

ESR

(Experimental Storage Ring)

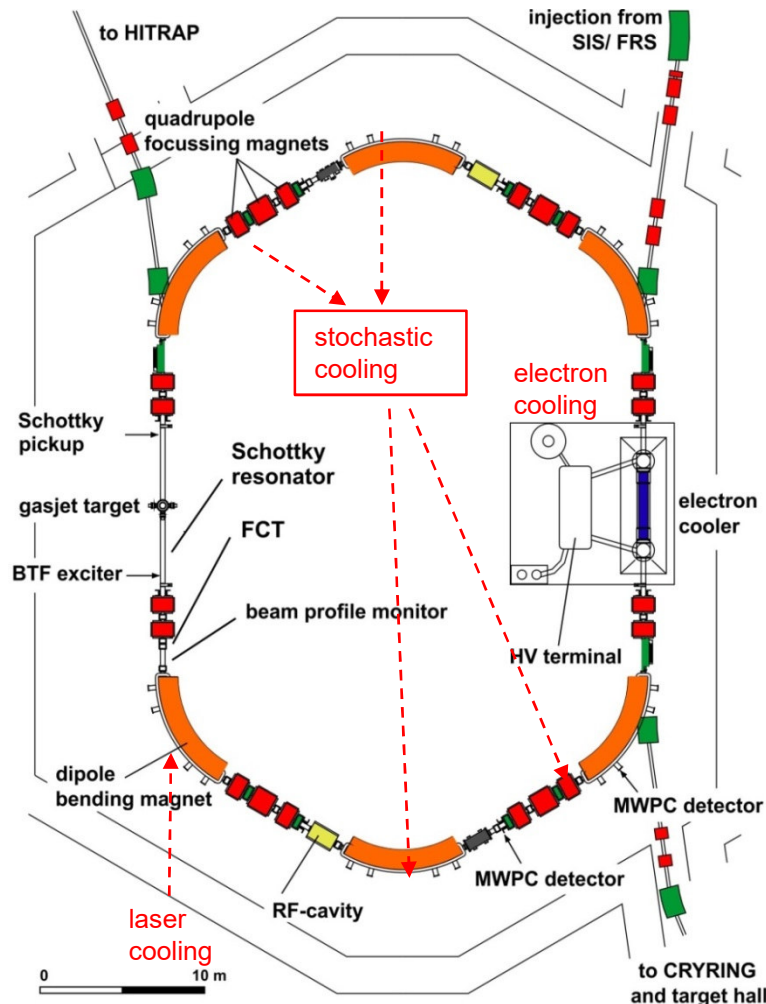
an Extremely versatile Storage Ring

M. Steck

Storage Rings

GSI Accelerator Operations

- ion species** p ,..., Li^{1+} , $\text{C}^{3+,6+}$, ..., $\text{Xe}^{52+,53+,54+}$, ..., $\text{U}^{89+,90+,91+,92+}$, RIBs
- energy** 3 400(550) MeV/u (for U^{92+})
- acceptance (design)** $A_h \sim 400 \mu\text{m}$, $A_v \sim 150 \mu\text{m}$, $\Delta p/p \sim \pm 3 \%$
- intensity** single ions some 10^9 cooled ions
- beam emittance** $10^{-7} \dots 10 \mu\text{m}$
- beam momentum spread** $10^{-7} \dots 10^{-3}$
- lifetime** milliseconds to hours (depending on energy, species)



circumference 108.36 m
bending power 10 Tm

Storage of Highly Charged Ions (HCI) and Rare Isotope Beams (RIB) after fast injection e.g. U^{92+} : in the energy range 3 - 550 MeV/u

Stochastic cooling (≥ 400 MeV/u)

Electron cooling (3 - 430 MeV/u)

Laser cooling (C^{3+} 120 MeV/u)

Beam accumulation

Deceleration (down to 3 MeV/u)

Fast extraction (HITRAP/CRYRING)

Slow (resonant) extraction

Ultralow extraction (charge change)

Multi charge state/multi component operation

Schottky Mass Spectrometry (SMS) of RIBs

Isochronous mode (TOF detector)

Internal gas jet target

stochastic pre-cooling on the injection orbit

energy 400 (-550) MeV/u
 bandwidth 0.8 GHz (range 0.9-1.7 GHz)
 $\delta p/p = \pm 0.35 \%$ \rightarrow $\delta p/p = \pm 0.01 \%$
 $\varepsilon = 10 \mu\text{m}$ \rightarrow $\varepsilon = 2 \mu\text{m}$



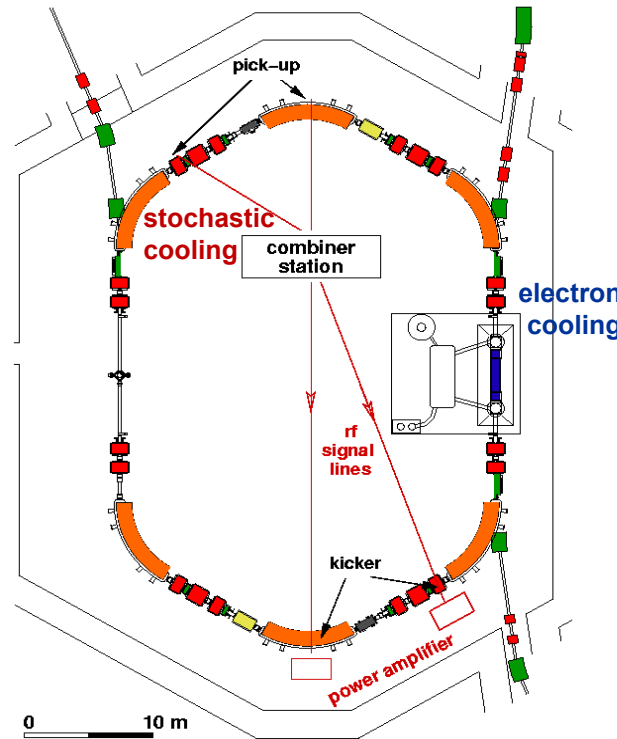
electrodes



combiner station



power amplifiers



electron cooling



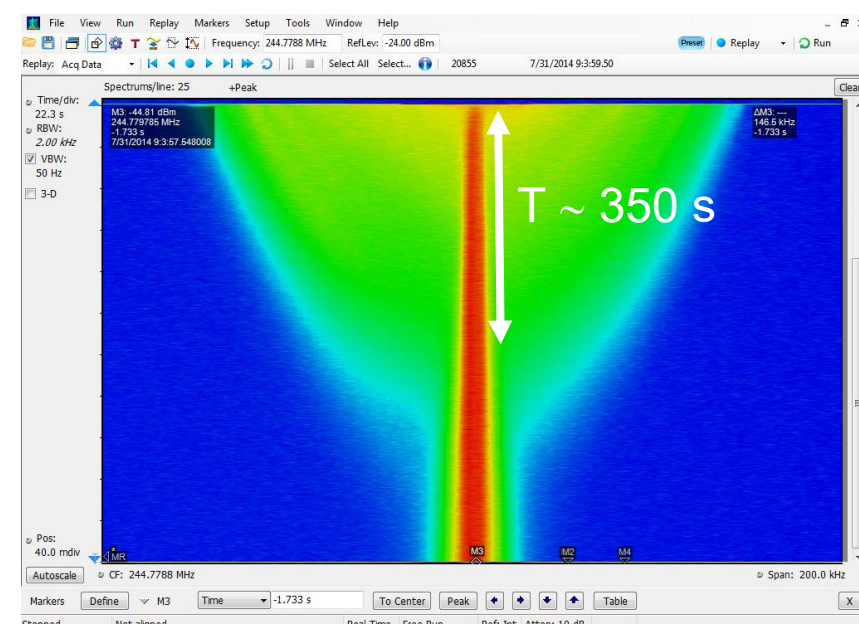
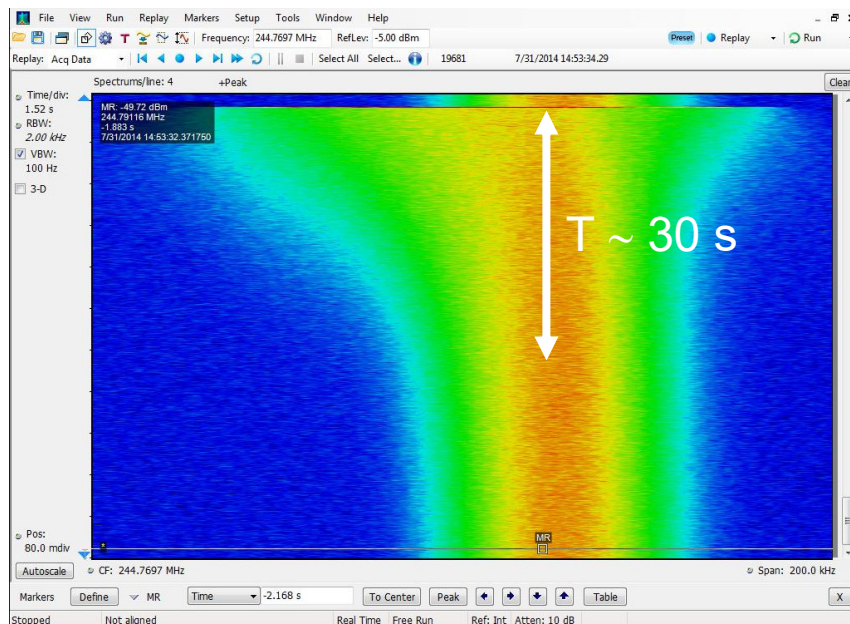
energy	1.6 – 250 keV
current	0.001 – 1 A
diameter	50.8 mm
magnetic field	0.01 - 0.2 T
collection efficiency	0.9998
transverse temperature	0.1 eV
longitudinal temperature	~0.1 meV
vacuum	1×10^{-11} mbar

- The injection energy of the ESR is flexible (30 - 500 MeV/u)
- electron cooling can be applied for any energy ≤ 430 MeV/u
- stochastic cooling is presently available at 400 MeV/u, but could be extended to higher energies (up to 550 MeV/u)
- at 400 MeV/u stochastic cooling is faster than electron cooling

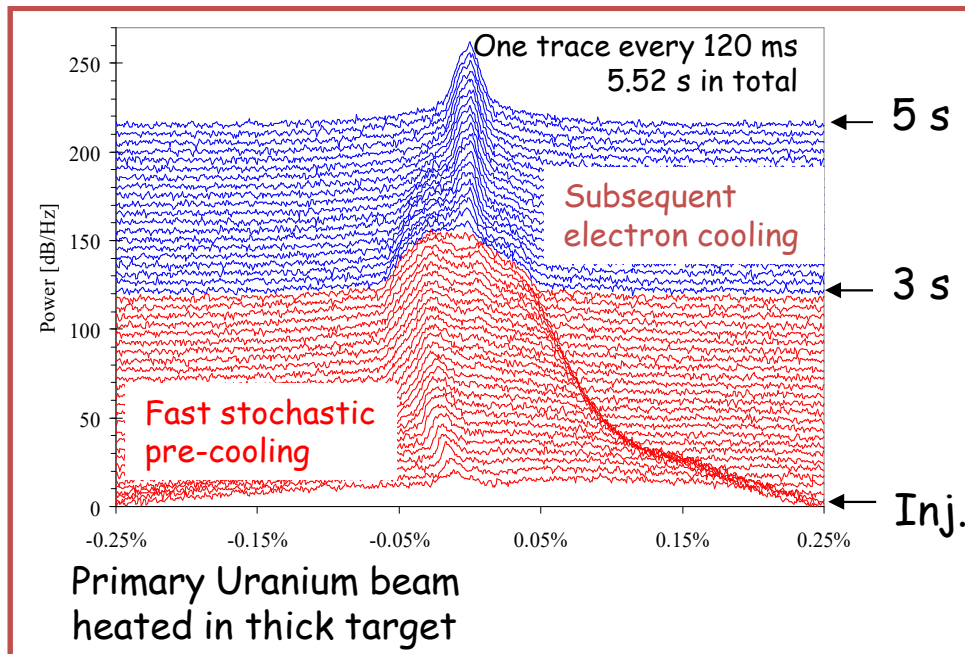
stochastic cooling

protons 400 MeV

electron cooling ($I_{e1}=0.25$ A)



