

Toward black hole entropy in chiral loop quantum supergravity

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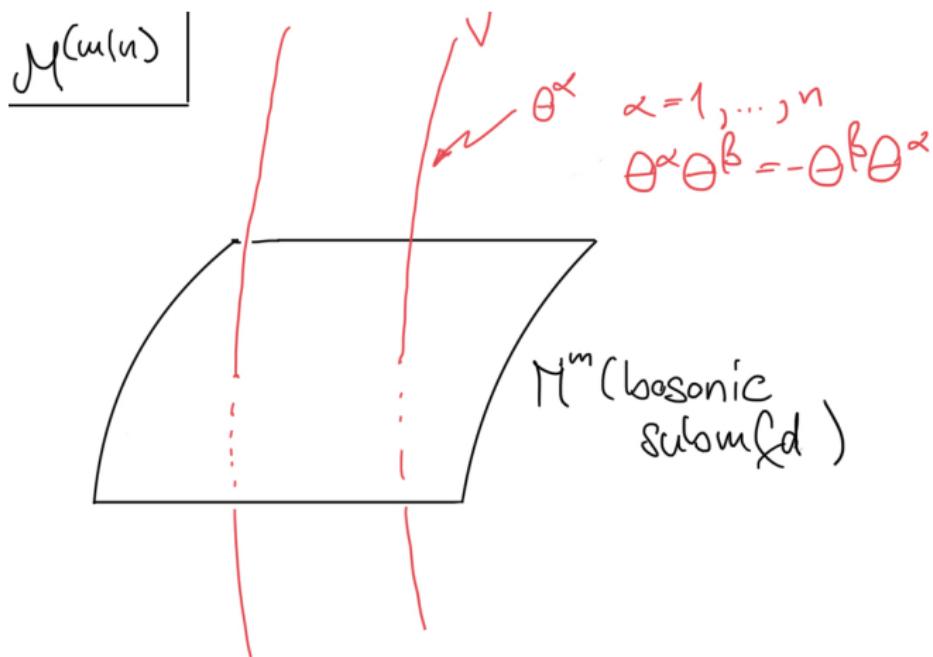
(joint collaboration with Hanno Sahlmann)

Section 1

Introduction

Supermanifold

Phase space of field theories containing both bosonic & fermionic d.o.f.
has structure of **supermanifold**



Super Lie group & algebra

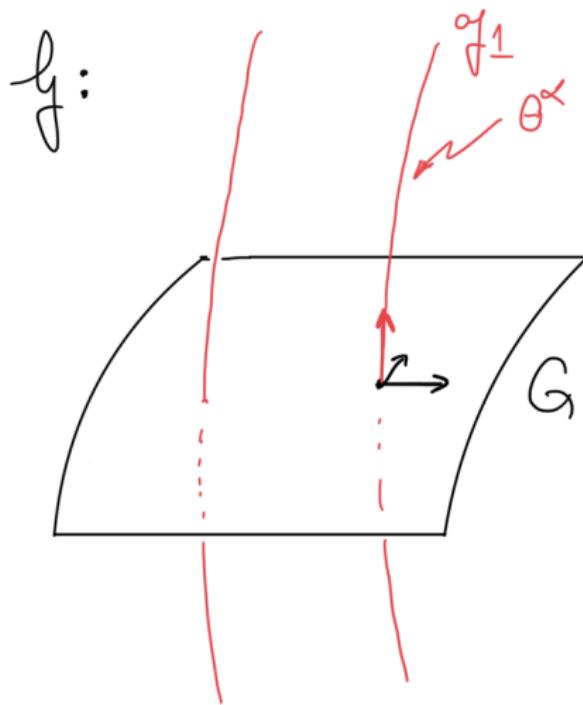
SUSY field theories: Field theories with underlying symmetry group given by a **supergroup**

- Lie superalgebra:

$$\mathfrak{g} := T_e \mathcal{G} = \mathfrak{g}_0 \oplus \mathfrak{g}_1$$

- Lie bracket:

$$[X, Y] := XY - (-1)^{|X||Y|} YX$$



Supersymmetry

Super anti-de Sitter group $\mathrm{OSp}(\mathcal{N}|4)$

$$\mathfrak{osp}(1|4) = \underbrace{\mathbb{R}^{1,3} \ltimes \mathfrak{so}(1,3)}_{\mathfrak{g}_0} \oplus \underbrace{\mathcal{S}_{\mathbb{R}}}_{\mathfrak{g}_1}$$

