

Patrick Fox Fermilab

Based on JHEP **1604** (2016) 135 (arXiv:1512.06853)

LHC DM WG meeting Nov 30, 2021

The story so far

- •DM makes up 23% of the universe
- •Gravitates like ordinary matter, but is non-baryonic
- •Is dark i.e. neutral under SM (not coloured, or charged)
- Does not interact much with itself
- •Does not couple to massless particle
- •Was no relativistic at time of CMB
- Is long lived
- IF DM is a thermal relic:
- •A weak scale annihilation x-sec gives correct abundance •Mass range is $10~{\rm MeV}\lesssim m_\chi\lesssim 70~{\rm TeV}$

What representation of the Lorentz group?

Fermionic DM (eg neutralino)

[Berlin, Hooper, McDermott 1404.0022]

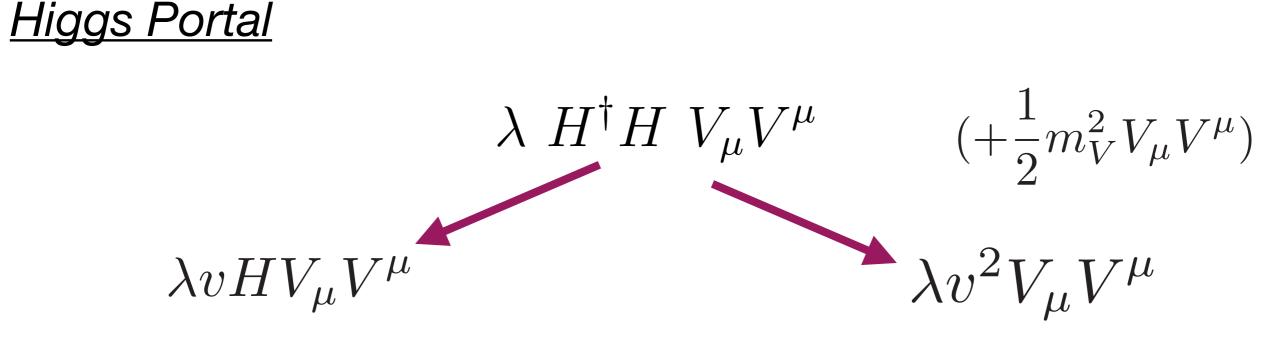
DM bilinear	SM fermion bilinear				
fermion DM	$\bar{f}f$	$ar{f}\gamma^5 f$	$ar{f}\gamma^\mu f$	$ar{f}\gamma^\mu\gamma^5 f$	
$\bar{\chi}\chi$	$\sigma v \sim v^2, \sigma_{\rm SI} \sim 1$	$\sigma v \sim v^2, \sigma_{\rm SD} \sim q^2$	_	_	
$\bar{\chi}\gamma^5\chi$	$\sigma v \sim 1, \sigma_{ m SI} \sim q^2$	$\sigma v \sim 1, \sigma_{ m SD} \sim q^4$	_	_	
$\left[\bar{\chi} \gamma^{\mu} \chi \text{ (Dirac only)} \right]$	—	_	$\sigma v \sim 1, \sigma_{\rm SI} \sim 1$	$\left \sigma v \sim 1, \sigma_{ m SD} \sim v_{\perp}^2 ight $	
$\bar{\chi}\gamma^{\mu}\gamma^{5}\chi$			$\sigma v \sim v^2, \sigma_{\rm SI} \sim v_\perp^2$	$\sigma v \sim 1, \sigma_{ m SD} \sim 1$	

Scalar DM (eg sneutrino, axion/ALP)

DM bilinear	SM fermion bilinear				
scalar DM	$\overline{f}f$	$ar{f}\gamma^5 f$	$ar{f}\gamma^\mu f$	$ar{f}\gamma^{\mu}\gamma^5 f$	
$\phi^\dagger \phi$	$\sigma v \sim 1, \sigma_{\rm SI} \sim 1$	$\sigma v \sim 1, \sigma_{ m SD} \sim q^2$		—	
$\phi^{\dagger} \overset{\leftrightarrow}{\partial_{\mu}} \phi \text{ (complex only)}$	_	_	$\sigma v \sim v^2, \sigma_{\rm SI} \sim 1$	$\sigma v \sim v^2, \sigma_{\rm SD} \sim v_\perp^2$	

