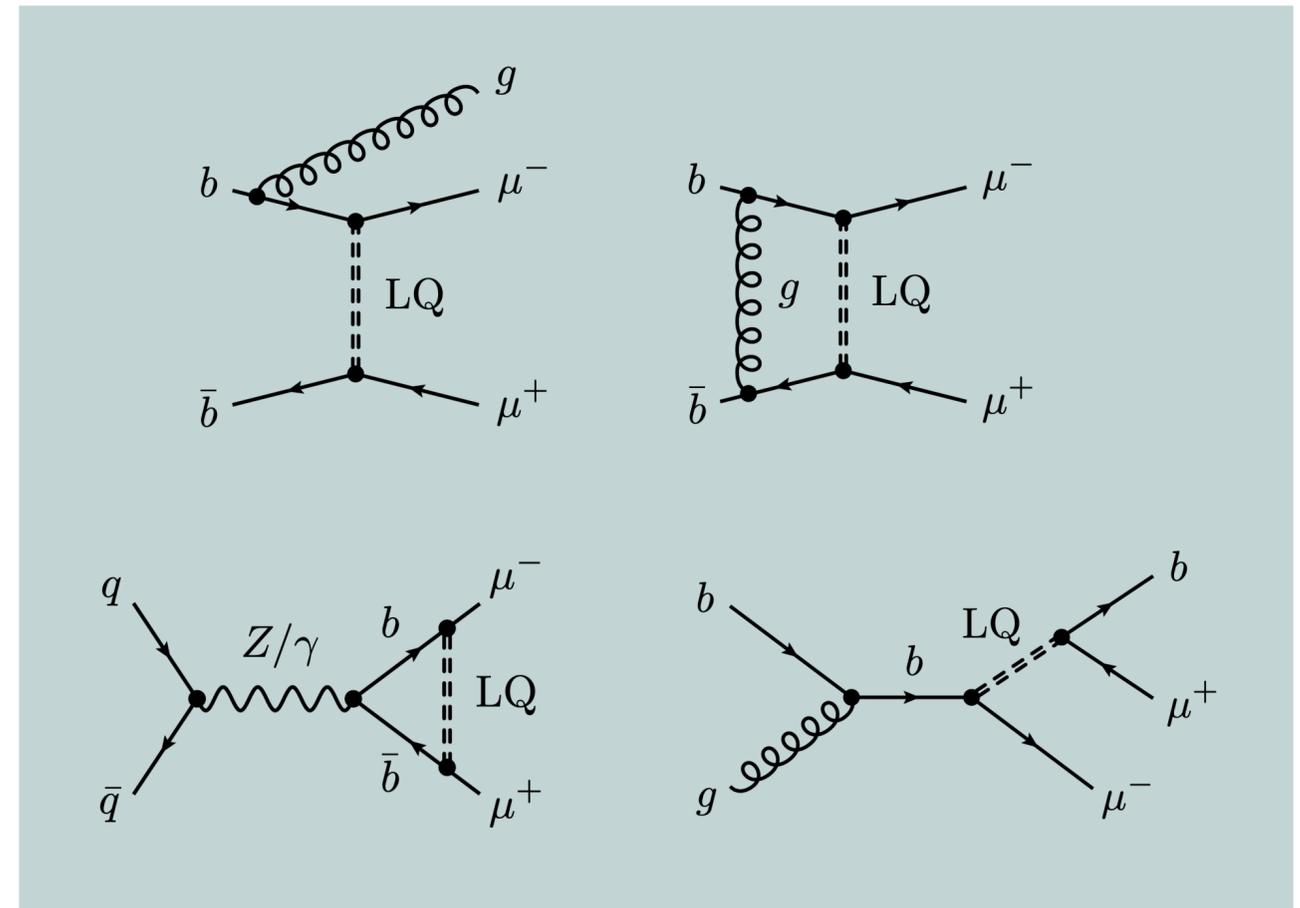
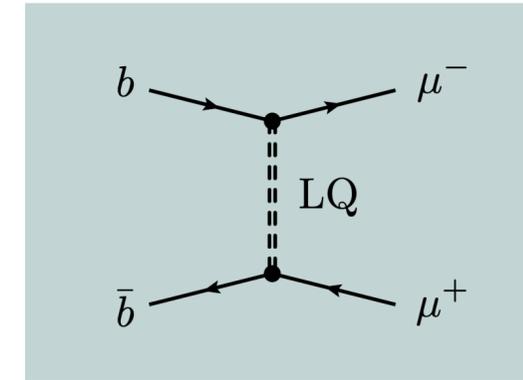


On Drell-Yan Production of Scalar Leptoquarks coupling to Heavy Quark Flavours

Luc Schnell
LF(U)V Workshop
July 4, 2022



1. Introduction

1.1 Motivation

1.2 Drell-Yan

1.3 Going beyond the leading order

1. Introduction

1.1 Motivation

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Hints for lepton
flavour universality
violation (LFUV)

$b s \mu \mu$

$b c \tau \nu$

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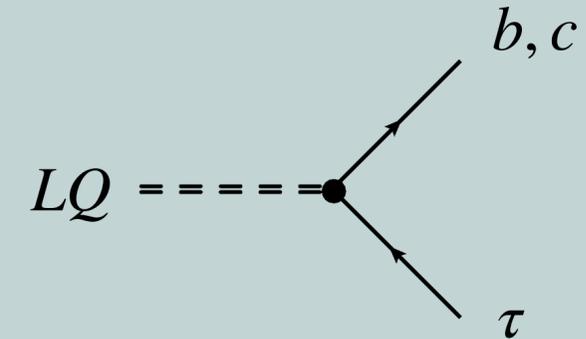
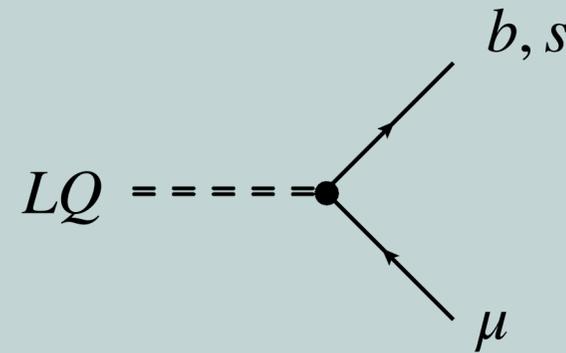
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- Can be explained with **leptoquarks (LQs)** couplings to **heavy quark flavours**.



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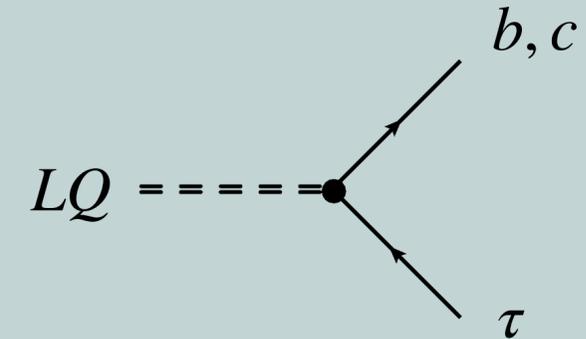
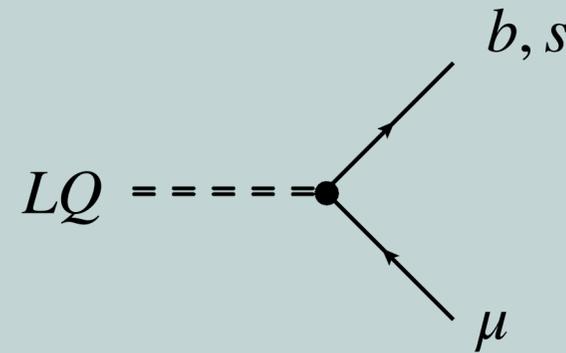
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B physics, LFV, meson mixing, Z couplings to fermions, AMMs, EDMs, PV and many more.

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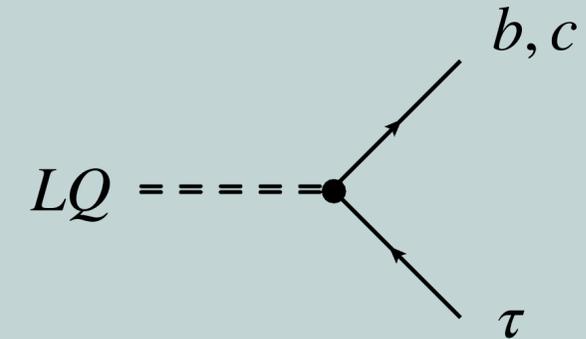
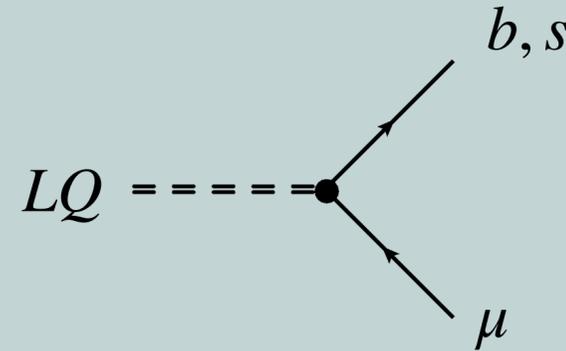
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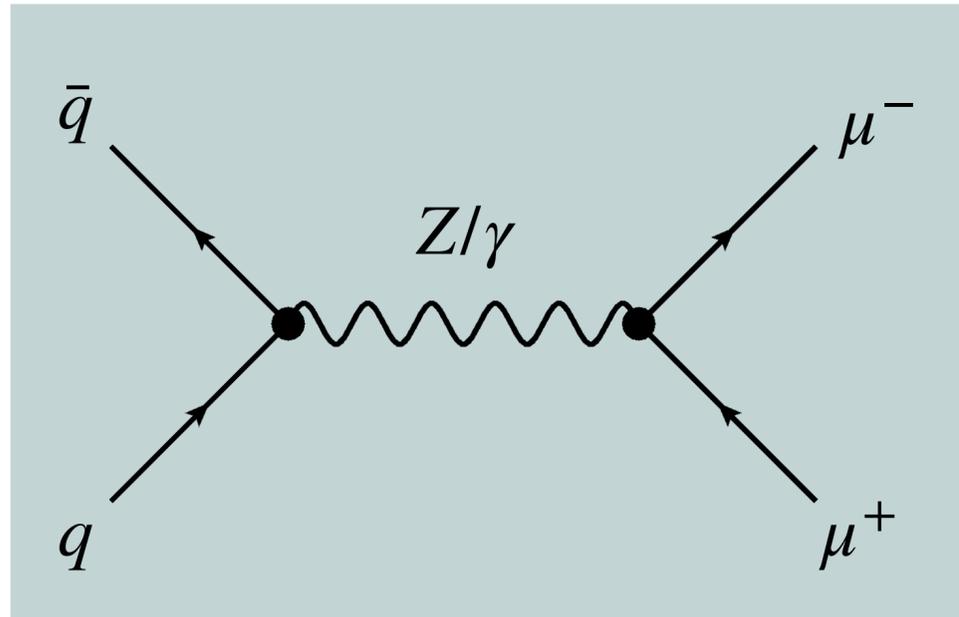
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High-energy searches at LHC

Pair production, single (resonant) production and **Drell-Yan**.

