CP violation in D decays to two pseudoscalars: A SM-based calculation

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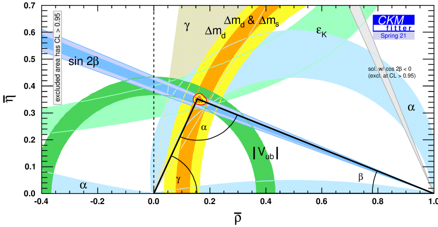
CP violation in D decays: just a SM system or gateway to New Physics?

$$\Delta A_{CP}^{exp} \equiv A_{CP}(D^0 \to K^+K^-) - A_{CP}(D^0 \to \pi^+\pi^-) = (-1.54 \pm 0.29) \cdot 10^{-3}$$

$$\Delta A_{CP}^{dir,exp} = (-1.57 \pm 0.29) \cdot 10^{-3}$$ [LHCb 2019]

- Is the SM theoretical prediction in agreement?
- Weak sector (CKM parameters) already probed by kaons, B mesons



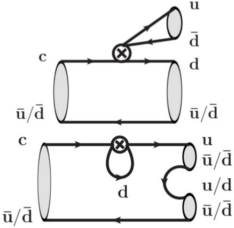
- Strong sector (hadronic uncertainties) problematic

CPV in D: the strong sector

- Does a beyond-naive estimation of hadronic effects matter?

$$\mathscr{A}(D^0 \to f) = A(f) + i r_{CKM} B(f)$$
$$\mathscr{A}(\overline{D^0} \to f) = A(f) - i r_{CKM} B(f)$$
$$a_{CP}^{dir} \approx 2 r_{CKM} \frac{|B(f)|}{|A(f)|} \cdot \sin \arg \frac{A(f)}{B(f)}$$

$$(r_{CKM} = Im \frac{V_{cb}^* V_{ub}}{V_{cd}^* V_{ud}})$$



Tree topology

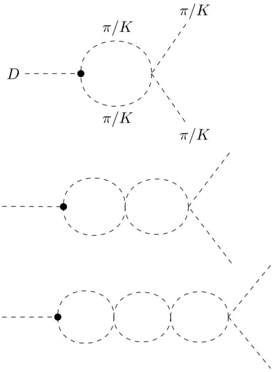
Penguin topology

[Other approaches: 1506.04121, 1706.07780]

In K decays: Chiral Perturbation Theory. In B decays: HQET.
$\Lambda_{\chi PT} \approx m_\rho < m_D = 1865$ MeV, $\frac{\Lambda_{QCD}}{m_c} = \mathscr{O}(1)$
$\to$ neither approach would be valid in charm!

A way to look at the problem: rescattering

- Strong process, blind to A or B
- Isospin (u$\leftrightarrow$d) is a good symmetry of strong interactions
- In I=0, two channels:



$$S_{strong} = \begin{pmatrix} \pi\pi \to \pi\pi & \pi\pi \to KK \\ KK \to \pi\pi & KK \to KK \end{pmatrix}$$

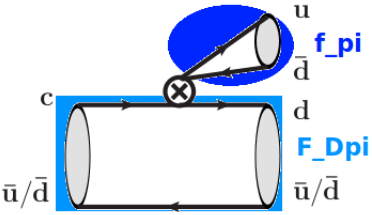
Rescattering & what we learn about strong phases

- S matrix is **unitary**, as well as strong sub-matrix
- For I=0: $\begin{pmatrix} A(D\to\pi\pi) \\ A(D\to KK) \end{pmatrix} = \underbrace{\begin{pmatrix} \eta e^{i2\delta_1} & i\sqrt{1-\eta^2}e^{i(\delta_1+\delta_2)} \\ i\sqrt{1-\eta^2}e^{i(\delta_1+\delta_2)} & \eta e^{i2\delta_2} \end{pmatrix}}_{S_{strong}} \cdot \begin{pmatrix} A^*(D\to\pi\pi) \\ A^*(D\to KK) \end{pmatrix}$
- The phases are related to the rescattering phases **which are known from data/other experiments**
- Watson's theorem (elastic rescattering limit): $argA(D\to\pi\pi) = argA(\pi\pi\to\pi\pi) mod \pi$
- With inelasticities: more complicated, phase-shifts dependent on magnitudes of the amplitudes too

What about magnitudes?