Spin 1

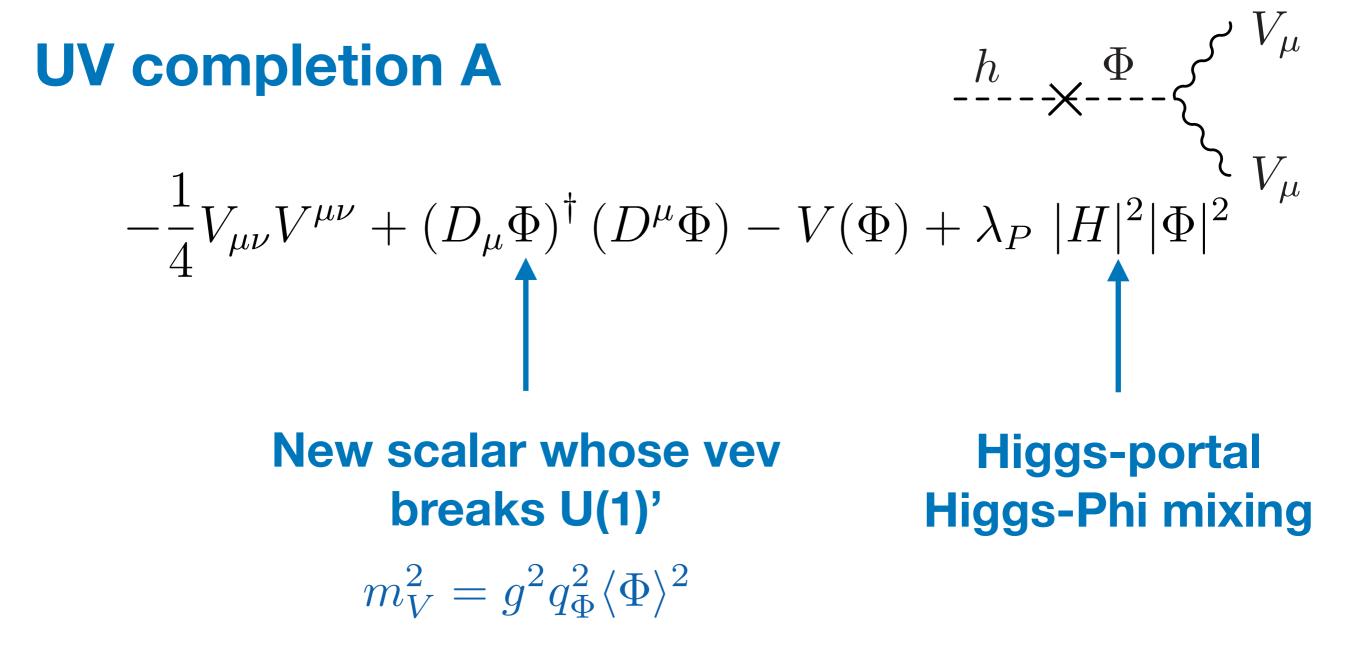
(Massive) Vector must be associated with a (broken) gauge symmetry e.g. U(1)' In what ways can it couple to SM?



Invisible Higgs width

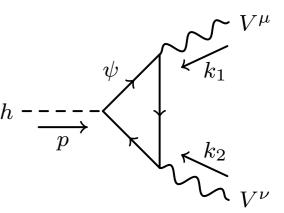
Vector Mass

Not U(1)' gauge invariant



DM phenomenology determined by theory with new scalar and a new vector, see eg Baek, Ko, Park (1212.2131), and other speakers