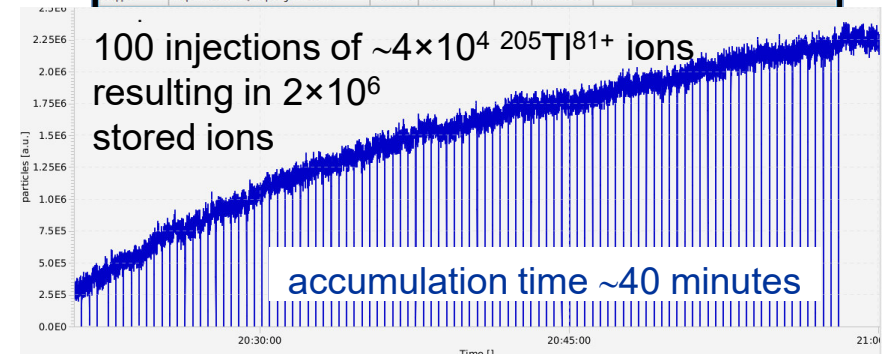
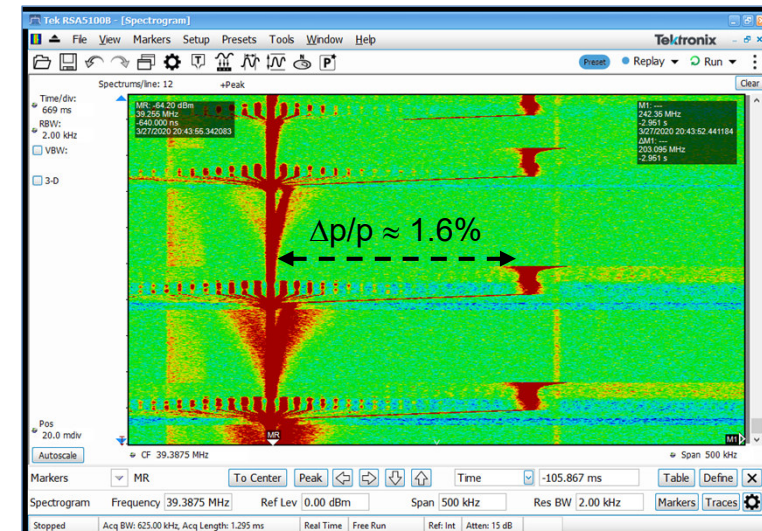
stochastic pre-cooling + **final electron cooling**
immediately after injection

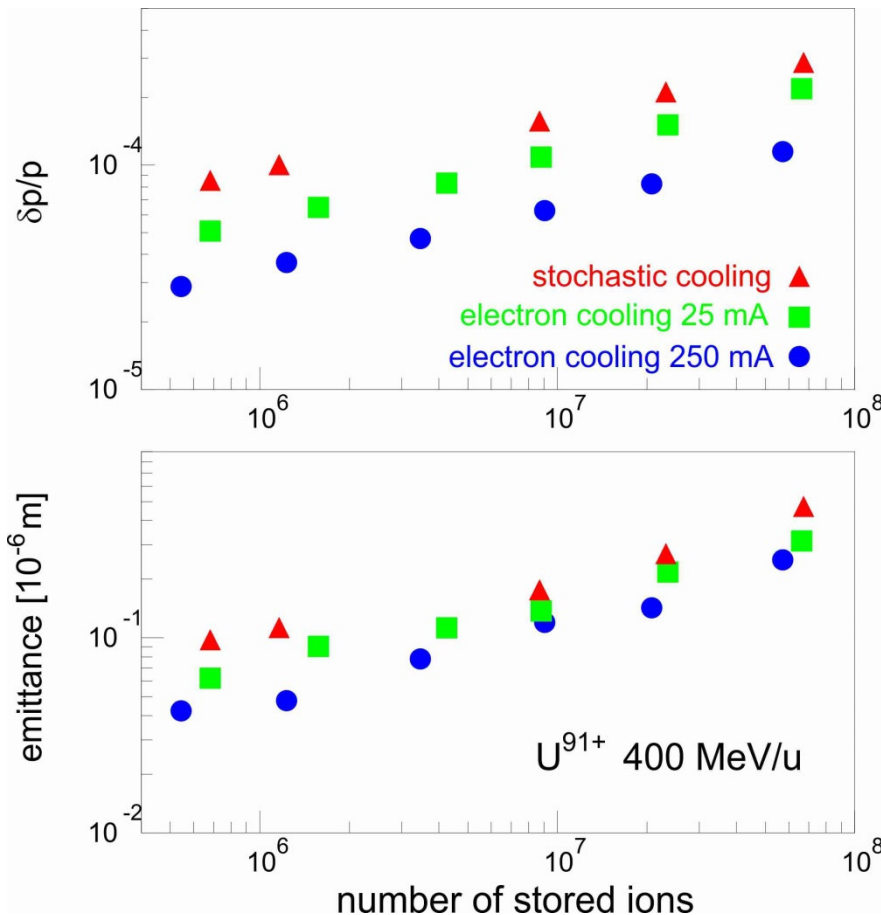


Stochastic pre-cooling reduces the total cooling time to a few seconds, electron cooling only takes some 10 s

Accumulation of secondary beams

- 1) stochastic cooling on injection orbit
- 2) rf stacking
- 3) electron cooling of stack





equilibrium values limited by intrabeam scattering (IBS)

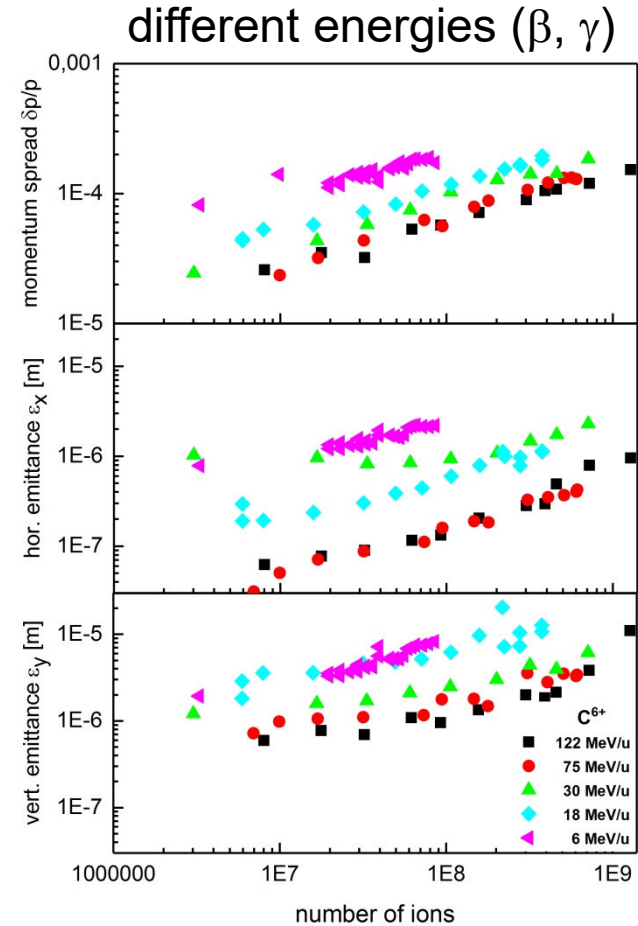
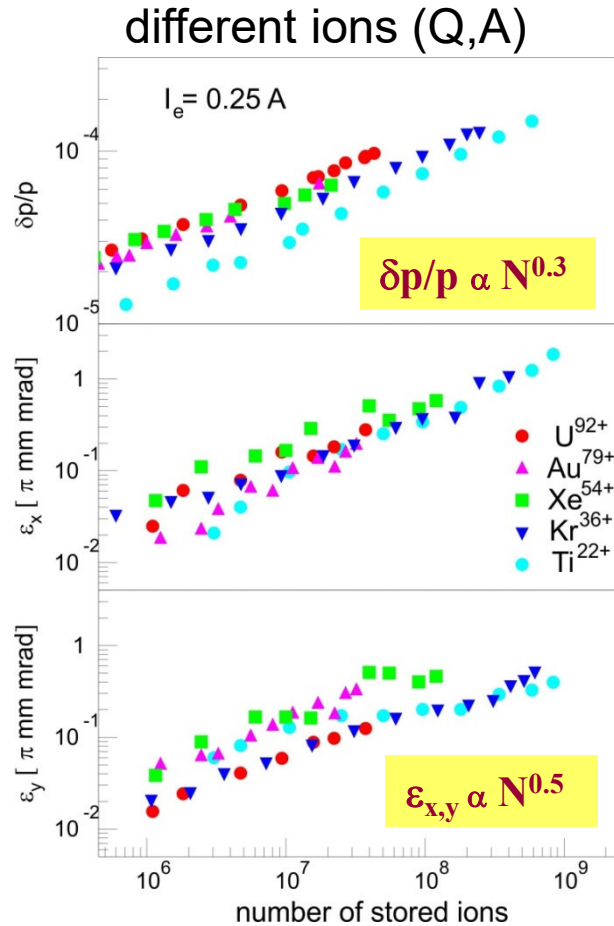
Electron cooling results in smaller momentum spread and emittance compared to stochastic cooling.

The equilibrium is a balance between the cooling rate and the heating rate by intrabeam scattering.

calculated IBS-heating/cooling rate [s⁻¹]

	longit.	transv.
stoch. cool. (Palmer)	0.9 - 2.2	0.5 - 1.3
elec. cool. [25 mA]	2.0 - 6.0	1.4 - 3.3
elec. cool. [250 mA]	18 - 58	7 - 10

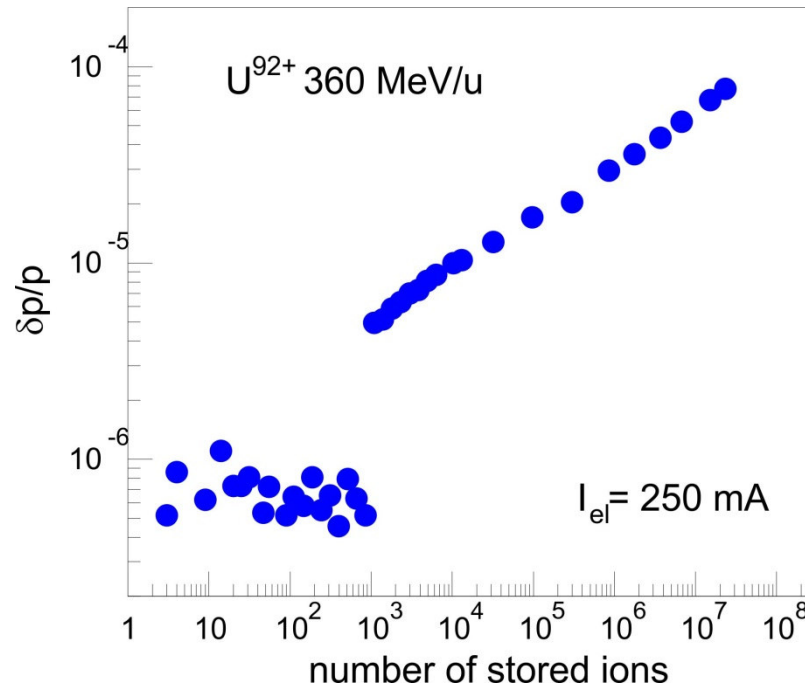
⇒ Electron cooling is more powerful for cold beams.



