

Top-quark mass related studies

LHC top Working Group Open Meeting, Dec 1- 3

Javier Aparisi and Matteo Defranchis,

on behalf of the ATLAS and CMS collaborations

Material covered in this talk

- *“Towards a precise interpretation of the top quark mass parameter in the ATLAS MC samples”*

[ATL-PHS-PUB-2021-034](#)

- “Running of the top quark mass from proton-proton collisions at $\sqrt{s} = 13 \text{ TeV}$ ”

[CMS-TOP-19-007](#)

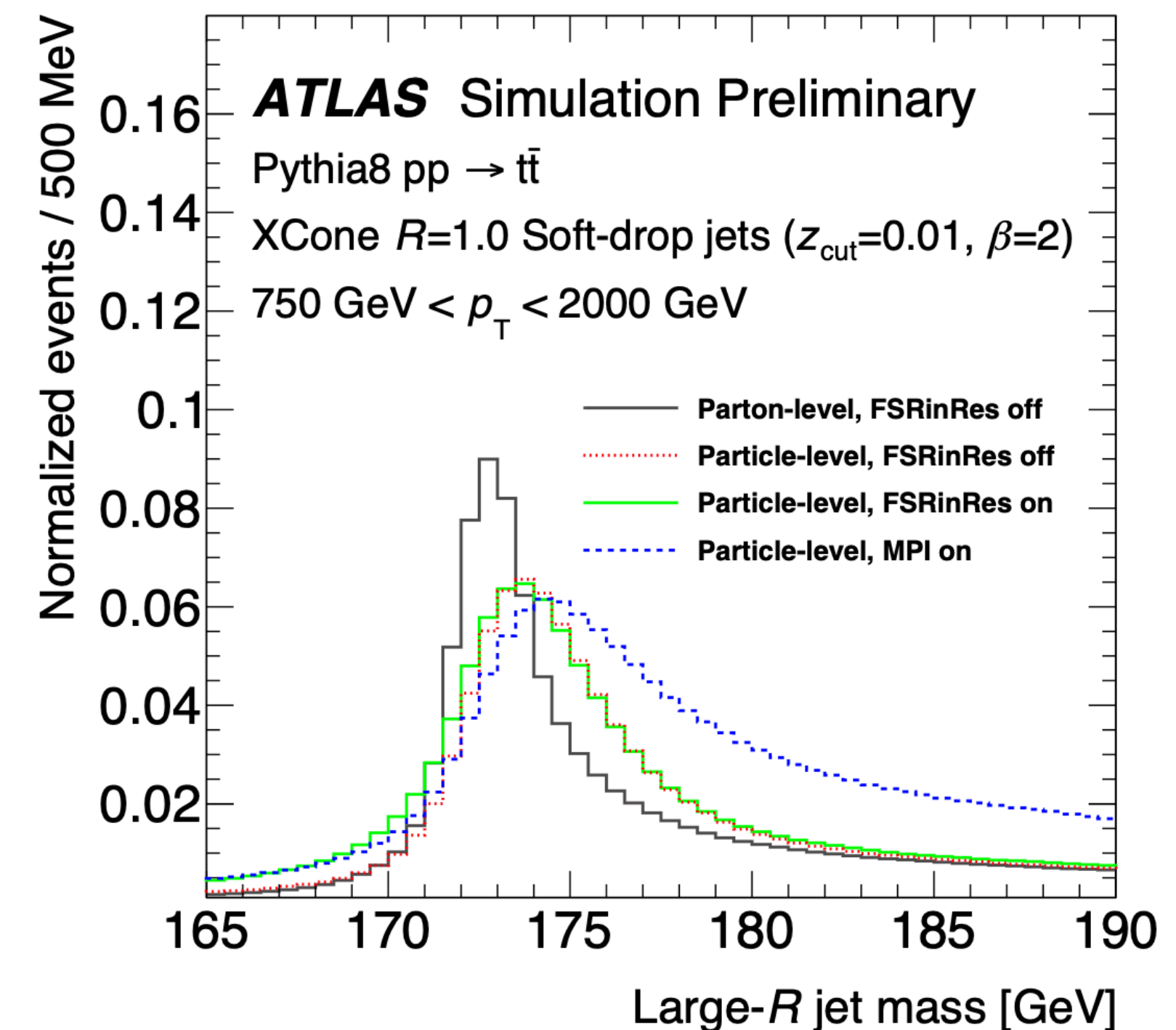
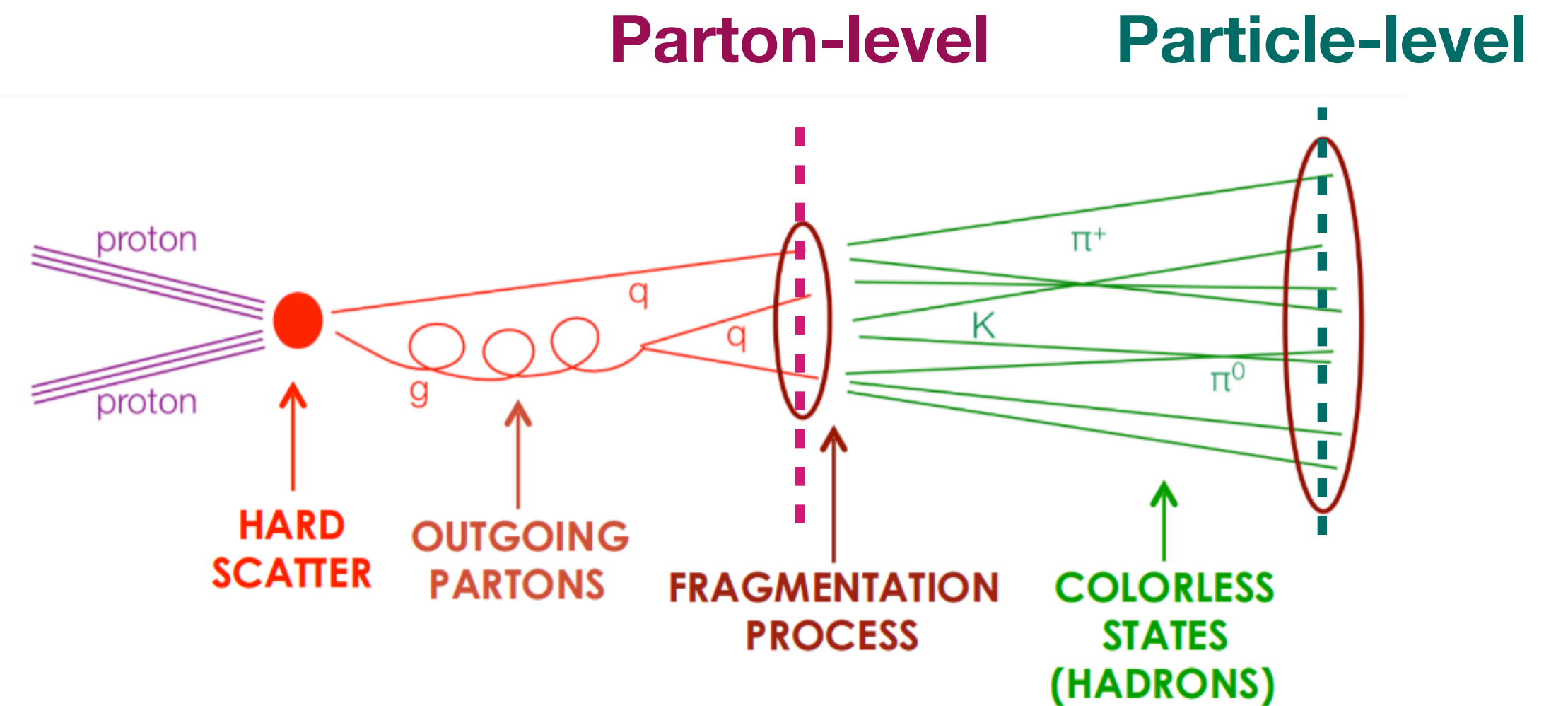
Towards a precise interpretation of the top quark mass parameter in the ATLAS MC samples

1. Motivation

○ Top quark mass experimentally determined in **direct** and **indirect** measurements:

- **DM:** Monte Carlo (MC) fits to data at detector level
 $\rightarrow m_{top}^{MC}$ with O(600) MeV precision
- **IM:** fixed-order theory fits to data at parton level
 $\rightarrow m_{top}^{pole, \overline{MS}}$ with O(1) GeV precision

○ Can the m_{top}^{MC} be identified with a well-defined mass scheme m_{top} with a precision better than 500 MeV?



2. Goal

- The main goal of the analysis is

the interpretation of the top mass in official ATLAS MC generators in light of a renormalized mass in the MSR scheme*:

$$m_{top}^{MC} = m_{top}^{MSR}(R) + \Delta m^{MSR}$$

The top MSR mass is a short-distance mass that depends on an R parameter, above which top self-interactions are absorbed. For $R = 1 \text{ GeV}$, it is conceptually close to the top MC mass ($\Lambda^{PS} \sim 1 \text{ GeV}$):

$$m_{top}^{MSR}(R = 0 \text{ GeV}) \rightarrow m_{top}^{pole}, \quad m_{top}^{MSR}(R = m_t) \rightarrow m_{top}^{\overline{MS}}$$

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$$m_{top}^{MC} = m_{top}^{MSR}(R) + \Delta m^{MSR}$$

- A *calibration* is performed by comparing ATLAS MC predictions to a NLL calculation for two MC models: Pythia8 (P8) and Herwig7 (H7).

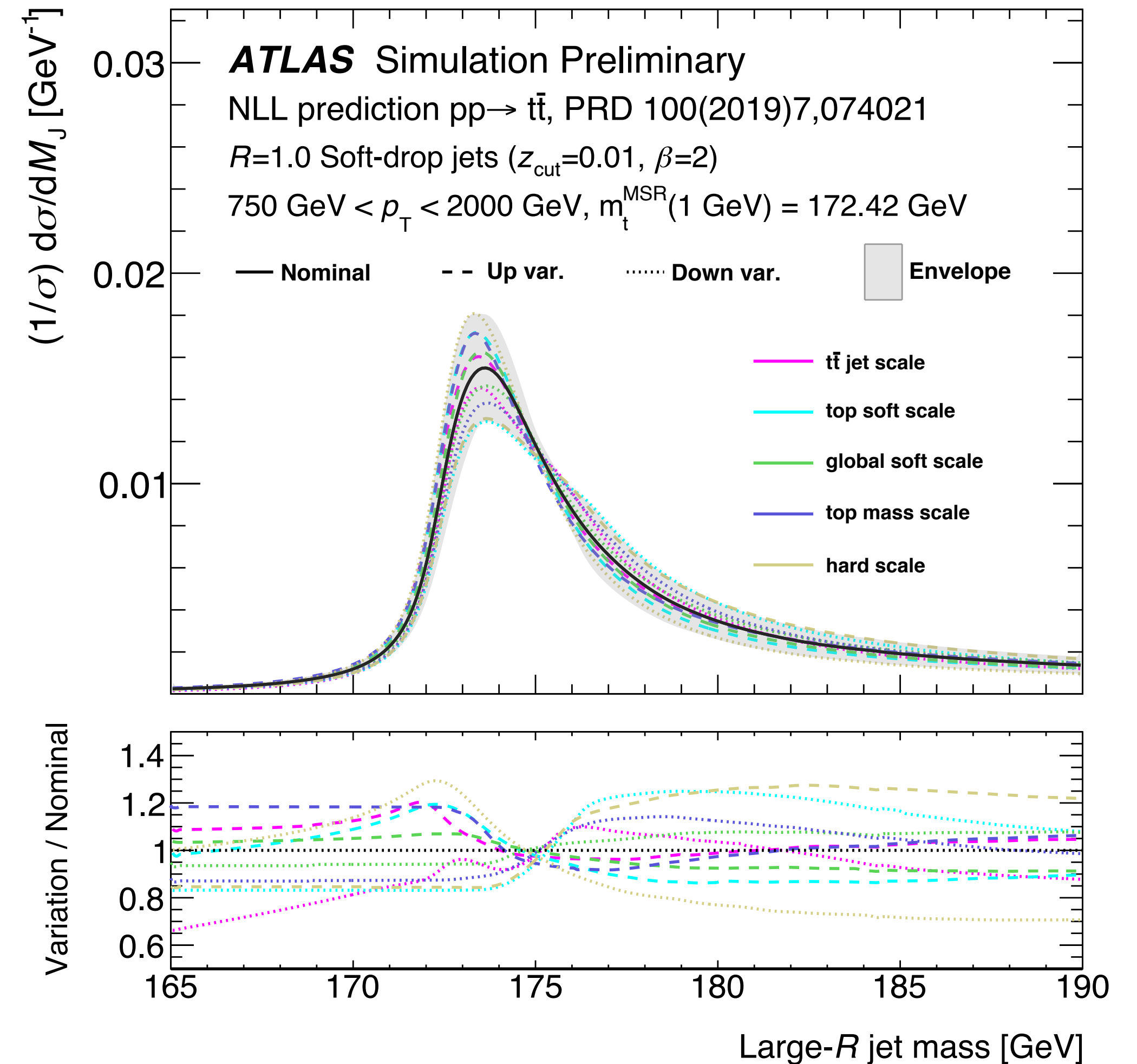
The m_{top}^{MC} is set to 172.5 GeV.

