tWZ production at the LHC in SMEFT

[2111.03080]

in collaboration with
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Outline

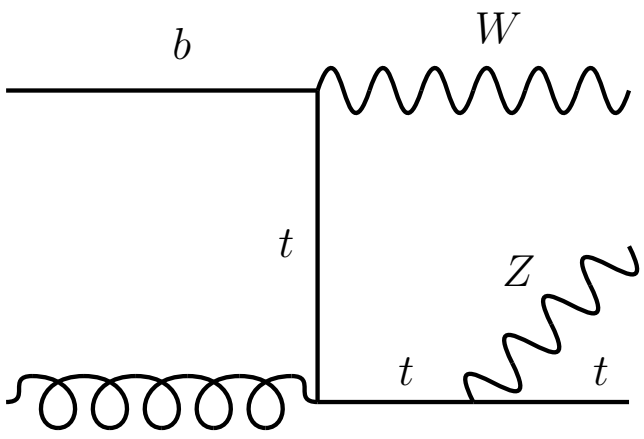
- Introduction
- Defining *tWZ* at NLO
- *tWZ* in the SM and SMEFT
- On NLOPS (NLO matched to Parton Shower)

Introduction

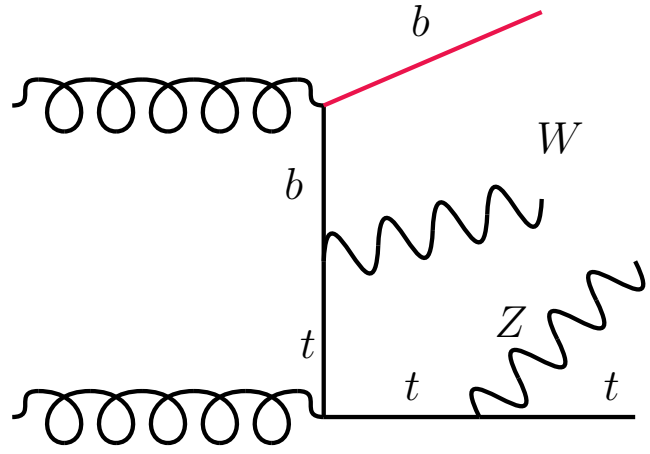
- *tWZ* is a rare EW process $\rightarrow$ **potential probe of EW couplings that are not well measured, so far**
- *tWZ* is sensitive to **unitarity-violating behaviour induced in its sub-amplitudes** via modified EW interactions [1904.05637] $\rightarrow$ can potentially serve as an avenue for new physics searches
- New physics phenomena can be encapsulated in a higher dimensional operators $\rightarrow$ **SMEFT**
- Accurate theoretical predictions$\rightarrow$ **study *tWZ* at NLO**
- ***tWZ* at NLO is non-trivial due to its overlap with other processes**

tWZ from LO to NLO in QCD

At NLO in QCD, the *tWZb* appears as a real emission final state

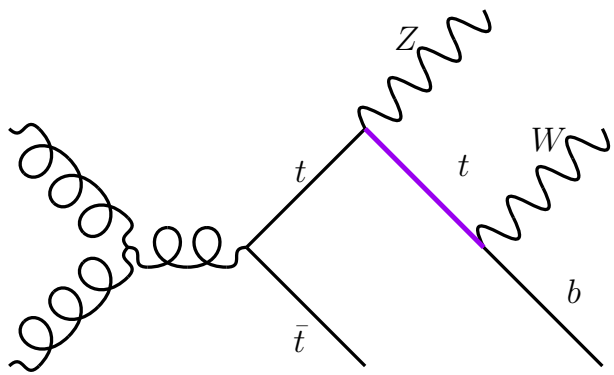


LO

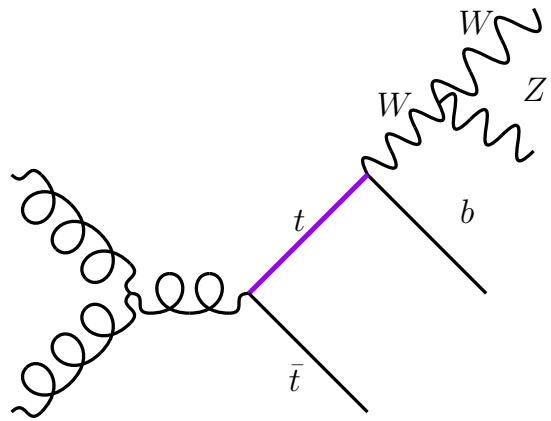


NLO

Defining *tWZ* at NLO



ttZ



tt

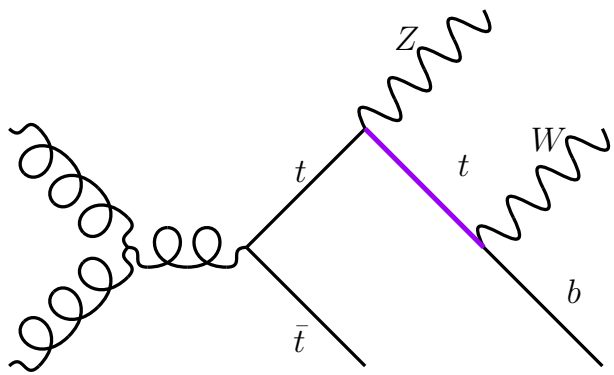
The ***tWZb*** final state can also have resonant contributions, from

$$\boldsymbol{ttZ, t \to Wb}, \quad \text{or} \quad \boldsymbol{tt, t \to WZb}$$

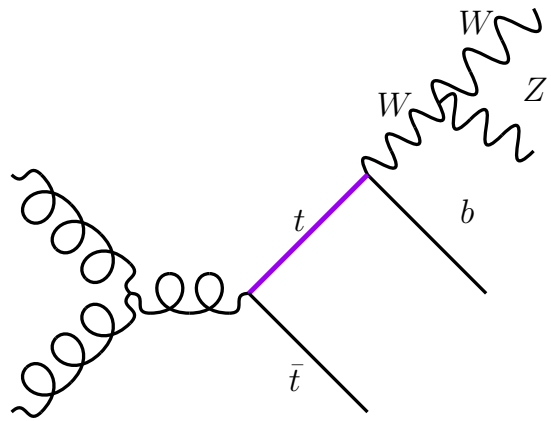
and not necessarily the non-resonant ***tWZb***

- These topologies **do not belong to the genuine *tWZ* final state process** but to the leading order ***ttZ*** and ***tt*** processes
- The underlying resonant structure can spoil the perturbative behaviour of the NLO expansion

Defining *tWZ* at NLO cont'd



ttZ



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Suppress contributions from the resonant amplitude!

