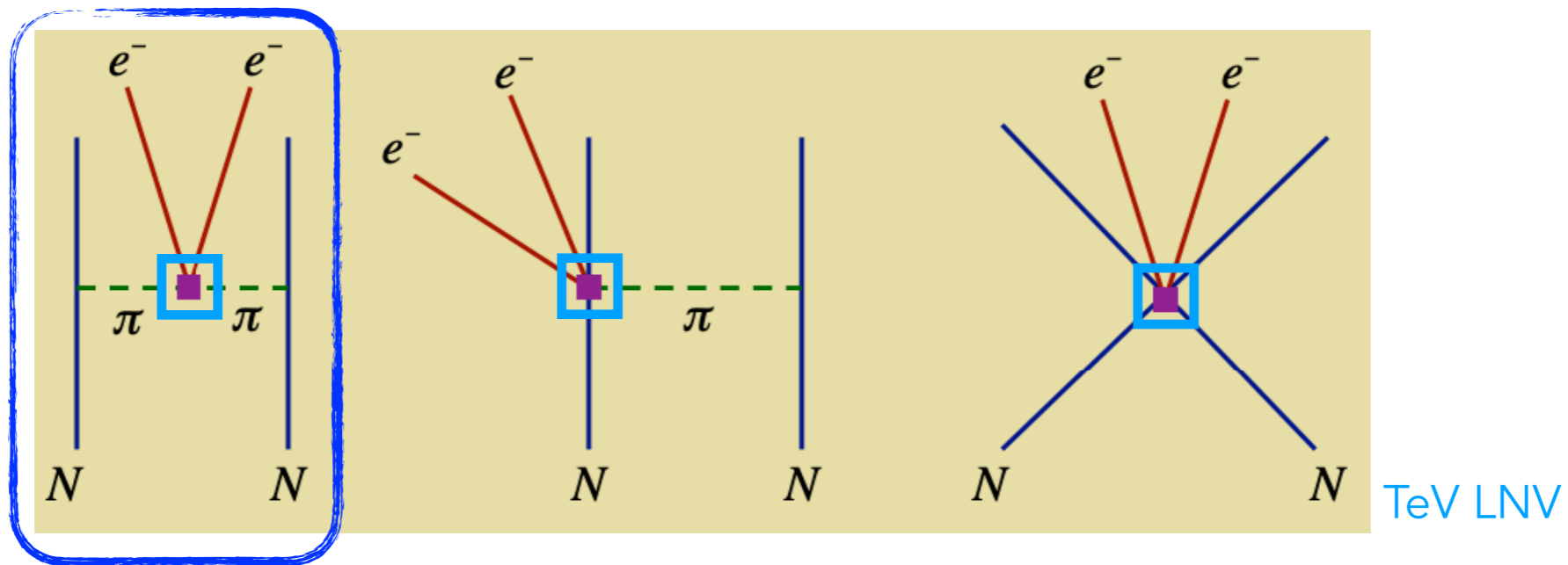
Cirigliano et al. (2017, 2018)