IBS: total phase space volume increases with ion beam intensity and ion charge

heating rate $\tau_{IBS}^{-1} = \frac{Q^4 e^4}{(Am_i)^2} \cdot \frac{N}{C \epsilon_h \epsilon_v \delta p / p} \cdot \frac{1}{(\gamma^4 \beta^3 c^3)} \cdot 4\pi L_C^{IBS}$

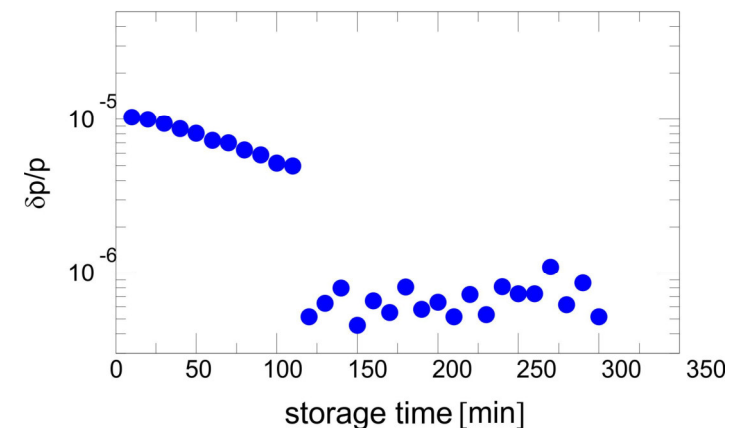
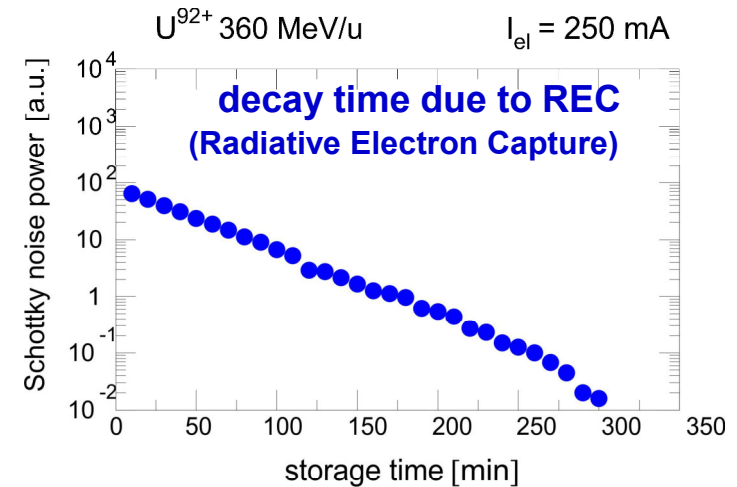
Reduction of momentum spread



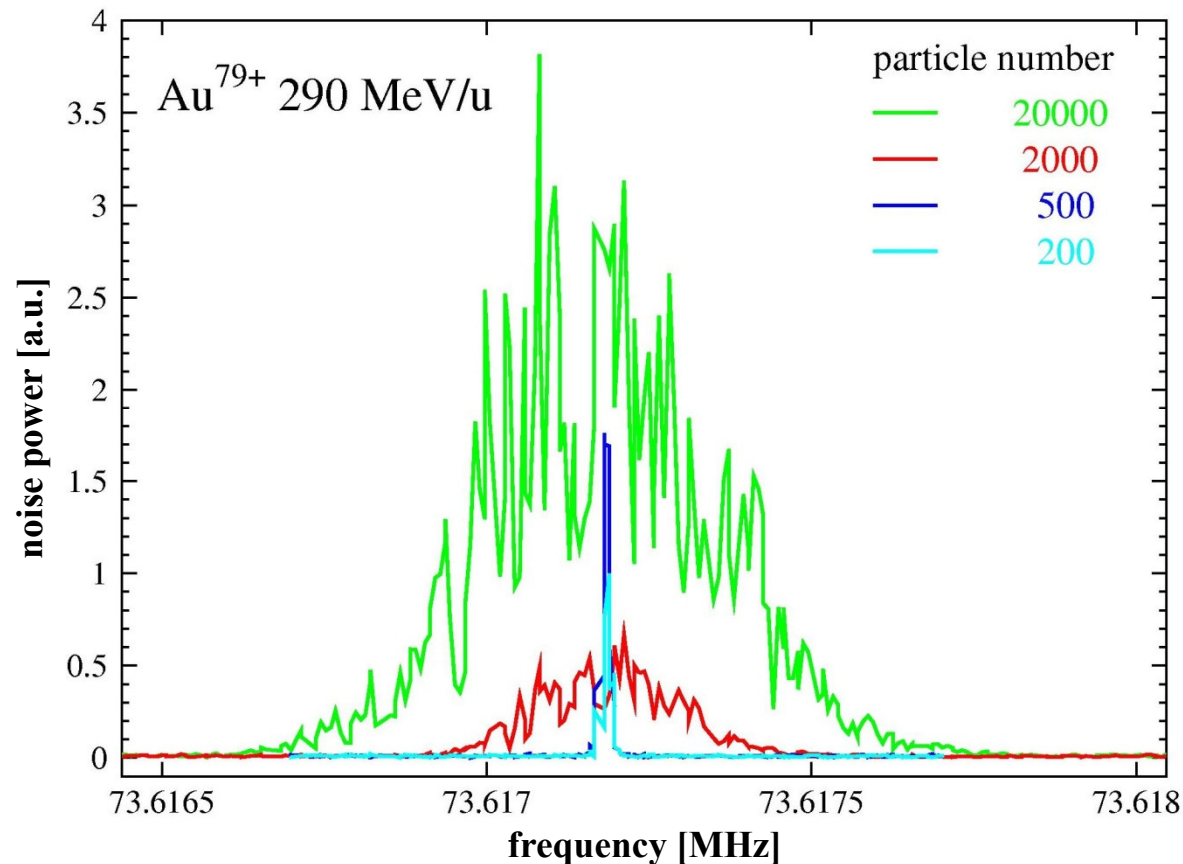
sudden reduction of the momentum spread for less than about one thousand stored ions

⇒ linear ordering in ion string

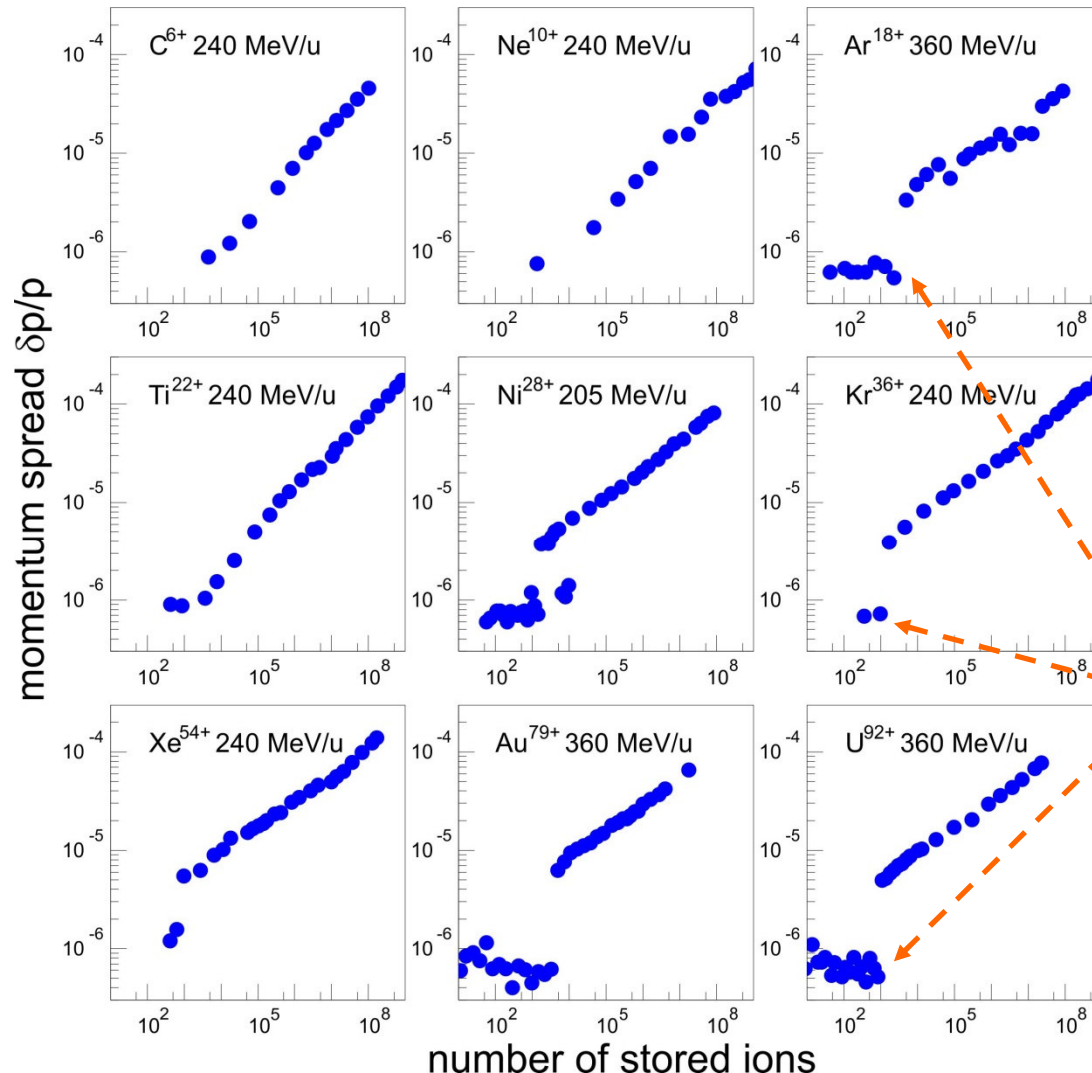
temporal evolution of Schottky noise allows independent determination of particle number



Schottky Noise Signal at Transition to Ordered Beam



The frequency (momentum) spread can be derived from the distribution.
The integrated noise power is proportional to the particle number.



Measurement of momentum spread for a fixed electron current of 0.25 A and variation of stored ion number:

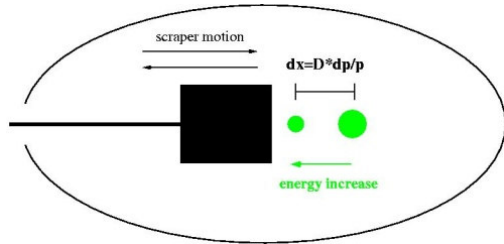
Transition to ordered state is observed around 1000 stored ions.

