

# Quarkonium azimuthal anisotropy at forward rapidity in Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV with ALICE



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148th LHCC Meeting, virtual poster session, CERN, Nov. 18, 2021

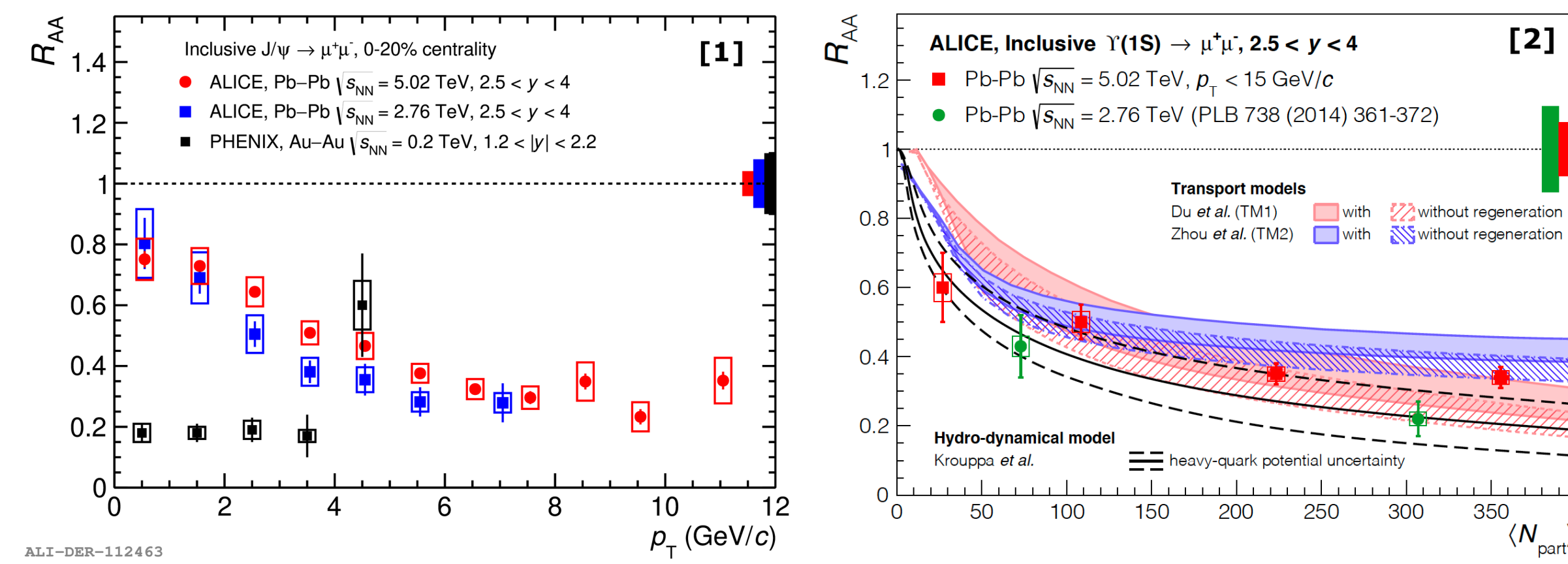
## 1 – Physics motivations

Charm and beauty created early in collisions: excellent probes of the deconfined medium. **Quarkonium** states (like  $J/\psi$ :  $c\bar{c}$ , or  $\Upsilon(1S)$ :  $b\bar{b}$ ) undergo:

- **suppression**, originally due to color screening with surrounding charges
- **regeneration**, recombination of thermalized heavy quarks in the medium