Defining *tWZ* at NLO cont'd

The amplitude associated to the ***tWZ*** process can be expressed as

$$\mathcal{A}_{tWZ} = \mathcal{A}_{tWZ}^{r\not{e}s} + \mathcal{A}_{tWZ}^{res} \tag{1}$$

and thus the matrix element,

$$|\mathcal{A}_{tWZ}|^2 = |\mathcal{A}_{tWZ}^{r\not{e}s}|^2 + 2\Re\left(\mathcal{A}_{tWZ}^{r\not{e}s}\mathcal{A}_{tWZ}^{\dagger res}\right) + |\mathcal{A}_{tWZ}^{res}|^2 \tag{2}$$

Two **Diagram Removal (DR)** schemes to handle the resonant part of the matrix element:

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Two **Diagram Removal (DR)** schemes to handle the resonant part of the matrix element:

- Keep ONLY the non resonant contribution $\longrightarrow$ **DR1**

Defining *tWZ* at NLO cont'd

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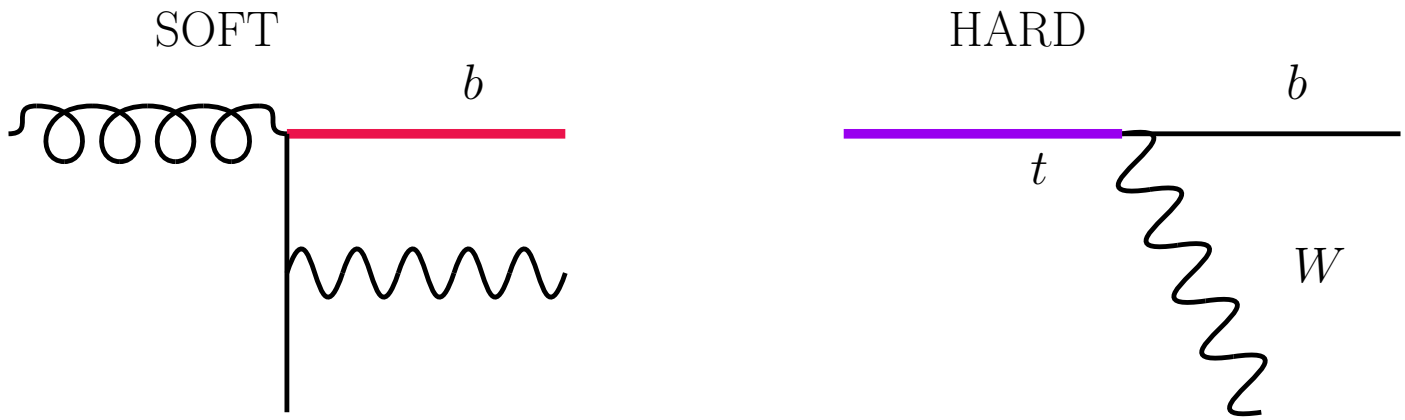
Two **Diagram Removal (DR)** schemes to handle the resonant part of the matrix element:

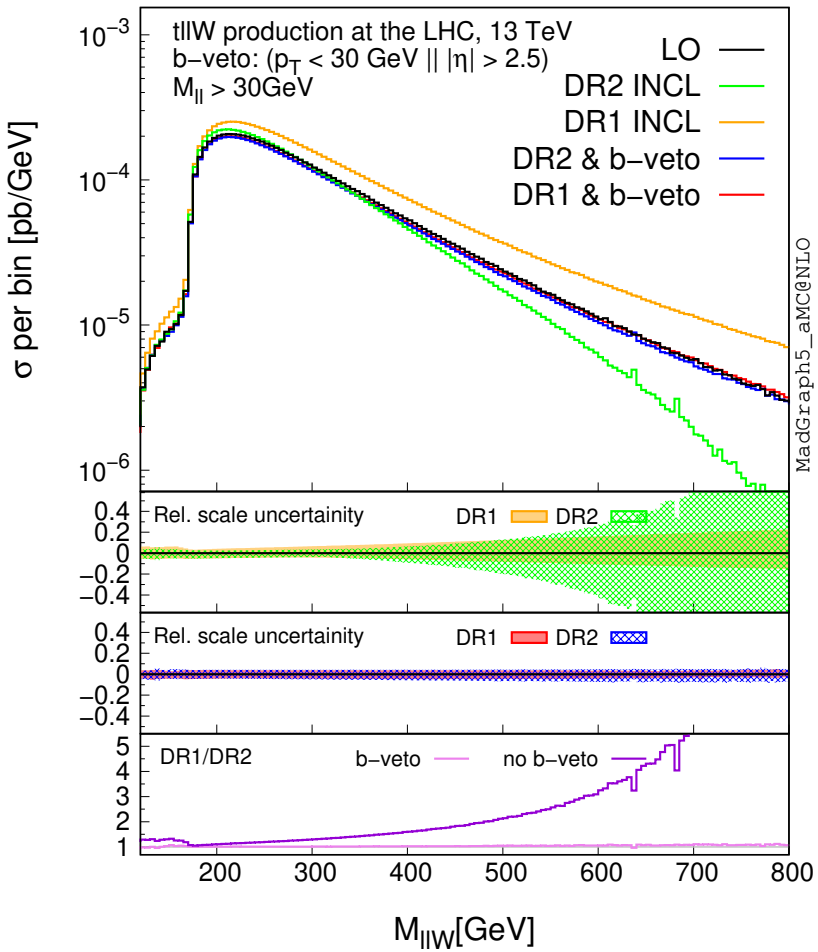
- Keep ONLY the non resonant contribution $\longrightarrow$ **DR1**
- Keep ALSO interference term $\longrightarrow$ **DR2**

