Università degli Studi di Milano

INFN Istituto Nazionale di Fisica Nucleare

Exact mixed NNLO QCD-EW corrections to the Neutral Current Drell-Yan process

Alessandro Vicini

University of Milano, INFN Milano

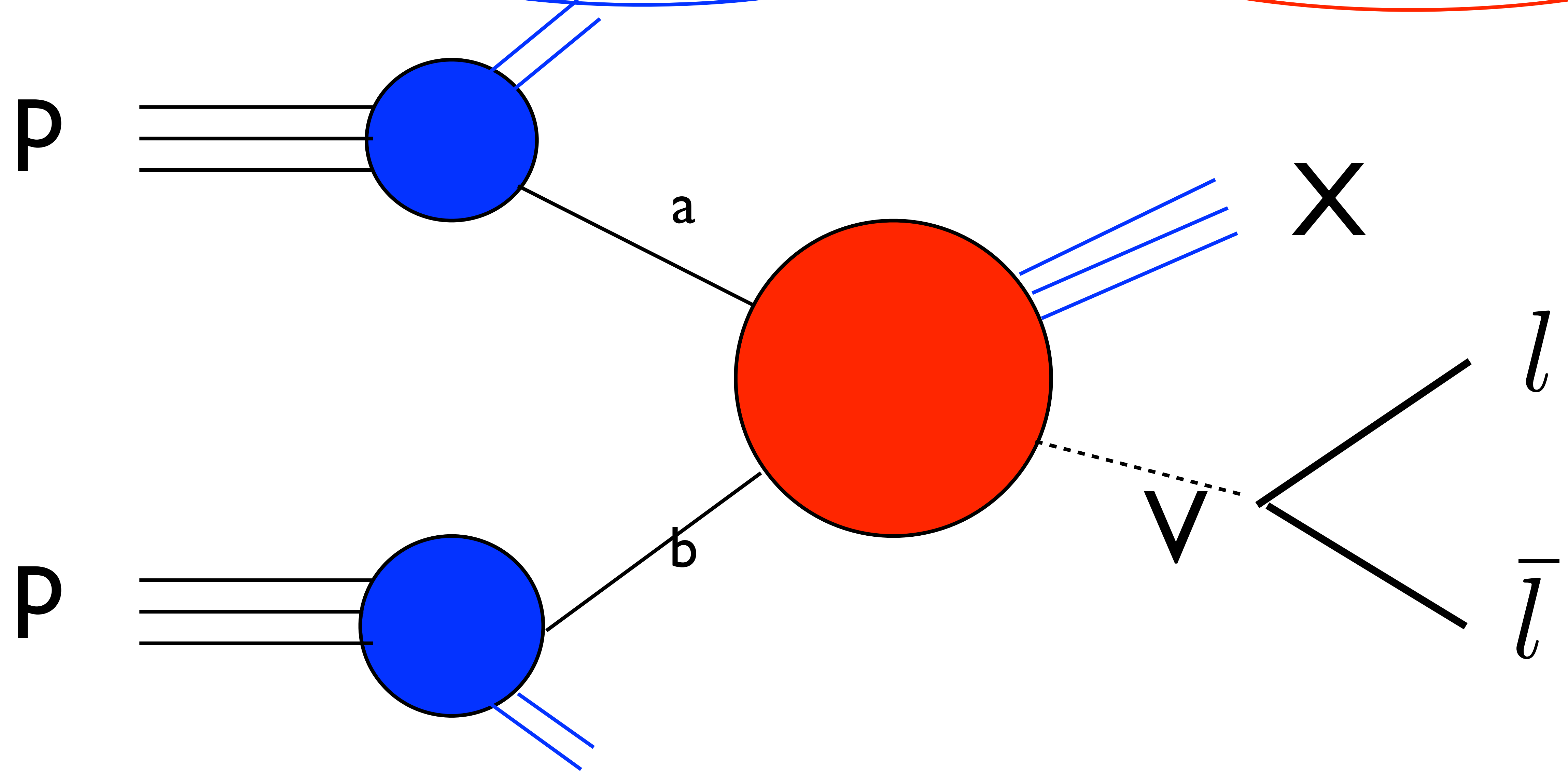
SM@LHC 2022, April 11th 2022

in collaboration with: T.Armadillo, R.Bonciani, F.Buccioni, L.Buonocore, S.Devoto, M.Grazzini, S.Kallweit, N.Rana, F.Tramontano,

Drell-Yan results: arXiv:2106.11953, arXiv:2201.01754 on-shell Z results arXiv:2007.06518, arXiv:2111.12694

Lepton-pair Drell-Yan production at hadron colliders

$$\sigma(P_1,P_2;m_V)=\sum_{a,b}\int_0^1 dx_1dx_2\, f_{h_1,a}(x_1,M_F)f_{h_2,b}(x_2,M_F)\,\hat\sigma_{ab}(x_1P_1,x_2P_2,\alpha_s(\mu),M_F)$$

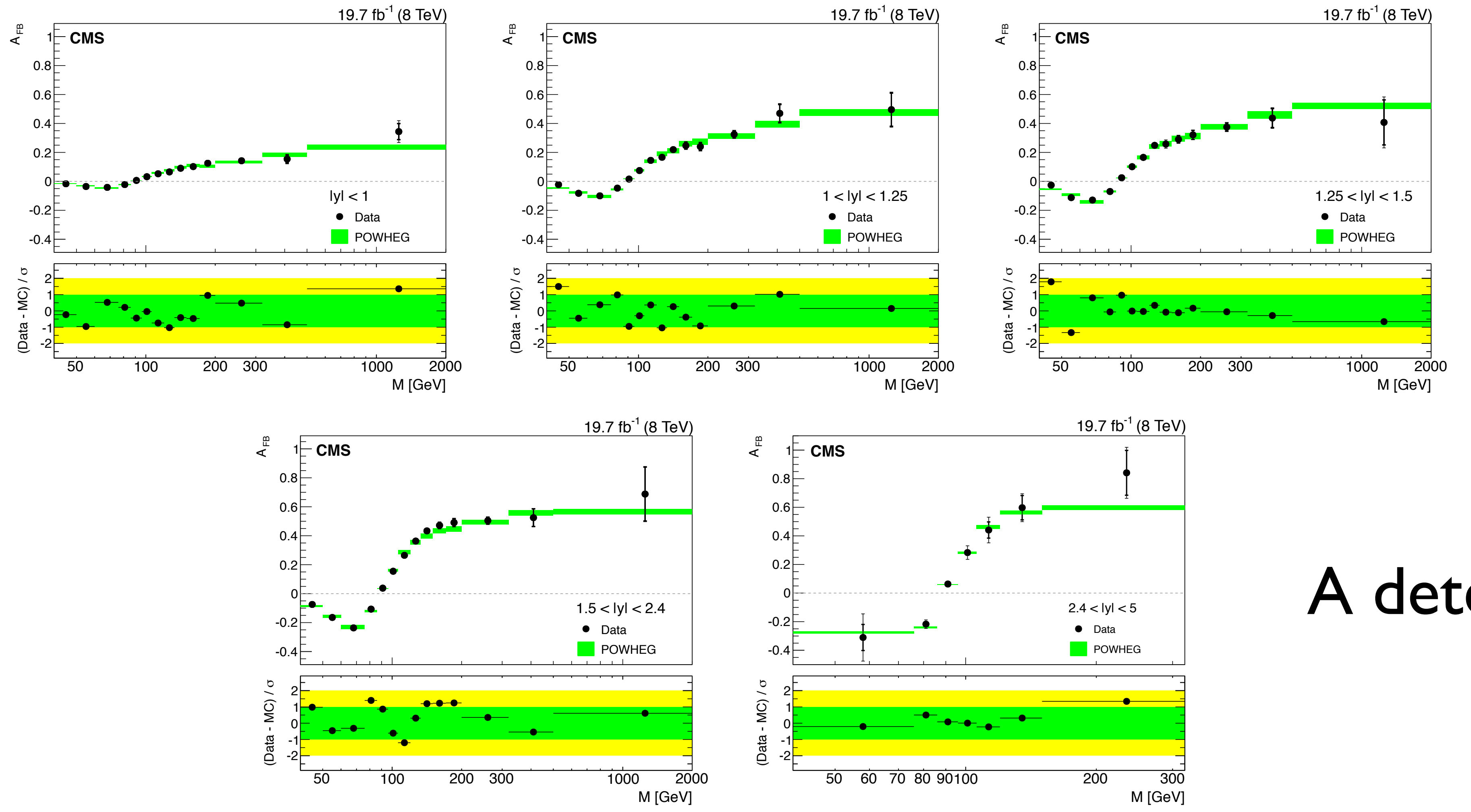


The factorisation theorems guarantee the validity of the above picture up to power correction effects

The interplay of QCD and EW interactions appears both in the partonic cross section and in the proton PDFs

Motivations: towards per mille physics 😲 at the LHC ?

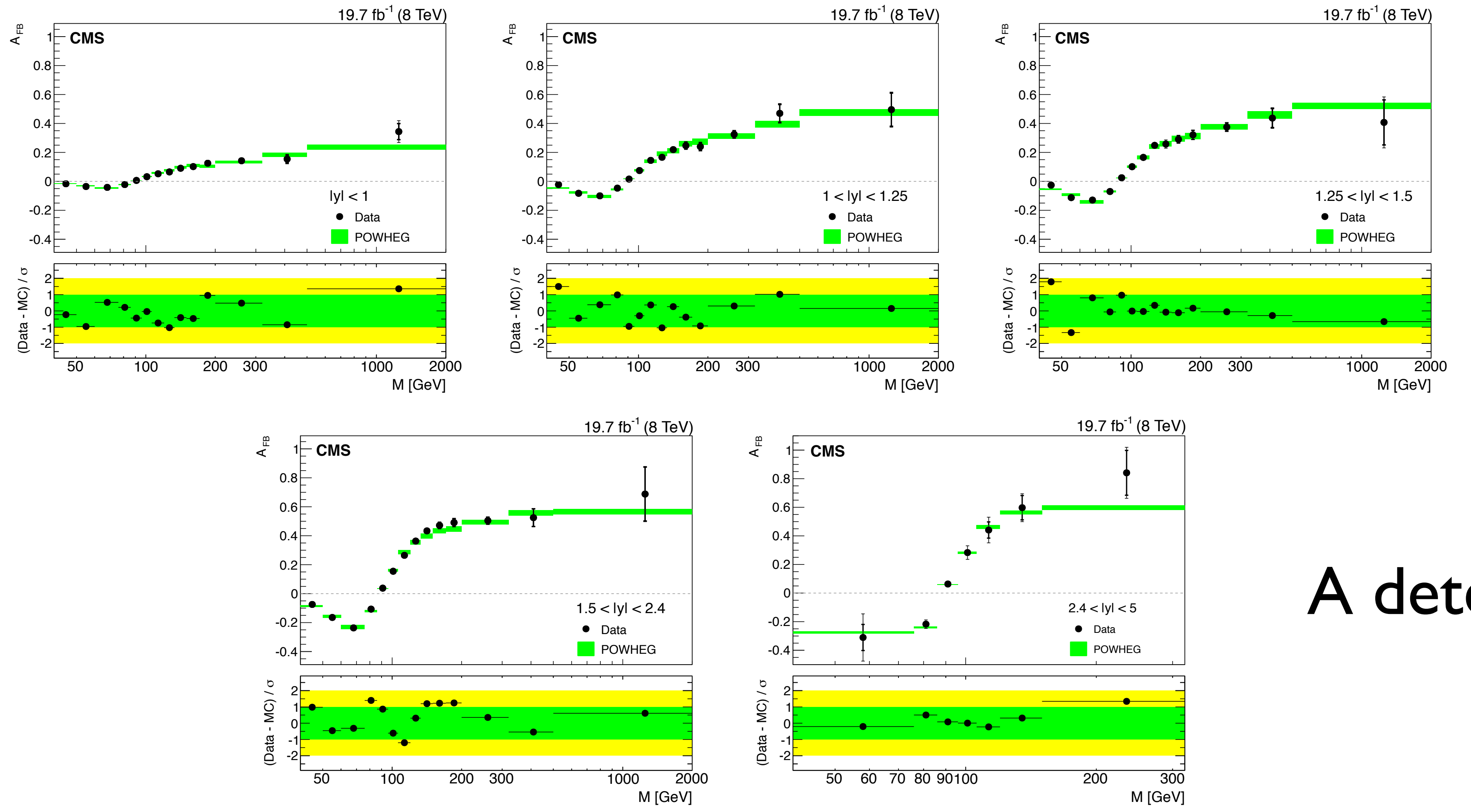
Motivations: towards per mille physics 😲 at the LHC ?