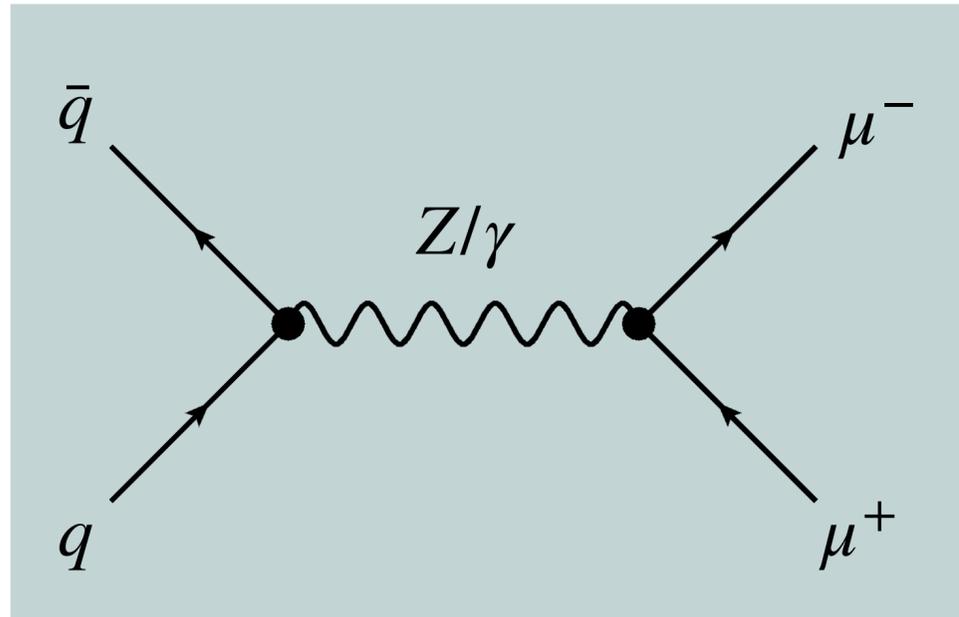
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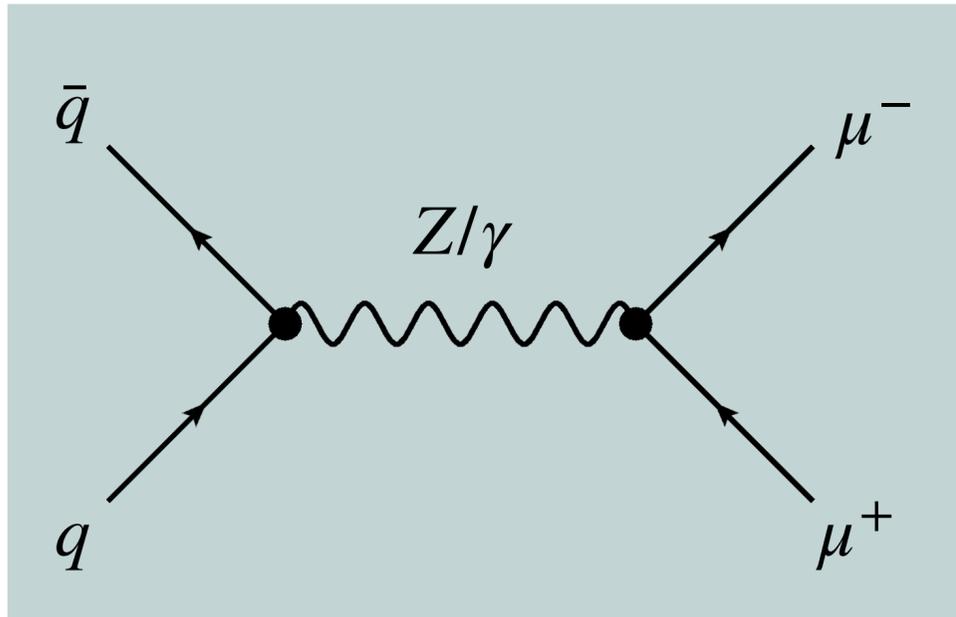
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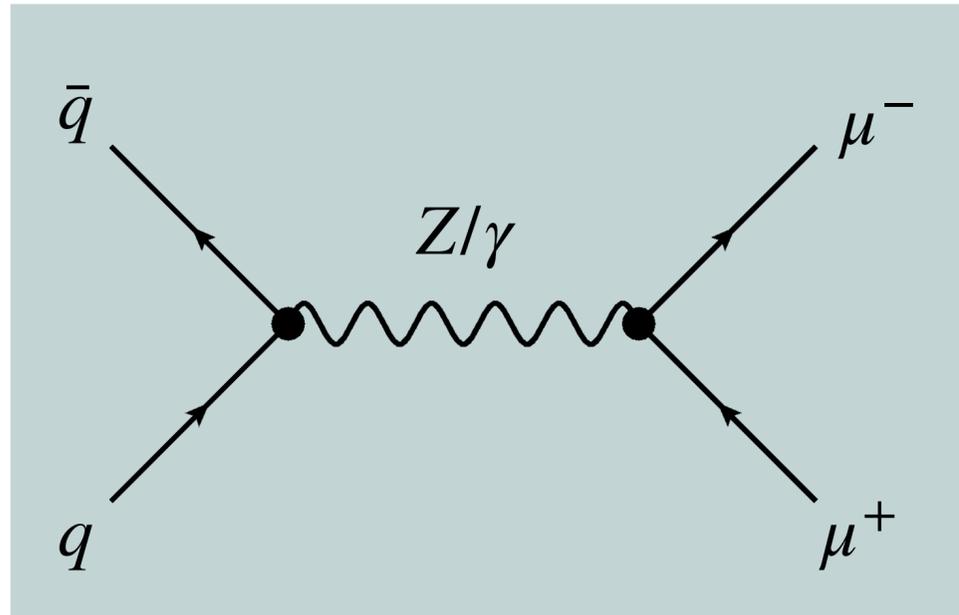
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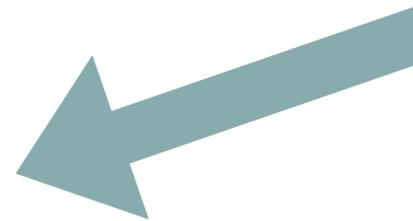
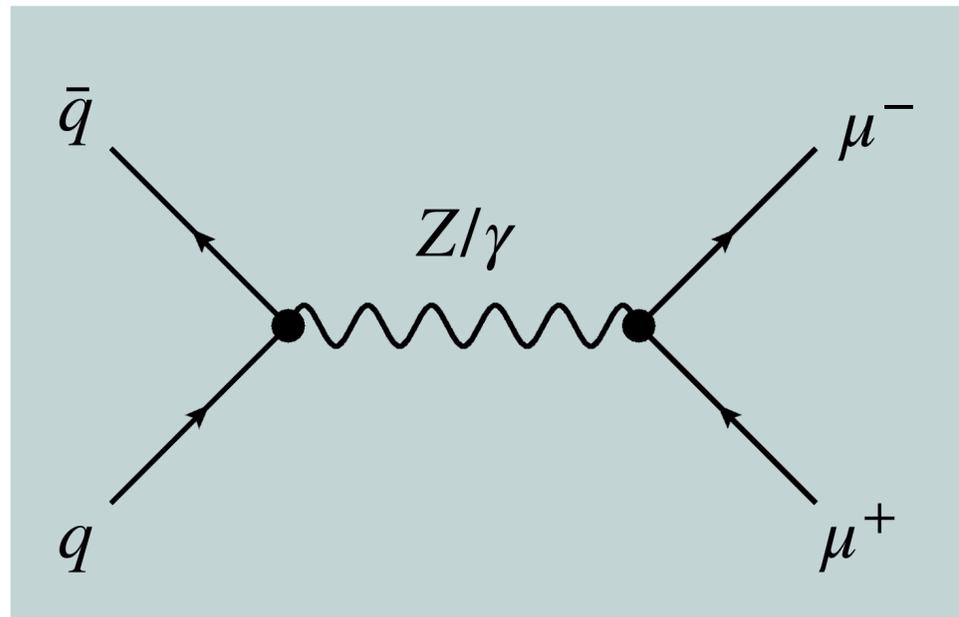
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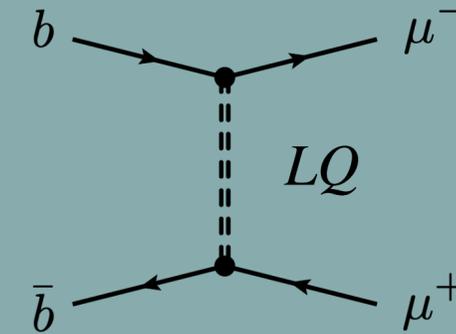
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1. Introduction

1.2 Drell-Yan



- LQs contribute at leading order via **non-resonant t -channel exchange**.

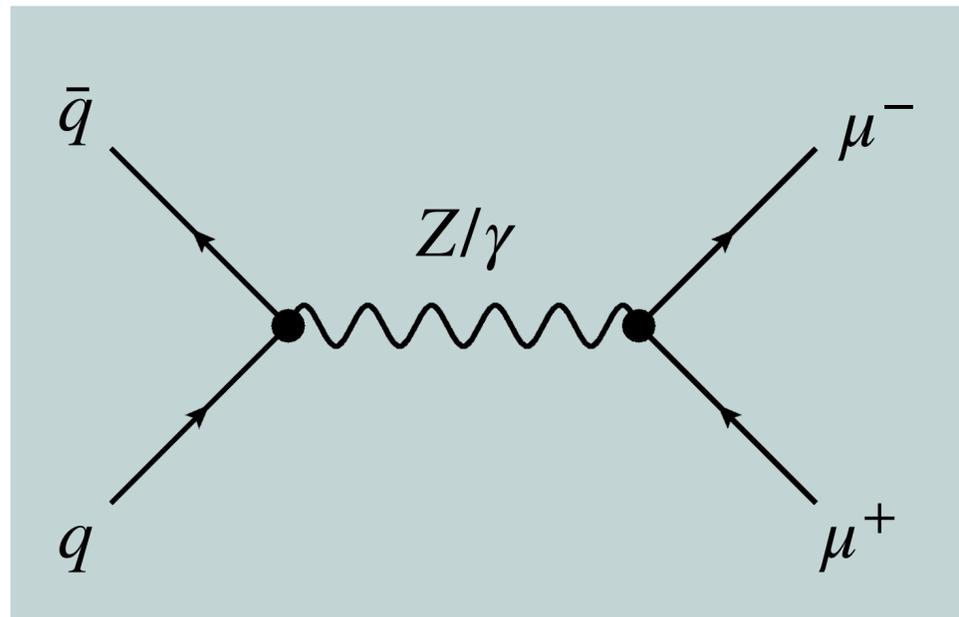


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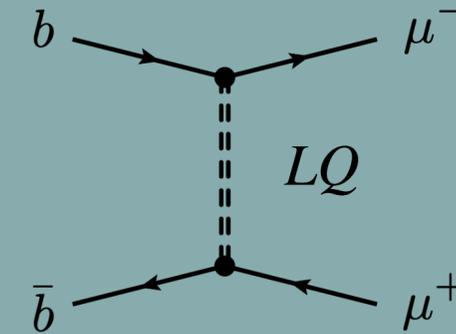
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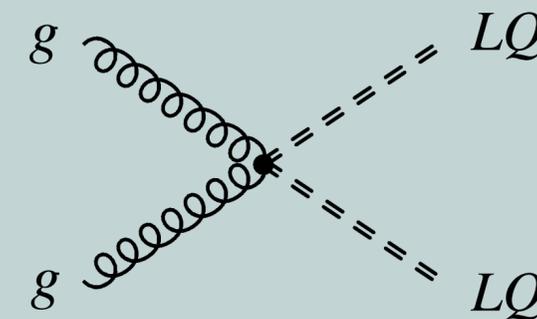


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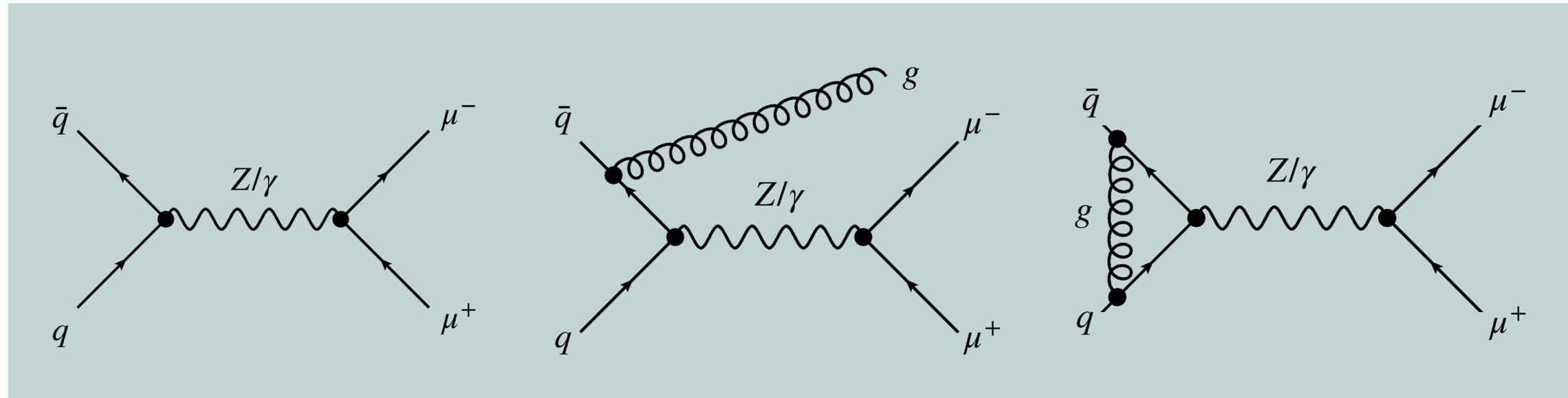


1. Introduction

1.3 Going beyond the leading order

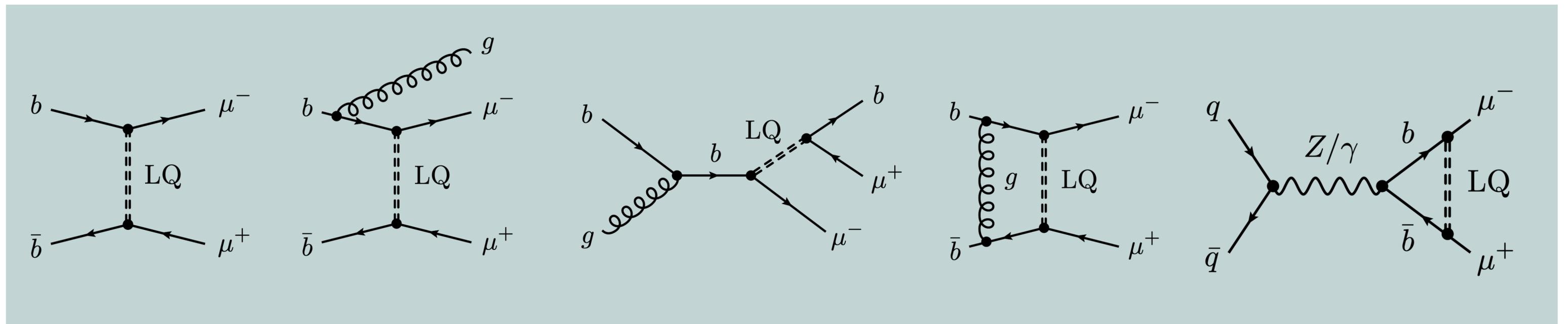
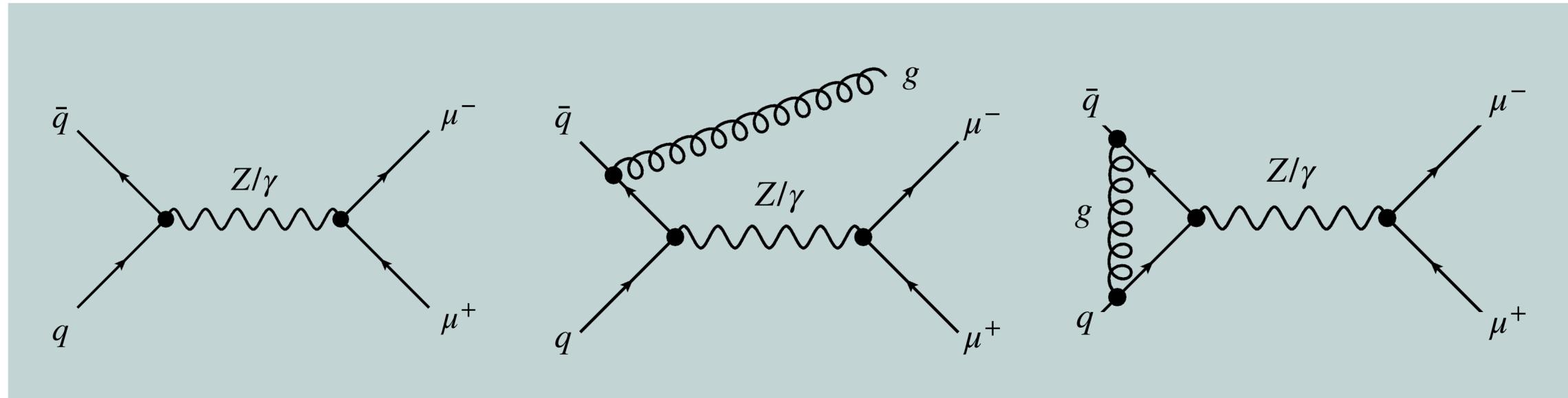
1. Introduction

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1.3 Going beyond the leading order



2. Theoretical framework and calculations

2.1 Simplified LQ model

2.2 Matrix elements

2.3 POWHEG-Box (in a nutshell)

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- We consider the **simplified LQ model**

$$\mathcal{L} \supset Y_{ul} \bar{u}^c \ell S_1^\dagger + Y_{dl} \bar{d}^c \ell \tilde{S}_1^\dagger + \text{h.c.},$$

involving the two $SU(2)_L$ singlets $S_1 \sim (\mathbf{3}, \mathbf{1}, -1/3)$ and $\tilde{S}_1 \sim (\mathbf{3}, \mathbf{1}, -4/3)$ coupling to right-handed fermions.

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- The size of the modifications in $pp \rightarrow \ell^+ \ell^-$ primarily depends on the **flavour structure** and the **magnitude** of $Y_{u\ell}, Y_{d\ell}$.
- The representation of the LQ plays a role only in the **interference** with the SM. S_1 (\tilde{S}_1) leads to destructive (constructive) interference, these interactions can therefore be used as a template to cover the **full space of scalar LQ models**.