Internal variations of the nominal MC (Pythia8) are studied as well in order to check their impact in the MSR-MC mass relation determined in the nominal setup

3. Theoretical calculation

- SCET-based theory at NLL accuracy
- Differential jet mass cross section at particle-level
- Strong **sensitivity to m_{top} in the jet mass peak** region
- Main ingredients:
 1. Inclusive treatment of hadronic top quark decays
 2. Light soft-drop grooming to remove soft-wide radiation
 3. Three free parameters: m_{top} and Ω_{1q}^{had} and x_2 for **non-perturbative corrections**
- **Does not account for UE effects**

$$\frac{\partial \sigma^{tot}}{\partial M_J} = \frac{\partial \sigma^{pQCD}}{\partial M_J} \otimes \frac{\partial \sigma^{npQCD}}{\partial M_J}$$



4. Jet reconstruction & selection

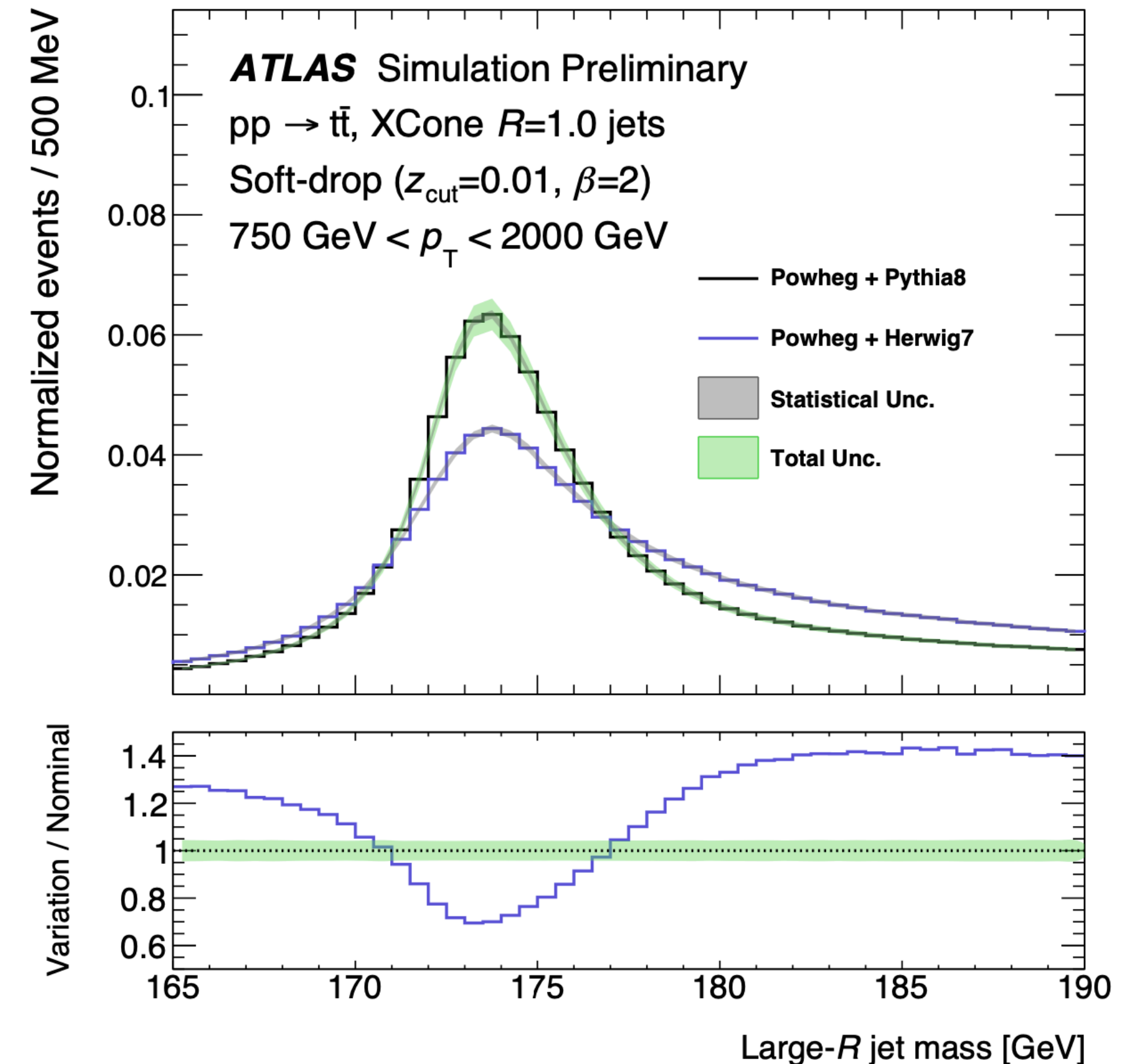
- The observable is a **particle-level groomed jet mass**, built from stable particles in $pp \rightarrow t\bar{t}$ events*
 - I) **XCone** jet algorithm with $R=1$
 - II) **Light soft-drop grooming** applied ($z_{cut} = 0.01, \beta = 2$)
- **Boosted regime**: needed to capture all top decay products in a single jet.

Three orthogonal p_T bins considered:

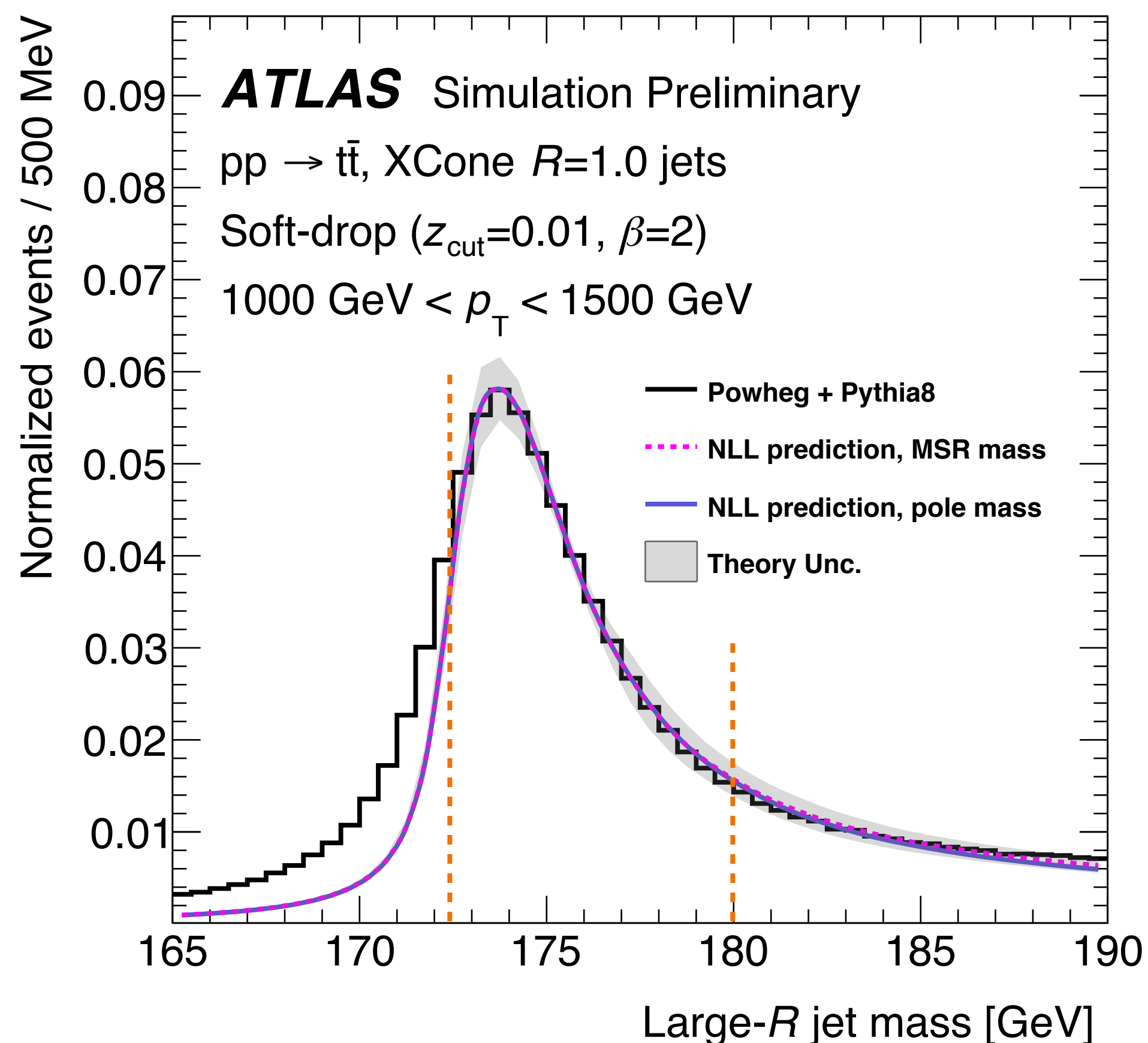
$$p_T^{jet} \in [750, 1000], [1000, 1500], [1500, 2000] \text{ GeV}$$

Three free parameters of theory scale differently with the jet p_T
- Large-R jet to hadronically decaying top parton matching applied: $\Delta R(jet, top) < 1$

*Events generated with MPI setting switched off in order to not include UE effects



5. Results (I)



Fit range: 172.5 - 180.0 GeV

- χ^2 fit and minimization to find the theory prediction that best describes MC:

Ph+P8:

$$m_{top}^{MSR}(1 \text{ GeV}) = 172.42 \pm 0.10 \text{ (stat.) GeV}$$

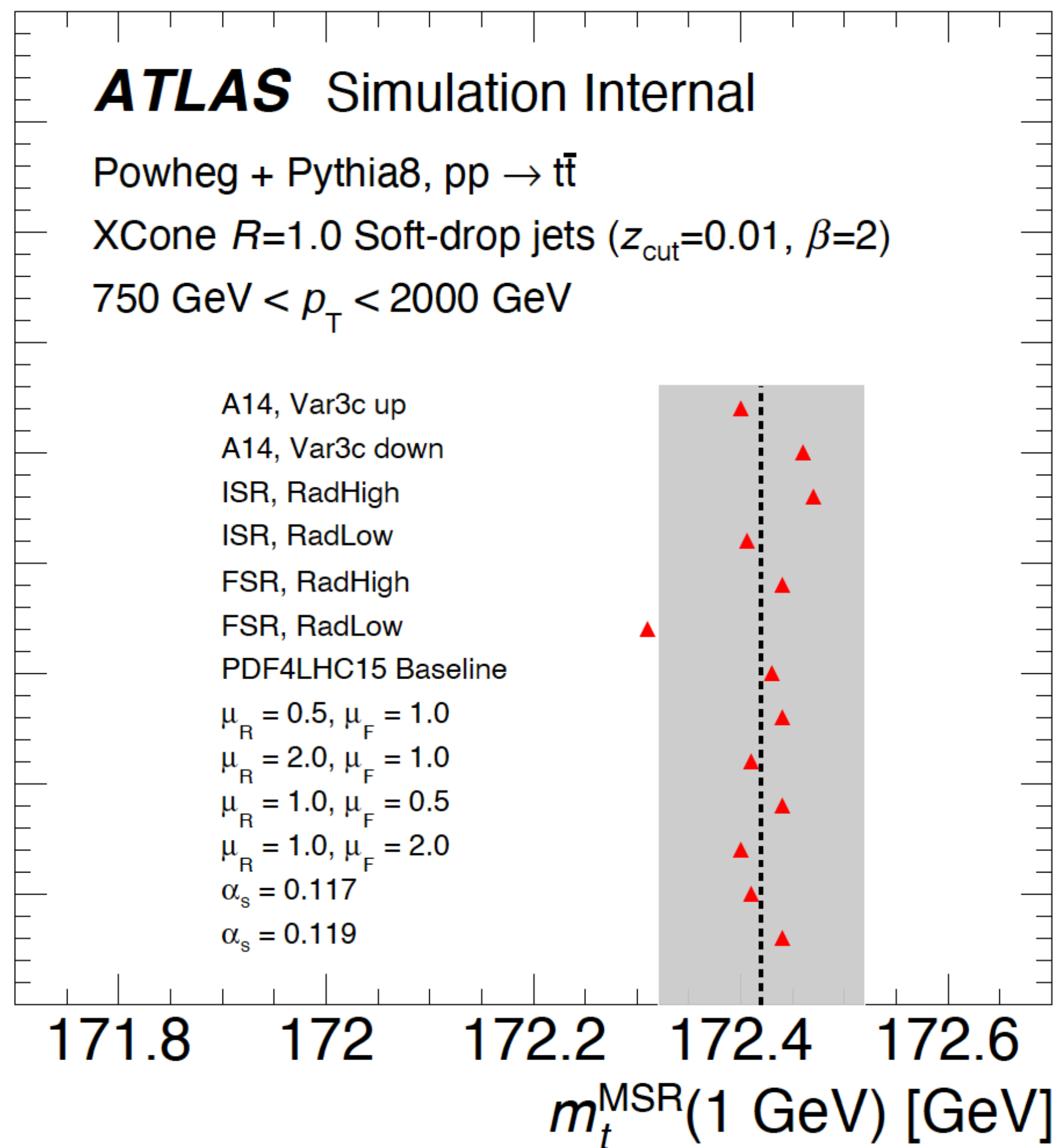


$$m_{top}^{MC} = m_{top}^{MSR}(1 \text{ GeV}) + 80^{+350}_{-410} \text{ MeV}$$

with **theory** component **dominating** total uncertainty:

Source	Size [MeV]	Comment
Theory (higher-order corrections)	+230/-310	Envelope of NLL scale variations
Fit methodology	± 190	Choice of fit range, p_T bins
Underlying Event model	± 155	A14 eigentune variations, CR models
Total Systematic	+340/-340	
Statistical Uncertainty	± 100	
Total Uncertainty	+350/-410	

5. Results (II)



- χ^2 fit and minimization to find the theory prediction that best describes MC:

Ph+P8:

$$m_{top}^{\text{MSR}}(1 \text{ GeV}) = 172.42 \pm 0.10 \text{ (stat.) GeV}$$

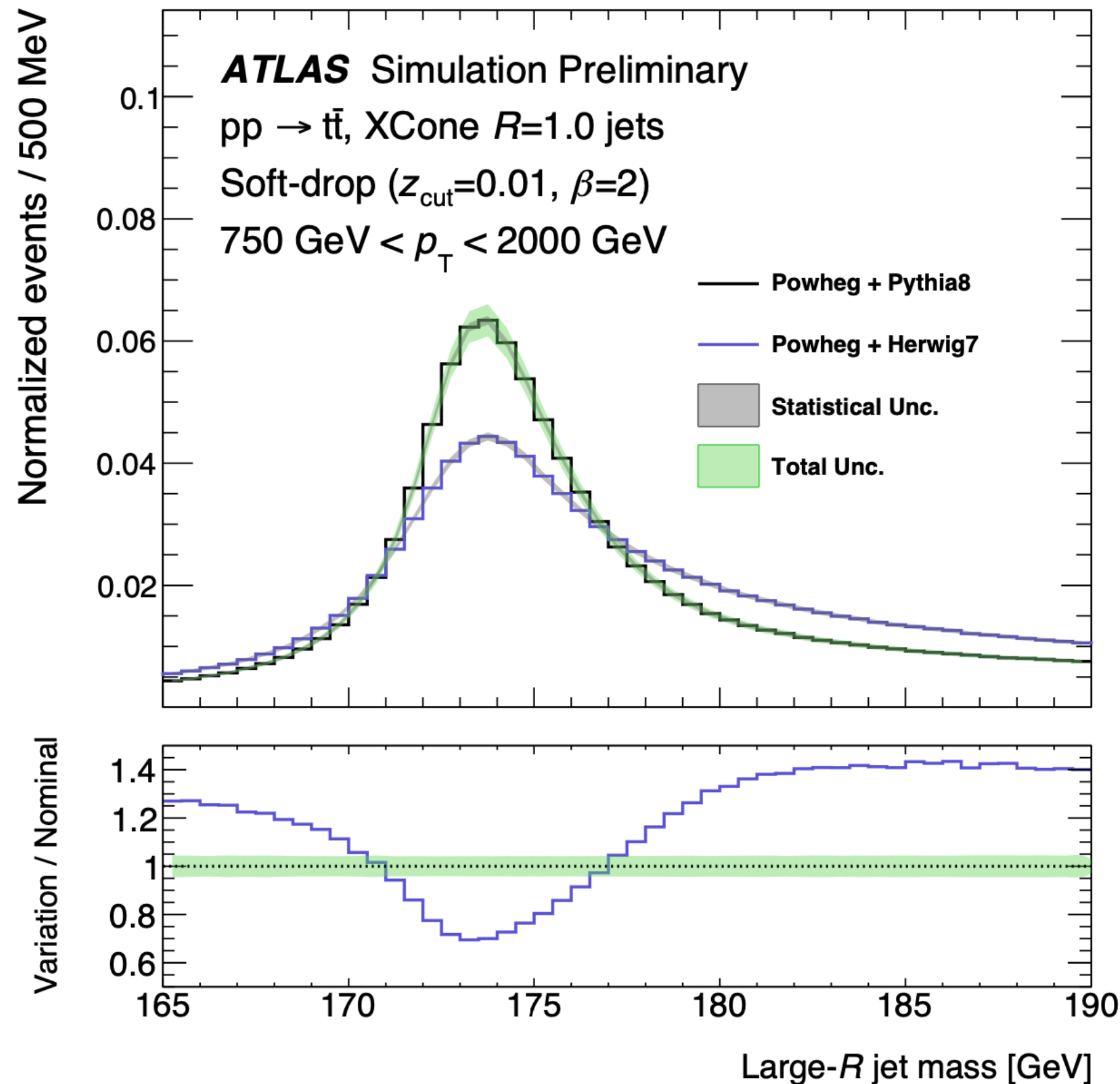


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5. Results (III)



