

# *STATUS OF NEGATIVE COUPLING MODIFIERS FOR EXTENDED HIGGS SECTORS*

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hep-ph: arXiv:2111.02533

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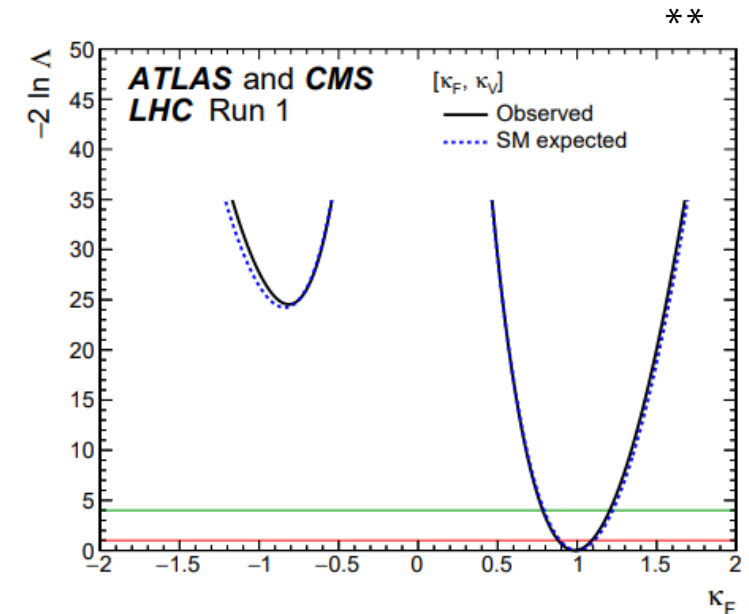


# OUTLINE

- Introduction
- Custodial Symmetry in BSM
- Accidentally Custodial triplet
- Generalization for different multiplets
- Scans
- Conclusion

# INTRODUCTION

- The LHC shows that corrections from new physics appear to be small.
- This could be an artifact of the way we are accessing information. Lose information on the sign of couplings!
- In this work we focus on the ratio of coupling modifier
$$\lambda_{WZ} = \frac{\kappa_W}{\kappa_Z}.$$
- This observable probes custodial violation on the Higgs sector.
- Current CMS fits\* gives a preference for **negative**  $\lambda_{WZ}$ .



# CUSTODIAL SYMMETRY IN BSM

- Custodial symmetry is an accidental global  $SU(2) \times SU(2)$  symmetry of the Higgs sector.
- In the SM we have the  $\lambda_{WZ} = 1$  and there is also the  $\rho$  parameter:  $\rho = \frac{m_W^2}{c_W^2 m_Z^2}$ .
- Naively one would think that both observables are connected in any model, which is not true.
- $\rho$  is connected to the vacuum while  $\lambda_{WZ}$  is connected to the potential.
- We can have custodial symmetric vacuum with custodial violating potentials (Accidentally custodial models)!

# CUSTODIAL SYMMETRY IN BSM

- If we assume that the vacuum preserves the symmetry, the custodial violation inside  $\lambda_{WZ}$  comes only from the diagonalization of the different multiplets:
- $h = R_0\varphi_0^R + R_1\psi_1^R + R_2\psi_2^R + \dots$
- The kappas have the form:
- $\kappa_V^h = R_0\kappa_V^{\text{doublet}} + R_1\kappa_V^{\text{multiplet 1}} + R_2\kappa_V^{\text{multiplet 2}} + \dots$
- If we want to explore if a region is excluded, we do not need to resolve the source of  $\vec{R}$  in terms of the potential, we can treat them as random variables and independent of the vevs!

# ACCIDENTALLY CUSTODIAL TRIPLET

- The simplest model\* that can generate negative  $\lambda_{WZ}$  involves two triplets: We have the usual Higgs doublet  $(\varphi^+, \varphi^0)$  with  $Y = 1$ , a complex triplet  $(\chi^{++}, \chi^+, \chi^0)$  with  $Y = 2$  and a real triplet  $(\xi^+, \xi^0, \xi^-)$  with  $Y = 0$ .