Generators: P_I , M_{IJ} , Q_α (Majorana spinor)

$$[P_I, Q_\alpha] = -\frac{1}{2L} Q_\beta (\gamma_I)^\beta{}_\alpha$$

$$[M_{IJ}, Q_\alpha] = \frac{1}{2} Q_\beta (\gamma_{IJ})^\beta{}_\alpha$$

$$\{Q_\alpha, Q_\beta\} = \frac{1}{2} (C\gamma_I)_{\alpha\beta} P^I + \frac{1}{4L} (C\gamma^{IJ})_{\alpha\beta} M_{IJ}$$

Further generators $Q_\alpha^r \rightarrow \mathcal{N}\text{-extended SUSY!}$

$$AdS_4 := \{x \in \mathbb{R}^5 \mid -(x^0)^2 + (x^1)^2 + \dots + (x^3)^2 - (x^4)^2 = -L^2\}$$

What has been done?

- Canonical SUGRA adapted to LQG [Jacobson '88]
- Hidden $\mathfrak{osp}(1|2)$ -symmetry in constraint algebra, construction of $\mathfrak{osp}(1|2)$ -valued connection [Fülöp '94, Gambini+Obregon+Pullin '96]
- Quantization [Gambini+Obregon+Pullin '96, Ling+Smolin '99]

- Considerations remain formal

But:

- Origin of hidden symmetry unclear
- Generalizations?

- Canonical theory for higher D and quantization [Bodendorfer +Thiemann+Thurn '11]

But:

- SUSY not manifest
- Formalism/constraints quite complicated

My goals:

Follow [L+S]: Keep SUSY manifest as much as possible!

- Origin of $\mathfrak{osp}(1|2)$ -symmetry? → Understand the geometric origins with a view towards generalizations:
 - Immirzi parameters
 - Higher $\mathcal{N} > 1$
 - Boundary theory → **black holes**
- Mathematically rigorous formulation, both classically and in quantum theory:
 - Construction of super spin-nets (\rightarrow *super parallel transport*)
 - Structure of Hilbert spaces \leftrightarrow relation to standard quantization of fermions in LQG

Section 2

Supergravity & boundary theory

Supergravity

AdS (Holst-)Supergravity as 'constrained' gauge theory with gauge group $\text{OSp}(\mathcal{N}|4)$ [D'Auria-Fré-Regge '80, Castellani-D'Auria-Fré '91, KE '21+'22]

Gauge field ($\mathcal{N} = 1$)

$$\mathcal{A} : \mathcal{M} \rightarrow \mathfrak{osp}(1|4), \quad \iota : M \hookrightarrow \mathcal{M}$$

Decomposition

$$\mathcal{A} = \underbrace{\text{pr}_{\mathfrak{g}_1} \circ \mathcal{A}}_{\psi} + \underbrace{\text{pr}_{\mathbb{R}^{1,3}} \circ \mathcal{A}}_e + \underbrace{\text{pr}_{\mathfrak{spin}(1,3)} \circ \mathcal{A}}_{\omega}$$

- e : co-frame (\rightarrow metric)
- ω : spin connection
- ψ : (spin-3/2) Rarita-Schwinger field

Supergravity and LQG

super Holst-MacDowell-Mansouri action [KE+HS '21]

$$S_{\text{sH-MM}}[\mathcal{A}] = \int_M \iota^* \text{STr}(F[\mathcal{A}] \wedge \mathbf{P}_\beta F[\mathcal{A}])$$

Super Holst-operator

$$\mathbf{P}_\beta := \mathbf{0} \oplus \mathcal{P}_\beta \oplus \mathcal{P}_\beta : \mathfrak{osp}(1|4) \rightarrow \mathfrak{osp}(1|4)$$

with

$$\mathcal{P}_\beta := \frac{\mathbb{1} + i\beta\gamma_5}{2\beta} \quad \beta : \text{Immirzi}$$

- → yields Holst action of $D = 4, \mathcal{N} = 1$ AdS-SUGRA *+ bdy terms*
- **Can show:** Boundary theory **unique** by SUSY-invariance!
[Andrianopoli+D'Auria '14, KE+HS '21]

