

Electroweak Effects in Boosted Top Jet Production

in collaboration with A. H. Hoang and M. Procura

FAKT Workshop 2022 (Bruck a.d. Mur)

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Particles and Interactions

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- Observable and pure QCD
- Electroweak effects
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Top quark mass determinations from $t\bar{t}$ -production

- Top group @ Uni Wien: precision **top mass determinations**
- **Top pair-production**: precision studies at future linear colliders
 - Threshold scans -> *non-relativistic tops*
 - Direct reconstruction from top decay products -> *boosted tops*
- Usual tools:
 - Multipurpose Monte Carlo (MMC)
 - Perturbative fixed-order (FO)
- Our work: focus on **analytic calculations** in **factorization approaches**
 - Systematically improvable (order and/or expansion/logarithmic accuracy)
 - Controlled uncertainties
 - Analytic resummation of large logarithms

$$m_t^{\text{pole}} = 172.5 \pm 0.7 \text{ GeV} \quad (\text{cross section})$$

$$m_t^{\text{MC}} = 172.76 \pm 0.30 \text{ GeV} \quad (\text{direct reconstruction})$$

[PDG, 20]

$$\Delta m_t = 50 \text{ MeV} \quad (\text{threshold scan})$$

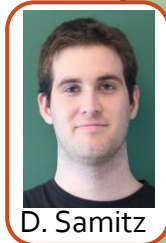
$$\Delta m_t = 100 \text{ MeV} \quad (\text{direct reconstruction})$$

Projected uncertainties for ILC500 [Zarnecki et al, 11]

- $t\bar{t}$ (stable): $\text{QCD}_{\text{NNLO}}^{\text{NNLL}} + \text{EW}_{\text{NLO}}$
[Czakon et al, 16, 17, 18]
- $t\bar{t} \times \text{top decay (NWA)}$: $\text{QCD}_{\text{NNLO}} + \text{EW}_{\text{NLO}}$
[Gao et al, 13, 17]
- $bW^+(vl^+)\bar{b}W^-(\bar{v}l^-)$: $\text{QCD}_{\text{NLO}} + \text{EW}_{\text{NLO}}$
[Heinrich et al, 14] [Denner et al, 18]

Top physics group @ Uni Wien

Threshold
Boosted tops



D. Samitz

MC/Parton shower aspects

- Shower effects in m_t^{MC} vs. m_t^{QFT}
- Soft and hadronization impact on m_t^{MC}
- Single-purpose MC for $e^+e^- \rightarrow t\bar{t}$



A. Widl

Top production at threshold

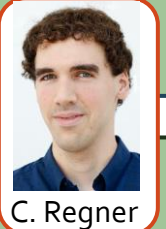
- Implementation of MSR mass scheme
- Improving interpolation between kinematic regimes

supported by

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S. Plätzer



C. Regner

Differential top

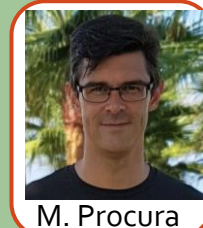
- Top decay-sensitive observables
- Unstable W
- Lepton spectra and (b-)jet observables
- QCD corrections in differential distributions



A.H. Hoang

Electroweak effects

- Systematic inclusion of EW effects
- Coherent resummation of EW and QCD logs
- Finite-lifetime effects



M. Procura



D. Lechner



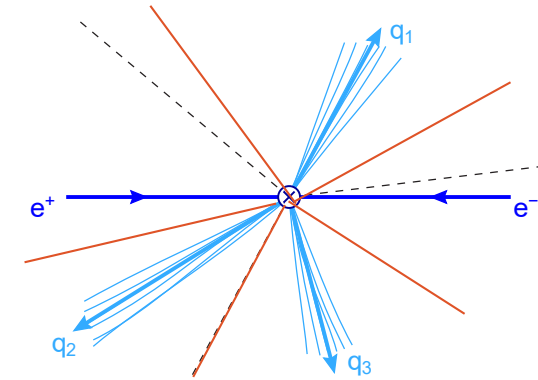
I. Ruffa

2-jettiness: a case study

- N -jettiness τ_N [Stewart et al, 10]
 - Field-theoretically well-defined
 - Simple, yet general in its IR-structure
 - **Top mass-sensitive !**

