

# Revisiting Loop Quantum Gravity with selfdual variables

Loops' 22, Lyon

Joint work with Hanno Sahlmann, to be published soon

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General relativity in Ashtekar variables: [Ashtekar 1986]

$A_a^i$   $SL(2, \mathbb{C})$  connection: pullback of spacetime connection  $\omega_\mu^I{}_J$  with  
$$*\omega = i\omega$$

$E_i^a$  densitised triads: pullback of  $e^I \wedge e^J$ , with  
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Benefits & expectations:

- no ambiguous Immirzi parameter
- simple Hamilton constraint
- Bekenstein-Hawking entropy [Frodden, Geiller, Noui, and Perez 2014; Han 2014; Achour and Noui 2015]
- coupling of chiral fermions to selfdual connection
- Ashtekar's variables for supergravity [Eder and Sahlmann 2021a]

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In the literature (far from being exhaustive)

- ideas/fundamentals: [Ashtekar 1986; Ashtekar 1987; Jacobson and Smolin 1987; Samuel 1987; Jacobson and Smolin 1988; Ashtekar, Romano, and Tate 1989]
- model systems: [Thiemann and Kastrup 1993; Wilson-Ewing 2015; Eder and Sahlmann 2021b]
- Wick rotation/coherent state transform/complexifier: [Thiemann 1995; Ashtekar, Lewandowski, Marolf, Mourão, and Thiemann 1996; Thiemann 1996; Varadarajan 2019]
- Ashtekar variables for real LQG: [Wieland 2012]
- emergent space-time [Ashtekar and Varadarajan 2021] (Ashtekar's talk)
- spinfoams [Borissova and Dittrich 2022] (Dittrich's talk)

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Classical theory:

- Hamilton formulation of selfdual part of complex Palatini action

$$S = \frac{1}{2\kappa} \int_{\mathcal{M}} \epsilon_{IJKL} e^I \wedge e^J \wedge F^{KL}[\omega]$$

- Poisson relation (“additional factor of i”)

$$\{A_a^i(x), E_j^b(y)\} = \frac{i\kappa}{2} \delta_a^b \delta_j^i \delta(x, y)$$

- spatial metric:  $qq^{ab} = \eta^{ij} E_i^a E_j^b$  with real internal metric  $\eta_{ij} = \eta^{kij}$
- Constraints (real  $N$  and  $N^a$ )

$$NC = -N \frac{\sqrt{2\eta^3} s_e}{i\eta\sqrt{q}\kappa s_E} \epsilon_{kl}^j \eta^{li} E_i^a E_j^b F_{ab}^k$$

Hamilton

$$N^a C_a = \frac{2}{i\kappa} N^a E_i^b F_{ab}^i$$

Diffeo

$$(A \cdot t)^i G_i = -(A \cdot t)^i \frac{2}{i\kappa} D_a(E_i^a)$$

Gauß

## Symplectic structure degenerate

- $\{ \cdot, \cdot \}$  not defined on  $\text{Re}A, \text{Im}A, \text{Re}E, \text{Im}E$
- holomorphic phase space description
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- choice of appropriate measure

## Reality conditions: [Ashtekar, Romano, and Tate 1989]

- spatial metric is real:  $q^{ab} \in \mathbb{R}$  (RCI)
- and stays real:  $\dot{q}^{ab} \in \mathbb{R}$  (RCII)
- adjointness relations / conditions on measure

RCI: reality of the spatial metric

$$qq^{ab} = \eta^{ij} E_i^a E_j^b \in \mathbb{R} \text{ and pos. definite}$$

→ allows two different embeddings of the real theory

$$\eta_{ij} \text{ positive definite} \rightarrow E_i^a \in \mathbb{R}$$

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## RCII: spatial metric stays real

$$(qq^{ab}) \cdot = 2\eta^{ij} E_i^{(a} \dot{E}_j^{b)} \in \mathbb{R} \quad \text{with derivative} \quad \{\cdot, C[N]\}$$

