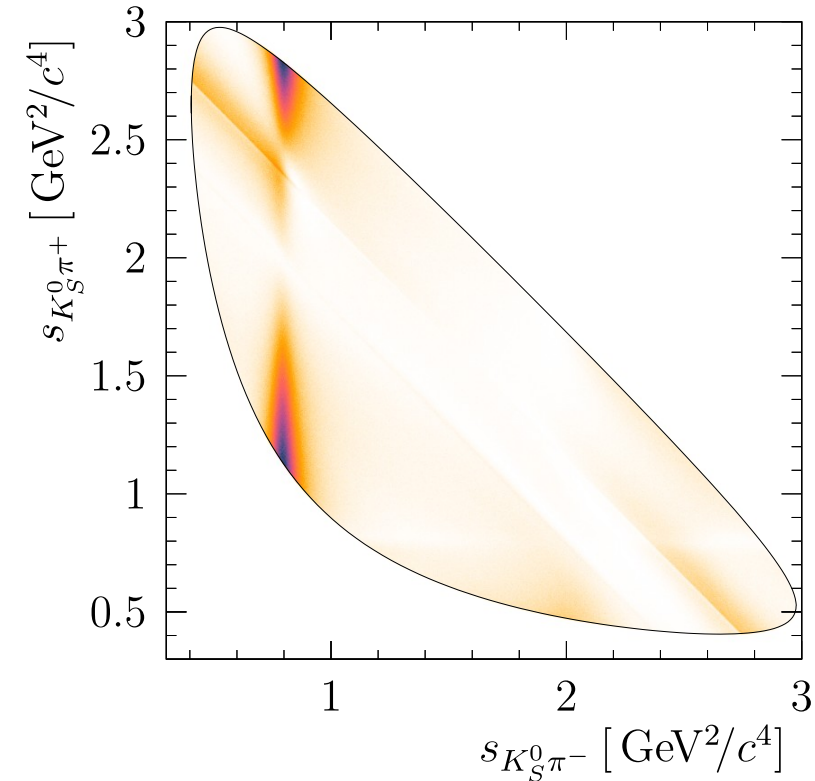


# AN INTRODUCTION TO AMPLITUDE FITTING IN LHCb

T.EVANS, A. MATHAD, M. MIKHASENKO, N. SKIDMORE, Z. XU



$$D^0 \rightarrow K_S^0 \pi^+ \pi^-$$



- Aim is to describe dynamics of multibody processes:
  - Measure CPV and mixing
  - Study intermediate hadronic systems (exotic + conventional)
  - As inputs to precision measurements
  - ...
- A multidimensional problem:
  - $\Rightarrow$  For a spinless  $N$ -body decay:  $d = 3N - 7$
- For a three body decay  $d = 2$  - the famous ‘Dalitz plot’
- Can have fairly enormous data sets (millions of signal candidates)

# AMPLITUDE ANALYSIS: IN A NUTSHELL

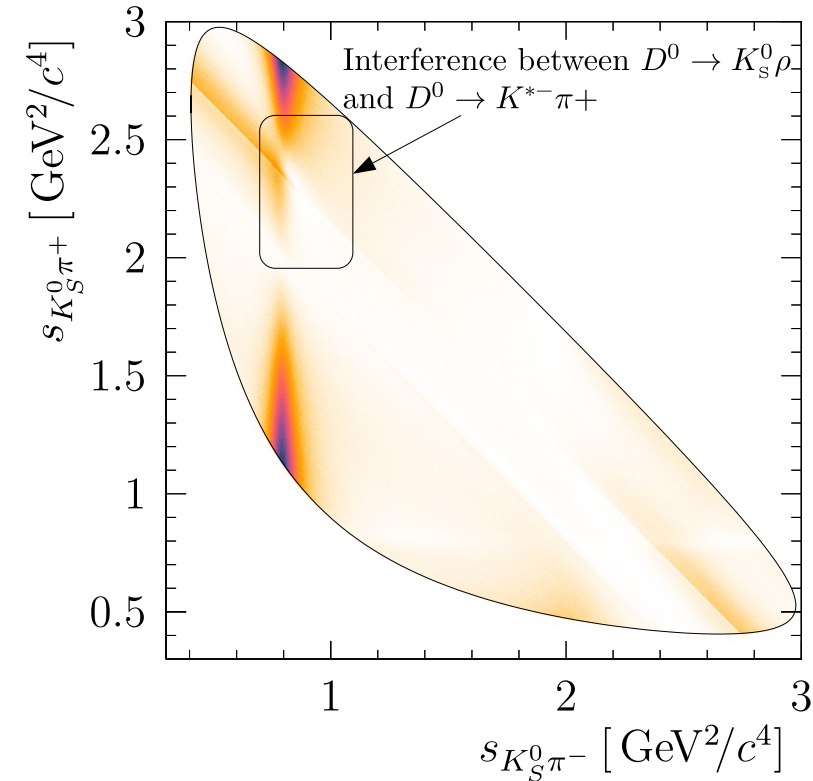
Hadronic systems are strongly coupled quantum mechanical entities  
⇒ Interference and phase behaviour are key

- The observed distribution of signal events for some process  $X \rightarrow Y_1 Y_2 \dots$

