# Corrections to the hadron resonance gas from lattice QCD and their effect on fluctuation-ratios at finite density

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#### in collaboration with

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# The Hadron Resonance Gas (HRG) model

### Importance:

- ▶ parameters of chemical freeze-out in experiments,
- non-critical baseline.

**HRG**: interacting gas of hadrons  $\cong$  non-interacting gas of hadrons and resonances.

$$\frac{p}{T^4} = \frac{1}{T^4} \sum_{\mathbf{h}} p_{\mathbf{h}} = \frac{1}{VT^3} \sum_{\mathbf{h}} \log \mathcal{Z}_{\mathbf{h}} \big( T, \pmb{\mu} = (\mu_B, \mu_Q, \mu_S) \big)$$

with

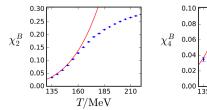
$$\begin{split} \log \mathcal{Z}_{\rm h} &= \mp \frac{d_{\rm h}}{2\pi^2} \frac{V}{T^3} \int\limits_0^\infty {\rm d}p \; p^2 \log \left[ 1 \mp z_{\rm h} \exp \left( -\frac{\sqrt{m_{\rm h}^2 + p^2}}{T} \right) \right] \\ &= V T m_{\rm h}^2 \frac{d_{\rm h}}{2\pi^2} \sum_{n=1}^\infty \frac{(\pm 1)^{n+1}}{n^2} z_{\rm h}^n K_2 \left( \frac{n m_{\rm h}}{T} \right) \end{split}$$

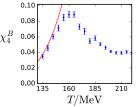
and fugacity factors:  $z_h = \exp[\beta(B_h\mu_B + Q_h\mu_Q + S_h\mu_S)].$ 

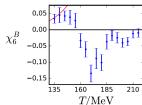


# Discrepancies of HRG and the lattice

Mostly good match for  $\mu = (0, 0, 0)$  in the hadronic phase:







[WB: 1805.04445] [Pisa: 1611.08285] [HotQCD: 2001.08530]

#### Possible reasons for the deviation:

- incomplete resonance list,
- finite widths are missing,
- lack of non-resonant and repulsive interactions,
- restoration of chiral symmetry.

### Generalised susceptibilities:

$$\chi^{BQS}_{lmn} = \frac{\partial^{l+m+n}(p/T^4)}{\partial \hat{\mu}^l_B \partial \hat{\mu}^m_Q \partial \hat{\mu}^n_S}$$

with  $\hat{\mu}_i = \mu_i/T$ .

# Fugacity expansion of the QCD free energy

#### Fugacity expansion:

$$\frac{p(T, \hat{\mu}_B, \hat{\mu}_S)}{T^4} = \sum_{j,k} P_{jk}^{BS}(T) \cosh(j\hat{\mu}_B - k\hat{\mu}_S)$$

with  $\mu_Q = 0$ .

## Sector coefficients: $P_{ik}^{BS}$ ,

- $\triangleright$  contributions from Hilbert subspaces with fixed B = j, S = k;
- e.g. (at  $\mu = 0$ )

$$\chi_4^B = \sum_k P_{1k}^{BS} + \sum_k 32 P_{2k}^{BS} + \sum_k 81 P_{3k}^{BS} + \dots$$

(B,S) $  $ $($	$(0,0) \mid (1,0)$	(0,1)	(1,1)	(1,2)	(2,0)	(0,2)	(2,1)
hadrons $   \pi$	$,\eta, ho \mid p,\Delta$	K	$\Lambda, \Sigma$	Ξ	p-p	K-K	$p-\Lambda, p-p-K$

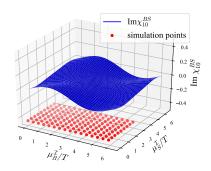
- $\blacktriangleright$  (main) contributions to ideal-HRG from  $P_{00}^{BS}, P_{01}^{BS}, P_{10}^{BS}, P_{11}^{BS}, P_{12}^{BS}$  and  $P_{13}^{BS},$
- ▶ for  $B \ge 2$  coefficients are close to zero  $\sim$  deviation!

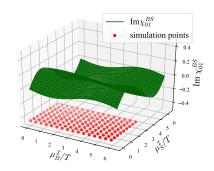


### Lattice setup

- ▶  $N_{\rm f} = 2 + 1$ ; 4stout-improved staggered action. [WB: 1507.04627]
- $\mu = i\mu^{\mathcal{I}} \sim \cosh(ix) = \cos(x) \sim \text{Fourier series}.$
- ▶  $N_{\tau} = 8, 10, 12; LT \approx 3; T = 145 \text{ MeV}, 150 \text{ MeV}, 155 \text{ MeV}$ and 160 MeV.

$$\mathrm{Im}\chi_{10}^{BS} = \sum_{j,k} j P_{jk}^{BS}(T) \sin(j\hat{\mu}_B^{\mathcal{I}} - k\hat{\mu}_S^{\mathcal{I}}) \qquad \quad \mathrm{Im}\chi_{01}^{BS} = \sum_{j,k} (-k) P_{jk}^{BS}(T) \sin(j\hat{\mu}_B^{\mathcal{I}} - k\hat{\mu}_S^{\mathcal{I}})$$

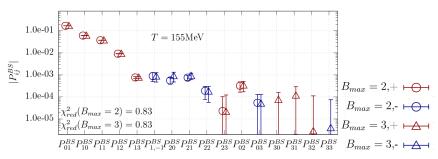




Example: T = 155 MeV.