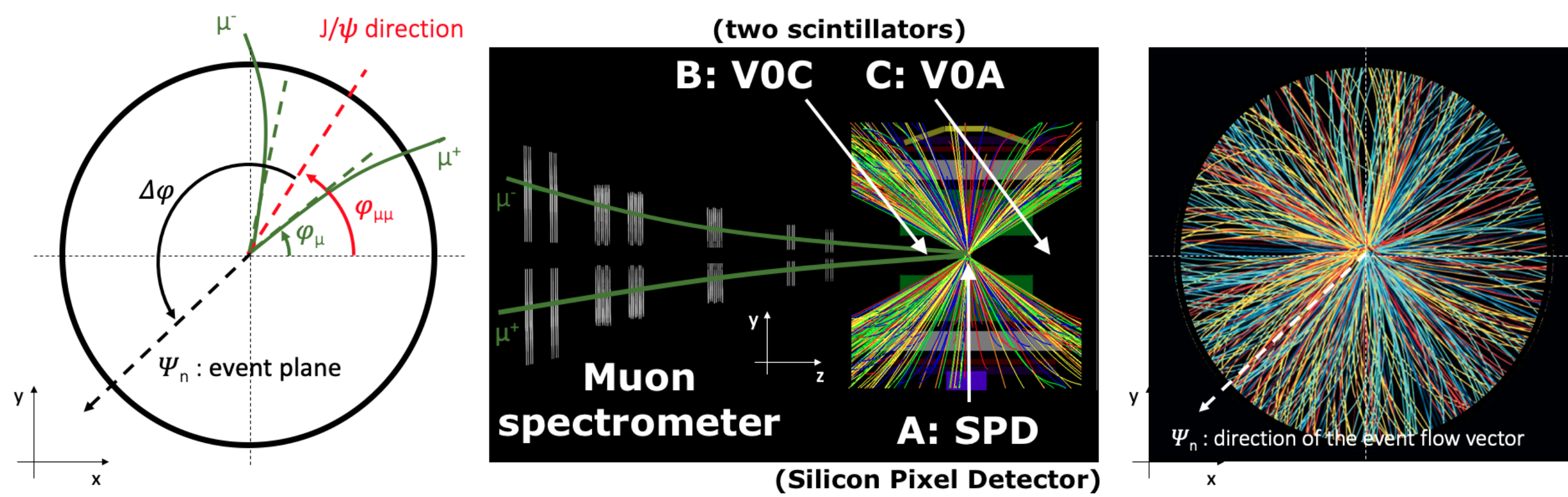
Quantify by  $R_{AA}$  ( $> 1$  enhancement,  $< 1$  suppression, Pb–Pb w.r.t. pp collisions)

- Inclusive  $J/\psi$  mainly regenerated at low  $p_T$  while strongly suppressed at high  $p_T$  [1]
- $\Upsilon(1S)$  suppressed in Pb–Pb (but large fraction from feed-down) [2]



## 3 – ALICE experimental setup, muon and dimuon tracks reconstruction

- V0A, V0C, SPD: multiplicity, event flow vector  $Q_n = \sum_j e^{in\varphi_j}$ , vertex
- **muon spectrometer**: quarkonia via dimuon decay channel (forward rapidity)



