

# Directly Probing the Higgs-top Coupling at High Scales

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# Why

# Motivation

## Top Yukawa, $y_t$

- $(y_t)$ , at  $\sim 1$ , is the **strongest** interaction of the Higgs boson in SM and hence most **sensitive** to BSM physics.
- It is crucial to the stability of SM vacuum during EWSB; has the dominant contribution to quantum corrections to Higgs mass etc.
- Precise measurement of  $y_t$  can be fundamental to pin down possible NP.
- HL-LHC projects measurement of  $y_t$  to an accuracy of  $\delta y_t \leq \mathcal{O}(4)\%$ .

## Probing at High Scales

- Current measurements are at EW scale,  $Q \sim v$ .
- BSM effects scale as  $\left(\frac{Q}{\Lambda}\right)^{n>0}$ ;  $\Lambda = \text{NP scale}$ .
- **NP effects can be enhanced by exploring top Yukawa at high scales.**

# How

# Direct probe of top Yukawa at high scales

- Previously some proposals were made to study off-shell Higgs in the  $gg \rightarrow h^* \rightarrow VV$  channel to probe Higgs physics at high scales<sup>2</sup>.
- In this work we **directly** probe Higgs-top coupling at **high scales** using **on-shell Higgs production with high**  $p_{T,h}$ .
- We look at the  $pp \rightarrow t\bar{t}h$  channel, where at high scales we can simultaneously enhance NP effects and suppress backgrounds.
- The new physics sensitivity is parametrized in terms of the effective field theory framework, and a non-local Higgs-top coupling form-factor.

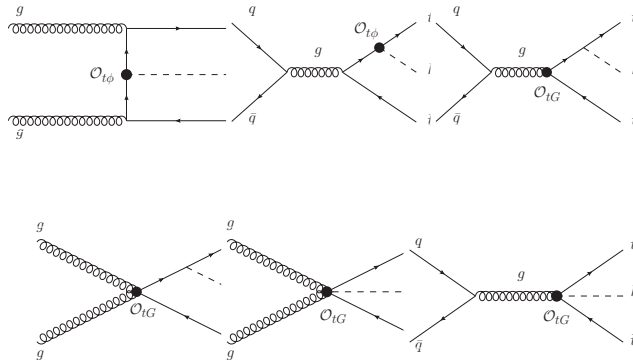
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<sup>2</sup>Gonçalves, Han, Mukhopadhyay (2018)

Gonçalves, Han, Leung, Qin (2020)

# EFT framework


- EFT is usually parameterized as,  $\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$
- Focusing on 2 fermion operators, we study the following operators.
  - $\mathcal{O}_{t\phi} = (H^\dagger H)(\bar{Q}t)\tilde{H} + \text{h.c.}$   
which simply **rescales the SM top Yukawa coupling**, and
  - $\mathcal{O}_{tG} = ig_s(\bar{Q}\tau^{\mu\nu}T_A t)\tilde{H}G_{\mu\nu}^A + \text{h.c.}$   
the **chromo-dipole moment of top quark**. It modifies  $g\bar{t}t$  vertex and introduces new vertices  $gg\bar{t}t$ ,  $gt\bar{t}h$ ,  $gg\bar{t}th$ .



# Higgs-top coupling form-factor

- The top-quark Yukawa coupling has a special role in the naturalness problem, displaying the dominant quantum corrections to the Higgs mass.
- Many well motivated scenarios consider the **top quark and Higgs as composite particles** arising from strongly interacting new dynamics at a scale  $\Lambda$ .<sup>3</sup>
- In such scenarios, **top Yukawa has momentum-dependent form-factor rather than a point-like interaction.**
- Motivated by nucleon form-factor we adopt the ansatz.  
$$\Gamma(Q^2/\Lambda^2) = \frac{1}{(1+Q^2/\Lambda^2)^n}$$
, with  $n=2$  the dipole form-factor (corresponding to exponential spatial distribution).