- The EW symmetry breaking has the form:

$$\varphi^0 = \frac{v_\varphi}{\sqrt{2}} + \frac{1}{\sqrt{2}}(\varphi_R^0 + i\varphi_I^0), \quad \xi^0 = v_\xi + \xi_R^0, \quad \chi^0 = v_\chi + \frac{1}{\sqrt{2}}(\chi_R^0 + i\chi_I^0).$$

- The  $\rho$  parameter in this model is:

$$\rho = \frac{m_W^2}{c_W^2 m_Z^2} = \frac{v_\varphi^2 + 4v_\chi^2 + 4v_\xi^2}{v_\varphi^2 + 8v_\chi^2}$$

- $v_\xi = v_\chi \rightarrow$  vacuum custodially symmetric

# ACCIDENTALLY CUSTODIAL TRIPLET

- $h = R_0\varphi_R^0 + R_1\chi_R^0 + R_2\xi_R^0$
- $\lambda_{WZ} = \frac{2\sqrt{2}\nu_\chi R_1 + 4\nu_\chi R_2 + \nu_\varphi R_0}{4\sqrt{2}\nu_\chi R_1 + \nu_\varphi R_0}$

- Assume  $\lambda_{WZ} = -1$ , write the other observables:

$$\lambda_{fZ} = \frac{\kappa_f}{\kappa_Z} \text{ and } \kappa_{fZ} = \frac{\kappa_f \kappa_Z}{\kappa_h}$$

- Then, solve for  $\vec{R}$ , if we assume  $\lambda_{fZ} = \pm 1$  we get:

$$\kappa_{fZ} = R_0 \frac{\nu_\varphi}{\nu} \text{ for } \lambda_{fZ} = 1$$

$$\kappa_{fZ} = -1.39 R_0 \frac{\nu_\varphi}{\nu} \text{ for } \lambda_{fZ} = -1$$

# GENERALIZATION FOR DIFFERENT MULTIPLETS

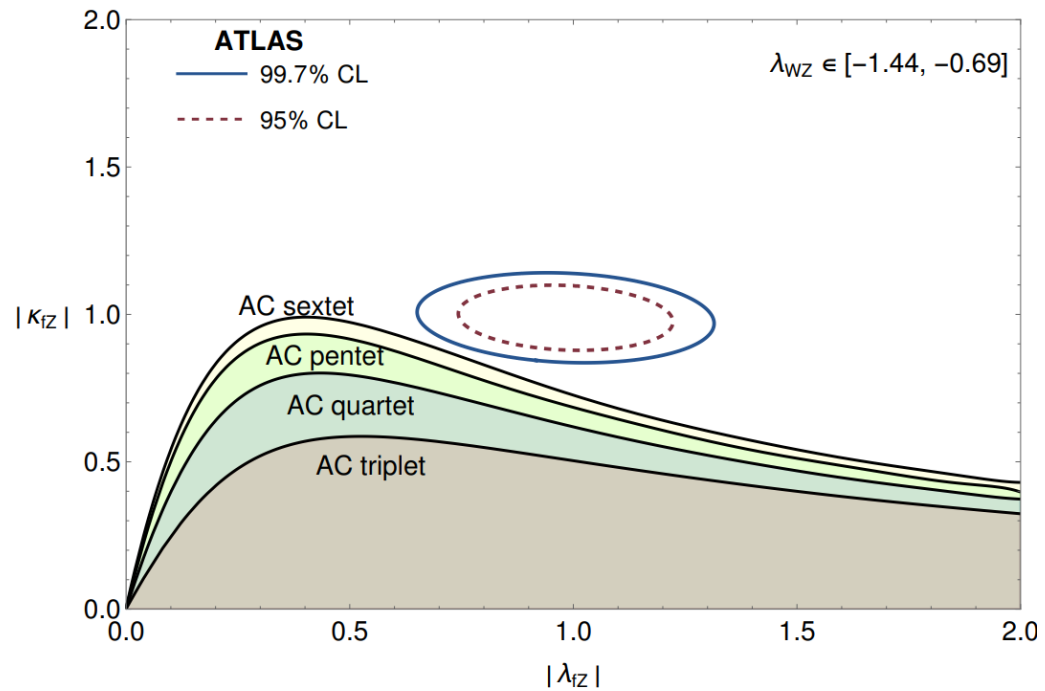
- Any model with a custodial limit can contribute to  $\lambda_{WZ}$  while avoiding the  $\rho$  parameter.
- The particle content of those models can be constructed from the generalized Georgi-Machacek models\*, which we can break down into  $SU(2) \times U(1)$  quantum numbers:
- AC triplets: one field with  $(1,2)$  and one with  $(1,0)$
- AC quartets: one field with  $(3/2,3)$  and one with  $(3/2,1)$
- AC pentets: one field with  $(2,4)$ , one with  $(2,2)$  and one with  $(2,0)$
- AC sextets: one field with  $(5/2,5)$ , one with  $(5/2,3)$  and one with  $(5/2,1)$