- χ^2 fit and minimization to find the theory prediction that best describes MC:

Ph+P8:

$$m_{top}^{MSR}(1 \text{ GeV}) = 172.42 \pm 0.10 \text{ (stat.) GeV}$$

Ph+H7:

$$m_{top}^{MSR}(1 \text{ GeV}) = 172.27 \pm 0.09 \text{ (stat.) GeV}$$

- The calibration can be performed on any MC sample. Surprisingly similar results for **Powheg+Pythia8** and **Powheg+Herwig7**, even if these two Monte Carlo setups predict very different jet mass distributions.
- These **MSR-MC mass relations are found to be stable within 200 MeV**, when repeating this exercise considering related observables accessible with the underlying theory

Related observables: Anti-kt and softer/harder SD configuration

Running of the top quark mass from proton-proton collisions at $\sqrt{s} = 13 \text{ TeV}$

1. Introduction

- In the \overline{MS} scheme, the **evolution** of any **quark mass** with the energy scale μ is **governed by the RGE**:

$$\mu^2 \frac{dm(\mu)}{d\mu^2} = -\gamma(\alpha_s(\mu)) m(\mu)$$

- The running of the **b-quark and the charm-quark mass** have been **already studied** using data from CERN LEP, SLAC SLC and DESY HERA.
- The **top-quark mass running is first studied** in $t\bar{t}$ events with 39.5 fb^{-1} of data collected by the **CMS detector** at $\sqrt{s} = 13 \text{ TeV}$.

2. Strategy

- Measure the **differential $t\bar{t}$ cross-section** as a function of the $t\bar{t}$ **invariant mass** $m_{t\bar{t}}$ in the following bins:

Bin	$m_{t\bar{t}}$ [GeV]	Fraction [%]	μ_k [GeV]
1	<420	30	384
2	420–550	39	476
3	550–810	24	644
4	>810	7	1024

so **the scale of the process is chosen to be given by the center-of-gravity of bin k , μ_k .**

- Compare with fixed-order **NLO prediction in QCD** at parton level to extract the **top \overline{MS} mass** simultaneously at each $m_{t\bar{t}}$ bin.
- Test the **running hypothesis** considering that at each k bin, $m_t(\mu) = m_t(\mu_k)$

3. Event selection and categorization

- $t\bar{t}$ events selected in the **dileptonic channel**:

- At least one e and one μ with opposite charge where:

$$p_T^{lead} > 25 \text{ GeV}, p_T^{slead} > 20 \text{ GeV} \text{ and } |\eta| < 2.4$$

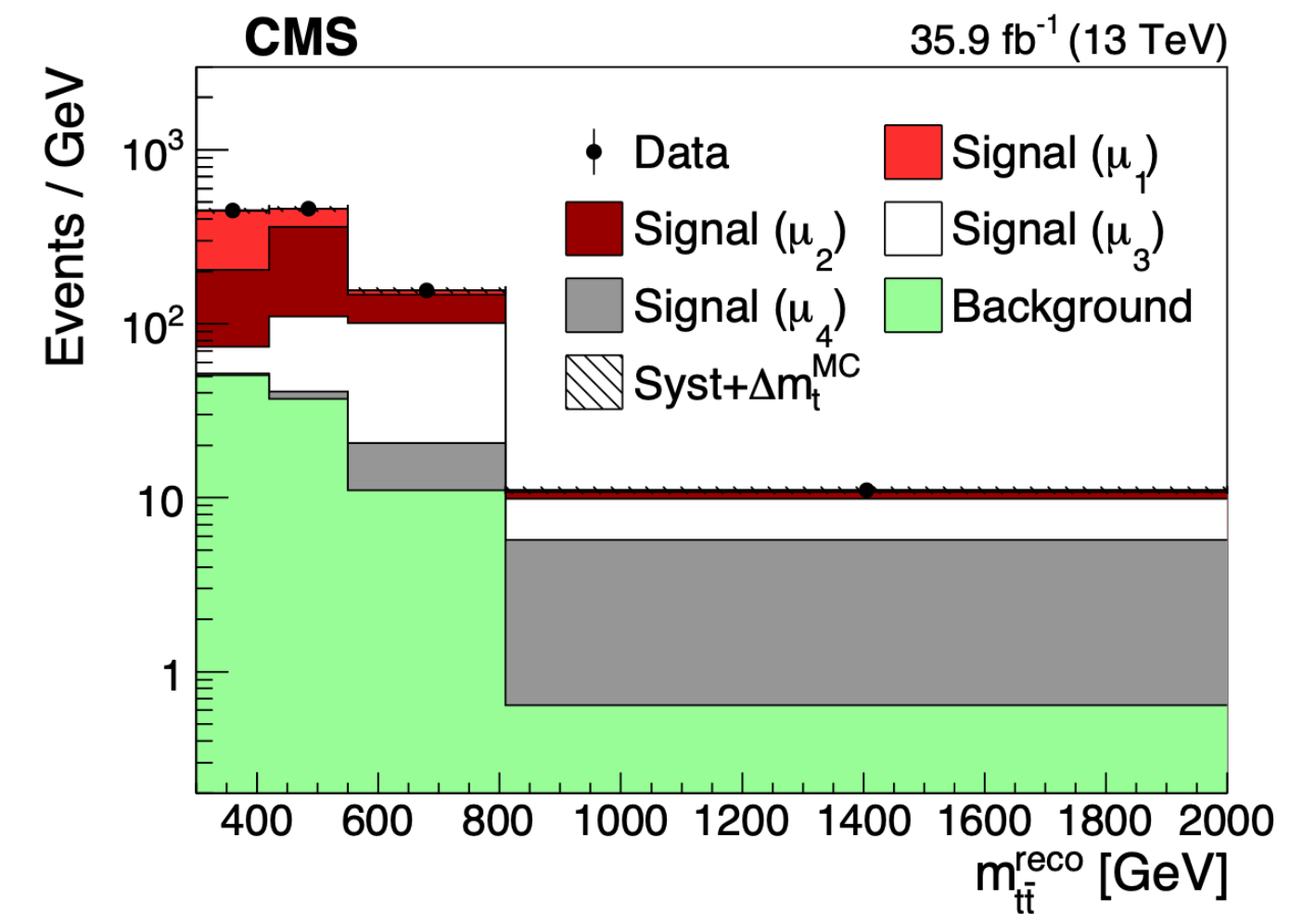
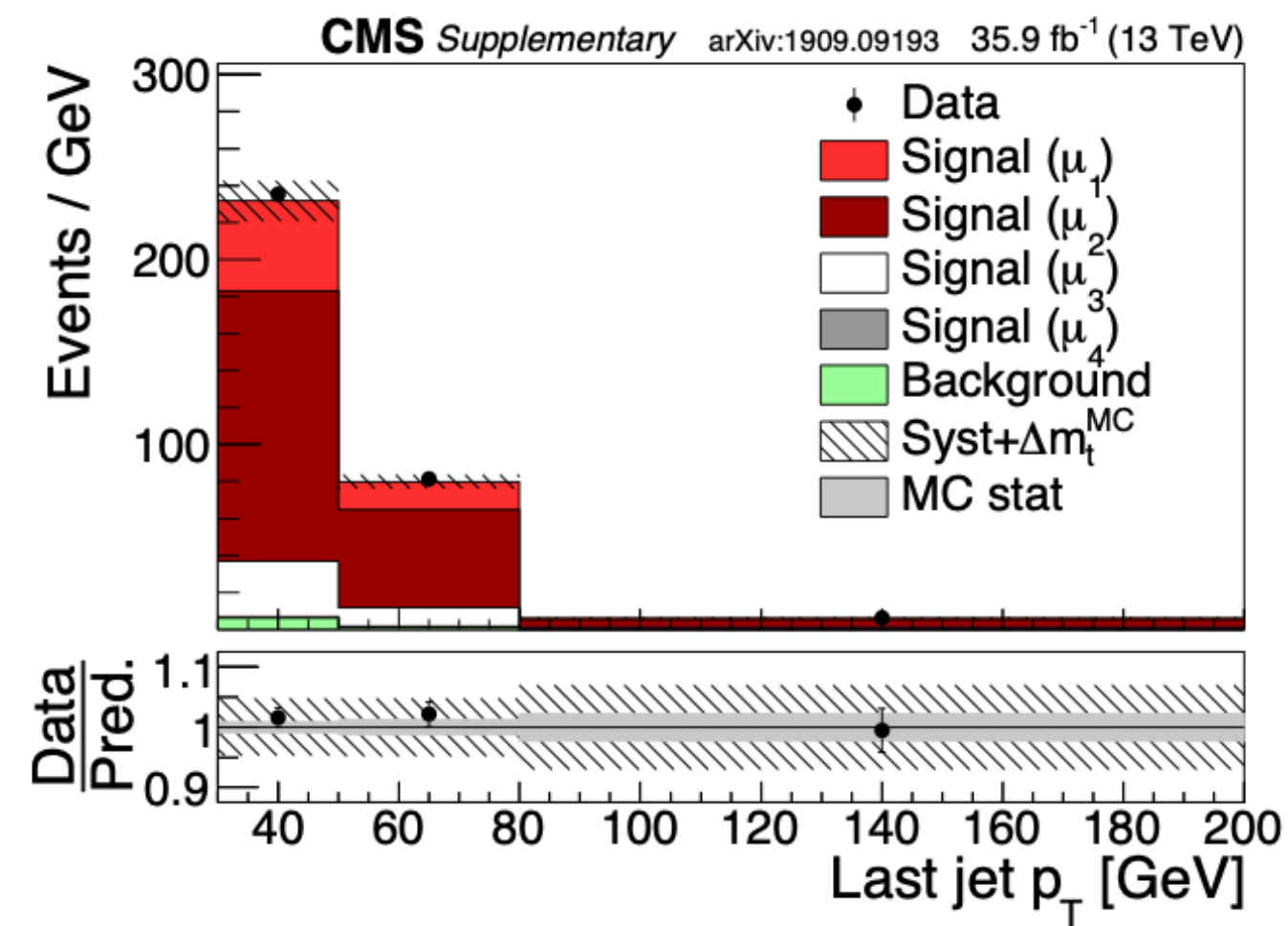
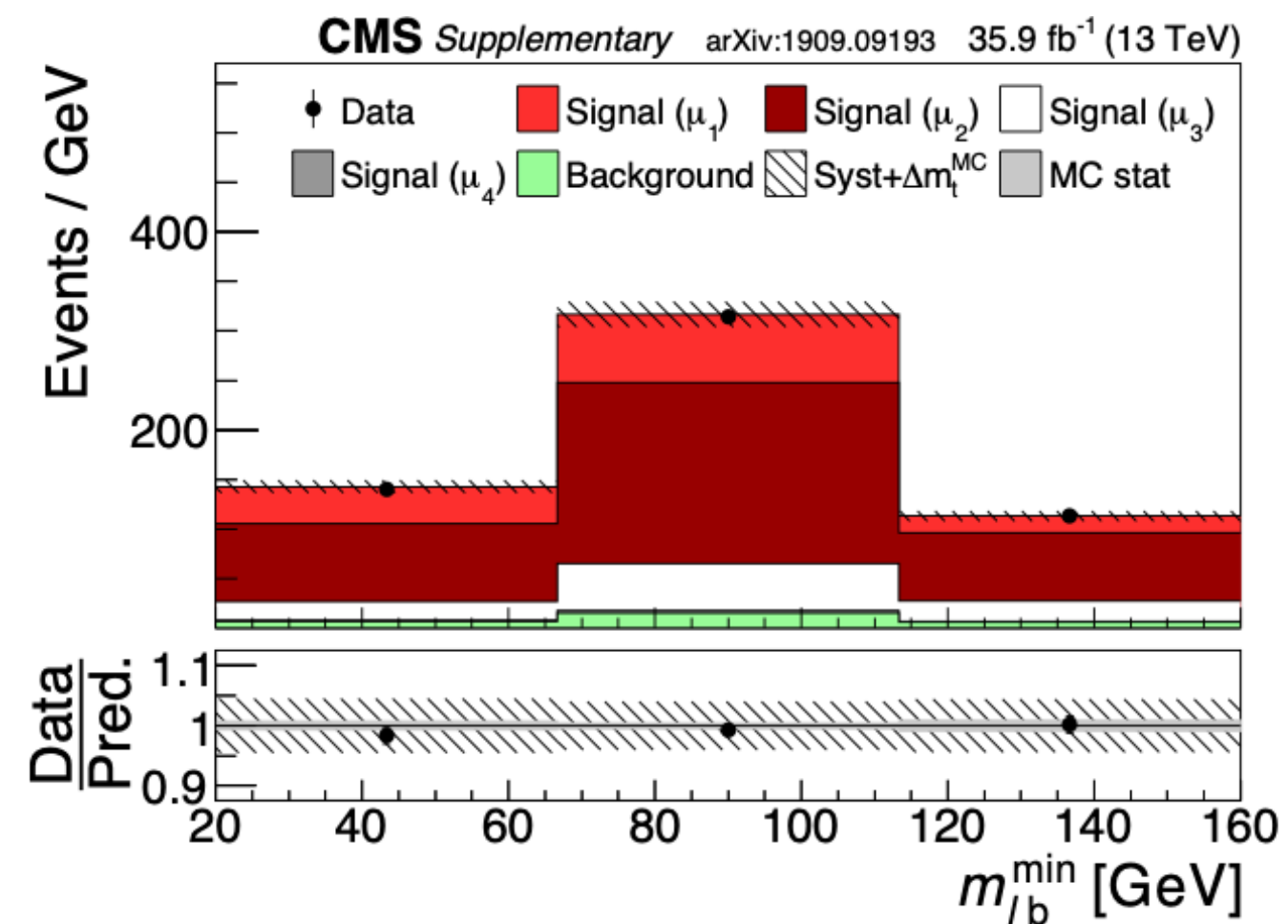
- Jets with $p_T^{jet} > 30 \text{ GeV}$ and $|\eta| < 2.4$.

