Sensitivity of future linear colliders to processes of dark matter production with light mediator exchange

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Introduction

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Probing Dark Matter with e⁺e⁻

Mono-photon signature

The mono-photon signature is considered to be the most general way to look for DM particle production in future e^+e^- colliders.



DM can be pair produced in the e^+e^- collisions via exchange of a new mediator particle, which couples to both electrons (SM) and DM states

This process can be detected, if additional hard photon radiation from the initial state is observed in the detector...





Heavy mediator approximation, generator level

arXiv:2103.06006

Mono-photon signature:

high energy, isolated photon no other "hard" activity in the detector

Highest sensitivity to DM production from the ratio of the photon energy distributions measured for the two electron beam polarisations





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WIMP Dark Matter at the 500 GeV ILC



arXiv:2001.03011

Heavy mediator (EFT limit), full simulation

Mono-photon signature:

high energy, isolated photon no other "hard" activity in the detector





Heavy mediator (EFT limit), full simulation

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Mono-photon signature:

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"Irreducible" background from radiative neutrino pair-production events $e^+e^- \rightarrow \nu\nu + \mathrm{N}\gamma$ dominates after selection and bg suppression cuts





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Comparison of mass scale limits calculated in the EFT framework



New analysis approach

Most of the studies performed so far focused on heavy mediator exchange (EFT limit) and coupling values O(1) \Rightarrow extracted were limits on DM or mediator masses

In our study:

- focus on light mediator exchange (DM even lighter)
- consider very small mediator couplings to SM, $\Gamma_{SM} \ll \Gamma_{tot}$



"Experimental-like" approach \Rightarrow focus on cross section limits as a function of mediator mass and width





Mono-photon study

Publications

- J. Kalinowski et al., *Simulating hard photon production with WHIZARD*, Eur. Phys. J. C 80 (2020) 634, arXiv:2004.14486
- J. Kalinowski et al., Sensitivity of future linear e⁺e⁻ colliders to processes of dark matter production with light mediator exchange, Eur. Phys. J. C in print, arXiv:2107.11194

Conference presentations

- P.Sopicki, Simulating hard photon production with WHIZARD, ICHEP'2020
- A.F.Żarnecki, Dark matter searches with mono-photon signature at future e⁺e⁻ colliders, DIS'2021 & Dark Matter 2021
- A.F.Żarnecki, *Probing Dark Matter with ILC*, SUSY'2021 & PANIC'2021
- A.F.Żarnecki, Sensitivity to dark matter production at future e⁺e⁻ colliders, Corfu'2021
- J.Kalinowski, Sensitivity of future e^+e^- colliders to processes of dark matter production with light mediator exchange, SUSY'2021
- A.F.Żarnecki, Sensitivity of future e⁺e⁻ colliders to processes of dark matter production with light mediator exchange, Matter To The Deepest 2021
- A.F.Żarnecki, *The use of beam polarization in the search for dark matter at the ILC*, SPIN'2021



Outline



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- 3 Simulating mono-photon events
 - 4 Analysis framework
- 5 Results
- 6 Conclusions



Colliders and Experiments

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Colliders



International Linear Collider



Technical Design (TDR) completed in 2013

arXiv:1306.6328

- superconducting accelerating cavities
- 250 500 GeV c.m.s. energy (baseline), 1 TeV upgrade possible
- footprint 31 km
- polarisation for both e^- and e^+ (80%/30%)

Colliders



Compact Linear Collider



Conceptual Design (CDR) presented in 2012

CERN-2012-007

- high gradient, two-beam acceleration scheme
- staged implementation plan with energy from 380 GeV to 3 TeV
- footprint of 11 to 50 km
- e⁻ polarisation (80%)

For details refer to arXiv:1812.07987



Running scenarios

Staged construction assumed for both ILC and CLIC. In our study we focused on the highest energy stages.