No dependence on ion species is evident.

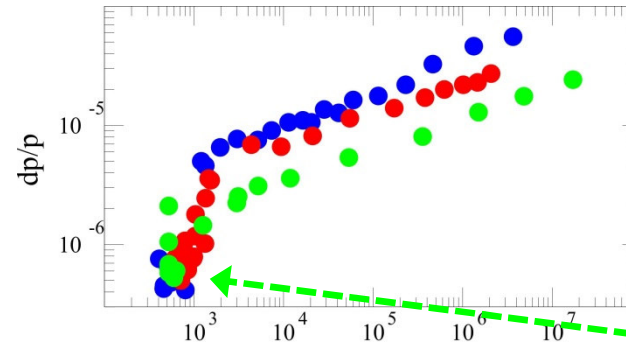
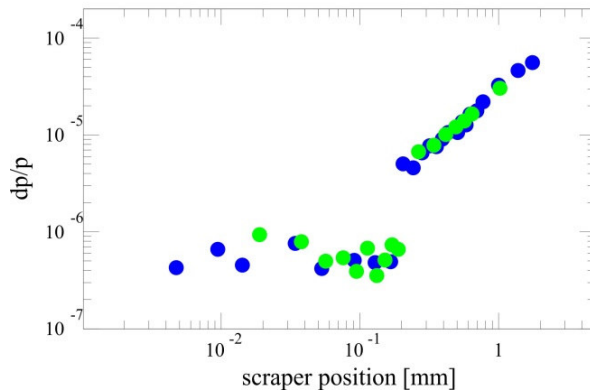
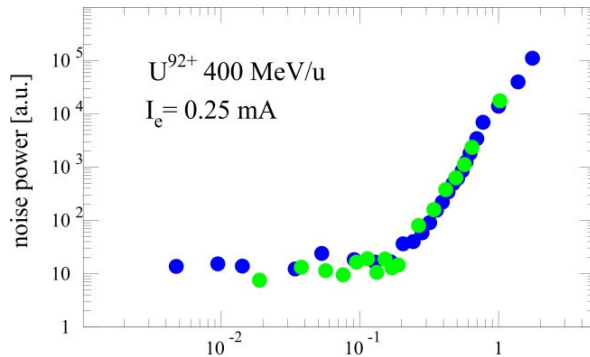
The minimum momentum (frequency) spread is the same for all ion species.

It is determined by the stability of magnet power supplies.

stability of power converters limits the frequency stability

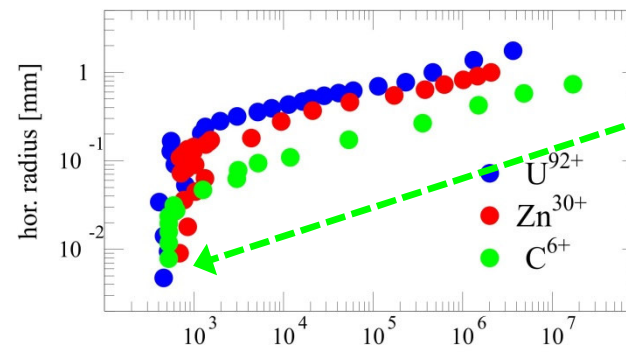


good reproducibility of the scraper measurement

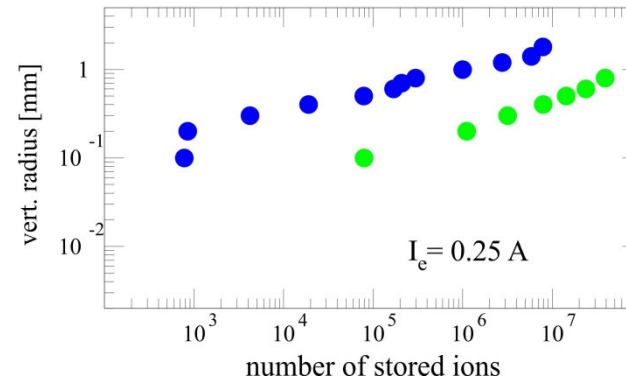


lowest temperature for C⁶⁺ at 4800 MeV

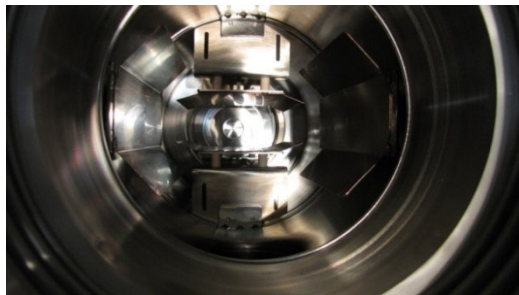
$kT_{||} = 0.26 \text{ meV}$



$kT_x = 0.14 \text{ meV}$



minimum ion temperature of the order of the longitudinal electron temperature \Rightarrow magnetized cooling

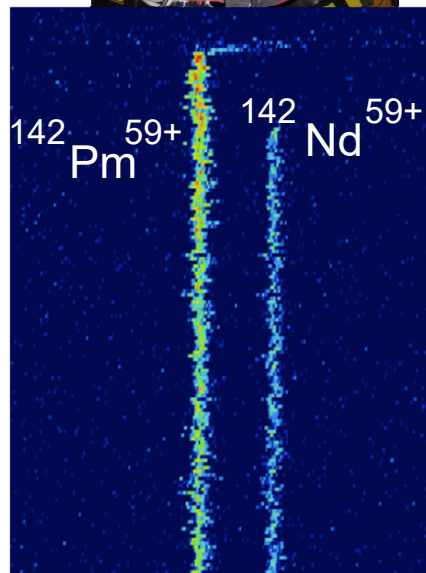
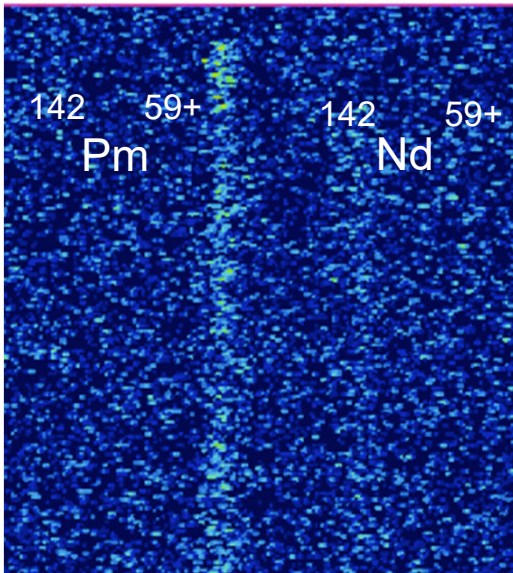
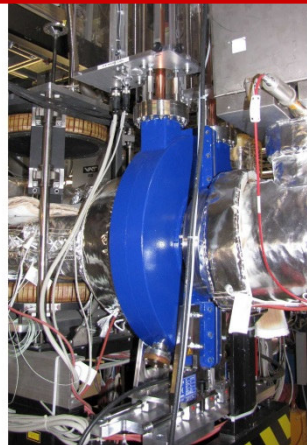


Broad-band
Schottky pickup

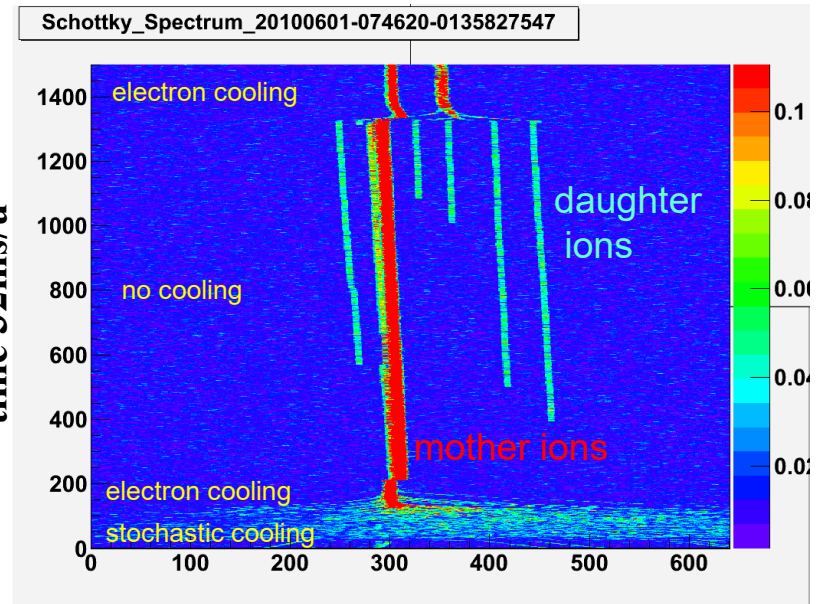
Resonant circuit
at 30th harmonic

Resonator Cavity

124th harmonic



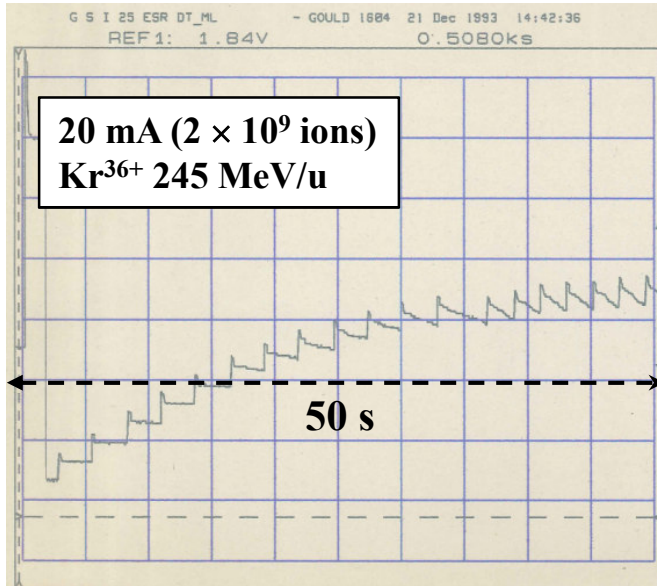
single ion detection (without cooling)



