Prospects of same sign trilepton search in the Type I two Higgs doublet model at the LHC

**Prasenjit Sanyal** 

Asia Pacific Center for Theoretical Physics, Korea

1

### Outline

- Overview of Type I two Higgs doublet model (2HDM).
- Limits on the charged Higgs parameters  $(m_{H^{\pm}} \tan \beta)$ .
- Electro Weak production of charged Higgs.
- Fermiophobic nature of charged Higgs and heavy Higgs in Type I 2HDM.
- Signatures of same sign trilepton in Type I 2HDM.

- 2HDM is the minimal but phenomenologically rich extension of SM under the same gauge symmetry.
- The scalar sector of the 2HDM consists of two SU(2) Higgs doublets  $\Phi_i$ , i = 1, 2.

$$\begin{array}{rcl} \Phi_i & = & \left(\begin{array}{c} \phi_i^+ \\ \frac{\nu_i + \rho_i + i \eta_i}{\sqrt{2}} \end{array}\right) \\ v_i = \langle \rho_i \rangle & v & = & \sqrt{\nu_1^2 + \nu_2^2} = 246 \ \text{GeV} \end{array}$$

 Mostly studied: CP conserving 2HDM with softly broken Z<sub>2</sub> (to avoid Higgs mediated FCNC) symmetry.

$$V_{\text{2HDM}} \supset m_{12}^2 \Phi_1^{\dagger} \Phi_2 + h.c.$$

- After EWSB, the scalar sector consists of two CP even Higgses (h and H), one CP odd scalar A and a pair of charged Higgs H<sup>±</sup>.
- Parameters:  $m_h^2$ ,  $m_H^2$ ,  $m_H^2$ ,  $m_{H^\pm}^2$ ,  $m_{12}^2$ , v,  $\tan \beta (= v_2/v_1)$ ,  $\sin(\beta \alpha)$
- In 2HDM, h is identified as the observed 125 GeV Higgs boson.

- Based on the Z<sub>2</sub> charge assignments for the fermions, there are four possible Yukawa structures – Type I, II, X and Y.
- After EW symmetry breaking the Yukawa Lagrangian in terms of the mass eigenstates is:

$$\begin{split} \mathcal{L}_{\text{Yuk, I}}^{\text{2HDM}} & = -\sum_{f=u,d,\ell} \frac{m_f}{v} \left( \xi_h^f \overline{f} h f + \xi_H^f \overline{f} H f - i \xi_A^f \overline{f} \gamma_5 A f \right) \\ & - \left\{ \frac{\sqrt{2} V_{ud}}{v} \overline{u} \left( \xi_A^u m_u P_L + \xi_A^d m_d P_R \right) H^+ d + \frac{\sqrt{2} m_l}{v} \xi_A^I \overline{v}_L H^+ I_R + \text{h.c.} \right\} \end{split}$$

$\xi_h^u$	$\xi_h^d$	$\xi_h^\ell$	$\xi_H^u$	$\xi_H^d$	$\xi_H^\ell$	$\xi_A^U$	$\xi^d_A$	$\xi_{A}^{\ell}$
$c_{lpha}/s_{eta}$	$c_{lpha}/s_{eta}$	$c_{lpha}/s_{eta}$	$s_lpha/s_eta$	$s_lpha/s_eta$	$s_lpha/s_eta$	$\cot \beta$	$-\cot \beta$	$-\cot \beta$

### Other couplings:

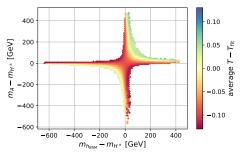
(A)  $hVV : \sin(\beta - \alpha)g_{hVV}^{SM}$ ,  $HVV : \cos(\beta - \alpha)g_{hVV}^{SM}$ , AVV : 0 where V = Z,  $W^{\pm}$ .

$$(\mathsf{B}) \ \ \mathit{hAZ}_{\mu} : \ \tfrac{g}{2C_W} \cos(\beta - \alpha)(\mathit{p} + \mathit{p}')_{\mu}, \quad \mathit{HAZ}_{\mu} : \ -\tfrac{g}{2C_W} \sin(\beta - \alpha)(\mathit{p} + \mathit{p}')_{\mu}$$

(C)  $H^{\pm}hW_{\mu}^{\mp}: \mp i\frac{g}{2}\cos(\beta-\alpha)(p+p')_{\mu}, \quad H^{\pm}HW_{\mu}^{\mp}: \pm i\frac{g}{2}\sin(\beta-\alpha)(p+p')_{\mu}, \quad H^{\pm}AW_{\mu}^{\mp}: \frac{g}{2}(p+p')_{\mu}$ 

4

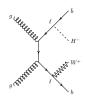
- Unitarity and vacuum stability bounds can be satisfied by proper choice of  $m_{12}^2 \in [0, m_H^2 \sin \beta \cos \beta]$ .
- The EWPO, T parameter depends strongly on the mass spliting of the charged Higgs and the neutral scalars.