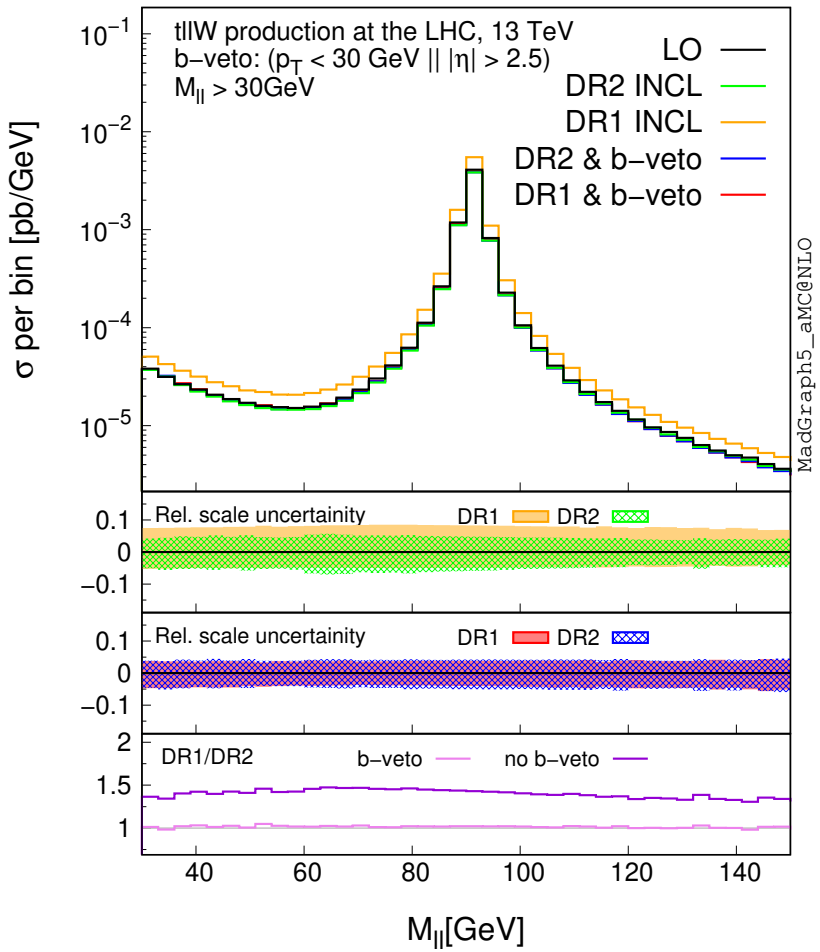
tWZ in the SM

- The current MG5 diagram removal plugin (MadSTR) does not handle $1 \to N$ decays where $N > 2 \to$ **DR1 and DR2 schemes are implemented by hand**
- The **resonant part of the phase space is suppressed by vetoing hard *b*-quarks** as they tend to have come from the decay of a top
- For SM predictions $\to$ both the ***ttZ*** and the ***tt*** overlaps are removed
- A good agreement between the DR1 and the DR2 schemes $\to$ **the non-resonant part of *tWZ* dominates the phase space**