ILC

Total of 4000 fb^{-1} assumed at 500 GeV (H-20 scenario)

- $\bullet~2{\times}1600\,fb^{-1}$ for LR and RL beam polarisation combinations
- $2 \times 400 \text{ fb}^{-1}$ for RR and LL beam polarisation combinations

assuming polarisation of $\pm 80\%$ for electrons and $\pm 30\%$ for positrons

arXiv:1903.01629

CLIC

Total of $5000 \, \text{fb}^{-1}$ assumed at $3 \, \text{TeV}$

- 4000 fb^{-1} for negative electron beam polarisation
- 1000 fb⁻¹ for positive electron beam polarisation

assuming polarisation of $\pm 80\%$ for electrons

Detector Requirement

"Particle Flow" concept:

High calorimeter granularity ⇒ single particle reconstruction/ID

Precise momentum measurement \Rightarrow best energy for charged particles \Rightarrow dominates jet energy resolution

High precision vertex detector \Rightarrow very efficient flavour tagging

Instrumentation down to smallest angles \Rightarrow hermecity, missing energy tagging









Detector Requirements same for ILC and CLIC

- Track momentum resolution: $\sigma_{1/p} < 5 \cdot 10^{-5} \text{ GeV}^{-1}$
- Impact parameter resolution: $\sigma_d < 5\mu m \oplus 10\mu m \frac{1 \text{ GeV}}{n \sin^{3/2} \Theta}$
- Jet energy resolution: $\sigma_E/E = 3 4\%$ (for highest jet energies)
- Hermecity: $\Theta_{min} = 5 \text{ mrad}$

Detailed detector concepts for ILC and CLIC:



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CLICdet



Simulating mono-photon events

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Simulating mono-photon events in $\operatorname{WHIZARD}$



For proper estimate of the mono-photon signature sensitivity consistent simulation of BSM processes and of the SM backgrounds is crucial.

"Irreducible" background comes from radiative neutrino pair-production



Detector acceptance & reconstruction efficiency

 \Rightarrow significant contribution from radiative Bhabha scattering

WHIZARD provides the ISR structure function option that includes all orders of soft and soft-collinear photons as well as up to the third order in high-energy collinear photons.

However, WHIZARD ISR photons are not ordinary final state photons: they represent all photons radiated in the event from a given lepton line.

Simulating mono-photon events in $\operatorname{WHIZARD}$



ISR structure function can not account for hard non-collinear photons \Rightarrow all "detectable" photons generated on Matrix Element level

Dedicated procedure developed to avoid double-counting of ISR and ME For details: J. Kalinowski et al., Eur. Phys. J. C 80 (2020) 634, arXiv:2004.14486

Two variables, calculated separately for each emitted photon:

$$\begin{array}{lll} q_{-} & = & \sqrt{4E_{0}E_{\gamma}}\cdot\sin\frac{\theta_{\gamma}}{2} \ , \\ q_{+} & = & \sqrt{4E_{0}E_{\gamma}}\cdot\cos\frac{\theta_{\gamma}}{2} \end{array}$$

are used to separate "soft ISR" emission region from the region described by ME calculations.





Validation of the procedure

WHIZARD predictions were compared to the results from the KKMC code for $e^+e^- \rightarrow \nu\bar{\nu} + N\gamma$ 3 TeV CLIC



 \Rightarrow very good agreement observed (both for shape and normalisation)

For more details:

J. Kalinowski et al., Eur. Phys. J. C 80 (2020) 634, arXiv:2004.14486



Simplified DM model

UFO model covering most popular scenarios of DM pair-production Possible mediators: \Rightarrow Feynrules

- scalar
- pseudo-scalar
- vector
- pseudo-vector
- V-A coupling
- V+A coupling
- Possible DM candidates:
 - real or complex scalar
 - Majorana or Dirac fermion
 - real vector

Cross section for $e^+e^- \rightarrow \chi\chi$ for $M_{\chi} = 50 \text{ GeV}$ and $M_{Y} = 300 \text{ GeV}$





ISR rejection probability

Fraction of events generated by WHIZARD **removed** in merging procedure (ISR photons emitted in the phase-space region covered by ME)



CLIC @ 3 TeV





Detector simulationfor ILC running at 500 GeVHundreds of signal scenarios to consider \Rightarrow DELPHES fast simulation

Modeling of the forward detector performance crucial for the analysis

Results on BeamCal efficiency from full simulation:

ILD IDR

Moritz Hebermehl PhD Thesis





Detector simulation for ILC running at 500 GeV

 ILCgen model for $\operatorname{DelPHeS}$ includes proper modelling of forward detectors

BeamCal geometry

Parameterisation of the full simulation results



Included in the official DELPHES repository as delphes_card_ILCgen.tcl



Detector simulation for ILC running at 500 GeV

 ${\sf ILCgen}$ model for ${\rm DelPHes}$ includes proper modelling of forward detectors

BeamCal geometry

Reconstruction results for $e^+e^-
ightarrow e^+e^-$



Included in the official DELPHES repository as delphes_card_ILCgen.tcl



Tagging efficiency

based on $\operatorname{DELPHES}$ simulation

Mono-photons reconstructed only in a fraction of generated signal event

 $\sigma\left(\textit{e}^{+}\textit{e}^{-} \rightarrow \chi \; \chi \; \gamma_{\text{tag}}\right) \; = \; \textit{f}_{\text{mono-photon}} \cdot \sigma\left(\textit{e}^{+}\textit{e}^{-} \rightarrow \chi \; \chi \; (\gamma) \;\right)$

ILC @ 500 GeV

CLIC @ 3 TeV



Analysis framework

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Analysis framework



Event selection

- On generator level:
 - 1, 2 or 3 ME photons nonradiative events for signal only (for normalisation)
 - all ME photons with $q_{\pm} > 1 \text{ GeV}$ & $E^{\gamma} > 1 \text{ GeV}$ rejected are events with $q_{\pm} > 1 \text{ GeV}$ & $E^{\gamma} > 1 \text{ GeV}$ for any of the ISR photons
 - at least one ME photon with $p_T^{\gamma} > 2 \text{ GeV } \& 5^{\circ} < \theta^{\gamma} < 175^{\circ}$ $p_T^{\gamma} > 5 \text{ GeV } \& 7^{\circ} < \theta^{\gamma} < 173^{\circ}$

DELPHES framework used for detector simulation and event reconstruction.

Require:

• single photon with

 $p_T^{\gamma} > 3 \ GeV \ \& \ |\eta^{\gamma}| < 2.8 \ (ILC) \ p_T^{\gamma} > 10 \ GeV \ \& \ |\eta^{\gamma}| < 2.6 \ (CLIC)$

- no other activity in the detector other reconstructed objects
 - no electrons

• no jets (ILC 500 GeV)

(CLIC 3 TeV)

- no LumiCal photons
- no BeamCal photons



Background distributions

Two SM backgrounds considered: with up to 3 ME photons Bhabha scattering and (radiative) neutrino pair production





Background vs Signal distributions



For mono-photon events, two variables fully describe event kinematics \Rightarrow use 2D distribution of (p_T^{γ}, η) to constrain DM production Background Signal



- ji **F**...

Histogram binning

Cross section limits for vector mediator exchange, M_Y =2TeV, Γ/M =0.1 as a function of the number of histogram bins:



Signal MC statistics: 100k (open symbols) and 1M (full symbols)



Systematic uncertainties

following ILD study: Phys. Rev. D 101, 075053 (2020), arXiv:2001.03011

Considered sources of uncertainties:

- Integrated luminosity uncertainty of 0.26% uncorrelated between polarisations
- Luminosity spectra shape uncertainty correlated between polarisations
- Uncertainty in neutrino background normalisation of 0.2% (th+exp) correlated between polarisations
- Uncertainty in Bhabha background normalisation of 1% (th+exp) correlated between polarisations
- Uncertainty on beam polarisation of 0.02–0.08% (ILC)/0.2% (CLIC) correlated for runs with same beam polarisation at ILC

 \Rightarrow nuisance parameters in the RooStat model fit (11 for ILC, 7 for CLIC)

Systematic uncertainties



Cross section limits for radiative events (with tagged photon) Vector Mediator $\Gamma/M = 0.03$ with and without systematics ILC @ 500 GeV CLIC @ 3 TeV [fb] [fb] Combined Combined (H-20) ILC CLIC Nea 4000 fb _R 1600 fb 500 GeV 3 TeV 10 10 RL 1600 fb⁻¹ Pos 1000 fb⁻¹