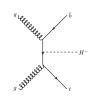


Henning Bahl, Tim Stefaniak, Jonas Wittbrodt, JHEP 06 (2021) 183

• Alignment limit:  $sin(\beta - \alpha) \rightarrow 1$  implies that the couplings of h is like SM Higgs boson.

## $H^{\pm}$ production and decay channels







- The QCD induced production of charged Higgs depends on its mass with respect to top quark and can be divided into
  - (A) Light Scenario:  $m_{H^\pm} \lesssim 160~{\rm GeV}$
  - (B) Heavy Scenario: ( $m_{H^\pm} \gtrsim 200~{\rm GeV}$ )
  - (C) Intermediate Scenario:  $(m_{H^\pm} \sim m_t)$
- In Type I 2HDM, the strong production of charged Higgs is  $\tan \beta$  suppressed.

### Conventional decay channels:

(A) 
$$H^+ \rightarrow \tau^+ \nu$$

(B) 
$$H^+ \rightarrow t\bar{b}$$

(C) 
$$H^+ \rightarrow c\bar{s}$$

(D) 
$$H^+ \rightarrow c\bar{b}$$

### Exotic decay channels:

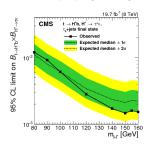
(A) 
$$H^{\pm} \rightarrow hW^{\pm}$$
 :  $\frac{\mp ig}{2}\cos(\beta - \alpha)(p_{\mu} - p_{\mu}^{\mp})$ 

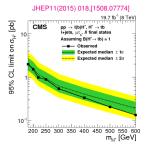
(B) 
$$H^{\pm} \rightarrow HW^{\pm}: \frac{\mp ig}{2} \sin(\beta - \alpha)(p_{\mu} - p_{\mu}^{\mp})$$

(C) 
$$H^\pm o AW^\pm : rac{g}{2}(p_\mu - p_\mu^\mp)$$

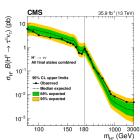
### CMS 8 TeV and 13 TeV results: ( $\tau \nu$ and tb channels)

#### JHEP11(2015) 018,[1508.07774]

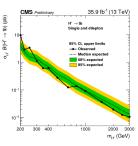




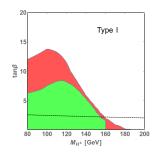
#### JHEP 07 (2019) 142,[1903.04560]

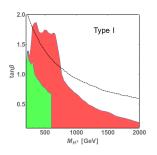


#### CMS PAS HIG-18-004



## **Constraints on charged Higgs parameter space**





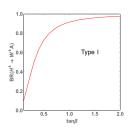
P.Sanyal, Eur.Phys.J.C 79 (2019) 11, 913

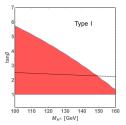
#### Exclusions:

- (A) Green Region --- 8 TeV CMS results
- (B) Red Region → 13 TeV CMS results
- (C) Dashed Line  $\longrightarrow \mathcal{BR}(B \to X_s \gamma)$

## Charged Higgs exotic decay channel ( $H^{\pm} \rightarrow W^{\pm} A$ )

- Minimal mass splitting together with alignment limit restricts the exotic decay channels H<sup>±</sup> → h/H/AW<sup>±</sup>.
- But once open, H<sup>±</sup> can dominantly decay to these channels.
- CMS collaboration put upper bounds on  $\mathcal{BR}(t \to H^+ b)$  at 95% CL assuming  $\mathcal{BR}(H^+ \to W^+ A) \to 1$  and  $\mathcal{BR}(A \to \mu^+ \mu^-) \to 3 \times 10^{-4}$ .
- Mass difference  $m_{H^\pm}-m_A=85$  GeV with  $m_{H^\pm}\in[100-160]$  GeV is considered by CMS group.
- T parameter is satisfied by choosing  $m_{H^\pm} \sim m_H$ .
- In Type I scenario,  $\mathcal{BR}(H^{\pm} \to W^{\pm}A) \to 1$  for  $\tan \beta \gtrsim 1$  and  $\mathcal{BR}(A \to \mu^{+}\mu^{-}) \sim 2.6 \times 10^{-4}$  for  $m_{A} \in [15-75]$  GeV.