- Does rescattering also affect the *magnitudes* of amplitudes, apart from the *phases*?

- An estimate for magnitudes: factorisation/large number-of-colors (N_C)



CKM $\times$ Wilson coefficient $\times$factorisation

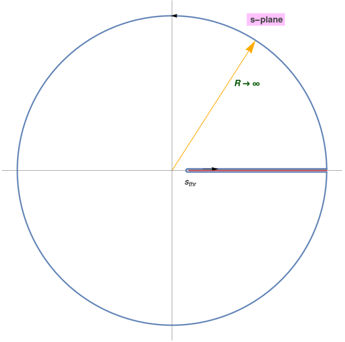
- Does not take (all) rescattering into account
- Decay constant and form factor come from data

Basic property of scattering amplitudes: analyticity

- Fundamental, model-independent property related to **causality**
- Cauchy's theorem: $A(s) = \frac{1}{2\pi i} \oint_C ds' \frac{A(s')}{s'-s}$ leads to

$$ReA(s) = \frac{1}{\pi} \int_{s_{thr}}^{\infty} ds' \frac{ImA(s')}{s'-s}$$

(Dispersion relation)



- Unitarity of S-matrix & dispersion relation:

$$\underbrace{ReA(s)}_{\text{Re at a point}} = \frac{1}{\pi} \underbrace{\int_{s_{thr}}^{\infty} ds' \frac{\tan \delta(s')}{s'-s} ReA(s')}_{\text{integral of Re along the physical region}}$$

Analyticity & what we learn about magnitudes

- Integral equation, studied by Muskhelishvili-Omnes
- One subtraction: needs one piece of physical information
- Single channel case (& one subtraction at s_0), **physical** solution:

$$|A_I(s)| = A_I(s_0)\,\underbrace{exp\{\frac{s-s_0}{\pi} PV \int_{4M_\pi^2}^{\infty} dz \frac{\delta_I(z)}{(z-s_0)(z-s)}\}}_{\text{Omnes factor } \Omega}$$

We need more than just large N_C!

$|A_I(s=m_D^2)| = (\text{large } N_C \text{ result}) \times (\text{Omnes factor})_I$

- More channels: Equally more solutions. No analytical solution

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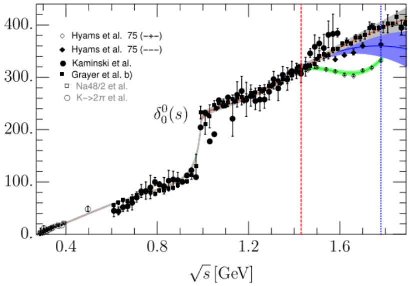
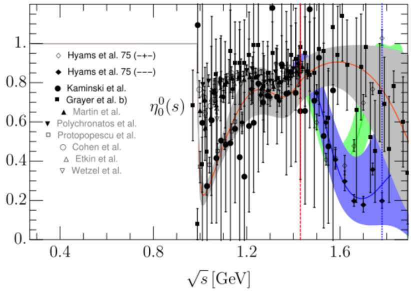
Summary of our method

- Separate CP-even (A) and odd (B) part
- Flavour basis to isospin
- Isospin blocks:
  - I=0 with 2 channels: $\pi\pi$ and KK
  - I=1 with KK elastic rescattering
  - I=2 with $\pi\pi$ elastic rescattering
- Isospin amplitudes treated with dispersion relations calculated numerically (based on Moussallam et al. [hep-ph/9909292])
- Physical input: unitarity (for integrand), large N_C limit (for polynomial ambiguity/subtraction point)