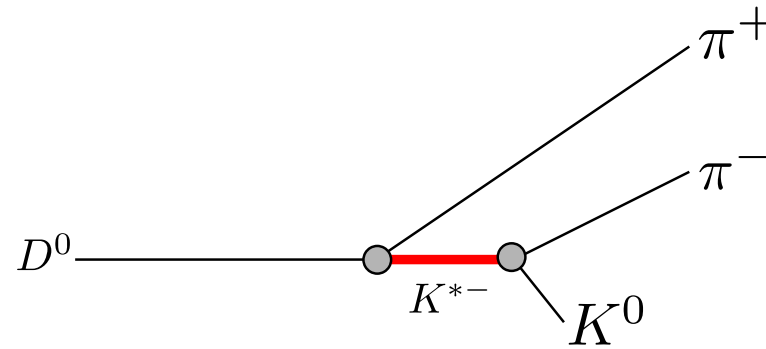
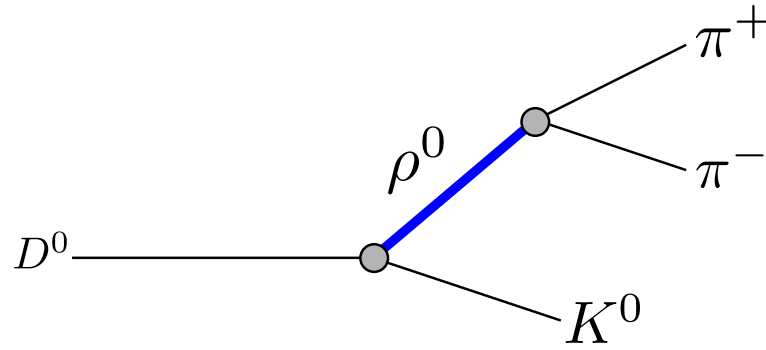
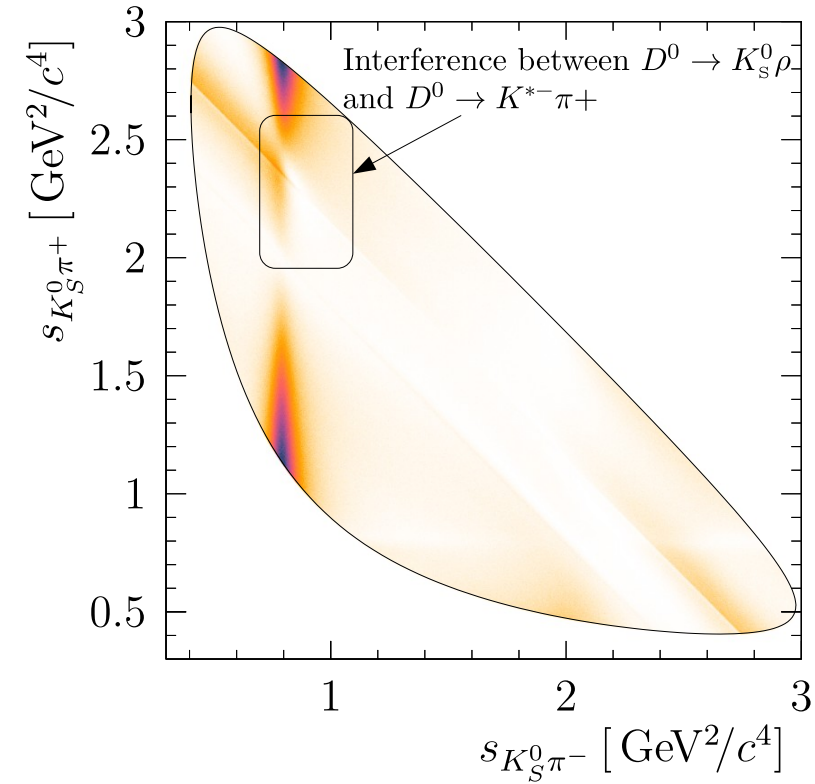
$$\mathcal{P}(\psi) d\psi \propto \rho_N(\psi) \left| \mathcal{M}_{X \rightarrow Y_1 Y_2 \dots} \right|^2 d\psi$$

- Often work within the ‘isobar’ approximation:

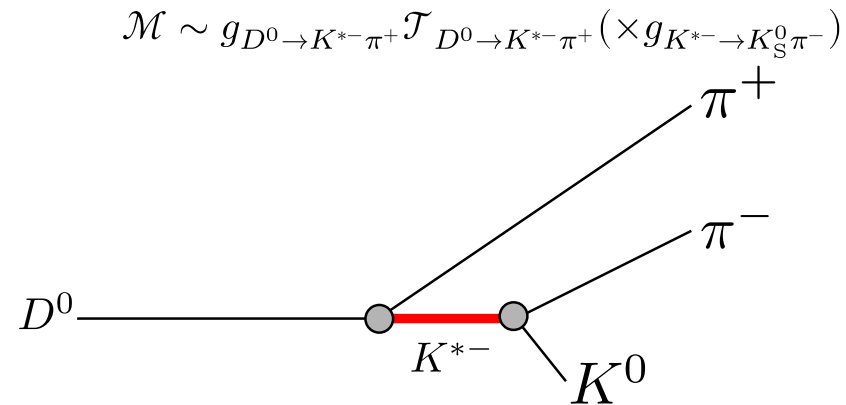
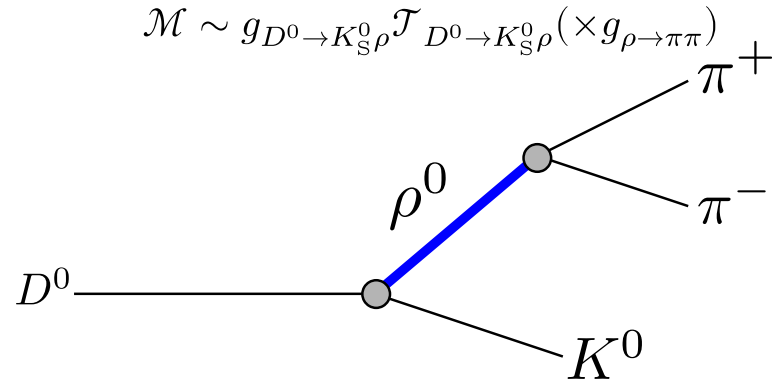
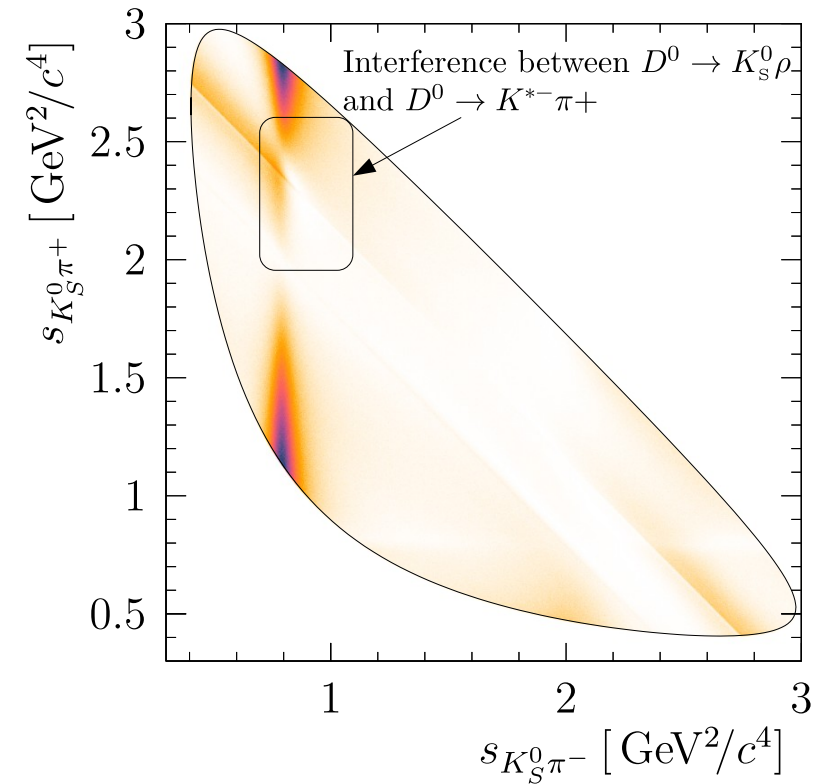
$$\mathcal{M}_{X \rightarrow Y_1 Y_2 \dots} = \sum g_{X \rightarrow R(Y_1 Y_2) \dots} \mathcal{T}_{X \rightarrow R(Y_1 Y_2) \dots}$$



# AMPLITUDE ANALYSIS: IN A NUTSHELL

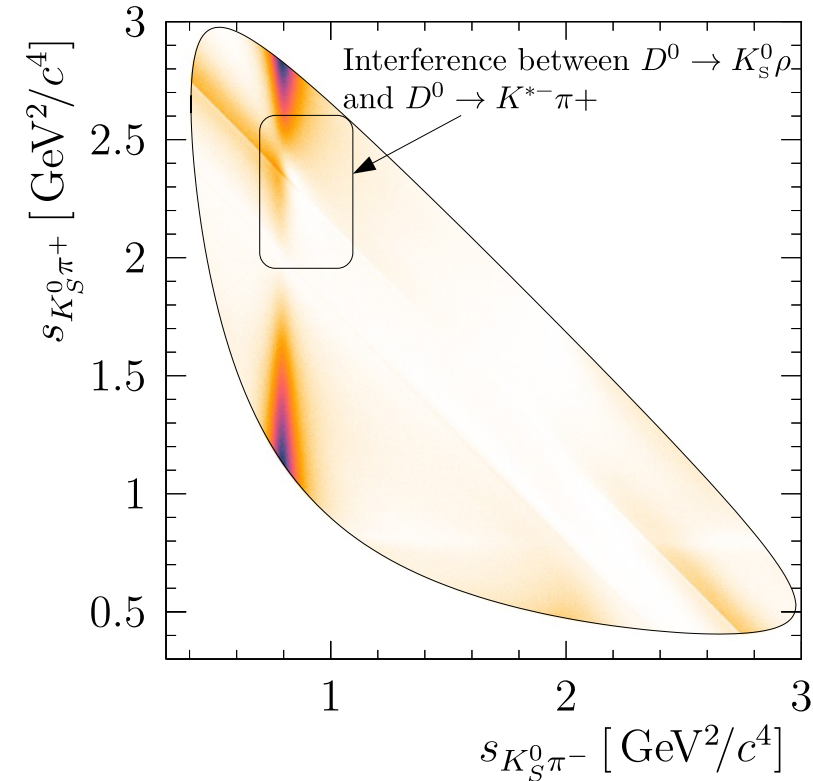


# AMPLITUDE ANALYSIS: IN A NUTSHELL



# AMPLITUDE ANALYSIS: IN A NUTSHELL

Hadronic systems are strongly coupled quantum mechanical entities  
 $\Rightarrow$  Interference and phase behaviour are key



- The observed distribution of signal events for some process  $X \rightarrow Y_1 Y_2 \dots$

$$\mathcal{P}(\psi) d\psi \propto \rho_N(\psi) \left| \mathcal{M}_{X \rightarrow Y_1 Y_2 \dots} \right|^2 d\psi$$