$$\tau_N \equiv \sum_k \min_i \frac{2q_i \cdot p_k}{Q_i}$$

Jet directions (pre-determined) Particle momenta



Jet sector boundaries without/with beam

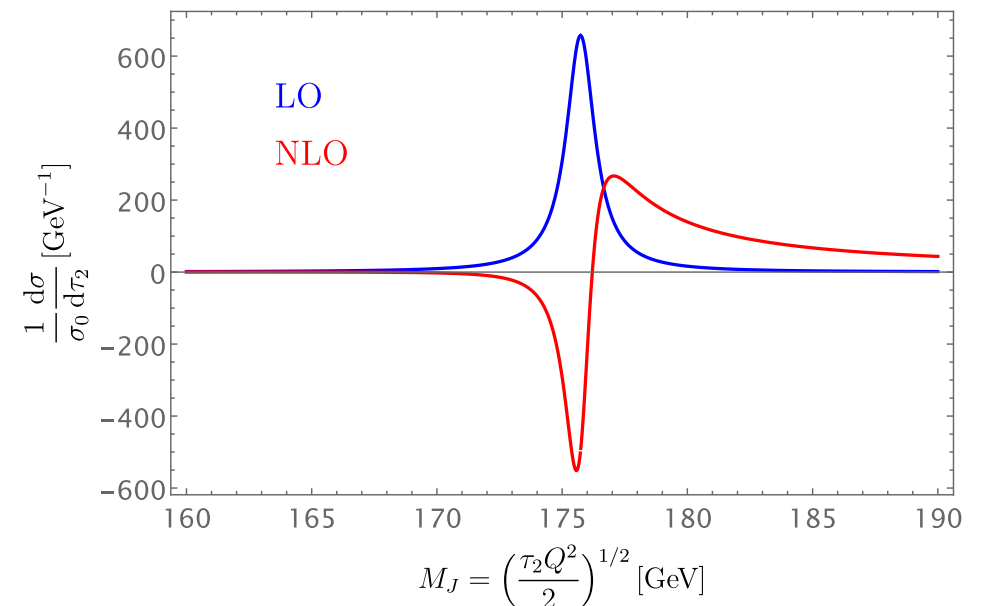
- Important:
 - Treating **EW (QED!) effects** necessitates including the beam (jets)
 - Not very realistic concerning beam treatment
 - Important future direction: more differential observables!

- Tree-level: Breit-Wigner

$$\frac{1}{\sigma_0} \left(\frac{d\sigma}{d\tau_2} \right)^{(0)} = f_{\text{BW}} \left(\tau_2 - \tau_{\text{min}}, \frac{2\Gamma_t}{Q\varrho_t} \right)$$

$$\begin{aligned}
 Q &= \sqrt{\hat{s}} && \dots \text{ CM energy} \\
 \varrho_t &= Q/m_t && \dots \text{ boost factor} \\
 \beta_t &= \sqrt{1 - 4/Q^2} && \dots \text{ quark velocity} \\
 \tau_{\text{min}} &= 1 - \beta_t && \dots \text{ stable threshold}
 \end{aligned}$$

- Radiative corrections (QCD NLO and beyond) sizable



Radiative corrections to 2-jettiness for $t\bar{t}$ production

- Of interest: **boosted jet** regime: $E_{jet}^2 \gg p_{jet}^2 \sim m_t^2$, i.e. **small τ_2**
- Complex beyond LO because of multiple scales
- Due to **scale separation**, e.g. in the **peak region**

scale	magnitude [GeV]
Q	1000
m_t	173
Γ_t	1.4
Λ_{QCD}	< 1

$$Q^2 \gg m_t^2 \gg Q^2 \tau_2 - m_t^2 \gtrsim m_t \Gamma_t \gg \rho_t^{-1} m_t \Gamma_t \gtrsim \Lambda_{\text{QCD}}$$

Top physics = multiscale problem

- Large **QCD/QED** logs: $\alpha_s^n \ln^k \tau_2$, $\alpha_{em}^n \ln^k \tau_2$
- (EW) **Sudakov** logs: $\alpha_s^n \ln^k Q/m_t$, $\alpha_w^n \ln^k Q/M_w$



Resummation necessary

- **Factorization** properties of N -jettiness:

$$\hat{\tau}_2 = \hat{\tau}_c + \hat{\tau}_s + \mathcal{O}(\tau_2)$$

Indicates power-suppressed/
non-singular terms

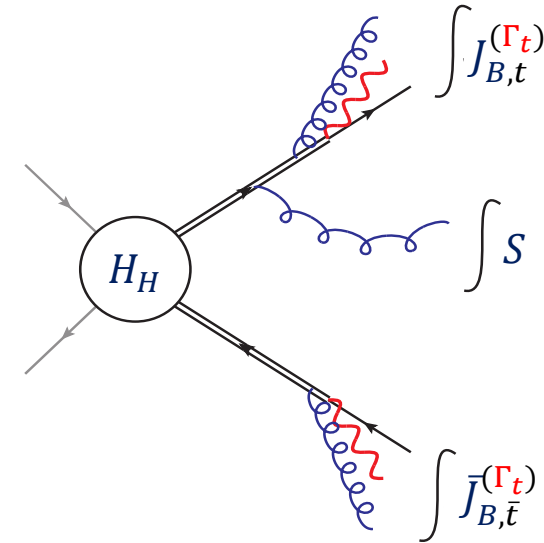
- **Effective Field Theory** (EFT) approach:
 - Soft-Collinear Effective Theory (**SCET**) [Stewart et al, 01]
 - boosted Heavy Quark Effective Theory (**bHQET**) [Fleming et al, 07]

Factorization for 2-jettiness in pure QCD

$$\frac{1}{\sigma_0} \frac{d\sigma}{d\tau_2} = H_Q(Q, \mu) \times H_M(\varrho_t, m_t, \mu)$$

[Fleming et al, 07]

$$\times J_{B,t}^{(\Gamma_t)}(Q^2 \tau_2, m_t, \mu) \otimes \bar{J}_{B,\bar{t}}^{(\Gamma_t)}(Q^2 \tau_2, m_t, \mu) \otimes S(Q\tau_2, \mu) + \mathcal{O}\left(\lambda, \varrho_t^{-1}, \frac{\Gamma_t}{m_t}\right)$$