Channel	Not constraining PDFs	Constraining PDFs
Muons	0.23125 ± 0.00054	0.23125 ± 0.00032
Electrons	0.23054 ± 0.00064	0.23056 ± 0.00045
Combined	0.23102 ± 0.00057	0.23101 ± 0.00030

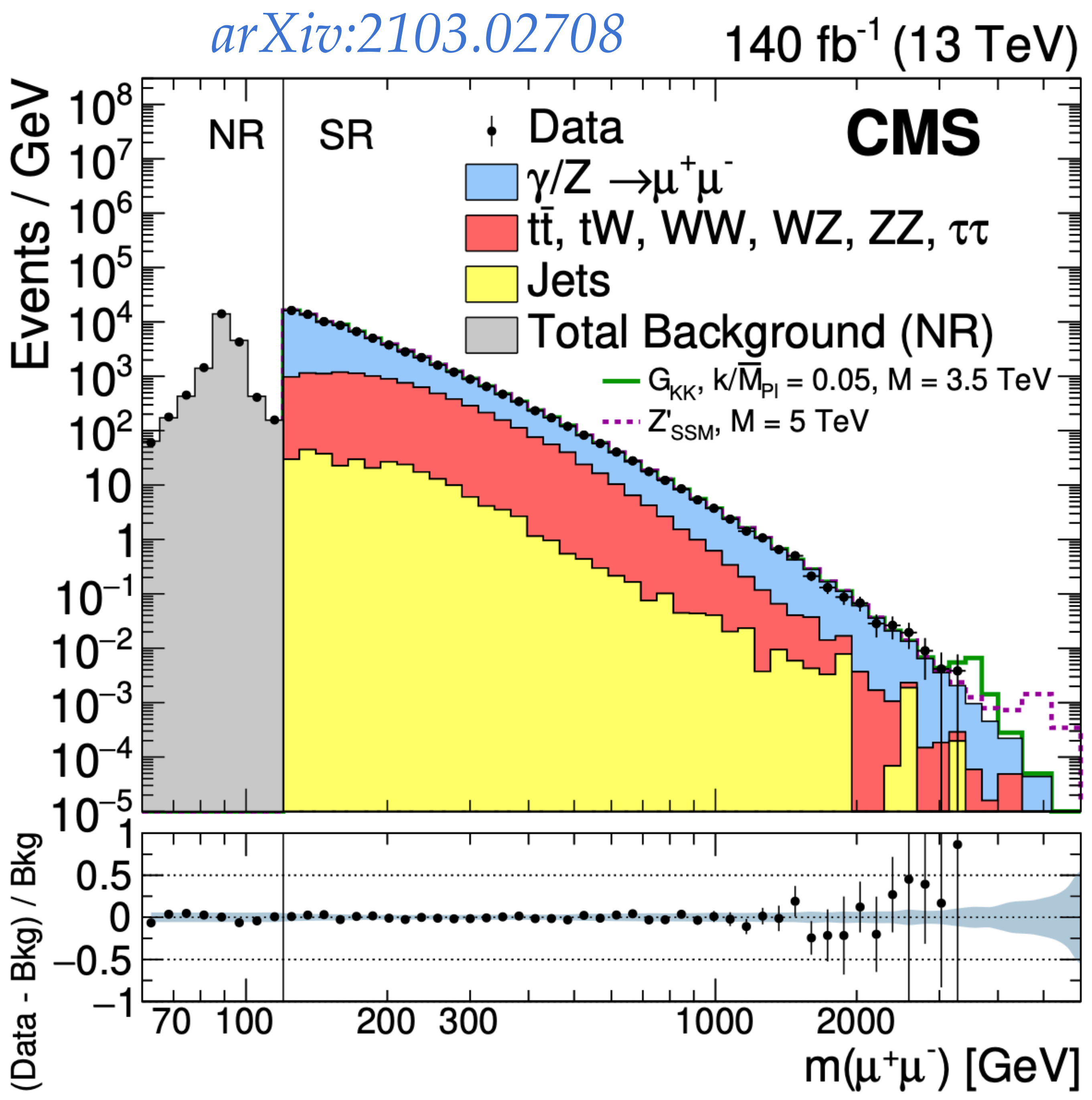
A determination of $\sin^2\theta_{eff}^{lep}$ competitive with the LEP results (0.23152(16)) is becoming possible

Motivations: towards per mille physics 😲 at the LHC ?



Channel	Not constraining PDFs	Constraining PDFs
Muons	0.23125 ± 0.00054	0.23125 ± 0.00032
Electrons	0.23054 ± 0.00064	0.23056 ± 0.00045
Combined	0.23102 ± 0.00057	0.23101 ± 0.00030

A determination of $\sin^2\theta_{eff}^{lep}$ competitive with the LEP results (0.23152(16)) is becoming possible

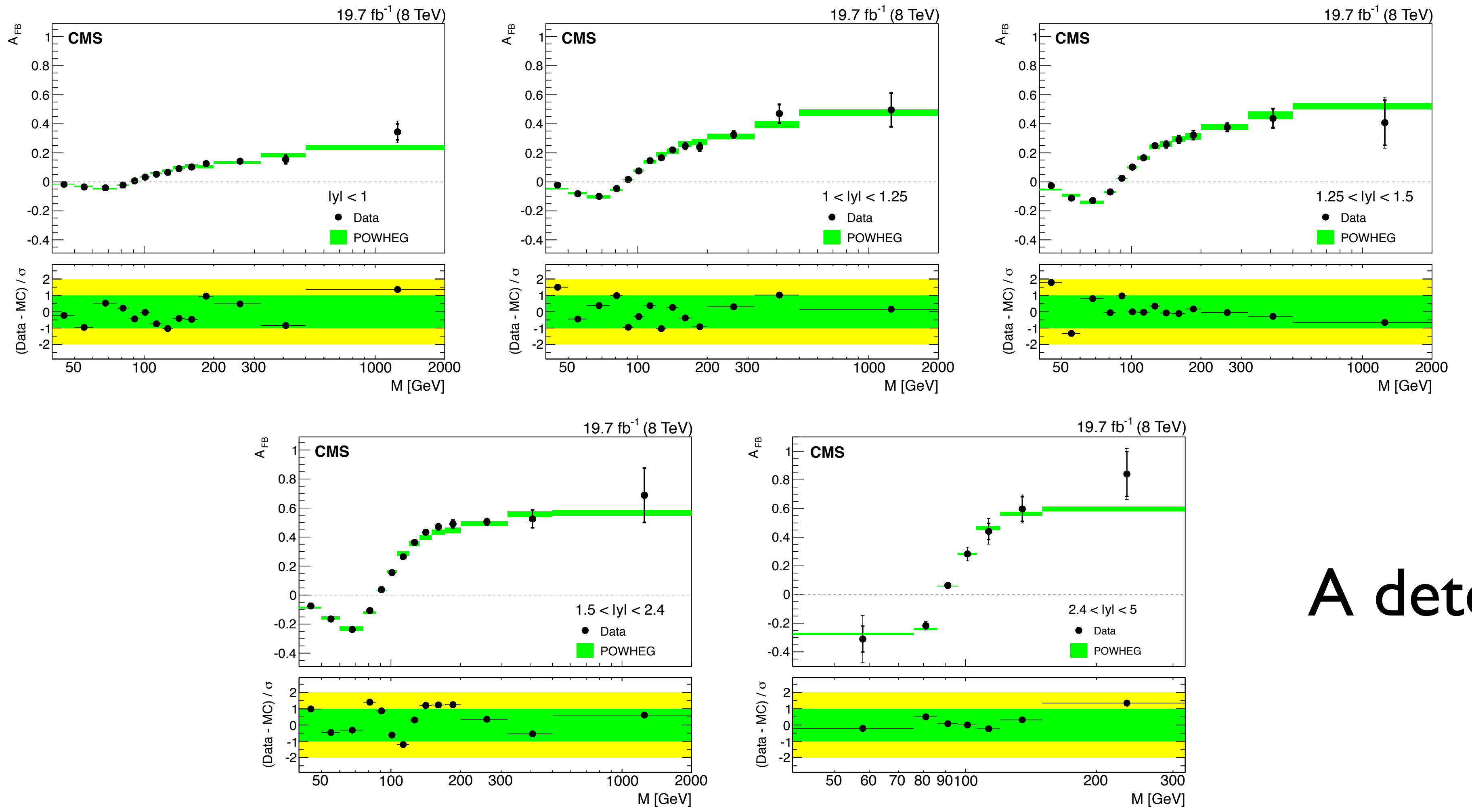


mass window [GeV]	stat. unc. 140fb^{-1}	stat. unc. 3ab^{-1}
$600<m_{\mu\mu}<900$	1.4%	0.2%
$900<m_{\mu\mu}<1300$	3.2%	0.6%

A deviation from the SM prediction can point towards New Physics

Is the SM prediction under control at the O(0.5%) level in the TeV region of the $m_{\ell\ell}$ distribution ?

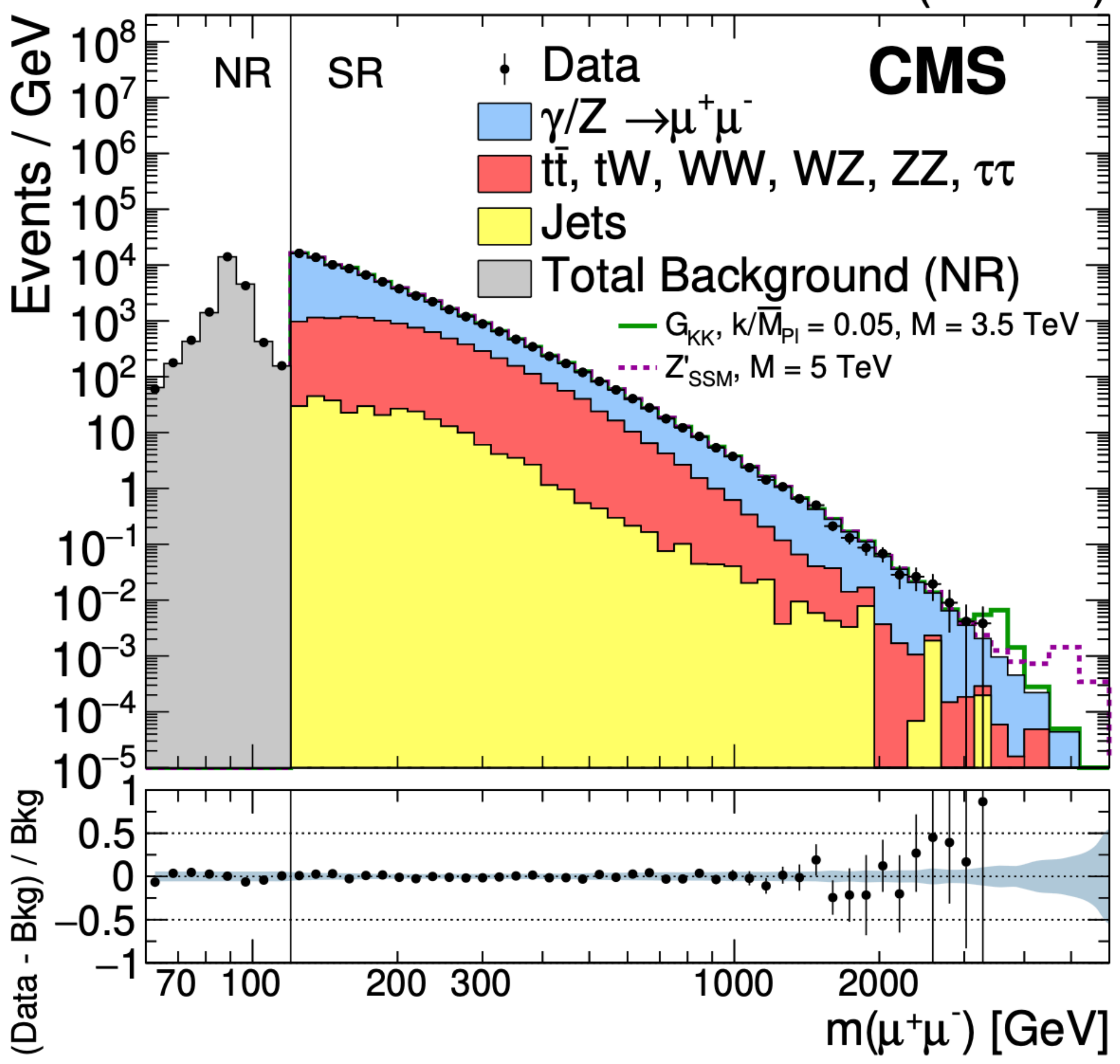
Motivations: towards per mille physics 😲 at the LHC ?



Channel	Not constraining PDFs	Constraining PDFs
Muons	0.23125 ± 0.00054	0.23125 ± 0.00032
Electrons	0.23054 ± 0.00064	0.23056 ± 0.00045
Combined	0.23102 ± 0.00057	0.23101 ± 0.00030

A determination of $\sin^2\theta^{lep}_{eff}$ competitive with the LEP results (0.23152(16)) is becoming possible

arXiv:2103.02708

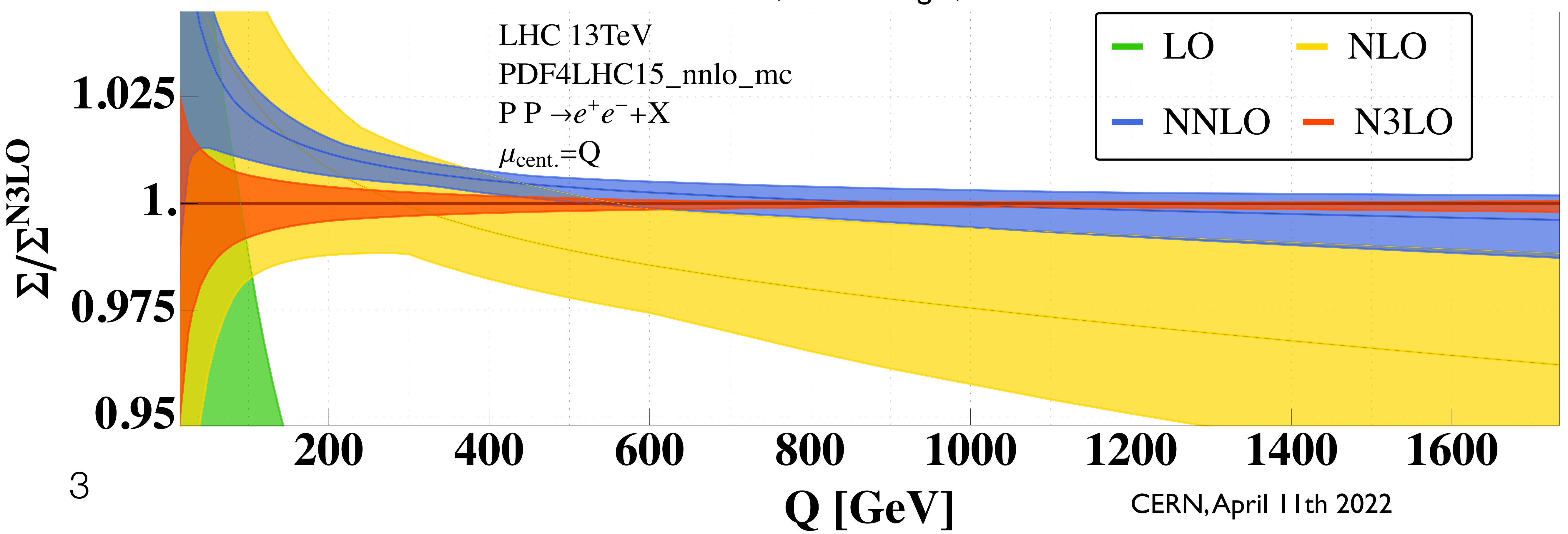


mass window [GeV]	stat. unc. $140 fb^{-1}$	stat. unc. $3 ab^{-1}$
$600<m_{\mu\mu}<900$	1.4%	0.2%
$900<m_{\mu\mu}<1300$	3.2%	0.6%

A deviation from the SM prediction can point towards New Physics

Is the SM prediction under control at the O(0.5%) level in the TeV region of the $m_{\ell\ell}$ distribution ?

