#### Search for critical point of strongly interacting matter

(Intermittency analysis by NA61/SHINE at CERN SPS)

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#### Outline

The goal of this study is to locate the critical point of the strongly interacting matter

measuring scaled factorial moments of multiplicity distribution  $\mbox{from a selection of}$ 

Pb+Pb at 13A GeV/ $c(\sqrt{s}_{NN} \approx 5.1$  GeV)

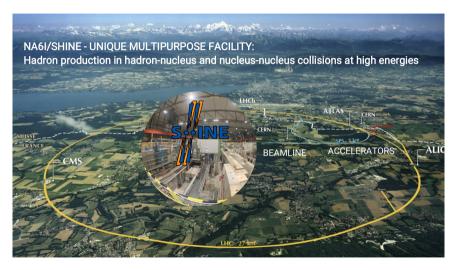
Pb+Pb at 30A GeV/ $c(\sqrt{s}_{NN}\approx 7.5 \text{ GeV})$ 

Ar+Sc at 150A GeV/ $c(\sqrt{s}_{NN}\approx 17 \text{ GeV})$ 

using statistically independent points and cumulative variables

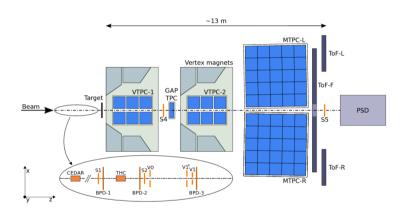
- Introduction
- QCD critical point search strategies
  - QCD critical point
  - Exploring the phase diagram with heavy-ion collisions
- Experimental measures to search QCD critical point
  - Fluctuations in large momentum bins
    - Multiplicity fluctuations
    - Multiplicity-transverse momentum fluctuations
  - Fluctuation as a function of momentum bin size
    - Proton intermittency anlysis
    - h<sup>-</sup> intermittency analysis
- Exclusion plot and intermittency analysis result
- Summary

# NA61/SHINE at CERN SPS



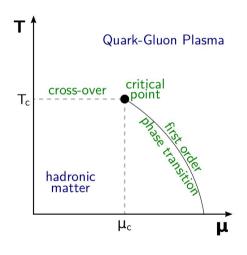
NA61/SHINE (SPS Heavy Ion and Neutrino Experiment) is a particle physics fixed-target experiment at CERN SPS

# NA61/SHINE detector



- VTPC-1, VTPC-2 are placed in the magnetic field
- TPC system: track reconstruction and particle identification based on specific energy loss
- Projectile Spectator Detector (PSD): hadronic calorimeter, measures projectile spectators energy

## Critical point of strngly interacting matter



- Critical Point (CP): a hypothetical end point of first order phase transition line (QGP-HM) that has properties of second order phase transition
- $\bullet$  Second order phase transition  $\to$  scale invariance  $\to$  power-law form of correlation function
- These expectations are for fluctuations and correlations in the configuration space
- They are expected to be projected to the momentum space via quantum effect and/or collective flow
- Predictions on the CP existence, its location and what and how should fluctuate are model-dependent

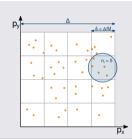
Asakawa, Yazaki NPA 504 (1989) 668

#### Scaled factorial moments of order r

Second order phase transition  $\to$  scale invariance  $\to$  power-law form of correlation function  $\to$  enhanced multiplicity fluctuations that can be revealed by scaled factorial moments of order r:

$$F_r(\delta) = rac{\left\langle rac{1}{M^2} \sum\limits_{i=1}^{M^2} n_i (n_i-1)...(n_i-r+1) 
ight
angle}{\left\langle rac{1}{M^2} \sum\limits_{i=1}^{M^2} n_i 
ight
angle^r}$$

 $\delta$ : size of each of the M sub-division intervals of  $\Delta$   $n_i$ : number of particles in i-th bin



When the system is a simple fractal and  $F_r(\delta)$  follows a power law dependences:

$$F_r(\delta) = F_r(\Delta) \cdot (\frac{\Delta}{\delta})^{D \cdot \phi_r}$$

Additionally, the exponent (intermittency index)  $\phi_r$  obeys the relation:

$$D \cdot \phi_r = (r-1) \cdot d_r$$

Where the anomalous fractal dimension  $d_r$  is independent of r

Wosiek, APPB 19 (1988) 863 Bialas, Hwa, PLB 253 (1991) 436 Bialas, Peschanski, NPB 273(1986) 703 Antoniou. Diakonos. Kapovannis. Kousouris. PRL 97 (2006) 032002

# Intermittency analysis

In NA61/SHINE, intermittency analysis is performed at mid-rapidity and particle fluctuations are studied in transverse momentum plane to locate the critical point of the strongly interacting matter measuring scaled factorial moments of multiplicity distribution

