

# Gravitational Waves from an Inflation Triggered First-Order Phase Transition

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Based on work in collaboration with Haipeng An,  
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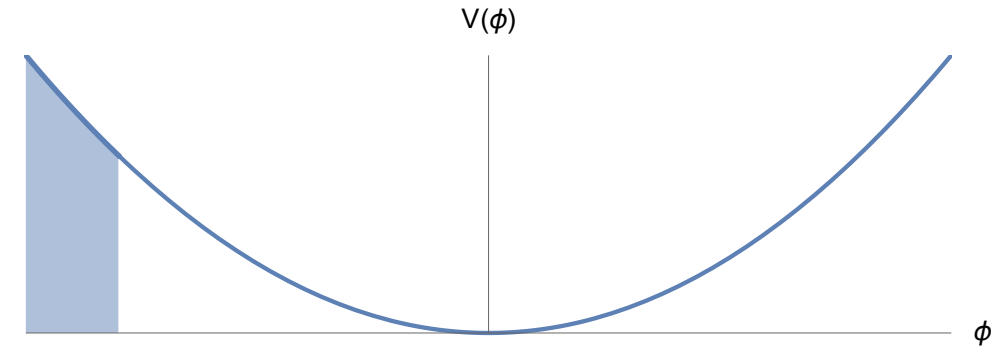
PHENO 2022, May 10<sup>th</sup>

# GW as a tool to probe the Early Universe

- Stochastic gravitational waves carry the information of its origin and the evolution history of the universe.
- It may be detected by the ground-based (LIGO, Virgo, etc.), space-based interferometers (LISA, BBO, etc.), radio telescope (SKA, etc.) or CMB in different frequency ranges.
- One interesting source is the strong first order phase transition during inflation.

# Phase Transition induced by spectator field during inflation

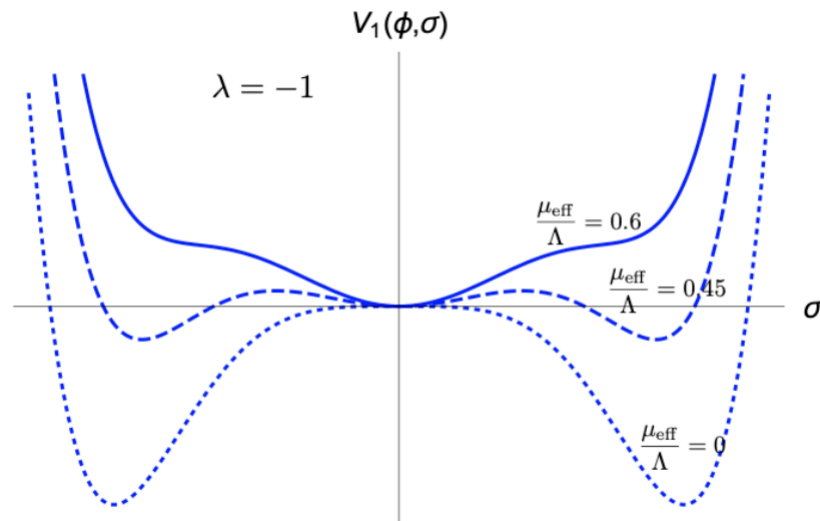
- The inflaton field would roll over a wide range during the inflation era.
- If there is spectator sector weakly coupled to inflaton, the variation of inflaton value would lead to some significant phenomena in the spectator sector.