Extended SUGRA and boundary theory

What about **extended** SUSY? → Consider $\mathcal{N} = 2$ [KE+HS '21]

Super Cartan connection

$$\mathcal{A} = e^I P_I + \frac{1}{2} \omega^{IJ} M_{IJ} + \Psi_r^\alpha Q_\alpha^r + \hat{A} T$$

\hat{A} : U(1) gauge field (*graviphoton*)

super Holst-MacDowell-Mansouri action ($\mathcal{N} = 2$)

$$S_{\text{sH-MM}}^{\mathcal{N}=2}[\mathcal{A}] = \int_M \iota^* \text{STr}(F[\mathcal{A}] \wedge \mathbf{P}_\beta F[\mathcal{A}])$$

$$\mathbf{P}_\beta := \mathbf{0} \oplus \mathcal{P}_\beta \oplus \mathcal{P}_\beta \oplus \mathcal{P}_\beta \oplus \frac{1}{2\beta} (1 + \beta \star)$$

⇒ β literally has interpretation as **θ -ambiguity!**

Section 3

The chiral theory

Chiral Theory

In general: \mathbf{P}_β destroys manifest SUSY-invariance!

→ This changes in **chiral theory** ($\beta = \mp i$)

Holst projection [KE+HS '21]

$$\mathbf{P}_{-i} : \mathfrak{osp}(\mathcal{N}|4)_{\mathbb{C}} \rightarrow \mathfrak{osp}(\mathcal{N}|2)_{\mathbb{C}}$$

$$M_{IJ} \mapsto T_i^+ = \frac{1}{2}(J_i + iK_i)$$

$$Q_\alpha \mapsto Q_A$$

Q_A : left-handed Weyl-fermion

Chiral Theory

Super Ashtekar connection [KE+HS '21]

$$\mathcal{A}^+ := \mathbf{P}_{-i}\mathcal{A} = A^{+i}T_i^+ + \psi_r^A Q_A^r + \frac{1}{2}\hat{A}_{rs}T^{rs}$$

Chiral action

$$S_{\text{sH-MM}}^{\beta=-i} = \frac{i}{\kappa} \int_M \langle \mathcal{E} \wedge F(\mathcal{A}^+) \rangle + \frac{1}{4L^2} \langle \mathcal{E} \wedge \mathcal{E} \rangle + \underbrace{S_{\text{CS}}^{\text{OSP}(\mathcal{N}|2)}(\mathcal{A}^+)}_{\text{boundary term}}$$

\mathcal{E} : super electric field

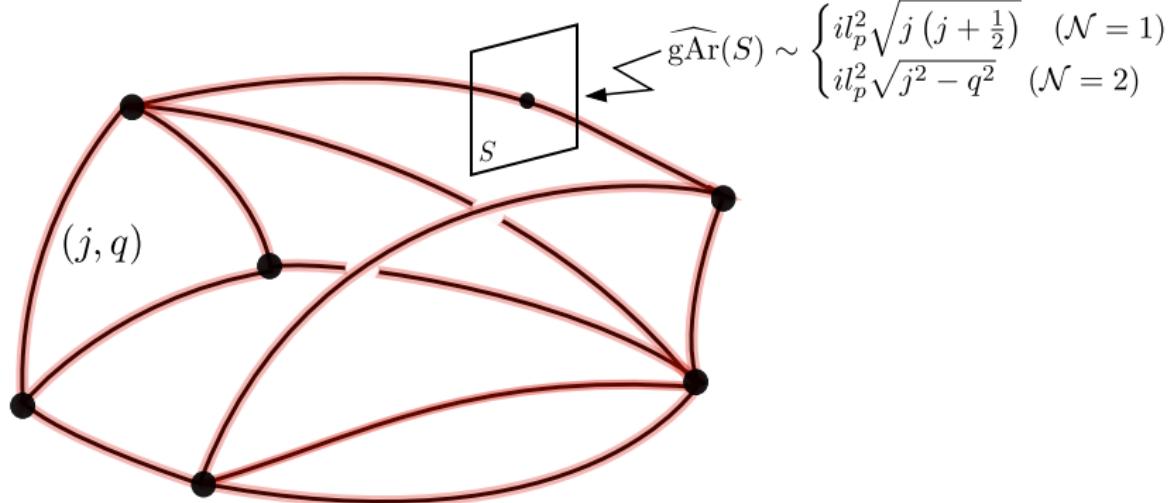
SUSY-invariance: Coupling bulk \leftrightarrow boundary