# Simplified models $\mapsto$ Effective operators - II

## 2. Match quark-electron operators onto two-nucleon operators

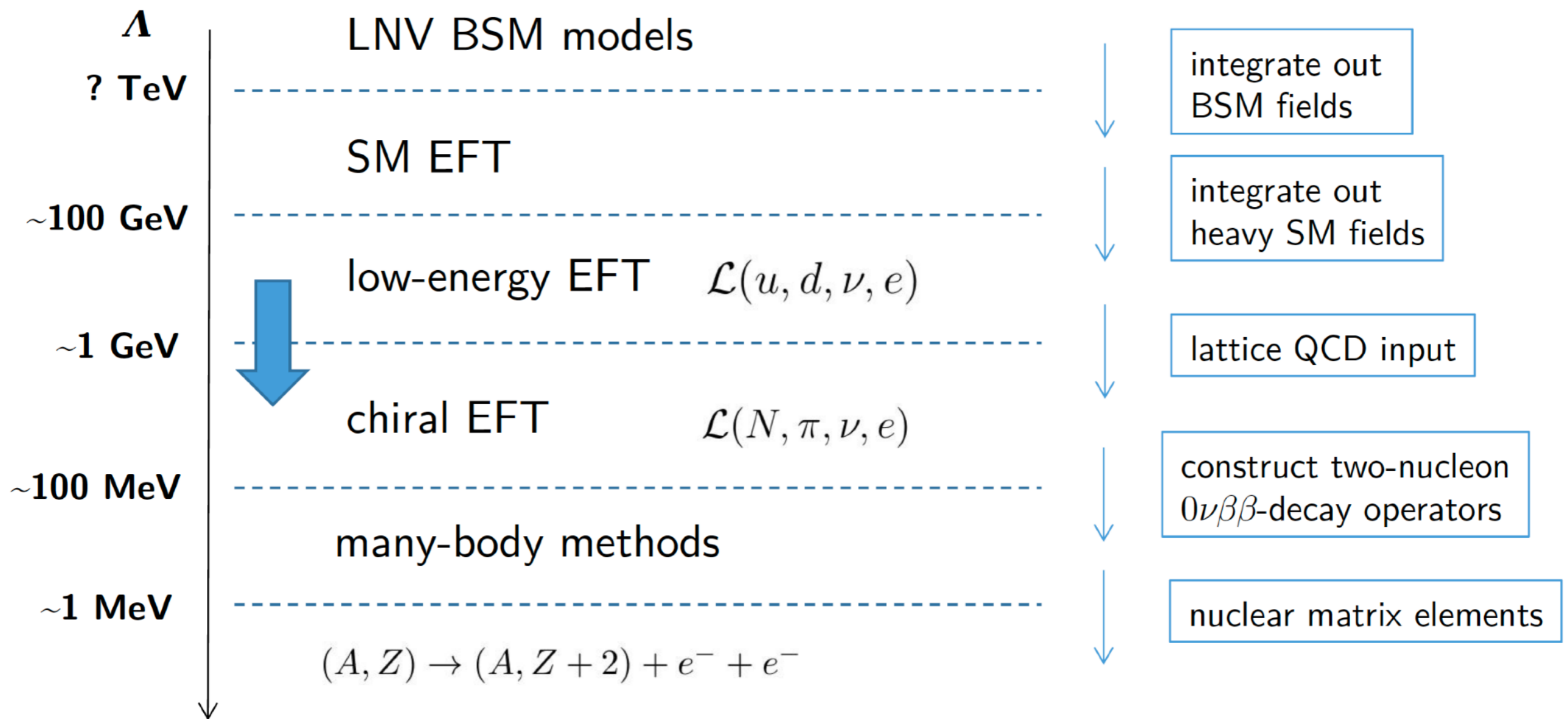
$$\mathcal{L}(q, e) = \frac{G_F^2}{\Lambda_{\beta\beta}^2} \sum_i C_i \mathcal{O}_i \bar{e} \Gamma_i e^c + \text{h.c.} \quad \text{Chiral symmetry: } \mathcal{O}_i \mapsto \mathcal{L}(\pi, N)$$



**Model O2:** LO + counterterm Cirigliano et al. (2018)

(Thanks! M. Ramsey-Musolf)

# Simplified models $\mapsto$ Effective operators - III

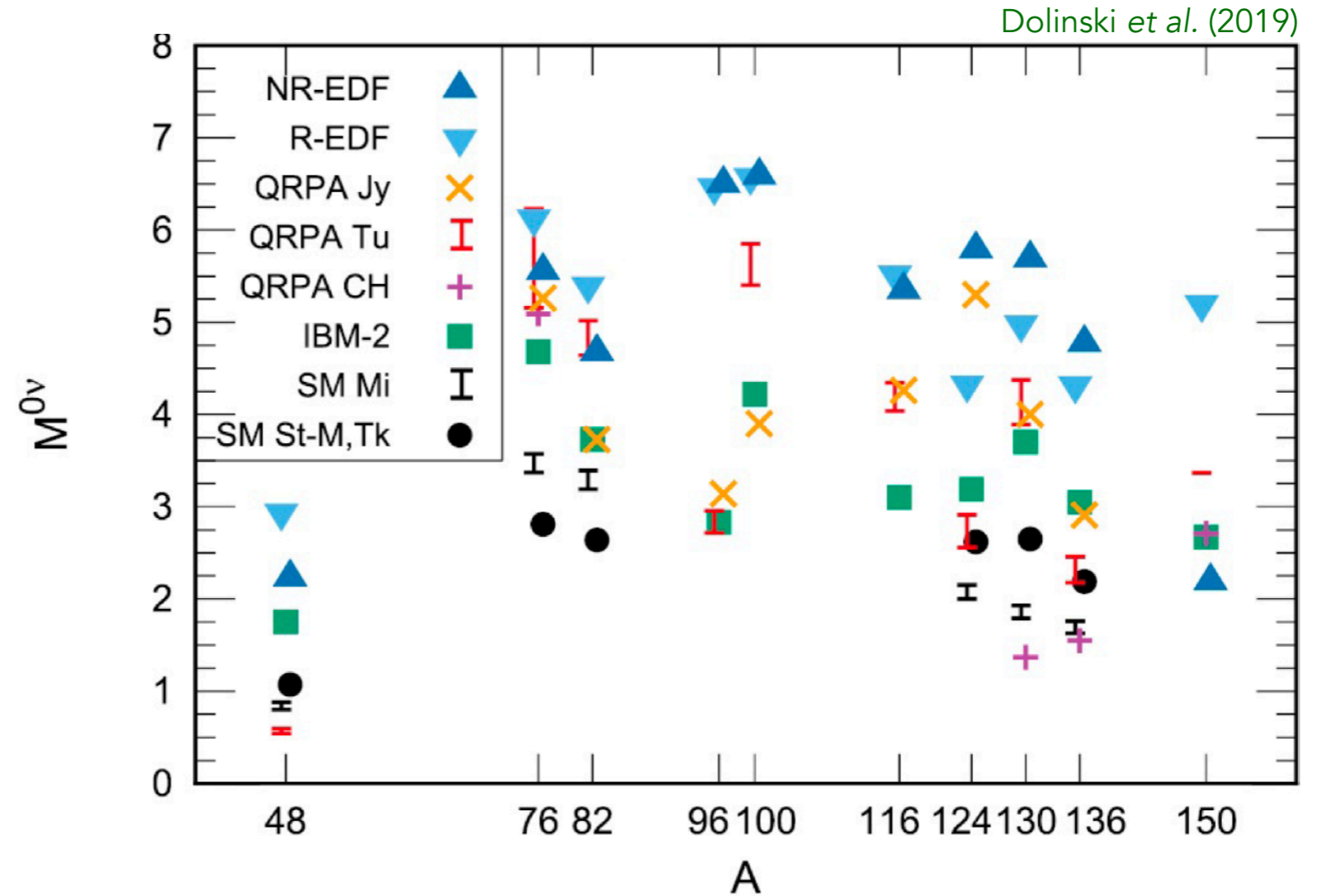
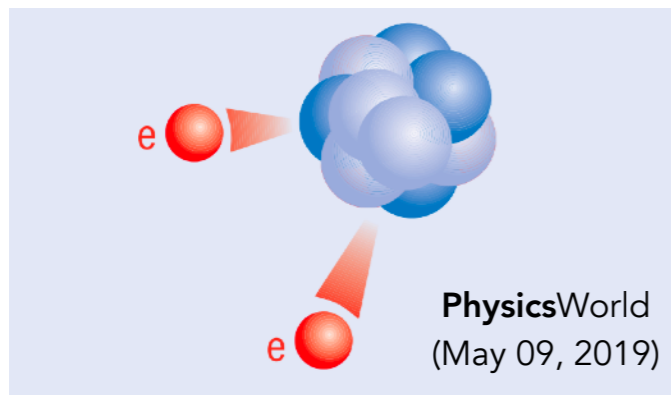