Measured by  
CMB and LSS

No measurement

## Spectator sector



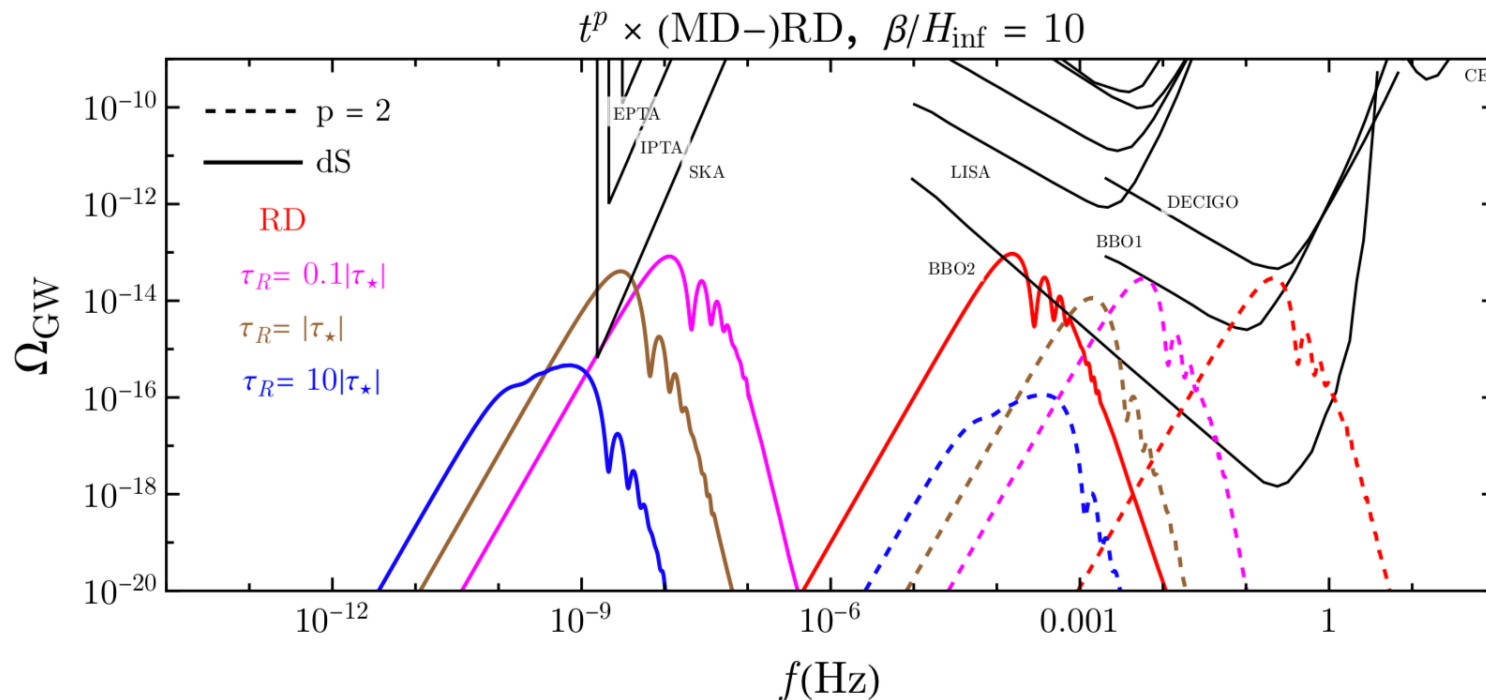
$$V_1(\phi, \sigma) = -\frac{1}{2}(\mu^2 - c^2\phi^2)\sigma^2 + \frac{\lambda}{4}\sigma^4 + \frac{1}{8\Lambda^2}\sigma^6$$

- The mass or coupling in the spectator sector may be changed due to the large excursion of the inflaton.
- It may induce

**Strong first order phase transition in spectator sector.**

# Strong First Order Phase Transition

- Bubble of true vacuum would be nucleated and expand if its size is larger than critical value.
- If the supercritical size bubbles nucleation rate  $\beta \gg H$ , the first order phase transition would complete
- Due to inflation, sound waves and magnetohydrodynamic turbulence can be ignored, we focus on bubble collision.