Systematic effects reduced for on-shell production of narrow mediator

Systematic uncertainties



Cross section limits for radiative events (with tagged photon) Vector Mediator $\Gamma/M = 0.5$ with and without systematics ILC @ 500 GeV CLIC @ 3 TeV [fb] [fb] Combined Combined (H-20) ILC CLIC Neg 4000 fb LR 1600 fb 500 GeV 3 TeV 10 Pos 1000 fb⁻¹



Systematic effects reduced for on-shell production of narrow mediator



Cross section limits for radiative events (with tagged photon) Impact of beam polarisation assuming 4 ab^{-1} for ILC @ 500 GeV

Vector mediator

Scalar mediator



Polarisation significantly reduces impact of systematic uncertainties...

Impact of constraints

How important are the external constraints on the systematic effects? Eg. precision of the lumionosity measurement or theoretical calculations...

Radiative cross section limits for Vector mediator $M_Y=2$ TeV, $\Gamma/M=0.03$

Single uncertainty only

All systematics (one varied)



Most of the systemactic effects are constrained by the data itself!!!

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Experimental-like approach

Cross section for DM pair-production via mediator exchange depends on

- mediator mass, M_{γ} , and DM mass m_{χ}
- SM-mediator coupling value, g_{eeY} and coupling structure \mathcal{O}_{eeY}
- DM-mediator coupling value, $g_{\chi\chi\gamma}$ and coupling structure $\mathcal{O}_{\chi\chi\gamma}$



Experimental-like approach

Cross section for DM pair-production via mediator exchange depends on

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The cross section can also be expressed in terms of the widths

$$\sigma(e^+e^- \to Y \to \chi\chi) = \frac{12\pi}{M_Y^2} \frac{s \,\Gamma_{ee} \,\Gamma_{\chi\chi}}{(s - M_Y^2)^2 + M_Y^2 \Gamma_Y^2}$$

In the limit $\Gamma_{ee} \ll \Gamma_{\chi\chi} \approx \Gamma_Y$ the cross section depends only on M_Y , Γ_Y , g_{eeY} and \mathcal{O}_{eeY} (dependence on m_{χ} , $g_{\chi\chi Y}$ and $\mathcal{O}_{\chi\chi Y}$ "absorbed" in Γ_Y)

⇒ study limits on the DM pair-production cross section (or g_{eeY}) as a function of the mediator mass and width (for given \mathcal{O}_{eeY})

Results

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Light mediator exchange



Cross section limits for total DM production cross section Corrected for probability of hard photon tagging!

Combined limits for Vector mediator

ILC @ 500 GeV

CLIC @ 3 TeV



Radiation suppressed for narrow mediator with $M_Y \sim \sqrt{s} \Rightarrow$ weaker limits

Light mediator exchange



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Cross section limits for total DM production cross section Corrected for probability of hard photon tagging!

Combined limits for mediators with $\Gamma/M = 0.03$

ILC @ 500 GeV

CLIC @ 3 TeV



- **F**

Coupling limits

Combined limits for Vector mediator



Almost uniform sensitivity to mediator coupling g_{eeY} up to kinematic limit.

ILC @ 500 GeV

Coupling limits

Combined limits for mediators with $\Gamma/M = 0.03$

g_{ee7} (95%CL) (95%CL ILC 10 CLIC 4 ab⁻¹ @ 500 GeV 5 ab⁻¹ @ 3 TeV /-A coupling V-A coupling $g_{\rm ee\gamma}$ V+A coupling V+A coupling 10⁻¹ 10⁻² 10⁻² 10⁻³ 10³ 10² 10² 10³ 10⁴ M_v [GeV] M_v [GeV]

CLIC @ 3 TeV

Almost uniform sensitivity to mediator coupling g_{eeY} up to kinematic limit.