$$F(\mathcal{A}^+) = -\frac{1}{2L^2} \mathcal{E}$$

Quantum theory: Super spin nets

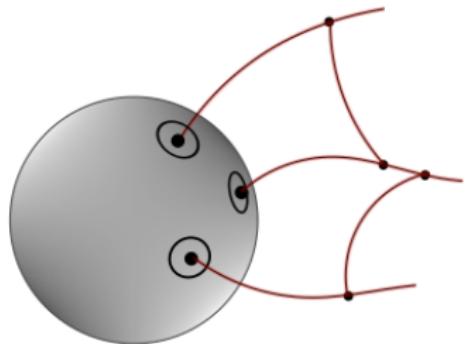
What do we need for quantum theory?

- Holonomies (parallel transport map) ✓
- Hilbert spaces (✓)
- Spin network states ✓



SUSY black holes in LQG

Holst-MM in the presence of inner boundary: [KE '22, KE+HS '22]



- Geometric theory induces super CS on inner boundary
- Gauge group: $G = \text{OSp}(\mathcal{N}|2)$

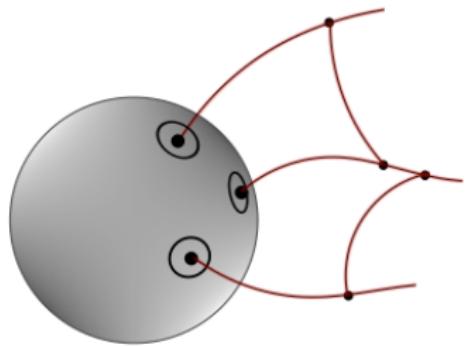
Hilbert space full theory

$$\mathcal{H}_\gamma^{\text{full}} = \mathcal{H}_\gamma^{\text{bulk}} \otimes \mathcal{H}_\gamma^{\text{CS}}$$

Quantum boundary condition

$$1 \otimes \widehat{F}_{\underline{A}}(p) = -\frac{2\pi i}{\kappa k} \widehat{\mathcal{E}}_{\underline{A}}(p) \otimes 1$$

SUSY black holes in LQG



- For $\mathcal{N} = 1$: $G = \mathrm{OSp}(1|2)$
- **Issue:** G complex/non-compact \rightarrow adapt strategy of [Perez et al '14, Noui et al '15]
- Consider compactification $\mathrm{UOSp}(1|2)$ and analytically continue $j \rightarrow -\frac{1}{4} + is$

UOSp(1|2) state counting

$$N = \frac{1}{2\pi} \int_0^\pi d\theta \sin^2(2\theta) \left[4 - n + \sum_{i=1}^p n_i d_{j_i} \frac{\tan(d_{j_i}\theta)}{\tan \theta} \right] \prod_{l=1}^p \left(\frac{\cos(d_{j_l}\theta)}{\cos \theta} \right)^{n_l}$$

\rightarrow analytic continuation

Entropy

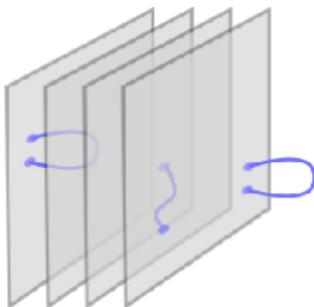
$$S = \ln N = \frac{a}{4} + \dots \rightarrow [\text{Bekenstein '73, Hawking '75}]$$

Section 4

Outlook

Outlook

- Generalization to $\mathcal{N} \geq 3$ (*in particular:* $\mathcal{N} = 4, 8$)
- Limit $L \rightarrow \infty$ (vanishing cosmological constant)
[Concha+Ravera+Rodríguez '19]
- Simplification dynamics extending generalized GCD to super setting
→ Abhay's talk, Madhavan's talk
- (Charged) supersymmetric BHs \leftrightarrow entropy calculations in string theory
[Strominger+Vafa '96, Cardoso et al. '96, KE+HS '22]
- \leftrightarrow Boundaries in string theory and **O $S\!p$** -super Chern-Simons theory
[Mikhailov+Witten '14]