C.Duhr, B.Mistlberger, arXiv:2111.10379



Neutral current Drell-Yan in fixed order

Drell-Yan (1970)

$$\sigma(h_1 h_2 \to \ell\bar{\ell} + X) = \sigma^{(0,0)} +$$

Baur, Brein, Hollik, Schappacher, Wackeroth (2001)

Altarelli, Ellis, Martinelli (1979)

$$\alpha_s\,\sigma^{(1,0)} + \alpha\,\sigma^{(0,1)} +$$

Hamberg, Matsuura, van Nerveen, (1991)
Anastasiou, Dixon, Melnikov, Petriello, (2003)
Catani, Cieri, Ferrera, de Florian, Grazzini (2009)

$$\alpha_s^2\,\sigma^{(2,0)} + \alpha\,\alpha_s\,\sigma^{(1,1)} + \alpha^2\,\sigma^{(0,2)} +$$

still missing
Sudakov high-energy approximations

$$\alpha_s^3\,\sigma^{(3,0)} + \ldots$$

C.Duhr, B.Mistlberger, arXiv:2111.10379

R.Bonciani, L.Buonocore, M.Grazzini, S.Kallweit, N.Rana, F.Tramontano, AV, arXiv:2106.11953
T.Armadillo, R.Bonciani, S.Devoto, N.Rana, AV, arXiv:2201.01754
F.Buccioni, F.Caola, H.Chawdhry, F.Devoto, M.Heller, A.von Manteuffel, K.Melnikov, R.Röntsch, C.Signorile-Signorile, arXiv:2203.11237

Because of the lack of time, the presentation will focus mostly on arXiv:2106.11953, arXiv:2201.01754
pointing out the main methodological differences adopted in arXiv:2203.11237

The availability of two completely independent calculations of such a complex set of corrections
will be crucial for technical validation (forthcoming)

Progress towards Drell-Yan simulations at NNLO QCD-EW

Strong boost of the activities in the theory community in the last 2 years!

→ mathematical and theoretical developments and computation of universal building blocks

- 2-loop virtual Master Integrals with internal masses

U. Aglietti, R. Bonciani, arXiv:0304028, arXiv:0401193, R. Bonciani, S. Di Vita, P. Mastrolia, U. Schubert, arXiv:1604.08581, M.Heller, A.von Manteuffel, R.Schabinger arXiv:1907.00491, S.Hasan, U.Schubert, arXiv:2004.14908, M.Long,R,Zhang,W.Ma,Y,Jiang,L.Han,,Z.Li,S.Wang, arXiv:2111.14130, X.Liu, Y.Ma, arXiv:2201.11669

- Altarelli-Parisi splitting functions including QCD-QED effects

D. de Florian, G. Sborlini, G. Rodrigo, arXiv:1512.00612

- renormalization

G.Degrassi, AV, hep-ph/0307122, S.Dittmaier,T.Schmidt,J.Schwarz, arXiv:2009.02229, S.Dittmaier, arXiv:2101.05154

→ on-shell Z and W production as a first step towards full Drell-Yan

- pole approximation of the NNLO QCD-EW corrections

S.Dittmaier, A.Huss, C.Schwinn, arXiv:1403.3216, 1511.08016

- analytical total cross section including NNLO QCD-QED and NNLO QED corrections

D. de Florian, M.Der, I.Fabre, arXiv:1805.12214

- ptZ distribution including QCD-QED analytical transverse momentum resummation

L. Cieri, G. Ferrera, G. Sborlini, arXiv:1805.11948

- fully differential on-shell Z production including exact NNLO QCD-QED corrections

M.Delto, M.Jaquier, K.Melnikov, R.Roentsch, arXiv:1909.08428

- total Z production cross section in fully analytical form including exact NNLO QCD-EW corrections

R. Bonciani, F. Buccioni, R.Mondini, AV, arXiv:1611.00645, R. Bonciani, F. Buccioni, N.Rana, I.Triscari, AV, arXiv:1911.06200, R. Bonciani, F. Buccioni, N.Rana, AV, arXiv:2007.06518, arXiv:2111.12694

- fully differential on-shell Z and W production including exact NNLO QCD-EW corrections

F. Buccioni, F. Caola, M.Delto, M.Jaquier, K.Melnikov, R.Roentsch, arXiv:2005.10221, A. Behring, F. Buccioni, F. Caola, M.Delto, M.Jaquier, K.Melnikov, R.Roentsch, arXiv:2009.10386, 2103.02671,

→ complete Drell-Yan

- neutrino-pair production including NNLO QCD-QED corrections

L. Cieri, D. de Florian, M.Der, J.Mazzitelli, arXiv:2005.01315

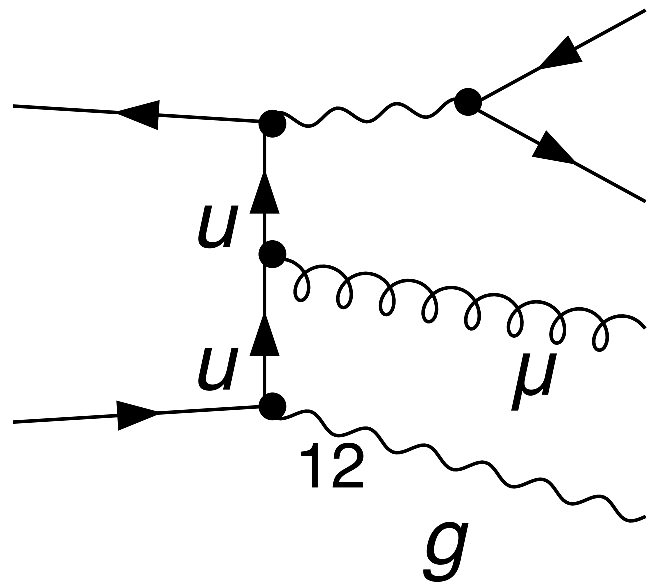
- 2-loop amplitudes

M.Heller, A.von Manteuffel, R.Schabinger, arXiv:2012.05918

- NNLO QCD-EW corrections to charged-current DY including leptonic decay (2-loop contributions in pole approximation).

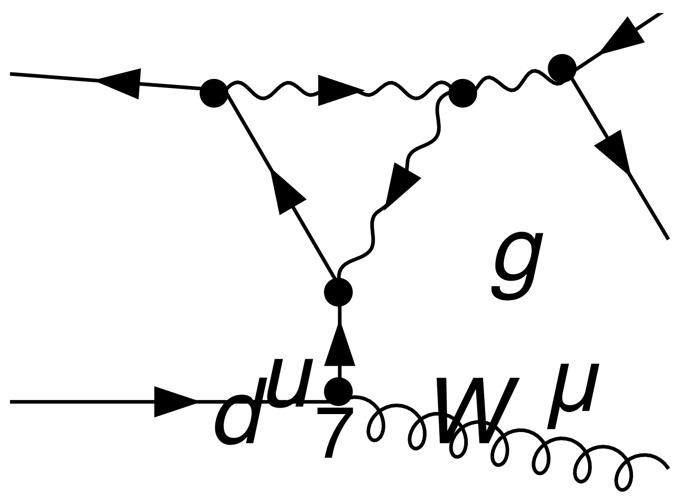
L.Buonocore, M.Grazzini, S.Kallweit, C.Savoini, F.Tramontano, arXiv:2102.12539

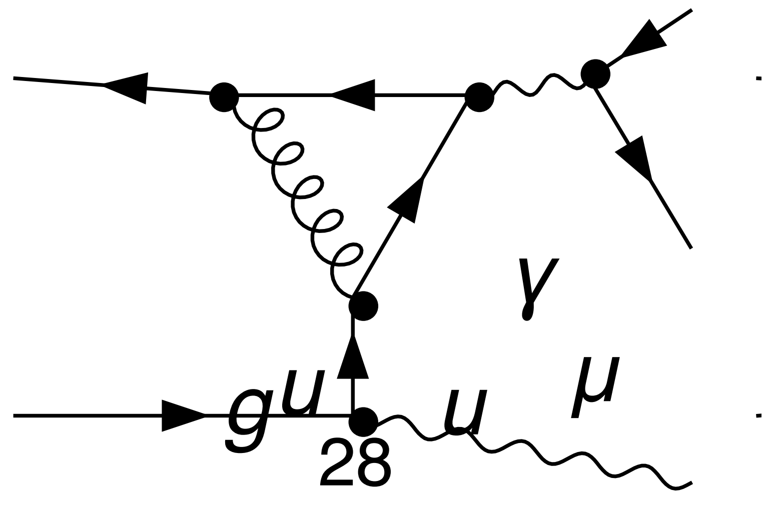
Different kinds of contributions at $\mathcal{O}(\alpha\alpha_s)$ and corresponding problems



double-real contributions

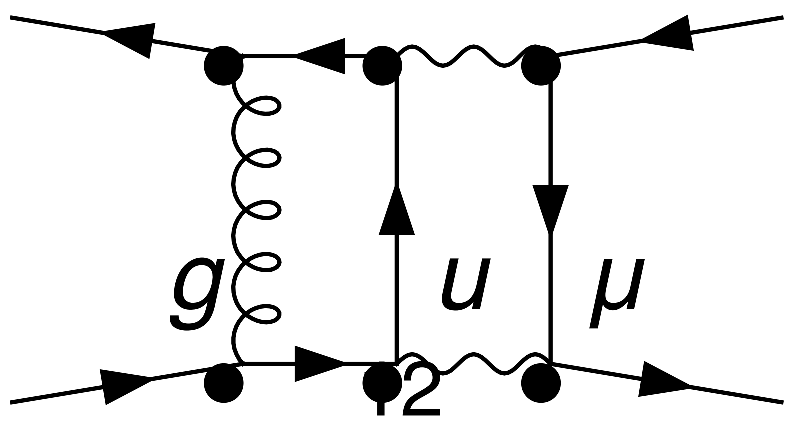
amplitudes are easily generated with OpenLoops
IR subtraction
care about the numerical convergence when aiming at 0.1% precision

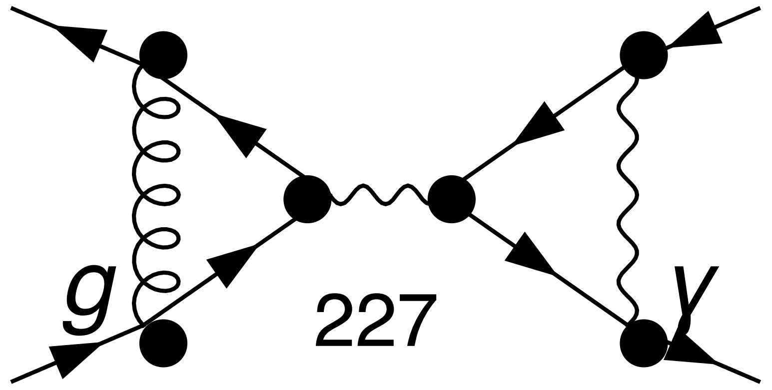




real-virtual contributions

amplitudes are easily generated with OpenLoops or Recola
1-loop UV renormalisation and IR subtraction
care about the numerical convergence when aiming at 0.1% precision





double-virtual contributions

generation of the amplitudes
γ_5 treatment
2-loop UV renormalization
subtraction of the IR divergences
solution and evaluation of the Master Integrals
numerical evaluation of the squared matrix element

The double-real and real-virtual corrections already known from studies of the large transverse momentum lepton pair final state
A.Denner, S.Dittmaier, T.Kasprzik, A.Muck, arXiv:1103.0914, A.Denner, S.Dittmaier, M.Hecht, C.Pasold, arXiv:1510.08742 J.Lindert et al., arXiv:1705.04664
Now we can consider the inclusive spectrum, also in the $q_T \to 0$ limit

The infrared structure of the calculation

Soft and/or collinear divergences
appear in the expression of all the different subprocesses contributing to the inclusive cross section.

Their cancellation in IR-safe sufficiently inclusive observables, after the combination of real and virtual corrections, guarantees the predictivity of the calculation.

At NNLO different methodologies have been devised to handle this problem,
making the individual contributions separately finite and ready for numerical evaluation.
The final IR-finite prediction must be independent of the subtraction technique

arXiv:2106.11953 qt-subtraction method, implemented in the Matrix framework
massive final-state leptons (muons) → bare (i.e. not recombined) definition
explicit logarithms of the muon mass in the final result, regularising the FSR collinear divergences
KLN cancellation of muon logs in inclusive observables

arXiv:2203.11237 nested soft-collinear subtraction
massless final-state leptons → lepton-photon recombination procedure
KLN cancellation of collinear divergences inside the cone of the e.m. jet

General structure of the inclusive cross section and the q_T-subtraction formalism in $\mathtt{Matrix}$

$$d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathscr{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right]_{q_T/Q > r_{cut}}$$

IR structure associated to the QCD-QED part derived from NNLO-QCD results via abelianisation
(de Florian, Rodrigo, Sborlini, 2016, de Florian, Der , Fabre, 2018)

the q_T-subtraction formalism has been extended to the case of massive final-state emitters (heavy quarks in QCD, leptons in EW)
(Catani, Torre, Grazzini, 2014, Buonocore,Grazzini, Tramontano 2019.)