UV completion B



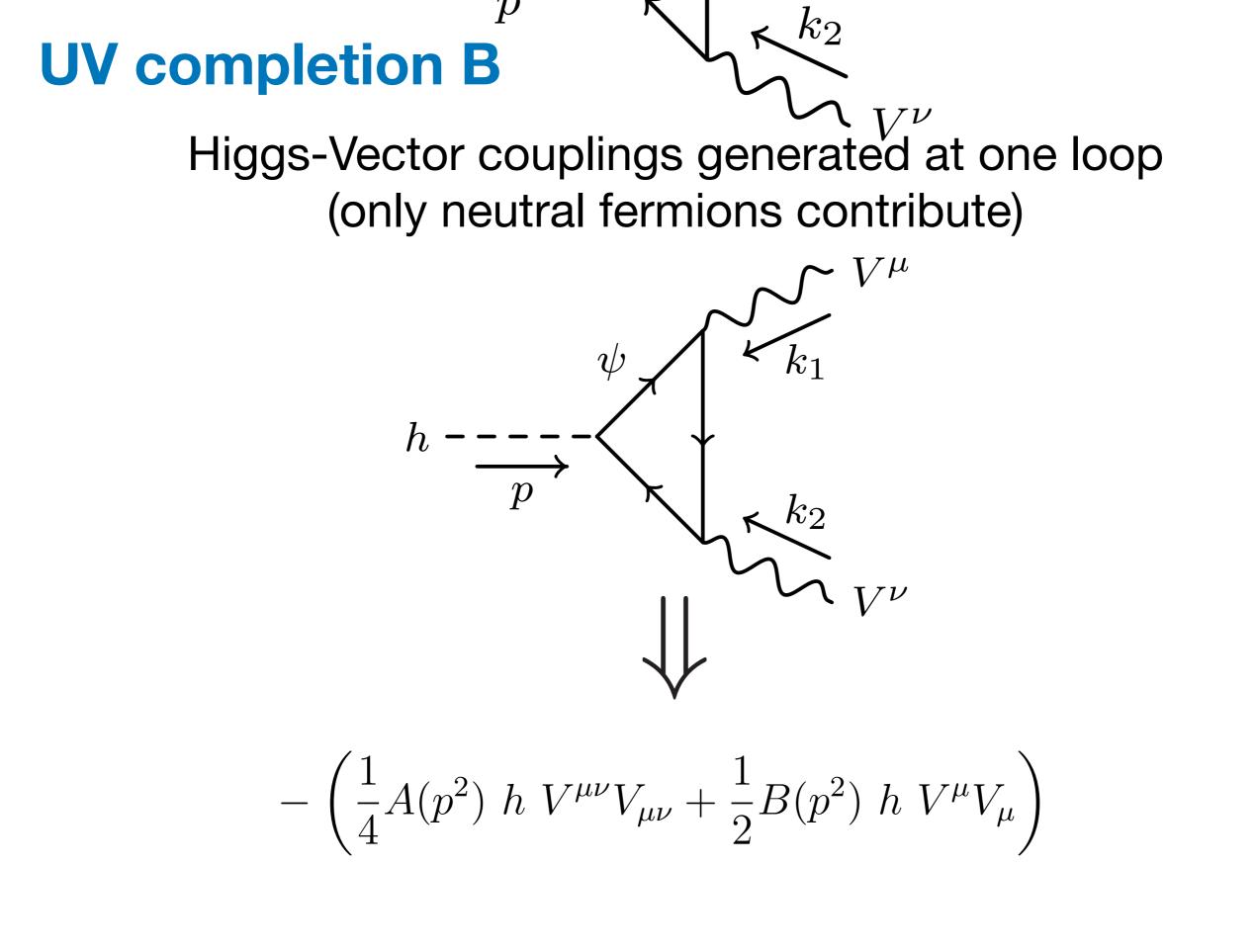
- Radiatively generated Higgs-Vector coupling
- Necessarily more complicated
- New and interesting phenomenology eg decouple Phi, new states at LHC

Field	$(SU(2)_W, U(1)_Y, U(1)')$	Field	$(SU(2)_W, U(1)_Y, U(1)')$
ψ_{1lpha}	(2, 1/2, 1) Vec (2, -1/2, -1) Vec	ctor- ψ_{2lpha}	(2, 1/2, -1)
χ_{1lpha}	(2, -1/2, -1) / like	pair χ_{2lpha}	(2, 1/2, -1) (2, -1/2, 1)
n_{1lpha}	(1, 0, -1)	\rightarrow $n_{2\alpha}$	$(1, \ 0, \ 1)$
Φ	$(1, \hspace{0.1cm} 0, \hspace{0.1cm} Q_{\Phi})$	H	(2, -1/2, 0)
	$Q_{\Phi} \neq \pm 1, \pm 2$		

U(1)' charge conjugation symmetry ($1 \leftrightarrow 2$) makes V stable

UV completion **B**

Four parameters (2 masses, 2 Yukawas) as well as gauge coupling and vector DM mass