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Quantum theory: Overview

- Quantization: study \mathcal{A}^+ and associated holonomies
- turns out \rightarrow need additional **parametrizing supermanifold \mathcal{S}** to incorporate **anticommutative** nature of fermionic fields

Super holonomies [KE '20+'21]

$$h_e[\mathcal{A}^+] = h_e[A^+] \cdot \mathcal{P}\exp\left(-\oint_e \text{Ad}_{h_e[A^+]^{-1}} \psi^{(\tilde{s})}\right) : \mathcal{S} \rightarrow \mathcal{G}$$

- $\mathcal{A}^+ = A^+ + \psi$ and $e : [0, 1] \rightarrow M \subset \mathcal{M}$
- $h_e[A^+]$: bosonic holonomy associated to A^+

Quantum theory: Overview

- $\mathcal{A}_{\mathcal{S},\gamma}$: set of generalized super connections

$$\begin{array}{ccc} \mathcal{A}_{\mathcal{S}',\gamma'} & \longrightarrow & \mathcal{A}_{\mathcal{S}',\gamma} \\ \downarrow & & \downarrow \\ \mathcal{A}_{\mathcal{S},\gamma'} & \longrightarrow & \mathcal{A}_{\mathcal{S},\gamma} \end{array}$$

- for fixed graph γ induces functor $\mathcal{A}_\gamma : \mathcal{S} \rightarrow \mathcal{A}_{\mathcal{S},\gamma}$
→ **Molotkov-Sachse-type smf.**
- → can study cylindrical functions and invariant measures on \mathcal{A}_γ
- covariance under reparametrization requires **Berezin-type integral** for fermionic d.o.f.
- → induces **Krein structure** on state space \leftrightarrow strong similarities to quantization of fermions in standard LQG

Symmetry reduction

- supersymmetric minisuperspace models [D'Eath + Hughes '88+'92, D'Eath + Hawking + Obregon '93]
- hybrid homogeneous isotropic ansatz for bosonic and fermionic d.o.f.
- in general: fermions not compatible with isotropy
- **But:** in (chiral) LQC can exploit enlarged gauge symmetry!
- → natural interpretation in terms of homogeneous isotropic super connection [KE '20, KE+HS '20]

Symmetry reduced connection

$$\mathcal{A}^+ = c \, \dot{\mathbf{e}}^i T_i^+ + \psi_A \dot{\mathbf{e}}^{AA'} Q_{A'}$$

$\dot{\mathbf{e}}^i$: fiducial co-triad

- **Also:** can show that this is the most general ansatz consistent with reality conditions (contorsion tensor isotropic)

Loop quantum cosmology

- Construction of state space via super holonomies $h_e[\mathcal{A}^+]$ along straight edges of a fiducial cell
- \Rightarrow motivates state space of quantum theory

Hilbert space

$$\mathcal{H} = \overline{H_{AP}(\mathbb{C})}^{\|\cdot\|} \otimes \Lambda[\psi_{A'}]$$

- reality condition in quantum theory can be solved exactly! [Wilson-Ewing '15, KE+HS '20]

Quantum right SUSY constraint

$$\widehat{S}_{A'}^R = \frac{3g^{\frac{1}{2}}}{2\lambda_m} |\nu|^{\frac{1}{4}} \left((\mathcal{N}_- - \mathcal{N}_+) - \frac{\kappa\lambda_m}{6g|\nu|} \widehat{\Theta} \right) |\nu|^{\frac{1}{4}} \widehat{\phi}_{A'}$$

- λ_m : quantum area gap (full theory)

Loop quantum cosmology

- Quantum algebra between left and right SUSY constraint closes and reproduces Poisson algebra exactly!

Quantum algebra

$$[\hat{S}^{LA'}, \hat{S}_{A'}^R] = 2i\hbar\kappa\hat{H} - \frac{\hbar\kappa}{6g^{\frac{1}{2}}|v|^{\frac{1}{2}}} \hat{\pi}_\psi^{A'} \hat{S}_{A'}^R$$

- fixes some of the quantization ambiguities
- semiclassical limit: $\lambda_m \rightarrow 0$ i.e. corrections from quantum geometry negligible
- **Hartle-Hawking state** as solution of constraints \Leftrightarrow [D'Eath '98]

Hartle-Hawking state

$$\Psi(a) = \exp\left(\frac{3a^2}{\hbar}\right)$$