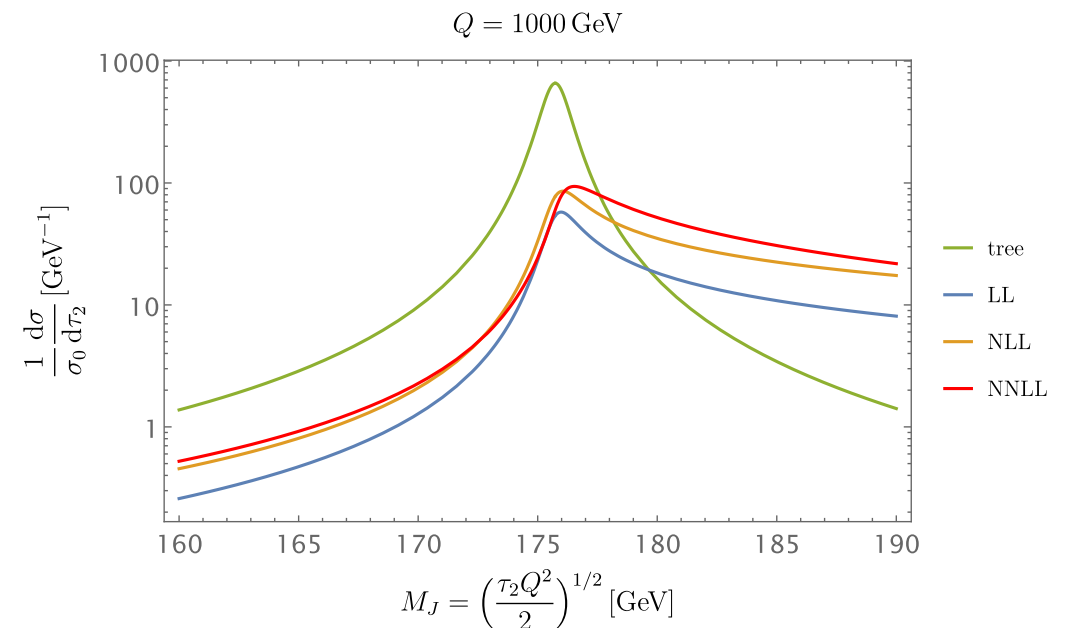
- Factorization \equiv product \times or convolution \otimes
- EFTs provide the possibility to resum large logarithmic corrections by **renormalization group evolution**

$$\mu \frac{d}{d\mu} F(t, \mu) = \gamma_F(t, \mu) F(t, \mu)$$

$$\gamma_F(t, \mu) = \Gamma_{\text{cusp}}[\alpha_s] \ln(i\mu e^{\gamma_E} t) + \gamma[\alpha_s]$$

- Formalism can be extended to resum **electroweak Sudakov logarithms** consistently with QCD logs

[Chiu et al, 07, 08, 09]



QCD state-of-the-art (peak region)

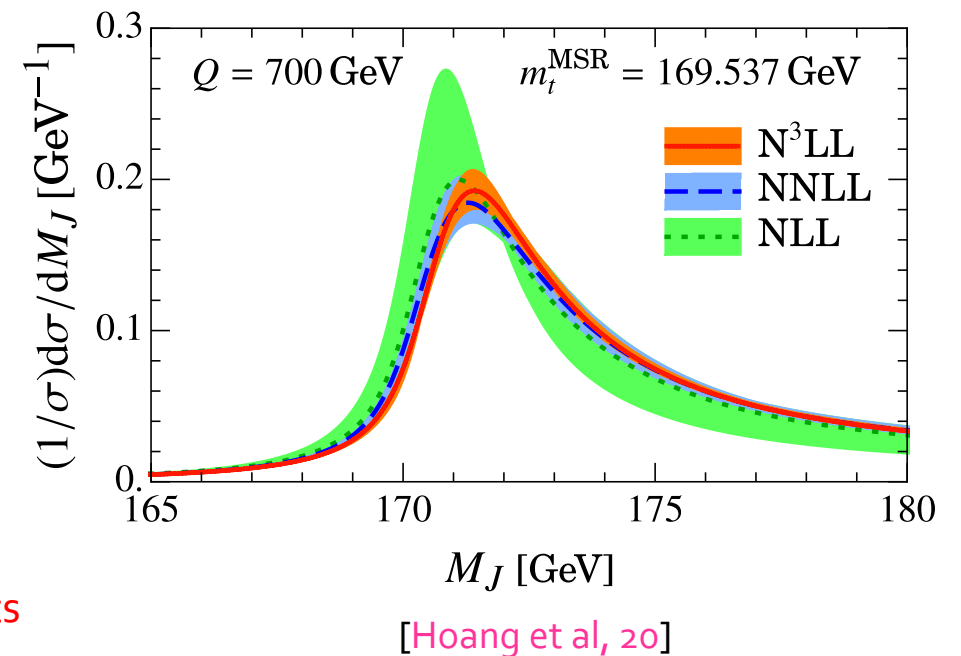
$$\frac{1}{\sigma_0} \frac{d\sigma}{d\tau_2} = H_Q(Q, \mu) \times H_M(\varrho_t, m_t, \mu)$$

$$\times J_{B,t}^{(\Gamma_t)}(Q^2\tau_2, m_t, \mu) \otimes \bar{J}_{B,\bar{t}}^{(\Gamma_{\bar{t}})}(Q^2\tau_2, m_t, \mu) \otimes S(Q\tau_2, \mu) + \mathcal{O}\left(\lambda, \varrho_t^{-1}, \frac{\Gamma_t}{m_t}\right)$$

	$F(\mu_F)$	$\Gamma_{\text{cusp}, \beta}$	γ
LL	0	1	0
NLL	0	2	1
NLL'	1	2	1
N ² LL	1	3	2
N ³ LL	2	4	3

Loop-orders of resummation ingredients

- State-of-the-art QCD:
 - NLL QCD + **LO EW** (width Γ_t in peak region) [Fleming et al, 07, 08]
 - 2-loop bHQET jet function [Stewart et al, 08]
 - 2-loop SCET-bHQET current matching [Pathak et al, 15]
 - **N³LL**-study (peak region) [Hoang et al, 20]
 - 2-loop SCET massive jet function [Lepenik et al, 18]
- QCD field-theoretic subtleties:
 - Non-perturbative corrections, Renormalons
 - Top mass schemes (MSR mass)
 - ...
- **EW** only **LO** so far
- Extend formalism: **systematically** incorporate **subleading EW effects**



Master formula (double-resonant case)

Formalism extension of
[Chiu et al, 07, 08, 09]

$$\frac{d\sigma_{e^+e^-t\bar{t}}(P_-, P_+)}{d\tau_2 d\Phi_2} = \sum_{\kappa, \rho} \frac{K_{\Phi}^{\kappa\rho}}{\mathcal{F}} \text{Tr}[\rho_{e^+e^-t\bar{t}} \mathcal{M}_M H_H \mathcal{M}_M^\dagger]^{\kappa\rho} * J_{B,t}^{(\Gamma_t)} \otimes \bar{J}_{B,t}^{(\Gamma_t)} \otimes B_{e^-}^{\kappa|P_-} \otimes B_{e^+}^{\kappa|P_+} \otimes S_{e^+e^-t\bar{t}}$$

Hard phase space

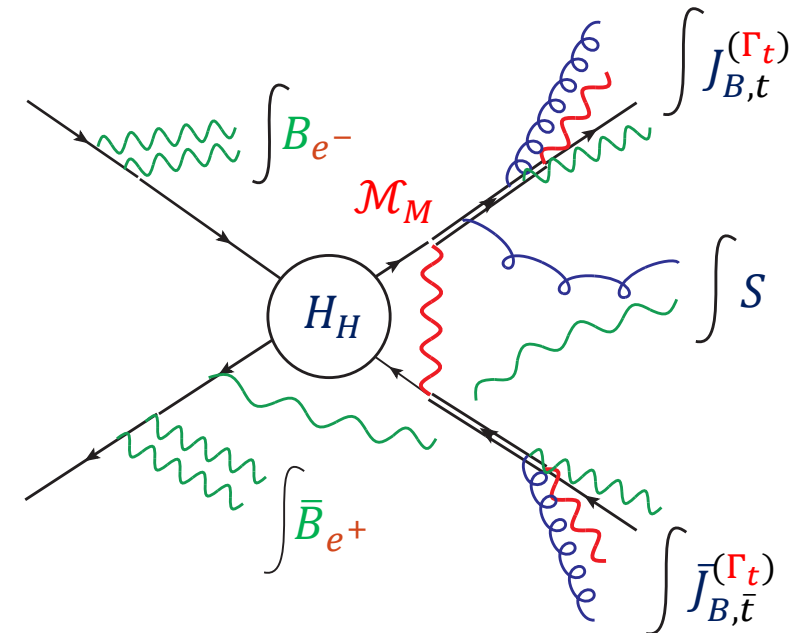
Beam polarizations, $P_i \in [-1, 1]$

Hard kinematics + flux

- Lepton chirality $\kappa = \pm 1$, quark chirality $\rho = \pm 1$
- Isospin density matrix (in isospin space):

$$\rho_{e^+e^-t\bar{t}} = \text{diag}(0, 0, 1, 0, 0, 0)$$

$$i \sim \begin{pmatrix} (\bar{\nu}\nu)(\bar{u}u) \\ (\bar{\nu}\nu)(\bar{d}d) \\ (\bar{e}e)(\bar{u}u) \\ (\bar{e}e)(\bar{d}d) \\ (\bar{e}\nu)(\bar{u}d) \\ (\bar{\nu}e)(\bar{d}u) \end{pmatrix}$$

