

# Addressing the CKM unitarity problem with a vector-like up quark

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Cofinanciado por:



# Overview

- 1 Motivation
- 2 Lagrangian and Parameterisation
- 3 New Physics and Results
- 4 Conclusions

# Motivation

The SM has unitary quark and leptonic mixing matrices. The  $\nu SM$  (SM +  $\nu_R$ ) has a non-unitary leptonic mixing matrix, in principle. What if  $V_{CKM}$  is also not unitary?

$$\Delta_1 = 1 - |V_{11}|^2 - |V_{12}|^2 - |V_{13}|^2$$

$$\Delta_2 = 1 - |V_{21}|^2 - |V_{22}|^2 - |V_{23}|^2$$

$$\Delta_3 = 1 - |V_{31}|^2 - |V_{32}|^2 - |V_{33}|^2$$

Let's define this measure for the deviations from unitarity (DU) of mixing matrices  $V \equiv V_{CKM}$ .

# Motivation

Experiments are pointing towards non-zero deviations from unitarity (DU) in the quark sector:

- **CKM Unitarity Problem/Cabibbo Angle Anomaly** : Improved measurements of  $V_{ud}$  seem to indicate  $\sqrt{\Delta_1} \sim 0.04$   
Solution(?): VLQs (one of many, arguably the simplest)

# Motivation: How to deal with DU?

Approximations are common in literature.

- VLQs: Common assumption that heavy VLQs only couple to third generation quarks.

What if you want to study a region... **where these approximations fail?**

What if you want to ... **perform a general scan of the parameter space**, without biases?

# An up VLQ as a solution



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## Addressing the CKM unitarity problem with a vector-like up quark

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In [1], one up VLQ was introduced to explain the CKM unitarity problem.

<sup>1</sup>Branco, G.C., Penedo, J.T., Pereira, P.M.F. et al. J. High Energ. Phys. 2021, 99 (2021). [https://doi.org/10.1007/JHEP07\(2021\)099](https://doi.org/10.1007/JHEP07(2021)099)

## VLQs I

$$\begin{aligned}\Delta &\equiv \Delta_1 = 1 - |V_{ud}|^2 - |V_{us}|^2 - |V_{ub}|^2 = |\mathcal{V}_{L41}^*|^2 = s_{14}^2 \\ \Delta_2 &= 1 - |V_{cd}|^2 - |V_{cs}|^2 - |V_{cb}|^2 = |\mathcal{V}_{L42}^*|^2 = c_{14}^2 s_{24}^2 \\ \Delta_3 &= 1 - |V_{td}|^2 - |V_{ts}|^2 - |V_{tb}|^2 = |\mathcal{V}_{L43}^*|^2 = c_{14}^2 c_{24}^2 s_{34}^2\end{aligned}$$

The experimental data suggests  $\sqrt{\Delta} \sim 0.04$  [2]

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<sup>2</sup>B. Belfatto, R. Beradze and Z. Berezhiani, Eur. Phys. J. C 80 (2020) 149  
[1906.02714]

## Lagrangian

$$\mathcal{L}_M = - \begin{pmatrix} \bar{d}_L^0 & \bar{D}_L^0 \end{pmatrix} \mathcal{M}_d \begin{pmatrix} d_R^0 \\ D_R^0 \end{pmatrix} - \begin{pmatrix} \bar{u}_L^0 & \bar{U}_L^0 \end{pmatrix} \mathcal{M}_u \begin{pmatrix} u_R^0 \\ U_R^0 \end{pmatrix} + \text{h.c.}$$

$$\mathcal{M}_q = \begin{pmatrix} m_q & \bar{m}_q \\ \bar{M}_q & M_q \end{pmatrix} \left. \vphantom{\begin{pmatrix} m_q & \bar{m}_q \\ \bar{M}_q & M_q \end{pmatrix}} \right\} \begin{matrix} 3 \\ n_q \end{matrix}$$

