# Corrections to the hadron resonance gas from lattice QCD and their effect on fluctuation-ratios at finite density 

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## The Hadron Resonance Gas (HRG) model

## Importance:

- parameters of chemical freeze-out in experiments,
- non-critical baseline.

HRG: interacting gas of hadrons $\cong$ non-interacting gas of hadrons and resonances.

$$
\frac{p}{T^{4}}=\frac{1}{T^{4}} \sum_{\mathrm{h}} p_{\mathrm{h}}=\frac{1}{V T^{3}} \sum_{\mathrm{h}} \log \mathcal{Z}_{\mathrm{h}}\left(T, \boldsymbol{\mu}=\left(\mu_{B}, \mu_{Q}, \mu_{S}\right)\right)
$$

with

$$
\begin{aligned}
\log \mathcal{Z}_{\mathrm{h}} & =\mp \frac{d_{\mathrm{h}}}{2 \pi^{2}} \frac{V}{T^{3}} \int_{0}^{\infty} \mathrm{d} p p^{2} \log \left[1 \mp z_{\mathrm{h}} \exp \left(-\frac{\sqrt{m_{\mathrm{h}}^{2}+p^{2}}}{T}\right)\right] \\
& =V T m_{\mathrm{h}}^{2} \frac{d_{\mathrm{h}}}{2 \pi^{2}} \sum_{n=1}^{\infty} \frac{( \pm 1)^{n+1}}{n^{2}} z_{\mathrm{h}}^{n} K_{2}\left(\frac{n m_{\mathrm{h}}}{T}\right)
\end{aligned}
$$

and fugacity factors: $z_{\mathrm{h}}=\exp \left[\beta\left(B_{\mathrm{h}} \mu_{B}+Q_{\mathrm{h}} \mu_{Q}+S_{\mathrm{h}} \mu_{S}\right)\right]$.

## Discrepancies of HRG and the lattice

Mostly good match for $\boldsymbol{\mu}=(0,0,0)$ in the hadronic phase:



[WB: 1805.04445]
[Pisa: 1611.08285]
[HotQCD: 2001.08530]

## Generalised susceptibilities:

Possible reasons for the deviation:

- incomplete resonance list,
- finite widths are missing,

$$
\chi_{l m n}^{B Q S}=\frac{\partial^{l+m+n}\left(p / T^{4}\right)}{\partial \hat{\mu}_{B}^{l} \partial \hat{\mu}_{Q}^{m} \partial \hat{\mu}_{S}^{n}},
$$

- lack of non-resonant and repulsive with $\hat{\mu}_{i}=\mu_{i} / T$. interactions,
- restoration of chiral symmetry.


## Fugacity expansion of the QCD free energy

## Fugacity expansion:

$$
\frac{p\left(T, \hat{\mu}_{B}, \hat{\mu}_{S}\right)}{T^{4}}=\sum_{j, k} P_{j k}^{B S}(T) \cosh \left(j \hat{\mu}_{B}-k \hat{\mu}_{S}\right)
$$

with $\mu_{Q}=0$.

Sector coefficients: $P_{j k}^{B S}$,

- contributions from Hilbert subspaces with fixed $B=j, S=k$;
- e.g. (at $\boldsymbol{\mu}=0$ )

$$
\chi_{4}^{B}=\sum_{k} P_{1 k}^{B S}+\sum_{k} 32 P_{2 k}^{B S}+\sum_{k} 81 P_{3 k}^{B S}+\ldots
$$

| $(B, S)$ | $(0,0)$ | $(1,0)$ | $(0,1)$ | $(1,1)$ | $(1,2)$ | $(2,0)$ | $(0,2)$ | $(2,1)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| hadrons | $\pi, \eta, \rho$ | $p, \Delta$ | $K$ | $\Lambda, \Sigma$ | $\Xi$ | $p-p$ | $K-K$ | $p-\Lambda, p-p-K$ |

- (main) contributions to ideal-HRG from $P_{00}^{B S}, P_{01}^{B S}, P_{10}^{B S}, P_{11}^{B S}, P_{12}^{B S}$ and $P_{13}^{B S}$,
- for $B \geq 2$ coefficients are close to zero $\sim$ deviation!