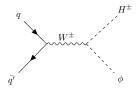




P.Sanyal, Eur.Phys.J.C 79 (2019) 11, 913

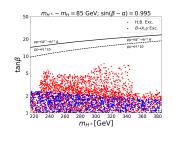
## Same sign trilepton search at the LHC in Type I 2HDM

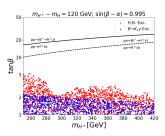
• For large  $\tan \beta$  the cross section  $pp \to W^{*\pm} \to H^{\pm} \phi$  dominates over the  $pp \rightarrow H^{\pm}tb$  channel in Type I 2HDM.



- For close to the alignment limit  $\phi \neq h_{SM}$ .
- Signal:
  - (A)  $pp \to W^{*\pm} \to H^{\pm}H \to (W^{\pm}H)(W^{+}W^{-}) \to (W^{\pm}W^{+}W^{-})(W^{+}W^{-}) \to 3\ell^{\pm}\cancel{E}_{T} + X$ (B)  $pp \to W^{*\pm} \to H^{\pm}A \to (W^{\pm}H)(ZH) \to (W^{\pm}W^{+}W^{-})(ZW^{+}W^{-}) \to 3\ell^{\pm}\cancel{E}_{T} + X$
- SM backgrounds: WZ+ jets,  $Z\ell^+\ell^-+$  jets and  $t\bar{t}W+$  jets.
- Parameter Choice:  $m_{H^{\pm}} m_H = 85$ , 120 GeV,  $m_H \in [130 300] \text{GeV}$ ,  $m_{H^{\pm}} \approx m_A$ ,  $\tan \beta \in [1, 50]$ ,  $\sin(\beta - \alpha) = 0.995$  and  $m_{12}^2 \in [0, m_H^2 \sin \beta \cos \beta]$ .

## **Current Limits in Type I 2HDM**





T. Mondal and P. Sanyal, arXiv:2109.05682

- Red Region: Exclusion regions from the LHC  $\sqrt{s}$  =7, 8, 13 TeV constraints on neutral Higgses and charged Higgs.
- Channels:

(1) 
$$H/A \rightarrow \tau \tau$$

(2)  $H/A \rightarrow \gamma \gamma$ 

(3)  $H \rightarrow VV (V = W^{\pm}, Z)$ 

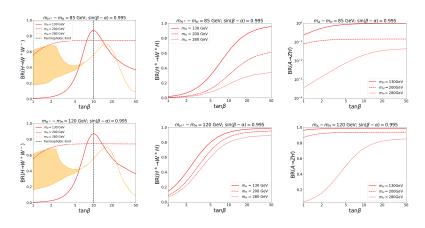
(4) 
$$A \rightarrow HZ$$

(5) 
$$A \rightarrow hZ$$

(6) 
$$H^{\pm} \rightarrow tb$$

Blue Region: Exclusion region coming from the  $\mathcal{BR}(B \to X_s \gamma)$  constraint.

## Bosonic decay modes of Higgs bosons



#### T. Mondal and P. Sanyal, arXiv:2109.05682

(A) 
$$\mathcal{BR}(H \to W^+W^-)$$

(B) 
$$\mathcal{BR}(H^{\pm} \rightarrow W^{\pm}H)$$

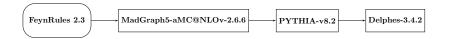
(C) 
$$\mathcal{BR}(A \rightarrow ZH)$$

$$HW^+W^-:\cos(eta-lpha)g_{hVV}^{SM}$$

$$H^{\pm}HW^{\mp}_{\mu}$$
:  $\pm irac{g}{2}\sin(eta-lpha)(p_{\mu}+p'_{\mu})$ 

$$AHZ:-rac{g}{2c_W}\sin(eta-lpha)(p_\mu+p'_\mu)$$

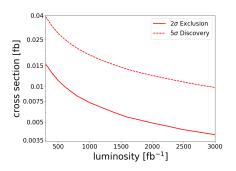
### **Collider Analysis**



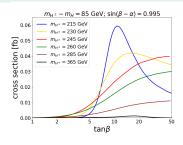
#### Selection cuts:

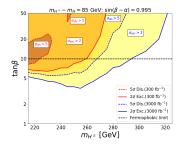
- (1) SS3L: Three isolated leading leptons  $(e, \mu)$  with same sign.
- (2) Momentum  $p_T$  cuts:  $p_T(\ell_1) > 30$  GeV,  $p_T(\ell_2) > 30$  GeV,  $p_T(\ell_3) > 20$  respectively and  $\not\not\!E_T > 30$  GeV.
- (3) Lepton and jet separation cuts: Lepton-lepton separation,  $\Delta R_{\ell\ell} > 0.4$  and lepton-jet separation cuts,  $\Delta R_{\ell j} > 0.4$  where  $\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2}$ .
- (4) Z-veto: Veto events with additional leptons oppositely charged to the tagged SS3L with same flavor satisfying the condition  $80 < m_{\ell^+\ell^-} < 100$  GeV.
- (5) b-veto: Veto events if there are any tagged b-jets.

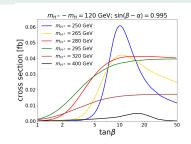
Background cross sections at $\sqrt{s}=$ 13 TeV								
Selection cuts	WZ+ jets [fb]	$Z\ell^+\ell^-+$ jets [fb]	$t\bar{t}W$ + jets [fb]					
MG5	1360.80	246.550	62.570					
SS3L	0.0543	0.00991	0.0878					
Lepton $p_T \& \mathcal{E}_T$	0.0122	0.00122	0.0118					
$\Delta R_{\ell\ell} \& \Delta R_{\ell j}$	0.0073	0.00083	0.0103					
Z-veto	0.0065	0.00065	0.0103					
<i>b</i> -veto	0.0061	0.00061	0.0018					

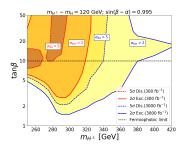


### Results









T. Mondal and P. Sanyal, arXiv:2109.05682

#### Conclusions

- Large  $\tan \beta$  regions in Type I 2HDM are allowed due to the fermiophobic nature of BSM Higgses.
- Two most conventional decay modes  $H^\pm \to au^\pm 
  u$  and  $H^+ \to t \bar b$  are studied using the CMS 13 TeV results and compared with 8 TeV results.
- CMS collaboration for the first time studied the exotic channel,  $H^{\pm} \to W^{\pm}A$ ,  $A \to \mu^{+}\mu^{-}$  to put upper limits on  $\mathcal{BR}(t \to H^{\pm}b)$ .
- The EW production of charged Higgs dominates over the strong production in Type I scenario for large tan β.
- Mass splitting of H<sup>±</sup> and H along with the fermiophobic limit of H in Type I scenario are used to study the same sign trilepton final state at the LHC.
- The exclusion and discovery limits are discussed for  $300 {\rm fb^{-1}}$  and  $3000 {\rm fb^{-1}}$  luminosities at  $\sqrt{s} = 13 {\rm TeV}$ .



#### Conclusions

- • Large  $\tan \beta$  regions in Type I 2HDM are allowed due to the fermiophobic nature of BSM Higgses.
- Two most conventional decay modes  $H^\pm \to \tau^\pm \nu$  and  $H^+ \to t \bar b$  are studied using the CMS 13 TeV results and compared with 8 TeV results.
- CMS collaboration for the first time studied the exotic channel,  $H^{\pm} \to W^{\pm} A$ ,  $A \to \mu^+ \mu^-$  to put upper limits on  $\mathcal{BR}(t \to H^{\pm} b)$ .
- The EW production of charged Higgs dominates over the strong production in Type I scenario for large tan β.
- Mass splitting of H<sup>±</sup> and H along with the fermiophobic limit of H in Type I scenario are used to study the same sign trilepton final state at the LHC.
- The exclusion and discovery limits are discussed for  $300 {\rm fb^{-1}}$  and  $3000 {\rm fb^{-1}}$  luminosities at  $\sqrt{s}=13 {\rm TeV}$ .

# Thank You