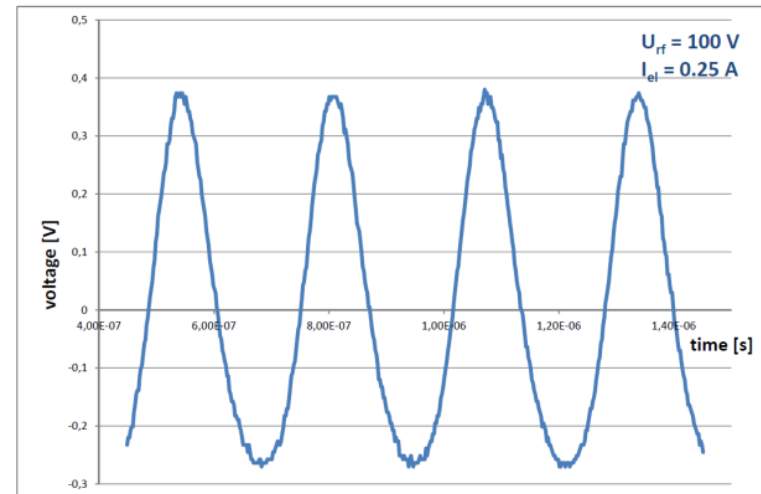
studies of nuclear decays

improved sensitivity
and time resolution

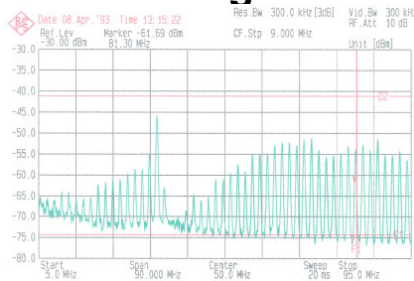
no sensitivity at low beam energy



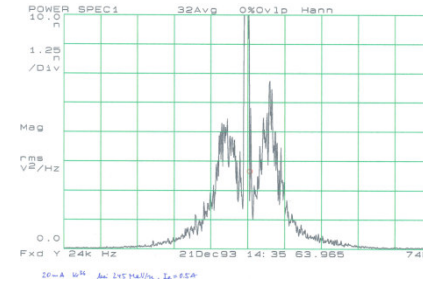
18 mA Xe⁵⁴⁺ 350 MeV/u bunched (h=2, B=0.40) corresponding to over 40 mA coasting beam



high intensity observations



transverse coherent Schottky signal 0-90 MHz



longitudinal Schottky signal with central coherent peak

what matters?
cooling \Rightarrow emittance, momentum spread
orbit, tune, chromaticity, impedances

no feedback system for regular operation, only some basic studies in the past

J. Eichler, T. Stöhrker / Physics Reports 439 (2007) 1–99

relativistic Doppler transformation

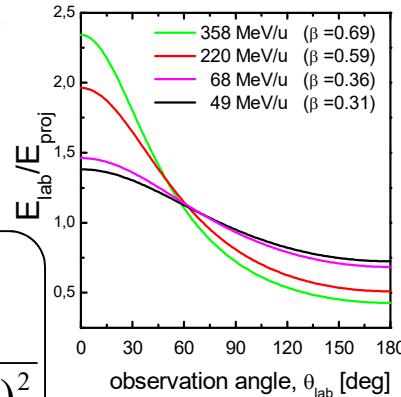
$$E_{\text{lab}} = \frac{E_{\text{proj}}}{\gamma \cdot (1 - \beta \cdot \cos \theta_{\text{lab}})}$$

relativistic transformation of solid angle:

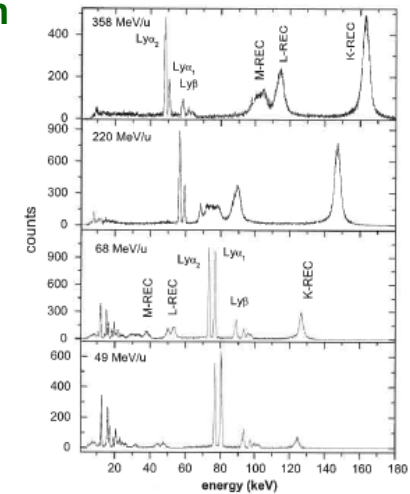
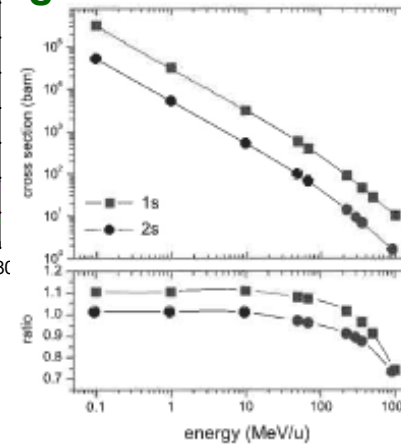
$$\Delta\Omega_{\text{lab}} = \frac{\Delta\Omega_{\text{proj}}}{\gamma^2 \cdot (1 - \beta \cdot \cos \theta_{\text{lab}})^2}$$

E_{lab} : photon energy in lab system

E_{proj} : photon energy in emitter system



radiative recombination U^{92+}



X-rays of REC $U^{92+} \rightarrow N_2$

Benefits for Atomic Physics

deceleration reduces:

- dependence on velocity and observation angle θ_{lab}
- influence of energy and intensity

deceleration:

- provides increased reaction cross sections
- new reaction channels (transitions, polarization)

Benefits for Nuclear Physics

- access to the Gamov window (nuclear physics of stars)
- studies in inverse kinematics

Ni²⁸⁺ 400 → 30 → 4 MeV/u

1100 μA → 180 μA → 25 μA

45%

37%

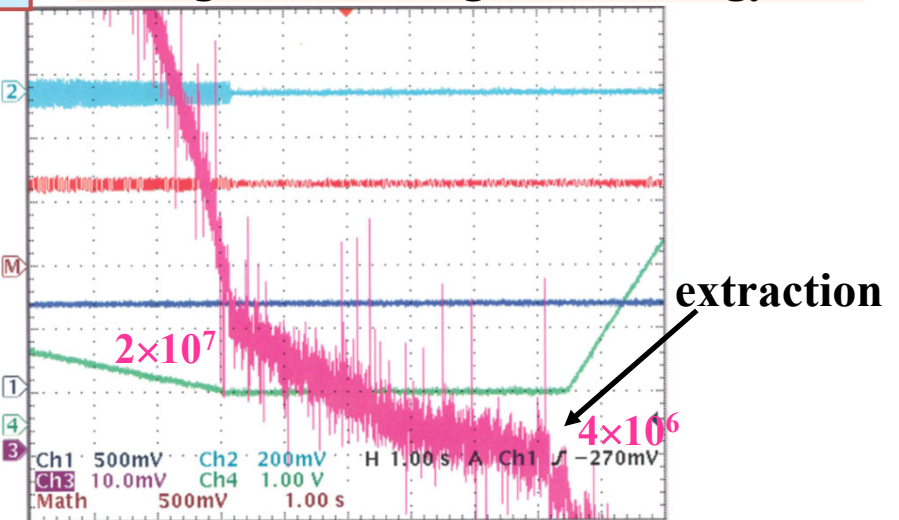
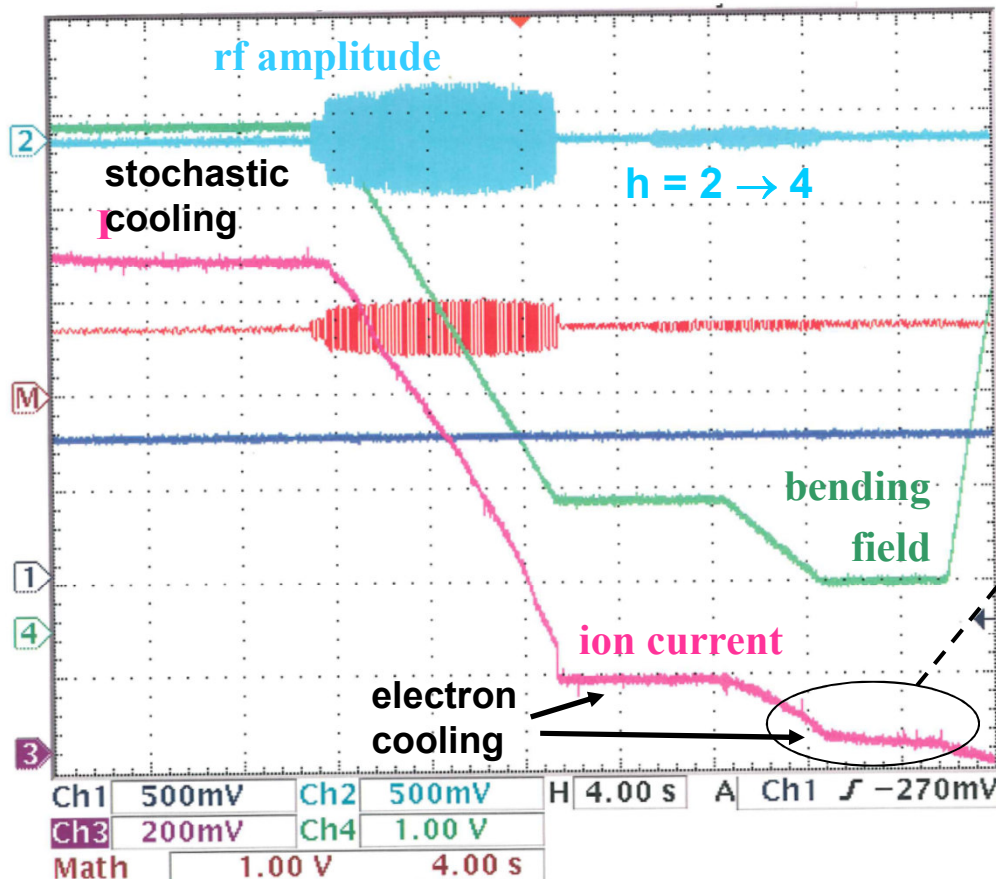
cycle time

45 s

Main losses:

End of ramp

Storage and cooling at low energy



beam half life (vacuum dominated):

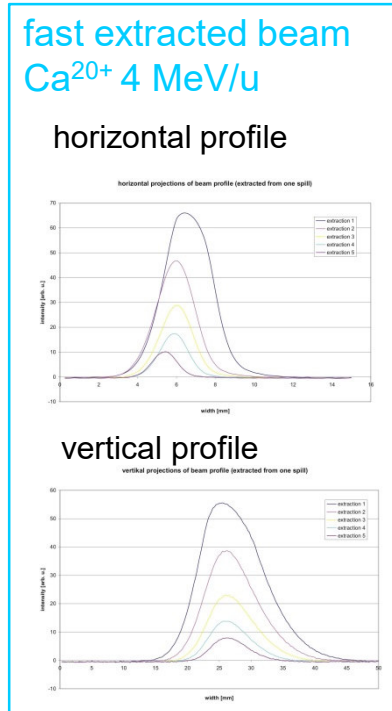
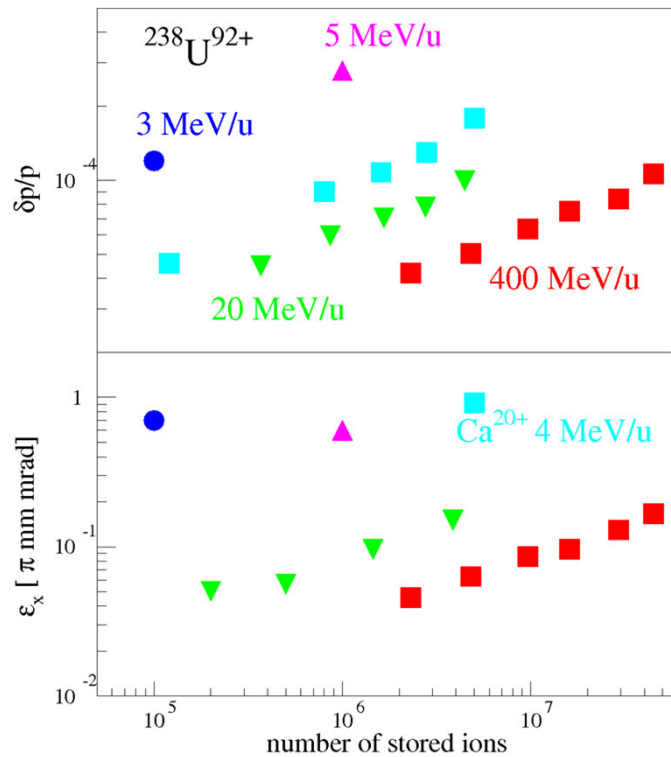
30 MeV/u: $T_{1/2} \approx 480$ s