## Lattice setup

- $N_{\mathrm{f}}=2+1$; 4stout-improved staggered action. [wB: 1507.04627]
- $\boldsymbol{\mu}=i \boldsymbol{\mu}^{\mathcal{I}} \sim \cosh (i x)=\cos (x) \sim$ Fourier series.
- $N_{\tau}=8,10,12 ; L T \approx 3 ; T=145 \mathrm{MeV}, 150 \mathrm{MeV}, 155 \mathrm{MeV}$ and 160 MeV .
$\operatorname{Im} \chi_{10}^{B S}=\sum_{j, k} j P_{j k}^{B S}(T) \sin \left(j \hat{\mu}_{B}^{\mathcal{I}}-k \hat{\mu}_{S}^{\mathcal{I}}\right)$

$$
\operatorname{Im} \chi_{01}^{B S}=\sum_{j, k}(-k) P_{j k}^{B S}(T) \sin \left(j \hat{\mu}_{B}^{\mathcal{I}}-k \hat{\mu}_{S}^{\mathcal{I}}\right)
$$




Example: $T=155 \mathrm{MeV}$.

## Sectors

$$
N_{\tau}=12
$$



- Hierarchy: $P_{01}^{B S}>P_{02}^{B S}>P_{03}^{B S} ; P_{10}^{B S}>P_{20}^{B S}>P_{30}^{B S} ; P_{10}^{B S}>P_{12}^{B S}>P_{13}^{B S}$; etc.
- $B=2$ sector: $P_{2 k}^{B S}<0$, and heavily underestimated in ideal HRG.

Systematic errors $\sim B_{\max }=2$ or 3 .

## Subleading sectors - examples



## $P_{21}^{B S}$

- more negative for higher $T$,
- e.g. $N-\Lambda$ or $N-\Sigma$ interactions,
- ideal HRG deviates,
- repulsive mean field HRG, [1708.00879]
- excluded volume HRG. [1708.02852]

$P_{1-1}^{B S}$
- e.g. $N-K^{+}$scattering,
- $S$-matrix formalism applicable. [1806.02177]


## Fluctuation ratios at finite density

Extrapolation to finite $\mu_{B}$ using truncated

$$
\frac{p}{T^{4}}=\sum_{j, k} P_{j k}^{B S}(T) \cosh \left[j \hat{\mu}_{B}-k \hat{\mu}_{S}\left(\hat{\mu}_{B}\right)\right]
$$

with strangeness neutrality:
$\chi_{1}^{S}=\sum_{j, k}(-k) P_{j k}^{B S}(T) \sinh \left(j \hat{\mu}_{B}-k \hat{\mu}_{S}\right) \stackrel{!}{=} 0$.

## Ratios:

- $\chi_{1}^{B} / \chi_{2}^{B} \sim$ proxy for $\mu_{B}$,
- $\chi_{3}^{B} / \chi_{1}^{B}, \chi_{4}^{B} / \chi_{2}^{B} \sim$ baryon thermometer,
- $\chi_{11}^{B S} / \chi_{2}^{S} \sim$ strangeness thermometer
and
experimental proxy: $\sigma_{\Lambda}^{2} /\left(\sigma_{\Lambda}^{2}+\sigma_{K}^{2}\right)$.
[WB: 1910.14592]

Crossover line: $T_{\mathrm{c}}\left(\mu_{B}\right) \approx T_{\mathrm{c}}^{0}\left(1-\kappa_{2} \hat{\mu}_{B}^{2}\right)$ [WB: 2002.02821]


## Comparison with experiment


$\rightarrow$ Consistent with: [HotQCD: 2001.08530], [WB: 1805.0444].

- Using $C_{i}$ net-proton cumulants from STAR Experiment.
$\rightarrow$ Assuming that the crossover and chemical freeze-out lines are close to each other.
- Ignoring ceaveats of comparison.


## Summary and conclusion

- Scanning of the QCD free energy in imaginary $\mu_{B}-\mu_{S}$ plane $\sim$ separation of sectors.
- Possible separation of processes like $K-K$ or $p-p$ scattering.
- Continuum estimates for $P_{j k}^{B S}$ sector coefficients and fluctuation ratios.
- Consistency with STAR data for small $\mu_{B}$.
- Deviation of $\chi_{4}^{B} / \chi_{2}^{B}$ from ideal HRG is due to $B=2$ effects.
- Our lattice results could provide help in

1. phenomenology,
2. construction of more realistic models.

## Thank you for your attention.

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