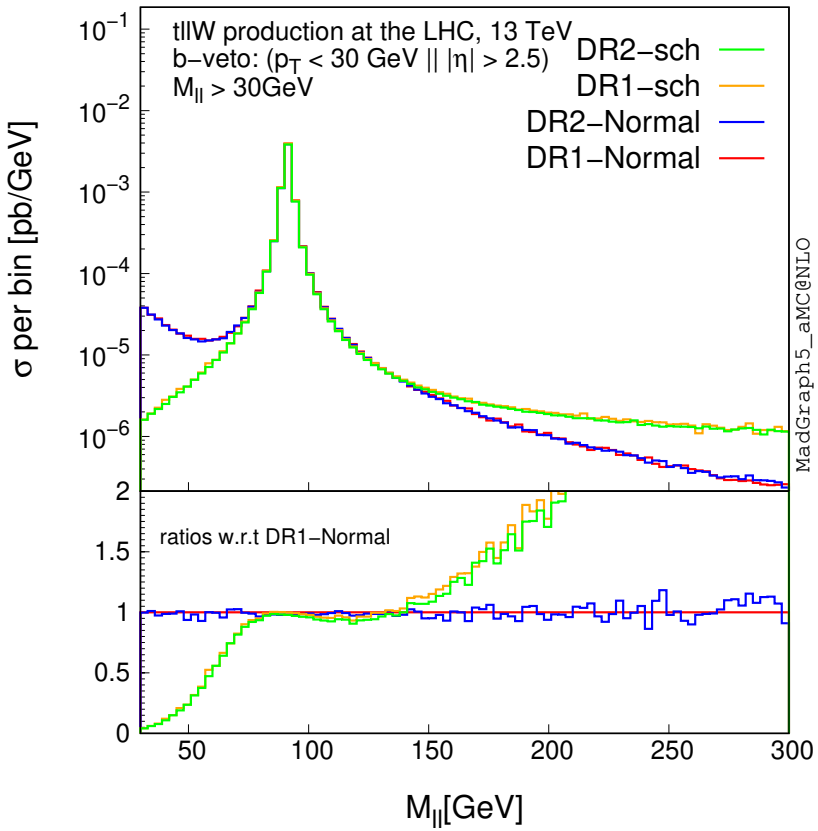
tWZ in the SM: differential predictions

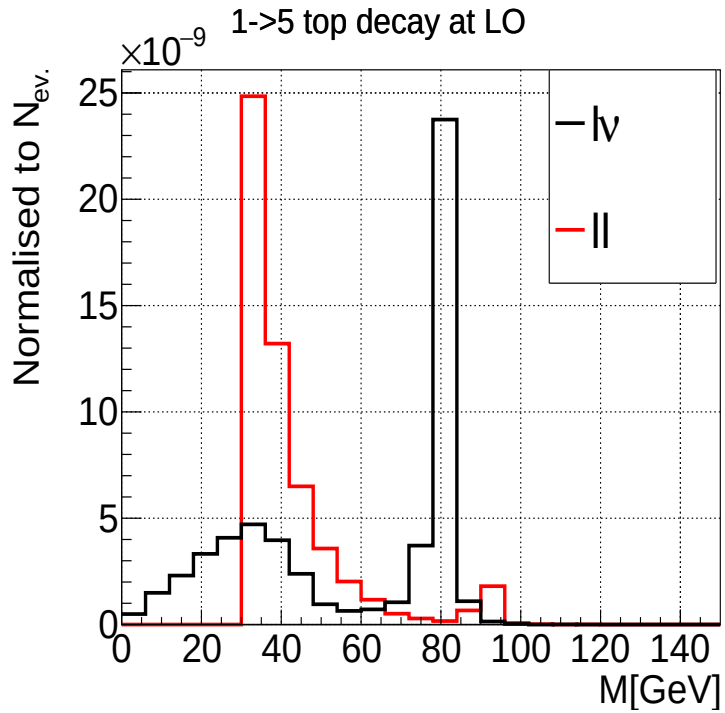




The *b*-veto significantly improves the DR1-DR2 agreement

Interlude: why can we keep the *Z* stable?

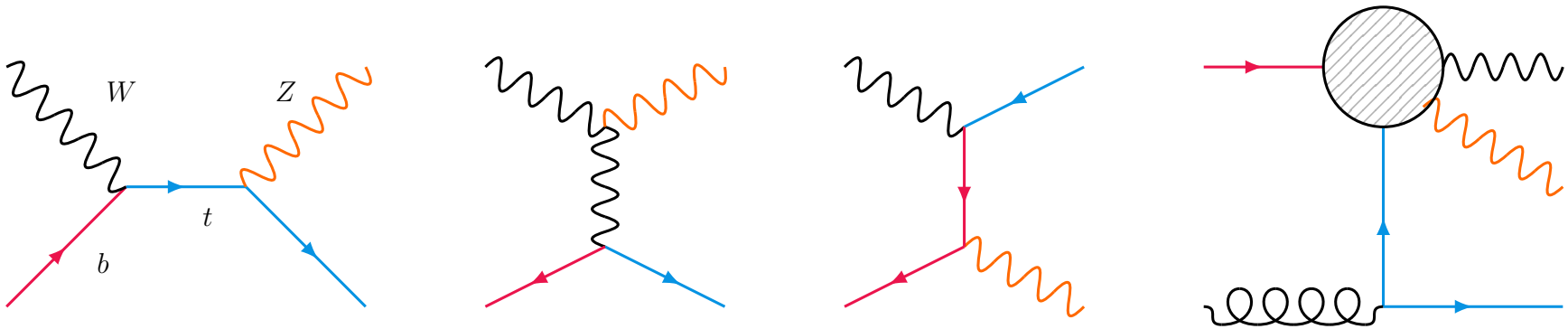




- Most of the *tWll* cross-section lies around the *Z*-peak
- $t\bar{t}$ overlap becomes irrelevant as $t \to WZb$ is suppressed
- In high-energy regions, $t\bar{t}$ overlap is non-resonant

tWZ in the SMEFT

The $bW \to tZ$ sub-amplitude:



and therefore the operators included in the analysis are

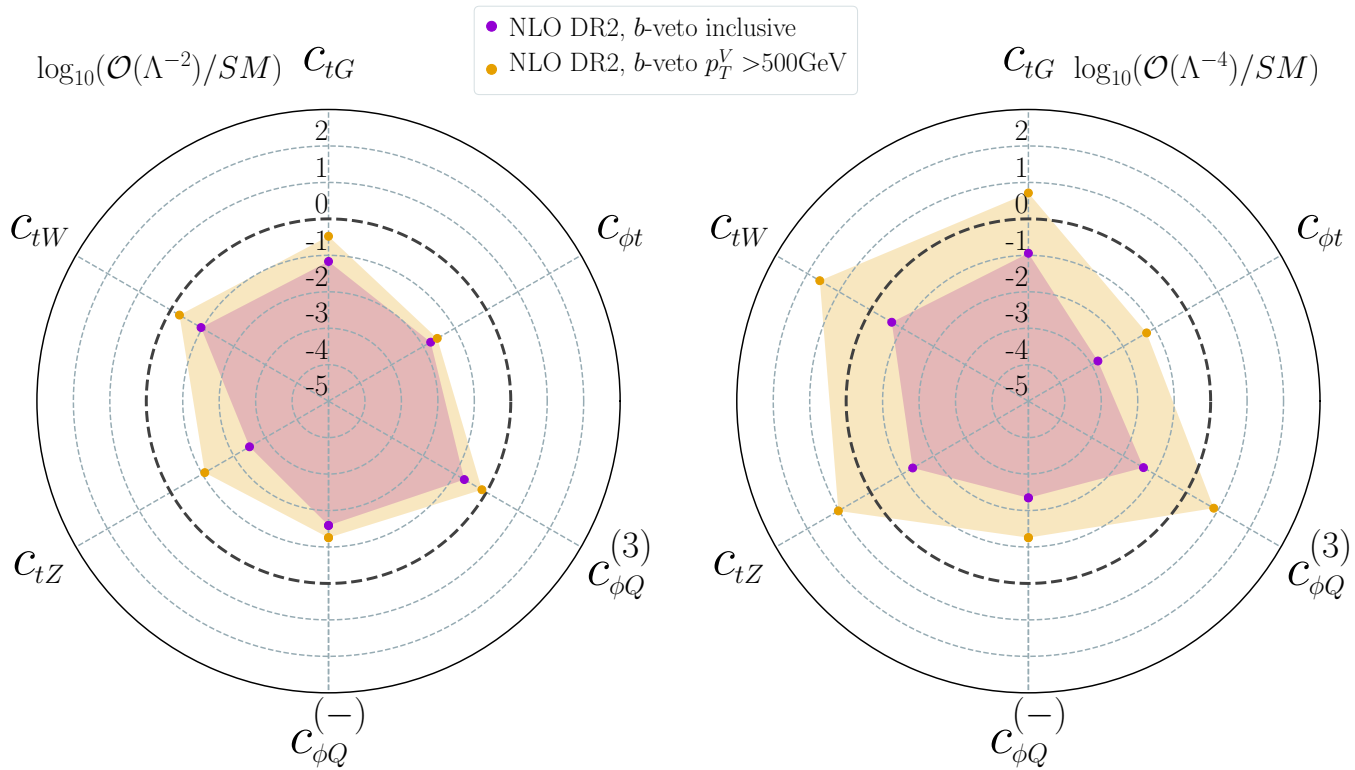
- $\mathcal{O}_{\varphi Q}^{(3)}$, $\mathcal{O}_{\varphi Q}^{(-)} \to$ interaction of two fermions fields with the gauge bosons
- $\mathcal{O}_{\varphi t} \to$ the right handed ***ttZ*** interaction
- $\mathcal{O}_{tW}$, $\mathcal{O}_{tZ} \to$ interaction of the top with the weak isospin and the weak hypercharge gauge fields
- ($\mathcal{O}_{tG} \to$ the gluon-top interaction)