## 4 – Analysis details, signal extraction and $v_n$ measurement

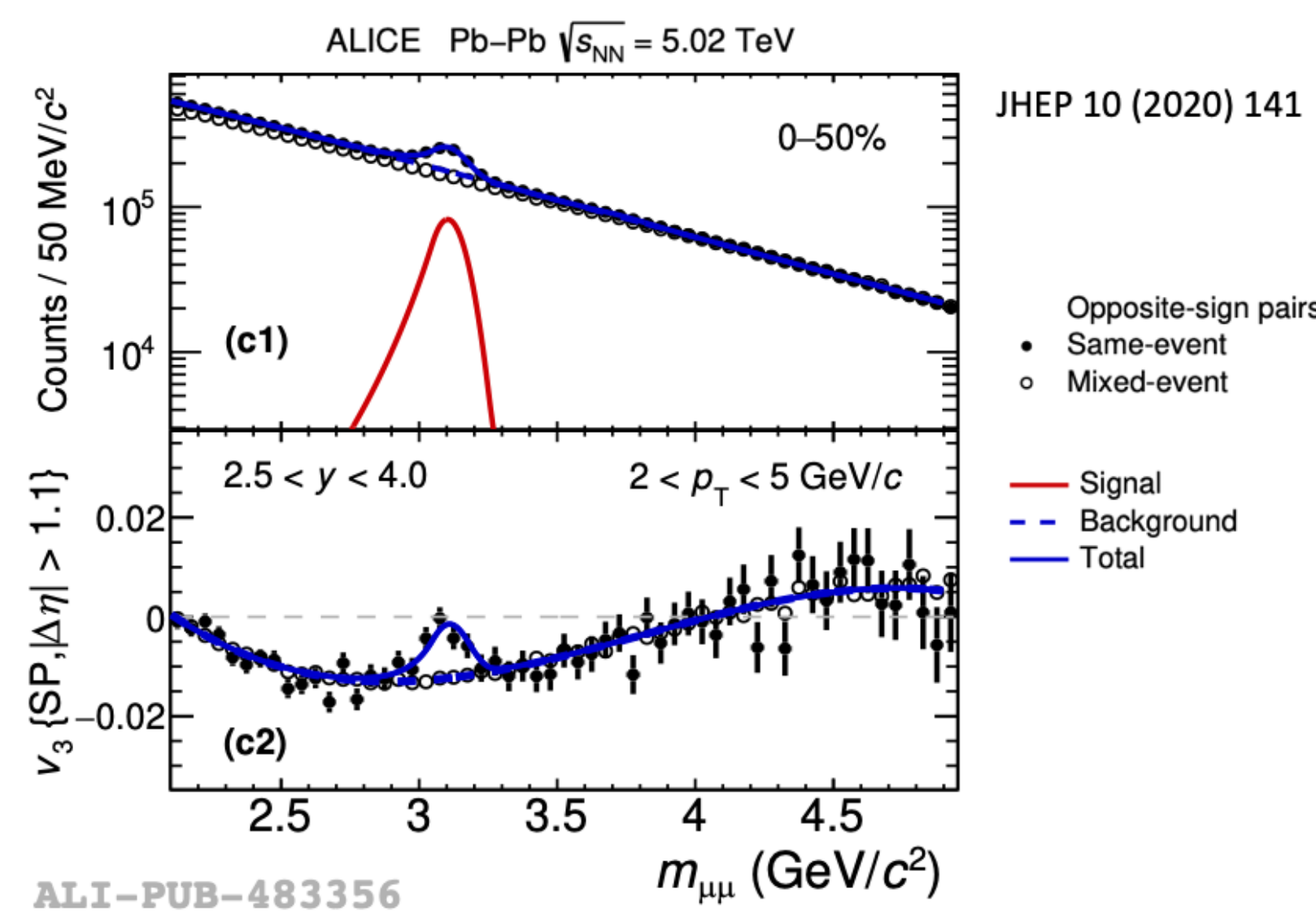
Scalar product (SP) method, three sub-event technique (A,B,C sub-detectors: SPD, V0A, V0C): suppress non-flow effects [4]:

$$v_n\{SP\} = \left\langle \frac{\mathbf{u}_n \mathbf{Q}_n^A}{\sqrt{\frac{\langle \mathbf{Q}_n^A \mathbf{Q}_n^{B*} \rangle \langle \mathbf{Q}_n^A \mathbf{Q}_n^{C*} \rangle}{\langle \mathbf{Q}_n^B \mathbf{Q}_n^{C*} \rangle}}} \right\rangle_{\mu\mu}$$

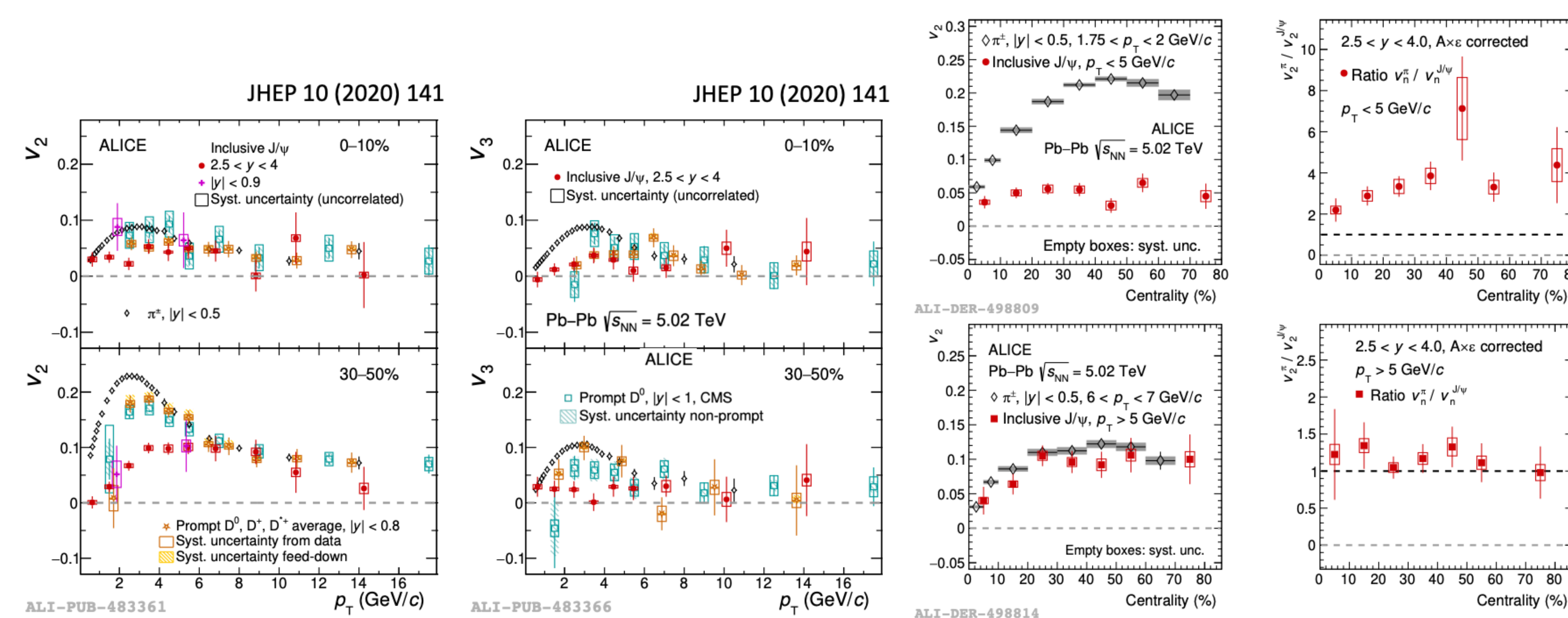
$\mathbf{u}_n = e^{in\varphi}$ ,  $\langle \dots \rangle$ : average over dimuons from all events. **signal** + background (dashed blue): extracted from invariant mass distribution, through  $\alpha = S/(S+B)$ .