- **Kinematic reconstruction of $m_{t\bar{t}}$** from combination of jets, leptons and MET ($m_W = 80.4 \text{ GeV}$) in events with at least two jets -> **dependence on m_t^{MC} mitigated in the fit**
- **Events** are **categorized according to b-tagged jets multiplicity** (1, 2 or (0 and >2)) and further divided into **subcategories of $m_{t\bar{t}}^{reco}$** .
- The **observables** used are the **p_T of the softest jet**, the invariant mass of the b-jet and lepton system m_{lb} and the **event yield**.

Bin	$m_{t\bar{t}}$ [GeV]	Fraction [%]	μ_k [GeV]
1	<420	30	384
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4. Fit procedure

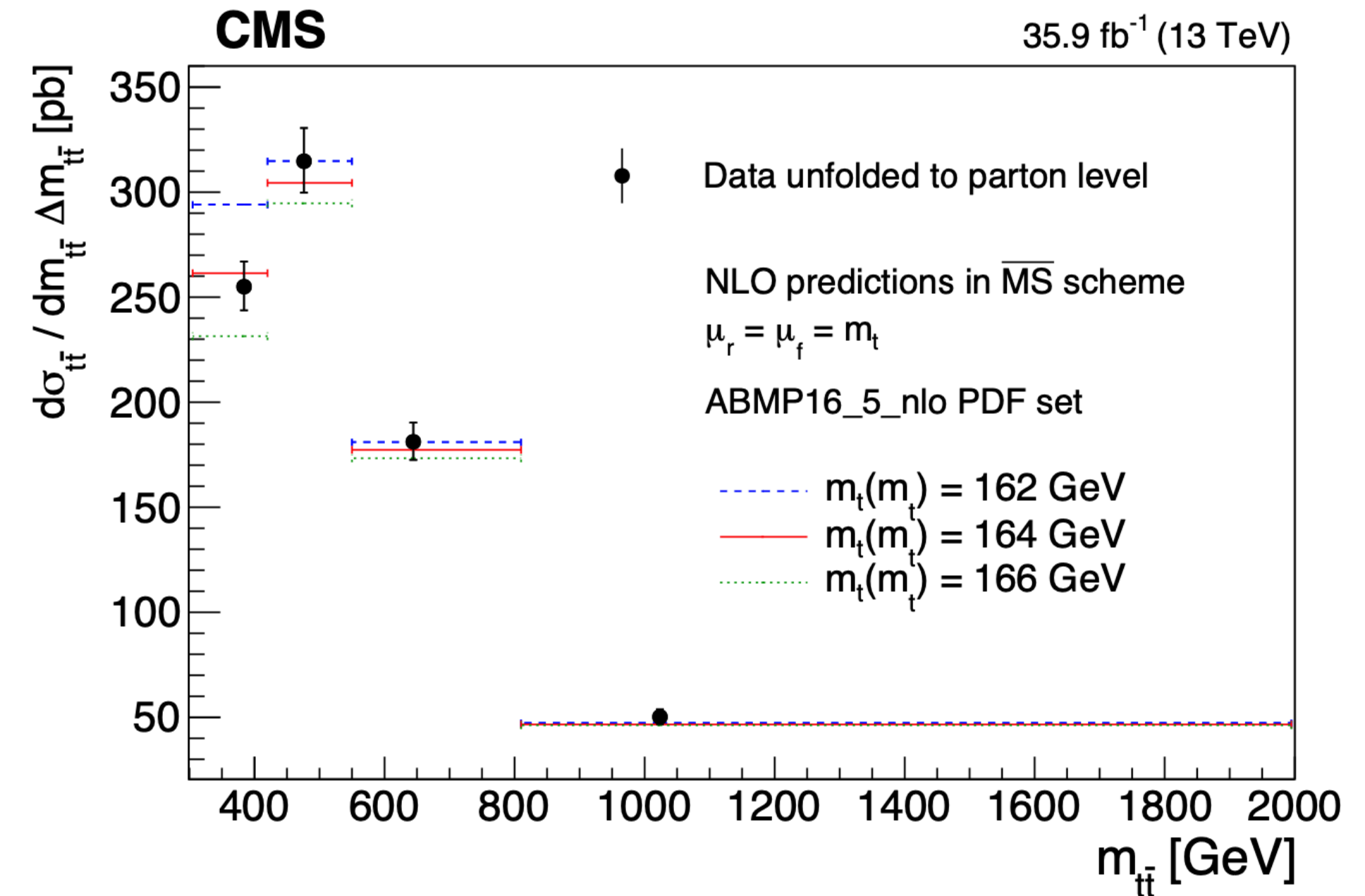
- $d\sigma_{t\bar{t}}/dm_{t\bar{t}}$ determined directly **at parton level with a maximum-likelihood fit** to distributions of final-states observables:
 - **Unfolding implicit** in the likelihood fit: $t\bar{t}$ simulated sample splitted into $m_{t\bar{t}}$ bins, treated as independent signals and interpreted as $t\bar{t}$ cross-section at the physical scale μ_k .
 - **Syst. unc.** treated as **nuisance parameters** → constrained with soft jet p_T and **b-jet categorization**
 - m_t^{MC} is a floating parameter → **fitted from m_{lb}** distribution



Illustrative figures of data and post-fit MCs for $m_{t\bar{t}}^{\text{reco}} \in \{420, 550\}$ GeV

5. Results (I): extraction of the top \overline{MS} quark mass

- Value of $m_t(m_t)$ is determined independently in each bin from χ^2 fit to NLO QCD calculation.
- It is evolved from $m_t(m_t)$ to $m_t(\mu_k)$ with RGE at one-loop precision ($n_f = 5$, $\alpha_s(m_Z) = 0.1191$)

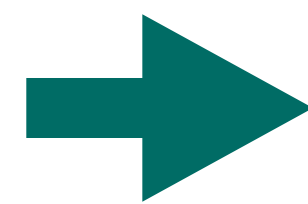


$$\sigma_{t\bar{t}}^{(\mu_1)} = 255 \pm 11 \text{ (syst)} \pm 2 \text{ (stat)} \text{ pb,}$$

$$\sigma_{t\bar{t}}^{(\mu_2)} = 315 \pm 15 \text{ (syst)} \pm 2 \text{ (stat)} \text{ pb,}$$

$$\sigma_{t\bar{t}}^{(\mu_3)} = 181 \pm 9 \text{ (syst)} \pm 1 \text{ (stat)} \text{ pb,}$$

$$\sigma_{t\bar{t}}^{(\mu_4)} = 50 \pm 3 \text{ (syst)} \pm 1 \text{ (stat)} \text{ pb.}$$



$$m_t(\mu_1) = 155.4 \pm 0.8 \text{ (fit)} \pm 0.2 \text{ (PDF}+\alpha_s) \pm 0.1 \text{ (extr)} \begin{matrix} +0.9 \\ -0.6 \end{matrix} \text{ (scale),}$$

$$m_t(\mu_2) = 150.9 \pm 3.0 \text{ (fit)} \begin{matrix} +1.1 \\ -0.7 \end{matrix} \text{ (PDF}+\alpha_s) \begin{matrix} +0.4 \\ -0.5 \end{matrix} \text{ (extr)} \begin{matrix} +3.9 \\ -4.3 \end{matrix} \text{ (scale),}$$

$$m_t(\mu_3) = 148.2 \pm 4.6 \text{ (fit)} \begin{matrix} +2.0 \\ -1.4 \end{matrix} \text{ (PDF}+\alpha_s) \begin{matrix} +0.9 \\ -1.0 \end{matrix} \text{ (extr)} \begin{matrix} +7.3 \\ -9.5 \end{matrix} \text{ (scale),}$$

$$m_t(\mu_4) = 136.4 \pm 9.0 \text{ (fit)} \begin{matrix} +3.8 \\ -3.0 \end{matrix} \text{ (PDF}+\alpha_s) \begin{matrix} +2.8 \\ -2.3 \end{matrix} \text{ (extr)} \begin{matrix} +9.6 \\ -16.1 \end{matrix} \text{ (scale).}$$

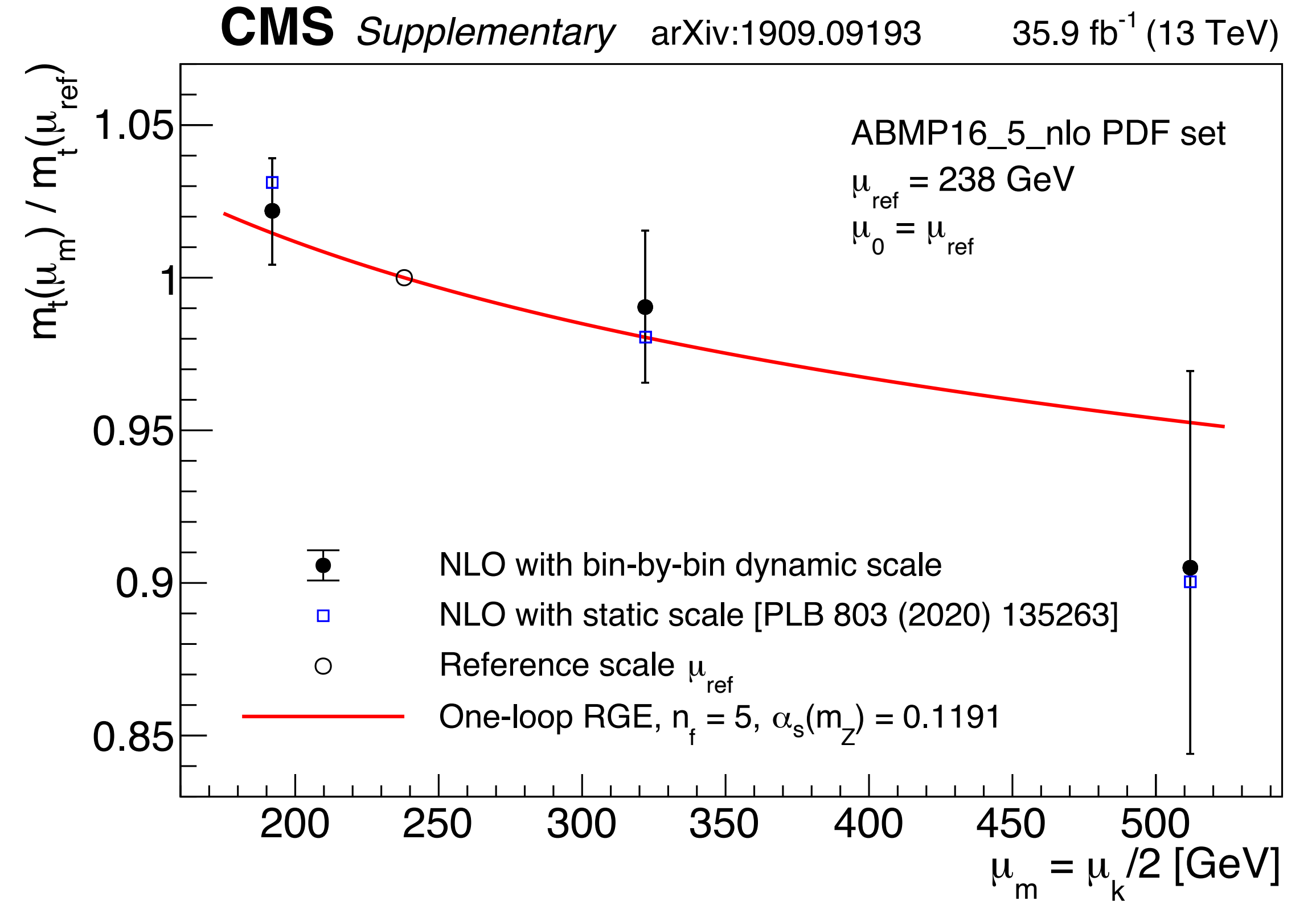
5. Results (II): testing the top quark mass running

- Finally, the **ratios** $r_{ij} = m_t(\mu_i)/m_t(\mu_j)$ are **considered**
- Taking $\mu_j = \mu_{ref} = \mu_2/2 = 238 \text{ GeV}$ (final result independent from μ_{ref} choice), the **running is parametrized with x** :

$$f(x, \mu) = x \underbrace{\left[r_{ij} - 1 \right]}_{\text{RGE evolution}} + 1 \begin{cases} x = 0 : \text{No running} \\ x = 1 : \text{SM running} \end{cases}$$

For fixed scale: χ^2

$$\hat{x} = 2.05 \pm 0.61 \text{ (fit)} \quad {}^{+0.31}_{-0.55} \text{ (PDF + } \alpha_s) \quad {}^{+0.24}_{-0.49} \text{ (extr)}$$



Results obtained with dynamic (filled circles) and static (empty squares) scales in the theory calculation.