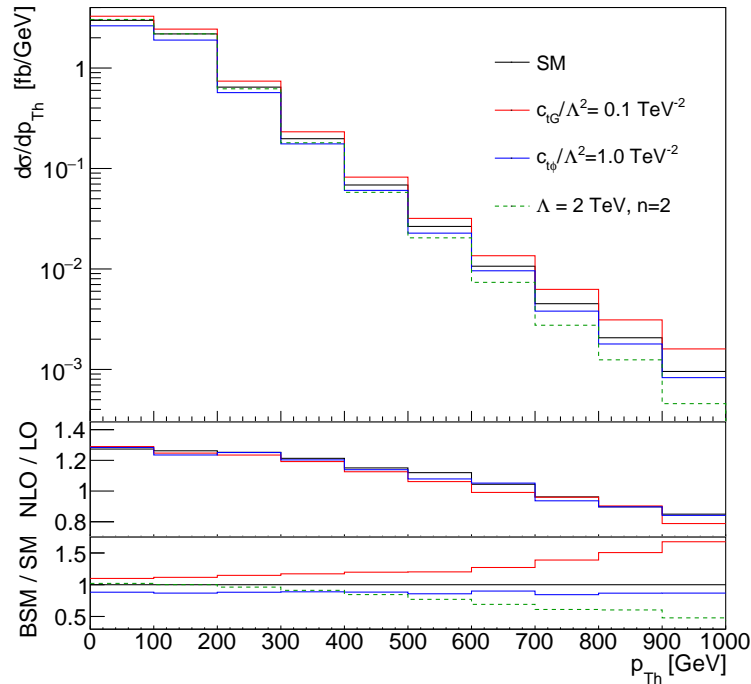
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<sup>3</sup>Pomarol, Riva (2012); Panico, Wulzer (2015); Liu, Low, Wagner (2017) etc. 

# New Physics Effects

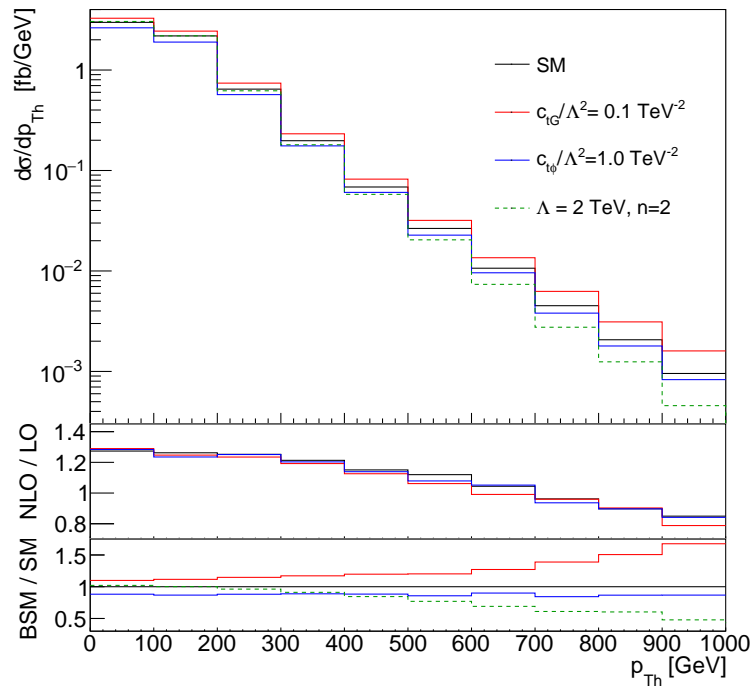


# Higher order and BSM effects



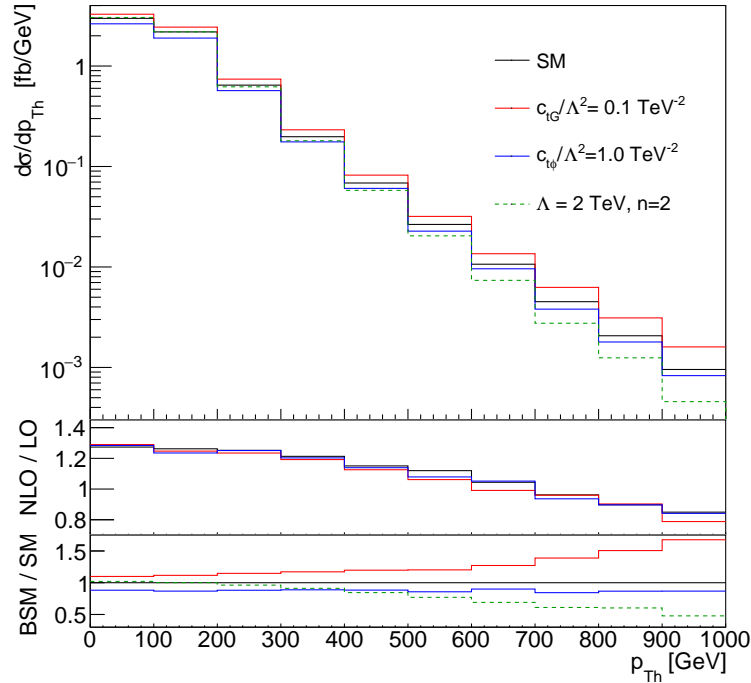
- The K-factor decreases with  $p_{T,h}$  and hence cannot be captured by a global NLO K-factor.

# Higher order and BSM effects



- We also see the enhancement arising from  $\mathcal{O}_{tG}$  operator is scale ( $p_{T,h}$ ) dependent.

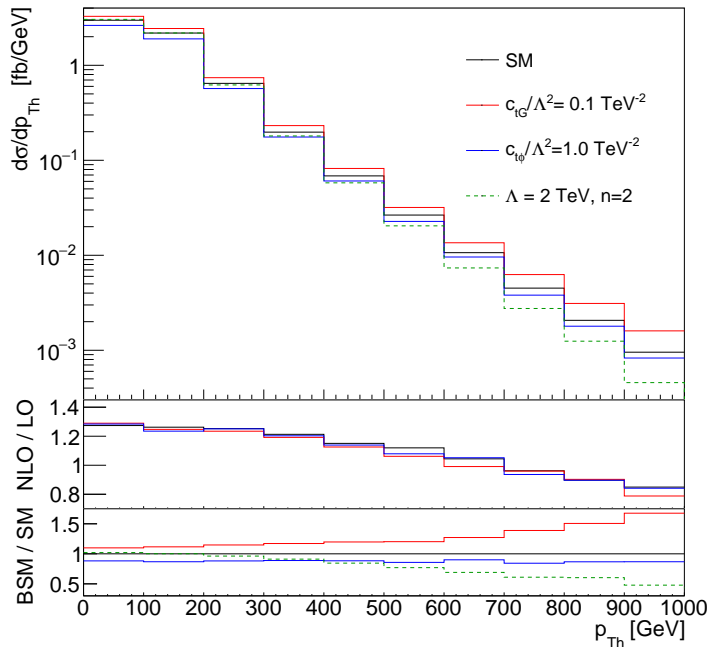
# Higher order and BSM effects



- In the  $t\bar{t}h$  process,  $\mathcal{O}_{t\phi}$  only contributes to a shift of the top Yukawa resulting in a flat rescaling w.r.t the SM cross-section.

# Higher order and BSM effects

$$\Gamma(Q^2/\Lambda^2) = \frac{1}{(1+Q^2/\Lambda^2)^n}, \quad Q = p_{T,h}$$



- The form-factor scenario displays a depletion in cross-section at higher  $p_{T,h}$ , due to the dipole suppression ( $n=2$ ).

# Full Analysis

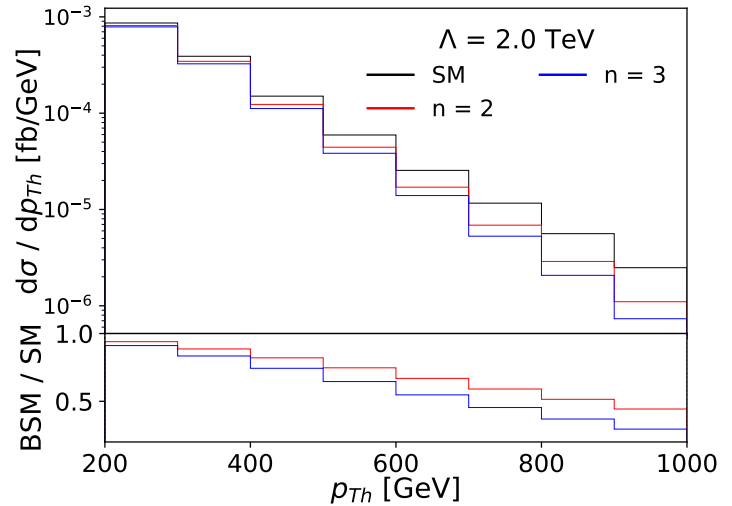
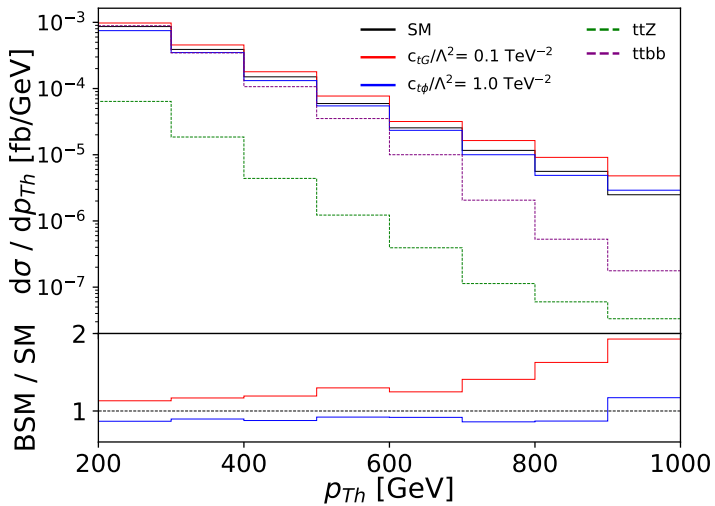
# Full Analysis - Using jet substructure

- Our signal is  $pp \rightarrow t\bar{t}h$  with the  $h \rightarrow b\bar{b}$  and top-quark pair decaying leptonically. Final state is 4 b-tagged jets and 2 opp. sign leptons.
- Leading backgrounds are  $t\bar{t}b\bar{b}$  and  $t\bar{t}Z (l^+l^-)$ .
- Use of jet sub-structure techniques can effectively suppress backgrounds<sup>4</sup>.
  - We require at least 1 boosted fat jet ( $R=1.2$ ) with  $p_{TJ} > 200$  GeV, Higgs tagged with the BDRS algorithm.
  - Outside the fat jet we require 2 b-tagged jets ( $R=0.4$ ) with  $p_{Tj} > 30$  GeV; and Higgs candidate has a mass close to 125 GeV.

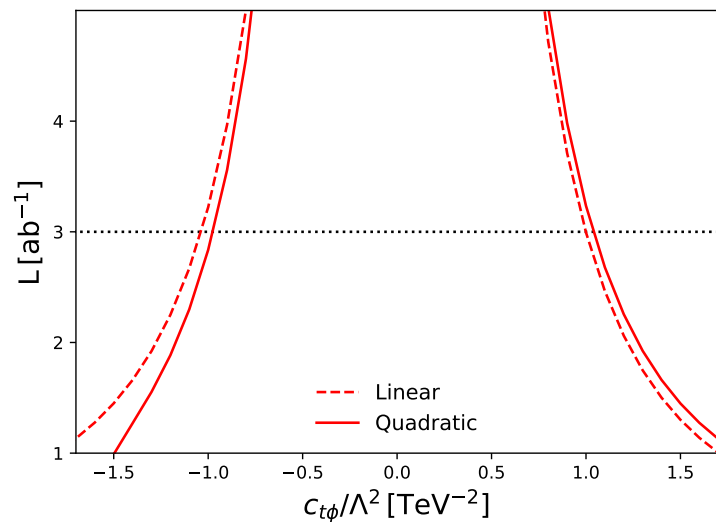
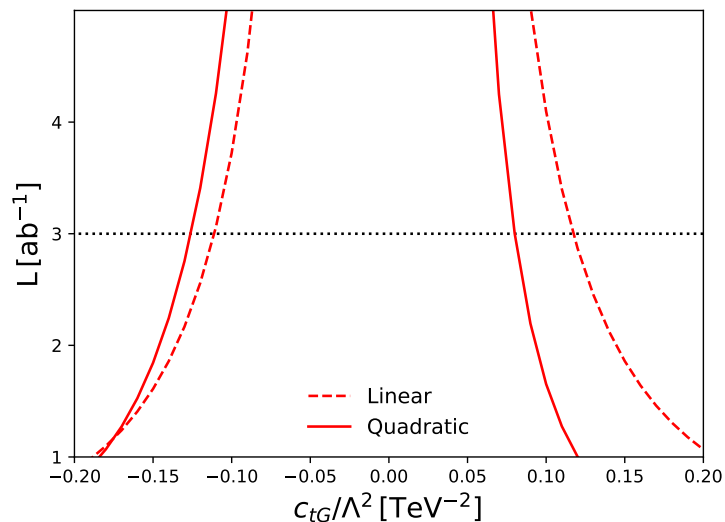
cuts	$t\bar{t}h$	$t\bar{t}b\bar{b}$	$t\bar{t}Z$
BDRS $h$ -tag, $p_{T\ell} > 10$ GeV, $ \eta_\ell  < 3$ , $n_\ell = 2$	3.32	6.35	1.02
$p_{Tj} > 30$ GeV, $ \eta_j  < 3$ , $n_j \geq 2$ , $n_b=2$	0.72	1.97	0.22
$ m_h^{\text{BDRS}} - 125  < 10$ GeV	0.15	0.14	0.009

<sup>4</sup>Buckley, Gonçalves (2015)

# Full Analysis - EFT and form-factor

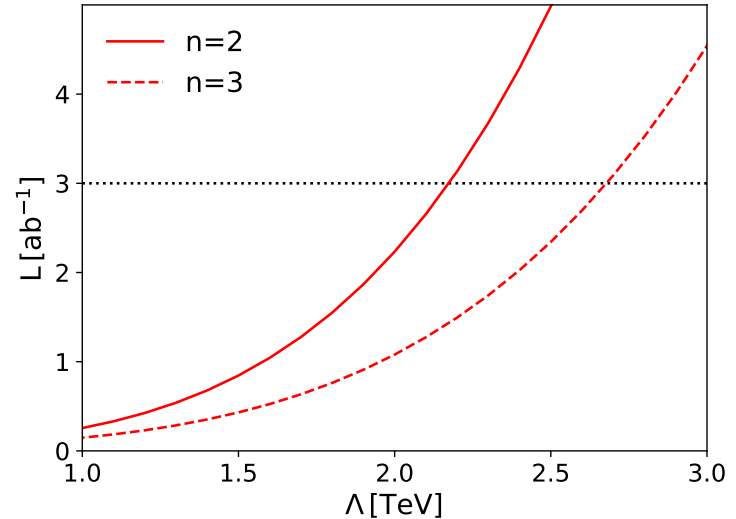
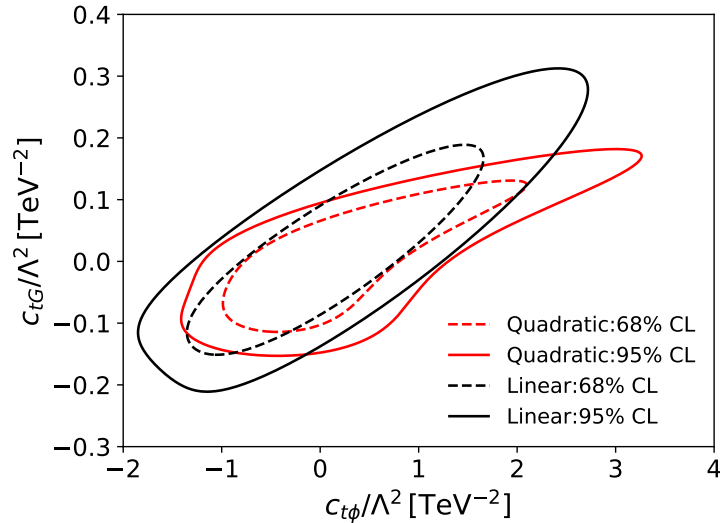


# Results - Individual



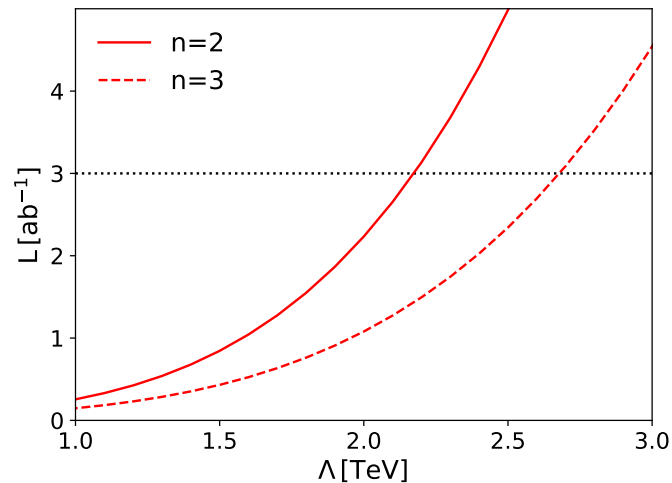
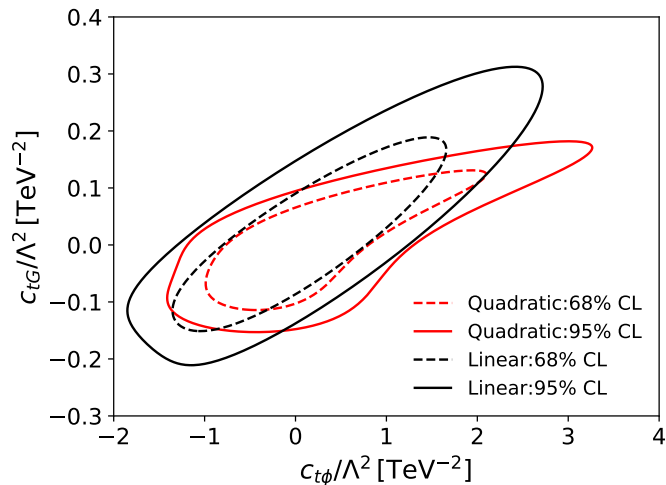


# Results - Combined and form-factor



Linear order calculation is robust to quadratic effects.

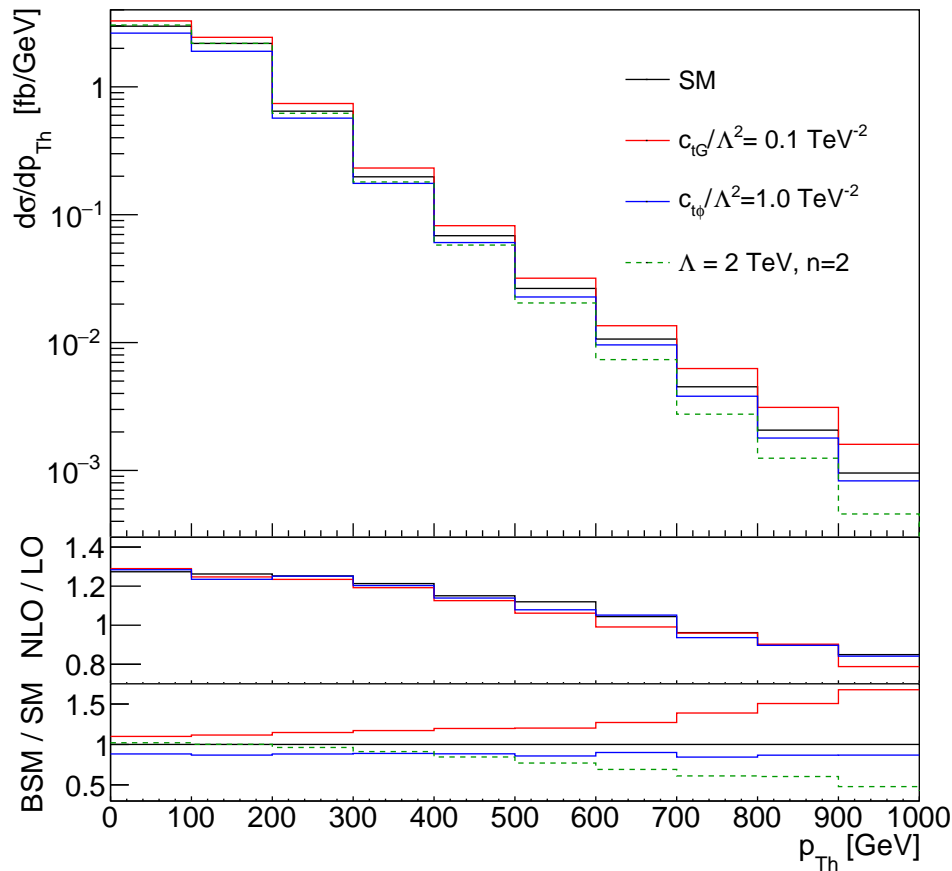
# Results - Combined and form-factor



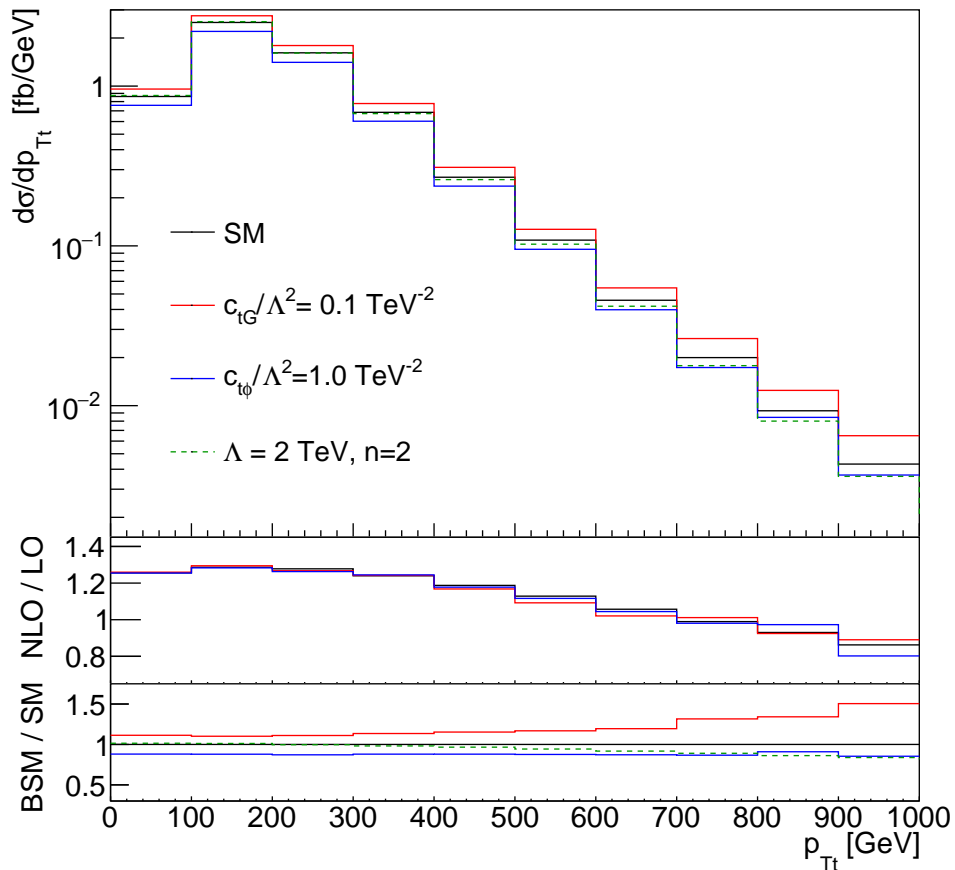
	$c_i/\Lambda^2$ [TeV $^{-2}$ ] 95% CL range	$\Lambda/\sqrt{c_i}$ [TeV] BSM scale
$c_{tG}$ (linear)	[-0.11 , 0.12]	2.9
$c_{tG}$ (quadratic)	[-0.13 , 0.08]	2.8
$c_{t\phi}$ (linear)	[-1.04, 1.00]	1.0
$c_{t\phi}$ (quadratic)	[-0.97 , 1.04]	1.0
n=2	*	2.1
n=3	*	2.7