Data deployed: phase-shifts & inelasticities





- Use inelasticity and phase-shift parametrisations [Pelaez et al., 1907.13162],[Pelaez et al., 2010.11222]
- Parametrisations available up to energies $\sim m_D$ - extrapolate for higher & vary relevant parameters for uncertainties
- For I=1 and 2, extract Omnes factors from Br's of $A(D^+ \to \pi^+\pi^0) \sim A_{I=2}, A(D^+ \to K^+\overline{K^0}) \sim A_{I=1}$

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Omnes factors

For the isospin=0 channels we **calculate** numerically the Omnes matrix at $s = m_D^2$:

$$\Omega_{I=0} = \begin{pmatrix} 0.56e^{1.8i} & 0.70e^{-1.7i} \\ 0.55e^{-1.4i} & 0.65e^{2.3i} \end{pmatrix}$$

(In data: inelasticity taken mainly from $\pi\pi$ rescattering - solution I from Pelaez et al. '19)

The **physical solution** is

$$\begin{pmatrix} \mathbf{A}(D \to \pi\pi) \\ \mathbf{A}(D \to KK) \end{pmatrix} = \Omega_{I=0} \cdot \begin{pmatrix} \mathbf{A}_{\text{factorisation}}(D \to \pi\pi) \\ \mathbf{A}_{\text{factorisation}}(D \to KK) \end{pmatrix}$$

(Same for **B** instead of **A**)

It turns out:

Isospin=0 amplitudes are modified significantly (roughly by 20%) compared to pure large-N_C/factorisation result!

Branching fraction predictions

Adjusting $\delta^{\pi\pi}_{I=2}, \delta^{KK}_{I=1}$ we find:

Decay channel	$\frac{Br_{theo}}{Br_{exp}}$	(preliminary)
$D^0 \to \pi^+\pi^-$	1.1	
$D^0 \to \pi^0\pi^0$	1.2	
$D^0 \to K^+K^-$	1.1	
$D^0 \to K^0\overline{K^0}$	1.2	

We *calculate* these based on our method!

CP asymmetries

$\Delta A_{CP}^{dir,exp} = (-1.57 \pm 0.29) \cdot 10^{-3}$ [LHCb 2019]

We get $\Delta A_{CP}^{dir,theo} \leq \mathcal{O}(10^{-4})$!! (preliminary)

$$\text{Recall: } a_{CP}^{dir} \approx 2 \underbrace{r_{CKM}}_{\sim 6 \cdot 10^{-4}} \underbrace{\frac{|B(f)|}{|A(f)|}}_{\text{needs be } \mathcal{O}(1)} \cdot \underbrace{\sin \arg \frac{A(f)}{B(f)}}_{\text{needs be close to 1}}$$

- We **do not** find a sufficient enough enhancement of B's
- Phase-shift differences between A and B are closer to π than to $\pi/2$

Summary

- SM approach deploying
  1. S-matrix unitarity, scattering amplitude analyticity, isospin symmetry and factorisation
  2. as much data as possible (rescattering, form factors and decay constants, Br's of D^+ decays)
- We succeed in calculating the branching fractions in reasonable agreement with experiment, from scratch
- We still estimate the CP asymmetry **an order of magnitude too small** compared to the experimental value!
- Preliminary results, but seems difficult to accommodate the current value in our SM calculation...

Thank you very much!