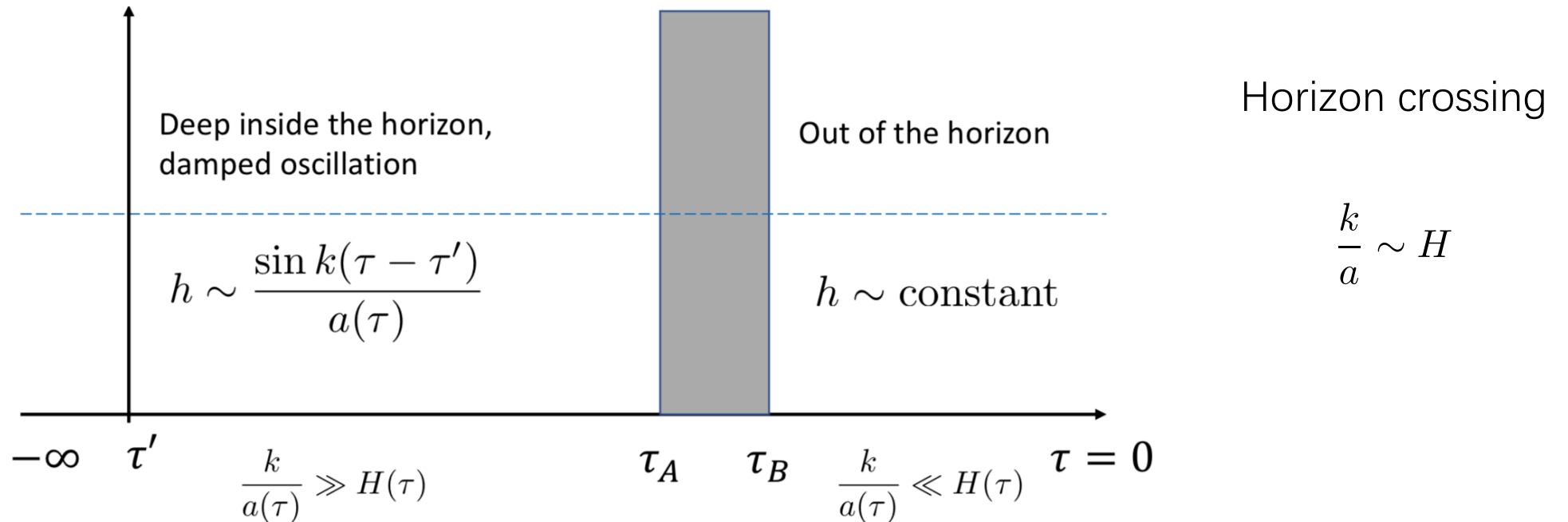


**Multiple peaks and oscillatory pattern!**

# GW is not diluted away during inflation.

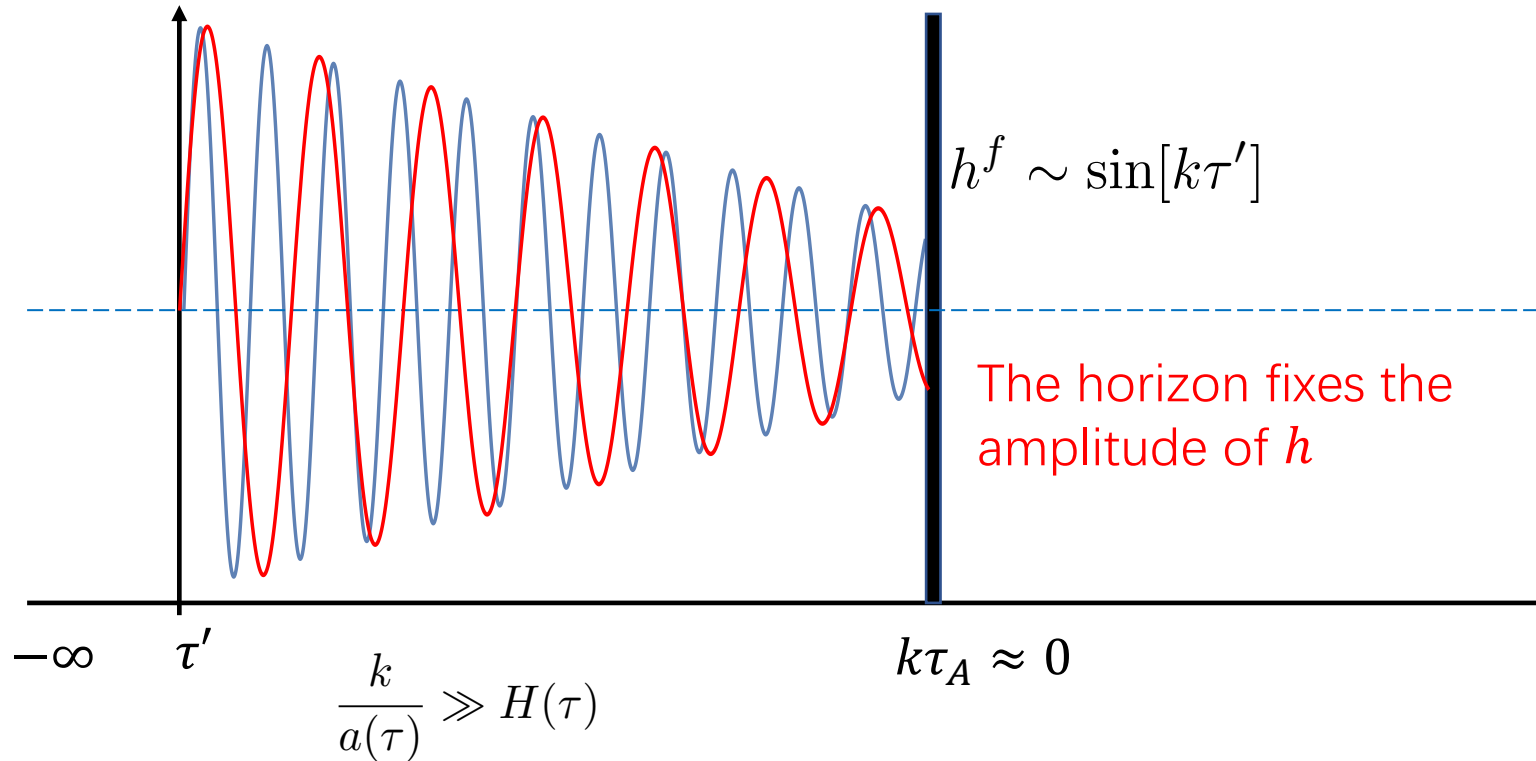
Solving the Green Function

$$h''_{ij} + \frac{2a'}{a} h'_{ij} - \nabla^2 h_{ij} = \frac{16\pi G_N}{a(\tau_*)} T_{ij}^{(0)} \delta(\tau - \tau_*)$$



- Almost all modes would become super-horizon at the end of inflation hence constant amplitude.