2. Theoretical framework and calculations

2.2 Matrix elements

2. Theoretical framework and calculations

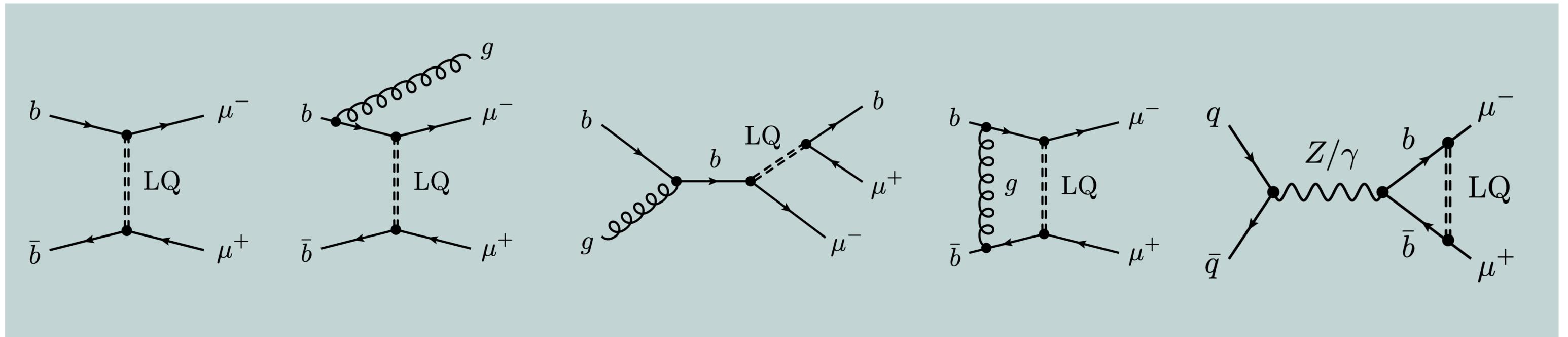
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- Calculation of the tree-level, real emission and virtual correction matrix elements relying on the full **FeynRules+FeynArts+FormCalc** machinery.
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1. Dealing with infrared and collinear divergences

- The virtual and real matrix elements have infrared and collinear divergences:

$$\sigma_{\text{NLO}} = \int d\Phi_n \mathcal{L} \left[\mathcal{B}(\Phi_n) + \mathcal{V}_b(\Phi_n) \right] + \int d\Phi_{n+1} \mathcal{L} \mathcal{R}(\Phi_{n+1})$$

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- These are cured by introducing appropriate counterterms and a clever subtraction of divergent terms in the singular regions of phase space (e.g. in the **FKS subtraction** scheme):

$$\begin{aligned} \bar{B}(\Phi_n) = & B(\Phi_n) + V(\Phi_n) \\ & + \left[\int d\Phi_{\text{rad}} [R(\Phi_{n+1}) - C(\Phi_{n+1})] + \int \frac{dz}{z} [G_{\oplus}(\Phi_{n,\oplus}) + G_{\ominus}(\Phi_{n,\ominus})] \right]_{\bar{\Phi}_n = \Phi_n}, \end{aligned}$$

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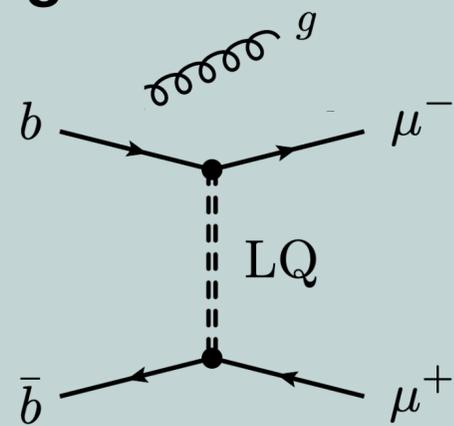
2. Avoiding double counting

2. Theoretical framework and calculations

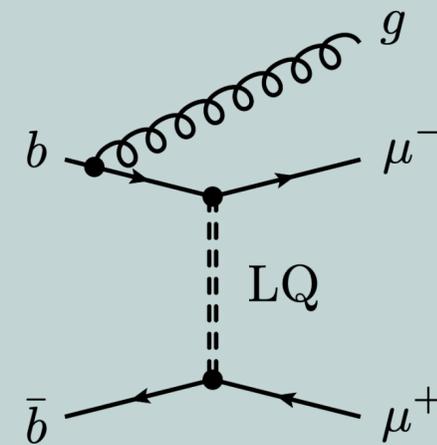
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LO + PS



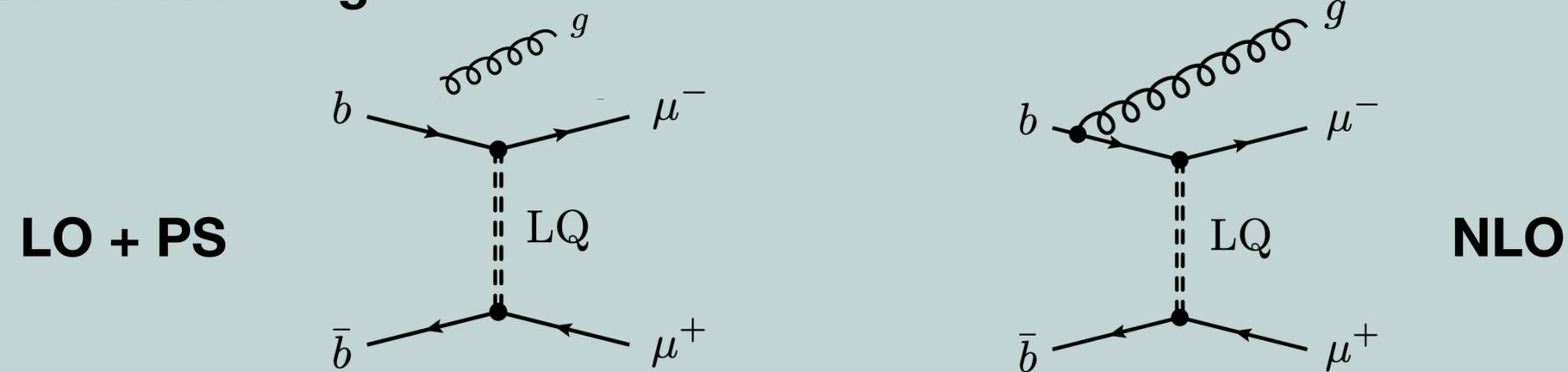
NLO



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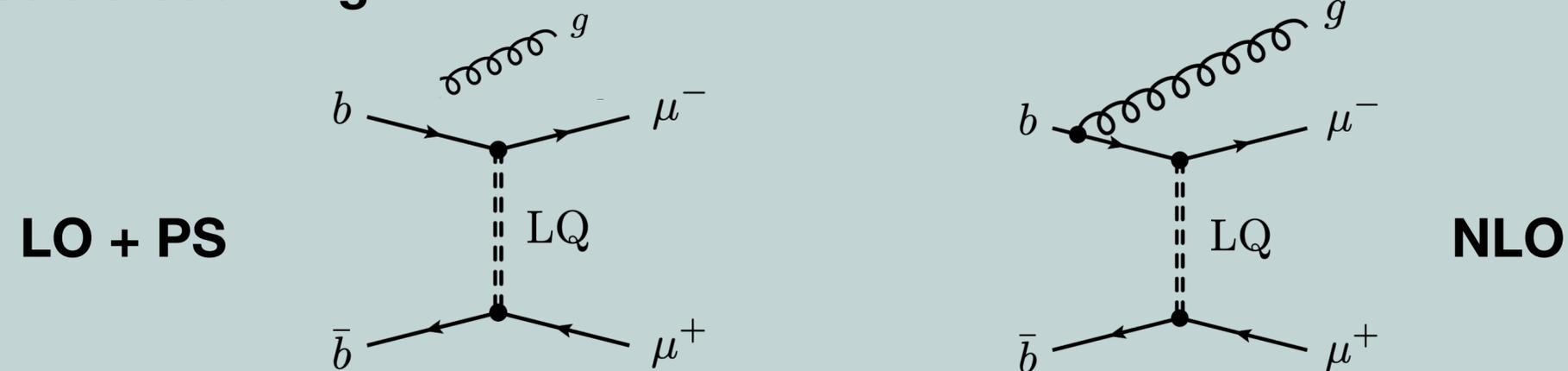


- **Jet-Matching:** generate tree-level and real emission events, run parton shower (PS) and reconstruct jets, veto jet if it cannot be matched to an underlying parton within some scale.
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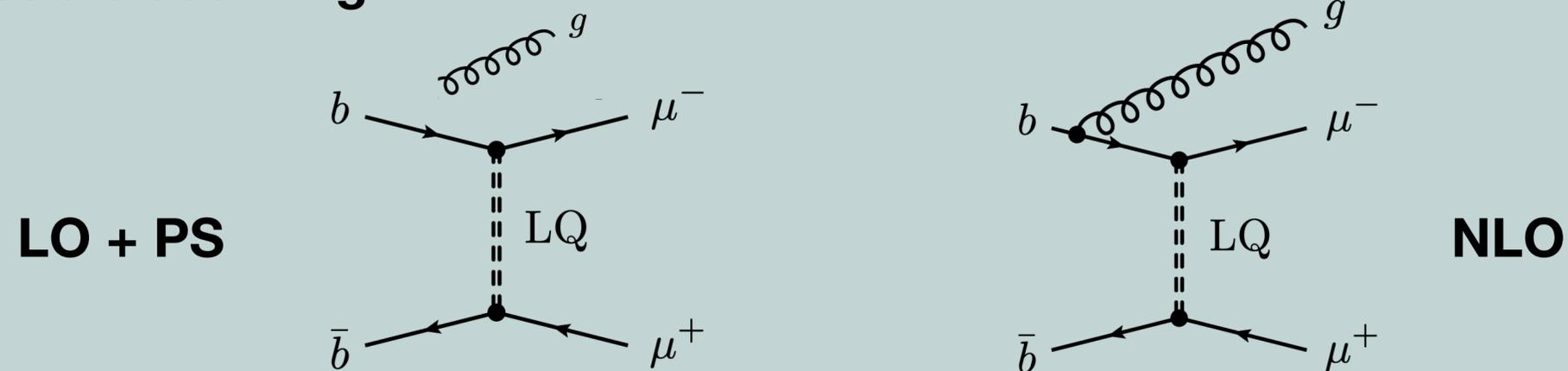


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2. Theoretical framework and calculations

2.3 POWHEG-Box (in a nutshell)

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- **Jet-Matching:** generate tree-level and real emission events, run parton shower (PS) and reconstruct jets, veto jet if it cannot be matched to an underlying parton within some scale.
→ **unphysical matching scale**
- The POWHEG framework uses a different strategy, **generating hardest emission events using full NLO accuracy.**
- The **PS** is then only needed for emissions below p_T^{min} .

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Matrix elements

Born.f
real.f
virtual.f

3. Phenomenology

3.1 Inclusive cross section

3.2 b -tags

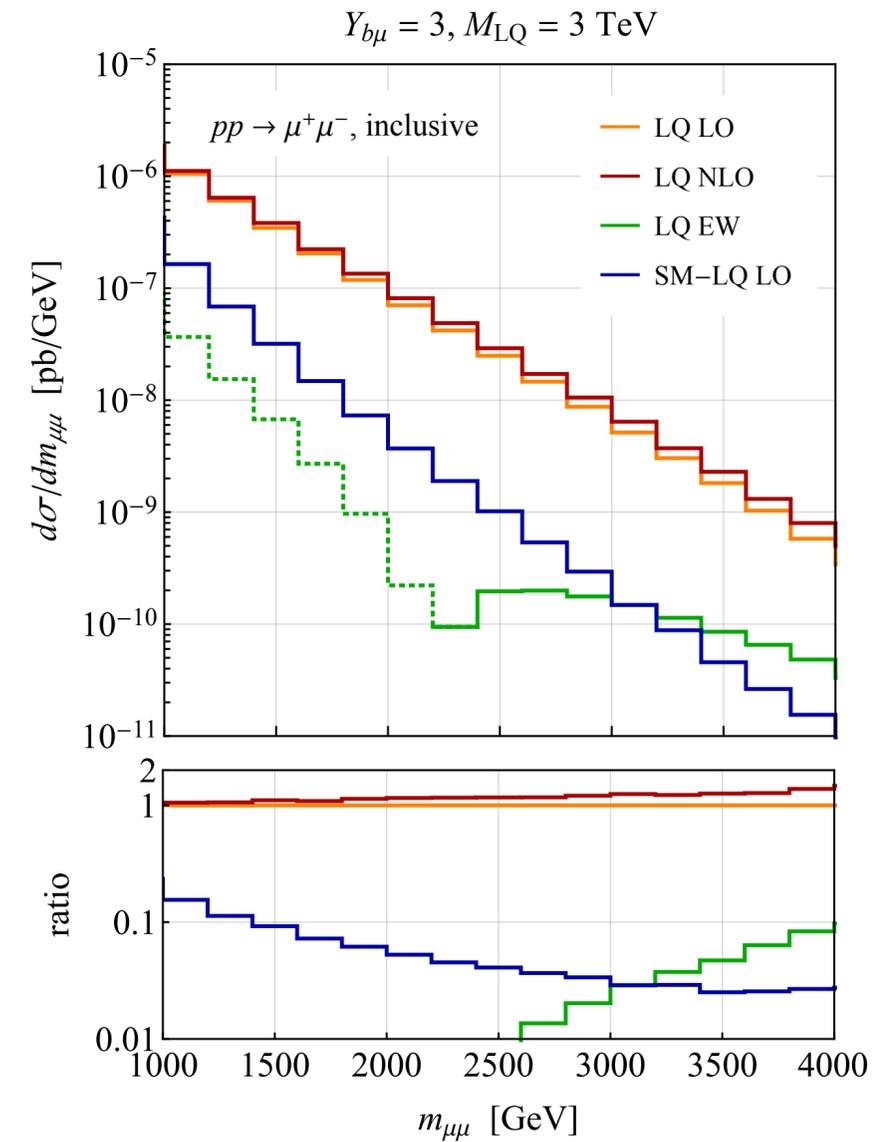
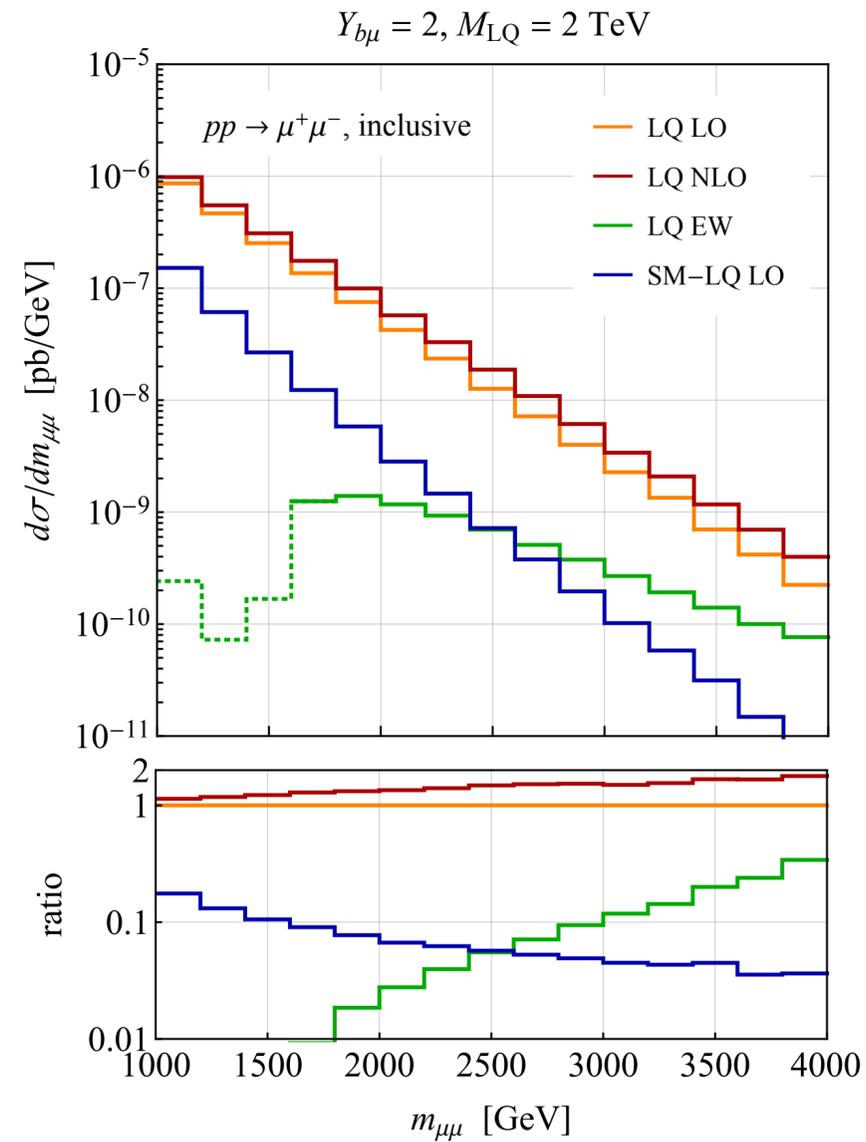
3.3 Exclusion limits