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Light mediator exchange

Prospects for SM decay observation

Limits on the SM-mediator couplings can be translated into limits on the mediator branching ratio to charged leptons

ILC @ 500 GeV

CLIC @ 3 TeV





Light mediator exchange

Prospects for SM decay observation

Limits on the SM-mediator couplings can be translated into expected numbers of produced lepton pairs



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Effective mass scale limits $\Lambda^2 = \frac{M_Y^2}{|g_{eeY}g_{yyY}|}$

Combined limits for Vector mediator

ILC @ 500 GeV

CLIC @ 3 TeV



Light mediator exchange



Effective mass scale limits

 $\Lambda^2 = \frac{\mathsf{M}_Y^2}{|\mathsf{g}_{eeY}\mathsf{g}_{\chi\chi Y}|}$

Combined limits for mediators with $\Gamma/M=0.03$

ILC @ 500 GeV

CLIC @ 3 TeV





arXiv:2001.03011

Comparison with ILD study Effective mass scale limits: $\Lambda^2 = \frac{M_Y^2}{|g_{eeY}g_{YYY}|}$

Limits from fast simulation (points) vs limits from full simulation (lines)



Very good agreement between full simulation and fast simulation results! ⇒ reliable extrapolation to low mediator mass domain...

Conclusions

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Sensitivity of future linear e^+e^- colliders to processes of dark matter production with light mediator exchange

arXiv:2107.11194, in print

Mono-photon signature: the most general way to look for DM production, EFT sensitivity extending to the O(10) TeV mass scales

New framework for mono-photon analysis developed focus on light mediator exchange and very small mediator couplings to SM

• $\mathcal{O}(1\,{
m fb})$ limits on the radiative production $e^+e^- o \chi\chi\gamma_{
m tag}$

• $\mathcal{O}(10 \text{ fb})$ limits on the DM pair-production $e^+e^- \rightarrow \chi \chi(\gamma)$ except for the resonance region $M_Y \sim \sqrt{s}$

• $\mathcal{O}(10^{-3}-10^{-2})$ limits on the mediator coupling to electrons up to the kinematic limit $M_Y \leq \sqrt{s}$

For light mediators limits more stringent than those expected from direct resonance search in SM decay channels

Other options for probing Dark Matter

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Probing Dark Matter with e⁺e⁻

Dark Matter

Many hints for existence of Dark Matter (DM), but its nature is unknown. Many possible scenarios, wide range of masses and couplings to consider.

High energy e^+e^- machines offer many options for DM searches:





Tomohiko Tanabe @ LCWS'2021

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First ILC running stage will clearly be focused on Higgs measurements **Production cross section**



At 250 GeV dominated by Higgs-strahlung (ZH production)



but we still profit from combining two production channels ⇒ model independent analysis

Invisible decays

High sensitivity to invisible Higgs boson decays with recoil mass technique



Expected 95% C.L. limit for 2 ab⁻¹ collected at 250 GeV ILC: 0.23% a factor of 10 better than the HL-LHC prospect. arXiv:2002.12048





Search for new scalars

Many BSM models introduce extended Higgs sectors. New scalars could be light, if their couplings to SM particles are small.



ILC search independent on the scalar decay: $e^+e^- \rightarrow Z \ S^0 \rightarrow \mu^+\mu^- + X$



Soft SUSY scenarios

Thanks to clean environment, sensitivity of e^+e^- colliders extends down to very small NLSP-LSP mass differences



arXiv:2105.06408

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Search for IDM scalar production

Production of IDM scalars at e^+e^- colliders dominated by two processes:

 $e^+e^- \rightarrow A H$

Search for *AH* production leptonic channel

 $e^+e^-
ightarrow H^+H^-$

Search for H^+H^- production semi-leptopnic channel





arXiv:2107.13803

arXiv:2002.11716

ILC beam dumps



ILC beam dumps



Concept of main beam dump experiments searching for axion-like particles or new scalars:







Looking for SM decays of new exotic particles produced in the beam dump arXiv:2009.13790



Axion-like particle model looking for $a \to \gamma\gamma$ $\mathcal{L} \ni -\frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu} + \frac{1}{2}(\partial_{\mu}a)^2 - \frac{1}{2}m_a^2a^2$

An order of magnitude better sensitivity than other experiments



Looking for SM decays of new exotic particles produced in the beam dump arXiv:2009.13790