### Sectors

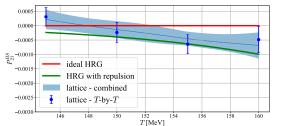
$$N_{\tau} = 12$$
:

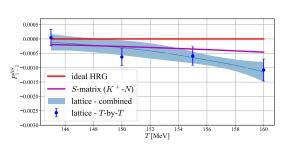


- $\qquad \qquad \text{Hierarchy: } P_{01}^{BS} > P_{02}^{BS} > P_{03}^{BS}; \, P_{10}^{BS} > P_{20}^{BS} > P_{30}^{BS}; \, P_{10}^{BS} > P_{12}^{BS} > P_{13}^{BS}; \, \text{etc.}$
- ▶ B = 2 sector:  $P_{2k}^{BS} < 0$ , and heavily underestimated in ideal HRG.

Systematic errors  $\sim B_{\text{max}} = 2 \text{ or } 3.$ 

# Subleading sectors – examples





# $P_{21}^{BS}$

- more negative for higher T,
- ▶ e.g. N- $\Lambda$  or N- $\Sigma$  interactions,
- ▶ ideal HRG deviates,
- repulsive mean field HRG,
- excluded volume HRG. [1708.02852]

# $P_{1-1}^{BS}$

- ▶ e.g. N-K<sup>+</sup> scattering,
- ► S-matrix formalism applicable. [1806.02177]

# Fluctuation ratios at finite density

Extrapolation to finite  $\mu_B$  using truncated

$$\frac{p}{T^4} = \sum_{j,k} P_{jk}^{BS}(T) \cosh \left[ j \hat{\mu}_B - k \hat{\mu}_S(\hat{\mu}_B) \right]$$

with strangeness neutrality:

$$\chi_1^S = \sum_{j,k} (-k) P_{jk}^{BS}(T) \sinh(j\hat{\mu}_B - k\hat{\mu}_S) \stackrel{!}{=} 0.$$

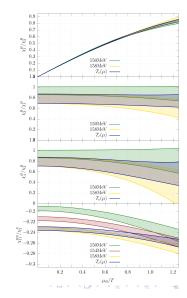
#### Ratios:

- $\chi_1^B/\chi_2^B \sim \text{proxy for } \mu_B$ ,
- $ightharpoonup \chi_3^B/\chi_1^B,\,\chi_4^B/\chi_2^B\sim {
  m baryon~thermometer},$
- $\triangleright \chi_{11}^{BS}/\chi_2^S \sim \text{strangeness thermometer}$

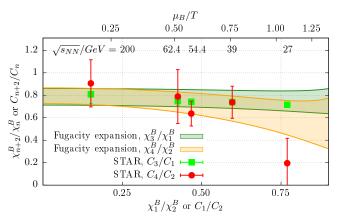
and

experimental proxy:  $\sigma_{\Lambda}^2/(\sigma_{\Lambda}^2+\sigma_K^2)$ . [WB: 1910.14592]

Crossover line:  $T_{\rm c}(\mu_B) \approx T_{\rm c}^0(1 - \kappa_2 \hat{\mu}_B^2)$ [WB: 2002.02821]



## Comparison with experiment



- Consistent with: [HotQCD: 2001.08530], [WB: 1805.0444].
- ▶ Using  $C_i$  net-proton cumulants from STAR Experiment.
- ▶ Assuming that the crossover and chemical freeze-out lines are close to each other.
- Ignoring ceaveats of comparison.

## Summary and conclusion

- ▶ Scanning of the QCD free energy in imaginary  $\mu_B$ - $\mu_S$  plane  $\sim$  separation of sectors.
- ▶ Possible separation of processes like *K-K* or *p-p* scattering.
- ightharpoonup Continuum estimates for  $P_{ik}^{BS}$  sector coefficients and fluctuation ratios.
- $\triangleright$  Consistency with STAR data for small  $\mu_B$ .
- ▶ Deviation of  $\chi_4^B/\chi_2^B$  from ideal HRG is due to B=2 effects.
- Our lattice results could provide help in
  - 1. phenomenology,
  - construction of more realistic models.

### Thank you for your attention.

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