$\underbrace{\hspace{10em}}_3 \quad \underbrace{\hspace{10em}}_{n_q}$



## Lagrangian II

$$\mathcal{V}_L^{q\dagger} \mathcal{M}_q \mathcal{V}_R^q = \mathcal{D}_q$$

$$\mathcal{V}_\chi^q = \begin{pmatrix} A_\chi^q \\ \text{-----} \\ B_\chi^q \end{pmatrix}$$

$$\chi = L, R \quad , \quad q = u, d$$

$$A = 3 \times (3 + n) \quad , \quad B = n \times (3 + n)$$

Charged Currents:

$$\mathcal{L}_W = -\frac{g}{\sqrt{2}} (\bar{u}_L \quad \bar{U}_L) V \gamma^\mu \begin{pmatrix} d_L \\ D_L \end{pmatrix} W_\mu^+ + \text{h.c.}$$

$$V = A_L^{u\dagger} A_L^d$$

# Lagrangian III

Neutral Interactions:

$$\mathcal{L}_H = -\frac{h}{v} \left[ (\bar{q}_L \quad \bar{Q}_L) F^q \mathcal{D}_q \begin{pmatrix} q_R \\ Q_R \end{pmatrix} \right] + \text{h.c.}$$

$$\mathcal{L}_H = -\frac{h}{v} \left[ (\bar{n}_L \quad \bar{N}_L) F \mathcal{D} \begin{pmatrix} n_L^c \\ N_L^c \end{pmatrix} \right] + \text{h.c.}$$

$$F^q = A_L^{q\dagger} A_L^q$$

## Parameterisation

$$\mathcal{V}_L^\dagger \mathcal{M}_u \mathcal{V}_R = \mathcal{D}_u$$

$$\mathcal{V}_L^\dagger = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & c_{34} & s_{34} \\ 0 & 0 & -s_{34} & c_{34} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & c_{24} & 0 & s_{24} e^{-i\delta_{24}} \\ 0 & 0 & 1 & 0 \\ 0 & -s_{24} e^{i\delta_{24}} & 0 & c_{24} \end{pmatrix} \begin{pmatrix} c_{14} & 0 & 0 & s_{14} e^{-i\delta_{14}} \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -s_{14} e^{i\delta_{14}} & 0 & 0 & c_{14} \end{pmatrix} \\ \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & c_{23} & s_{23} & 0 \\ 0 & -s_{23} & c_{23} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta_{13}} & 0 \\ 0 & 1 & 0 & 0 \\ -s_{13} e^{i\delta_{13}} & 0 & c_{13} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 & 0 \\ -s_{12} & c_{12} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

where  $c_{ij} = \cos \theta_{ij}$  and  $s_{ij} = \sin \theta_{ij}$ , with  $\theta_{ij} \in [0, \pi/2]$ ,  $\delta_{ij} \in [0, 2\pi]$ .

# Parameterisation II

$$V = A_L^\dagger = \begin{pmatrix} c_{12} c_{13} c_{14} & s_{12} c_{13} c_{14} & s_{13} c_{14} e^{-i\delta_{13}} \\ \dots & \dots & \dots \\ \dots & \dots & \dots \\ \dots & \dots & \dots \end{pmatrix} \equiv \begin{pmatrix} K_{\text{CKM}} \\ \dots \\ K_T \end{pmatrix}$$

In the weak basis where down quarks are diagonal.

$$K_{\text{CKM}} = H_L U_{\text{CKM}} \equiv (\mathbb{1} - \eta) U_{\text{CKM}}$$

## Parameterisation III

$$F^u = \left( \begin{array}{c|c} K_{\text{CKM}} K_{\text{CKM}}^\dagger & K_{\text{CKM}} K_T^\dagger \\ \hline K_T K_{\text{CKM}}^\dagger & K_T K_T^\dagger \end{array} \right)$$

$$= \begin{pmatrix} 1 - \Delta_1 & -\sqrt{\Delta_1}\sqrt{\Delta_2}e^{-i(\delta_{14}-\delta_{24})} & -\sqrt{\Delta_1}\sqrt{\Delta_3}e^{-i\delta_{14}} & -\sqrt{\Delta_1}e^{-i\delta_{14}} \\ -\sqrt{\Delta_1}\sqrt{\Delta_2}e^{i(\delta_{14}-\delta_{24})} & 1 - \Delta_2 & -\sqrt{\Delta_2}\sqrt{\Delta_3}e^{-i\delta_{24}} & -\sqrt{\Delta_2}e^{-i\delta_{24}} \\ -\sqrt{\Delta_1}\sqrt{\Delta_3}e^{i\delta_{14}} & -\sqrt{\Delta_2}\sqrt{\Delta_3}e^{i\delta_{24}} & 1 - \Delta_3 & -\sqrt{\Delta_3} \\ -\sqrt{\Delta_1}e^{i\delta_{14}} & -\sqrt{\Delta_2}e^{i\delta_{24}} & -\sqrt{\Delta_3} & \Delta_1 + \Delta_2 + \Delta_3 \end{pmatrix}$$