$v_n^{\text{sig}}$  ( $v_n^{J/\psi}$  or  $v_n^{\Upsilon(1S)}$ ): obtained from dimuon  $v_n(m_{\mu\mu})$  fit

$$v_n(m_{\mu\mu}) = \alpha(m_{\mu\mu}) v_n^{\text{sig}} + [1 - \alpha(m_{\mu\mu})] v_n^{\text{bkg}}(m_{\mu\mu}).$$

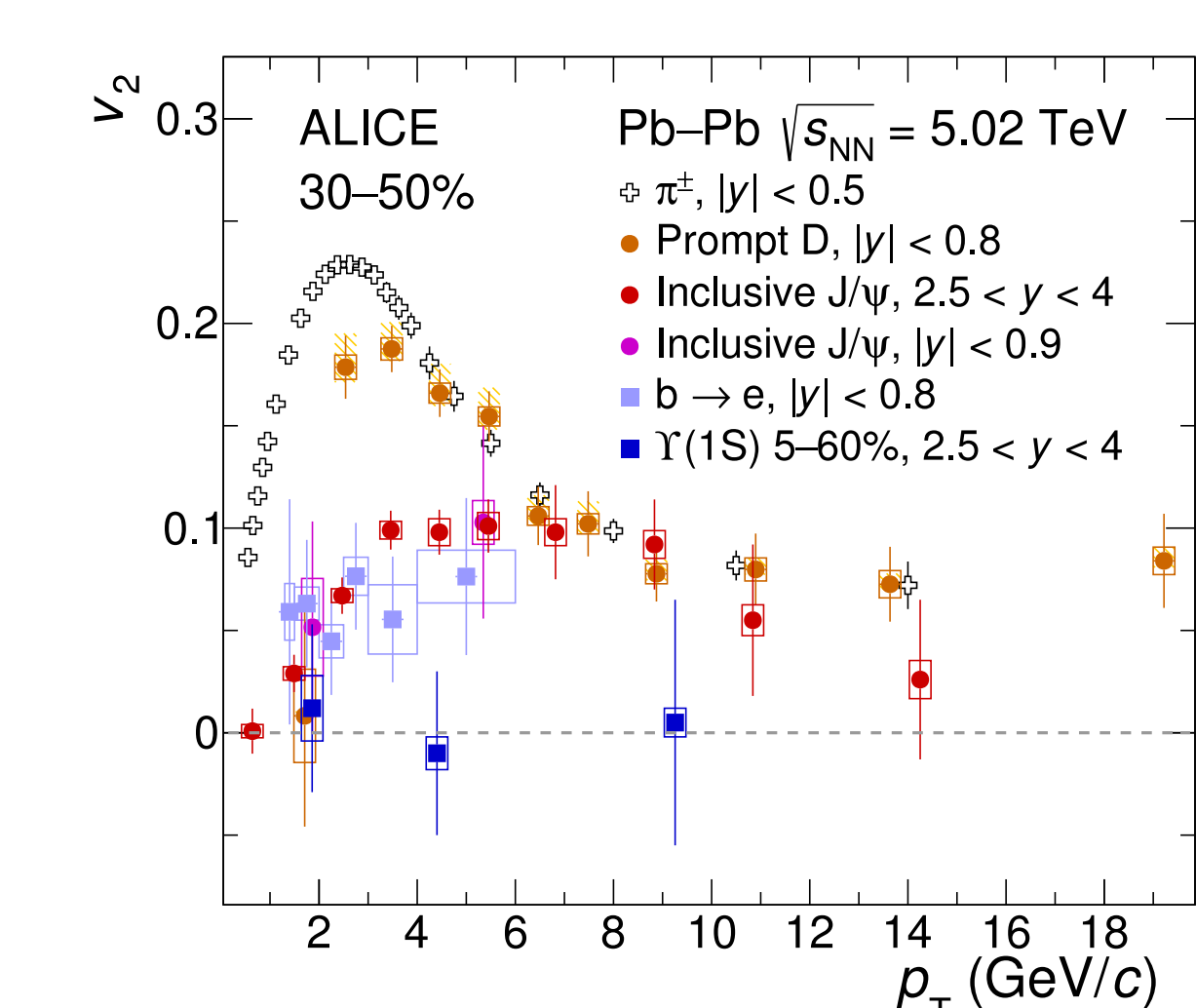


## 5 – Results: $p_T$ and centrality dependence



- Mass hierarchy at low  $p_T$  for  $\pi$ ,  $D$ , and  $J/\psi$   $v_n$ , while path-length dependent effects dominate at high  $p_T$
- Low  $p_T$   $v_2^\pi/v_2^{J/\psi}$  increases from unity toward peripheral collisions, while at high  $p_T$  it is compatible with 1

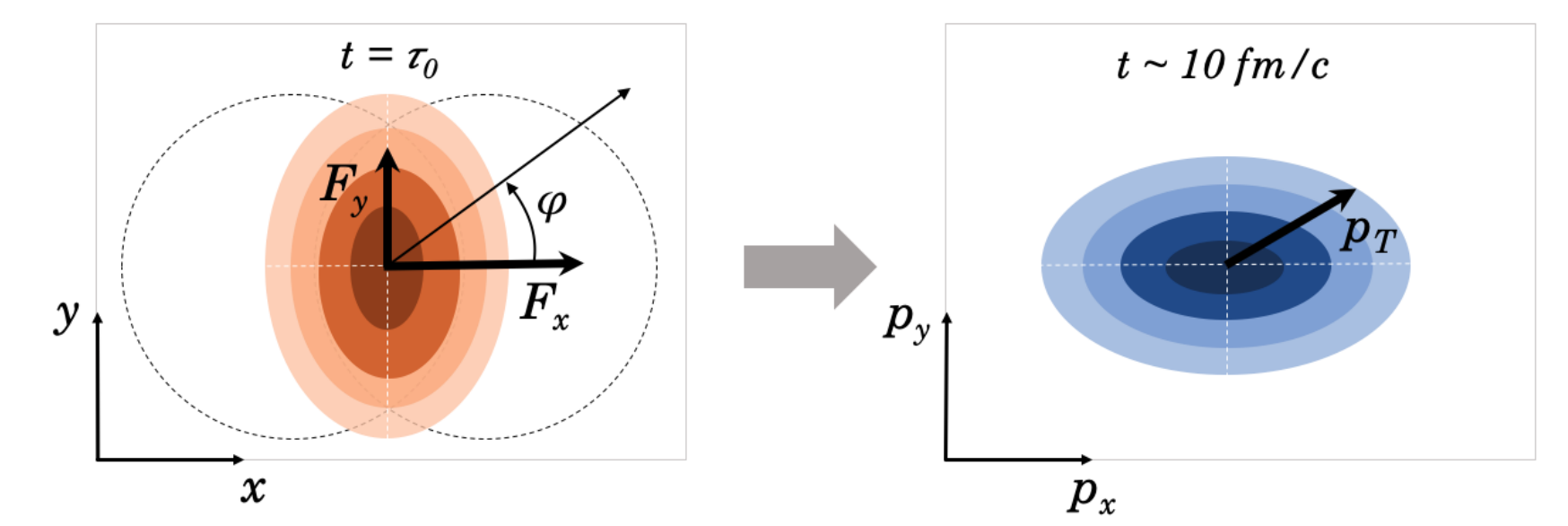
## 6 – Global picture: from light, to charm, to beauty flavor, and summary