# GENERALIZATION FOR DIFFERENT MULTIPLETS

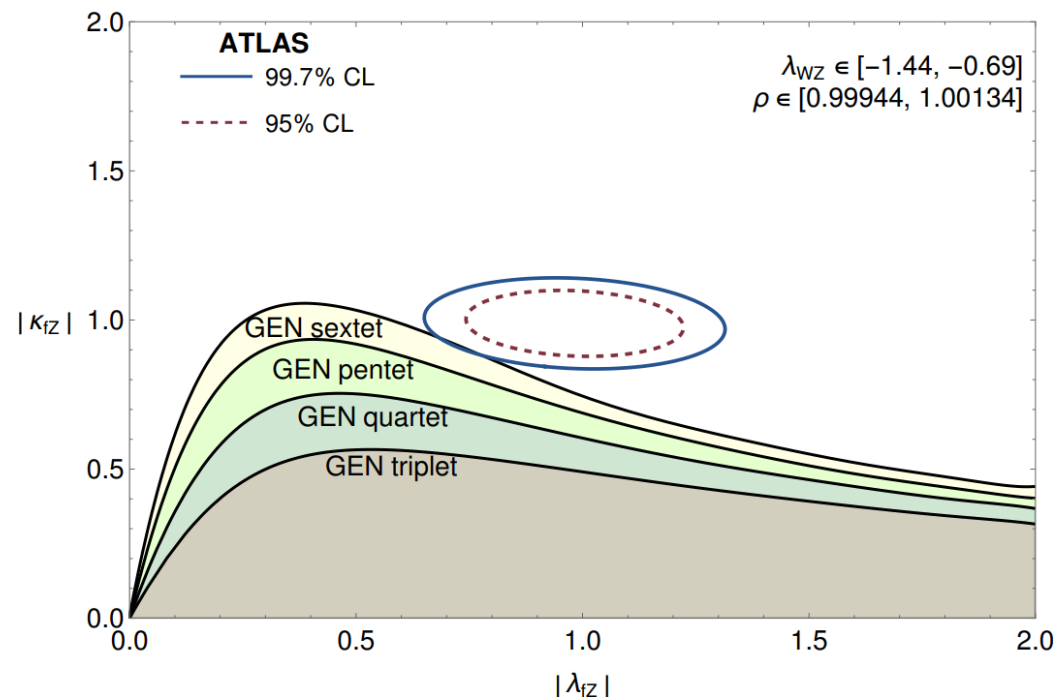
- Perturbative unitarity constraints the number of allowed models, assuming one doublet we can have 4487 possibilities.
- From these combinations, we can only have at most one AC sextet, four AC pentets, 23 AC quartets, or 145 AC triplets.
- The models that we study here are AC triplet, AC quadruplet, AC pentet, AC sextet, 2 AC triplets, AC pentet + AC sextet, and two AC pentet + AC sextet. Additionally, we also explore the case with general vev's for each of these models.
- We also explore the limiting cases of four AC pentets, 23 AC quartets, and 145 AC triplets.