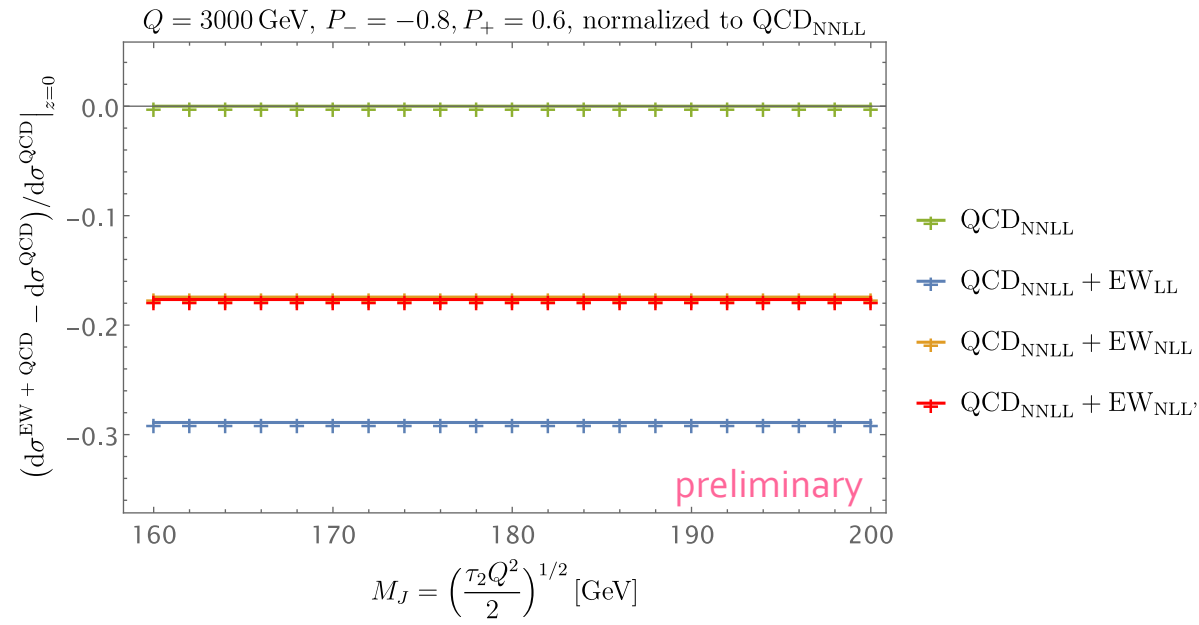
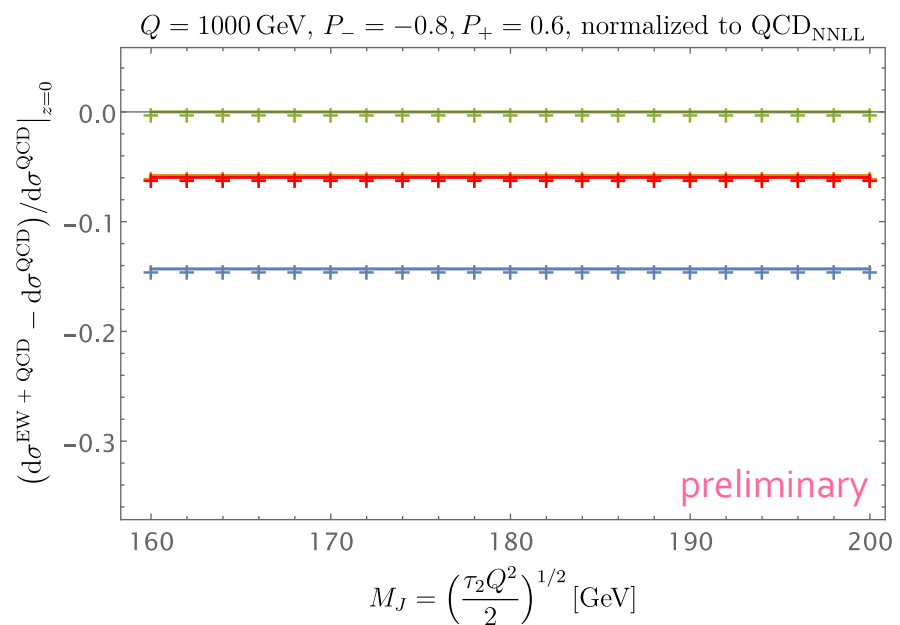
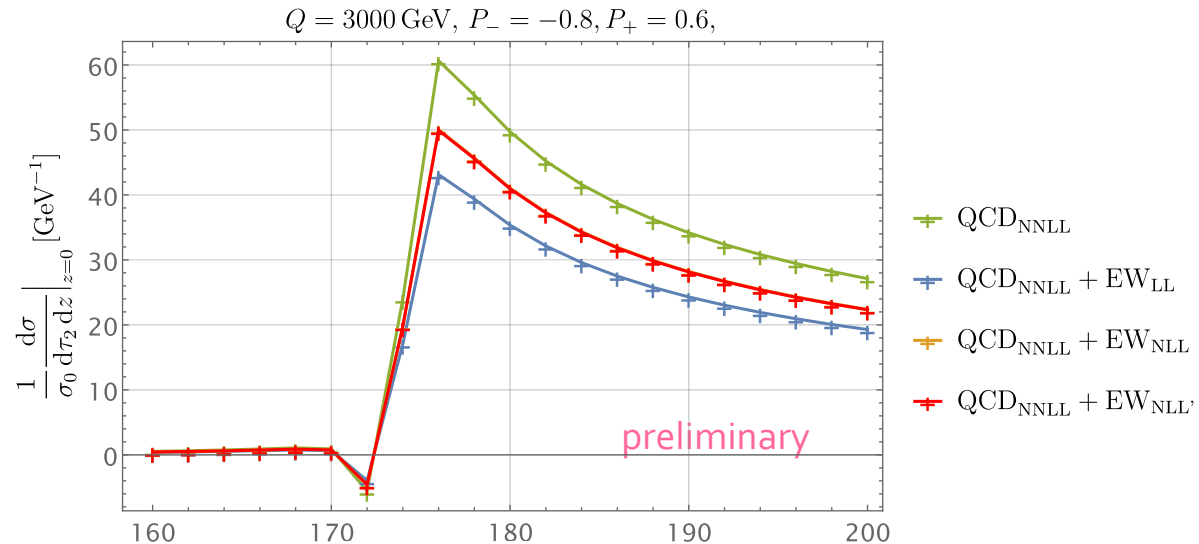
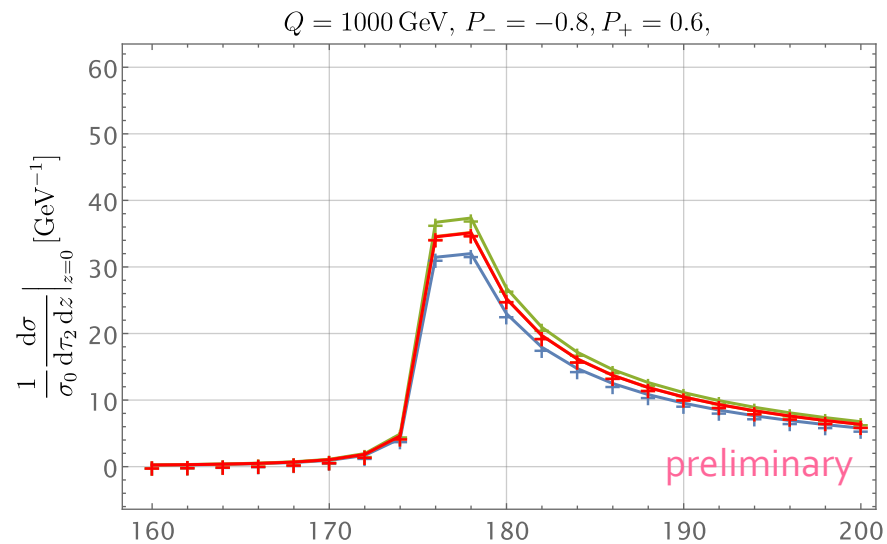