General structure of the inclusive cross section and the q_T-subtraction formalism

$$d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathcal{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right]_{q_T/Q > r_{cut}}$$

IR structure associated to the QCD-QED part derived from NNLO-QCD results via abelianisation
(de Florian, Rodrigo, Sborlini, 2016, de Florian, Der , Fabre, 2018)

the q_T-subtraction formalism has been extended to the case of massive final-state emitters (heavy quarks in QCD, leptons in EW)
(Catani, Torre, Grazzini, 2014, Buonocore, Grazzini, Tramontano 2019.)

$$\int d\sigma_R^{(1,1)} \sim \sum_{i=1}^{4} c_i \ln^i r_{cut} + c_0 + \mathcal{O}(r_{cut}^m) \qquad \rightarrow \qquad \int \left(d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right) \sim c_0 + \mathcal{O}(r_{cut}^m)$$

The counterterm removes the IR sensitivity to the cutoff variable
→ we need small values of the cutoff and explicit numerical tests to quantify the bias induced by the cutoff choice
we can fit the r_{cut} dependence and extrapolate in the $r_{cut} \to 0$ limit
(cfr. Buonocore, Kallweit, Rottoli, Wiesemann, arXiv:2111.13661, Camarda, Cieri, Ferrera, arXiv:2111.14509)

General structure of the inclusive cross section and the q_T-subtraction formalism

$$d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathscr{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)}\right]_{q_T/Q>r_{cut}}$$

IR structure associated to the QCD-QED part derived from NNLO-QCD results via abelianisation
(de Florian, Rodrigo, Sborlini, 2016, de Florian, Der , Fabre, 2018)

the q_T-subtraction formalism has been extended to the case of massive final-state emitters (heavy quarks in QCD, leptons in EW)
(Catani, Torre, Grazzini, 2014, Buonocore,Grazzini,Tramontano 2019.)

$$\int d\sigma_R^{(1,1)} \sim \sum_{i=1}^{4} c_i \ln^i r_{cut} + c_0 + \mathcal{O}(r_{cut}^m) \qquad \rightarrow \qquad \int \left(d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)}\right) \sim c_0 + \mathcal{O}(r_{cut}^m)$$

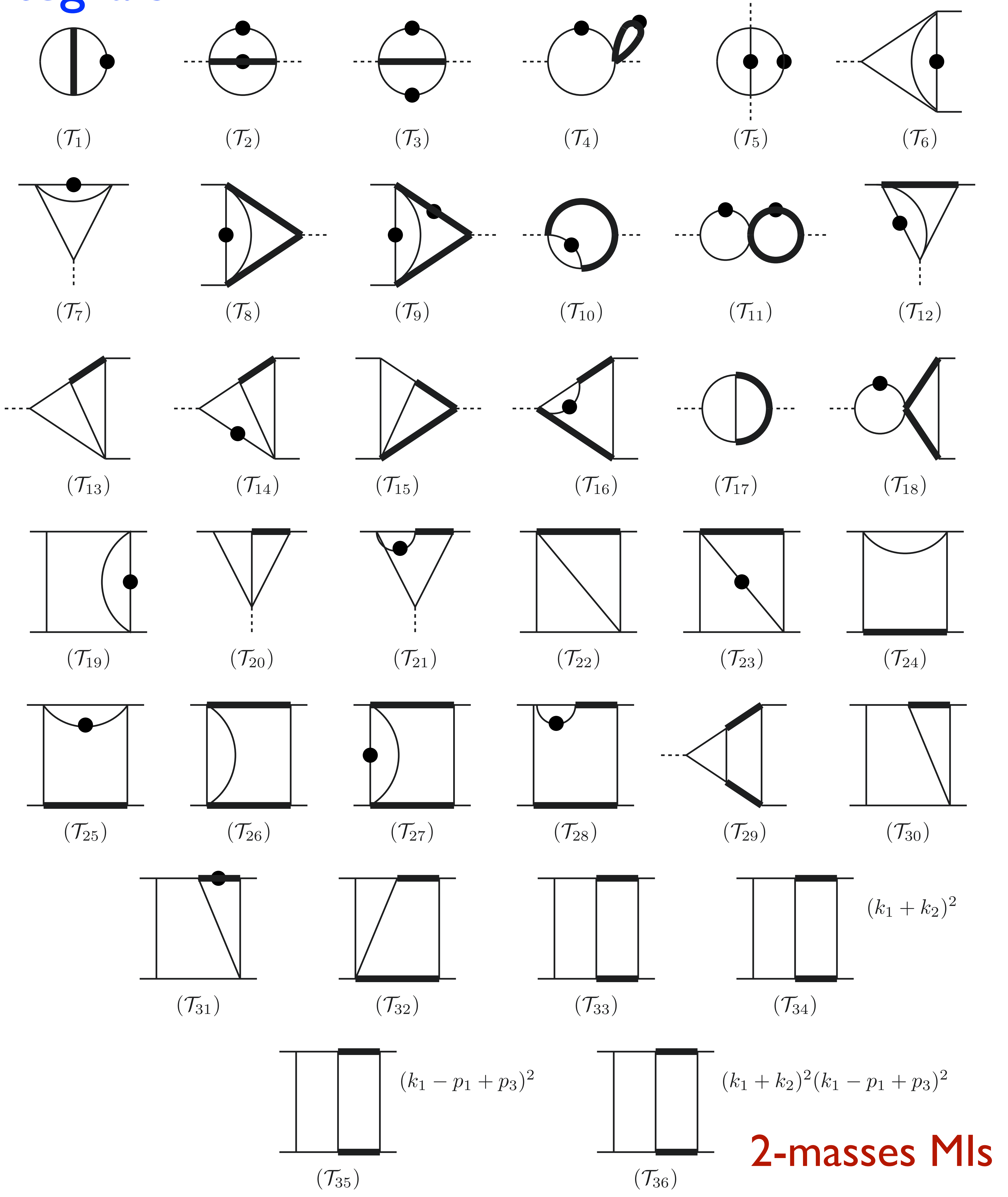
The counterterm removes the IR sensitivity to the cutoff variable
→ we need small values of the cutoff and explicit numerical tests to quantify the bias induced by the cutoff choice
we can fit the r_{cut} dependence and extrapolate in the $r_{cut} \to 0$ limit
(cfr. Buonocore, Kallweit, Rottoli, Wiesemann, arXiv:2111.13661, Camarda, Cieri, Ferrera, arXiv:2111.14509)

$$\mathscr{H}^{(1,1)} = H^{(1,1)}\, C_1\, C_2 \qquad 2\mathrm{Re}\langle \mathscr{M}^{(0,0)} | \mathscr{M}^{(1,1)} \rangle = \sum_{k=-4}^{0} \varepsilon^k f_i(s,t,m) \qquad |\mathscr{M}_{fin}\rangle \equiv (1-I)|\mathscr{M}\rangle \qquad H \propto \langle \mathscr{M}_0 | \mathscr{M}_{fin} \rangle$$

The IR poles are removed from the full 2-loop amplitude by means of a subtraction procedure (matching the real radiation one)

The double virtual amplitude: reduction to Master Integrals

$$2\mathrm{Re}\left(\mathscr{M}^{(1,1)}(\mathscr{M}^{(0,0)})^{\dagger}\right) = \sum_{i=1}^{N_{MI}} c_i(s,t,m;\varepsilon)\ \mathscr{T}_i(s,t,m;\varepsilon)$$



2-masses MIs

The double virtual amplitude: reduction to Master Integrals

$$2\mathrm{Re}\left(\mathcal{M}^{(1,1)}(\mathcal{M}^{(0,0)})^\dagger\right) = \sum_{i=1}^{N_{MI}} c_i(s,t,m;\varepsilon)\ \mathcal{T}_i(s,t,m;\varepsilon)$$

The coefficients c_i are rational functions of the invariants, masses and of ε
The size of the total expression can rapidly "explode"
→ careful work to identify the patterns of recurring subexpressions keeping the total size in the O(1-10 MB) range

The complexity of the MIs depends on the number of energy scales
MIs relevant for the QCD-QED corrections, with massive final state

Bonciani, Ferroglia, Gehrmann, Maitre, Studerus., arXiv:0806.2301, 0906.3671

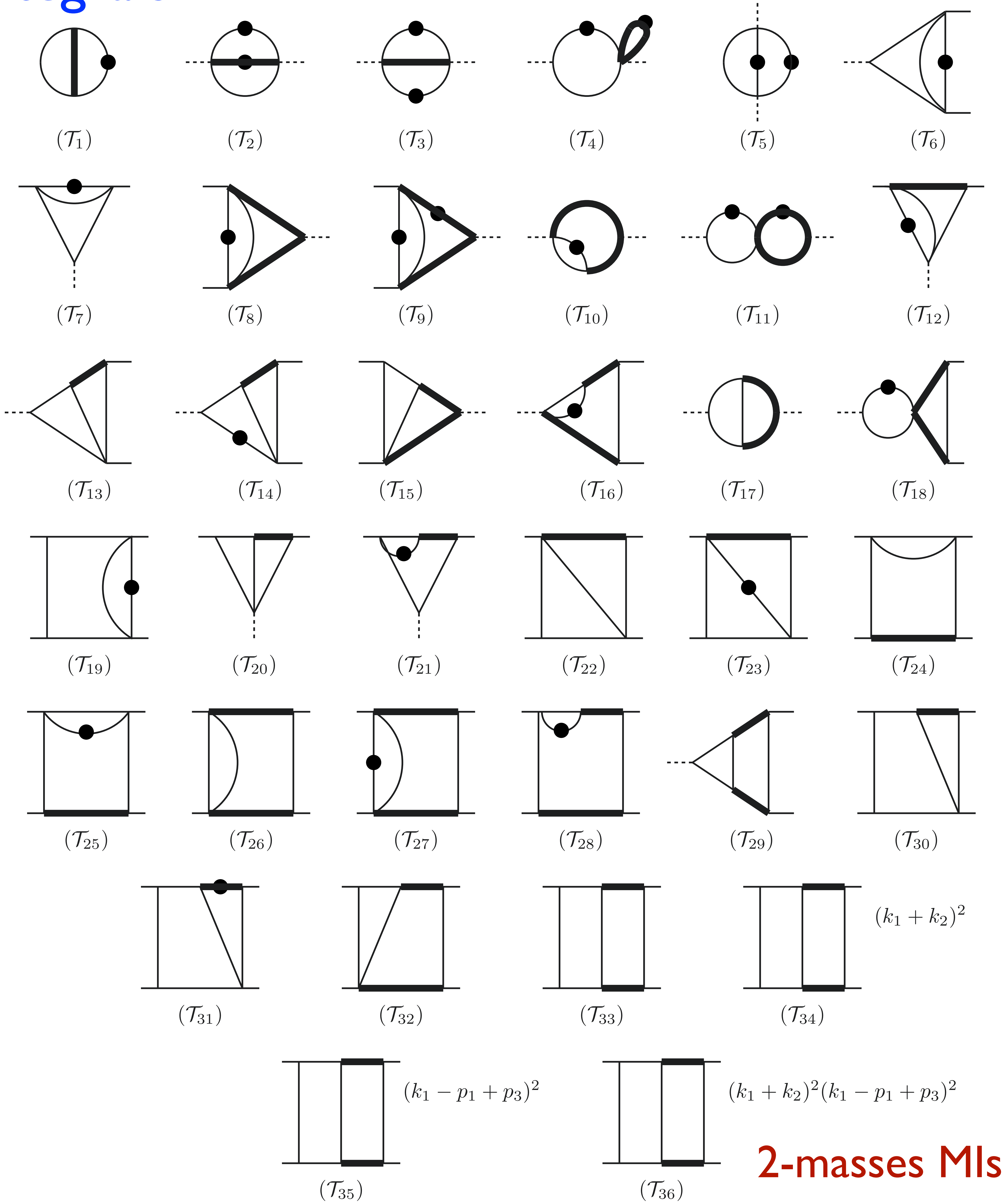
MIs with 1or 2 internal mass relevant for the EW form factor

Aglietti, Bonciani, hep-ph/0304028, hep-ph/0401193

31 MIs with 1 mass and 36 MIs with 2 masses including boxes, relevant for the QCD-weak corrections to the full Drell-Yan

Bonciani, Di Vita, Mastrolia, Schubert., arXiv:1604.08581

In the 2-mass case, 5 box integrals in Chen-Goncharov representation
→ problematic numerical evaluation → need an alternative strategy



2-masses MIs

cfr. also Heller, von Manteuffel, Schabinger, arXiv:1907.00491 for a representation of the MIs in terms of GPLs
arXiv:2012.05918 for a description of the 2-loop virtual amplitude

Evaluation of the Master Integrals by series expansions

The Master Integrals satisfy a system of differential equations.

The MIs are replaced by formal series with unknown coefficients → eqs for the unknown coefficients of the series.

The package `DiffExp` by M.Hidding, arXiv:2006.05510 implements this idea, for real valued masses, with real kinematical vars.

But we need complex-valued masses of W and Z bosons (unstable particles) → we wrote a new package (`SeaFire`)

Evaluation of the Master Integrals by series expansions

The Master Integrals satisfy a system of differential equations.

The MIs are replaced by formal series with unknown coefficients → eqs for the unknown coefficients of the series.

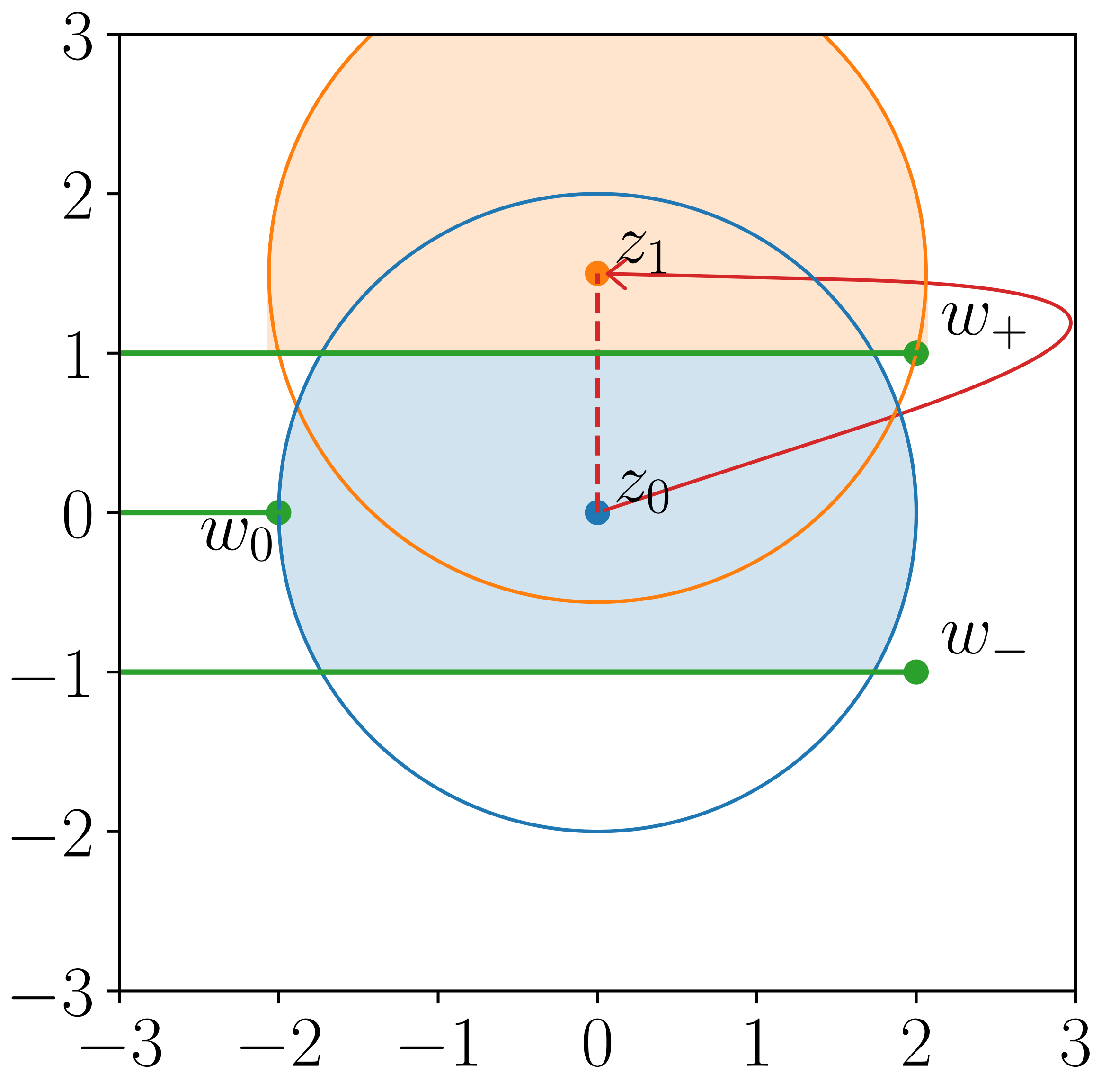
The package DiffExp by M.Hidding, arXiv:2006.05510 implements this idea, for real valued masses, with real kinematical vars.

But we need complex-valued masses of W and Z bosons (unstable particles) → we wrote a new package (SeaFire)

We implemented the same approach, for arbitrary complex-valued masses, working in the complex plane of each kinematical variable

Complete knowledge about the singular structure of the MI can be read directly from the differential equation matrix

The solution can be computed with an arbitrary number of significant digits, but not in closed form → semi-analytical



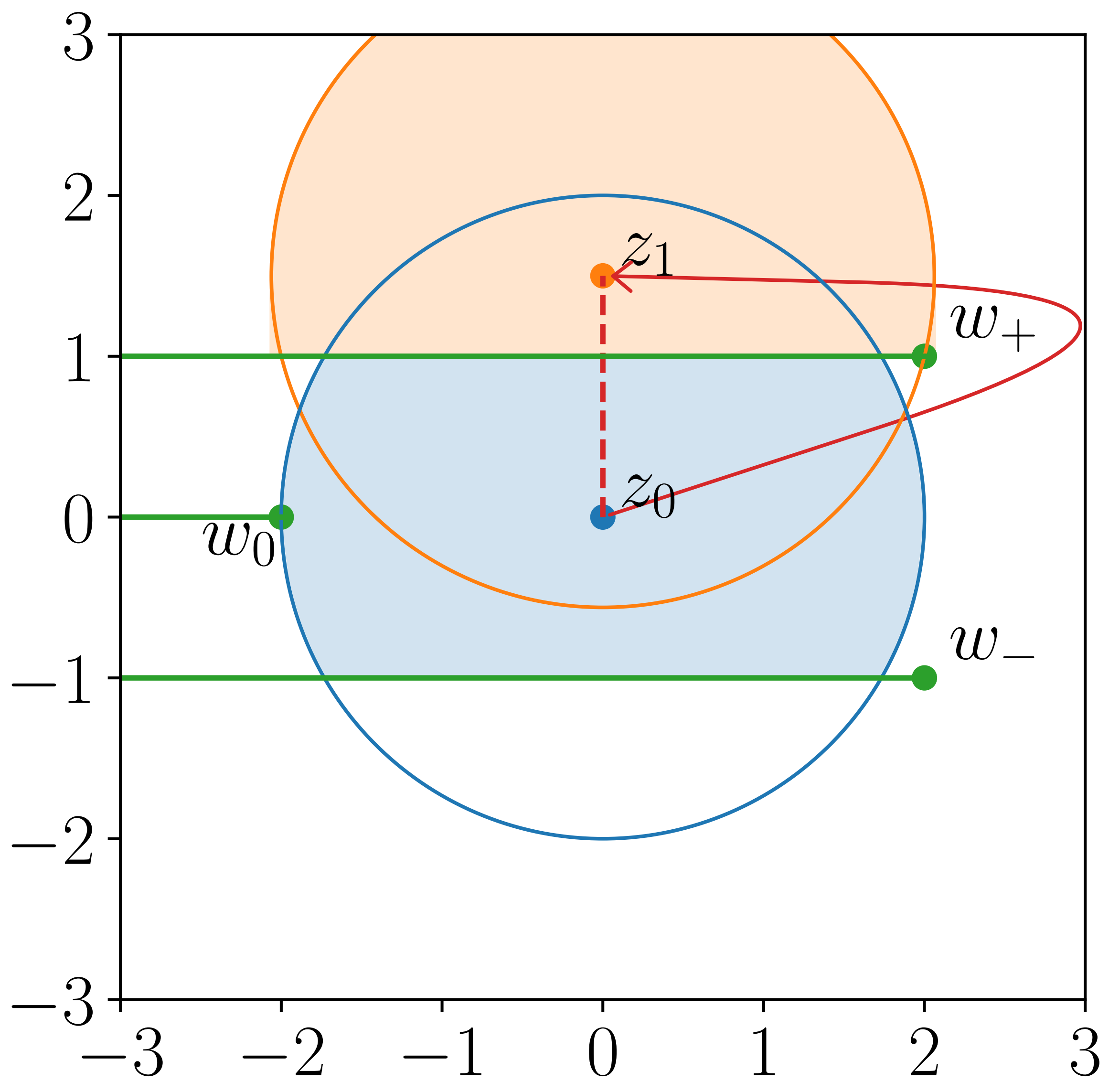
Evaluation of the Master Integrals by series expansions

The Master Integrals satisfy a system of differential equations.
The MIs are replaced by formal series with unknown coefficients → eqs for the unknown coefficients of the series.
The package DiffExp by M.Hidding, arXiv:2006.05510 implements this idea, for real valued masses, with real kinematical vars.
But we need complex-valued masses of W and Z bosons (unstable particles) → we wrote a new package (SeaFire)

We implemented the same approach, for arbitrary complex-valued masses,
working in the complex plane of each kinematical variable

Complete knowledge about the singular structure of the MI
can be read directly from the differential equation matrix

The solution can be computed with an arbitrary number of significant digits,
but not in closed form → semi-analytical



Numerical evaluation of the hard coefficient function

The interference term $2\mathrm{Re}\langle \mathcal{M}^{(1,1),fin} | \mathcal{M}^{(0,0)} \rangle$ contributes to the hard function $H^{(1,1)}$

After the subtraction of all the universal IR divergences, it is a finite correction

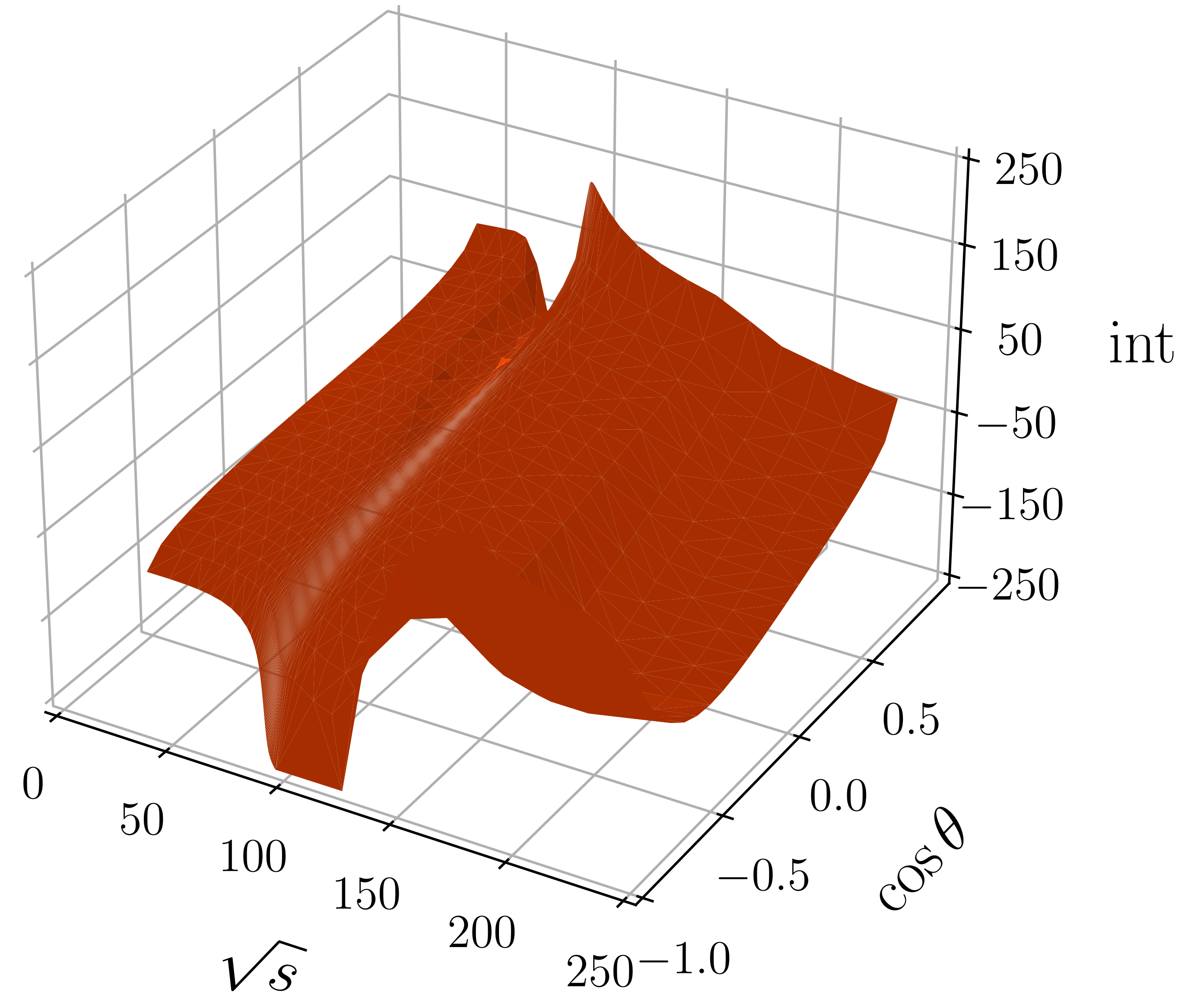
It has been published in arXiv:2201.01754 and is available as a Mathematica notebook

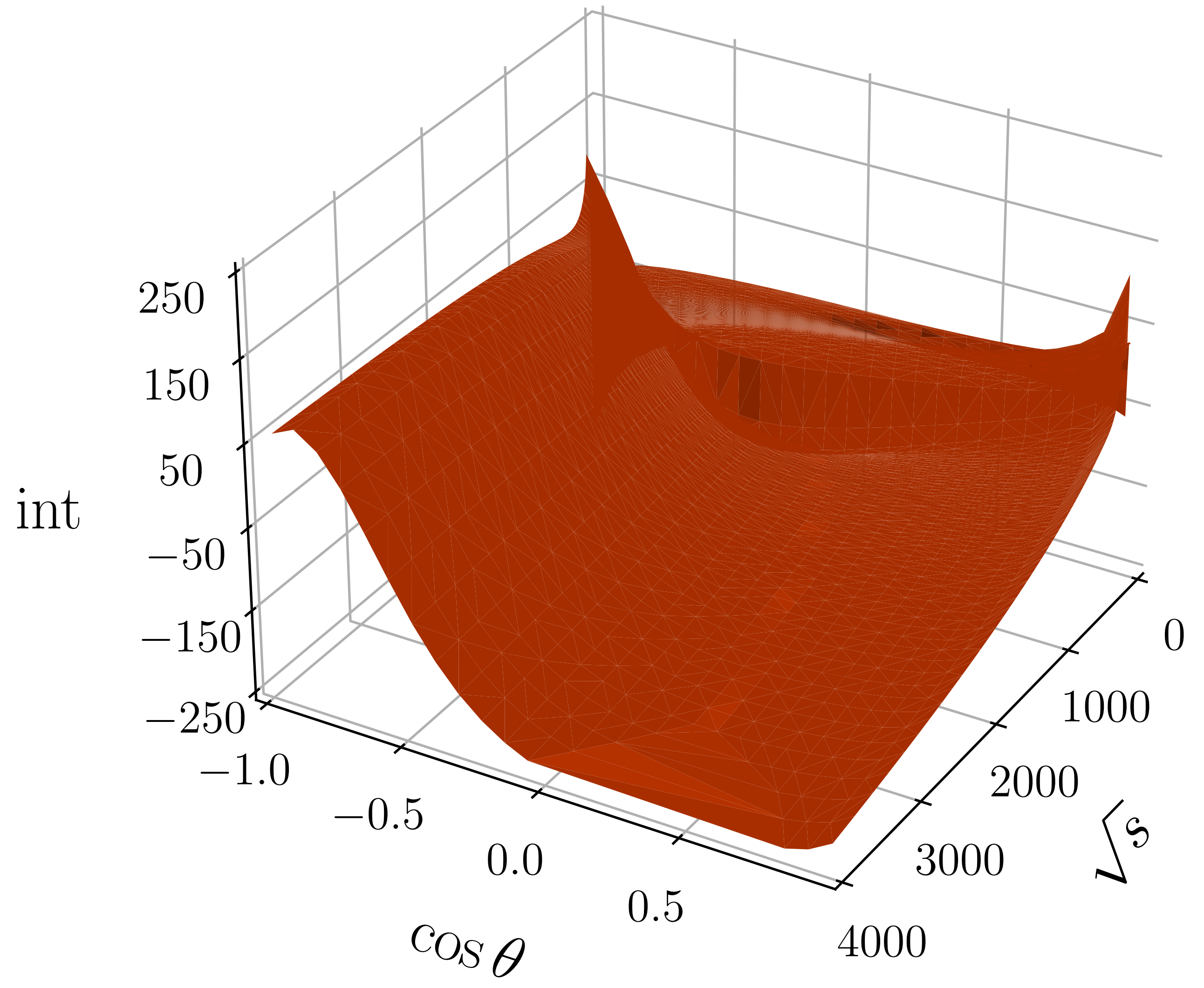
Several checks of the MIs performed with Fiesta and PySecDec

A numerical grid has been prepared for all the 36 MIs, with GiNaC and SeaFire (T.Armadillo et al, in preparation), covering the whole $2 \to 2$ phase space in (s,t), in O(12 h) on one 32-cores machine

→ a numerical grid for $2\mathrm{Re}\langle \mathcal{M}^{(1,1),fin} | \mathcal{M}^{(0,0)} \rangle$ has been prepared
values at arbitrary phase space points with excellent accuracy via interpolation, with negligible evaluation time

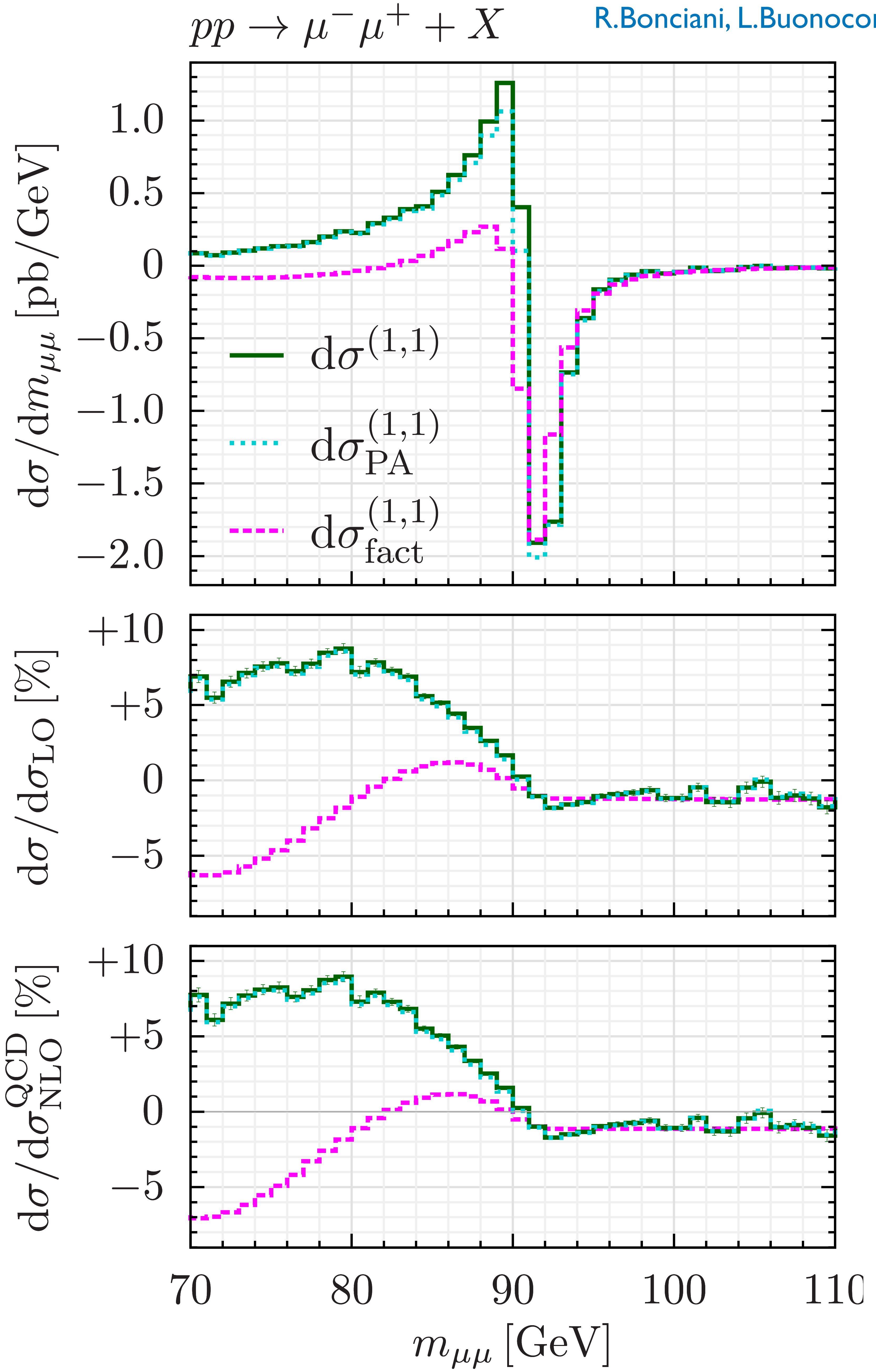
in units $\dfrac{\alpha}{2\pi}\dfrac{\alpha_s}{2\pi}\sigma_0$





The lepton-pair invariant mass distribution: QCD-EW corrections

R.Bonciani, L.Buonocore, M.Grazzini, S.Kallweit, N.Rana, F.Tramontano, AV, arXiv:2106.11953



below the Z resonance,
QCD and QED effects do not factorise

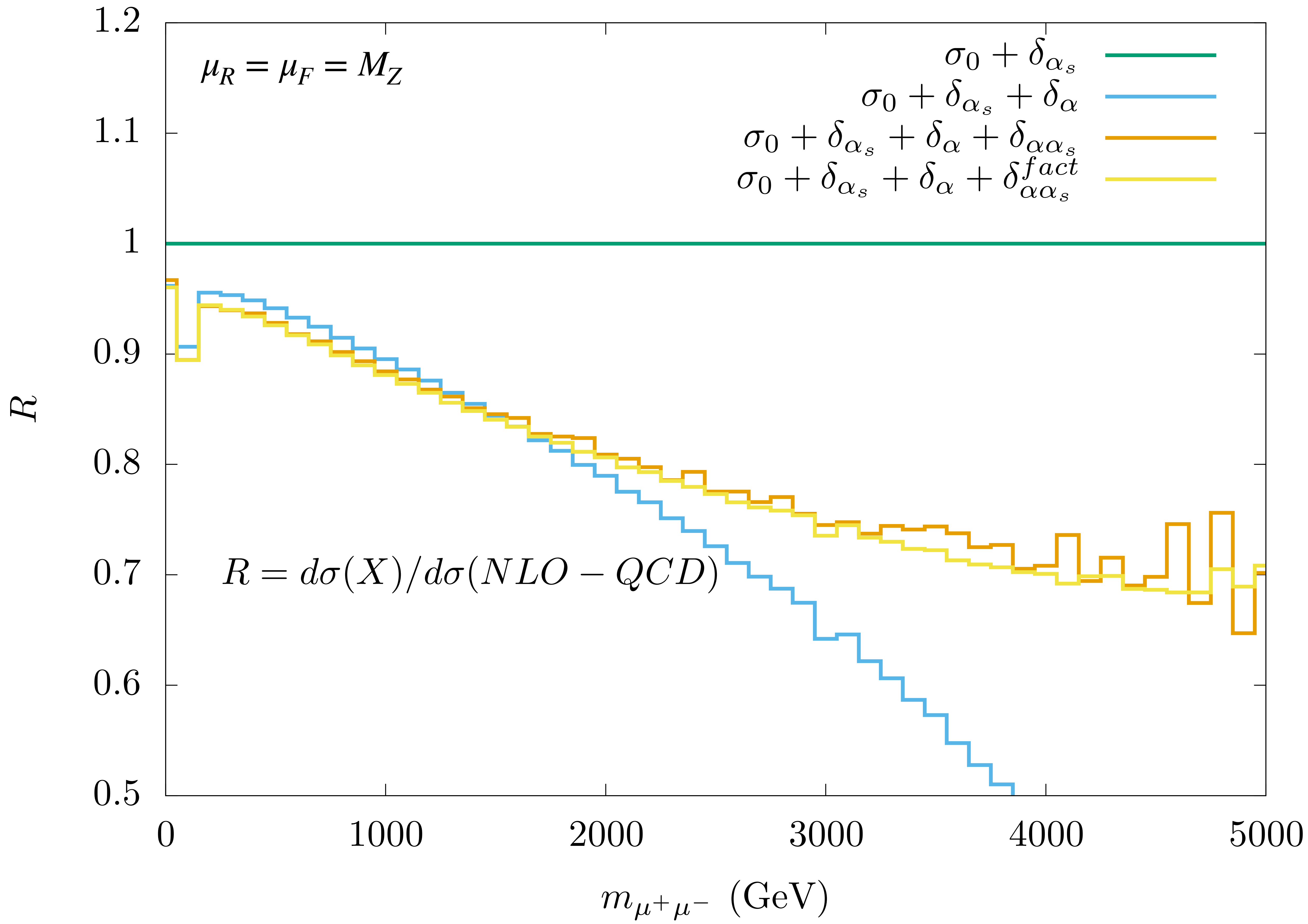
- a proper treatment of kinematics is needed (cfr. e.g. POWHEG QCD+EW)
- the full calculation has a non-trivial correction

the pole approximation provides an excellent description of the full calculation confirming S.Dittmaier, A.Huss, C.Schwinn, arXiv:1403.3216, 1511.08016

the distortion of the Z line shape has an impact on the determination of $\sin^2\theta^{lep}_{eff}$

The lepton-pair invariant mass distribution: QCD-EW corrections

R.Bonciani, L.Buonocore, M.Grazzini, S.Kallweit, N.Rana, F.Tramontano, AV, arXiv:2106.11953



below 1 TeV
the corrections are negative and of O(-1.5%)

the factorised approximation
catches the bulk of the QCD-EW correction
but
a residual O(1%) non-factorisable effect emerges

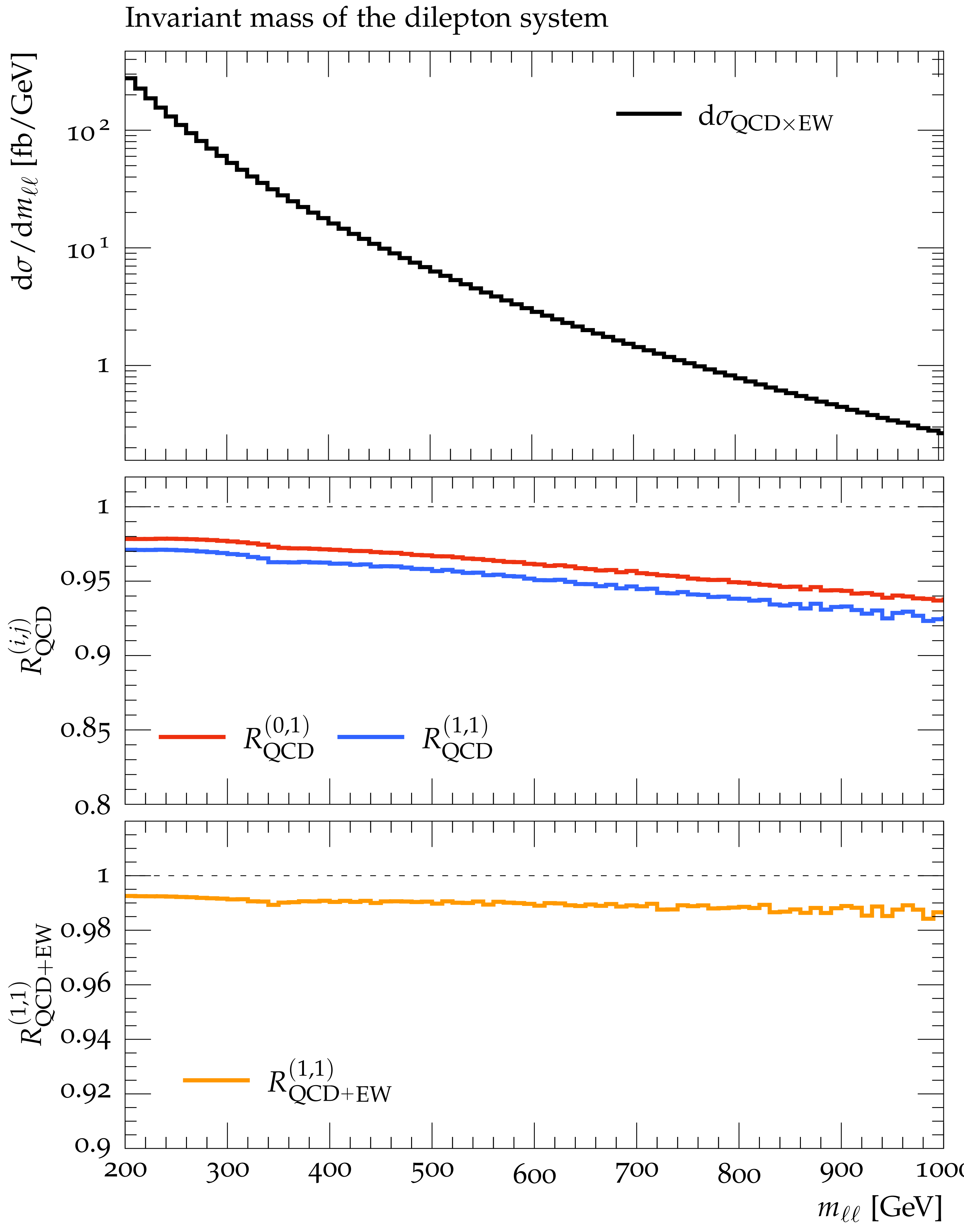
with fixed $\mu_R = \mu_F = M_Z$
in the very high invariant mass range,
QCD-EW effects are large and positive