4 MeV/u: $T_{1/2} \approx 2$ s

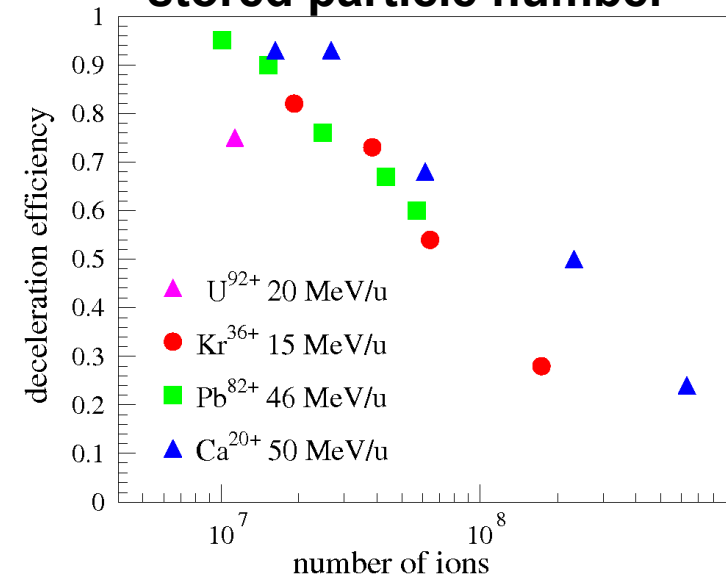
limits to faster cycle time:

- orbit control for fast ramp rate
- cooling times

Beam parameters due to equilibrium between IBS and electron cooling



Losses increase with stored particle number



Control of the beam parameters, emittance (cooling) and bunch length (rf voltage), is challenging

Space charge limit due to incoherent tune shift

$$\Delta Q_x = \frac{r_p Z^2 N g}{\pi A \beta^2 \gamma^3 B (\epsilon_x + \sqrt{\epsilon_x \epsilon_y} Q_x / Q_y)}$$

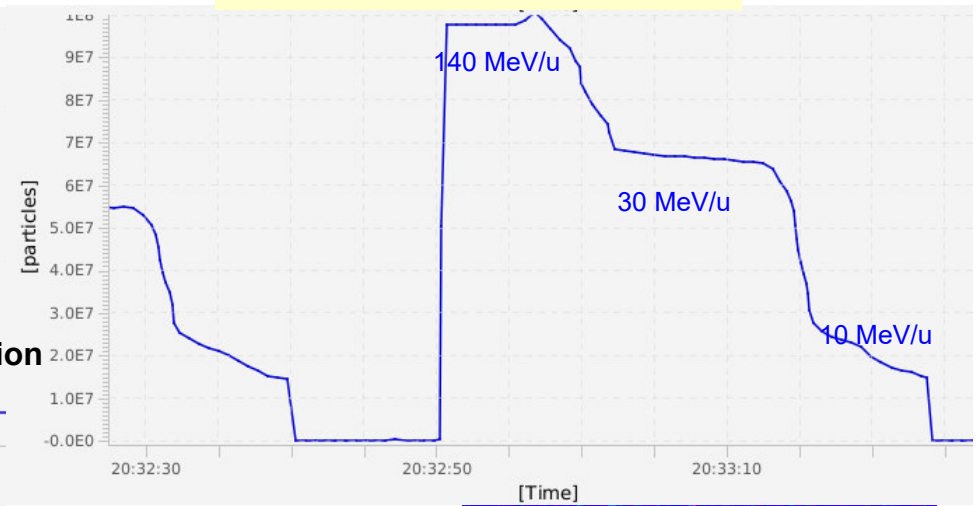
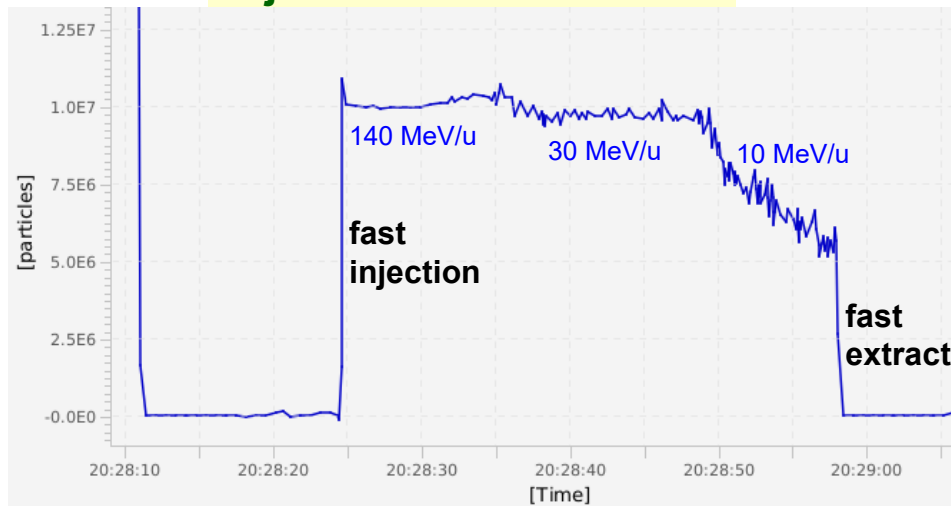
$$\Delta Q_x \simeq \frac{r_p Z^2 N}{2\pi A \beta^2 \gamma^3 B \epsilon_x}, \quad g = 1, \epsilon_x = \epsilon_y, Q_x = Q_y$$

space charge limit U^{92+} at 4 MeV/u for $\Delta Q_x = 0.1$:
 $2 \mu\text{s}$ bunch with $\epsilon_x = \epsilon_y = 10^{-6} \text{ m} \Rightarrow N \leq 1 \times 10^7$

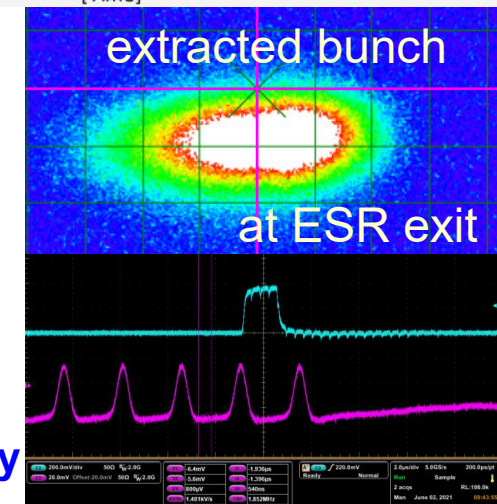
deceleration Ag^{47+} from 140 to 30 and 10 MeV/u and extraction to CRYRING

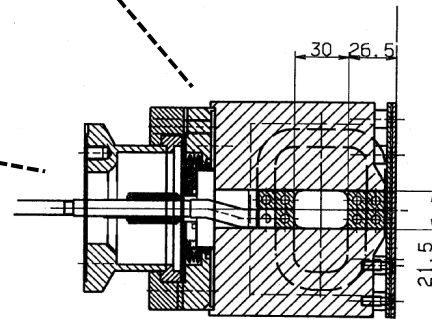
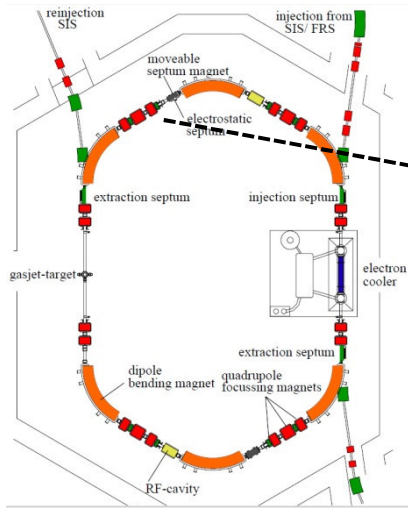
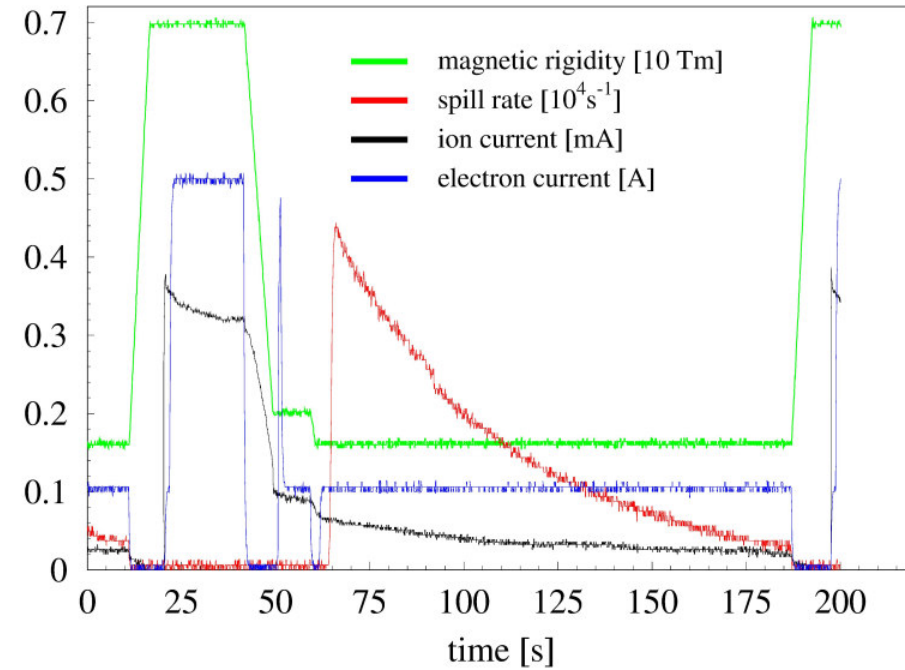
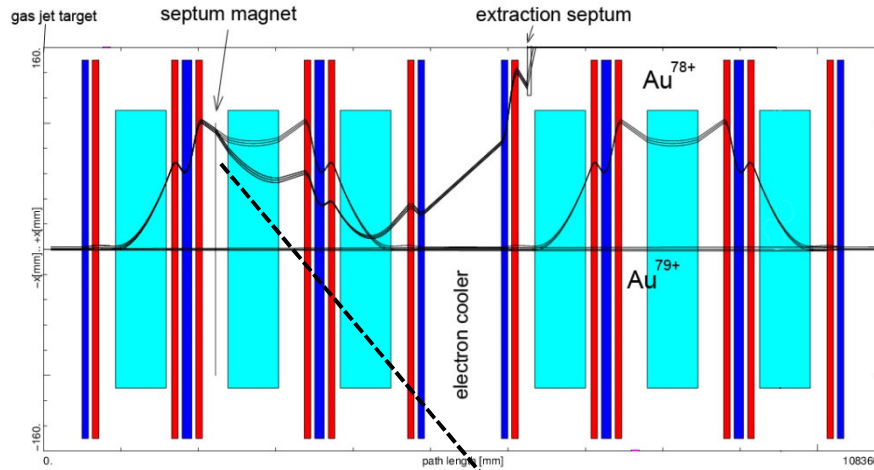
injection of 1×10^7 ions

injection of 1×10^8 ions



- cycle time of 35 second
- electron cooling at 140, 30 and 10 MeV/u
- poor lifetime (vacuum) with half life 12 s at 10 MeV/u
- extracted particle number $1-2 \times 10^7$ (cooled) ions per cycle
limited by slow regulation of electron cooler accel. high voltage
- **bunching into single bunch with new MA filled Barrier Bucket cavity**