# New Physics

## New Physics

- $D^0-\bar{D}^0$  (tree level)
- $K^0-\bar{K}^0$  (loop)
- $B_{d,s}^0-\bar{B}_{d,s}^0$  (loop)
- CP violation in  $K_L \rightarrow \pi\pi$
- Rare top decays  $t \rightarrow (u, c)Z$

**Conservative estimate: NP < Exp**

# New Physics II

We used this to restrain our fit, for instance

Observable	$m_T = 1 \text{ TeV}$	$m_T = 3 \text{ TeV}$
$\Delta m_K$	$ V_{Td}  V_{Ts}  < 7.4 \times 10^{-4}$	$ V_{Td}  V_{Ts}  < 2.7 \times 10^{-4}$
$\Delta m_B$	$ V_{Td}  V_{Tb}  < 6.7 \times 10^{-4}$	$ V_{Td}  V_{Tb}  < 3.4 \times 10^{-4}$
$\Delta m_{B_s}$	$ V_{Ts}  V_{Tb}  < 3.2 \times 10^{-3}$	$ V_{Ts}  V_{Tb}  < 1.6 \times 10^{-3}$
$ \epsilon_K $	$ V_{Td}  V_{Ts}  \sqrt{ \sin 2\Theta } < 8.8 \times 10^{-5}$	$ V_{Td}  V_{Ts}  \sqrt{ \sin 2\Theta } < 3.1 \times 10^{-5}$

where  $\Theta = \text{Arg}(V_{Ts}^* V_{Td})$

The mass dependence on the bounds comes from the Inami-Lim loop functions.

# Perturbativity

$$\text{Tr} \left( mm^\dagger + \bar{m} \bar{m}^\dagger \right) = p m_t^2$$

$$\text{Tr} \left( mm^\dagger + \bar{m} \bar{m}^\dagger \right) \simeq m_t^2 + (\Delta_1 + \Delta_2 + \Delta_3) m_T^2$$

$$\sqrt{\Delta_1 + \Delta_2 + \Delta_3} = \sqrt{p-1} \frac{m_t}{m_T}$$

For  $\Delta \equiv \Delta_1 \gg \Delta_2 \sim \Delta_3$  and since  $\sqrt{p-1}$  should be  $O(1)$

$$m_T \lesssim \frac{m_t}{\sqrt{\Delta}}$$



# Fit

## Procedure

- NP + Perturbativity to restrict parameter space
- Best fit region minimizing

$$\chi^2 = \sum_{ij} \left( \frac{V_{ij} - V_{ij}^c}{\sigma(V_{ij})} \right)^2 + \left( \frac{\gamma - \gamma^c}{\sigma(\gamma)} \right)^2$$

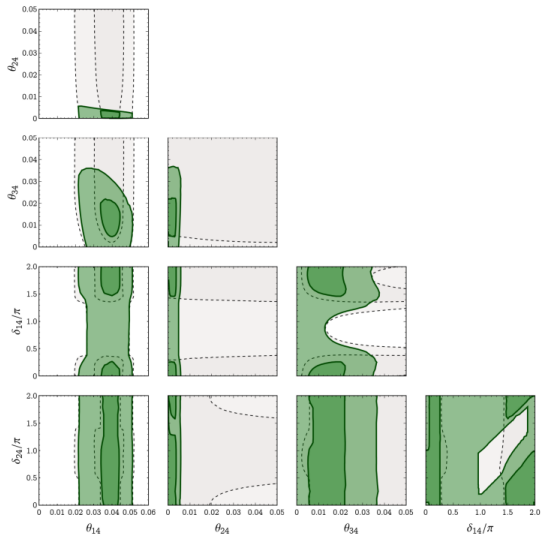
## Fit II

Using best-fit values of  $CKM$  without imposing unitarity,  $V_{ij}^c$  is given by

$$|K_{CKM}| = \begin{pmatrix} 0.97370 \pm 0.00014 & 0.2245 \pm 0.0008 & (3.82 \pm 0.24) \times 10^{-3} \\ 0.221 \pm 0.004 & 0.987 \pm 0.011 & (41.0 \pm 1.4) \times 10^{-3} \\ (8.0 \pm 0.3) \times 10^{-3} & (38.8 \pm 1.1) \times 10^{-3} & 1.013 \pm 0.030 \end{pmatrix}$$