- $v_2^{J/\psi}$  favors scenario that all initially created  $J/\psi$  dissociate in the medium, recombination of charm quarks dominates the  $J/\psi$  production [5]
- $v_3^{J/\psi} > 0$  in 0–50%,  $2 < p_T < 5$  GeV/c ( $5.1\sigma$ ), no centrality dependence
- Low  $p_T$   $v_2^\pi/v_2^{J/\psi}$  increases from central to peripheral collisions, compatible with the scenario where  $c$  quarks thermalize later than the light ones
- $v_2^{\Upsilon(1S)}$ : agreement with all scenarios:  $b\bar{b}$  bound state formation limited to first stages (transport [6]), anisotropic escape mechanism (hydro [7]), or based from a full thermalization (statistical hadronization model [8])
- Global ordering at low  $p_T$ :  $v_2^\pi > v_2^D > v_2^{J/\psi} \gtrsim v_2^{b \rightarrow e} > v_2^{\Upsilon(1S)}$

## 2 – Azimuthal anisotropy

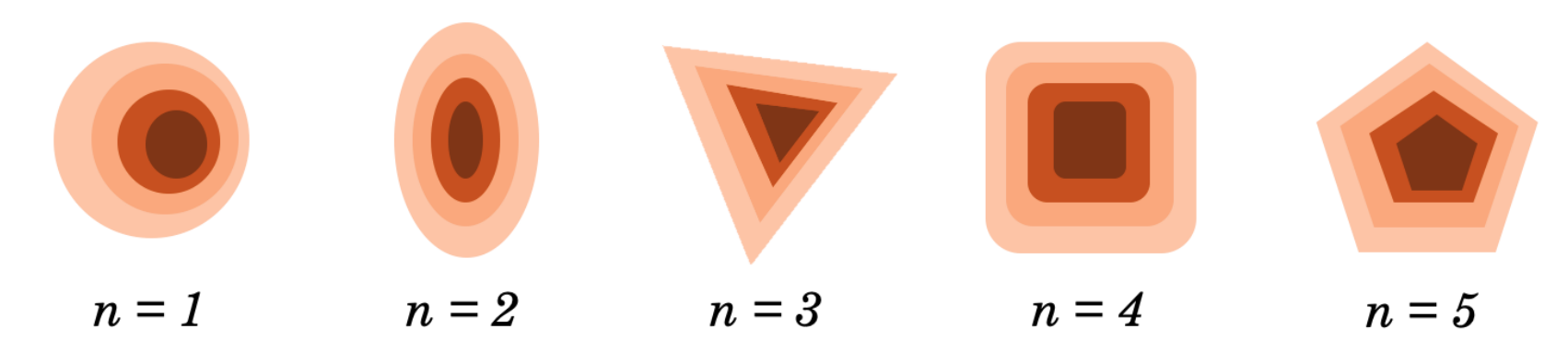
In non-central heavy-ion collisions, the nuclear overlap region has a **spatial anisotropy**. During system expansion, pressure gradients generated by parton interactions, transform the spatial anisotropy into a **momentum space anisotropy** of produced particles [3].



Azimuthal distributions of final state particles can be decomposed by Fourier series (with  $n$ : harmonics):

$$\frac{dN}{d\varphi} \propto 1 + 2 \sum_{n=1}^{+\infty} v_n \cos[n(\varphi - \Psi_n)],$$