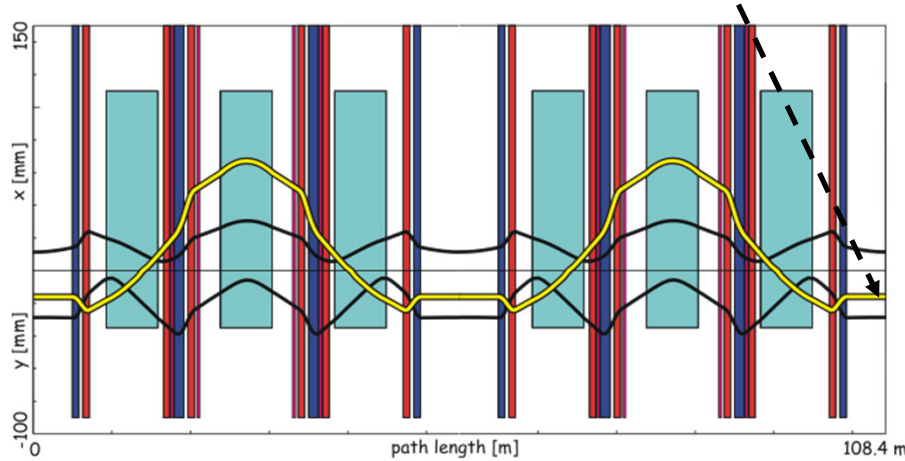


septum magnet installed in UHV

- extracted cooled beam at injection energy or after deceleration (so far: to $\geq 20 \text{ MeV/u}$)
- small emittance
- good control of extraction rate
- extraction time flexible: seconds to hours

Isochronous Mode $\gamma_t = \gamma$

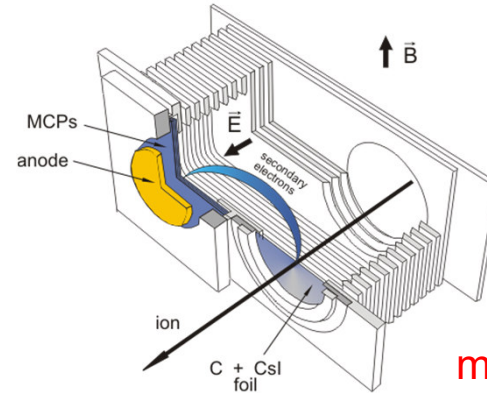
Isochronous Mode with large negative dispersion (- 8 m) at ToF detector



$$\frac{\Delta f}{f} = -\frac{\Delta T}{T} = \frac{1}{\gamma_t^2} \frac{\Delta(m/q)}{m/q} - \left(1 - \frac{\gamma^2}{\gamma_t^2}\right) \frac{\Delta v}{v}$$

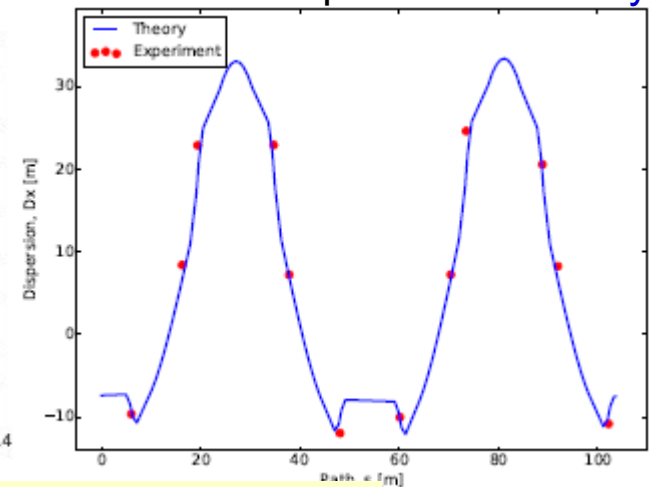
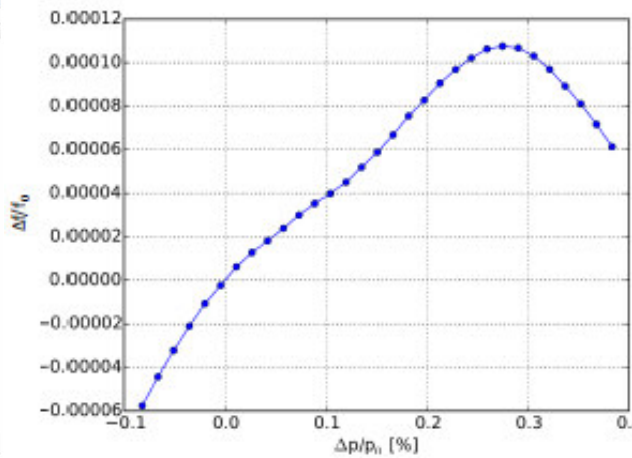
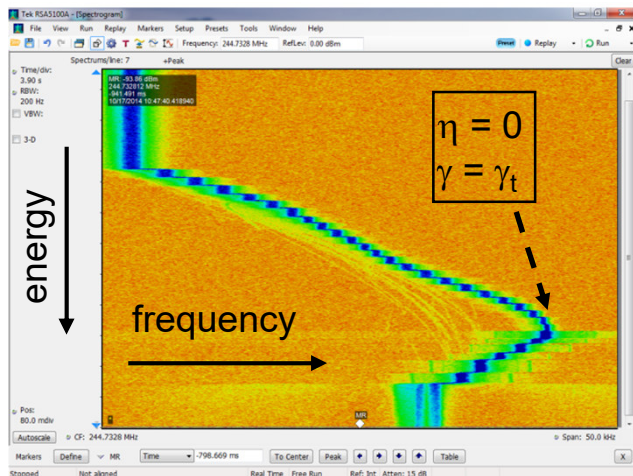
TOF Detector

(= 0, if $\gamma = \gamma_t$)

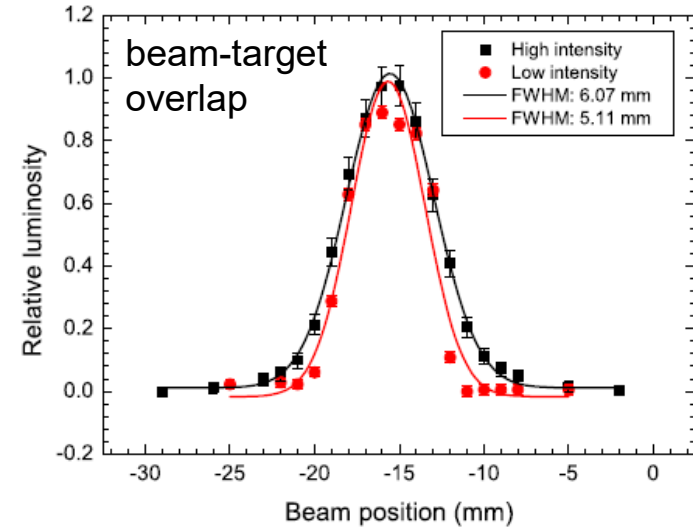
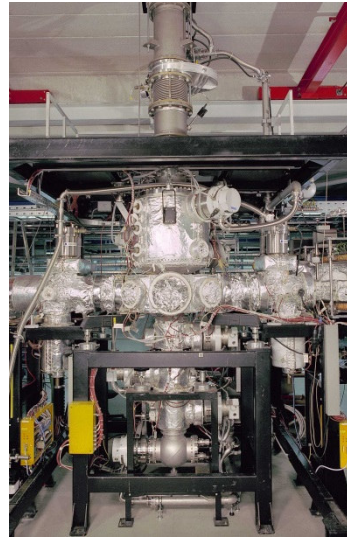
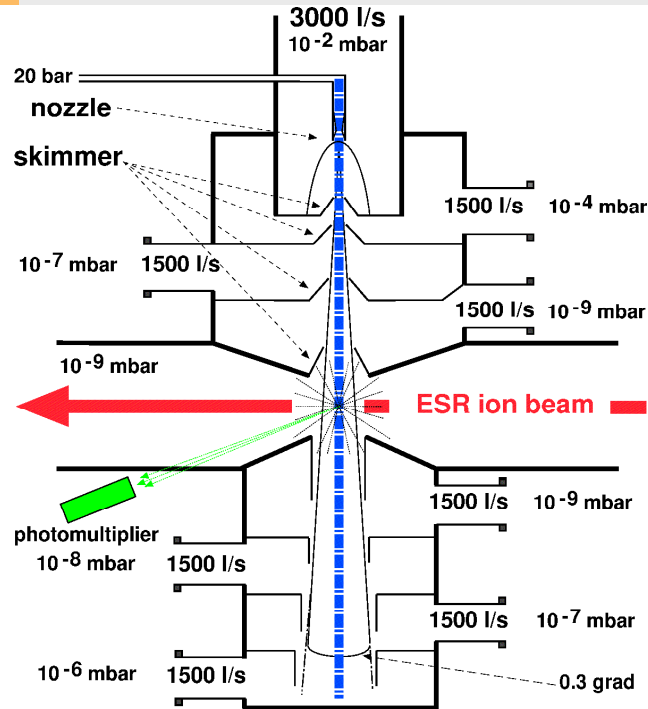


used in mass measurements with a ToF detector

measured dispersion value compared with theory



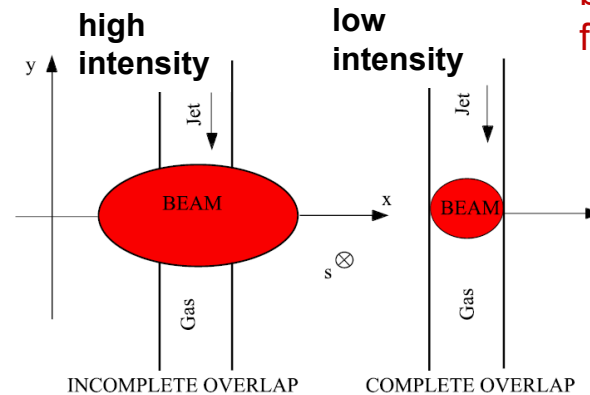
measurement of transition energy by variation of beam energy with the electron cooler