- Often work within the ‘isobar’ approximation:

$$\mathcal{M}_{X \rightarrow Y_1 Y_2 \dots} = \sum g_{X \rightarrow R(Y_1 Y_2) \dots} \mathcal{T}_{X \rightarrow R(Y_1 Y_2) \dots}$$

- (Complex) decay constants  $g_{X \rightarrow R(Y_1 Y_2) \dots}$  are often the primary output of an amplitude analysis
  - Can also measure properties of intermediate states (masses, widths,  $J^P$ , ...)
  - ‘Decay constants’ in most AmAn not really meaningful  $\Rightarrow$  should only be considered a parameter of a model

# A FIRST LOOK AT THE LIKELIHOOD FUNCTION

Multidim fits often doing by maximising likelihood (minimising  $-2LL$  of some observed data set  $\{\psi_i\}$ ):

$$\log \mathcal{L} = \sum_i \log \mathcal{P}(\psi_i)$$

# A FIRST LOOK AT THE LIKELIHOOD FUNCTION

Multidim fits often doing by maximising likelihood (minimising  $-2LL$  of some observed data set  $\{\psi_i\}$ ):

$$\log \mathcal{L} = \sum_i \log (f_{sig} \mathcal{P}_{sig}(\psi_i) + (1 - f_{sig}) \mathcal{P}_{bkg}(\psi_i))$$

- But there is also **background**



# A FIRST LOOK AT THE LIKELIHOOD FUNCTION

Multidim fits often doing by maximising likelihood (minimising  $-2LL$  of some observed data set  $\{\psi_i\}$ ):

$$\log \mathcal{L} = \sum_i \log (\varepsilon(\psi_i)(f_{sig}\mathcal{P}_{sig}(\psi_i) + (1 - f_{sig})\mathcal{P}_{bkg}(\psi_i)))$$

- But there is also **background**
- And the probability to reconstruct decays generally **varies** over  $\{\psi\}$

\* this view is biased: can also deal with these effects using event-weighting.

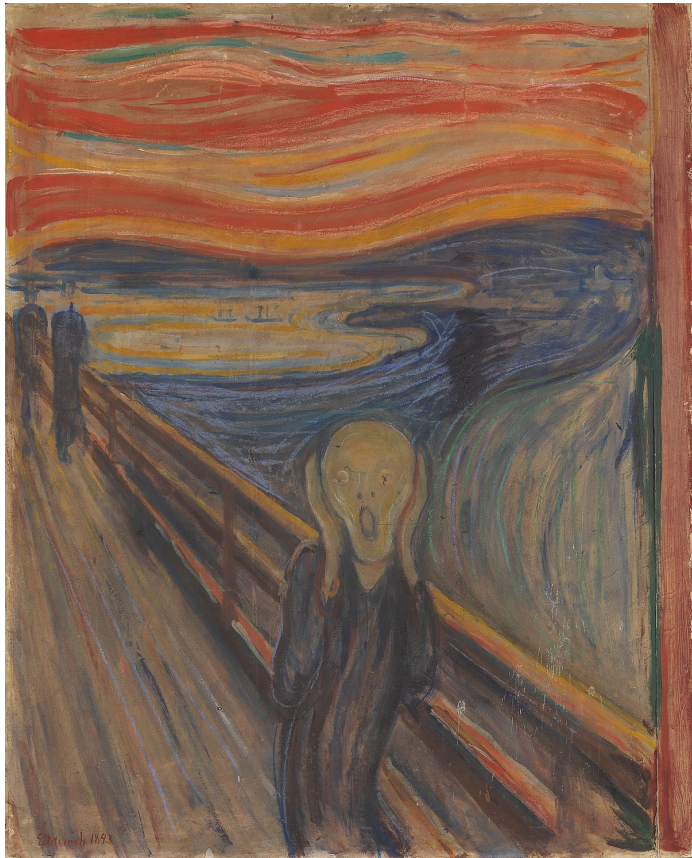
# A FIRST LOOK AT THE LIKELIHOOD FUNCTION

Multidim fits often doing by maximising likelihood (minimising  $-2LL$  of some observed data set  $\{\psi_i\}$ ):

$$\log \mathcal{L} = \sum_i \log \left( \varepsilon(\psi_i) \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi_i) + \frac{(1 - f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi_i) \right) \right)$$

- But there is also **background**
- And the probability to reconstruct decays generally **varies** over  $\{\psi\}$
- We also need to **normalise** our PDFs to use in a likelihood function

# A FIRST LOOK AT THE LIKELIHOOD FUNCTION



Multidim fits often doing by maximising likelihood (minimising  $-2LL$  of some observed data set  $\{\psi_i\}$ ):

$$\log \mathcal{L} = \sum_i \log \left( \int d\psi' \mathcal{G}(\psi', \psi_i) \varepsilon(\psi') \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi') + \frac{(1 - f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi') \right) \right)$$

- But there is also **background**
- And the probability to reconstruct decays generally **varies** over  $\{\psi\}$
- We also need to **normalise** our PDFs to use in a likelihood function.
- What we measure is also convolved with the **detector resolution**.