3. Phenomenology

3.1 Inclusive cross section

3. Phenomenology

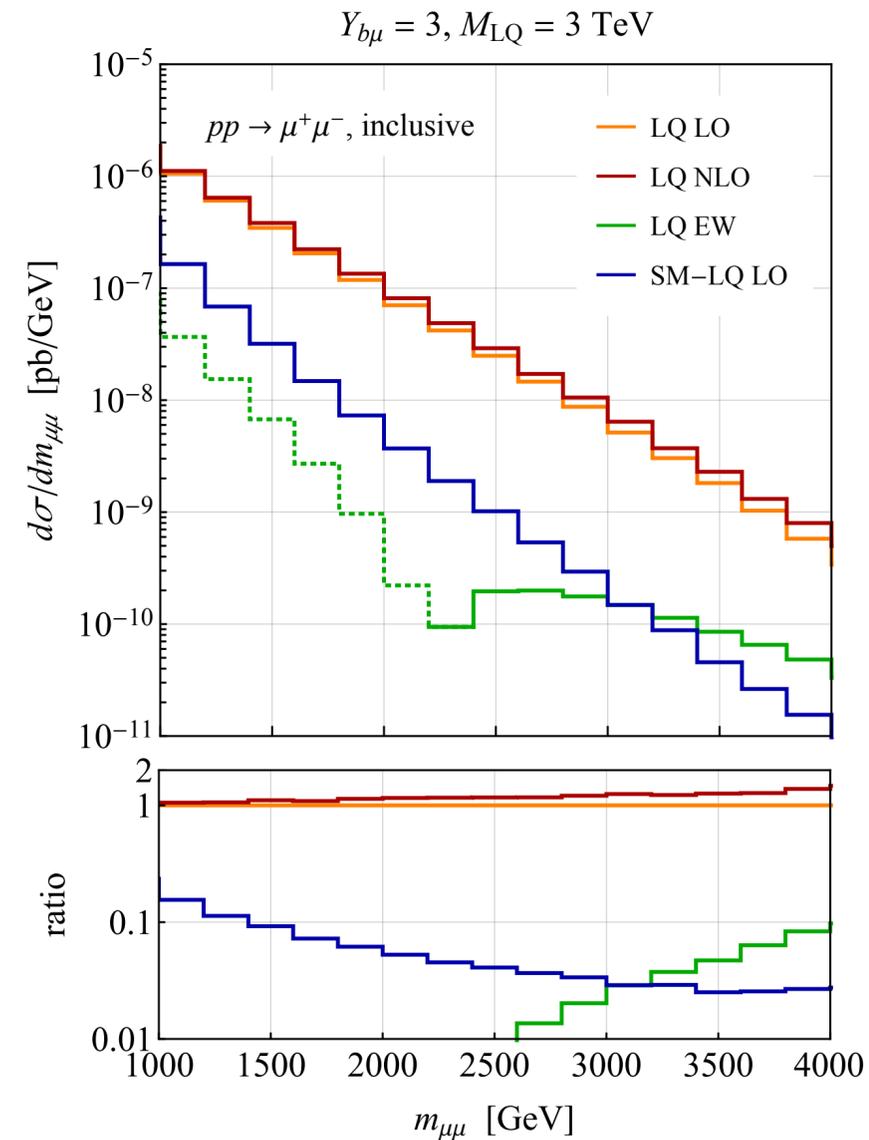
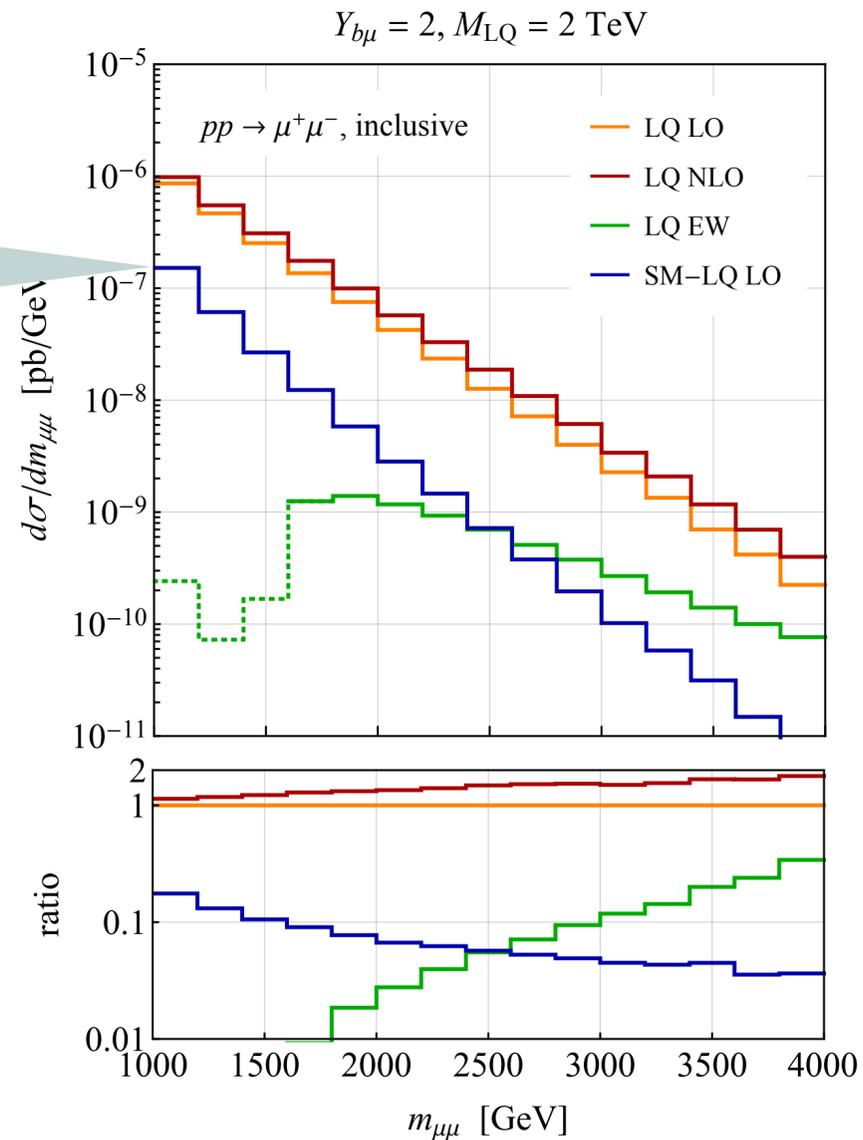
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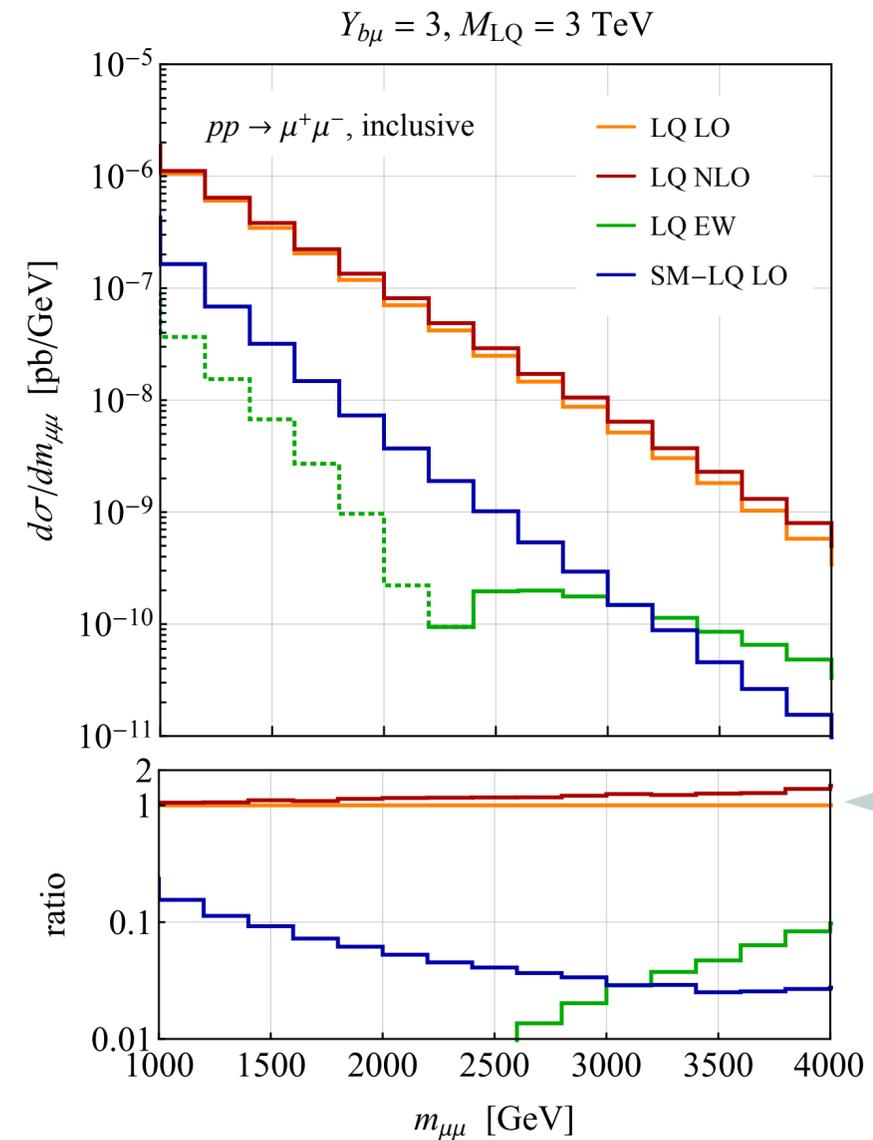
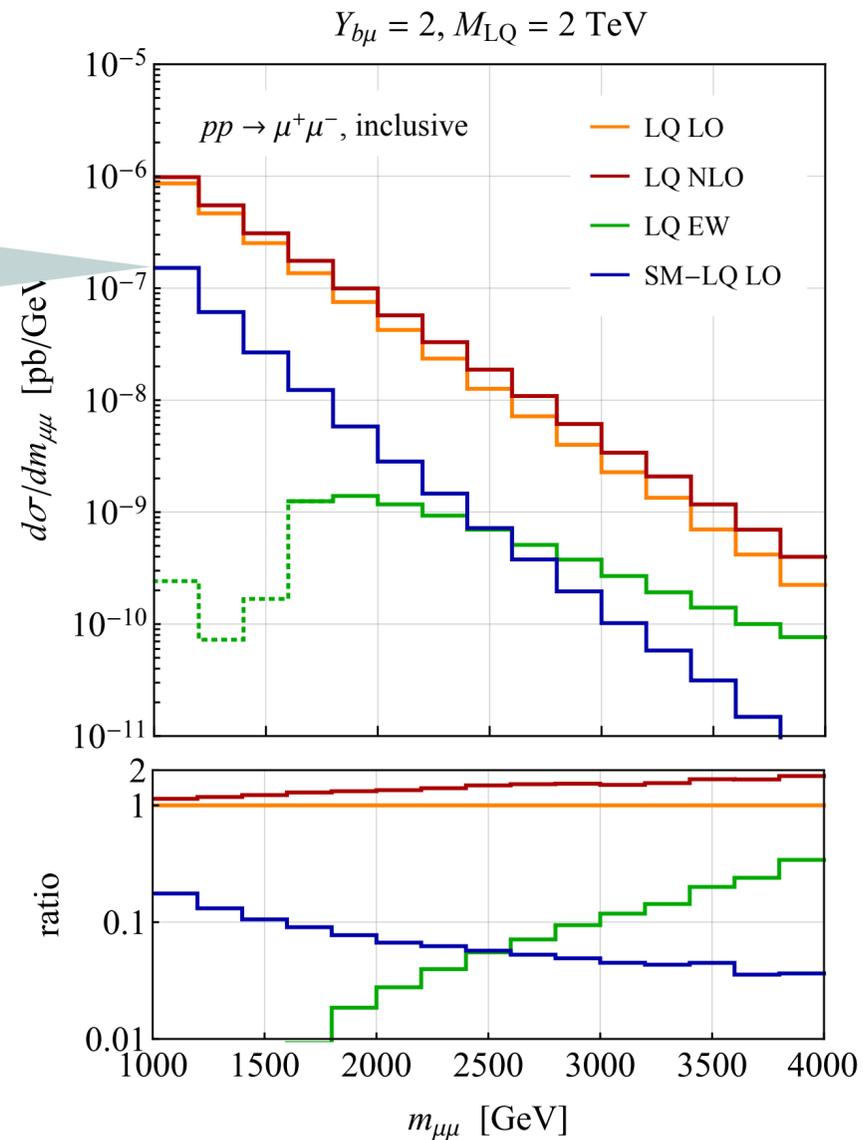
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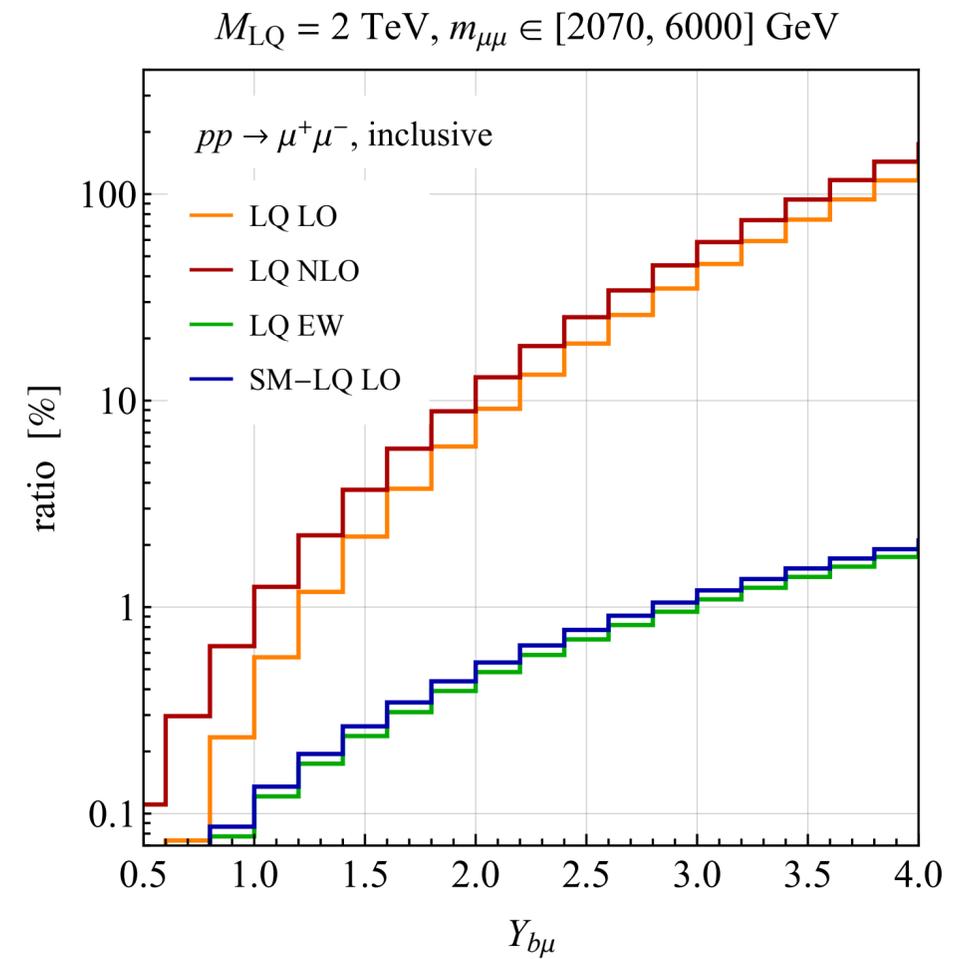
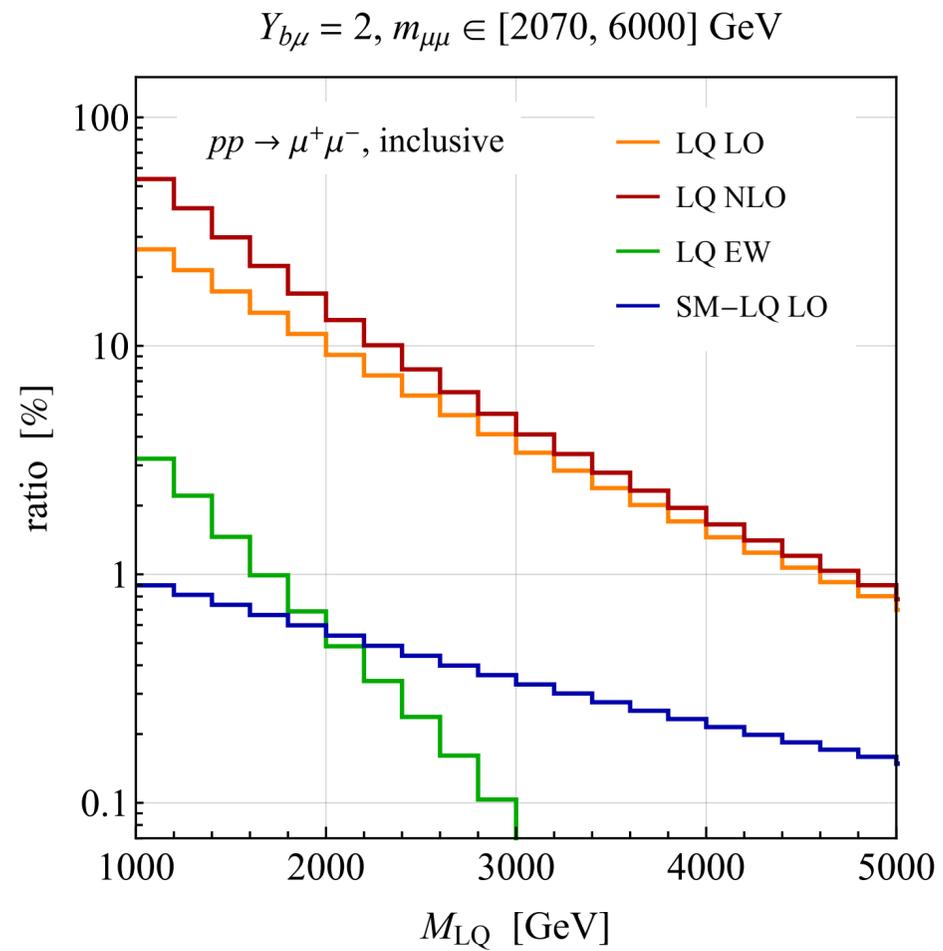
NLO corrections are important for obtaining precise predictions, especially at high $m_{\mu\mu}$.