beam diameter 3.3 mm
for 1.2×10^9 Ar¹⁸⁺ ions

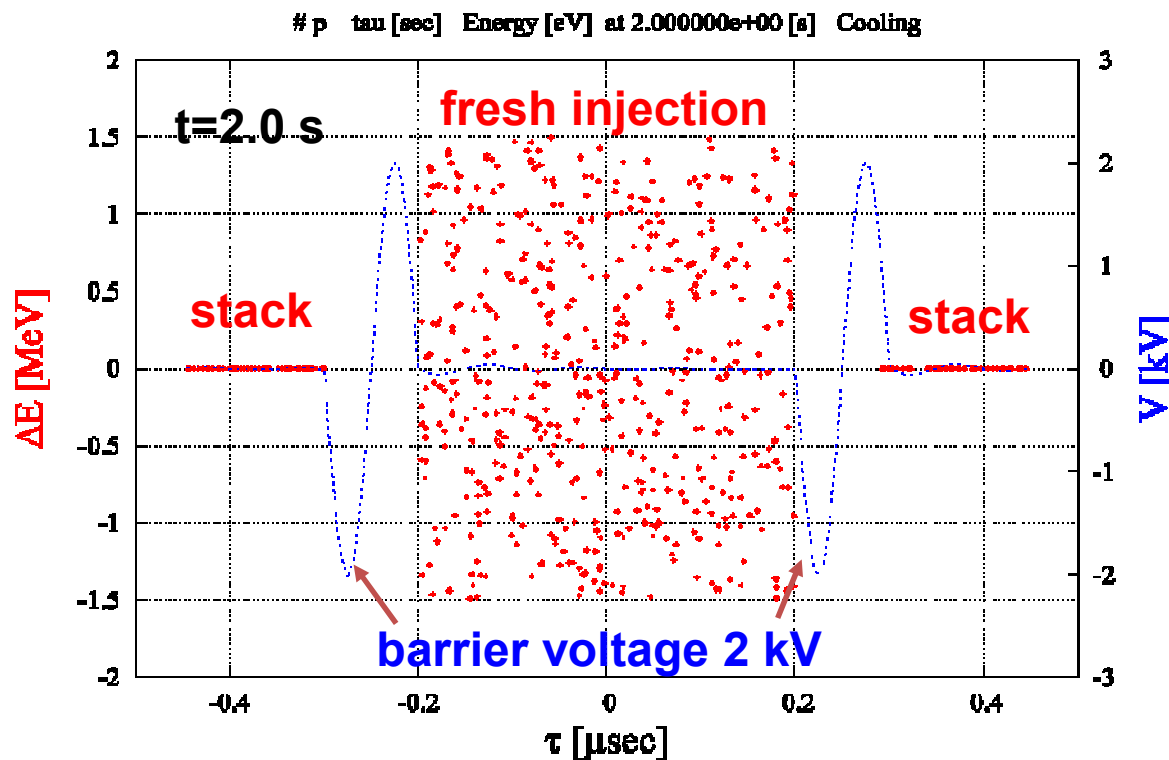
Target gas	Area density [cm ⁻²]	T ₀ [K]
Helium	1×10^{13}	20
Hydrogen	3×10^{13}	40
Nitrogen	8×10^{12}	130
Argon	3.5×10^{12}	300
Krypton	1.5×10^{12}	300
Xenon	6×10^{12}	300

presently available target densities



The luminosity in target experiments with cooled beams does not linearly increase with the beam intensity.

basic idea: confine stored beam to a fraction of the circumference, inject into gap
 apply strong electron cooling to merge the two beam components
 ⇒ fast increase of intensity (for low intensity RIBs)



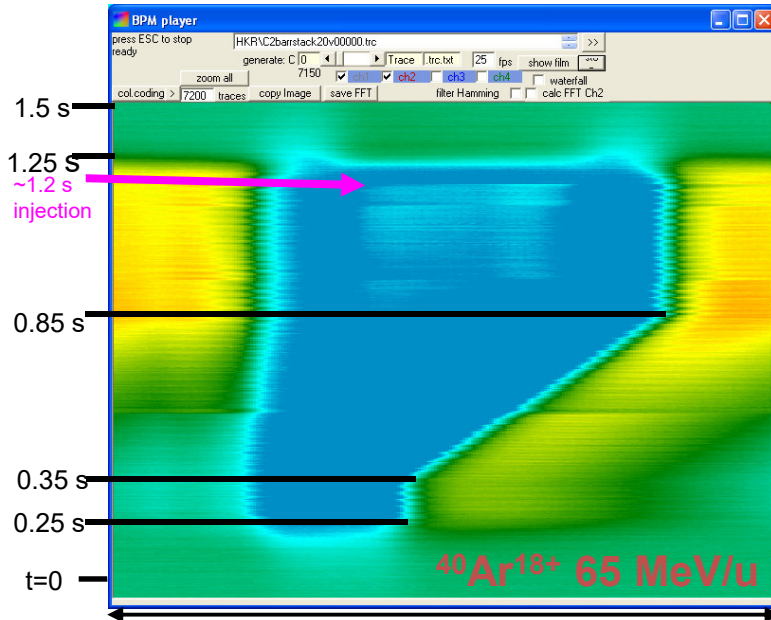
$^{132}\text{Sn}^{50+}$
 $E_k = 740 \text{ MeV/u}$

Longitudinal stacking
 with Barrier Buckets

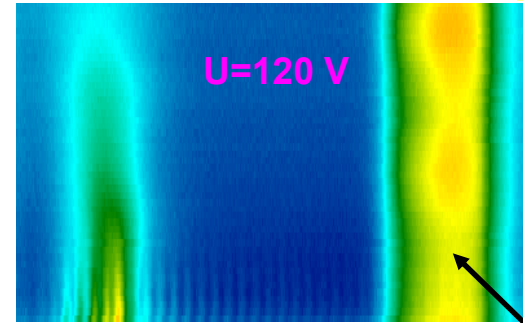
(simulation by T. Katayama)

revolution time $0.9 \mu\text{s}$

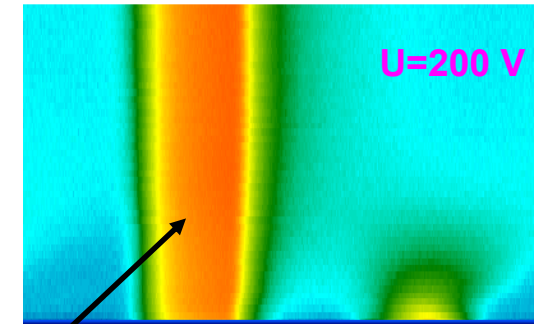
moving barriers



fixed barrier



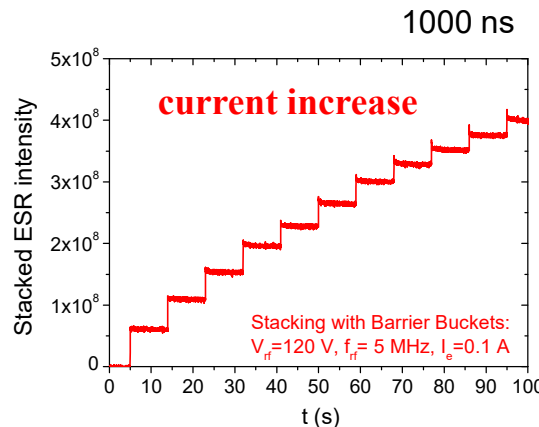
h=1 unstable fixed point



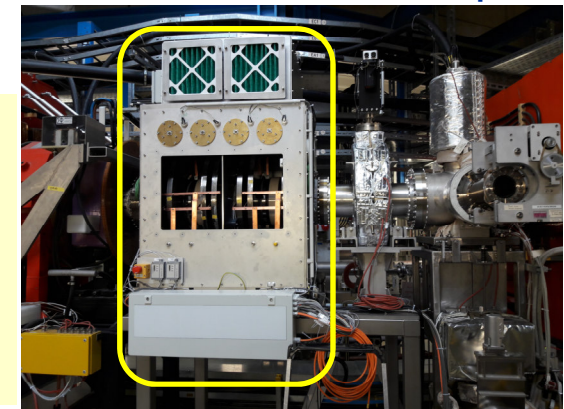
stack $^{124}\text{Xe}^{54+}$ 154 MeV/u

three schemes were successfully tested:

moving barriers, fixed barrier, h=1 unstable f. p.
efficient accumulation



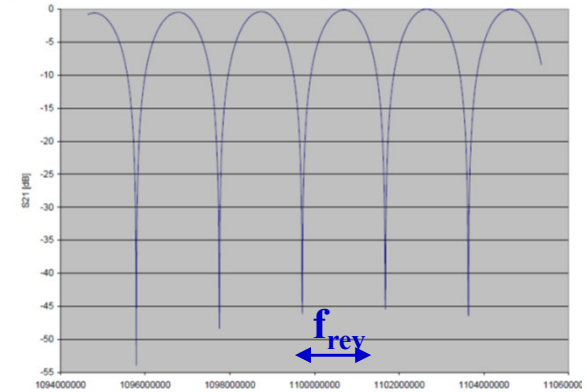
after successful proof of principle with modified ferrite cavity a dedicated magnetic alloy filled barrier bucket cavity was installed which allows flexible rf voltages



same accumulation scheme with stochastic cooling tested



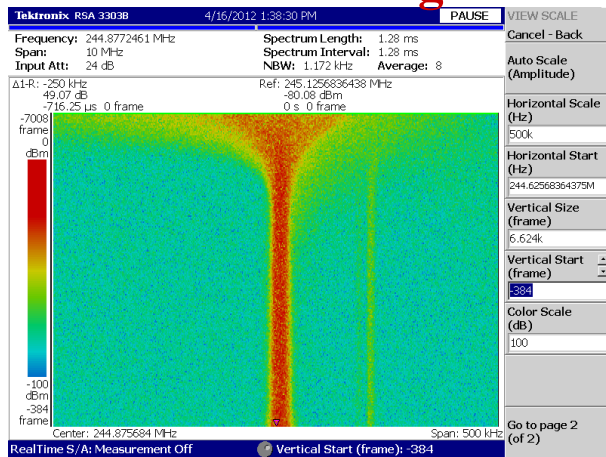
Optical delay line installed in the ESR for tests of TOF and notch filter cooling using existing electrodes designed for Palmer cooling



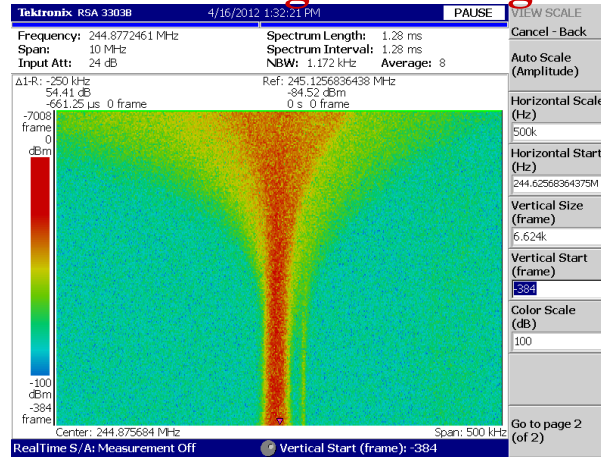
notches

Test of FAIR set-up

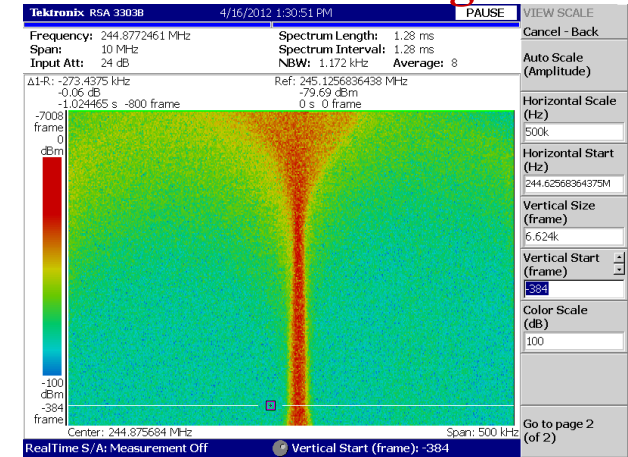
Palmer cooling



Time-of-Flight cooling



Notch filter cooling



Ar¹⁸⁺ 400 MeV/u

Various operation modes of the ESR have been developed over the years

Beam cooling is an indispensable prerequisite for ESR operation

ESR still serves a broad physics program in atomic and nuclear physics with highly charged ions which will continue into the operation of FAIR

ESR serves as a test bed for concepts required for FAIR

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