# IT'S NOT AS BAD AS IT LOOKS

$$\log \mathcal{L} = \sum_i \log \left( \int d\psi' \mathcal{P}(\psi', \psi_i) \varepsilon(\psi') \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi') + \frac{(1-f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi') \right) \right)$$

$$\log \mathcal{L} = \sum_i \log \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi_i) + \frac{(1-f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi_i) \right) + \log \varepsilon(\psi_i)$$

$$\log \mathcal{L} = \sum_i \log \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi_i) + \frac{(1-f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi_i) \right) + \log \varepsilon(\psi_i)$$

# IT'S NOT AS BAD AS IT LOOKS

$$\log \mathcal{L} = \sum_i \log \left( \int d\psi' \mathcal{P}(\psi', \psi_i) \varepsilon(\psi') \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi') + \frac{(1-f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi') \right) \right)$$

$$\log \mathcal{L} = \sum_i \log \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi_i) + \frac{(1-f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi_i) \right) + \log \varepsilon(\psi_i)$$

$$\log \mathcal{L} = \sum_i \log \left( \frac{f_{sig}}{N_{sig}} \mathcal{P}_{sig}(\psi_i) + \frac{(1-f_{sig})}{N_{bkg}} \mathcal{P}_{bkg}(\psi_i) \right) + \log \varepsilon(\psi_i)$$



These are PDFs that we want to use in likelihood fitting: needs normalisation

$$N = \int d\psi \varepsilon(\psi) \mathcal{P}(\psi) = \int d\psi \varepsilon(\psi) \rho_N(\psi) \left| \sum_i g_i \mathcal{T}_i(\psi) \right|^2$$

Everyone (mostly) uses the same trick\* :

$$\int d\psi \rho_N(\psi) \left| \sum_i g_i \mathcal{T}_i(\psi) \right|^2 = \sum_{ij} g_i g_j^* \int d\psi \rho_N(\psi) \mathcal{T}_i(\psi) \mathcal{T}_j(\psi)^*$$

Integral(s) are independent of  $g_{\{i\}}$ : can evaluate once (slow) and reuse each iteration of the fit.

\* slightly more involved for spinful initial / final states but idea is the same

For problems with  $d > 2$ , common to evaluate using MC integration:

$$\mathcal{J} = \int d\psi \rho_N(\psi) f(\psi) \sim \frac{1}{N} \sum_{j=0}^N f(\psi_j) \rho_N(\psi_j)$$

Suppose we sample the points with distribution  $g(\psi)$  w.r.t. phase space rather than uniformly:

$$\mathcal{J} \sim \frac{1}{N} \sum_{j=0}^N \frac{f(\psi_j)}{g(\psi_j)}$$

For problems with  $d > 2$ , common to evaluate using MC integration:

$$\mathcal{J} = \int d\psi \rho_N(\psi) f(\psi) \sim \frac{1}{N} \sum_{j=0}^N f(\psi_j) \rho_N(\psi_j)$$

Suppose we sample the points with distribution  $g(\psi)$  w.r.t. phase space rather than uniformly:

$$\mathcal{J} \sim \frac{1}{N} \sum_{j=0}^N \frac{f(\psi_j)}{g(\psi_j)}$$

Suppose we want to integrate  $\varepsilon(\psi) f(\psi)$  and have a simulated sample distributed as  $\varepsilon(\psi) g(\psi)$  ..

$$\mathcal{J} = \int d\psi \rho_N(\psi) f(\psi) \varepsilon(\psi) \sim \frac{1}{N} \sum_{j=0}^N \frac{f(\psi_j)}{g(\psi_j)}$$



Amplitude analysis might look a bit scary

- Multidimensional fits
- Lots of data
- Complex PDFs
- Complex likelihoods + experimental effects

But lots of tricks have been learnt to make it more manageable

- Aggressive caching
- Wise normalisation strategies
- ML tools for dealing with background

**Lots** of amplitude frameworks written to tackle these problems

- AmpGen
- cFit
- Craft
- GooFit
- Hydra
- Ipanema
- Laura++
- RooFit(+CUDA)
- Mint2
- TensorFlowAnalysis (TFA)
- zFit

+ Standalone code

**Lots** of amplitude frameworks written to tackle these problems

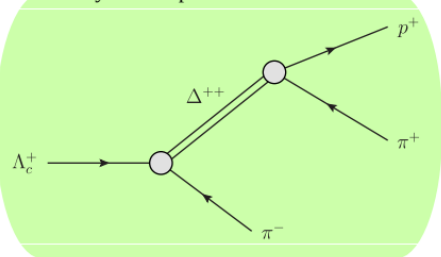
- [AmpGen](#)
- cFit
- Craft
- GooFit
- Hydra
- Ipanema
- Laura++
- [RooFit\(+CUDA\)](#)
- Mint2
- [TensorFlowAnalysis \(TFA\)](#)
- zFit

+ Standalone code

Decay descriptors:

$\Lambda_b(c) + \{\Delta(1232) + \{p^+, \pi^-\}, \pi^-\}$

Decay chain representation



High Level Code

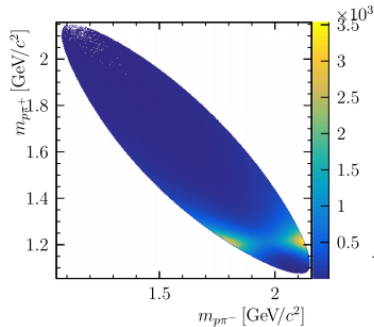
```

DEFINE_LINEShape( BW )
{
  auto s_cse = make_cte(s);
  auto groups = ParticlePropertiesList::get( particleName );
  const Expression mass = Parameter( particleName + "_mass", props->mass() );
  const Expression width = Parameter( particleName + "_width", props->width() );
  const Expression radius = Parameter( particleName + "_radius", props->radius() );
  const Expression q2 = make_cte( Abs( ( s_cse, s1, s2 ) ) );
  const Expression q2b = make_cte( Abs( ( q2 * mass * mass, s1, s2 ) ) );
  Expression FormFactor = sqrt( BlattWeisskopf_Norm( q2 * radius * radius, 0, L ) );
  if ( LineshapeModel == "BL" ) FormFactor = sqrt( BlattWeisskopf( q2 * radius * radius, L ) );
  Expression runTimeWidth = width * s_cse, s1, s2, mass, width, radius, L, dexpressions );
  const Expression BW = FormFactor / ( mass * mass - s_cse * mass * runTimeWidth );
  const Expression kf = kfactor( mass, width, dexpressions );
  ADD_DEBUG( FormFactor, dexpressions );
  ADD_DEBUG( runTimeWidth, dexpressions );
  ADD_DEBUG( BW, dexpressions );
  ADD_DEBUG( kf, dexpressions );
  return kf * BW;
}
    
```