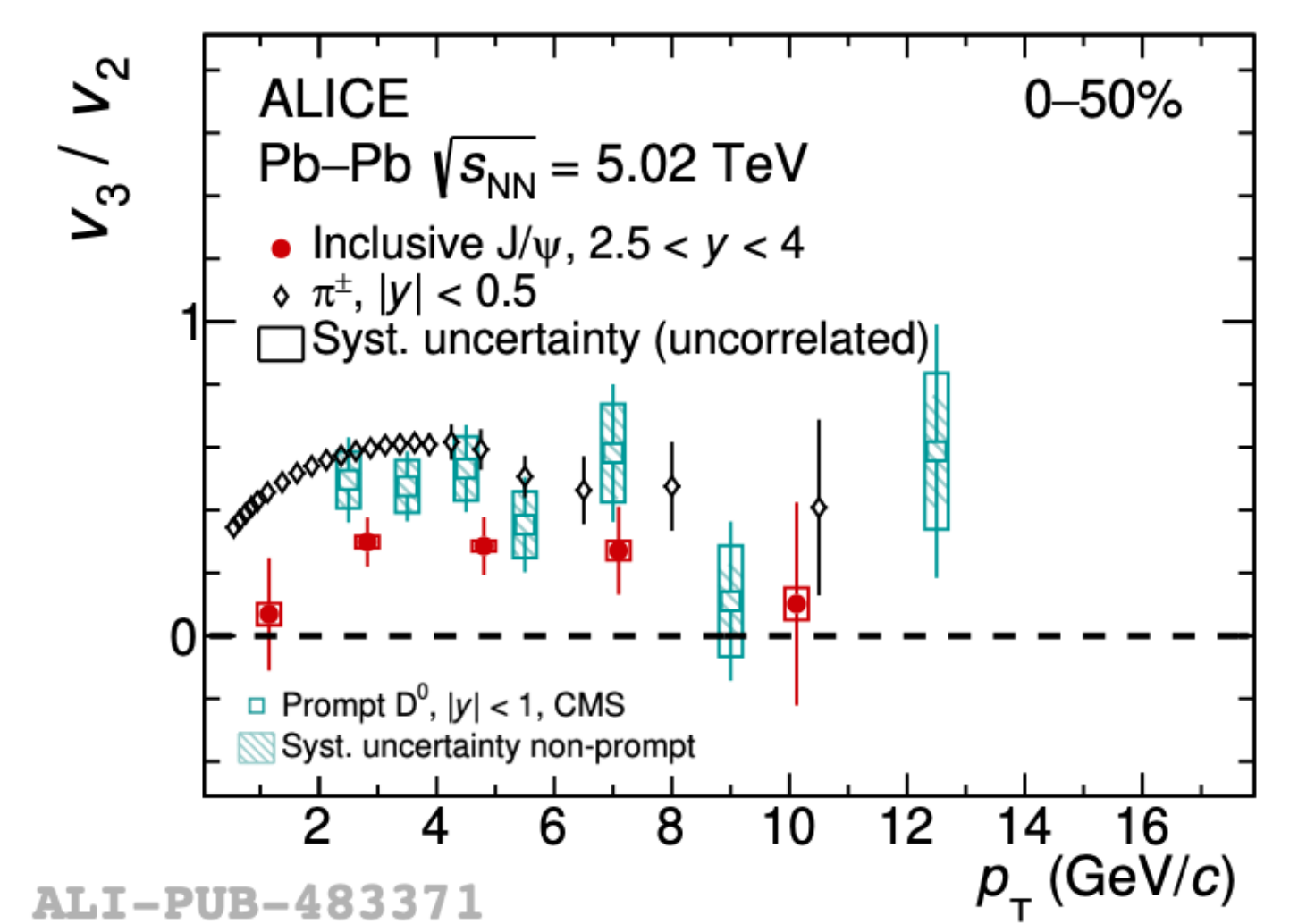
where  $\varphi$ : azimuthal angle, and  $\Psi_n$ : symmetry planes.  $v_n$  coefficients: sensitive to initial conditions, equation of state, and medium transport properties.



