

The Standard Model of Particle Physics, Lecture 1

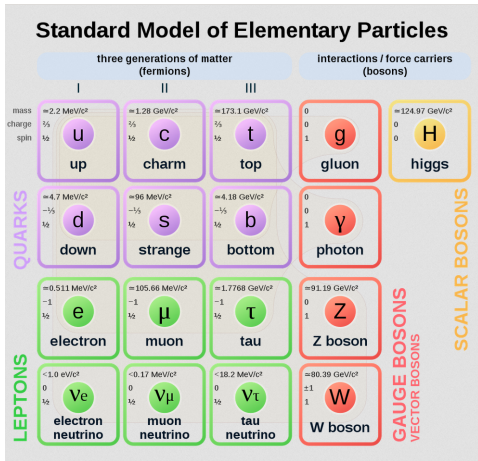
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Overview of Standard Model

- The standard model of particle physics involves understanding and categorizing **fundamental particles** and their **interactions**.



[courtesy:www.wikipedia.org]

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- It was observed that $M_N = 939.565$ MeV is strikingly close to the proton mass $M_P = 938.272$ MeV.
- This led to a proposal by **Heisenberg** that these are **two degenerate states of the same quantum number** → **Isospin !**

Isospin

- The quantum number corresponding to isospin is denoted by I . One can denote the proton and neutron states by the component of the total isospin given by I_3 . The wavefunctions can be denoted as,

$$|p\rangle \equiv |x, t, s, I, I_3\rangle, \quad |n\rangle \equiv |x, t, s, I, -I_3\rangle.$$

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$$|p\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |n\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- If we take the states $|n, p\rangle$ eigenfunctions of \hat{I}_3 , then denoting $I_3 = \frac{\sigma_3}{2}$ we get

$$\frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} |p, n\rangle = \pm \frac{1}{2} |p, n\rangle.$$

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- using which one can go from one component-state to the other

$$I_+|p\rangle = 0 , \quad I_+|n\rangle = |p\rangle , \quad I_-|p\rangle = |n\rangle , \quad I_-|n\rangle = 0$$

More on Isospin

- Using photo-emulsion techniques Cecil Powell and his group looked at the decay products of high energy cosmic rays. Pions were discovered in 1945.

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$$|\pi^\pm \rangle = |1, \pm 1 \rangle , \quad |\pi^0 \rangle = |1, 0 \rangle .$$

- It was observed that **strong interactions conserve isospin!** Interactions of two protons are of same strength as that between two neutrons. Mathematically $[H_{\text{strong}}, I] = 0$.

More on Isospin: Pion-nucleon scattering

- One can use this fact that strong interactions conserve isospin to calculate the ratio of cross sections of the decay process

$$\frac{\sigma(\pi^- p \rightarrow \pi^- p)}{\sigma(\pi^- p \rightarrow \pi^0 n)}$$

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- Step 1: One can write the states in terms of product states using addition of isospins $I = I^p + I^\pi$. The highest weight state, for example, is given as $|1, 1\rangle|\frac{1}{2}, \frac{1}{2}\rangle \equiv |\frac{3}{2}, \frac{3}{2}\rangle$.

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- Thus acting $I_- = I_-^p + I_-^\pi$ on the highest state one gets the next state with $I = \frac{3}{2}$, $I_3 = \frac{1}{2}$.

[Note that $I_\pm |I, I_3\rangle = \sqrt{I(I+1) - I_3(I_3 \pm 1)} |I, I_3 \pm 1\rangle$].

$$I_- |\frac{3}{2}, \frac{3}{2}\rangle = |\frac{3}{2}, \frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |1, 0\rangle |\frac{1}{2}, \frac{1}{2}\rangle + \sqrt{\frac{1}{3}} |1, 1\rangle |\frac{1}{2}, -\frac{1}{2}\rangle .$$

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- Step 2: Act I_- again on this state. We can show that both the states $\pi^- p$ and $\pi^0 n$ decay via $|\frac{3}{2}, -\frac{1}{2}\rangle$ with weights $\frac{1}{\sqrt{3}}$ and $\sqrt{\frac{2}{3}}$ respectively.

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- Since cross-sections are probability densities of the isospin states hence,

$$\frac{\sigma(\pi^- p \rightarrow \pi^- p)}{\sigma(\pi^- p \rightarrow \pi^0 n)} = \left| \frac{\frac{1}{\sqrt{3}}}{\sqrt{\frac{2}{3}}} \right|^2 = \frac{1}{2} .$$

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- In fact if you look into the Particle Data Group this predicted ratio of cross-sections are quite close to experimental values! **Small deviations due to the fact that $SU(2)$ isospin symmetry is near-exact.**

Strangeness

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- Resolution of this puzzle: Gell-Mann and Nishijima introduced a new quantum number called strangeness.
- They postulated that strangeness has to be conserved in strong and electromagnetic interactions. Thus assigned $S = 0$ to Δ , $S = +1$ to K^+ and $S = -1$ to the baryons Σ, Λ . This also explains why strange particles are produced in pairs.

References

- F. Halzen and A. D. Martin, "Quarks and Leptons", John Wiley & Sons (1984).
- T-P Cheng, L-F Li, "Gauge Theory of Elementary Particle Physics", Oxford University Press (1984).