#### modified equivalent formula

$$F_r(M) = \frac{r!(M^2)^{r-1}}{\langle N \rangle^r} \left\langle \sum_{m=1}^{M^2} {n_m \choose r} \right\rangle$$

for 
$$r = 2$$

$$F_2(M) = rac{2M^2}{\langle N 
angle^2} \, \langle N_{pp} 
angle$$

M: number of bins in  $p_x$  and  $p_y$ 

N: event multiplicity

 $n_i$ : numbers of particles in ith bin

< .... >: averaging over events

 $N_{pp}$ : total number of pairs in  $M^2$  bins in an event

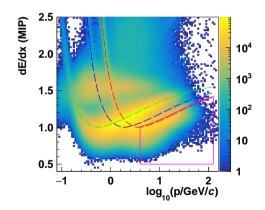
#### Error propagation:

$$\frac{\sigma_{F_r}}{|F_r|} = \sqrt{\frac{\textit{Var}(\textit{N}_r)}{\langle \textit{N}_r \rangle^2} + r^2 \frac{\textit{Var}(\textit{N})}{\langle \textit{N} \rangle^2} - 2r \frac{\textit{Cov}(\textit{N}_r,\textit{N})}{\langle \textit{N} \rangle \langle \textit{N}_r \rangle}}$$

#### Particle identification

#### Example: proton candidates selection

( Pb+Pb at  $30A \ GeV/c$  )



----: theoretical BB function for protons ----: theoretical BB function for kaons ----: theoretical BB function for pions

selection of protons is based on dEdx measurements in TPCs

- $0.60 < log_{10}(p/GeV/c) < 2.10$
- $0.5 \le dEdx \le BB_{proton} + 0.15(BB_{kaon} BB_{proton})$

around 60% protons are selected with few percent kaon contamination

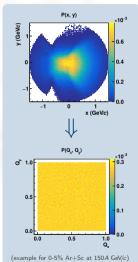
BB: theoretical Bethe-Bloch function

#### **Cumulative variables**

Instead of using  $p_x$  and  $p_y$ , one can use cumulative quantities:

$$Q_x = \int\limits_{x_{min}}^{x} 
ho(x) dx / \int\limits_{x_{min}}^{x_{max}} 
ho(x) dx \qquad Q_y = \int\limits_{y_{min}}^{y} P(x, y) dy / P(x)$$

- transform any distribution into uniform distribution (0,1)
- remove the dependence of F<sub>r</sub> on the shape of the single-particle distribution
- intermittency index of an ideal power-law correlation function system described in two dimensions in momentum space was proven to remain approximately invariant after the transformation



Bialas, Gazdzicki, PLB 252 (1990) 483

#### Preliminary result on intermittency analysis

Pb+Pb at 13A GeV/ $c(\sqrt{s_{NN}}\approx 5.1~{\rm GeV})$ 

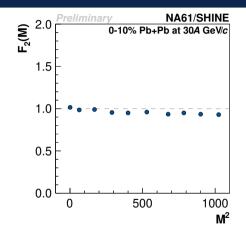
Pb+Pb at 30A GeV/c( $\sqrt{s_{NN}} pprox$  7.5 GeV)

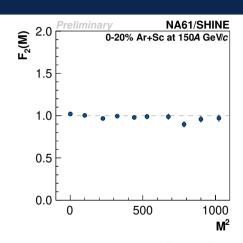
Ar+Sc at 150A GeV/ $c(\sqrt{s_{\it NN}}\approx$  17 GeV)

Cumulative variables Independent sub-sample for each M points

# Proton intermittency analysis result

( Experimental result on fluctuations as a function of momentum bin size )





CPOD 2021: https://indico.cern.ch/event/985460

ICNFP 2021: https://indico.cern.ch/event/1025480/

ISNP ERICE 2021: http://crunch.ikp.physik.tu-darmstadt.de/erice/2021

No indication for power-law increase with bin size

## Simple power-law model

A simple model that generates momentum of particles for a given number of events with a given multiplicity distribution.

It has two main parameters:

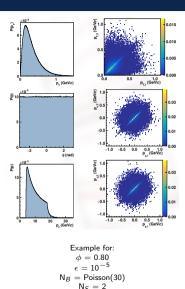
- ratio of correlated to uncorrelated particles
- power-law exponent

#### Uncorrelated particles (background)

$$\rho_{\rm B}({\rm p_T}) = {\rm p_T \cdot e^{-6p_T}}$$

#### Correlated pairs (signal)

$$\begin{split} \rho_{S}(\mathbf{p}_{\mathsf{T},1},\mathbf{p}_{\mathsf{T},2}) &= \rho_{B}(\mathbf{p}_{\mathsf{T},1}) \cdot \rho_{B}(\mathbf{p}_{\mathsf{T},2}) \\ & \cdot \left[ \left| \Delta \mathbf{p}_{\mathsf{x}} \right|^{\phi} + \epsilon \right]^{-1} \cdot \left[ \left| \Delta \mathbf{p}_{\mathsf{y}} \right|^{\phi} + \epsilon \right]^{-1} \end{split}$$



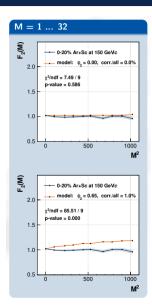
## Simple power-law model

Lots of model data sets are generated:

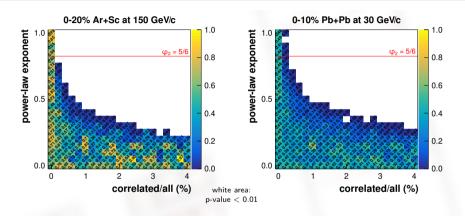
- correlated-to-all ratio: vary from 0.0 to 4.0% (with 0.2 steps)
- power-law-exponent: vary from 0.0 to 1.0 (with 0.05 steps)

and compared with the experimental data

For the construction of exclusion plots, statistical uncertainties were calculated using model with statistics corresponding to the data.