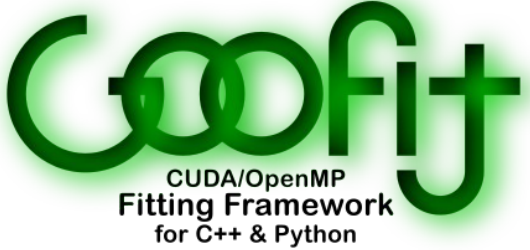
Low Level Code

```

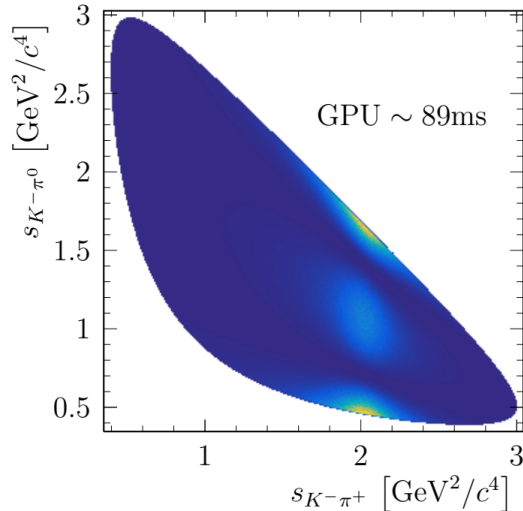
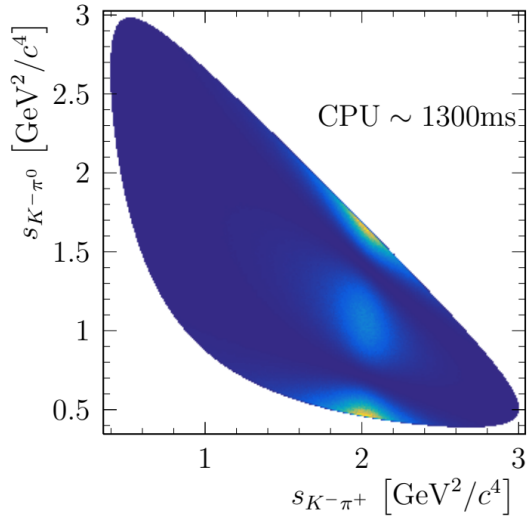
extern "C" std::complex Delta_1232_pp_pi0( double const* v8, double const* v11 )
{
  auto v2083369484 = sqrt( ( 0.02483 + v8[0]*v8[0] + v8[7]*v8[7] + v8[31]*v8[31] );
  auto v2033526506 = sqrt( ( 0.0083543 + v8[4]*v8[4] + v8[5]*v8[5] + v8[12]*v8[12] );
  auto v3330977563 = (-1.)*v8[2] + v8[31]*v8[2] + v8[33] + ( v1835526506 + v2083369484 ) * ( v1835526506 + v2033526506 );
  auto v2083369484 = (-1.)*v8[3] + v8[7]*v8[3] + v8[7] + (-1.)*v8[4] + v8[31]*v8[4] + v8[33];
  auto v2000052626 = std::abs( ( 0.25 ) * ( v3330977563 - ( 1.700709 ) * ( 0.03896 ) + ( ( 0.741105 ) / ( v3330977563 ) ) );
  auto v1843657272 = ( 0.25 ) * ( v3330977563 - ( 1.700709 ) * ( 0.03896 ) + ( ( 0.741105 ) / ( v3330977563 ) );
  auto v202454712 = sqrt( ( 0.009316 + v8[0]*v8[0] + v8[1]*v8[1] + v8[31]*v8[31] );
  sqrt( v8[0]*v8[0] + v8[1]*v8[1] + v8[31]*v8[31] );
  v2000052626 = sqrt( ( 0.009316 + v8[0]*v8[0] + v8[1]*v8[1] + v8[31]*v8[31] );
  std::abs( ( 0.25 ) * ( v8[0]*v8[0] - ( 1.700709 ) * ( 0.03896 ) + ( ( 0.741105 ) / ( v3330977563 ) ) );
  v2000052626 = sqrt( ( 0.009316 + v8[0]*v8[0] + v8[1]*v8[1] + v8[31]*v8[31] );
  std::abs( ( 0.25 ) * ( v8[0]*v8[0] - ( 1.700709 ) * ( 0.03896 ) + ( ( 0.741105 ) / ( v3330977563 ) ) );
  return v202454712;
}
    
```



- Started as a divergent branch of MINT, redesigned to deal with charm data sets (~ 10<sup>6</sup> signal candidates)
- Originally used for four-body decays, but now much more
  - Spin 1/2, 1 initial/final state(s) (including multifermion final states)
  - Indeterminate initial states (quantum-correlated decays, mixing, ... )
  - Supports both canonical and covariant spin formulations.
  - Very flexible support for different propagators (Breit-Wigners, multi-body running widths, K-matrices, ‘Dalitz’ models i.e.  $\eta^0 \rightarrow \pi^+ \pi^- \pi^0, \dots$  )
  - Can export models for publication + integration with other frameworks (*i.e.* EvtGen)