→ condition  $A_a^i = \Gamma_a^i + iK_a^i$  with

$$\Gamma_a^i(E) \text{ spin connection}$$

$$K_a^i = (A_a^i - \Gamma_a^i)/i \text{ real}$$

→ RCI  $\Rightarrow \Gamma_a^i$  is rotation part of  $A_a^i$

→ RCII:  $K_a^i$  boost part of  $A_a^i$

## Operators:

- $SL(2, \mathbb{C})$  holonomies  $h_e$
- fluxes  $E_f(S)$ : holomorphic invariant vector fields
- CCR: additional factor of „ $i$ “

$$L_i \Psi(g) = \left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \Psi(g e^{\epsilon \tau_i}) \quad \& \quad R_i \Psi(g) = \left. \frac{d}{d\epsilon} \right|_{\epsilon=0} \Psi(e^{-\epsilon \tau_i} g)$$
$$[h_e, E_f(S)] = \frac{\kappa}{4} \kappa(e, S) \begin{cases} h_e f(p) & p \text{ source of } e \\ -f(p) h_e & p \text{ target of } e \end{cases}$$

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**Holomorphic Hilbert space:**  $\mathcal{L}^2(SL(2, \mathbb{C}), \mu(g, \bar{g}) dg_H) \cap \mathcal{H}(SL(2, \mathbb{C}))$

- inner product

$$\langle \phi, \psi \rangle := \int_{SL(2, \mathbb{C})} dg_H \mu(g, \bar{g}) \overline{\phi(g)} \psi(g)$$

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- holom. derivatives in  $\{\cdot, \cdot\} \rightarrow$  holom. invariant vector fields in  $[\cdot, \cdot]$
- no adjointness relations  $\rightarrow$  not usual HF algebra  $\rightarrow$  no Fleischhack-LOST thm.

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**Idea:** use freedom in  $\mu(g, \bar{g})$  for RC implementation

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**Measure-ansatz:**  $\mu(g, \bar{g}) = \mu(gg^\dagger)$ , invariant under compact subgroup  $SU(2)$  & cyclic, e. g.  $\mu(g, \bar{g}) = \mu(\text{tr}(gg^\dagger))$



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ONB for  $\mathcal{HL}^2(SL(2, \mathbb{C}))$

- Weyl's unitary trick: holom.  $SL(2, \mathbb{C})$  irreps from  $SU(2)$  irreps  $\Pi_j(g)_{mn}$ ,  $g \in SL(2, \mathbb{C})$
- RCI +  $SU(2)$  orthogonality relation:

$$\langle \Pi_{jmn}, \Pi_{j'm'n'} \rangle = \frac{\mu_j}{2j+1} \delta_{jj'} \delta_{mm'} \delta_{nn'} \quad \text{with} \quad \mu_j := \int d|g| \mu(|g|^2) \text{tr}(\Pi_j(|g|^2))$$

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**Graph Hilbert space**  $\mathcal{H}_\gamma = \mathcal{L}^2(\bar{\mathcal{A}}_\gamma, d\mu_\gamma(A))$  + product measure

→  $SL(2, \mathbb{C})$  spin networks

→ for  $\mu_0 = 1$  cylindrically cons. measure on  $\mathcal{H} = \mathcal{L}^2(\bar{\mathcal{A}}, d\Omega)$

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**First reality condition**  $\Rightarrow$  new type of consistent measures on  $\bar{\mathcal{A}}$ , e. g.  $\mu(g, \bar{g}) \sim \exp\{-\text{tr}(gg^\dagger)\}$

# Reality condition II

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Reality condition for holonomies

- short edges
- adjointness relation for holonomies
- $\Gamma(E)$  complicated function of  $E$

$$\overline{A_a^j - \Gamma_a^j} = -(A_a^j - \Gamma_a^j) \quad \rightarrow \quad A_a^j + \overline{A_a^j} = 2\Gamma_a^j$$
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Consistency with CCR:  $E_i^\dagger = \alpha E_i \rightarrow 0 = [h_e(h_e^{-1})^\dagger, (E)^2] = (\alpha - 1)(\dots) ? \rightarrow$  singles out  $E_i^\dagger = E_i$

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## Known problem! [Wilson-Ewing 2015]

- LQC with selfdual variables
- RCI inconsistency: selfadjoint  $E$

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## Questions / Where to go from here?

- how strong is this obstruction?
- incompatibility of RC's and quantum theory based on holonomies?
- different formulation of RCII?

# What has been done:

## GR in selfdual variables & reality conditions

- degenerate symplectic structure, hence
  - holomorphic phase space
  - holomorphic quantisation
  - reality conditions are no constraints
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Selfdual variables are still interesting to look at and the reality conditions might even be implementable!

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# THANK YOU!

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