- Switching off QED below electroweak scale:

$$\frac{d\sigma_{e^+e^-t\bar{t}}(P_-, P_+)}{d\tau_2 dz} = \sum_{\kappa, \rho} \frac{K_{\Phi}^{\kappa\rho} \phi^{\kappa}(P_-, P_+)}{\mathcal{F}} \text{Tr}[\rho_{e^+e^-t\bar{t}} \mathcal{M}_M H_Q \mathcal{M}_M^\dagger]^{\kappa\rho} * J_{B,t}^{(\Gamma_t)} \otimes \bar{J}_{B,t}^{(\Gamma_t)} \otimes S_{e^+e^-t\bar{t}}$$

CM frame now: $z = \cos \theta_*$

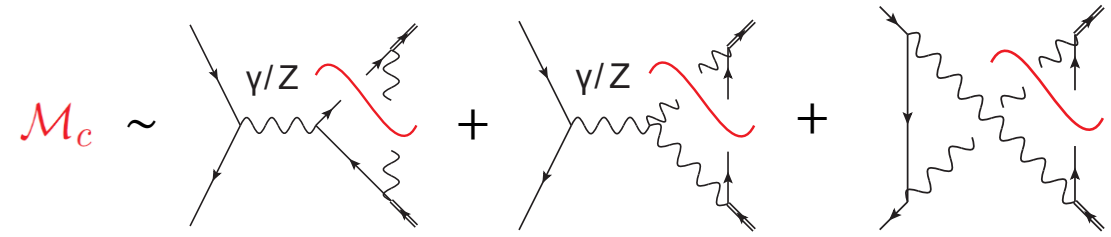
Beam polarization function



Single-resonant case

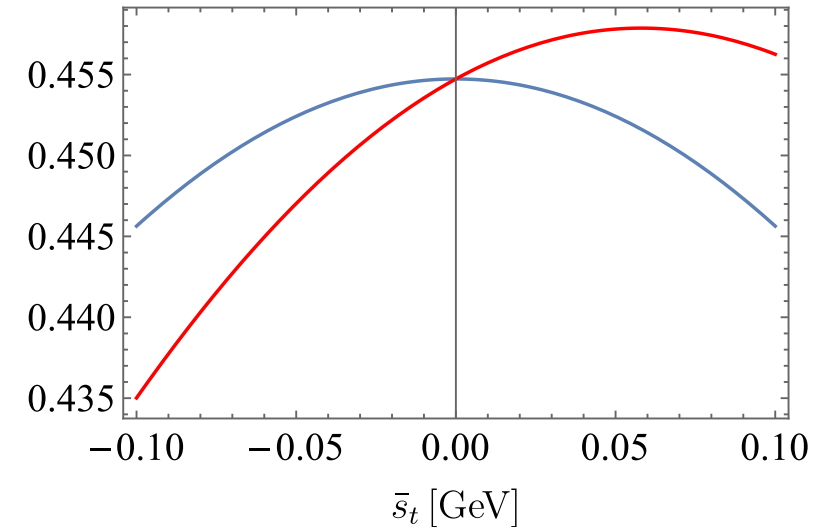
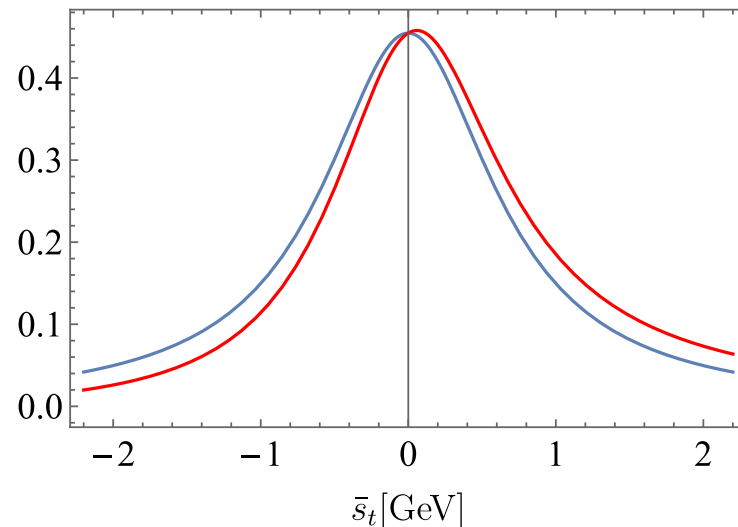
$$\frac{d\sigma_{e^+e^-t\bar{t}}(P_-, P_+)}{d\tau_2 dz} = \sum_{\kappa, \rho} \frac{K_{\Phi}^{\kappa\rho} \phi^{\kappa}(P_-, P_+)}{\mathcal{F}} \text{Tr}[\rho_{e^+e^-t\bar{t}} \mathcal{M}_c \mathcal{M}_s H_Q \mathcal{M}_s^\dagger \mathcal{M}_c^\dagger]^{\kappa\rho} * J_{B,t}^{(\Gamma_t)} \otimes \bar{J}_{B,t}^{(\Gamma_t)} \otimes S_{e^+e^-t\bar{t}}$$