# SCANS



EXCLUDED AT 99.7%CL

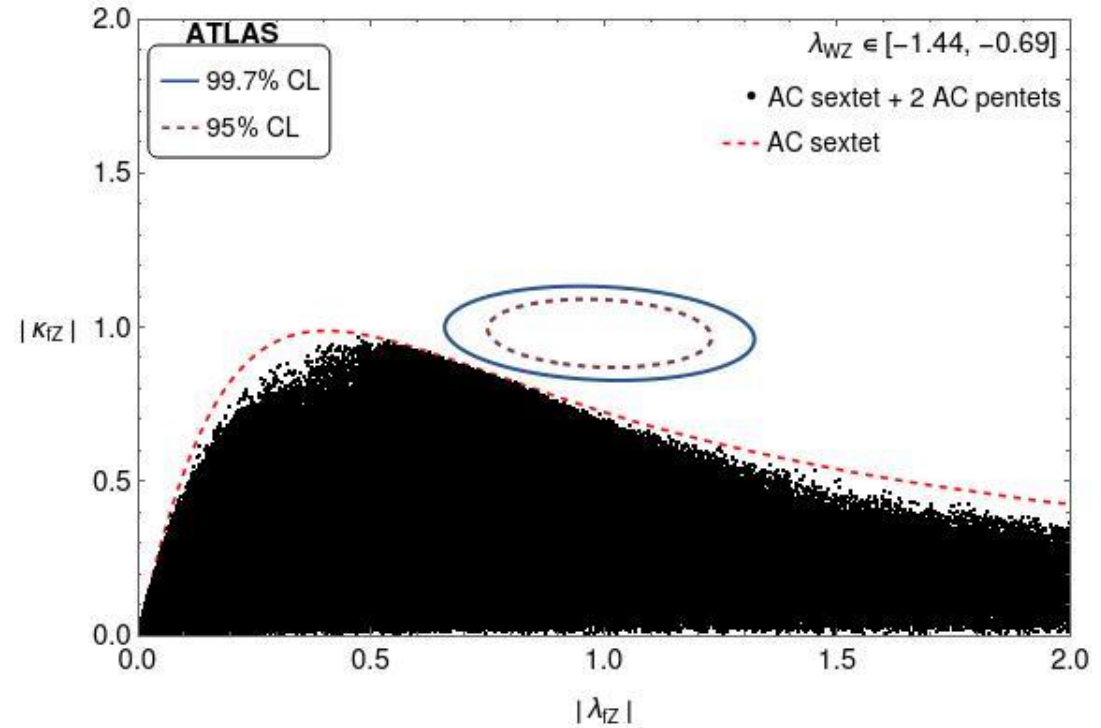
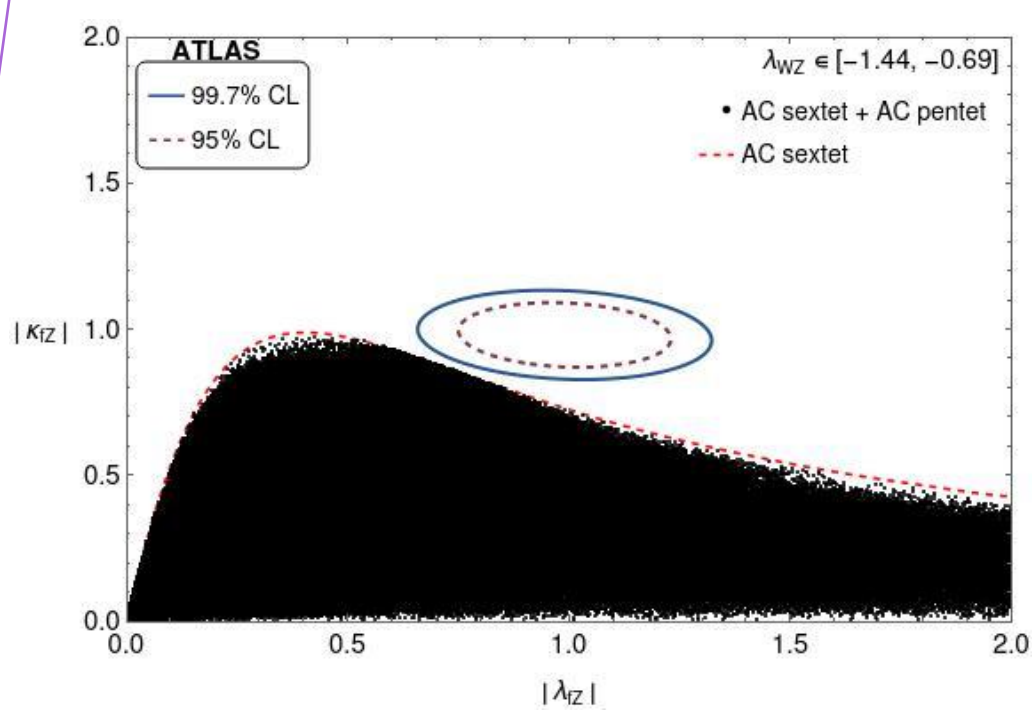
(a)