# Inflation: Oscillatory Pattern



- After horizon crossing, the amplitude is frozen and different modes exit the horizon at different time.
- The amplitude is **oscillating with  $k$**  at the horizon.

# Finite Duration of Phase Transition

- The phase transition happens during finite time interval.
- For  $\beta/H_{\text{inf}} \gg 1$ , we take the time integration with the source.
- Need detailed simulation for  $\beta/H_{\text{inf}} \sim O(1)$ .

**Convolved with Green Function!**

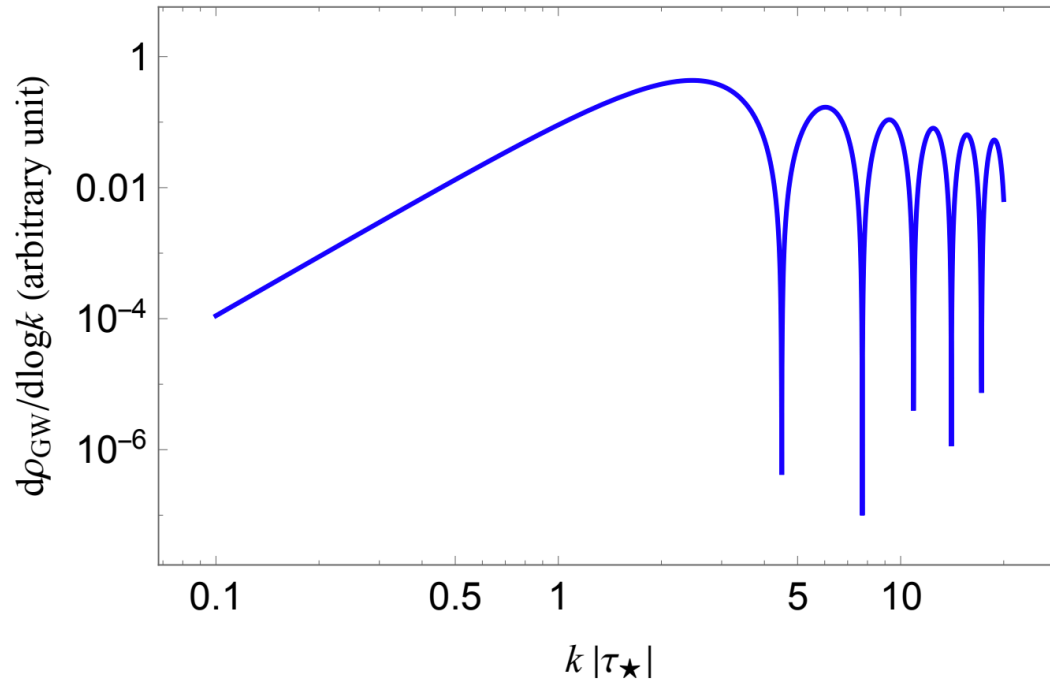
# De Sitter Inflation

$$a(\tau) = -\frac{1}{H_{\text{inf}}\tau}$$

$$\tilde{h}_{ij}(\eta, \mathbf{k}) = \frac{16\pi G_N}{k} \int d\eta' \frac{\tilde{\mathcal{G}}(\eta, \eta')}{a(\eta)} a^3(\eta') \tilde{\sigma}_{ij}^{\text{TT}}(\eta', \mathbf{k})$$

$$\tilde{h}_{\mathbf{k}}(\tau) \approx -\frac{16\pi G_N H_{\text{inf}} \tilde{T}^{(0)} \tau}{k} \left[ \left( \frac{1}{k\tau} - \frac{1}{k\tau_\star} \right) \cos k(\tau - \tau_\star) + \left( 1 + \frac{1}{k^2(\tau\tau_\star)} \right) \sin k(\tau - \tau_\star) \right]$$

$$\left. \frac{\mathcal{G}(\eta, \eta')}{a(\eta)} \right|_{\eta \rightarrow 0} = \frac{H_{\text{inf}}}{k} \left( \frac{\sin \eta'}{\eta'} - \cos \eta' \right)$$