(Thanks! G. Li)

# Simplified models $\mapsto$ Effective operators - III

## 3. Using nuclear physics, calculate the half-life

$$T_{1/2}^{0\nu}(g_L, g_Q, m_F, m_S)$$



### Computing methods:

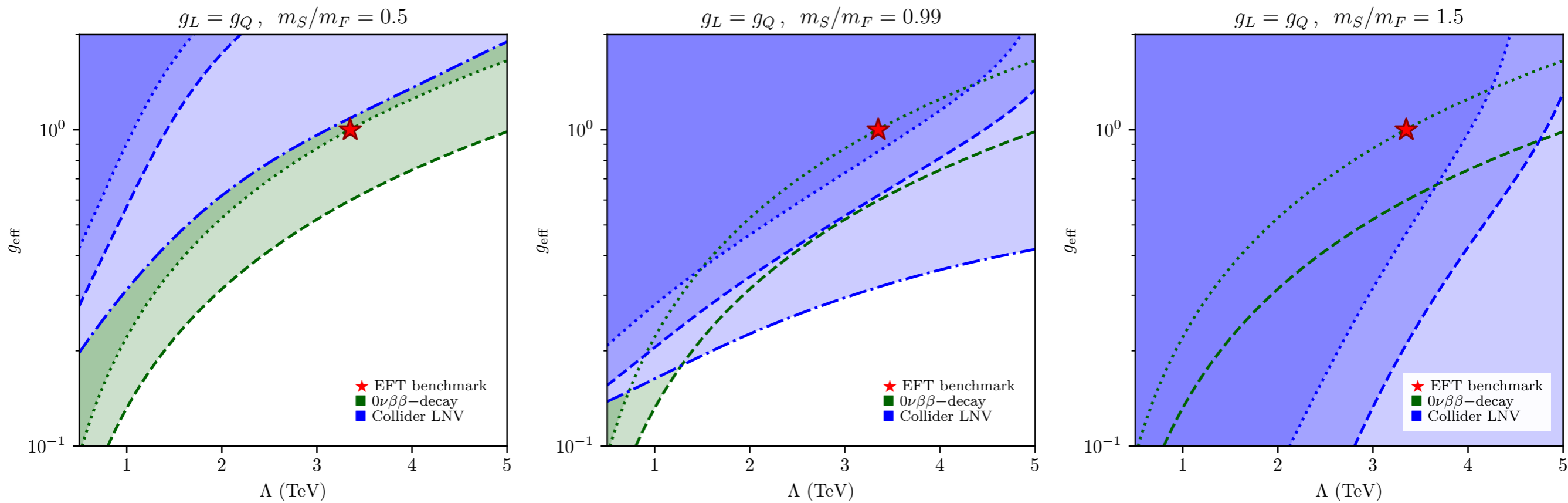
EDF: energy-density functional

QRPA: quasi-particle random phase approximation

SM: shell model

IBM: interacting boson model

# Comparison with EFT framework



$$\Lambda = (m_S^4 m_F)^{1/5}, \quad g_{\text{eff}} = g_L = g_Q$$