SMEFT predictions

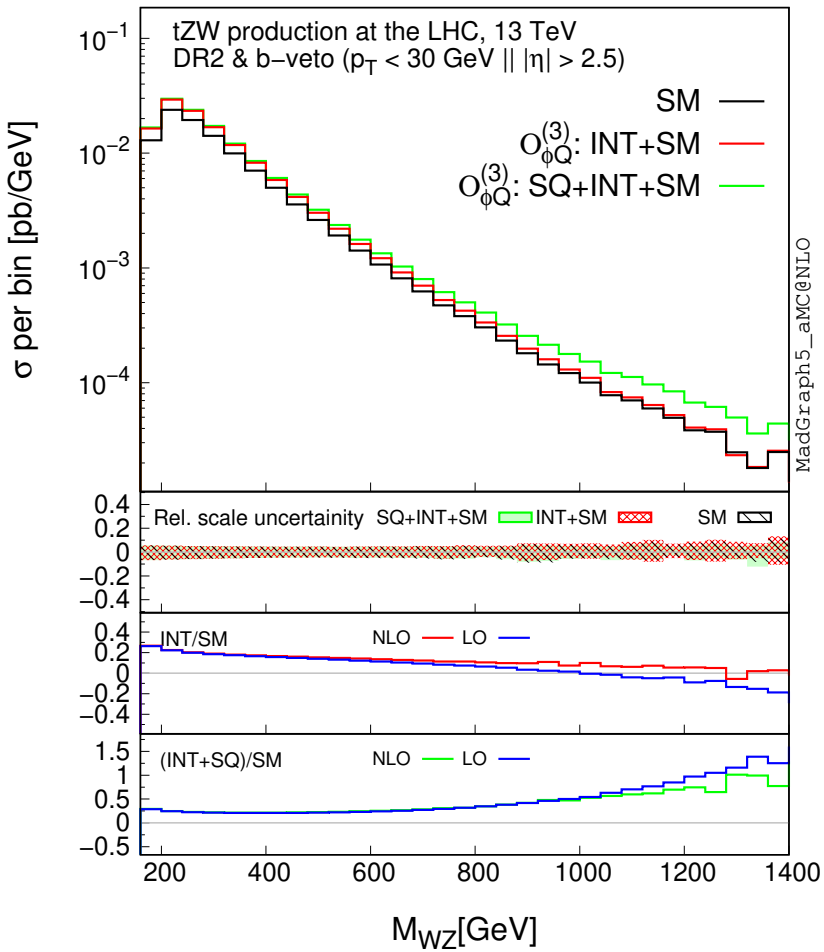
A generic observable in SMEFT can be expressed as

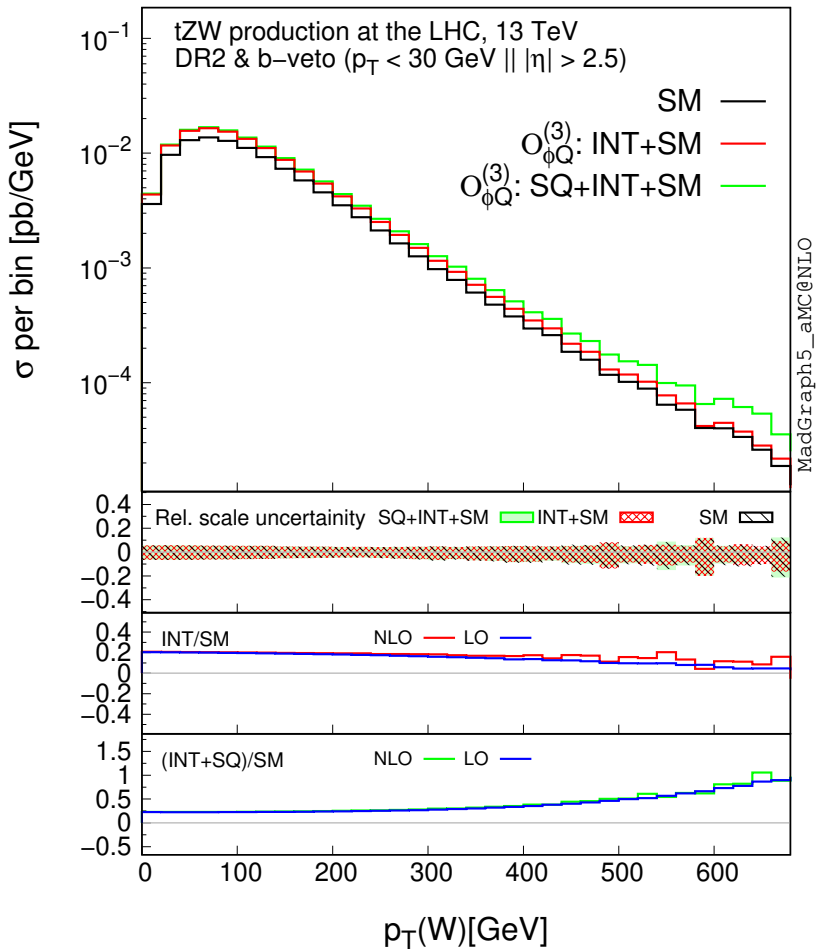
$$\sigma = \sigma_{SM} + \sum_i \frac{C_i^{(6)}}{\Lambda^2}\sigma_i + \sum_{ij} \frac{C_i^{(6)} C_j^{(6)}}{\Lambda^4}\sigma_{ij} \qquad (3)$$