EXCLUDED AT 99.5%CL

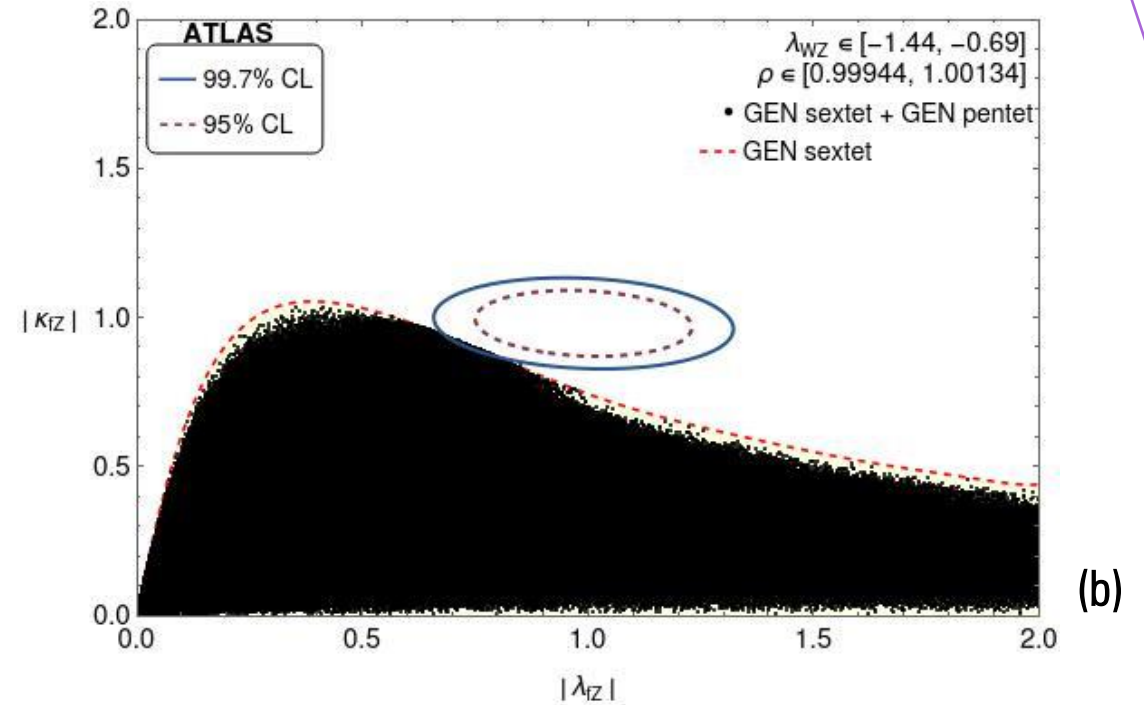
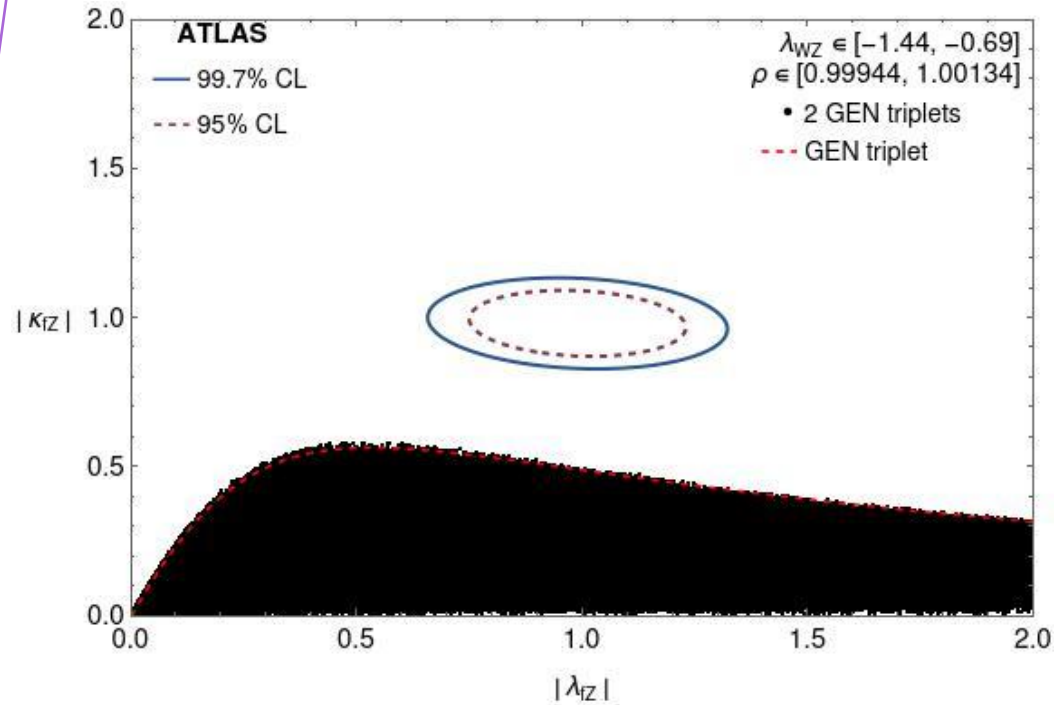
(b)