- There are **imaginary parts** in \mathcal{M}_c due to ***bW*-cuts**
- **Resonant/non-resonant interference** contributions are obtained by keeping these imaginary parts



[Hoang, Reisser, 04]

- Leads to **peak shift** O(30-50 MeV)
- WIP: including **QCD logs**

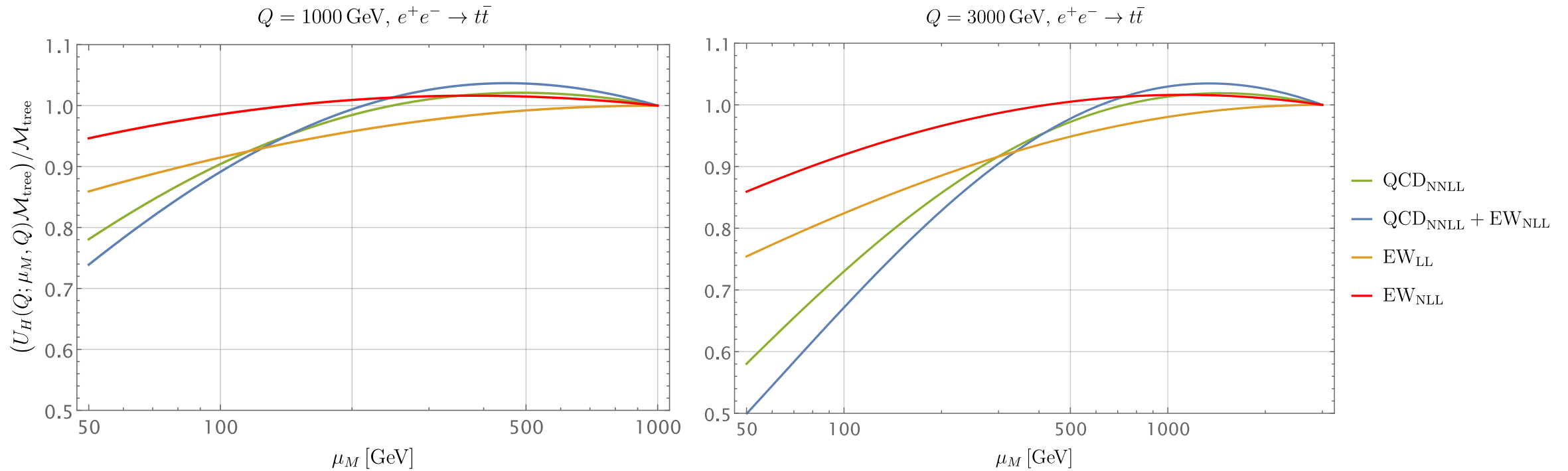


Summary and outlook

- Summary:
 - Extension of QCD factorization formalism to systematic inclusion of **subleading EW effects** for **inclusive boosted top jet pair-production**
 - Coherent resummation of QCD logs and EW Sudakov logarithms
 - 2-jettiness double-resonant cross section (**NNLL QCD** + **NLL' EW**) for linear collider environment (including beam polarization effects)
 - **Resonant/non-resonant interference**
- Outlook:
 - Coherent treatment of **resonant/non-resonant** interference and **QCD** effects
 - **QED** below electroweak scale
 - Realistic beam treatment (**QED! Weak?!**)
 - Extend factorization formalism to **more differential final-state**
 - Top decay-sensitive observables
 - Fragmentation
 - ...
 - Extension of the formalism to LHC observables