- second term $\rightarrow$ interference contributions of dim-6 operators
- last term $\rightarrow$ squared contributions of dim-6 operators



tWZ in the SMEFT: differential predictions

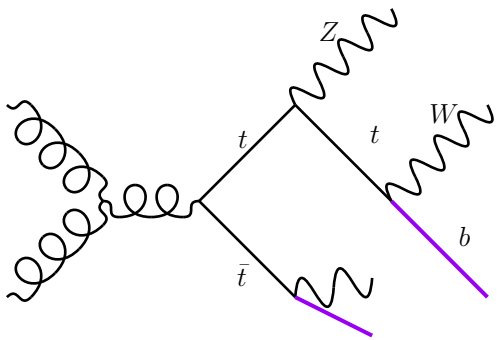




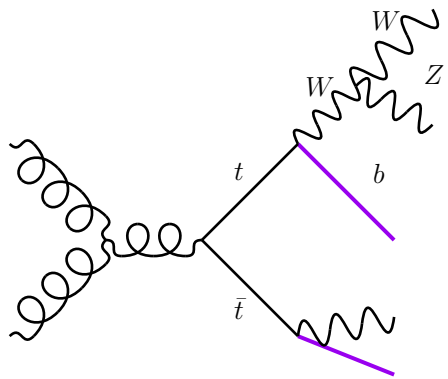
SMEFT impacts from LO to NLO accuracy are stable suggesting the DR treatment is correctly identifying the phase space of the *tWZ* process

How to do NLOPS?

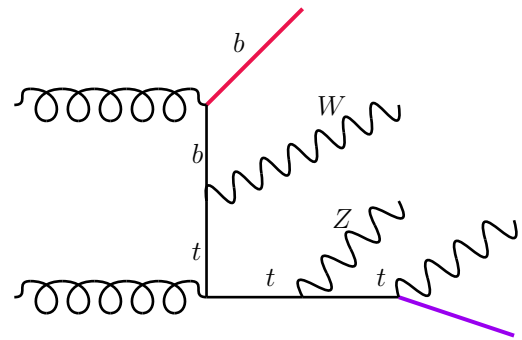
Decay the top and select events with **EXACTLY ONE** hard *b*-jet, the `1-bjet` scenario