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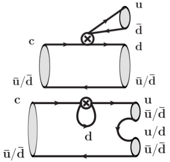
CPV in mesons

$$A_{CP}(f) = \frac{\Gamma(D^0 \to f) - \Gamma(\overline{D^0} \to \overline{f})}{\Gamma(D^0 \to f) - \Gamma(\overline{D^0} \to \overline{f})}$$

$$\approx \frac{A(D^0 \to f) - A(\overline{D^0} \to \overline{f})}{A(D^0 \to f) - A(\overline{D^0} \to \overline{f})} + \frac{< t_f >}{\tau_{D^0}} a_{CP}^{ind}$$

- $A_\Gamma = -a_{CP}^{ind} = (-2.8 \pm 2.8) \cdot 10^{-4}$
- For the decay $D^0 \to \pi^+\pi^-$: apply unitarity of the CKM matrix
  $A(D^0 \to \pi^+\pi^-) = \lambda_d A_d + \lambda_b A_b$
  $\to a_{CP}^{dir} \sim |\lambda_d||\lambda_b||A_d||A_b| \sin arg \frac{V_{cd}^* V_{ud}}{V_{cb}^* V_{ub}} \cdot \sin arg \frac{A_d}{A_b}$



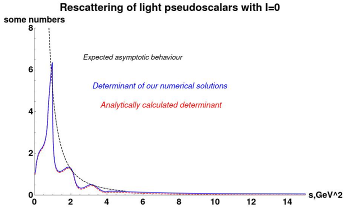
Isospin decomposition

- $\pi\pi$ states can have isospin=0,2. KK can have isospin=0,1.

$$\begin{pmatrix} A(\pi^+\pi^-) \\ A(\pi^0\pi^0) \\ A(K^+K^-) \\ A(K^0\overline{K}^0) \end{pmatrix} = \begin{pmatrix} -\frac{1}{2\sqrt{3}} & \frac{1}{\sqrt{6}} & 0 & 0 \\ \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{6}} & 0 & 0 \\ 0 & 0 & 1/2 & -1/2 \\ 0 & 0 & -1/2 & -1/2 \end{pmatrix} \begin{pmatrix} A_\pi^2 \\ A_\pi^0 \\ A_K^1 \\ A_K^0 \end{pmatrix}$$

Dispersion relations: Omnes solutions and properties

- Behaviour at ∞ determined by phase in integral: $A \to \frac{1}{s^n}$, $n = \frac{\delta(\infty)-\delta(s_{thr})}{\pi}$
- As many solutions as there are involved channels
- 2 channels: Omnes matrix $\begin{pmatrix} \Omega_{11} & \Omega_{12} \\ \Omega_{21} & \Omega_{22} \end{pmatrix}$ multiplied by appropriate polynomials



Large N_C limit & effective operators

- $T_{fac}(D^0 \to \pi^+\pi^-) = \lambda_d C_1 \frac{G_F}{\sqrt{2}} F_0^{D\pi}(m_\pi^2) f_\pi \cdot (m_D^2 - m_\pi^2)$
- $P_{fac}(D^0 \to \pi^+\pi^-) =$
  $\lambda_d(C_4 - 2C_6 \frac{M_\pi^2}{(m_u+m_d)(m_c+m_d)}) \frac{G_F}{\sqrt{2}} F_0^{D\pi}(m_\pi^2) f_\pi \cdot (m_D^2 - m_\pi^2)$
- $Q_1(i) = (\overline{d_i}c)_{V-A}(\overline{u}d_i)_{V-A}$, $Q_2(i) = (\overline{d_i}d_i)_{V-A}(\overline{u}c)_{V-A}$,
  $Q_{5,3} = (\overline{u}c)_{V-A}\sum_q(\overline{q}q)_{V\pm A}$,
  $Q_4 = \sum_q(\overline{u}q)_{V-A}(\overline{q}c)_{V-A}$, $Q_6 = -2\sum_q(\overline{u}q)_{S+P}(\overline{q}c)_{S-P}$
- $C_1 = 1.15$, $C_2 = -0.31$, $C_3 = 0.01$, $C_4 = -0.04$, $C_5 = 0.01$, $C_6 = -0.03$
- $\lambda_d = V_{cd}^* V_{ud} \approx 0.22$
- $\overline{m_c}(2GeV) = 1.045 GeV$
- Compare $m_D = 1865$ MeV to $\Lambda_{\chi PT} \approx m_\rho = 775$ MeV