3. Phenomenology

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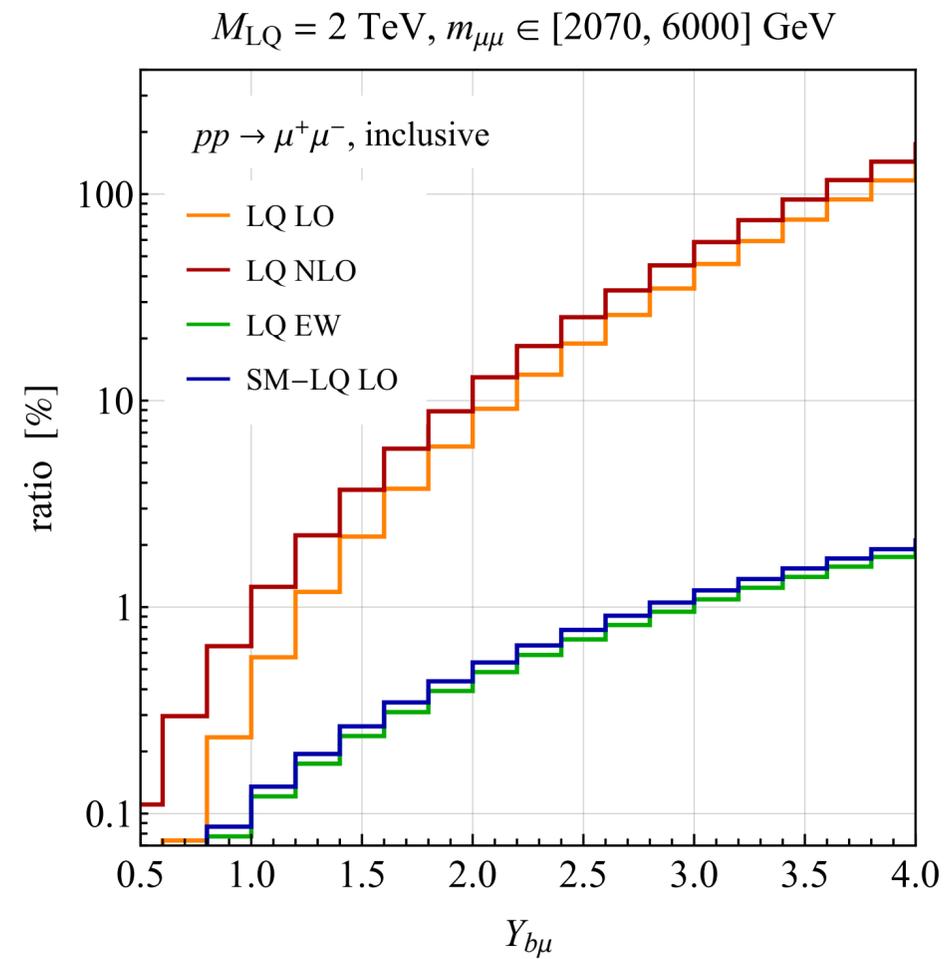
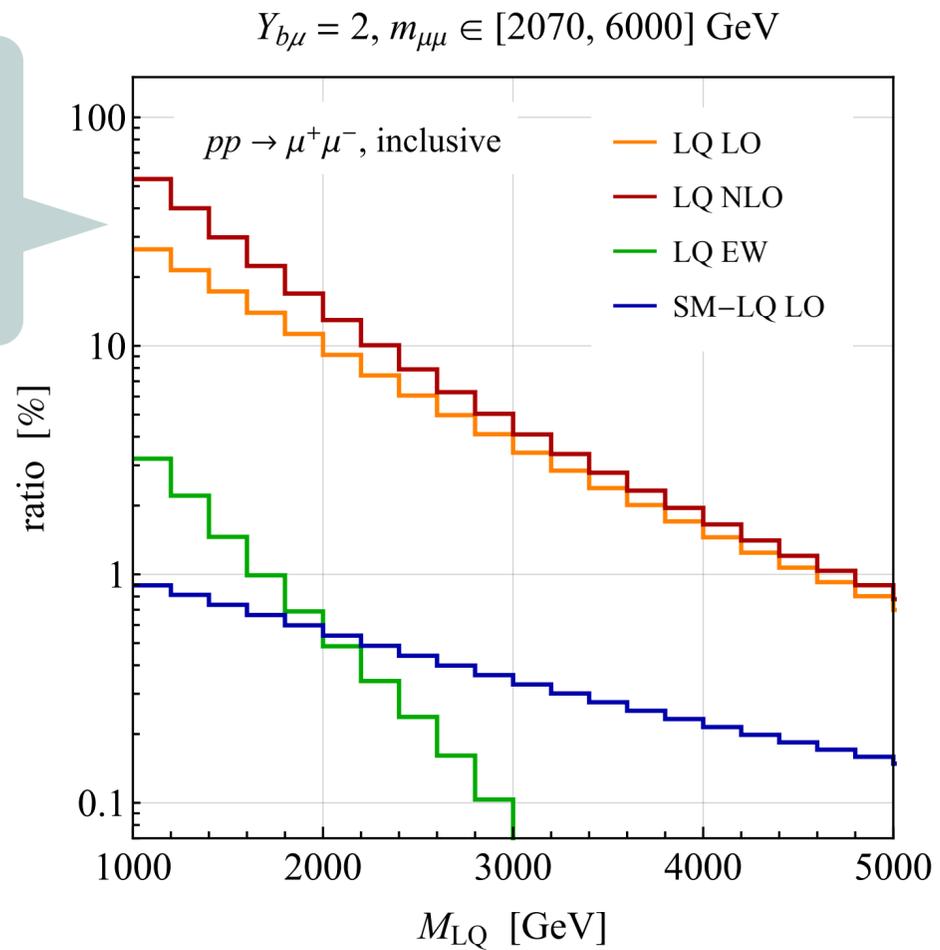
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3. Phenomenology

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NLO corrections decrease as the **single LQ production** channel decouples.

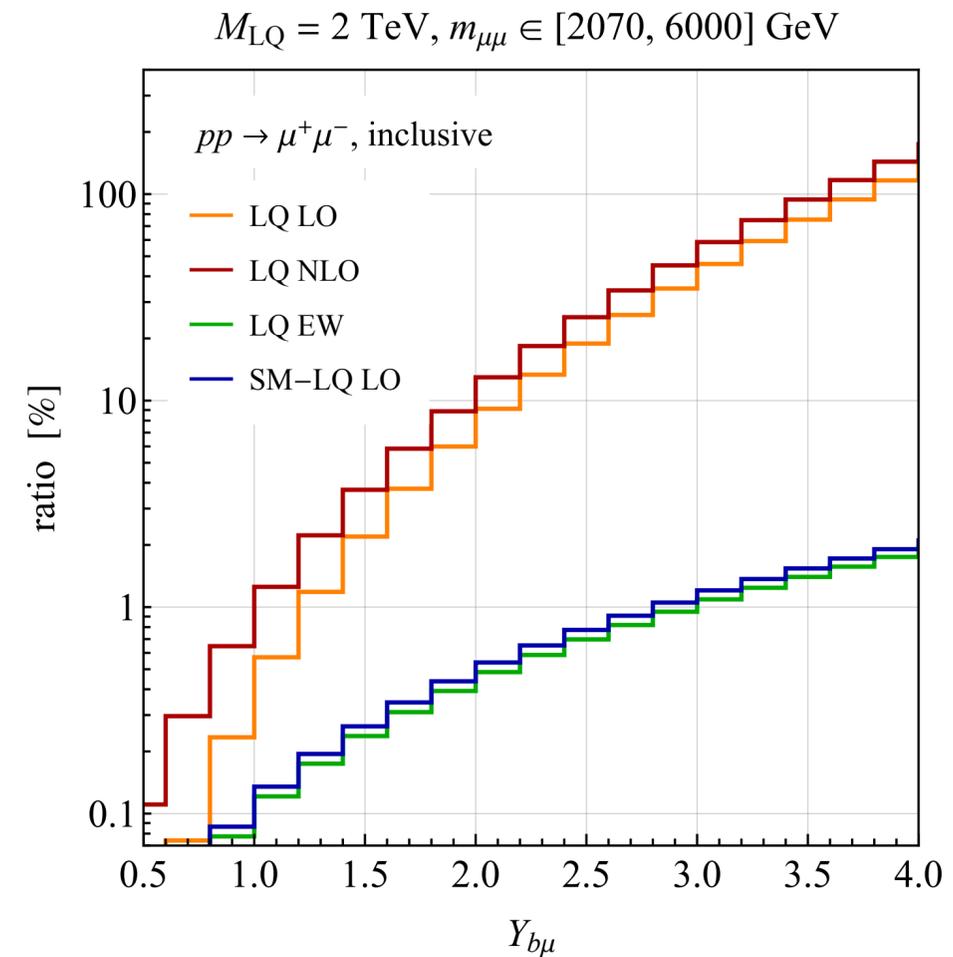
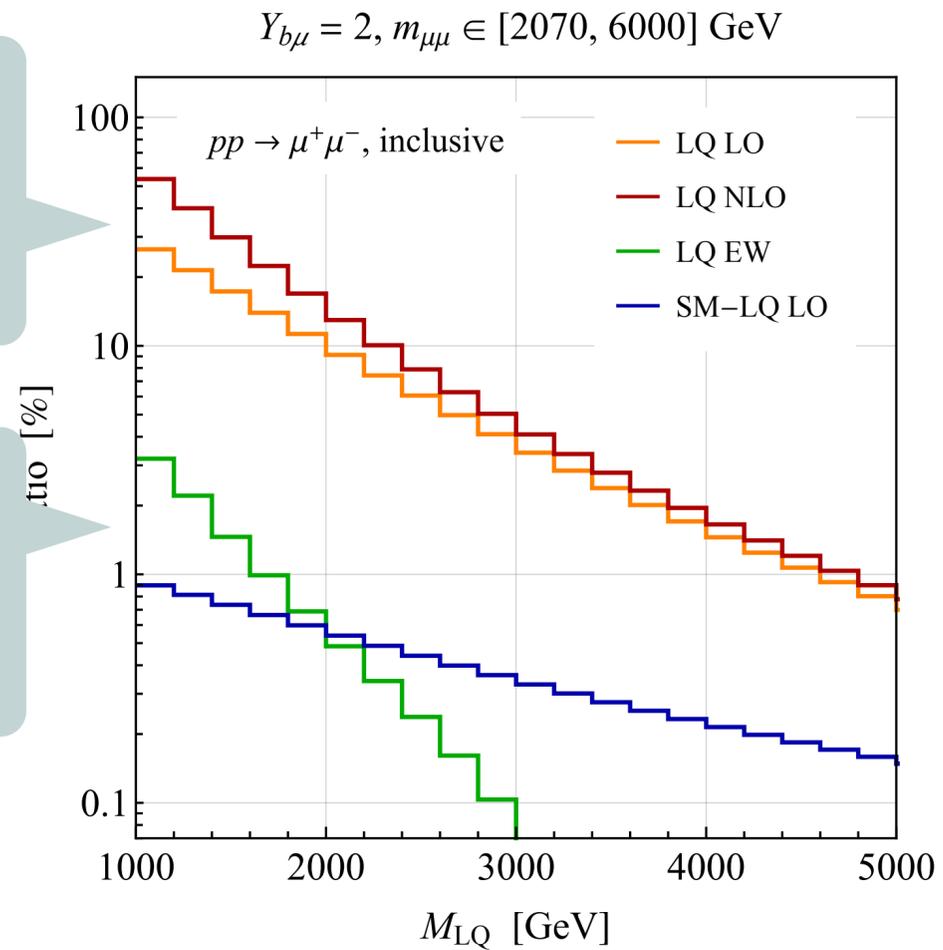


3. Phenomenology

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NLO corrections decrease as the **single LQ production** channel decouples.

Interference and EW effects are not relevant in the signal region considered by ATLAS.

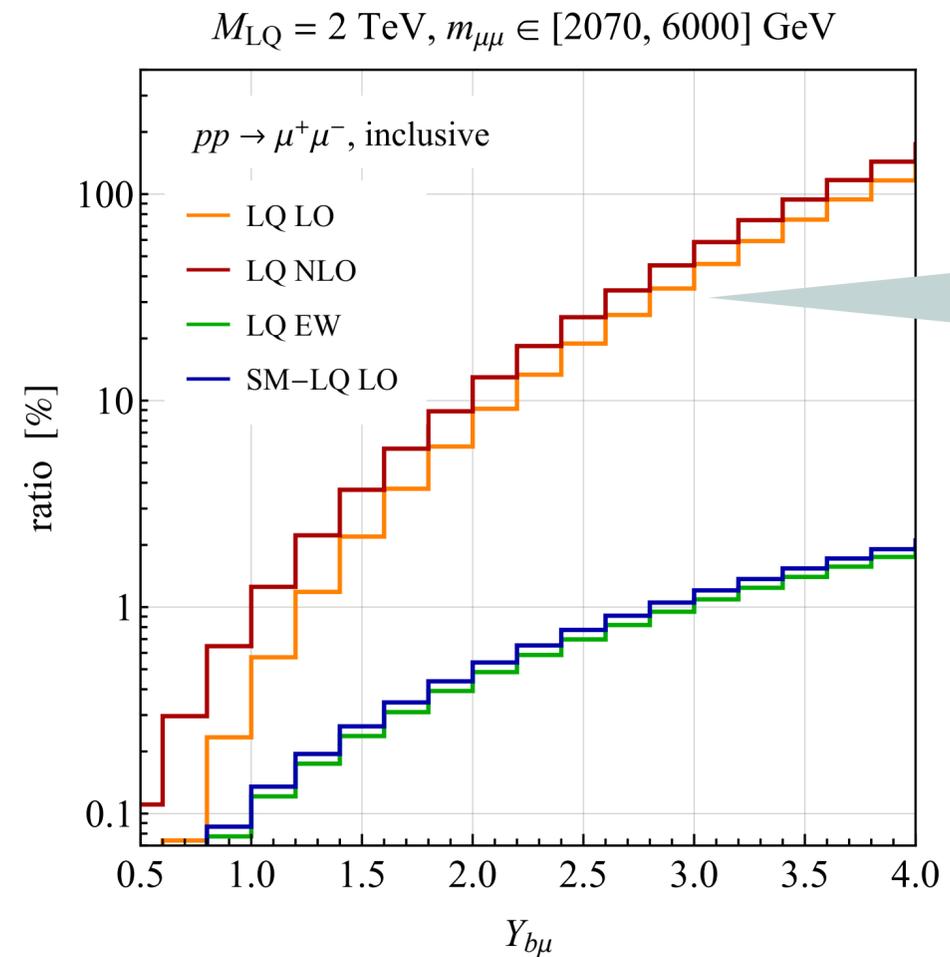
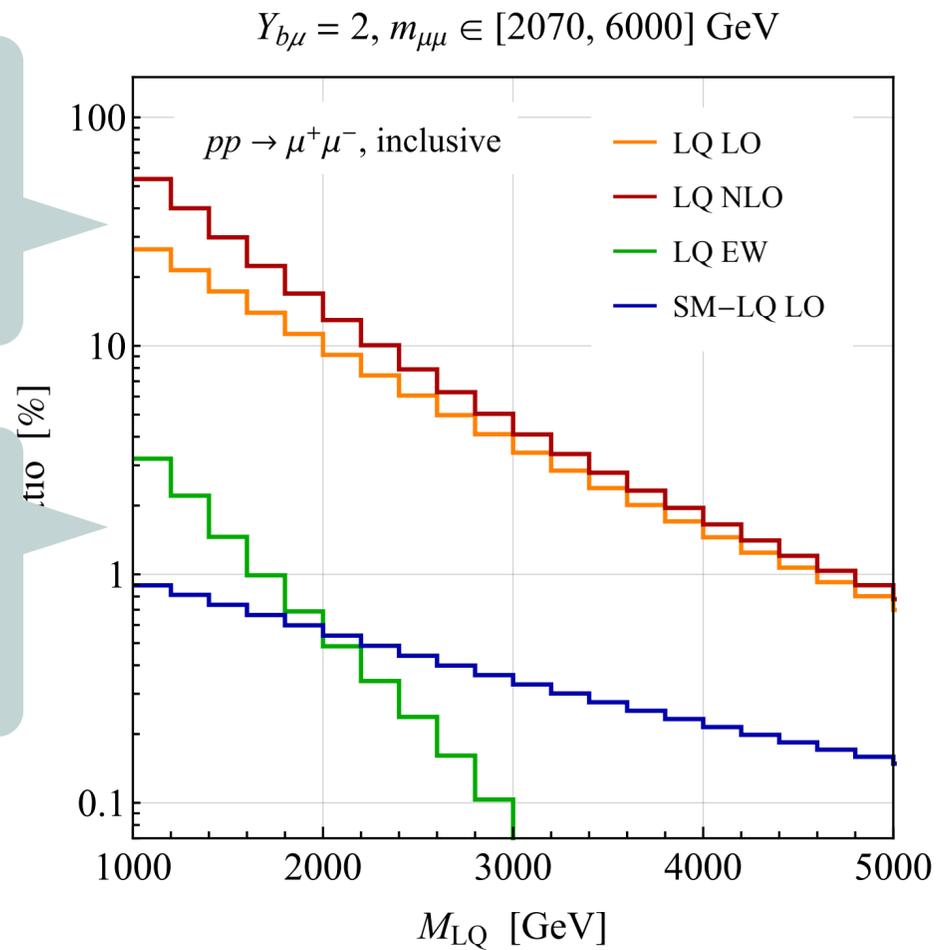


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Single LQ production channel becomes less important for larger $Y_{b\mu}$, since it scales as $|Y_{b\mu}|^2$, whereas the tree-level contribution scales as $|Y_{b\mu}|^4$.

3. Phenomenology

3.2 *b*-tags

3. Phenomenology

3.2 b -tags

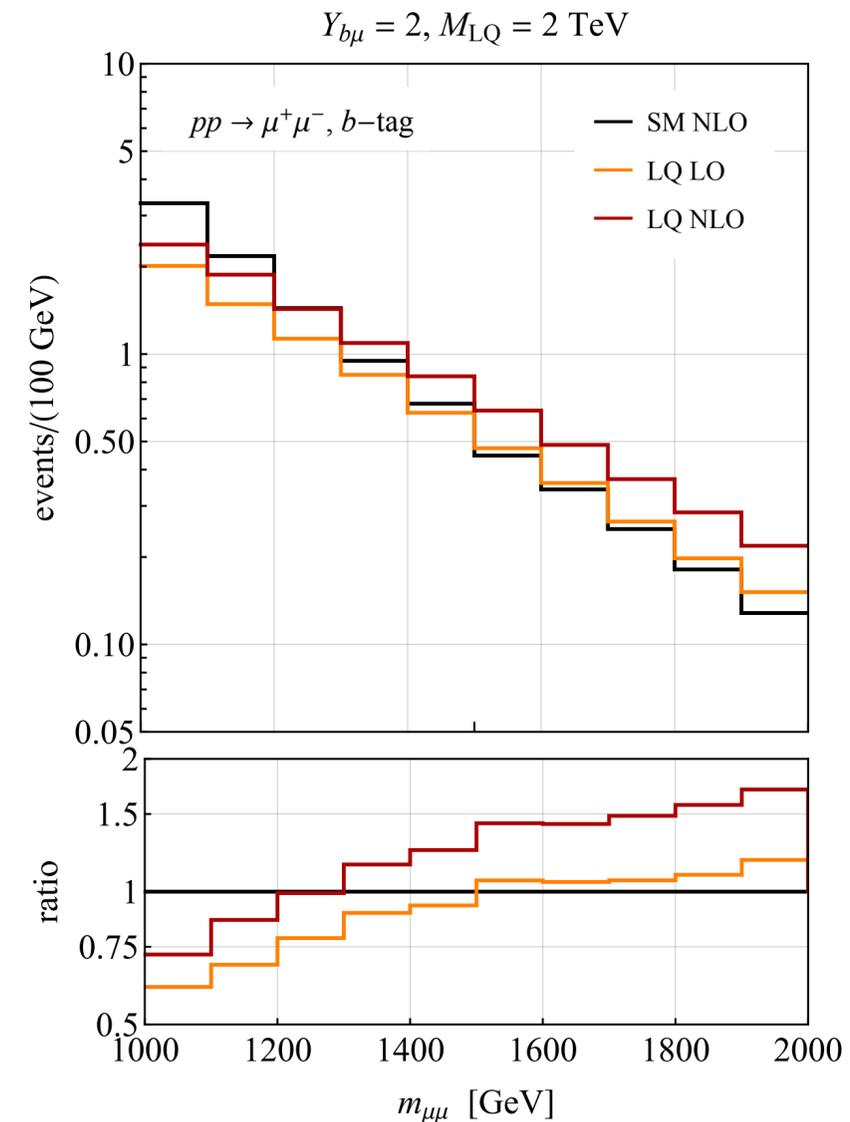
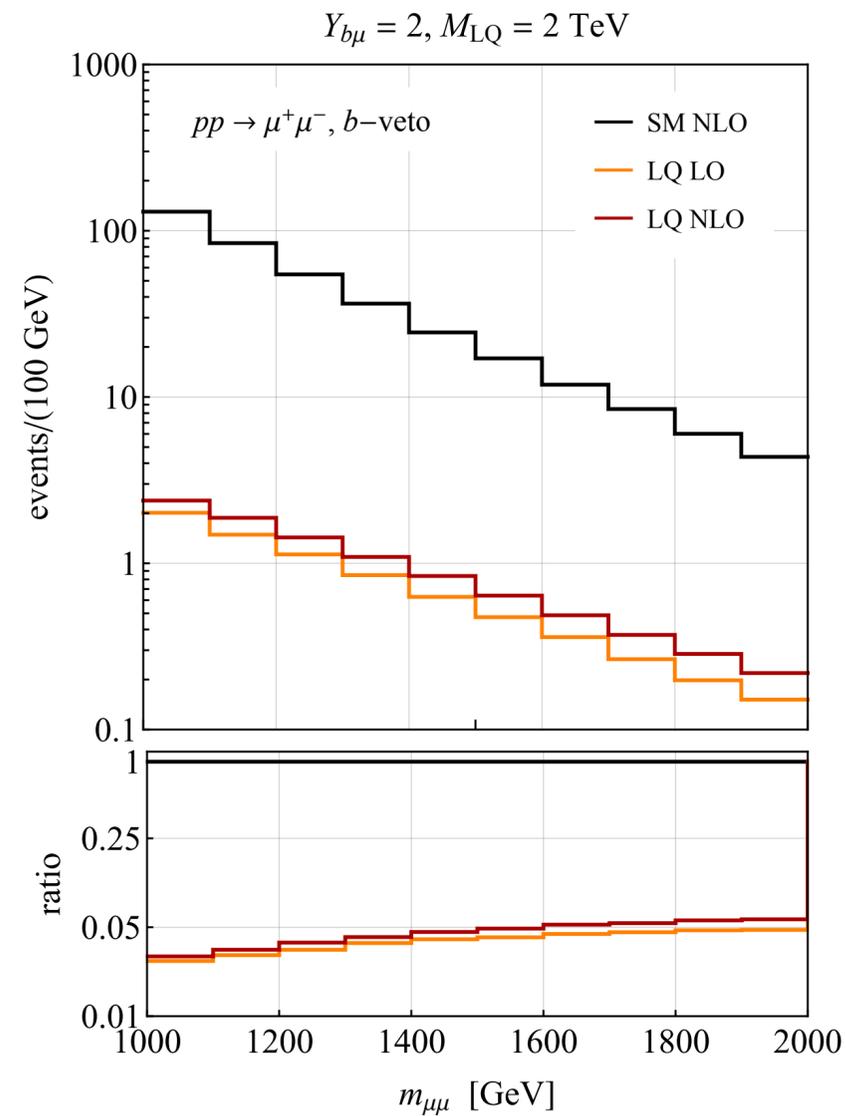
- ATLAS recently presented a **DY analysis for light-leptons** with and without b -tags.

3. Phenomenology

3.2 b -tags

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Recasting their analysis using the full **Madanalysis5 + Delphes** machinery:

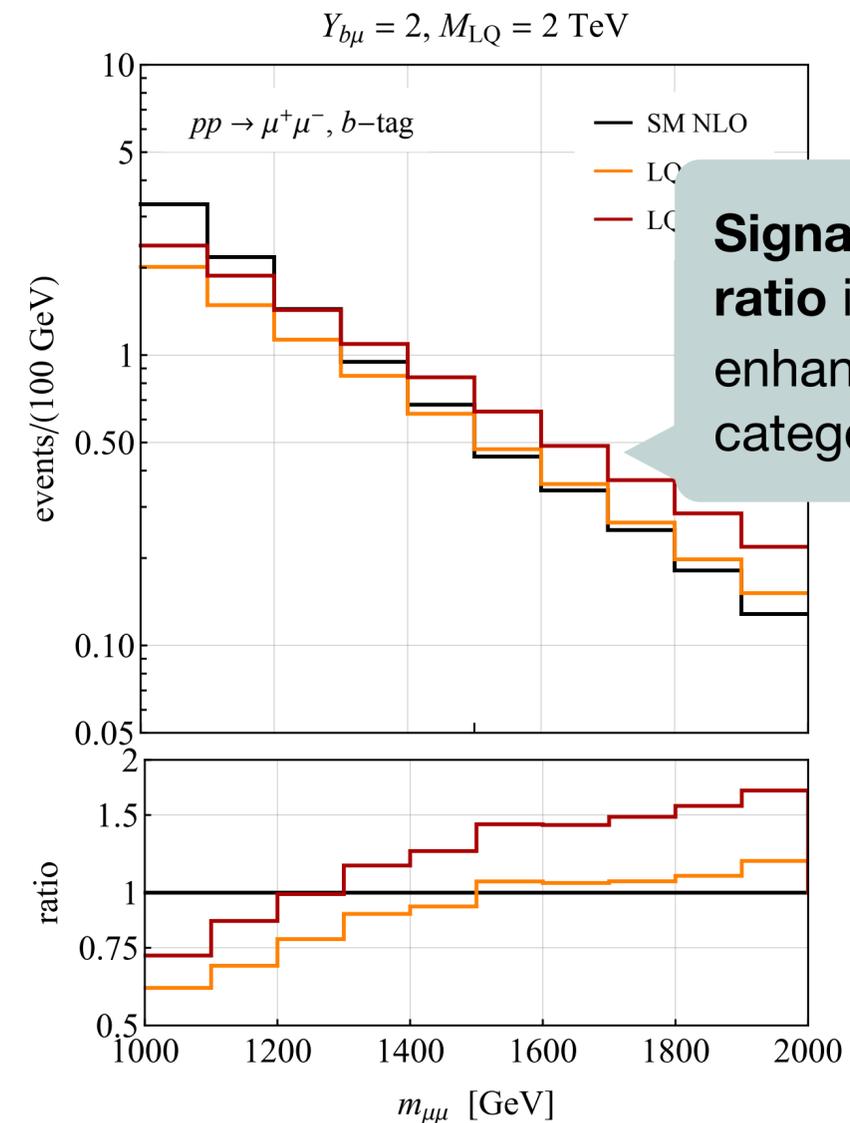
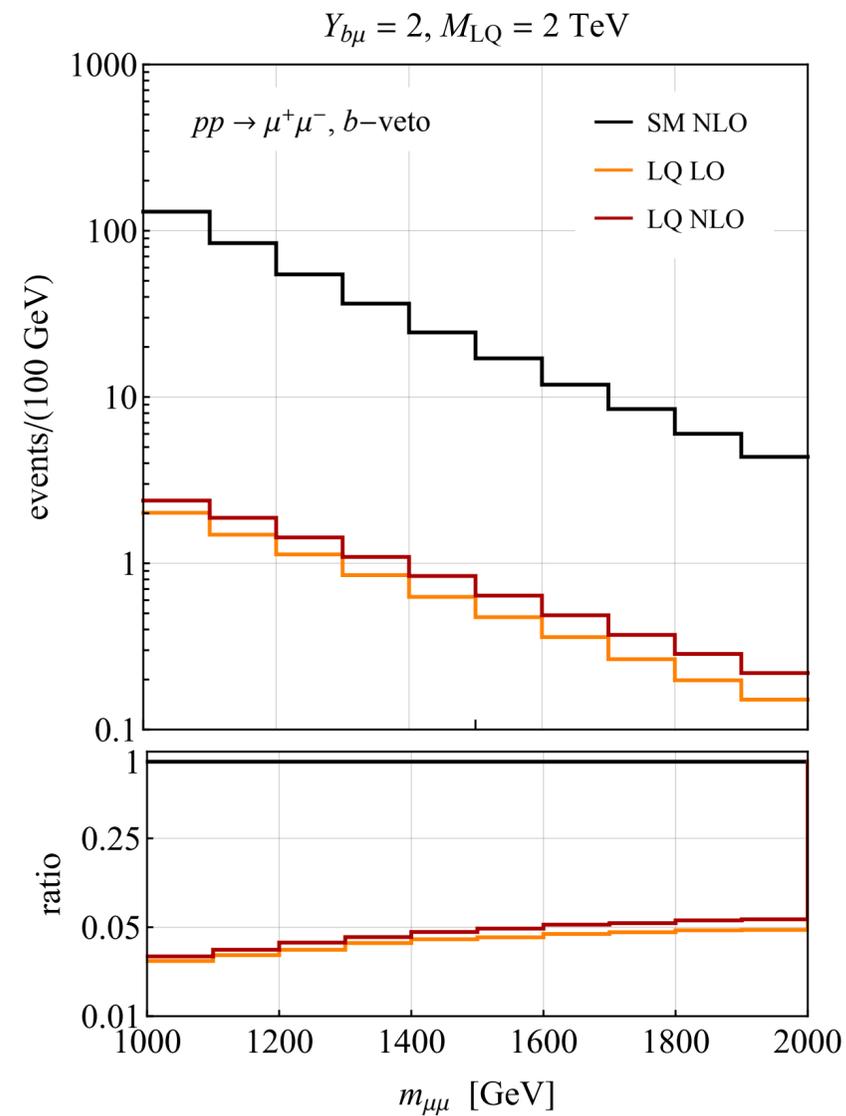


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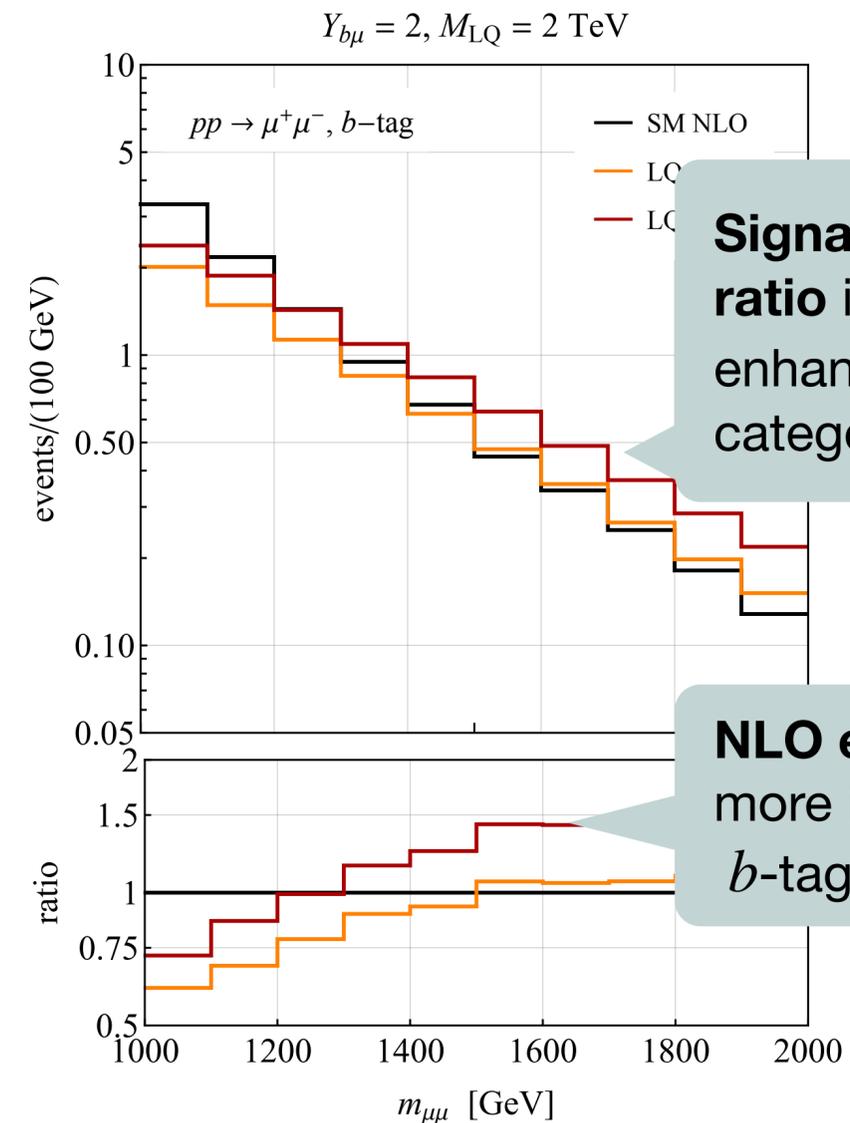
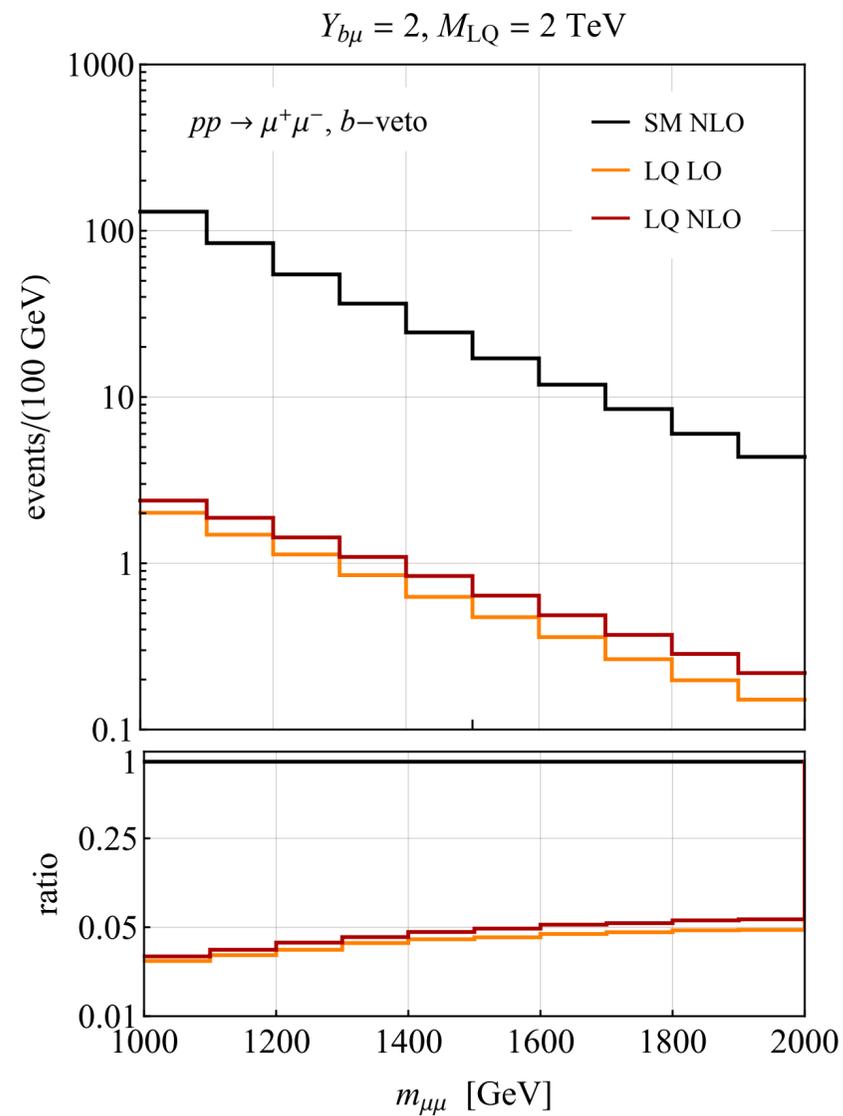
Signal-to-background ratio is greatly enhanced in the b -tag category.

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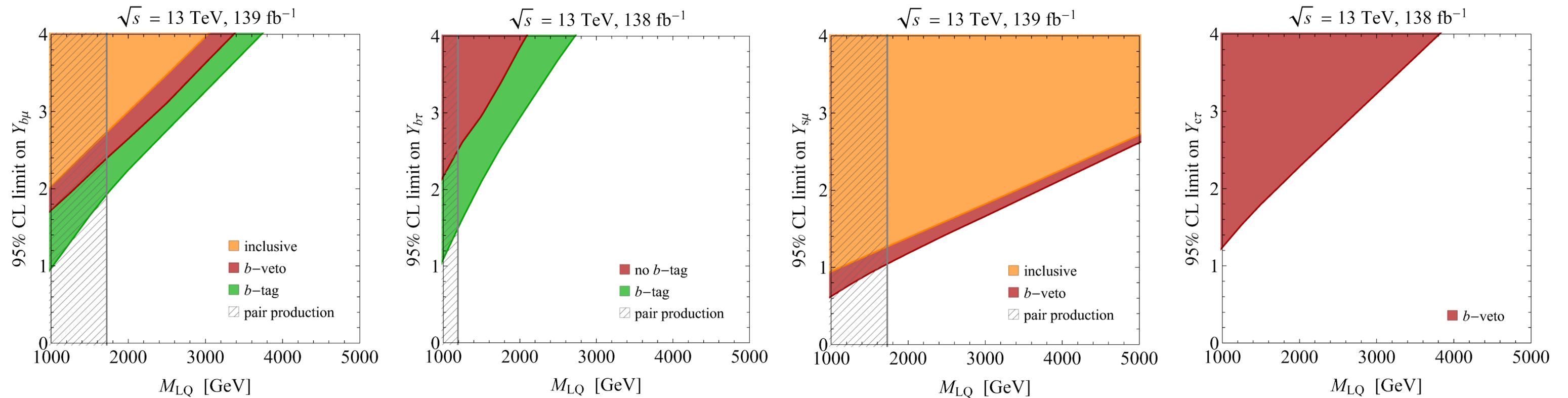
NLO effects are even more important in the b -tag category.

3. Phenomenology

3.3 Exclusion Limits

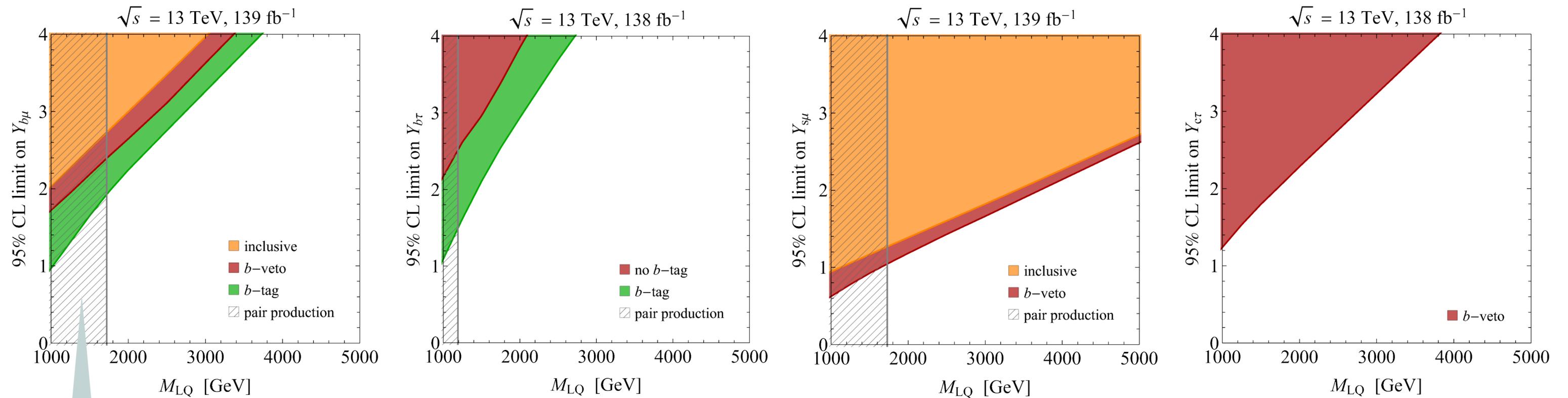
3. Phenomenology

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Improved exclusion limits due to NLO description and b -tagging.

4. Outlook

4. Outlook

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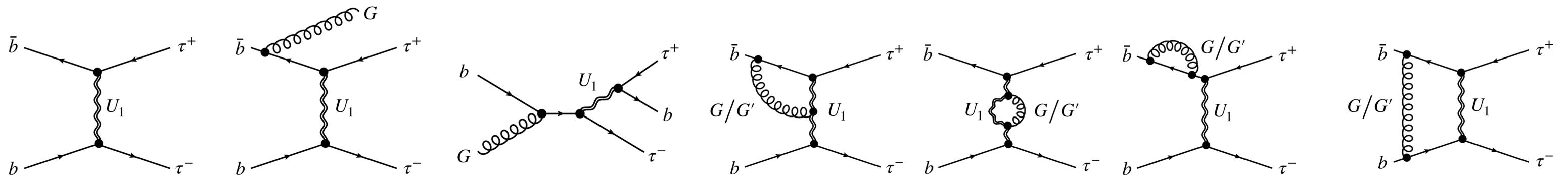
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4. Outlook

- We will make our POWHEG-BOX implementation available on the **POWHEG-BOX web page**.
- It can e.g. also be used to study the effects of LQs on angular observables such as the **forward-backward asymmetry** A_{FB} at full NLO accuracy. Note however that non-zero A_{FB} values can only arise from valence quarks but not from sea quarks at the LHC, which is why we did not include it in our study.