Summary

- *“Towards a precise interpretation of the top quark mass parameter in the ATLAS MC samples”*

$$m_t^{MC(P8)} = m_{top}^{MSR}(1 \text{ GeV}) + 80_{-410}^{+350} \text{ MeV}$$

_Theory error dominates. Significant improvement expected with NNLL calculation

_Compatible with m_t^{MC} calibration in $e^+e^- \rightarrow t\bar{t}$ events ([arXiv:1608.01318](https://arxiv.org/abs/1608.01318))

_Enables direct top mass measurement in boosted tops with cleaner interpretation

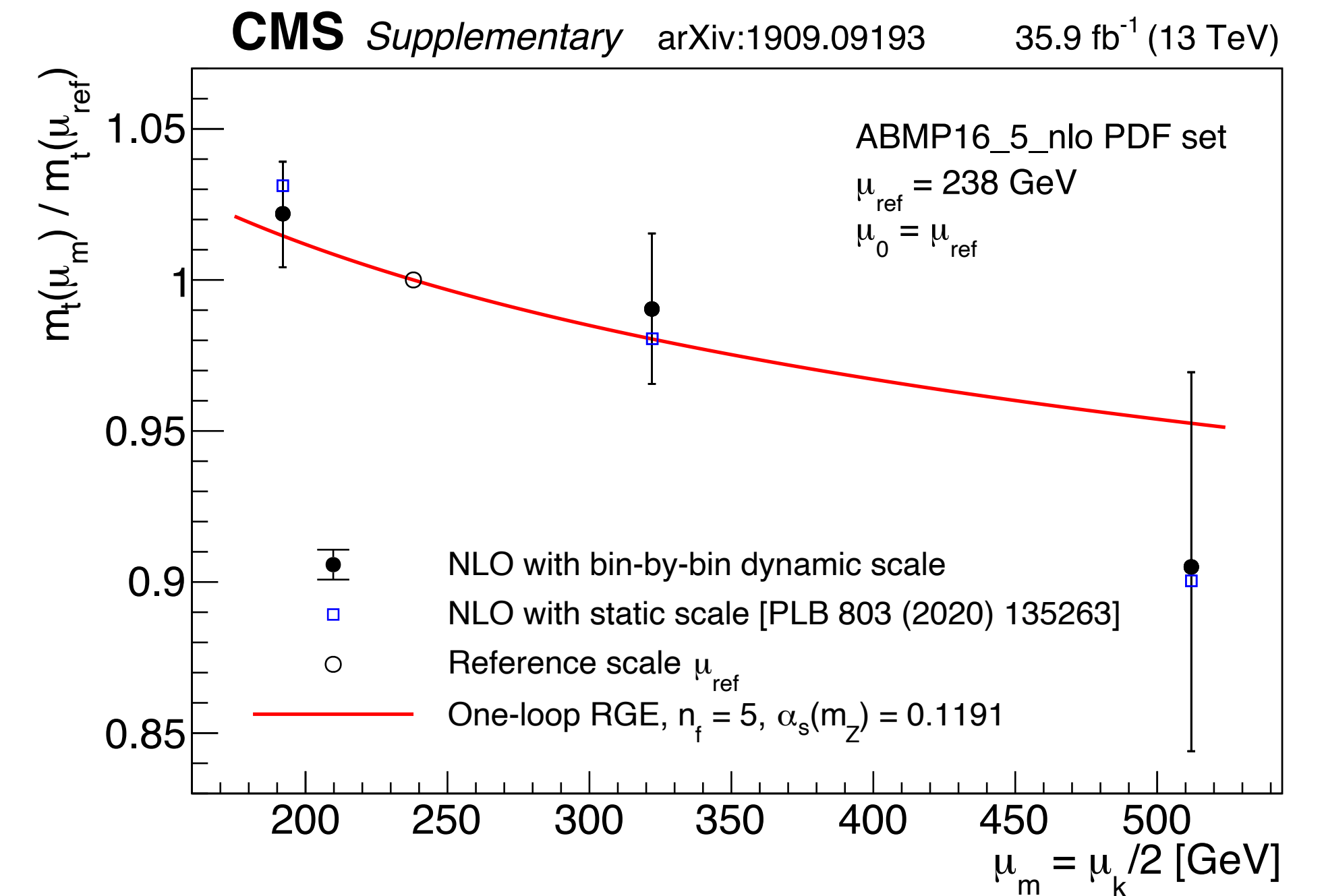
Summary

- “Running of the top quark mass from proton-proton collisions at $\sqrt{s} = 13 \text{ TeV}$ ”

Top \overline{MS} mass extracted from $d\sigma{t\bar{t}}/dm_{t\bar{t}}$ with NLO QCD calculation

_Top mass running probed up to $\sim 1 \text{ TeV}$. In agreement with RGE SM within 1.1σ

_ No-running scenario ruled out at above $\% 95 \text{ CL}$

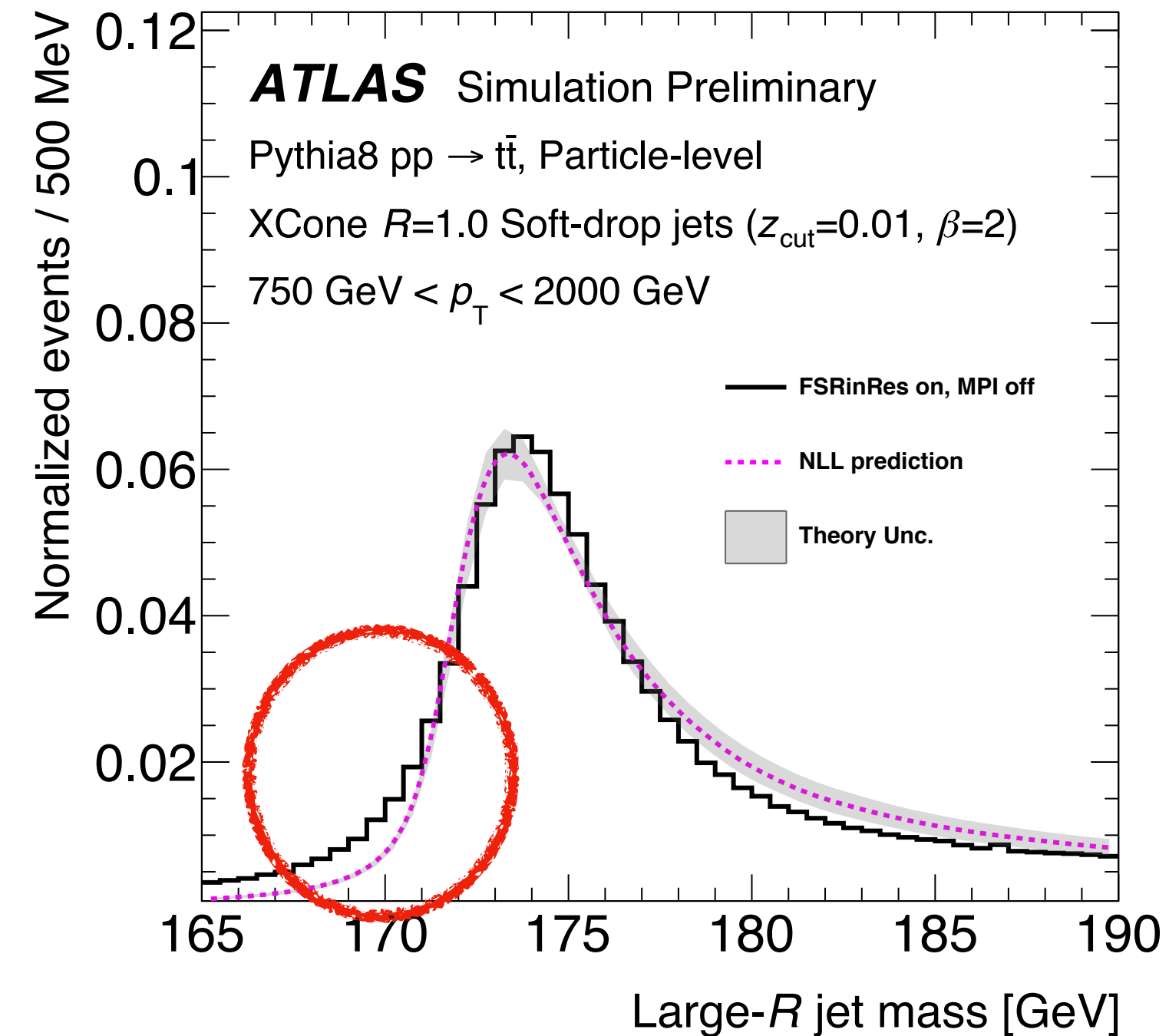
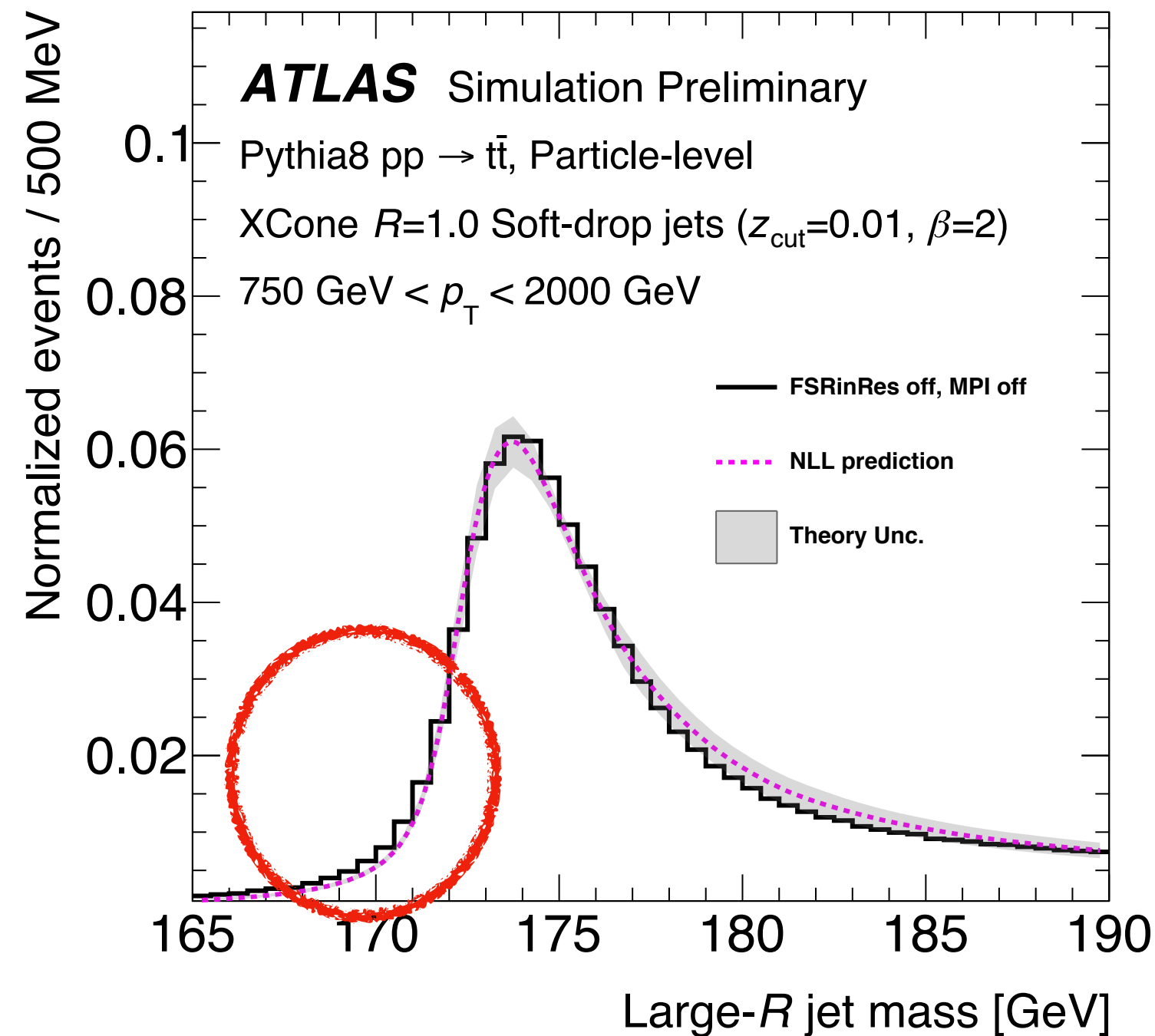


Thanks for your attention!

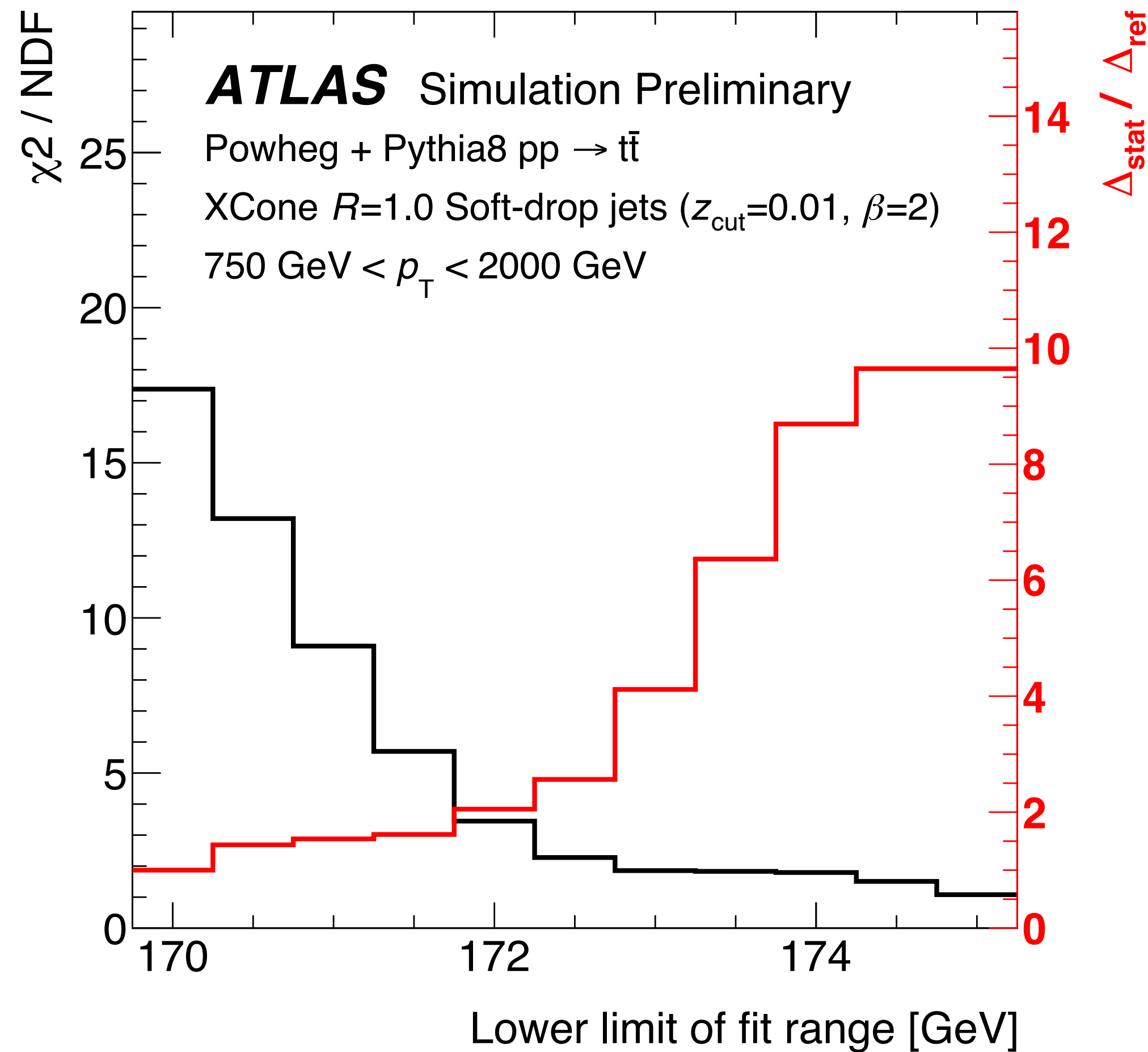
BACK-UP

Theory validity

- Theory gives an inclusive treatment of the top decay products
- In some MC events, FSR of decay products are groomed away (FSRinRes setting), lifting up the left tail of the jet mass peak
- This can be handled by carefully adjusting the fit range