## **Exclusion plot**



exclusion plots for parameters of simple power-law model

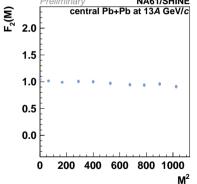
The intermittency index  $\varphi_2$  for a system freezing out at the QCD critical endpoint is expected to be  $\varphi_2=5/6$  assuming that the latter belongs to the 3-D Ising universality class.

# New proton intermittency analysis result

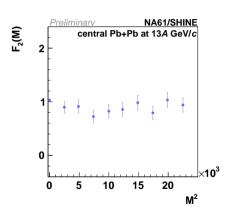
( using Pb+Pb at 13A GeV/c )

small M range (1,8,...32):

Preliminary NA61/SHINE



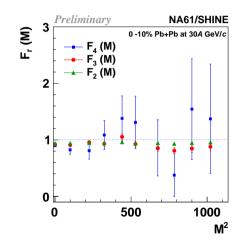
full M range (1,50,...,150):

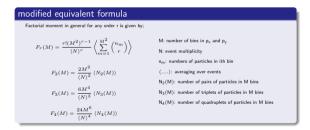


SSFMs are calculated using cumulative variable and independent sub-sample are used for each M points

## Negatively charged hadron intermittency analysis

( higher order moments in Pb+Pb at 30A GeV/c data)



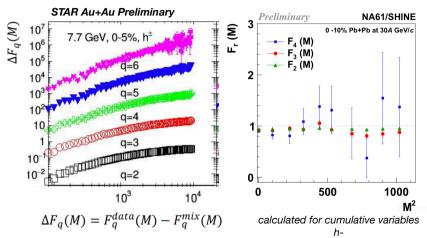


relationship between intermittency indices:

$$\phi_2 = \frac{\phi_3}{2} = \frac{\phi_4}{3}$$

No indication for power-law increase with bin size

#### STAR vs NA61/SHINE preliminary results on ΔFq(M)



\*\*This seems to be in tension with corresponding results by the STAR Collaboration\*\*

### Summary

- Results on the dependence of scaled factorial moments of multiplicity distribution on cumulative momentum bin size for:
  - protons in Pb+Pb at 30A GeV/c( $\sqrt{s_{NN}} \approx 7.5$  GeV)
  - protons in Pb+Pb at 13A GeV/ $c(\sqrt{s_{NN}} \approx 5.1 \text{ GeV})$
  - protons in Ar+Sc at 150A GeV/ $c(\sqrt{s_{NN}}\approx 17 \text{ GeV})$
  - negatively charged hadrons in Pb+Pb at 30A GeV/ $c(\sqrt{s_{NN}} \approx 7.5$  GeV)

show no indication of a power-law increase

- Exclusion plots for parameters of a simple model (ratio of correlated to background particles and power-law exponent) are shown
- We are continuing proton and negatively charged hadron intermittency analysis for Pb+Pb at 13A GeV/c and Ar+Sc at 13A - 75A GeV/c data

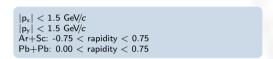
# **Acknowledgement**

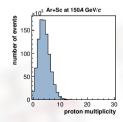
My work is supported by the National Science Centre, Poland under grant no. 2018/30/A/ST2/0026 & NA61/SHINE SPS CERN Collaboration

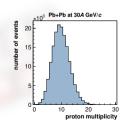


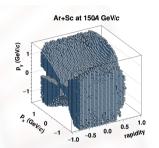
Back-ups!

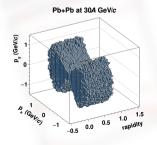
## Analysis acceptance











### mTTD cut

- Time Projection Chambers (TPCs) do not allow to reconstruct tracks too close to each other
- To reject those close tracks momentum based two track distance cut introduced for Pb+Pb at 30A GeV/c and Ar+Sc at 150A GeV/c data

#### Momentum based two-tracks distance cut

Coordinates of momentum will be used:

$$s_{x} = \frac{p_{x}}{p_{xz}} = cos(\Psi)$$

$$s_{y} = \frac{p_{y}}{p_{xz}} = sin(\lambda)$$

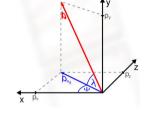
$$\rho = \frac{1}{p_{xz}}$$

For pairs:

$$\Delta s_x = s_{x_2} - s_{x_1}$$
  

$$\Delta s_y = s_{y_2} - s_{y_1}$$
  

$$\Delta \rho = \rho_2 - \rho_1$$



#### Effect of mTTD cut:

