INSTITUTO DE FÍSICA—UFRGS Programa de Pós-Graduação em Física



Holonomy Observables for the Turaev–Viro model

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Current goals (work in progress — ~ 2 months behind schedule):

- 1. Extend the definition of holonomy observables for the Ponzano–Regge model (Barrett and Hellmann 2012) to the Turaev–Viro model (\sim)
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Why?

- Quantum groups in LQG are not (yet) fully understood
- Work on the canonical formulation points to the origins of quantum groups (Dupuis et al. 2020; Noui, Perez and Pranzetti 2010, 2011, 2012, 2014), but the root of unity case is still murky (Rennert 2018)
- Developing a "holonomy representation" of the TV model is in parallel to the problem above (covariant vs. canonical)
- Adapt tools that use holonomies to the TV model, e.g., coupling of particles

The Turaev–Viro model as a TQFT

The Turaev–Viro model is

- \blacksquare a TQFT built from (the category of representations of) the quantum group $u_q(\mathfrak{sl}(2,\mathbb{C}))$
- \blacksquare a spinfoam model for 3d Euclidean LQG with $\Lambda>0$ and

$$q = \exp\left(i\frac{8\pi G\hbar\sqrt{\Lambda}}{c}\right)$$

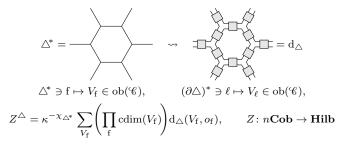
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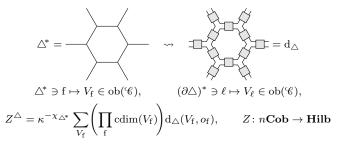
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TV model is the theory with $\mathscr{C} = {}_{u_q(\mathfrak{sl}(2,\mathbb{C}))}\mathcal{M}$ at q a root of unity

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When q is a root of unity, its representation theory is nontrivial \longrightarrow useful to deal with a reduced algebra

$$\begin{aligned} u_q(\mathfrak{sl}(2,\mathbb{C})) &= U_q(\mathfrak{sl}(2,\mathbb{C}))/I; \qquad K^{p'} = 1, \ E^{p'} = 0, \ F^{p'} = 0, \\ q^p = 1, \qquad p' = \begin{cases} p & \text{odd } p \\ p/2 & \text{even } p \end{cases} \end{aligned}$$

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Turaev-Viro model:

$$Z_{\rm TV} = \left(-\frac{2r}{(q-q^{-1})^2}\right)^{-P} \sum_{j_{\rm f}} \left(\prod_{\rm f} (-1)^{2j_{\rm f}} [2j_{\rm f}+1]_q\right) \prod_{\rm v} (-1)^{J_{\rm v}} \begin{cases} j_{\rm v1} & j_{\rm v2} & j_{\rm v3} \\ j_{\rm v4} & j_{\rm v5} & j_{\rm v6} \end{cases} ,$$
$$q = \exp\left(\frac{i\pi}{r}\right) = \exp(i\sqrt{\Lambda}), \qquad \bigsqcup_{V_{\rm f}}^{V_{\rm f}} = \Delta_{V_{\rm f}} \circ \left(\operatorname{id}_{V_{\rm f}} \otimes \int_{u_q(\mathfrak{sl}(2,\mathbb{C}))}\right)$$

The TV model has even $p = 2r \implies p' = r$

Holonomy observables

Holonomy observables $F(g_e)$ computed for the PR model in (Barrett and Hellmann 2012)

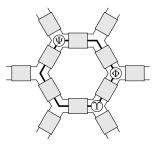
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These are $q\to 1$ limits of spin observables for the TV model: Wilson networks in the TQFT formalism

$$\langle \Gamma_W \rangle = \frac{Z^{\bigtriangleup}[\Gamma_W]}{Z^{\bigtriangleup}}, \quad (\Gamma_W, V_e, \Phi_v), \ e \mapsto V_e, \ v \mapsto \Phi_v$$



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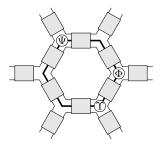
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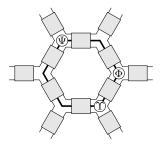
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How to extend to the TV model?



Extension to Turaev–Viro

The argument

 \blacksquare If $\mathscr{C}\cong \mathscr{C}',$ the TQFTs are isomorphic

If
$$H \cong H'$$
, ${}_{H}\mathcal{M} \cong {}_{H'}\mathcal{M}$

- For $\mathfrak{su}(2)$, $U(\mathfrak{su}(2)) \cong (U(\mathfrak{su}(2)))^* = \mathcal{O}(\mathrm{SU}(2))$ (Fourier transform between $\mathfrak{su}(2)$ and $\mathrm{SU}(2)$; between spins and holonomies)
- \blacksquare Finding a 1-to-1 duality for $u_q(\mathfrak{sl}(2,\mathbb{C}))$ allows for the definition of the holonomy observables
- Expectation: $\langle F \rangle = \langle \Gamma_W \rangle$ (gauge invariant case)

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For odd p ($q^p = 1$): (Aziz and Majid 2019; Glushchenkov and Lyakhovskaya 1997; Lyubashenko and Majid 1994):

$$\begin{split} (u_q(\mathfrak{sl}(2,\mathbb{C})))^* &= o_q(\mathrm{SL}(2,\mathbb{C})) = \mathcal{O}_q(\mathrm{SL}(2,\mathbb{C}))/J\\ a^p &= 1 = d^p, \quad b^p = 0, \quad c^p = 0\\ \mathscr{F} : o_q(\mathrm{SL}(2,\mathbb{C})) \to u_q(\mathfrak{sl}(2,\mathbb{C})) \text{ is invertible} \end{split}$$

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For even p (the TV case), things are (unfortunately) not so simple...

What is next?

- \blacksquare If possible, gain more information on the case of even p
- If impossible, work on simpler cases (e.g., $\Lambda < 0 \implies q \in \mathbb{R}$)
- Compute the observables explicitly
- Study the contributions to the canonical approach and compare results

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Thank you!

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