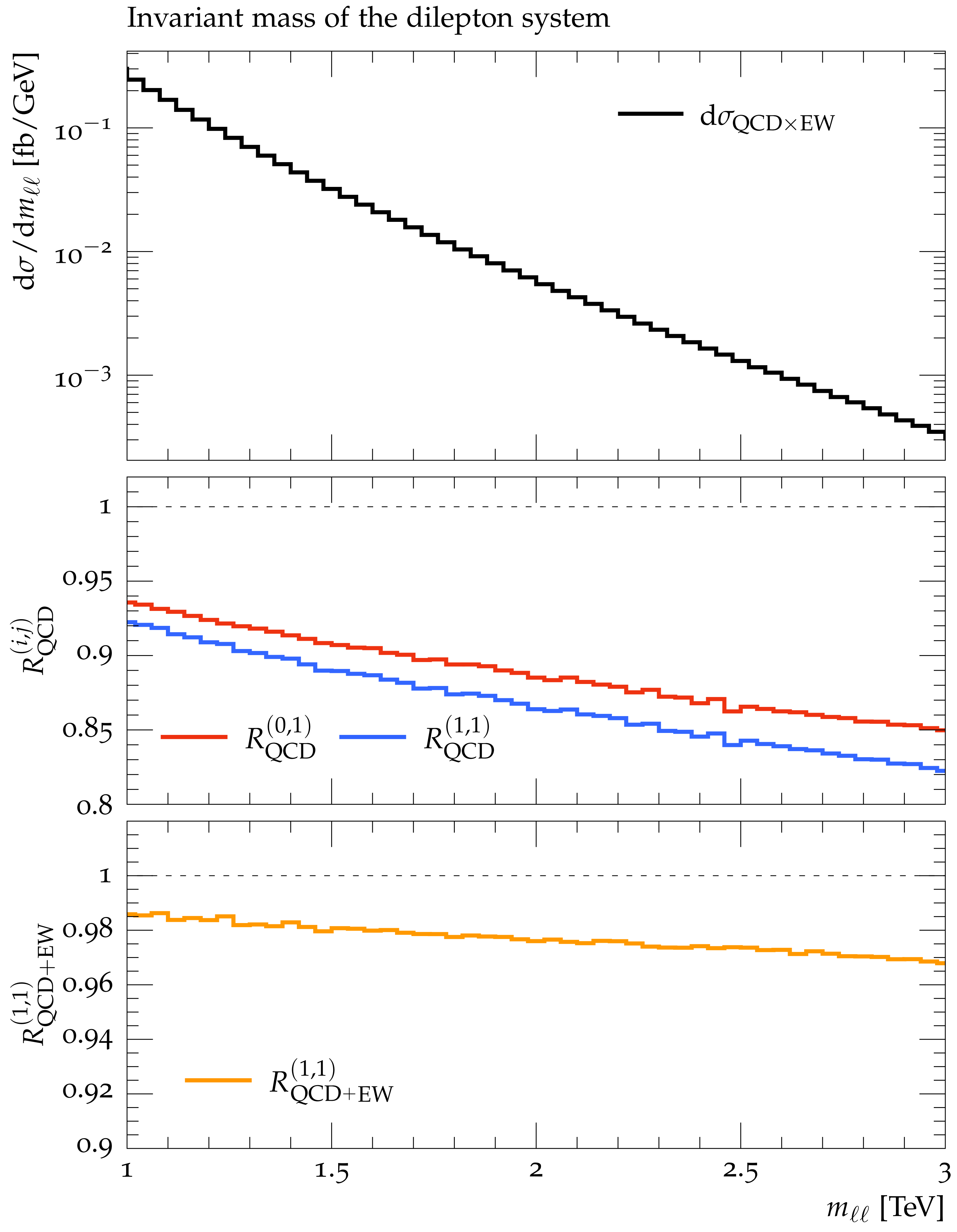
at 3 TeV they are O(+10%), still comparable
with the statistical uncertainty at the end of HL-LHC

using a dynamical choice for μ_R and μ_F
causes a non-trivial change in the size and shape
of the corrections

The lepton-pair invariant mass distribution: QCD-EW corrections

F.Buccioni, F.Caola, H.Chawdhry, F.Devoto, M.Heller, A.von Manteuffel, K.Melnikov, R.Röntsch, C.Signorile-Signorile, arXiv:2203.11237





large impact of the running of α_s on the total QCDxEW correction

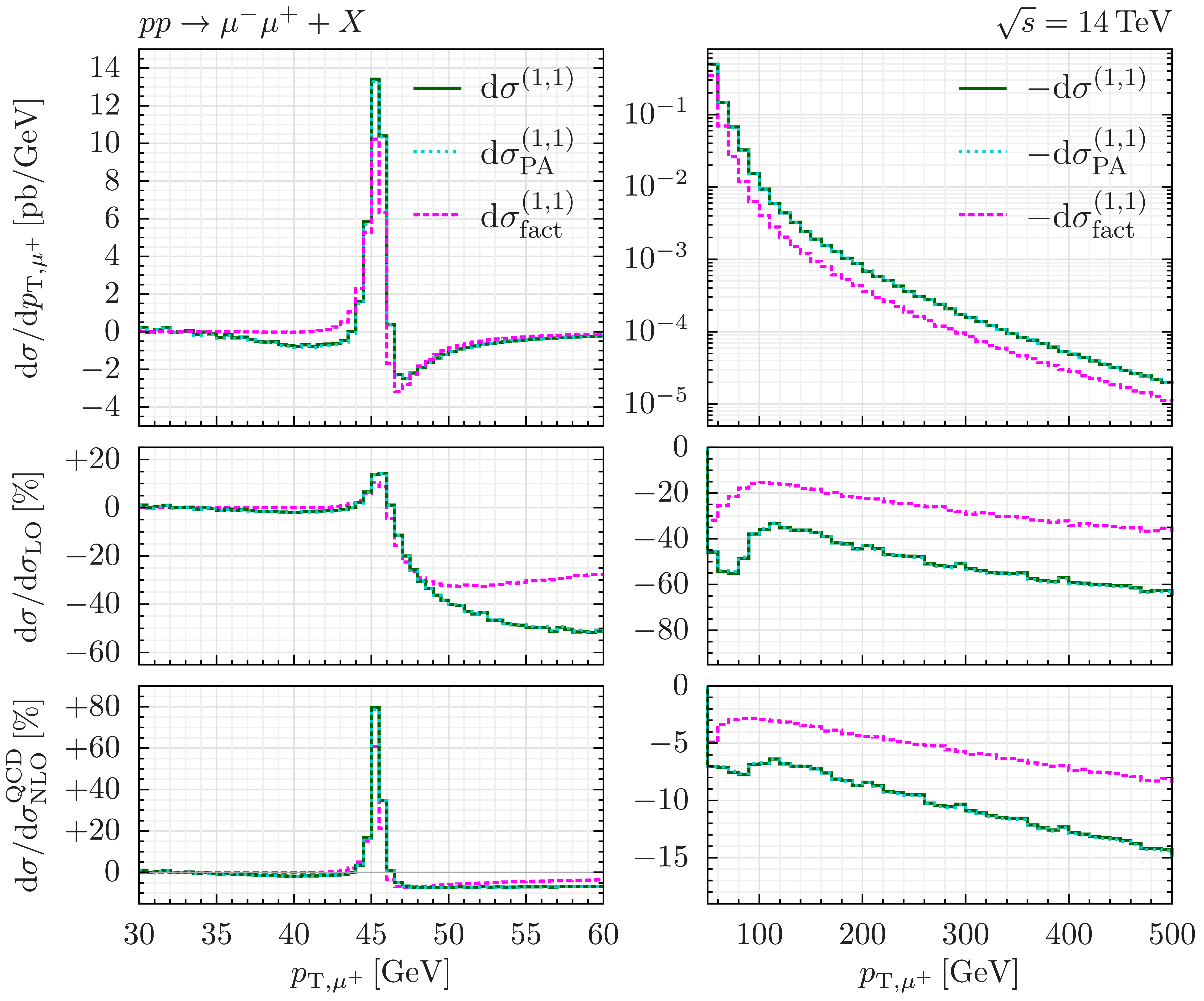
O(-1%) QCDxEW effects confirmed

$$R^{(0,1)}_{\mathrm{QCD}} = \frac{d\sigma^{(0,0)} + d\sigma^{(1,0)} + d\sigma^{(0,1)}}{d\sigma^{(0,0)} + d\sigma^{(1,0)}}, \quad R^{(1,1)}_{\mathrm{QCD}} = \frac{d\sigma^{(0,0)} + d\sigma^{(1,0)} + d\sigma^{(0,1)} + d\sigma^{(1,1)}}{d\sigma^{(0,0)} + d\sigma^{(1,0)}},$$

$$R^{(1,1)}_{\mathrm{QCD+EW}} = R^{(1,1)}_{\mathrm{QCD}}/R^{(0,1)}_{\mathrm{QCD}} = \frac{d\sigma^{(0,0)} + d\sigma^{(1,0)} + d\sigma^{(0,1)} + d\sigma^{(1,1)}}{d\sigma^{(0,0)} + d\sigma^{(1,0)} + d\sigma^{(0,1)}}$$

The lepton transverse momentum distribution: QCD-EW corrections

R.Bonciani, L.Buonocore, M.Grazzini, S.Kallweit, N.Rana, F.Tramontano, AV, arXiv:2106.11953



In the jacobian peak region the xsec is very much affected by the fixed-order nature of the results

In the large transverse momentum tail the dominant topology is Z+jet (qg-induced) which receives negative and large EW corrections

The exact calculation deviates from the factorised approximation because
the EW correction $d\sigma^{(0,1)}$ applies correctly to the $p_T^Z = 0$ bin
but
misses the large Sudakov logs which develop at $p_T^Z \gg 0$

The pole approximation on the 2-loop virtual corrections
affects only the $q\bar{q}$ process with $p_T^Z = 0$
which is negligible at large p_T^μ

Important impact on the Z+jet and W+jet generators

Estimate of the residual uncertainties: total cross section

The impact of the NNLO QCD-EW corrections is twofold: more accurate predictions (additional higher orders)
reduced uncertainties (scale, inputs, matching)

Ongoing phenomenological studies for full NC DY

A representative example from the results for the on-shell Z production total cross section
R.Bonciani, F.Buccioni, N.Rana, AV, arXiv:2007.06518, arXiv:2111.12694

→ dependence on the EW input-scheme choice

comparison of (G_μ, M_W, M_Z) and $(\alpha(0), M_W, M_Z)$ (very conservative choice that maximises the spread of the results)

order	G_μ	$\alpha(0)$	$\delta(G_\mu-\alpha(0))$ (%)
NNLO-QCD	55787	53884	3.53
NNLO-QCD+NLO-EW	55501	55015	0.88
NNLO-QCD+NLO-EW+ NNLO QCD-EW	55469	55340	0.23

the LO + NLO-EW result would suffer of only 0.55% spread;
the NLO-QCD and NNLO-QCD corrections are only LO-EW and reintroduce a dependence (→0.88%)
which is reduced by the NNLO QCD-EW (→0.23%)

Conclusions

The evaluation of the NNLO QCD-EW corrections is not yet a "pressing-just-one-button" game
but
the main obstacles to compute the 2-loop virtual corrections have been understood and solved for NC DY

- amplitude manipulation
- IR structure
- evaluation of Master Integrals with 2 internal masses

Two independent calculations are now available

The systematic automation of the progresses in the 2-loop virtual section is ongoing and will
allow the study of NNLO QCD-EW corrections to other scattering processes
be the starting point for the evaluation of NNLO-EW corrections

The phenomenological impact of mixed NNLO QCD-EW corrections is not negligible
in the precision physics program at the LHC

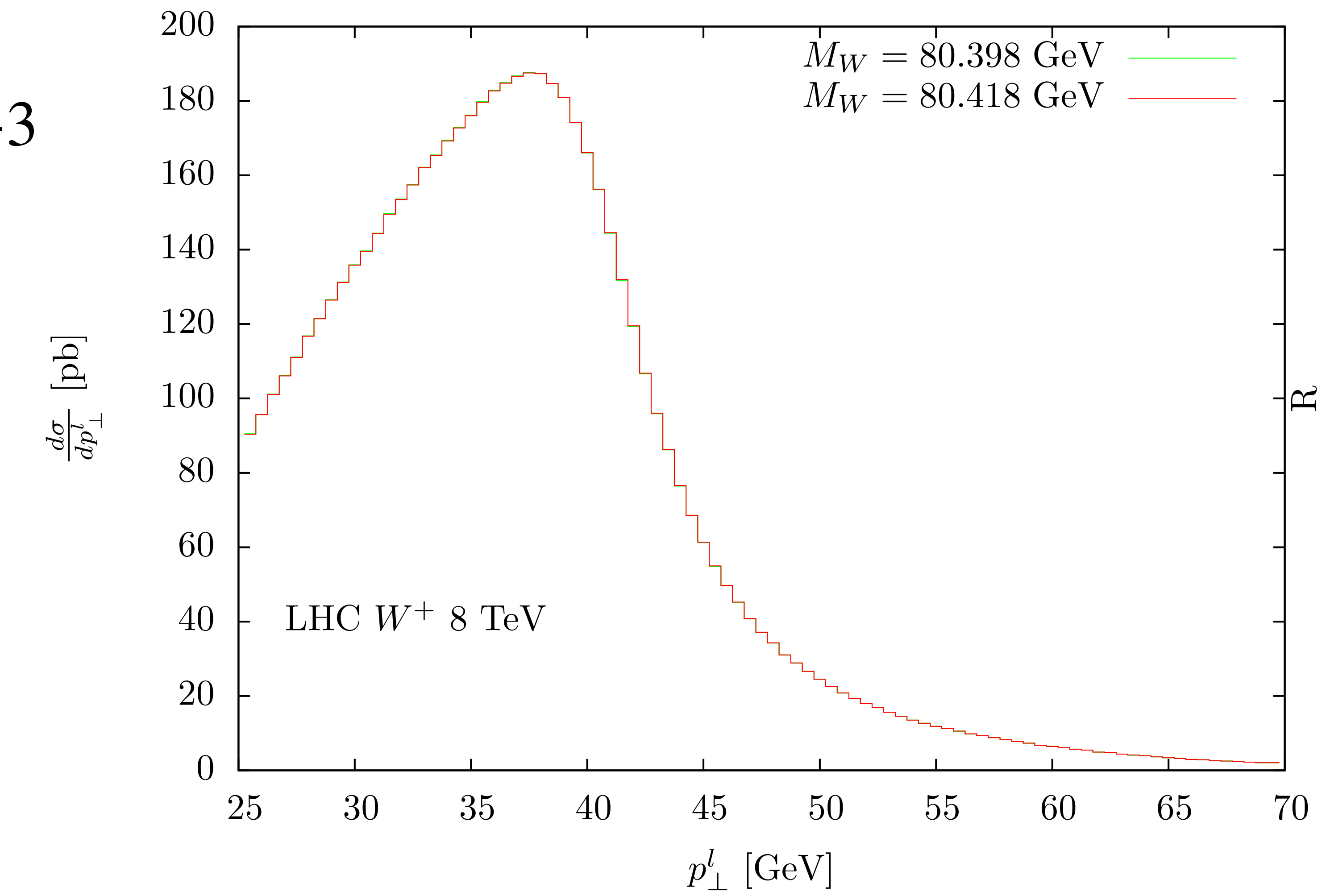
A precise SM prediction is the mandatory starting point for any SMEFT study or search in a UV-complete model