4. Outlook

- We will make our POWHEG-BOX implementation available on the **POWHEG-BOX web page**.
- It can e.g. also be used to study the effects of LQs on angular observables such as the **forward-backward asymmetry** A_{FB} at full NLO accuracy. Note however that non-zero A_{FB} values can only arise from valence quarks but not from sea quarks at the LHC, which is why we did not include it in our study.
- We are working on a full NLO implementation of the U_1 **vector leptoquark** from third family gauge models.



5. Conclusions

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- We implemented scalar LQ effects in DY production at NLO in perturbation theory, providing a dedicated POWHEG-BOX MC implementation that includes
 - The tree-level t-channel LQ exchange
 - All $\mathcal{O}(\alpha_s)$ corrections, including the resonant single-LQ production channel
 - The interference effects with the SM background
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- We derived improved exclusion limits on the LQ masses and Yukawa couplings from these analyses.

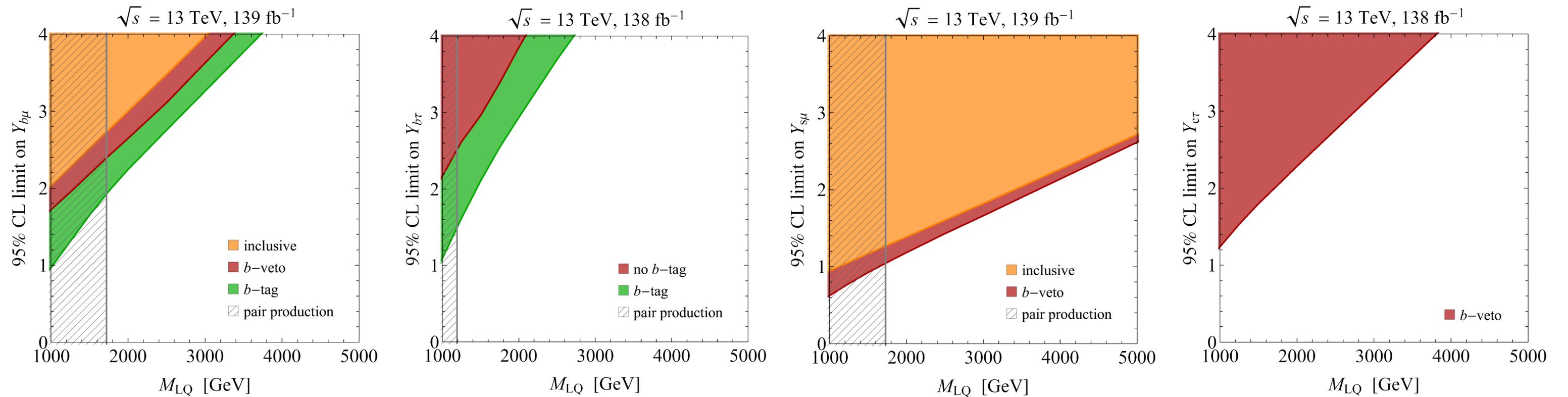
Thank you for your attention!

3. Phenomenology

3.3 Exclusion Limits

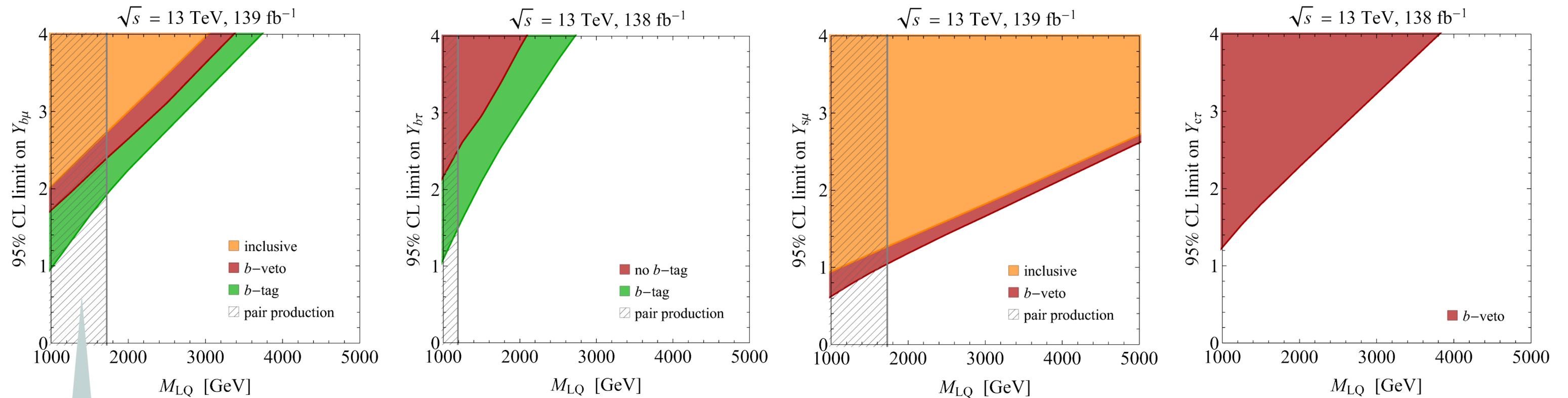
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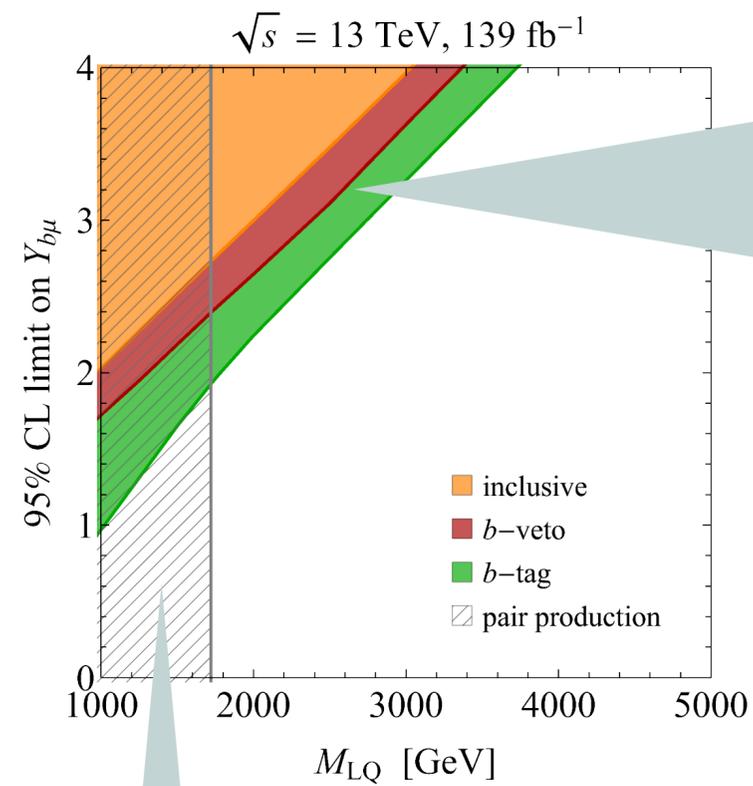
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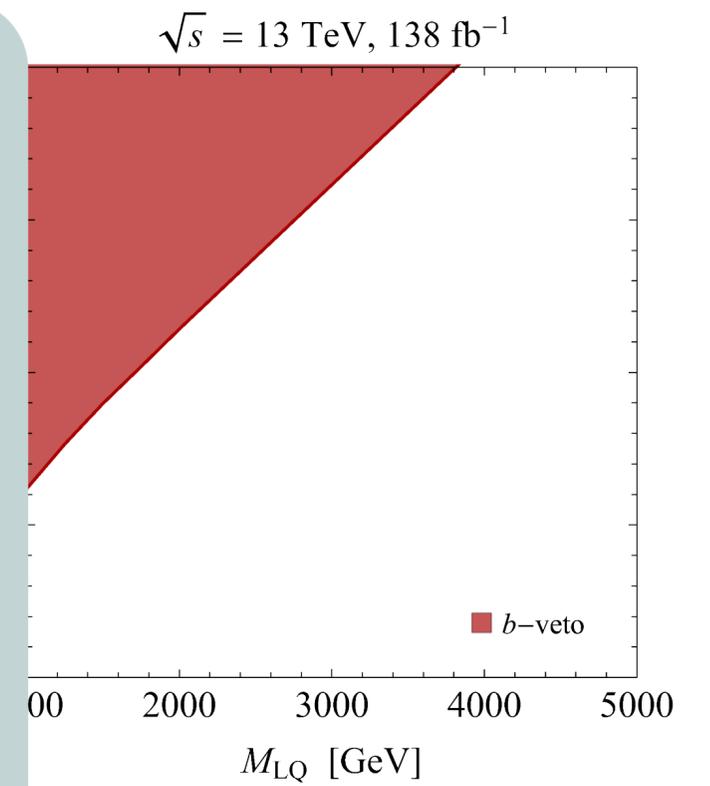
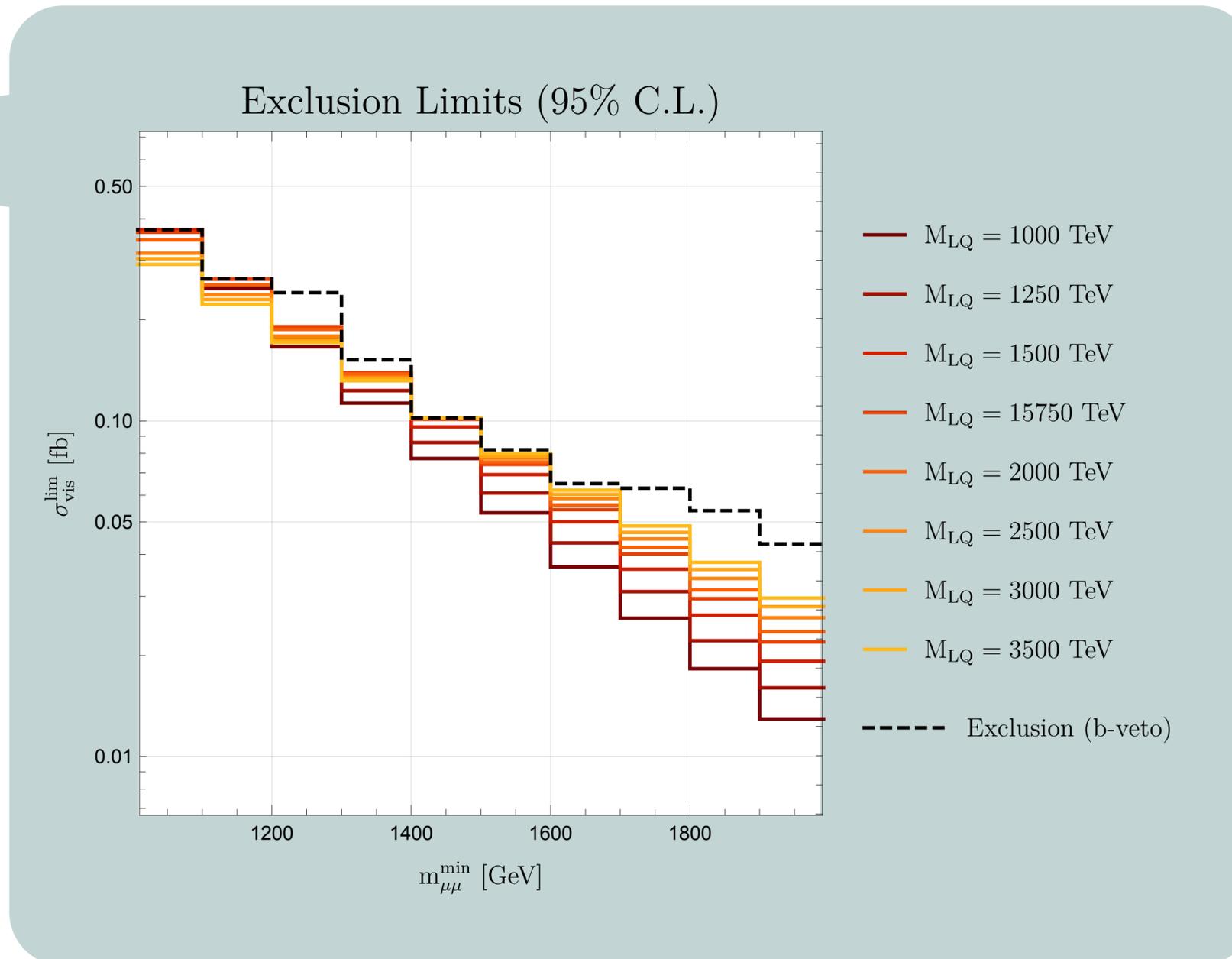
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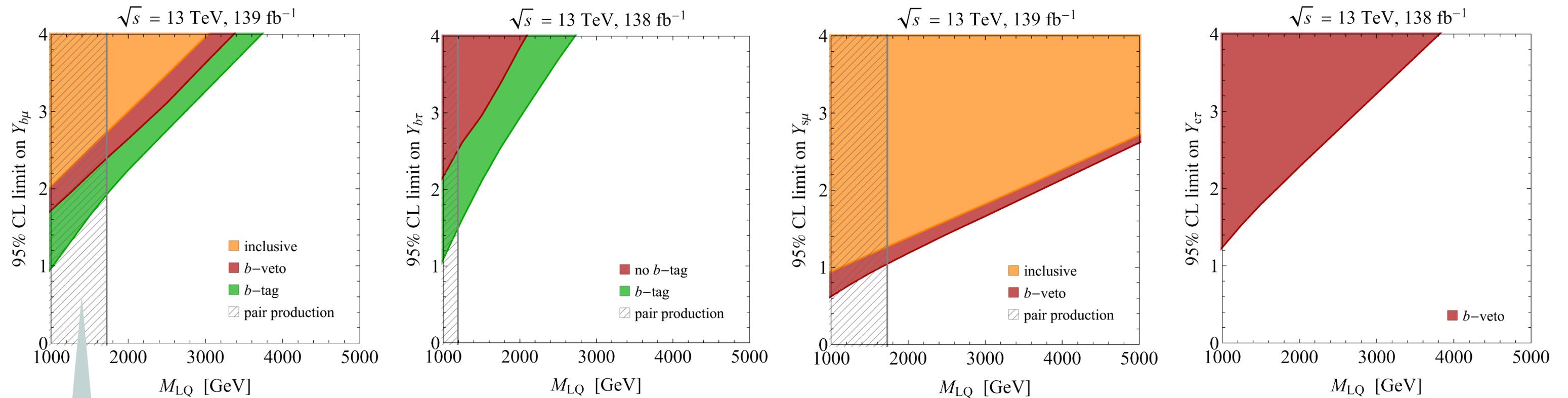


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2. Theoretical framework and calculations

2.3 POWHEG-Box (in a nutshell)

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- The POWHEG framework uses a different strategy, **generating hardest emission events using full NLO accuracy**:

$$\Delta(\Phi_n, p_T) = \exp \left\{ - \int \frac{[d\Phi_{\text{rad}} R(\Phi_{n+1}) \theta(k_T(\Phi_{n+1}) - p_T)]^{\bar{\Phi}_n = \Phi_n}}{B(\Phi_n)} \right\}.$$

$$d\sigma = \bar{B}(\Phi_n) d\Phi_n \left\{ \Delta(\Phi_n, p_T^{\text{min}}) + \Delta(\Phi_n, k_T(\Phi_{n+1})) \frac{R(\Phi_{n+1})}{B(\Phi_n)} d\Phi_{\text{rad}} \right\}_{\bar{\Phi}_n = \Phi_n};$$

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- The **PS** is then only needed for emissions below p_T^{\min} .