Higgs invisible width

$$\Gamma(h \to VV) = \frac{1}{64\pi m_h} \sqrt{1 - 4\frac{m_V^2}{m_h^2}} \left[\left| A(m_h^2) \right|^2 m_h^4 \left(1 - 4\frac{m_V^2}{m_h^2} + 6\frac{m_V^4}{m_h^4} \right) + 6\operatorname{Re}\left(A^*(m_h^2) B(m_h^2) \right) m_h^2 \left(1 - 2\frac{m_V^2}{m_h^2} \right) + \frac{1}{2} \left| B(m_h^2) \right|^2 \frac{m_h^4}{m_V^4} \left(1 - 4\frac{m_V^2}{m_h^2} + 12\frac{m_V^4}{m_h^4} \right) \right]$$

Finite as $m_V \to 0$

 k_1

 ψ

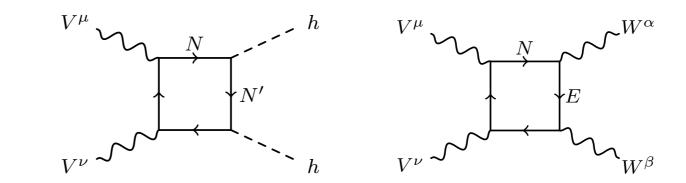
a

Direct detection cross section

+

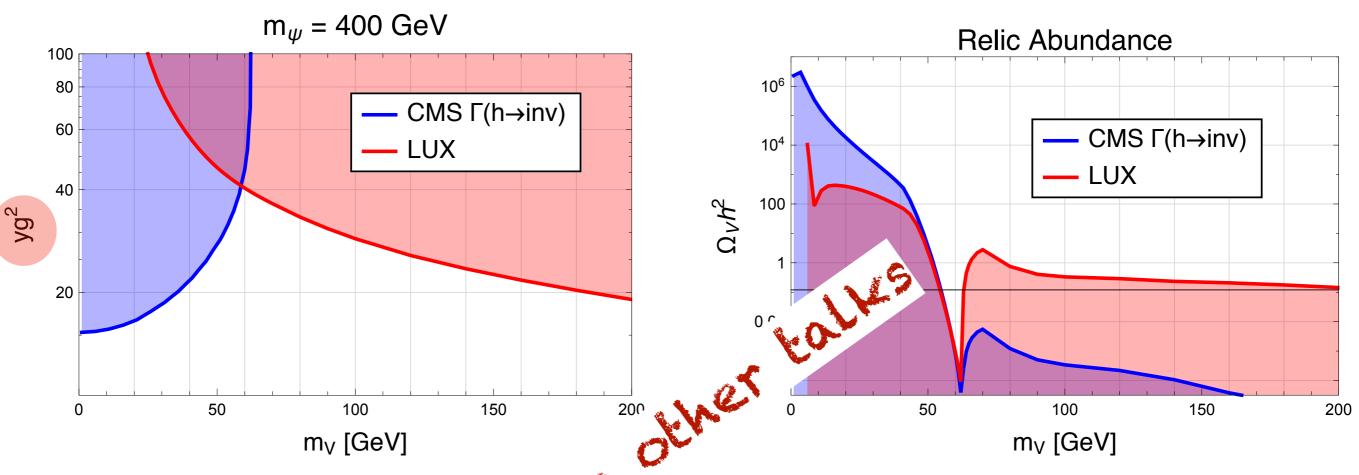
$$\sigma_{\rm SI} = \frac{1}{4\pi m_h^4} \left(\frac{f_n}{v}\right)^2 \left(\frac{m_n^2}{m_n + m_V}\right)^2 |B(0) - A(0) \ m_V^2|^2$$

Annihilation cross section $VV \rightarrow hh, ZZ, WW, \gamma\gamma, hZ, Z\gamma$