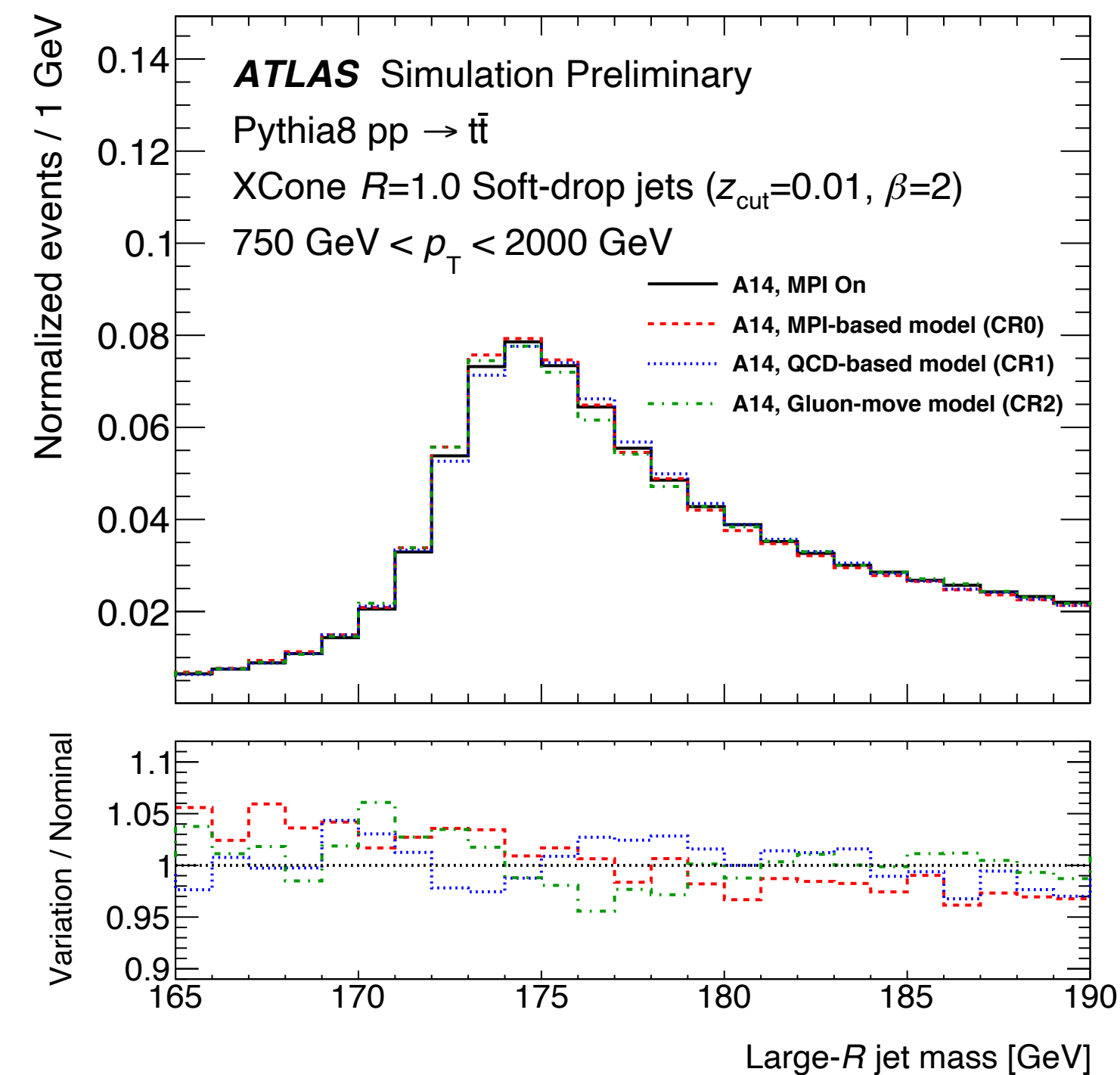
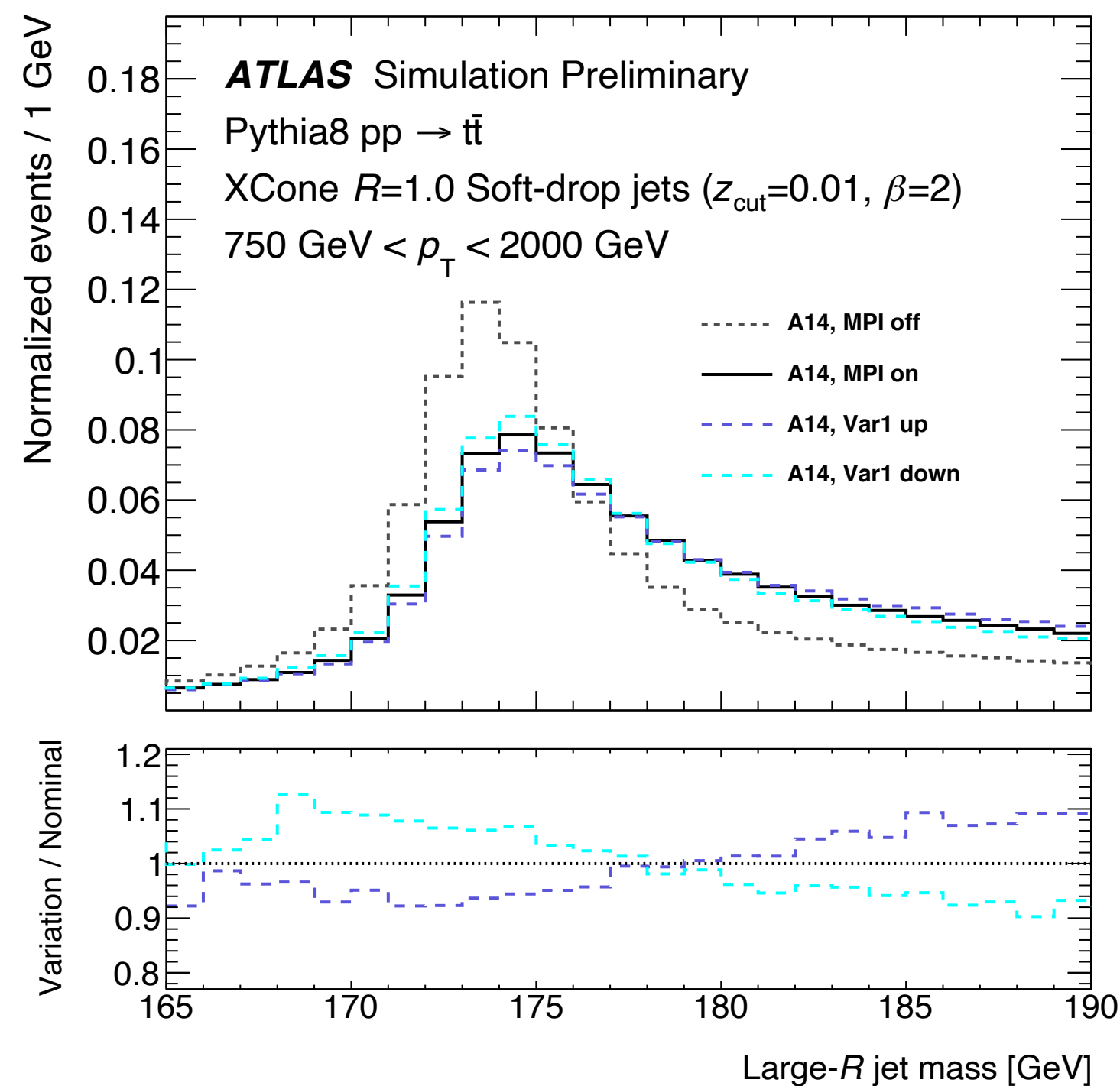
Fit range choice



- Fit range set to 172.5 - 180.0 GeV
- Lower limit is identified with the $m_{\text{top}}^{\text{MC}}$ the events were generated with.
- Fit is repeated with this value varied up and down by half GeV. Difference is registered as part of the methodological uncertainty.
- Variations of the upper limit have negligible impact on the fit result

Underlying Event treatment

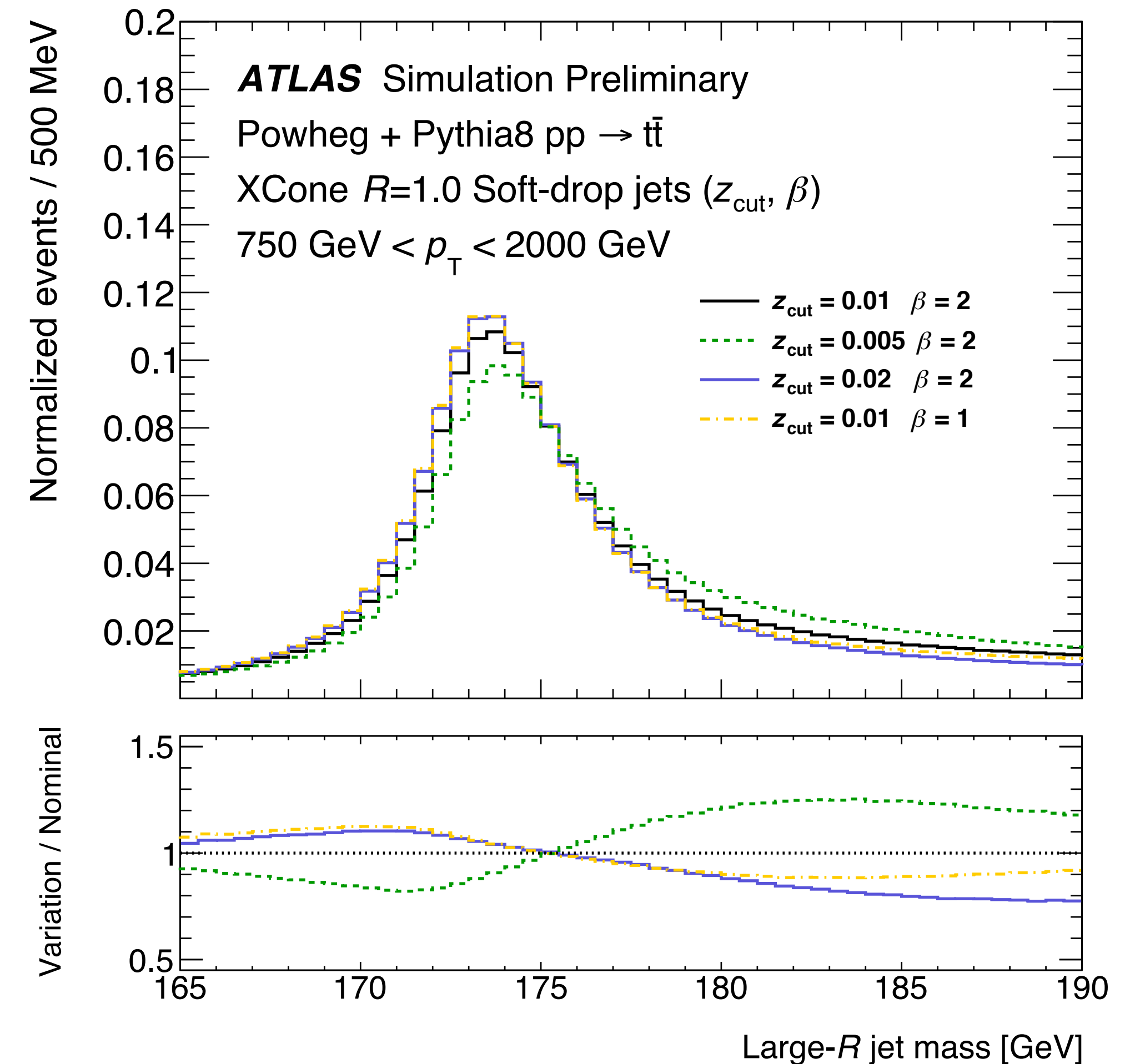
- Since theory cannot describe properly UE effects, another approach must be followed
- Var1 eigentune variations of A14 tune (devoted to provide coverage of UE modelling uncertainties) are considered along with alternative CR models
- MC-based templates used to estimate the shift in terms of the m_{top}^{MC} in the nominal MC sample due to such alternative samples



Stability

- The **MSR-MC mass relations** obtained could be **applied to direct measurements** in order to correct for universal effects (those independent from observables, kinematic regime, selection...)
- Stability of the **MC calibration** is studied for a number of related observables accessible with the first-principles calculation used