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Axion-like particle model looking for $a \rightarrow \gamma \gamma$

$$\mathcal{L} \ni -\frac{1}{4} g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{2} (\partial_{\mu} a)^2 - \frac{1}{2} m_a^2 a^2$$

An order of magnitude better sensitivity than other experiments

Light scalar coupled to charged leptons

$$\mathcal{L} \ni \frac{1}{2} (\partial_{\mu} S)^2 - \frac{1}{2} m_S^2 S^2 - \sum_{I=e,\mu,\tau} g_I S \overline{I} I$$

Model A: $g_I \propto m_I$ Sensitivity down to very small couplings



M.Perelstein @ LCWS'2021

Scenarios with Dark Photon (A') and Dirac fermion DM (χ)

$$\mathcal{L} \ni -\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \frac{1}{2} m_{\mathcal{A}'}^2 \mathcal{A}'_{\mu} \mathcal{A}'^{\mu} - \frac{\epsilon}{2} F'_{\mu\nu} F^{\mu\nu} + \bar{\chi} (\imath \not D - m_{\chi}) \chi$$

Resonant ($e^+e^- \rightarrow A'$), associated prod. ($e^+e^- \rightarrow A' \gamma$) or radiation ($e^{\pm} N \rightarrow e^{\pm} N A'$) \Rightarrow collimated stream of DM particles from A' decay (A' $\rightarrow \chi\chi$) \Rightarrow looking for elactic χ interactions in the detector

 \Rightarrow looking for elastic χ interactions in the detector



Approach used in SLAC Beam Dump Experiment E137 arXiv:1406.2698



M.Perelstein @ LCWS'2021

Scenarios with Dark Photon (A') and Dirac fermion DM (χ)

$$\mathcal{L} \ni -\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \frac{1}{2} m_{A'}^2 A'_{\mu} A'^{\mu} - \frac{\epsilon}{2} F'_{\mu\nu} F^{\mu\nu} + \bar{\chi} (\imath \not D - m_{\chi}) \chi$$

Resonant ($e^+e^- \rightarrow A'$), associated prod. ($e^+e^- \rightarrow A' \gamma$) or radiation ($e^{\pm} N \rightarrow e^{\pm} N A'$) \Rightarrow collimated stream of DM particles from A' decay ($A' \rightarrow \chi \chi$)

 \Rightarrow looking for elastic χ interactions in the detector



Non-collider experiments



Experiments with extracted beams

M.Perelstein @ LCWS'2021

Searching for Dark Photons with extracted positron beams

 ${\rm e^+e^-} \to {\rm A'}\,\gamma$

Missing energy reconstruction in thick active target



LDMX for SLAC: arXiv:1807.05884

Thin target, missing mass reconstruction in dedicated detector



PADME @ Frascati: arXiv:1910.00764

Sensitivity extending down to the minimum couplings allowed by relic density bounds



Probing Dark Matter with ILC

ILC will offer many complementary options for DM searches.

- Different scenarios can be constrained via precision Higgs studies.
- Clean environment and kinematic constraints of e^+e^- collisions result in high sensitivity to different DM production scenarios.
- Sensitivity extends to the TeV mass scales, order of magnitude higher than the collision energy.

The ILC will also offer highest energy electron and positron beams, with unprecedented intensities, for beam dump and extracted beam exp.

Fixed-target experiments offer many interesting opportunities for dark sector searches in the low mass domain and other science goals.

Thank you!

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Simplified DM model

Dark matter particles, X_i , couple to the SM particles via an mediator, Y_j .

Each simplified scenario is characterized by one dark matter candidate and one mediator from the set listed below:

	particle	mass	spin	charge	self-conjugate	type
DM	X _R	m_{X_R}	0	0	yes	real scalar
	X _C	m _{Xc}	0	0	no	complex scalar
	X_M	m_{X_M}	$\frac{1}{2}$	0	yes	Majorana fermion
	X _D	m _{X_D}	$\frac{\overline{1}}{2}$	0	no	Dirac fermion
	X_V	m_{X_V}	1	0	yes	real vector
mediator	Y _R	m _{Y_R}	0	0	yes	real scalar
	Y_V	m _{Yc}	1	0	yes	real vector
	T _C	m _{Tc}	0	1	no	charged scalar