And  $\gamma^c = (72.1^{+4.1}_{-4.5})^\circ$ , since the presence of an up-type VLQ does not affect the value of the phase  $\gamma = \arg(-V_{ud}V_{cb}V_{ub}^*V_{cd}^*)$ , which is obtained from SM tree-level dominated  $B$  decays.

## Results



# Results II

The selected point in parameter space corresponds to a mass  $m_T = 1.5$  TeV and is described by

$$\theta_{12} = 0.2265, \quad \theta_{13} = 0.003818, \quad \theta_{23} = 0.03998,$$

$$\theta_{14} \simeq \sqrt{\Delta_{(1)}} = 0.03951, \quad \theta_{24} \simeq \sqrt{\Delta_2} = 0.002078, \quad \theta_{34} \simeq \sqrt{\Delta_3} = 0.01271,$$

$$\delta_{13} = 0.396\pi, \quad \delta_{14} = 1.818\pi, \quad \delta_{24} = 0.728\pi,$$

corresponding to  $\chi^2 \simeq 3.2$ , to a perturbativity factor  $p - 1 \simeq 0.13$ .

## Results III

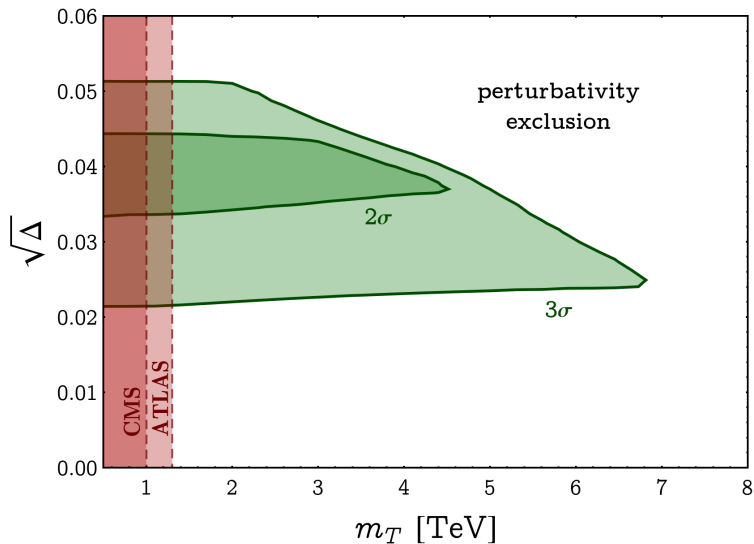
$$|\nu_L^\dagger| \simeq \begin{pmatrix} \boxed{\begin{matrix} 0.9737 & 0.2243 & 0.0038 \\ 0.2243 & 0.9737 & 0.0400 \\ 0.0081 & 0.0393 & 0.9991 \end{matrix}} & 0.0395 \\ 0.0390 & 0.0065 & 0.0126 & 0.9991 \end{pmatrix}$$

## Results IV

$$|F^u| = \begin{pmatrix} 0.99843969 & 0.00008203 & 0.00050179 & 0.03946672 \\ 0.00008203 & 0.99999569 & 0.00002638 & 0.00207496 \\ 0.00050179 & 0.00002638 & 0.99983863 & 0.01269224 \\ 0.03946672 & 0.00207496 & 0.01269224 & 0.001726 \end{pmatrix}$$

$$|\eta| = \begin{pmatrix} 0.78016 & 0.04102 & 0.25089 \\ 0.04102 & 0.00216 & 0.01319 \\ 0.25089 & 0.01319 & 0.08069 \end{pmatrix} \times 10^{-3}$$

## Results V



# Future Work

A Review on VLQs, in collaboration with J. M. Alves, A.L. Cherchiglia and C.C. Nishi.



# Conclusions

- DUs might play an important part in the discovery of New Physics.
- DUs can be the bridge to new heavy particles that might explain current anomalies.
- Vector-like quarks are one of the simplest solutions available.

# The End

**Thank You!**