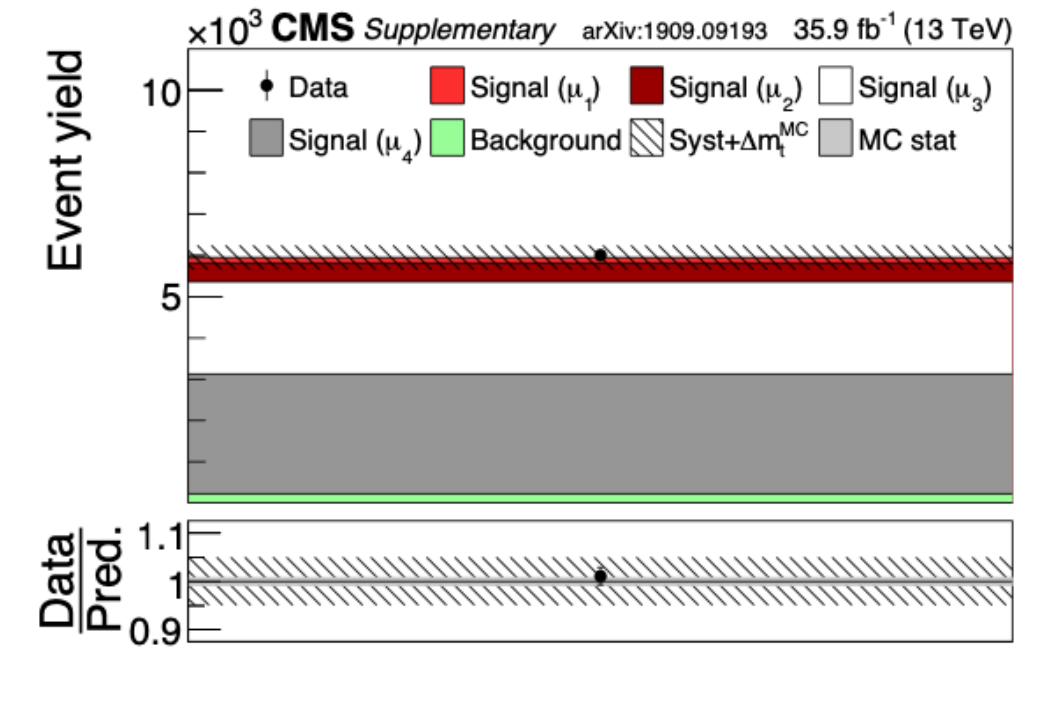
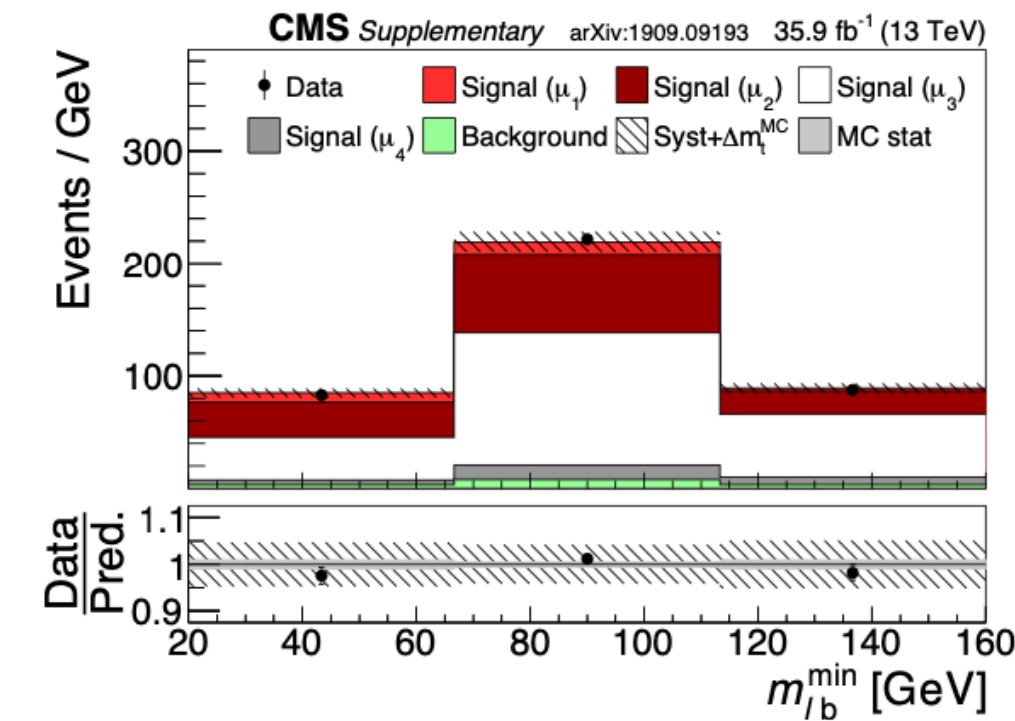
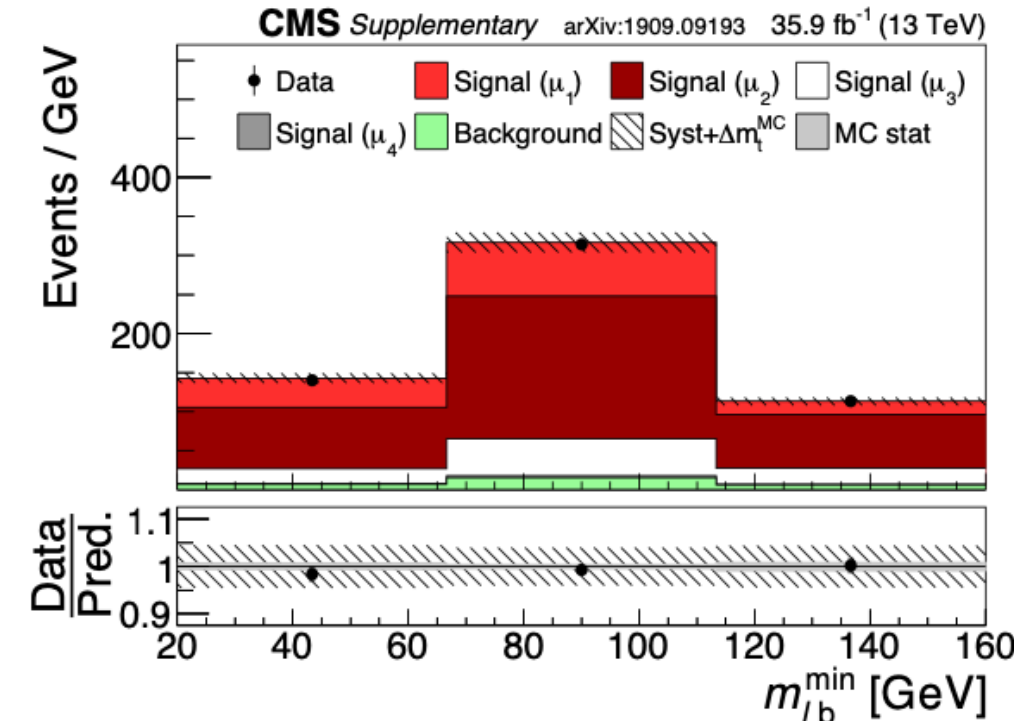
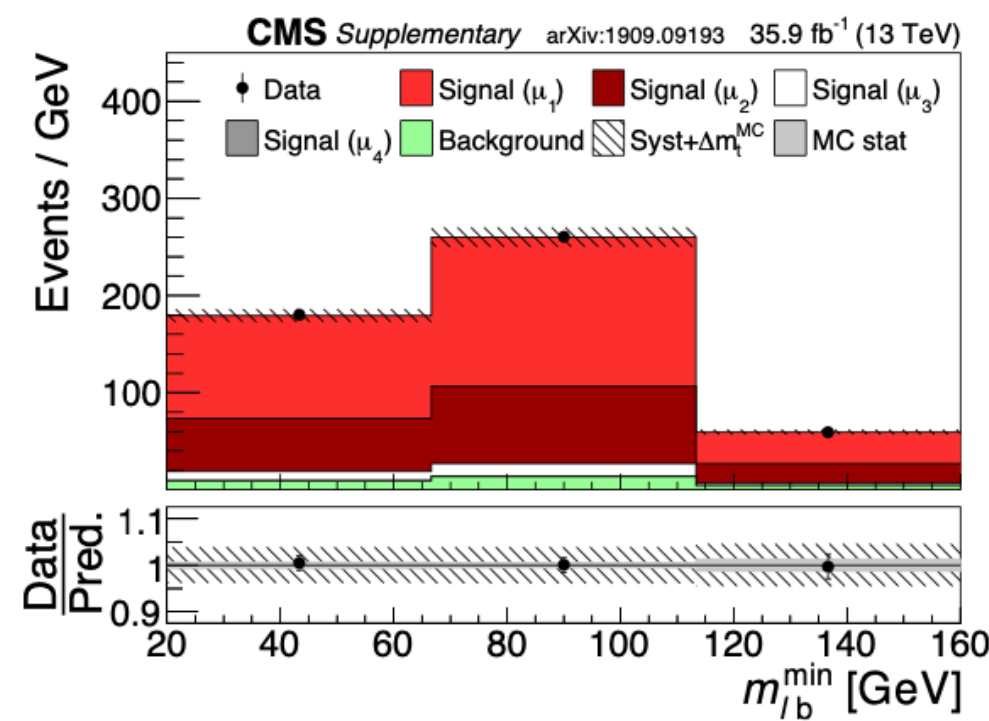
Aspects of the jet reconstruction and grooming technique are varied
- **The relation is found to be stable within 200 MeV**, within the restrictions of the underlying theoretical framework



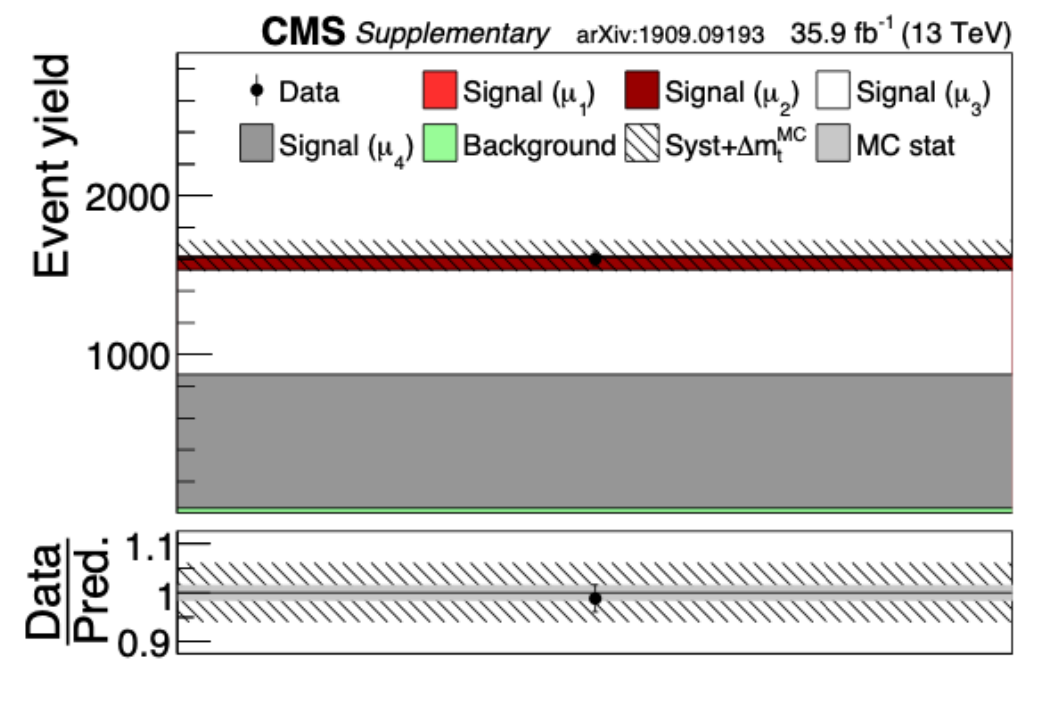
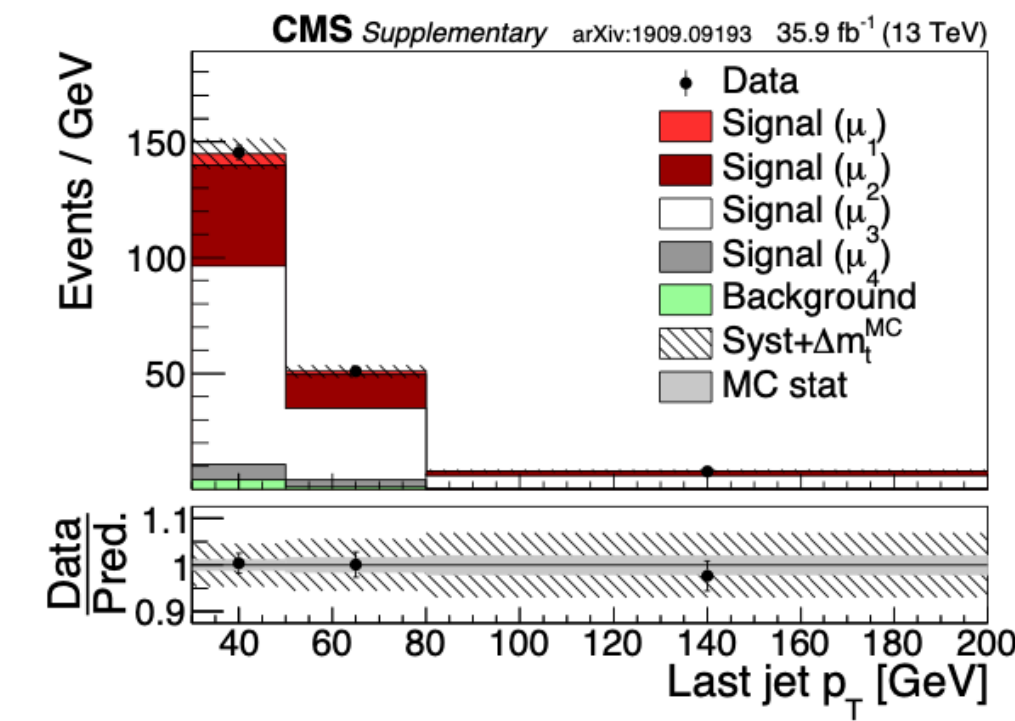
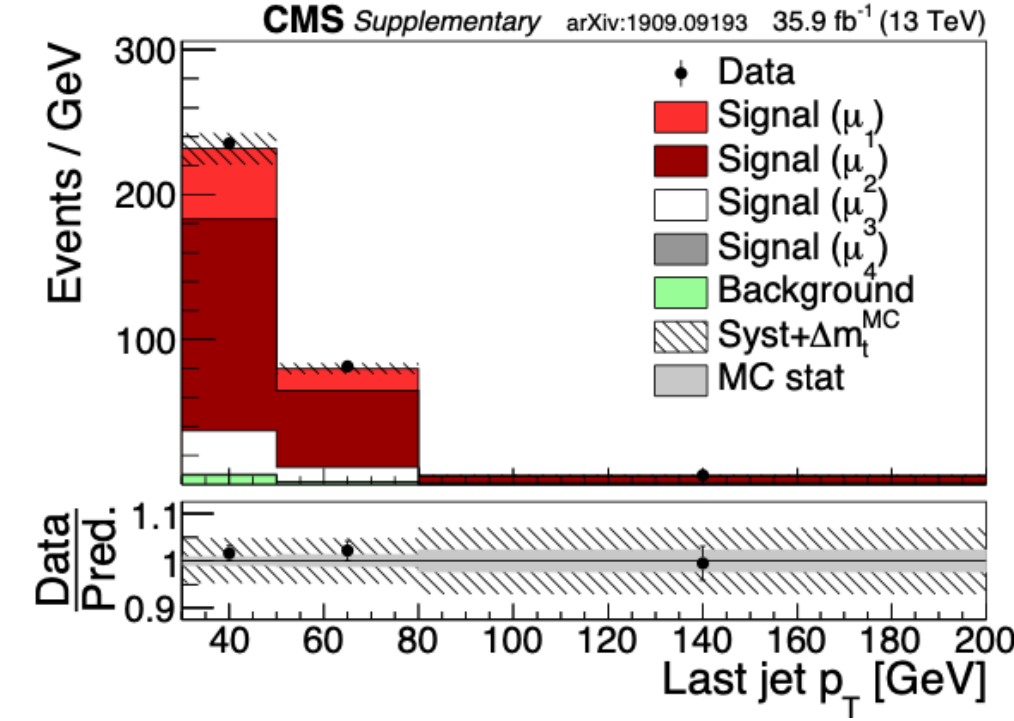
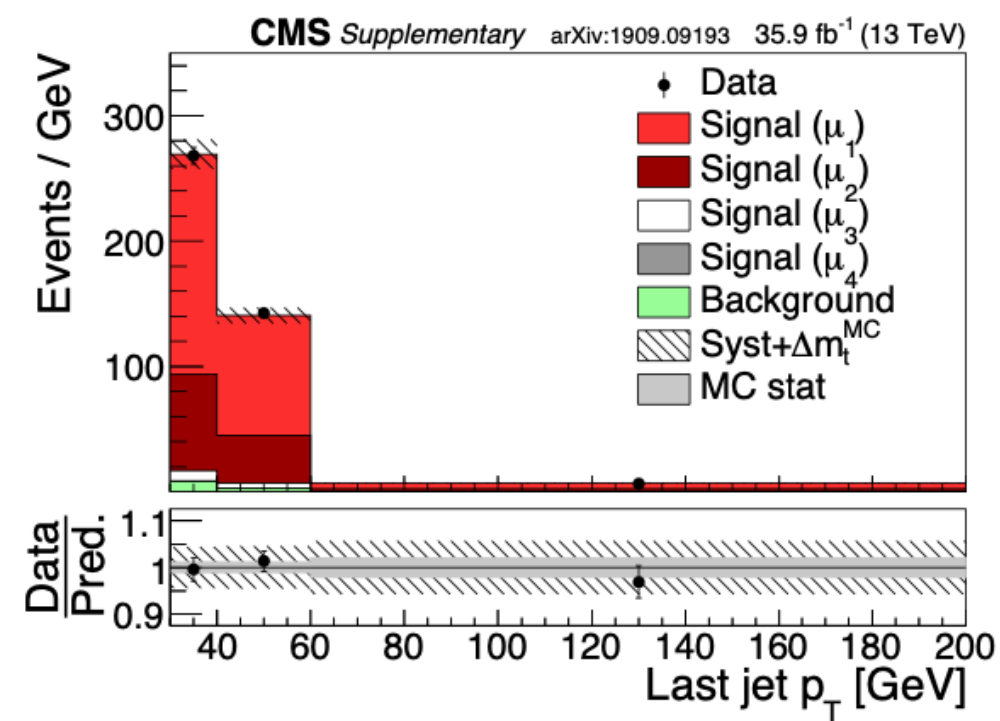
Post-fit distributions of final-state observables

$$m_{t\bar{t}}^{reco} \in \{0,420\} \text{ GeV} \quad \{420,550\} \text{ GeV} \quad \{550,810\} \text{ GeV} \quad \{810,\infty\} \text{ GeV}$$