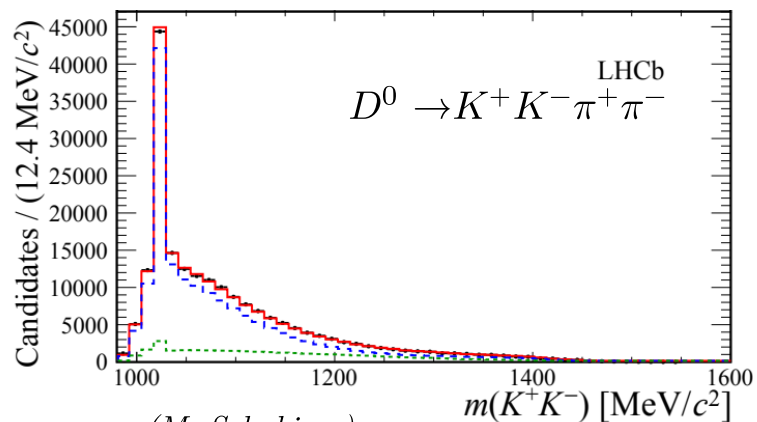


Part of the GooFit organisation  
<https://github.com/goofit/ampgen>

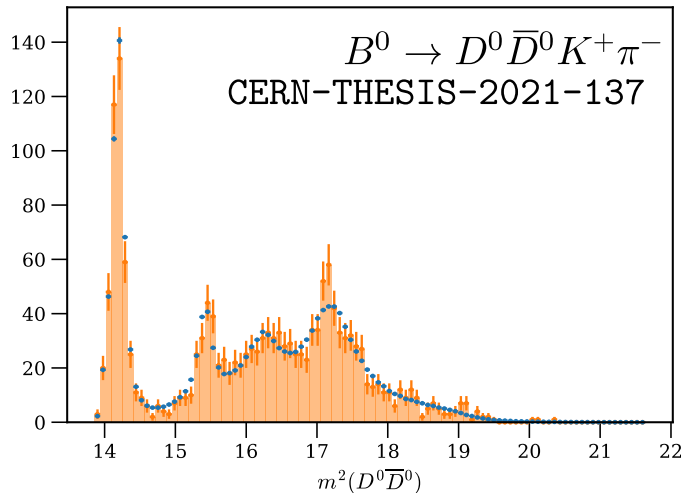


- Default target architecture is AVX2 + OpenMP
- Can also target other architectures:
  - Scalar, AVX512 ... (latter is untested)
  - Can target GPUs but not integrated into the project.
- This (+ design) means AmpGen is very fast
  - From  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  example, takes  $\sim 3$ s to do FCN evaluation ‘from scratch’ with 1 million signal + 2.5 million integration events.  $\implies$  about 1 minute to do the fit (20 amplitudes, 38 free parameters)
- Example compares CPU (scalar, single core) vs GPU for  $D^0 \rightarrow K^{*0} \pi^0$

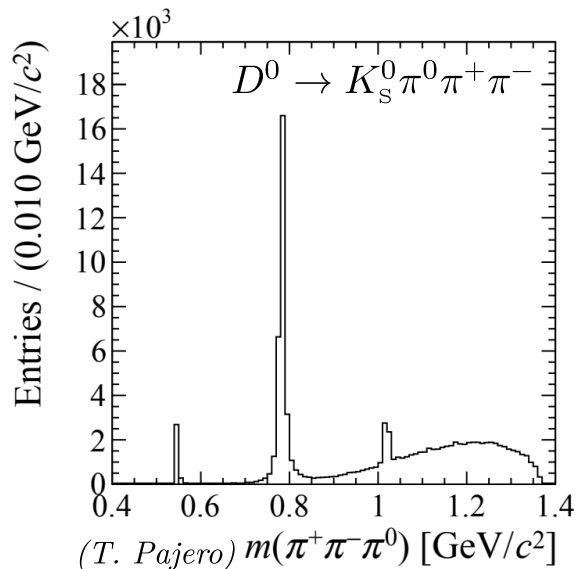
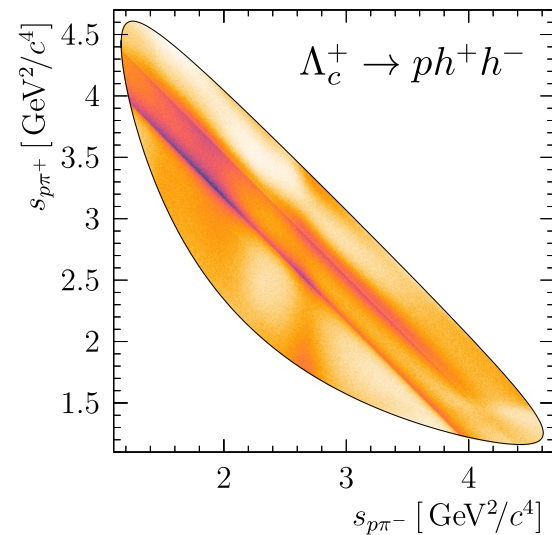
# SOME EXAMPLE ANALYSES