Single fermion limit

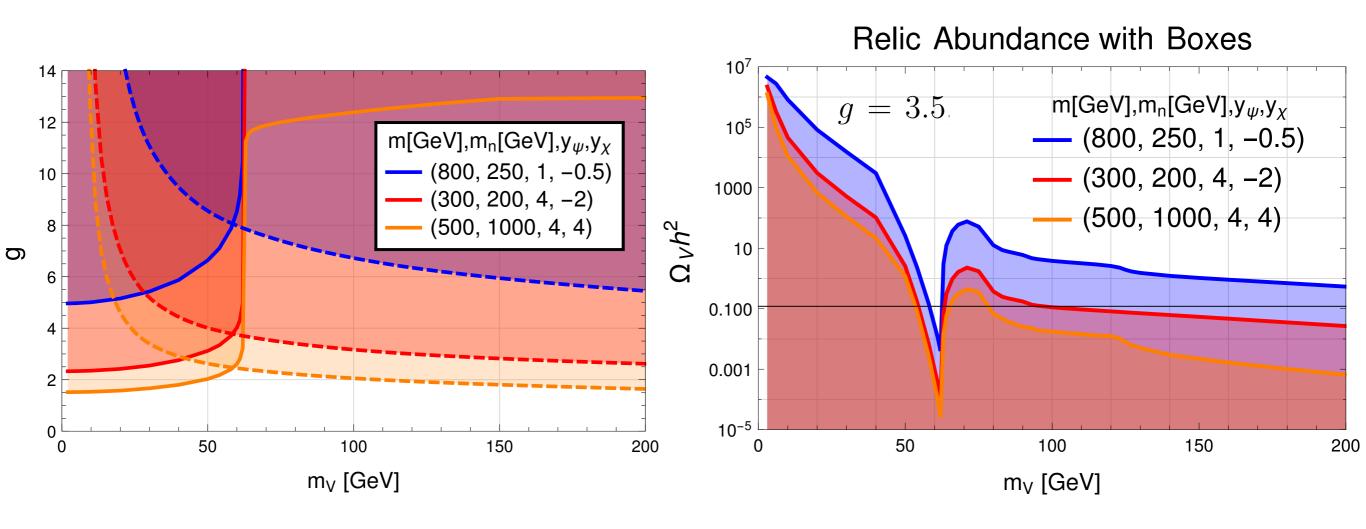
Everything decoupled except one light neutral fermion



- ن 20/fb. 8TeV (it's been a while) در الم CMS = VBF, invisible ' درمج ' 2
 LUX = 10 ton-day osure

Only Higgs-funnel region works for cosmology (but there are ways out)

All fermions



- Loop generated Higgs-V coupling, cf tree-level Higgs portal
- New sub-TeV electroweak states (decays involve MET)
- Collider perspective similar to simpler single fermion case

Conclusions

- DM could be a spin-1 state
- Could couple to SM through Higgs exchange
- Direct Higgs portal not gauge invariant, need a UV completion
- Generate through tree-level H-Phi mixing
- Generate through heavy new fermions running in loop
- Simple model with a few parameters, only subset relevant for collider physics
- LHC provides strongest constraint below mH/2
- Can decouple second higgs unlike tree level case
- New fermionic states to search for

Python/Mathematica code exists for this

$$A(p^{2}) = \sum_{i} \left(\frac{g^{2}y_{i}m_{i}}{2\sqrt{2}\pi^{2}}\right) F_{1}(p^{2}, m_{i}),$$
$$B(p^{2}) = \sum_{i} \left(\frac{g^{2}y_{i}m_{i}}{2\sqrt{2}\pi^{2}}\right) F_{2}(p^{2}, m_{i}).$$

$$F_{1}(p^{2},m) = \frac{1}{2bm^{2}(b-4a)^{2}} \Big\{ 2m^{2}(b-2a) \Big[4a(a-1) + b(1+6a-b) \Big] C_{0} \\ - 2a(2a+b)\Delta B_{0} + (b-2a)(b-4a) \Big\} \\ F_{2}(p^{2},m) = \frac{4a^{2}}{b(b-4a)^{2}} \Big\{ 2(b-a)\Delta B_{0} - 2m^{2} \Big[4a(a-1) + b(1-2a+b) \Big] C_{0} + 4a-b \Big\}$$

 $a \equiv \frac{m_V^2}{4m^2}, \qquad b \equiv \frac{p^2}{4m^2}$

$$C_{0} = \frac{1}{4m^{2}b\beta} \sum_{j,k=1}^{2} \left[2\text{Li}_{2} \left(\frac{1+(-1)^{j}\beta}{1+(-1)^{k}\text{X}\beta} \right) - \text{Li}_{2} \left(\frac{(1+(-1)^{j}\beta)^{2}}{1+(-1)^{k}2\text{Y}\beta+\beta^{2}} \right) \right]$$
$$\Delta B_{0} \equiv B_{0}(m_{V}^{2};m,m) - B_{0}(p^{2};m,m)$$
$$\sqrt{1-h} \left[\sqrt{h} \right] \sqrt{1-a} \left[\sqrt{a} \right]$$

$$= 2\sqrt{\frac{1-b}{b}} \arctan\left[\sqrt{\frac{b}{1-b}}\right] - 2\sqrt{\frac{1-a}{a}} \arctan\left[\sqrt{\frac{a}{1-a}}\right].$$