Oscillatory pattern:  
oscillating with  $k$



# Post Inflationary Evolution

$$\tilde{h}_{ij}(\tau, \mathbf{k}) = \tilde{h}_{ij}^f(\mathbf{k})\mathcal{E}(k\tau)$$

After entering deep RD  $\mathcal{E}(\eta) = \tilde{\mathcal{E}}_0^i a^{-1} \sin(\eta + \phi)$

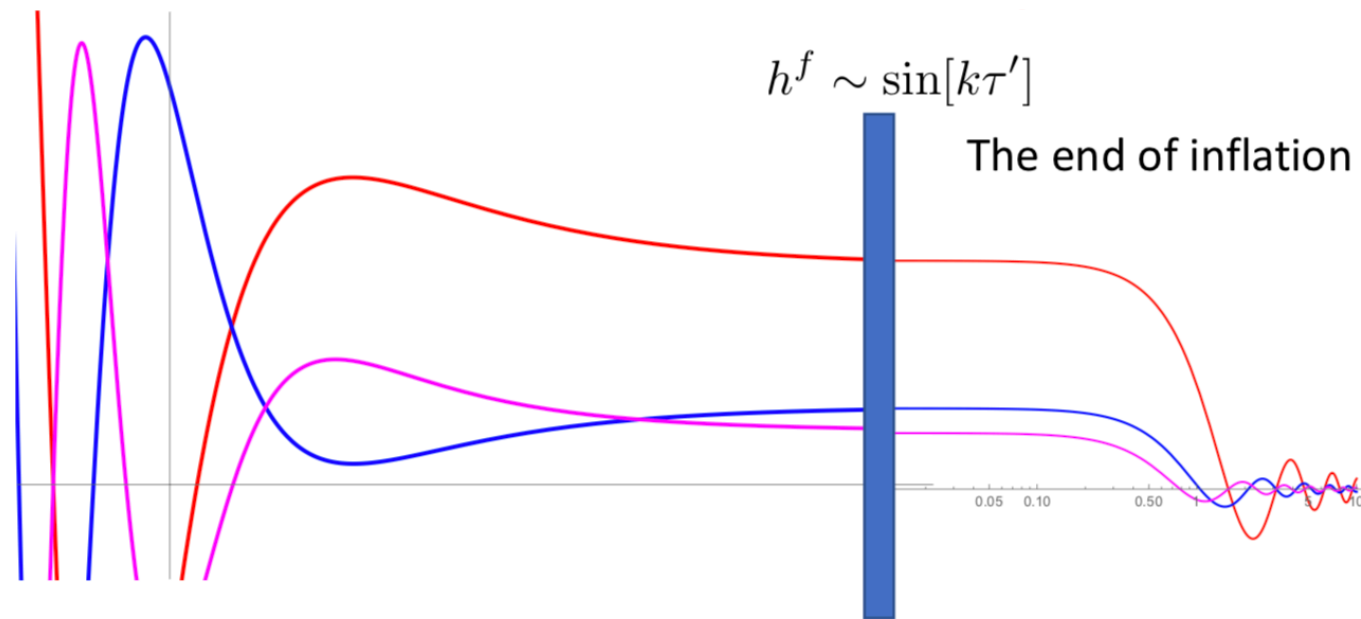
## Instantaneous Reheating

RD:  $\mathcal{E}(\eta) = \frac{\sin \eta}{\eta}$

Other cases:

Slow reheating, primordial black hole dominated intermediate stage

More complicated post-inflationary evolution.