ttZ

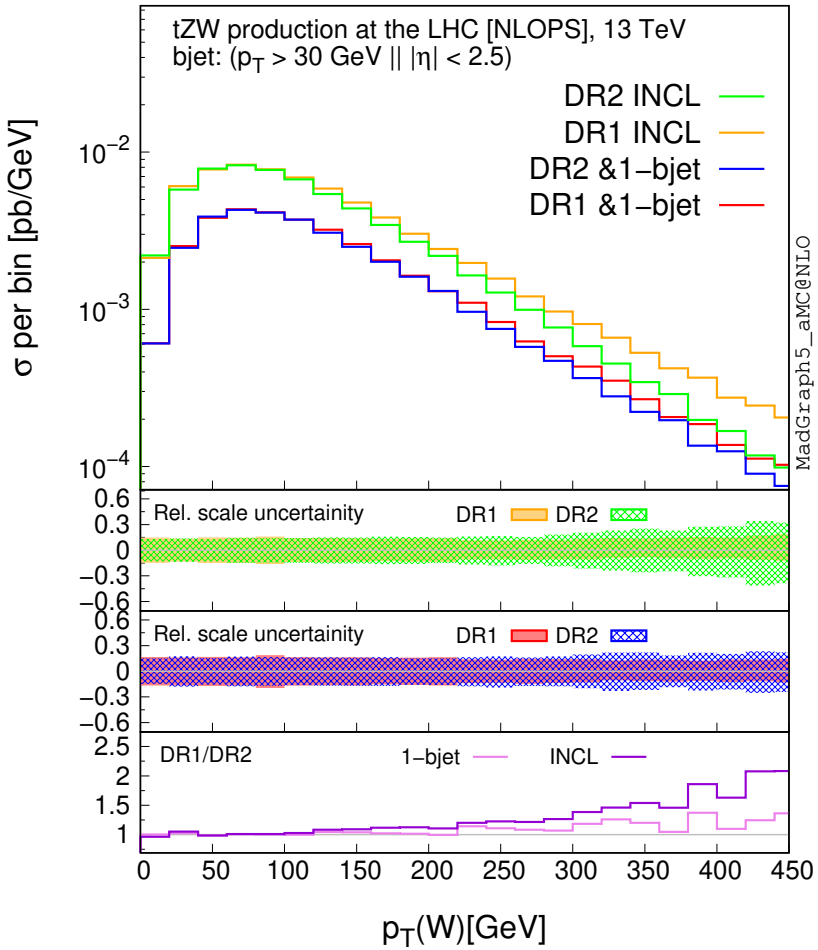


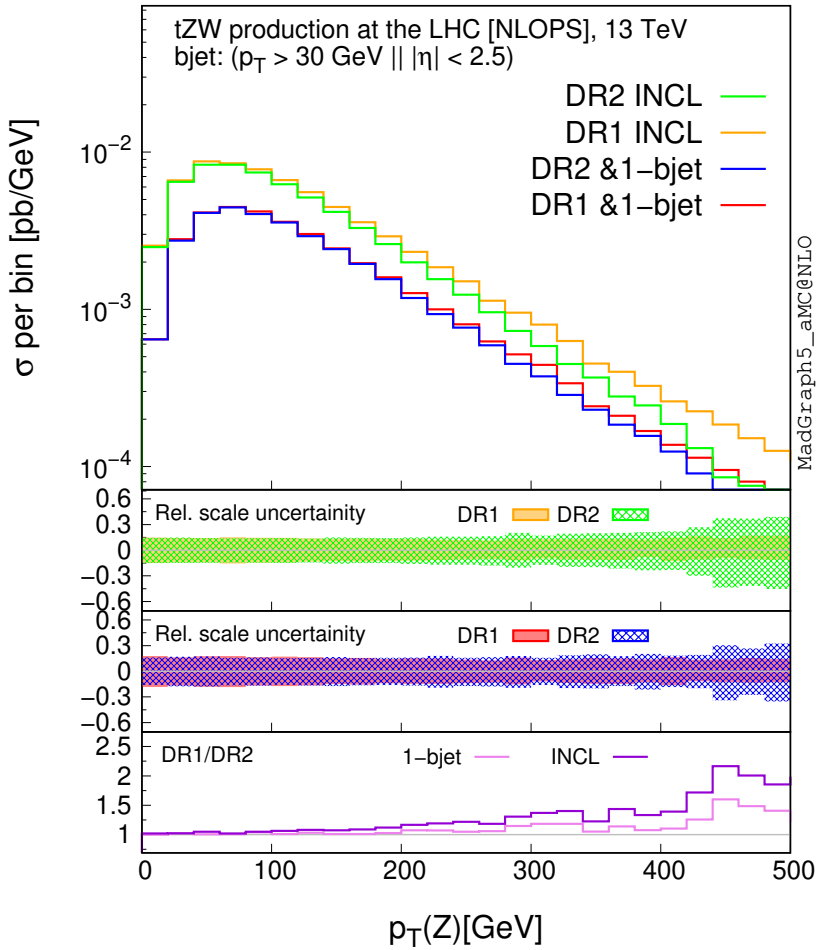
tt



tWZ

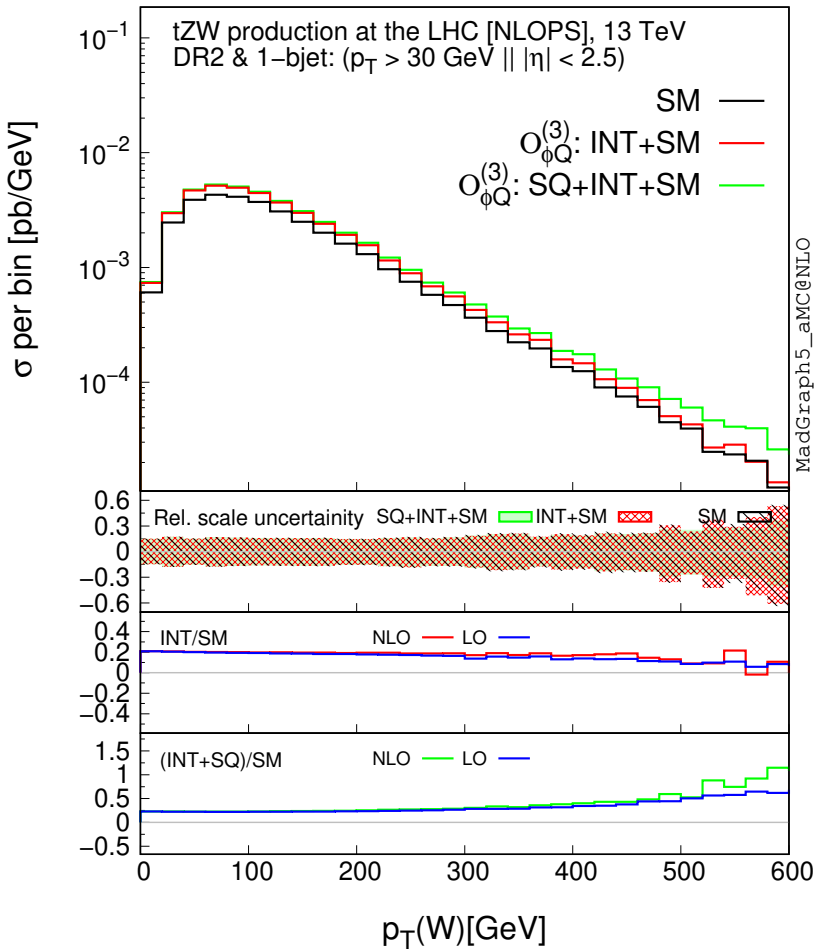
tWZ in the SM: NLOPS

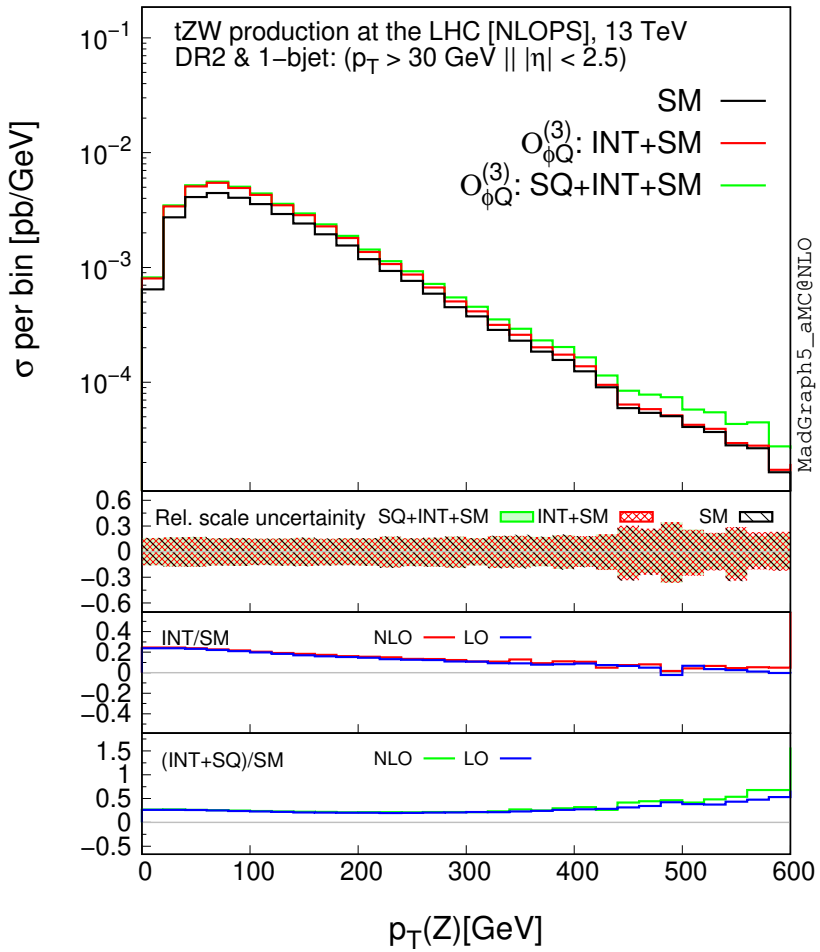




The `1-bjet` scenario recovers the DR1 and DR2 agreement

tWZ in the SMEFT: NLOPS



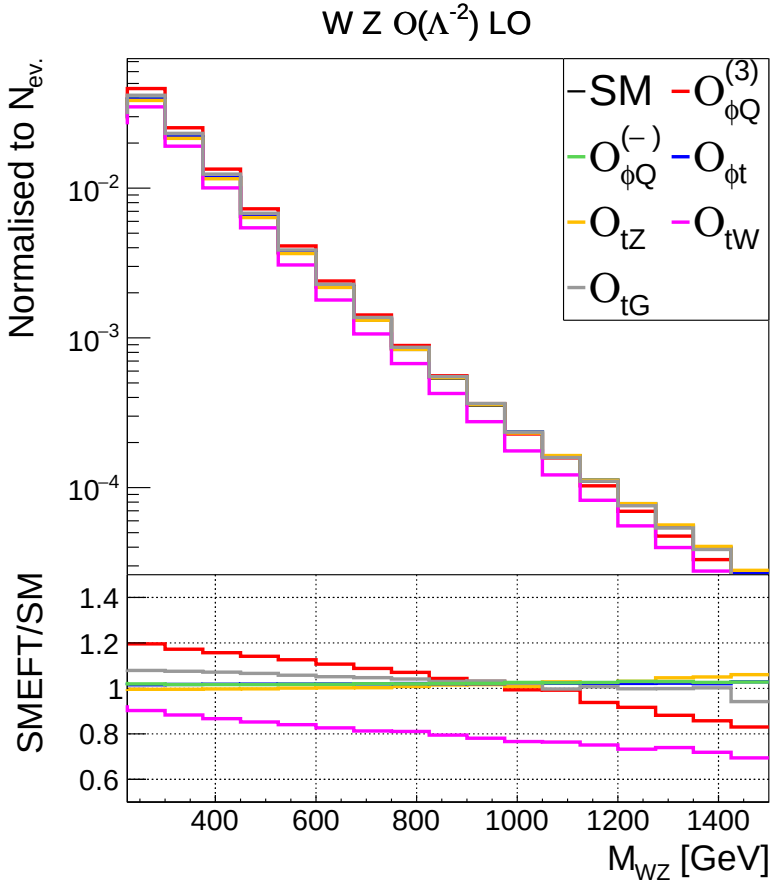


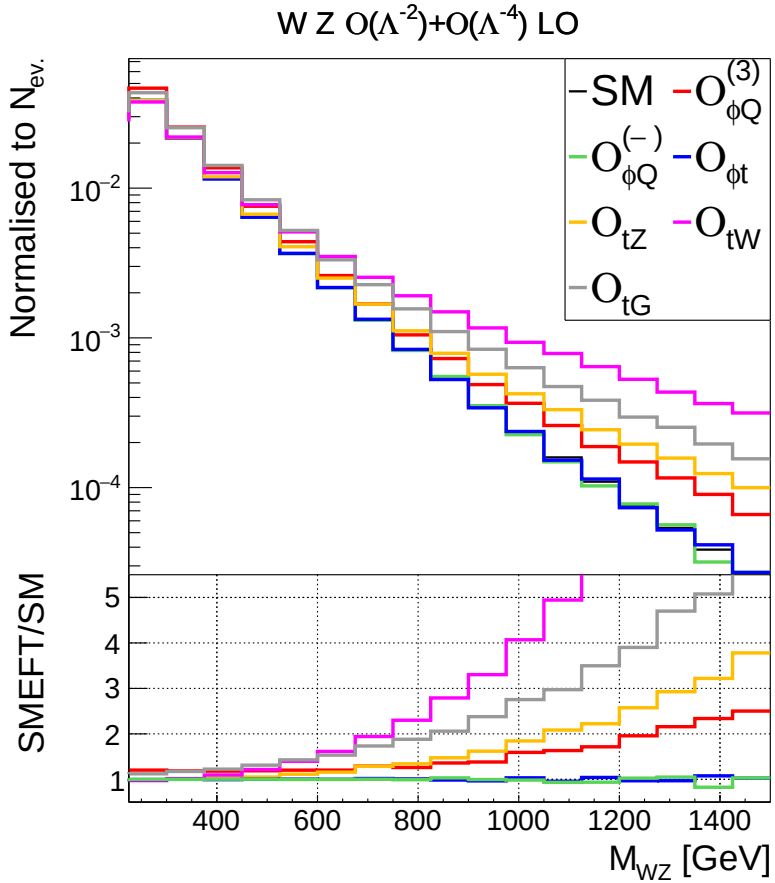
The behaviour of SMEFT operators resemble the corresponding one at fixed order

Summary

- A study of ***tWZ*** production at NLO in QCD is presented in the context of the SM and SMEFT
- The resonant overlap is handled using the DR schemes at the amplitude-level
- Vetoing hard ***b***-quarks ensures the ***tWZ*** process dominates the phase space after the diagram removal
- The differential results presented suggest the DR treatment correctly identifies the phase space of the ***tWZ*** process
- The work presented lays the foundation for precision LHC interpretations of ***tWZ*** data in the SMEFT framework

Backup
tWZ in the SMEFT: more differential predictions





M_{WZ} is a proxy for the $bW \to tZ$ sub-amplitude scattering energy