$N_{bjet} = 1$



$N_{bjet} = 2$



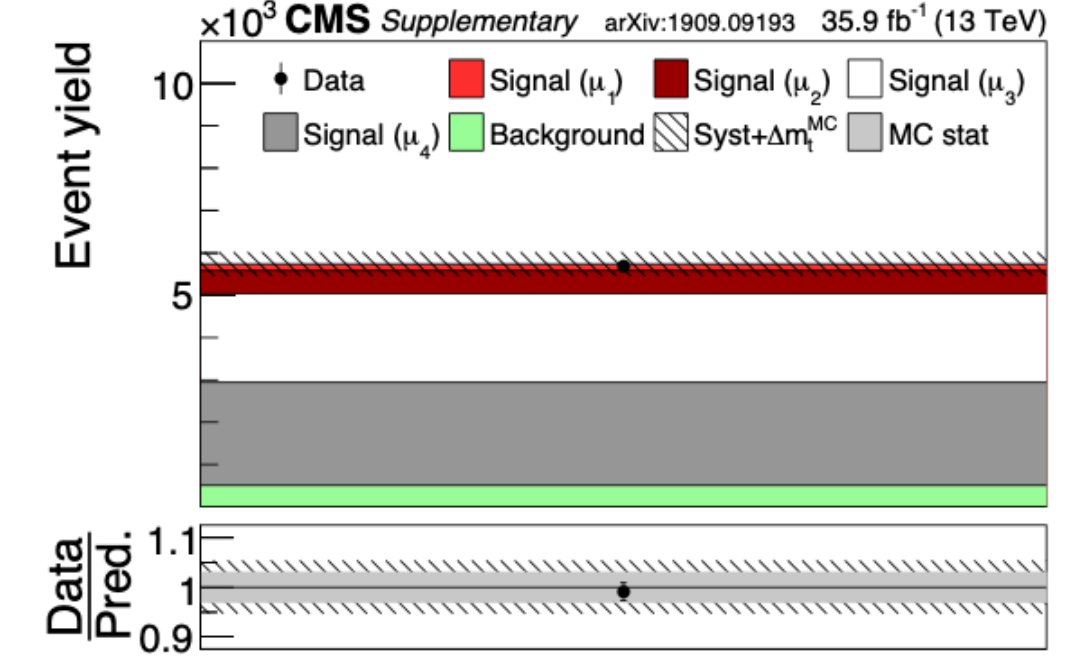
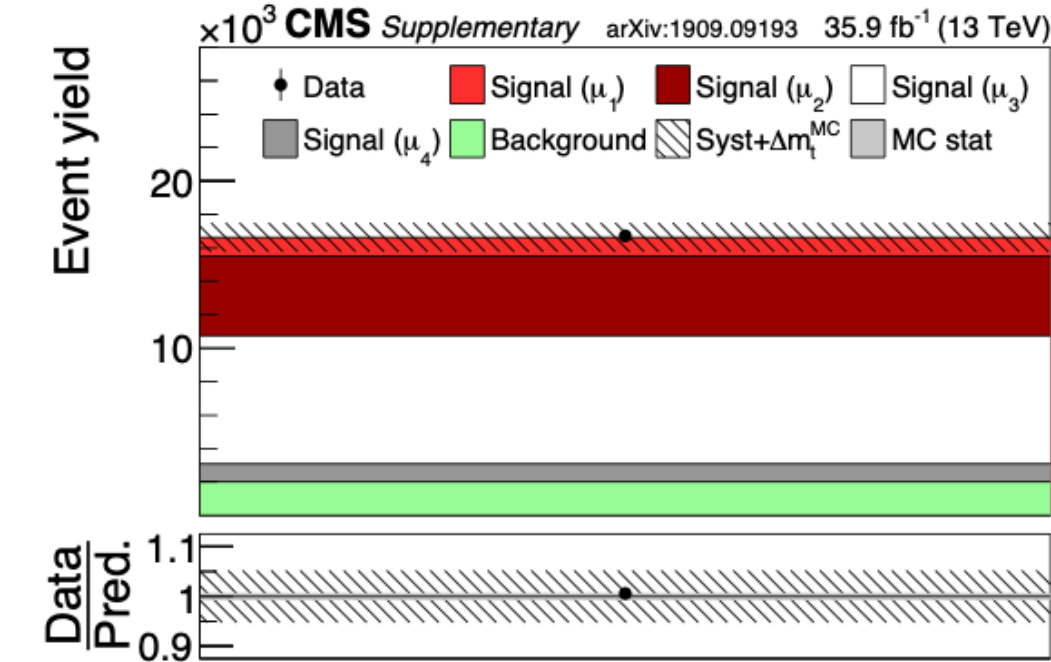
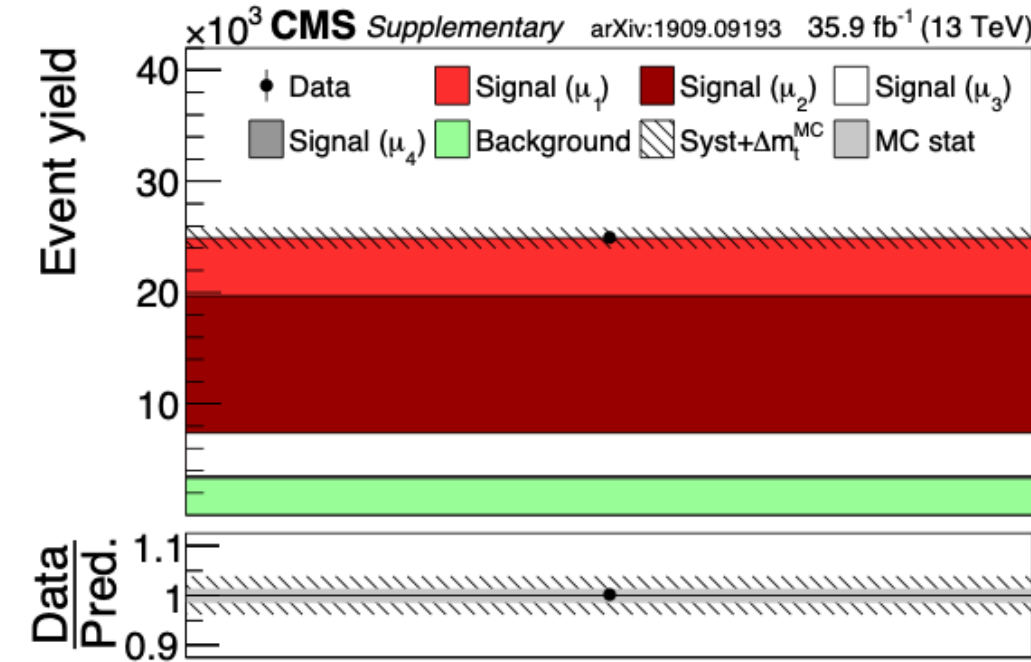
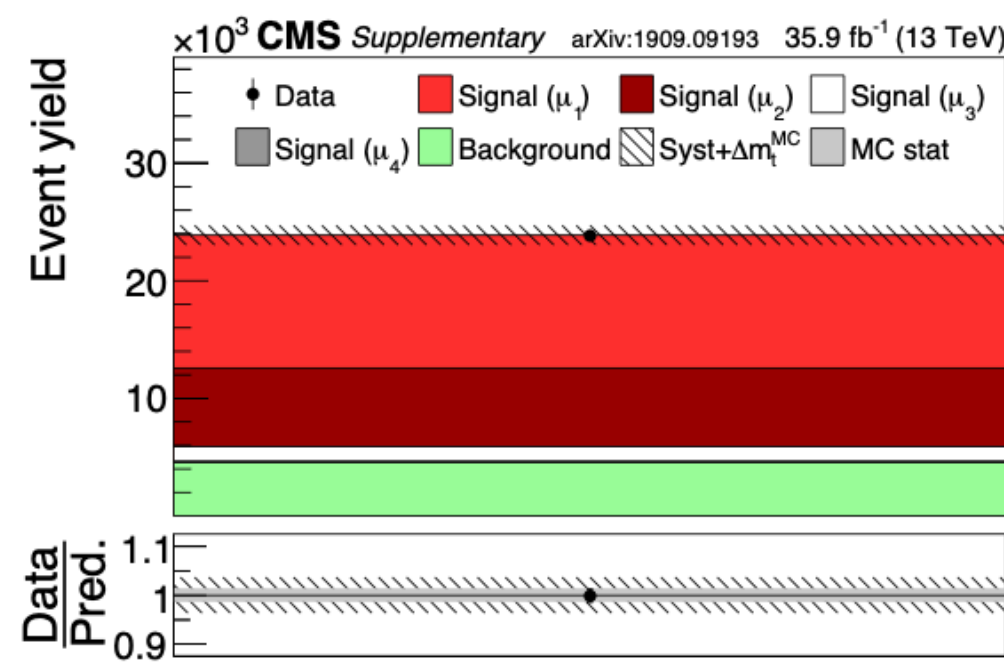
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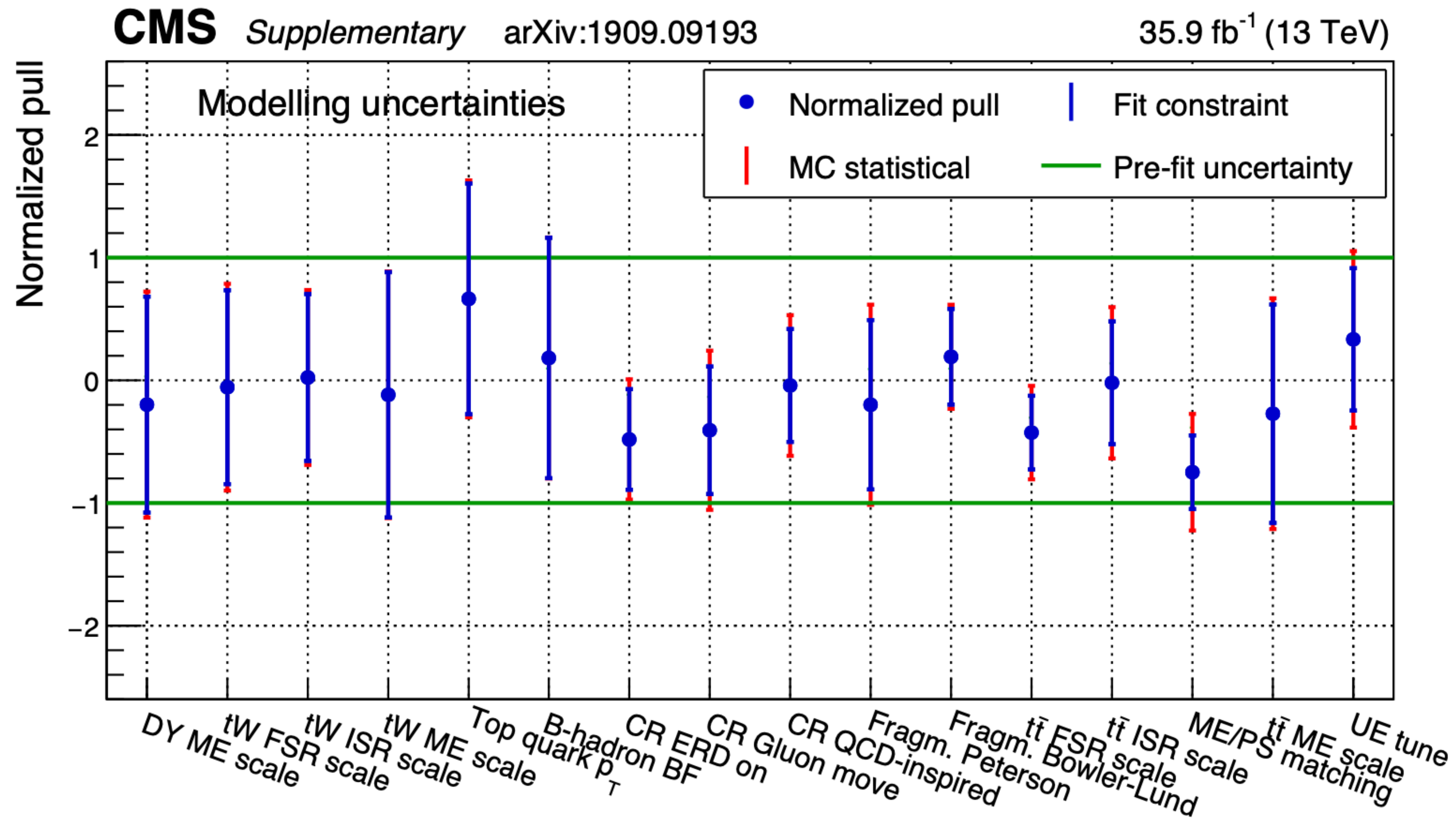
$$N_{bjet} = 0$$

&

$$N_{bjet} > 2$$



Pulls and constraints of nuisance parameters



Correlations of measured differential cross-section

$$\sigma_{t\bar{t}}^{(\mu_1)} = 255 \pm 11 \text{ (syst)} \pm 2 \text{ (stat) pb,}$$

$$\sigma_{t\bar{t}}^{(\mu_2)} = 315 \pm 15 \text{ (syst)} \pm 2 \text{ (stat) pb,}$$

$$\sigma_{t\bar{t}}^{(\mu_3)} = 181 \pm 9 \text{ (syst)} \pm 1 \text{ (stat) pb,}$$

$$\sigma_{t\bar{t}}^{(\mu_4)} = 50 \pm 3 \text{ (syst)} \pm 1 \text{ (stat) pb.}$$

	$\sigma_{t\bar{t}}^{(\mu_1)}$	$\sigma_{t\bar{t}}^{(\mu_2)}$	$\sigma_{t\bar{t}}^{(\mu_3)}$	$\sigma_{t\bar{t}}^{(\mu_4)}$
$\sigma_{t\bar{t}}^{(\mu_1)}$	1.00			
$\sigma_{t\bar{t}}^{(\mu_2)}$	0.64	1.00		
$\sigma_{t\bar{t}}^{(\mu_3)}$	0.72	0.60	1.00	
$\sigma_{t\bar{t}}^{(\mu_4)}$	0.32	0.65	0.47	1.00

Inclusive and differential measurements