# SCANS



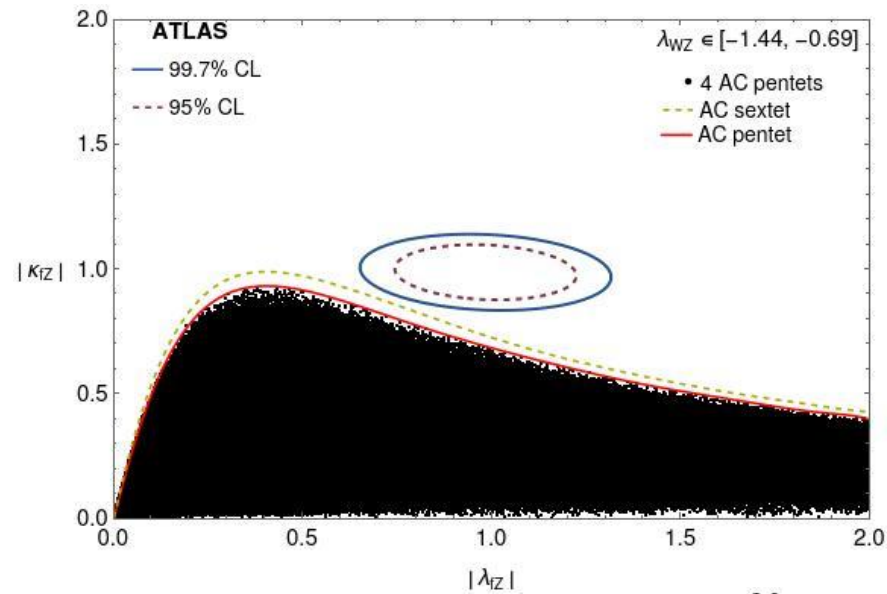
More multiplets is not better. Same behavior as the largest multiplet.

# SCANS

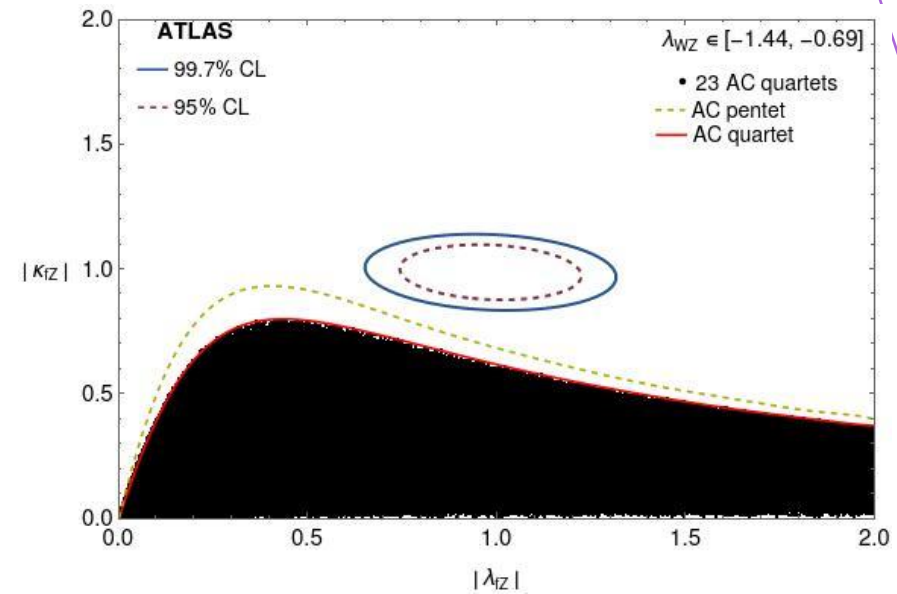


General vevs also does not help.  $\rho$  parameter constraint is too strong, even at  $5\sigma$ .  
Same behaviour as the situation with only the largest multiplet

# SCANS

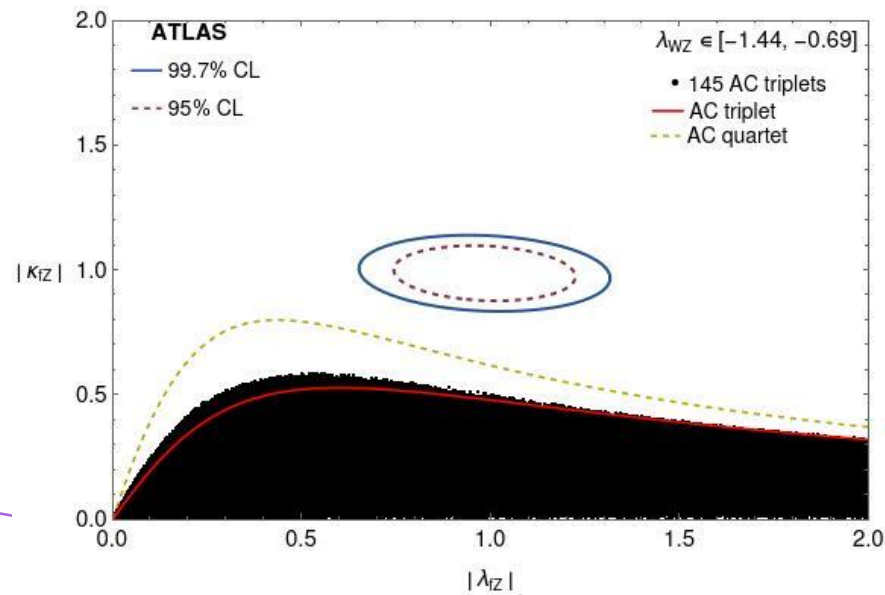


(a)



(b)

Extremal cases are also excluded with more than 99.7%CL!