- The HL-LHC promises unprecedented precision in the top Yukawa measurement allowing one to constrain NP.
- Using the boosted Higgs regime (in the  $t\bar{t}h$  channel) and jet substructure techniques, we show how Higgs-top coupling can be *directly* probed at high scales.
- Sensitivity to new physics is presented within the EFT framework and also the Higgs-top form-factor.

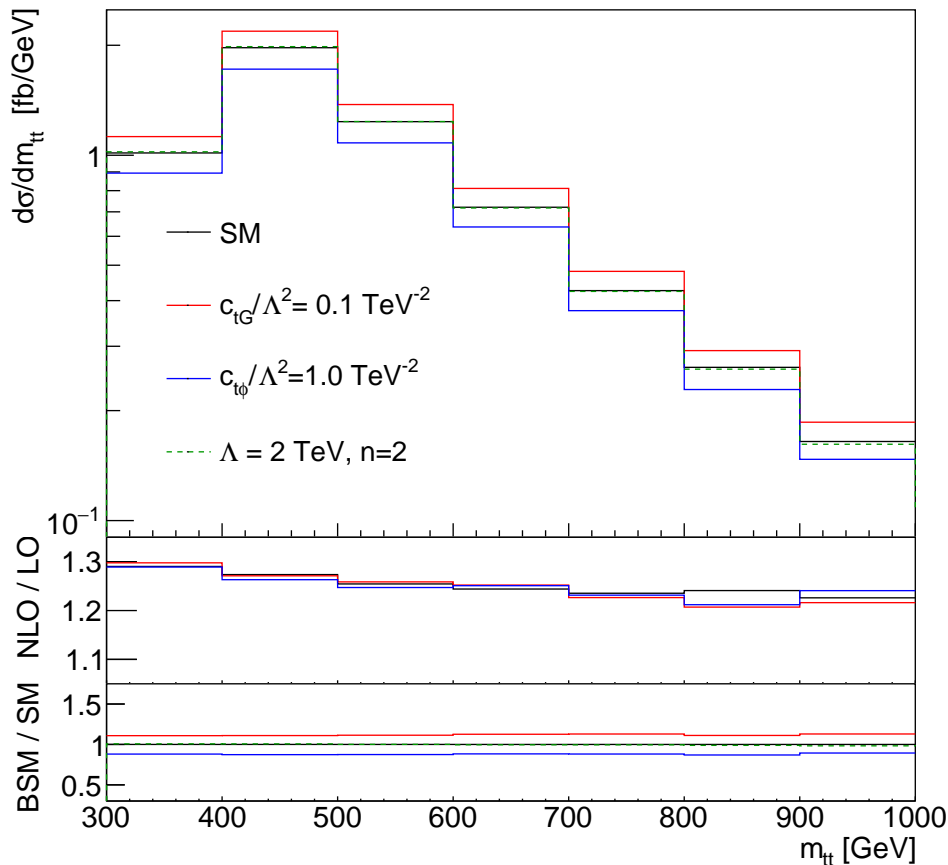
# Backup Slides - More kinematic distributions, $p_{T,h}$



# Backup Slides - More kinematic distributions, $p_{T,t}$



# Backup Slides - More kinematic distributions, $m_{tt}$



# Backup Slides - More kinematic distributions, $m_{th}$