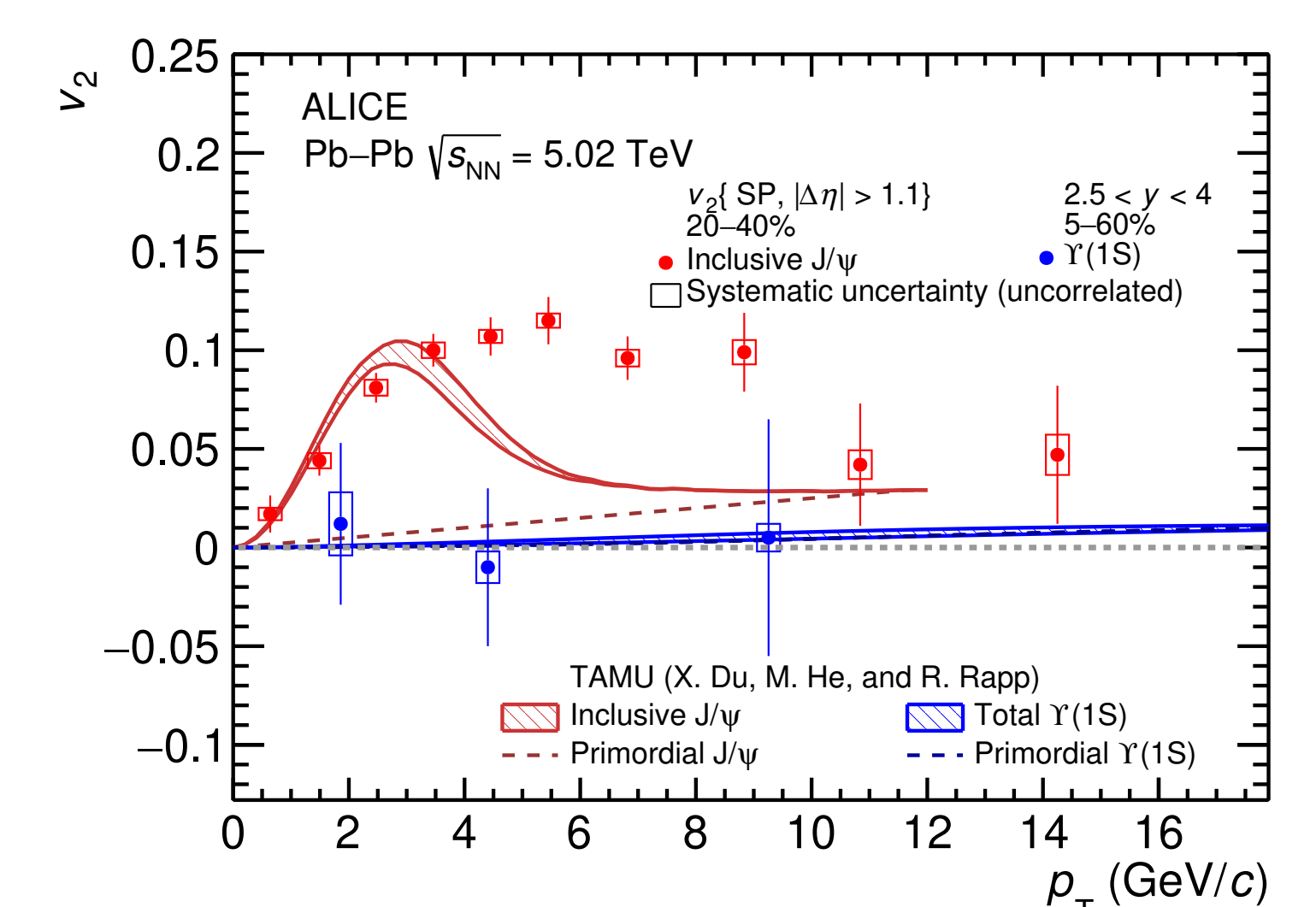
$v_2$  is called **elliptic flow**, and  $v_3$  **triangular flow**,...

## 5 – Results: $v_3/v_2$ and comparison with models

$v_3/v_2$  exhibits the same hierarchy than single flow coefficients indicating that higher harmonics are damped for charmonia compared to lighter particles



Precise  $v_2^{J/\psi}$  measurement which extend  $p_T$  coverage (w.r.t. previous data) and first  $v_2^{\Upsilon(1S)}$  measurement. Low  $p_T$ : agreement with microscopic transport models (strong regeneration for  $J/\psi$  [5], while small for  $\Upsilon(1S)$  [6]), missing mechanisms for  $J/\psi$  at intermediate  $p_T$



$v_2^{\Upsilon(1S)}$ : compatible with 0 and model predictions [6–8]

## References

- [1] Phys. Lett. B 766 (2017) 212-224
- [2] Phys. Lett. B 790 (2019) 89
- [3] Phys. Rev. D 46 (1992) 229
- [4] Phys. Rev. C 87 (2013) 044907
- [5] Nuclear Physics A, 943 (2015) 147-158
- [6] Phys. Rev. C 96 (2017) 054901
- [7] Phys. Rev. C 100 (2019) 051901
- [8] Phys. Rev. C 101 (2020) 064905