# Spectrum distortion

Flat space-time

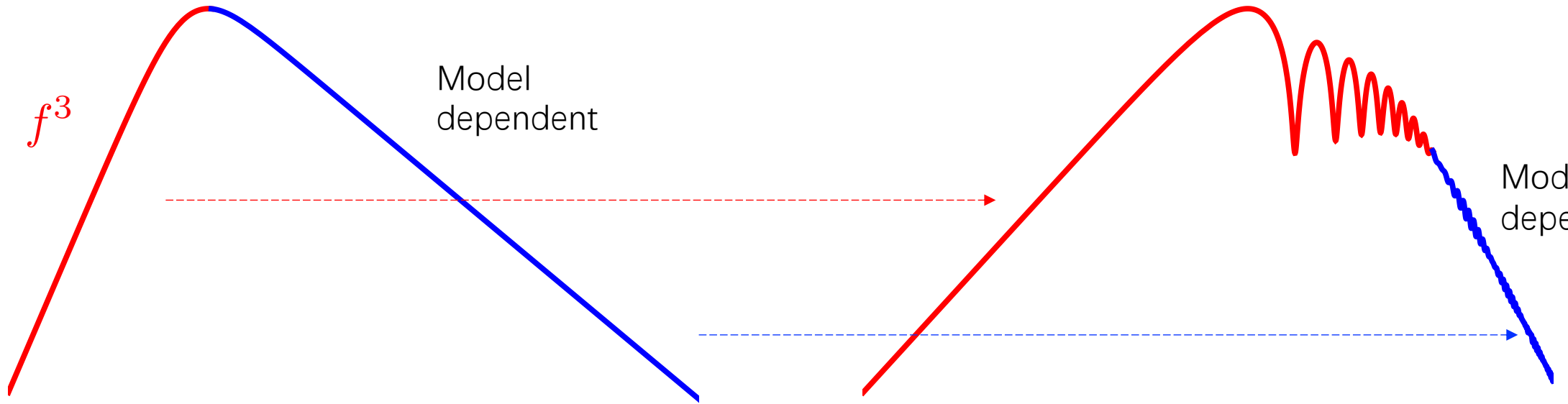
Quasi de Sitter

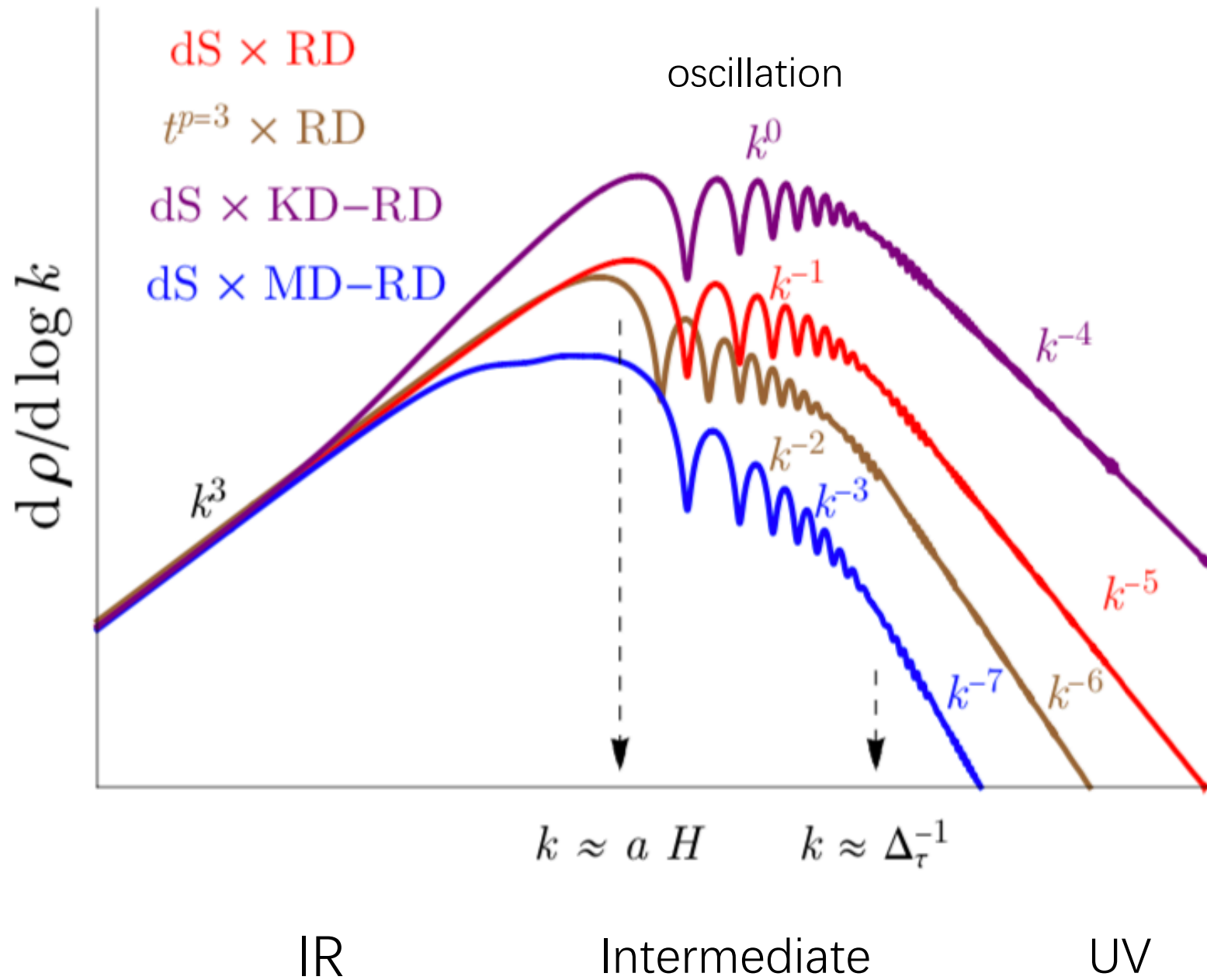
$f^3$

Model  
dependent

Model  
dependent

Cai, Pi, Sasaki, 1909.13728





# Observed GW signals today

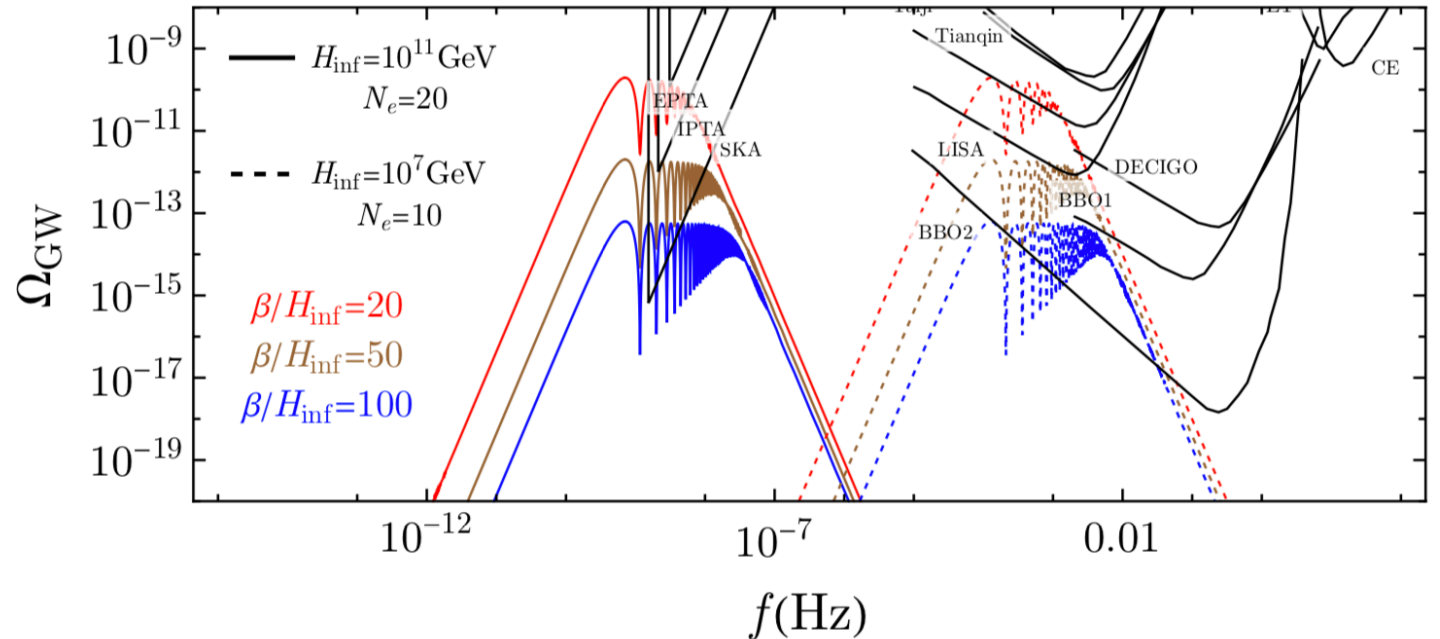
- Peak frequency today

$$f_{\text{today}} = f_{\star} \times \frac{a(\tau_{\star})}{a_r} \left( \frac{g_{*S}^{(0)}}{g_{*S}^{(R)}} \right)^{1/3} \frac{T_{\text{CMB}}}{\left[ \left( \frac{30}{g_{*}^{(R)} \pi^2} \right) \left( \frac{3H_r^2}{8\pi G_N} \right) \right]^{1/4}}$$

In dS- $t^{\tilde{p}}$ -RD

$$\tilde{f}_{\text{today}}^{\text{peak}} = 1.1 \times 10^{11} \text{ Hz} \times \left( \frac{H_{\text{end}}}{m_{\text{pl}}} \right)^{1/2} \left( \frac{a_r}{a_{\text{end}}} \right)^{-\frac{1}{2\tilde{\alpha}-1} - \frac{1}{2}} \frac{a(\tau_{\star})}{a_{\text{end}}}$$