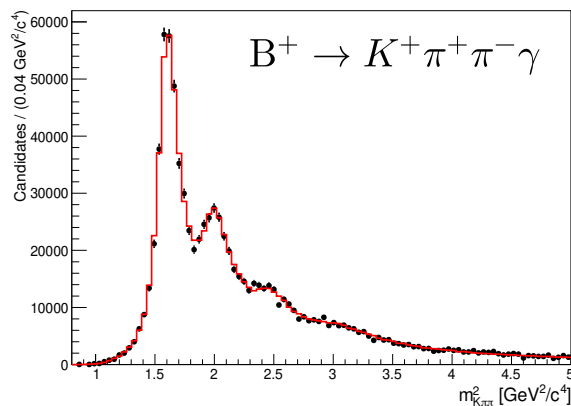
(M. Schubiger)



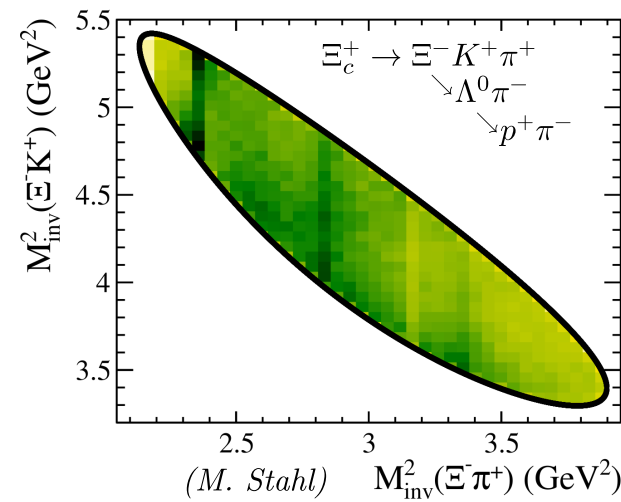
(S. Bhasin, J. Amey, M. Whitehead)



(T. Pajero)



(V. Bellee, M. Diaz, P. Pais)



(M. Stahl)

# UNIT TESTS AND PUBLICATION OF MODELS

It is **infamously** difficult to reproduce the results of amplitude analyses with code/framework other than whatever was used for the paper

- What we call a ‘decay constant’ vs what is encoded by the dynamic functions is *almost* entirely arbitrary

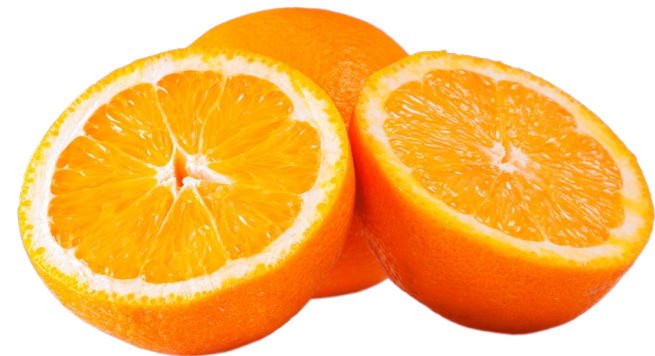
$$\mathcal{M} \sim g_{D^0 \rightarrow K_S^0 \rho} \mathcal{T}_{D^0 \rightarrow K_S^0 \rho} (\times g_{\rho \rightarrow \pi\pi})$$

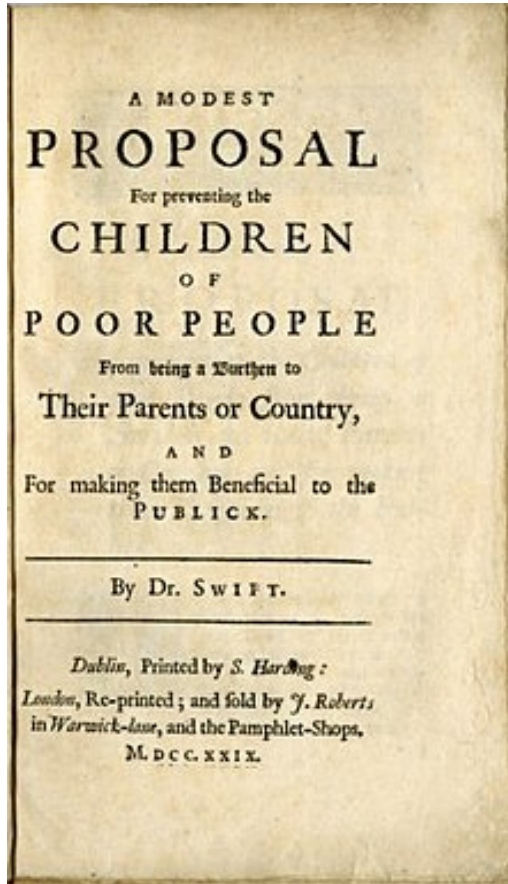
- (Almost) no-one agrees on a common set of conventions: tends to be reasonable choices for each analysis / framework
- Also makes it virtually impossible to have sensible ‘unit tests’ that are supposed to agree between frameworks

‘Solved’ problem for Dalitz models: publish high resolution histograms of amplitude / phase of models

⇒ Generally doesn’t work for  $> 2d$  problems due to curse-of-dimensionality

⇒  $> 2d$  requires some mechanism to publish code (preferably relatively independent of frameworks)





Can compare results from a few key decays (no specific proposals here)  
⇒ Less about comparing results / performance and more about showing some common examples for those starting an analysis to be able to compare different options

- Time to do fits is relevant *if* it becomes a bottleneck (or we start limiting what we do such that can be done in reasonable compute time)
- Readability
- Length of code
- Something else?

Basis can be a set of signal-only simulation samples different frameworks can try and fit in their own way

⇒ Critical that these are generated with whatever framework/code was used to write the model in the first place..