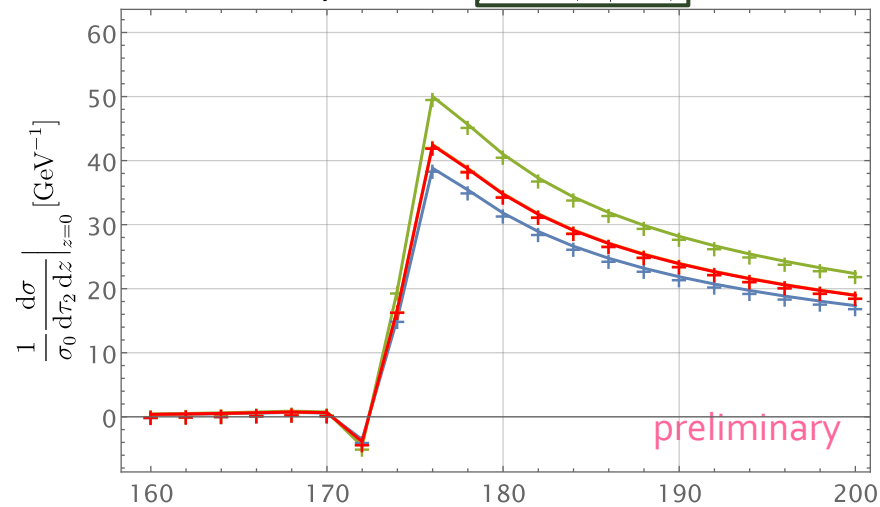
Backup

Sudakov evolution

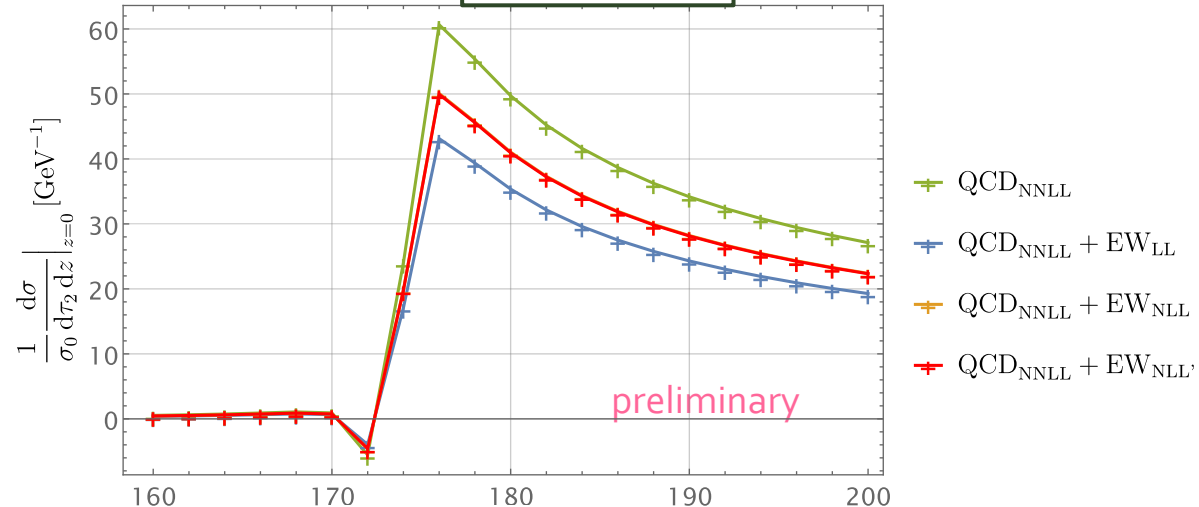


Chirality of operators shown: left-left (for right-left and left-right: SU(2)-parts roughly half the size)

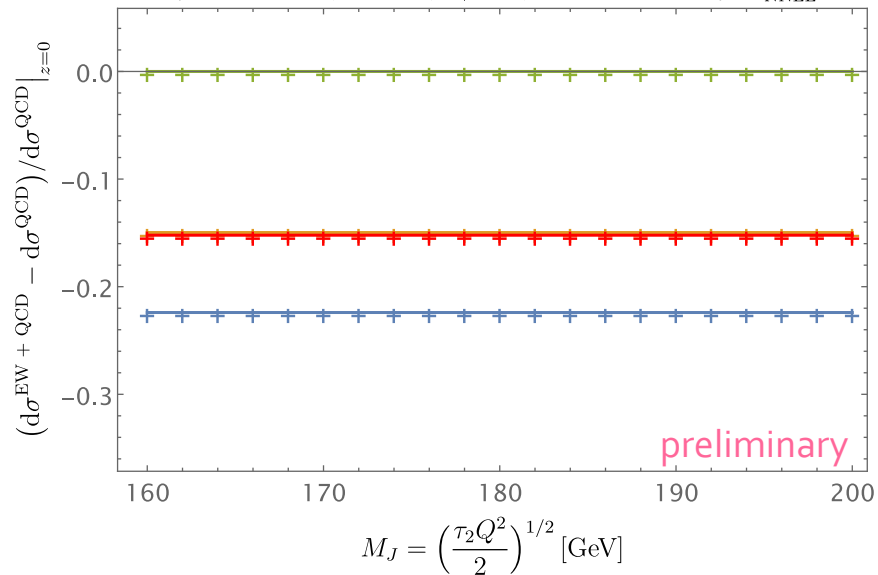
$Q = 3000 \text{ GeV}$, $P_- = 0, P_+ = 0$, unpolarized



$Q = 3000 \text{ GeV}$, $P_- = -0.8, P_+ = 0.6$, ILC polarization



$Q = 3000 \text{ GeV}$, $P_- = 0, P_+ = 0$, normalized to QCD_{NNLL}



$Q = 3000 \text{ GeV}$, $P_- = -0.8, P_+ = 0.6$, normalized to QCD_{NNLL}