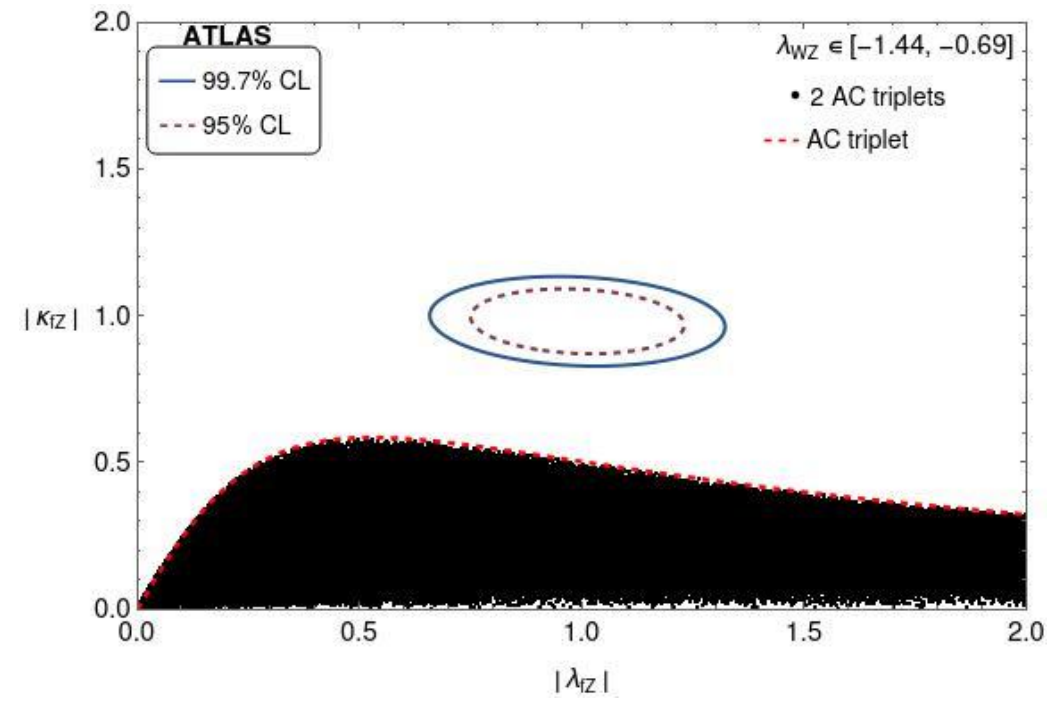
(c)

Equal small vevs ~ N copies

# CONCLUSION

- In this work, we studied the current status of negative coupling modifiers in the extended Higgs sector, with the focus on the observable  $\lambda_{WZ}$ .
- We present the analysis for the simplest case of AC triplets, and then we show how to generalize the procedure to different multiplets.
- The possibility of exploring this wide range of models lies in the fact that the coupling modifiers, in the end, depend only on the diagonalization matrix and the vevs.
- Under the analysis done we can see that all the models with one or more AC multiplets studied here are excluded under ATLAS results at 99.7%CL ( 99.5%CL for GEN sextet).
- What does this mean for negative  $\lambda_{WZ}$ ?
- We can then say that this region of parameter space is heavily disfavored for any weakly coupling extended scalar sector.
- In contrast, if the measured value for CMS stays to be negative and different experiments confirm this, we would not be able to describe the new physics using the current methods.
- It may be that new physics is hiding in plain sight, after all, only future experiments can tell.

# BACKUP



# BACKUP

Observables used:

$$\kappa_{fZ} = \frac{\kappa_f \kappa_Z}{\kappa_h}$$

$$\lambda_{WZ} = \frac{\kappa_W}{\kappa_Z}$$

$$\lambda_{fZ} = \frac{\kappa_f}{\kappa_Z}$$

$$\kappa_h = \sqrt{0.75\kappa_f^2 + 0.22\kappa_W^2 + 0.03\kappa_Z^2}$$

ATLAS FIT:

$$(\lambda_{fZ}, \kappa_{fZ}) = (0.99, 0.98)$$

$$COV = \begin{pmatrix} 0.0093 & -0.00054 \\ -0.00054 & 0.0020 \end{pmatrix}$$

CMS FIT (UNCORRELATED):

$$\kappa_{fZ} = 1.03 \pm 0.09$$

$$\lambda_{fZ} = 1.10 \pm 0.11$$

$$\lambda_{WZ} = -1.13^{+0.10}_{-0.11}$$



# BACKUP

Coupling modifiers for different multiplets:

DOUBLET:

$$\kappa_f^{(\frac{1}{2},1)} = \frac{v_{(\frac{1}{2},1)}}{v}, \kappa_W^{(\frac{1}{2},1)} = \frac{v_{(\frac{1}{2},1)}}{v} \kappa_Z^{(\frac{1}{2},1)} = \frac{v_{(\frac{1}{2},1)}}{v}$$

AC TRIPLET:

$$\kappa_f^{(1,2)} = 0, \kappa_W^{(1,2)} = 2\sqrt{2} \frac{v_{(1,2)}}{v} \kappa_Z^{(1,2)} = 4\sqrt{2} \frac{v_{(1,2)}}{v}$$

$$\kappa_f^{(1,0)} = 0, \kappa_W^{(1,0)} = 4 \frac{v_{(1,0)}}{v} \kappa_Z^{(1,0)} = 0$$

AC QUARTET:

$$\kappa_f^{(\frac{3}{2},3)} = 0, \kappa_W^{(\frac{3}{2},3)} = 3\sqrt{2} \frac{v_{(\frac{3}{2},3)}}{v} \kappa_Z^{(\frac{3}{2},3)} = 9\sqrt{2} \frac{v_{(\frac{3}{2},3)}}{v}$$

$$\kappa_f^{(\frac{3}{2},1)} = 0, \kappa_W^{(\frac{3}{2},1)} = 7\sqrt{2} \frac{v_{(\frac{3}{2},1)}}{v} \kappa_Z^{(\frac{3}{2},1)} = \sqrt{2} \frac{v_{(\frac{3}{2},1)}}{v}$$

AC PENTET:

$$\kappa_f^{(2,4)} = 0, \kappa_W^{(2,4)} = 4\sqrt{2} \frac{v_{(2,4)}}{v} \kappa_Z^{(2,4)} = 16\sqrt{2} \frac{v_{(2,4)}}{v}$$

$$\kappa_f^{(2,2)} = 0, \kappa_W^{(2,2)} = 10\sqrt{2} \frac{v_{(2,2)}}{v} \kappa_Z^{(2,2)} = 4\sqrt{2} \frac{v_{(2,2)}}{v}$$

$$\kappa_f^{(2,0)} = 0, \kappa_W^{(2,0)} = 12 \frac{v_{(2,0)}}{v} \kappa_Z^{(2,0)} = 0$$

AC SEXTET:

$$\kappa_f^{(\frac{5}{2},5)} = 0, \kappa_W^{(\frac{5}{2},5)} = 5\sqrt{2} \frac{v_{(\frac{5}{2},5)}}{v} \kappa_Z^{(\frac{5}{2},5)} = 25\sqrt{2} \frac{v_{(\frac{5}{2},5)}}{v}$$

$$\kappa_f^{(\frac{5}{2},3)} = 0, \kappa_W^{(\frac{5}{2},3)} = 13\sqrt{2} \frac{v_{(\frac{5}{2},3)}}{v} \kappa_Z^{(\frac{5}{2},3)} = 9\sqrt{2} \frac{v_{(\frac{5}{2},3)}}{v}$$

$$\kappa_f^{(\frac{5}{2},1)} = 0, \kappa_W^{(\frac{5}{2},1)} = 17 \frac{v_{(\frac{5}{2},1)}}{v} \kappa_Z^{(\frac{5}{2},1)} = \sqrt{2} \frac{v_{(\frac{5}{2},1)}}{v}$$

# BACKUP

N copies limit:

$$\begin{aligned}\kappa_V^h &= R_0 \kappa_V^{\text{doublet}} + (R_1 \kappa_V^{\text{mult } 1} + R_2 \kappa_V^{\text{mult } 2}) + (R_3 \kappa_V^{\text{mult } 1} + R_4 \kappa_V^{\text{mult } 2}) + \dots = \\ &= R_0 \kappa_V^{\text{doublet}} + (R_1 + R_3 + \dots) \kappa_V^{\text{mult } 1} + (R_2 + R_4 + \dots) \kappa_V^{\text{mult } 2} = \\ &= R_0 \kappa_V^{\text{doublet}} + \widetilde{R}_1 \kappa_V^{\text{mult } 1} + R_2 \widetilde{\kappa}_V^{\text{mult } 2}\end{aligned}$$

Cauchy inequality:

$$R_0^2 + \frac{\widetilde{R}_1^2}{N} + \frac{\widetilde{R}_2^2}{N} \leq 1$$