KD intermediate stage, dS  $\times t^{\tilde{p}=1/3}$ -RD



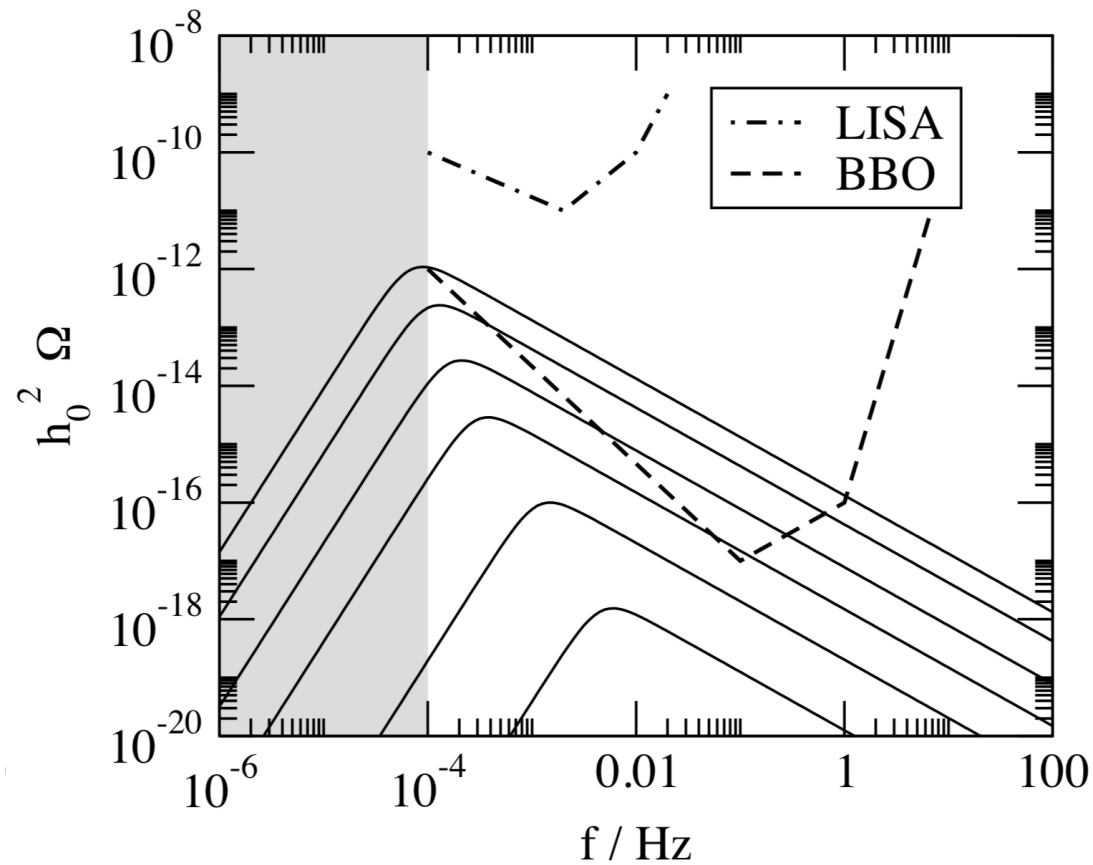
# Conclusion

- The power spectrum we get at the end of inflation shows oscillatory feature.
- The GWs might be observed on the space or ground based GW observatories.
- The observed GWs today produced by strong first order phase transition would carry the information of
  - (a) the details of the phase transition
  - (b) the inflation
  - (c) the post-inflationary evolution history

*Thank you!*

Backup

Simulation result under the envelope approximation done in flat space



$$\beta \equiv - \left. \frac{dS_4}{dt} \right|_{t_*}$$

determine the size and time interval of the phase transition



# GW Power spectrum

- Assume quasi-dS inflation, RD re-entering and fast reheating

$$\Omega_{\text{GW}}(k_{\text{today}}) = \underbrace{\Omega_R \frac{H_{\text{inf}}^4}{k_p^4}}_{\text{Dilution factor}} \left[ \frac{1}{2} + \underbrace{\mathcal{S}(k_p \beta^{-1})}_{\text{Smearing}} \cos\left(\frac{2k_p}{H_{\text{inf}}}\right) \right] \underbrace{\frac{\Delta\rho_{\text{vac}}}{\rho_{\text{inf}}}}_{\text{Suppressed by the energy fraction}} \underbrace{\frac{d\rho_{\text{GW}}^{\text{flat}}}{\Delta\rho_{\text{vac}} d \log k_p}}_{\text{Flat space spectrum}}$$

Redshift

$$\frac{f_{\text{today}}}{f_{\star}} = \frac{a(\tau_{\star})}{a_1} \left( \frac{g_{\star S}^{(0)}}{g_{\star S}^{(R)}} \right)^{1/3} \frac{T_{\text{CMB}}}{\left[ \left( \frac{30}{g_{\star}^{(R)} \pi^2} \right) \left( \frac{3H_{\text{inf}}^2}{8\pi G_N} \right) \right]^{1/4}}$$

$e^{-N_e}$

$N_e$ : e-folds before the end of inflation