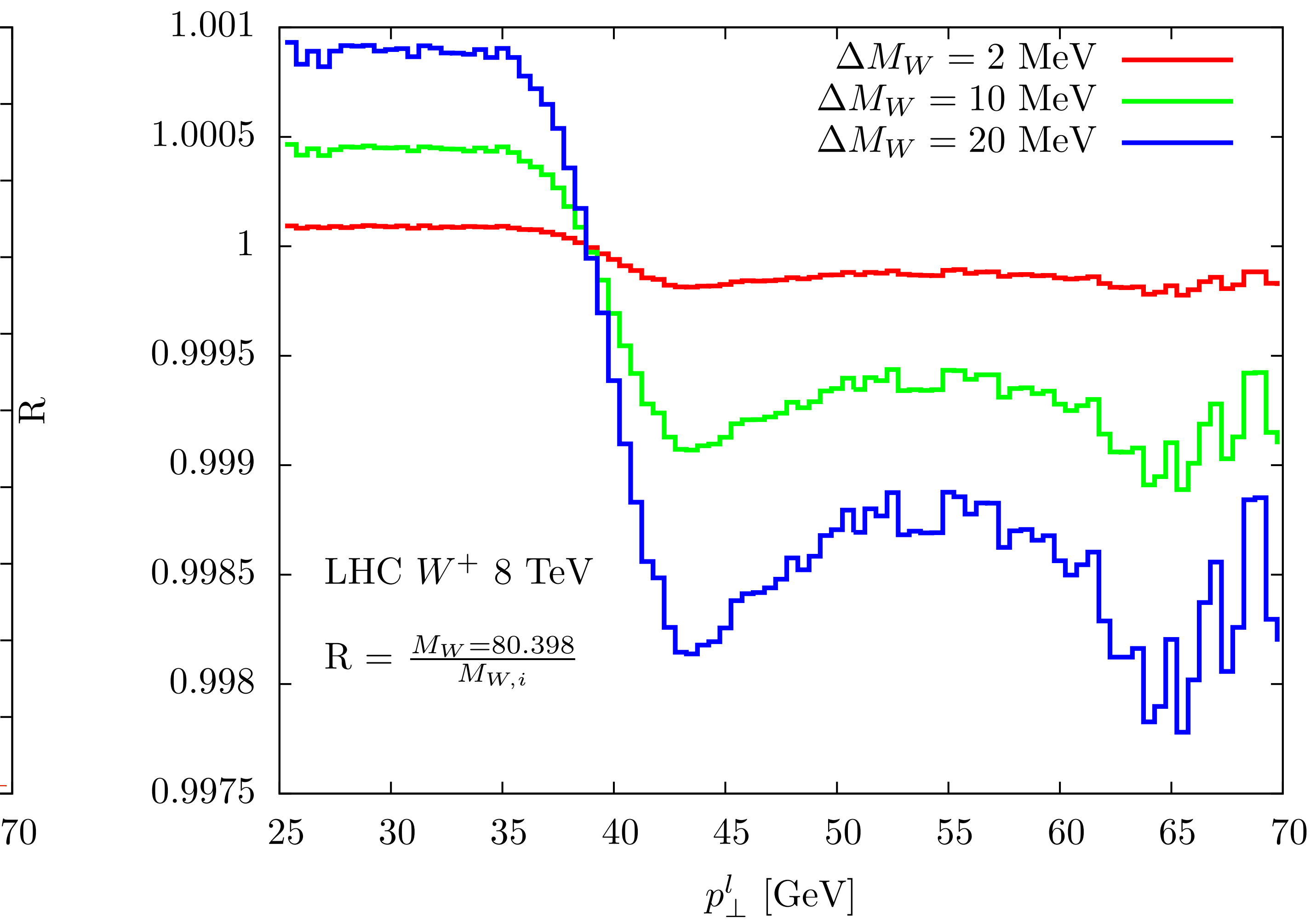
Back-up

High precision determination of the SM parameters at the LHC

The SM parameters are extracted from the data via template fitting. Templates = theoretical histograms of the kinematical distr.
The template theoretical uncertainties propagate as systematic errors on the determination of $(\alpha, G_\mu, m_W, m_Z, \sin^2\theta_{eff}, \dots)$

Given the very high precision goal
$$\delta m_W/m_W \sim 1\cdot 10^{-4}, \qquad \delta \sin^2\theta_{eff}/\sin^2\theta_{eff} \sim 1\cdot 10^{-3}$$
control on the shape of the distributions
at the sub-percent level is needed,
at a hadron collider…





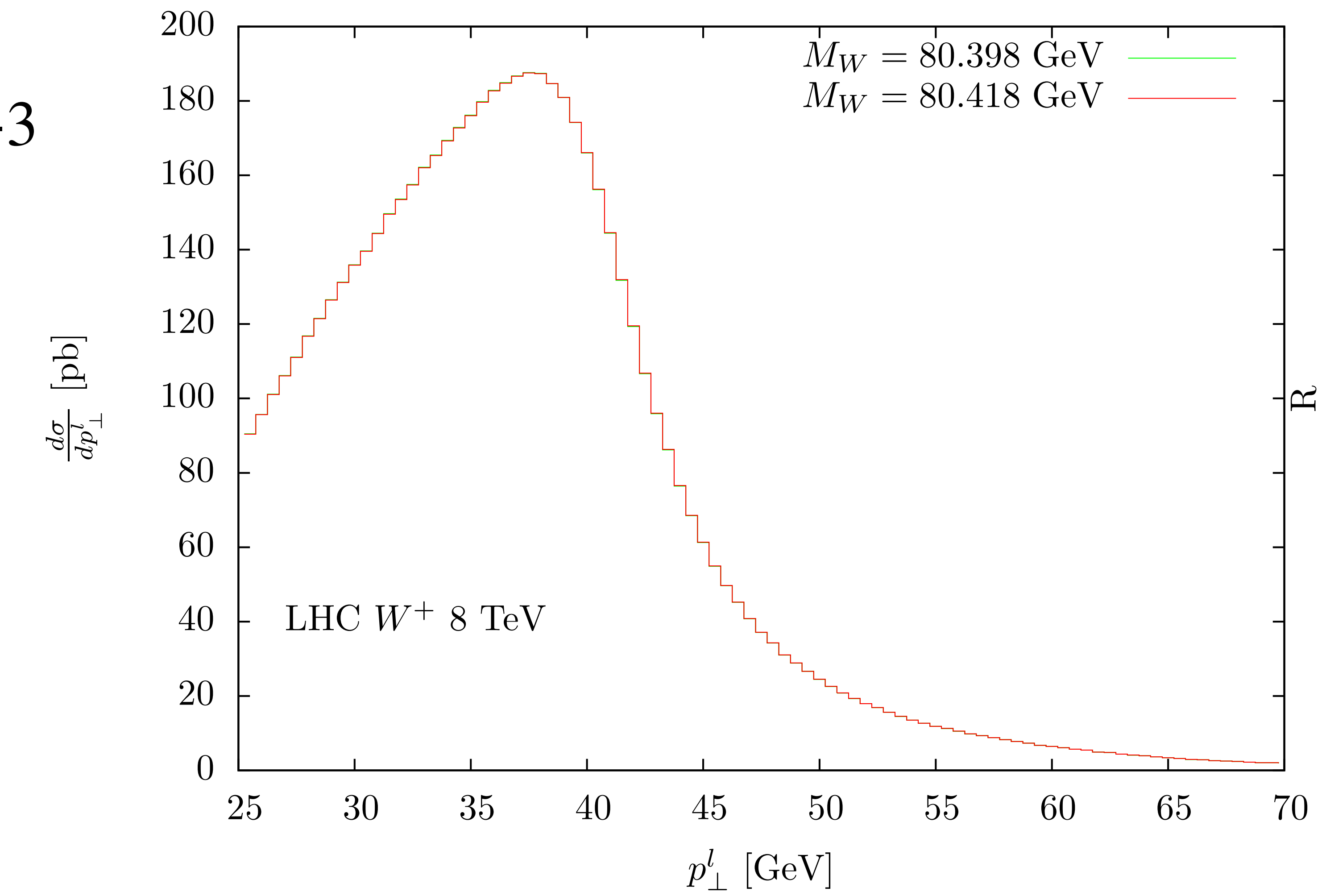
High precision determination of the SM parameters at the LHC

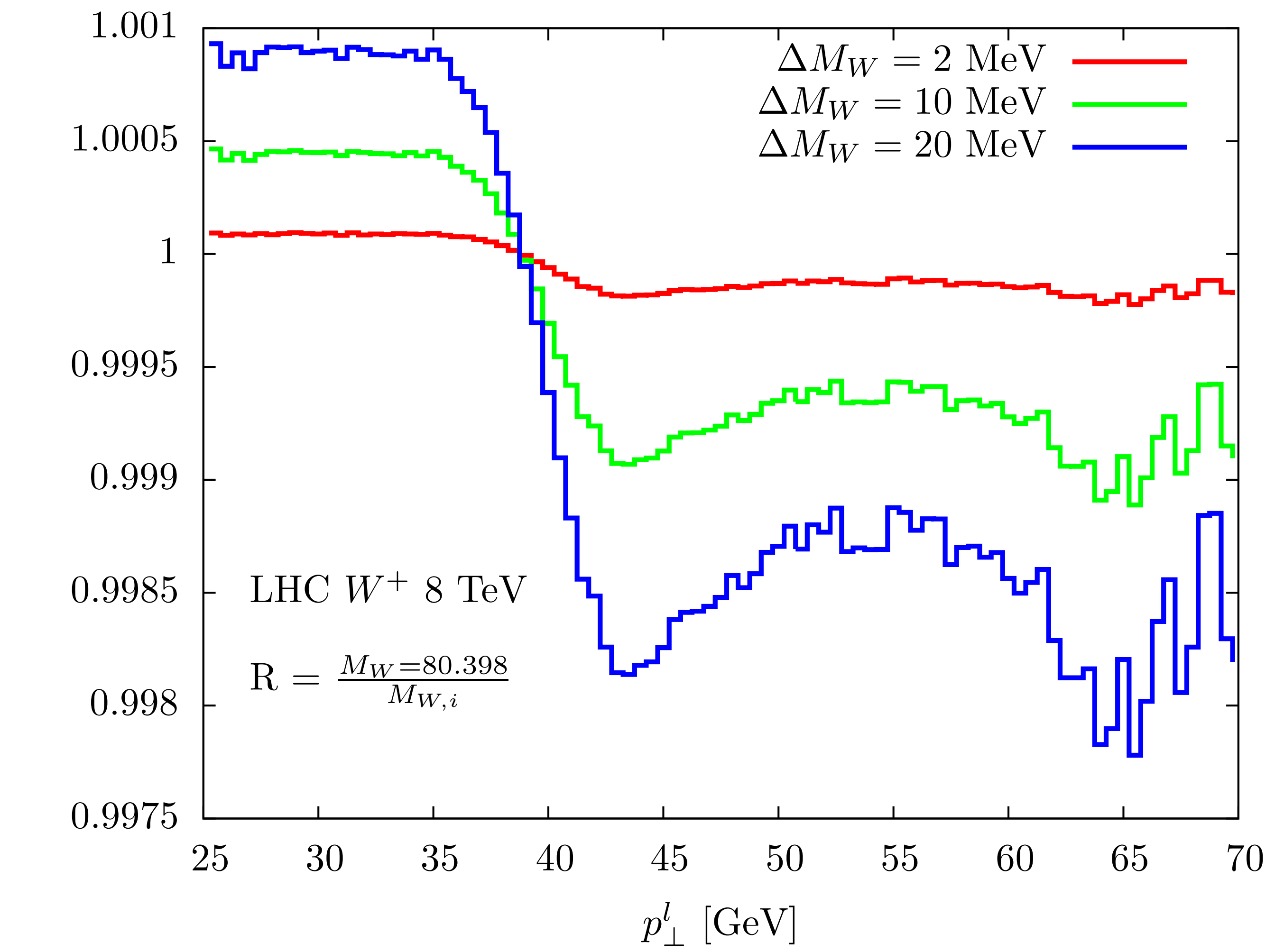
The SM parameters are extracted from the data via template fitting. Templates = theoretical histograms of the kinematical distr.
The template theoretical uncertainties propagate as systematic errors on the determination of $(\alpha, G_\mu, m_W, m_Z, \sin^2\theta_{eff}, \ldots)$

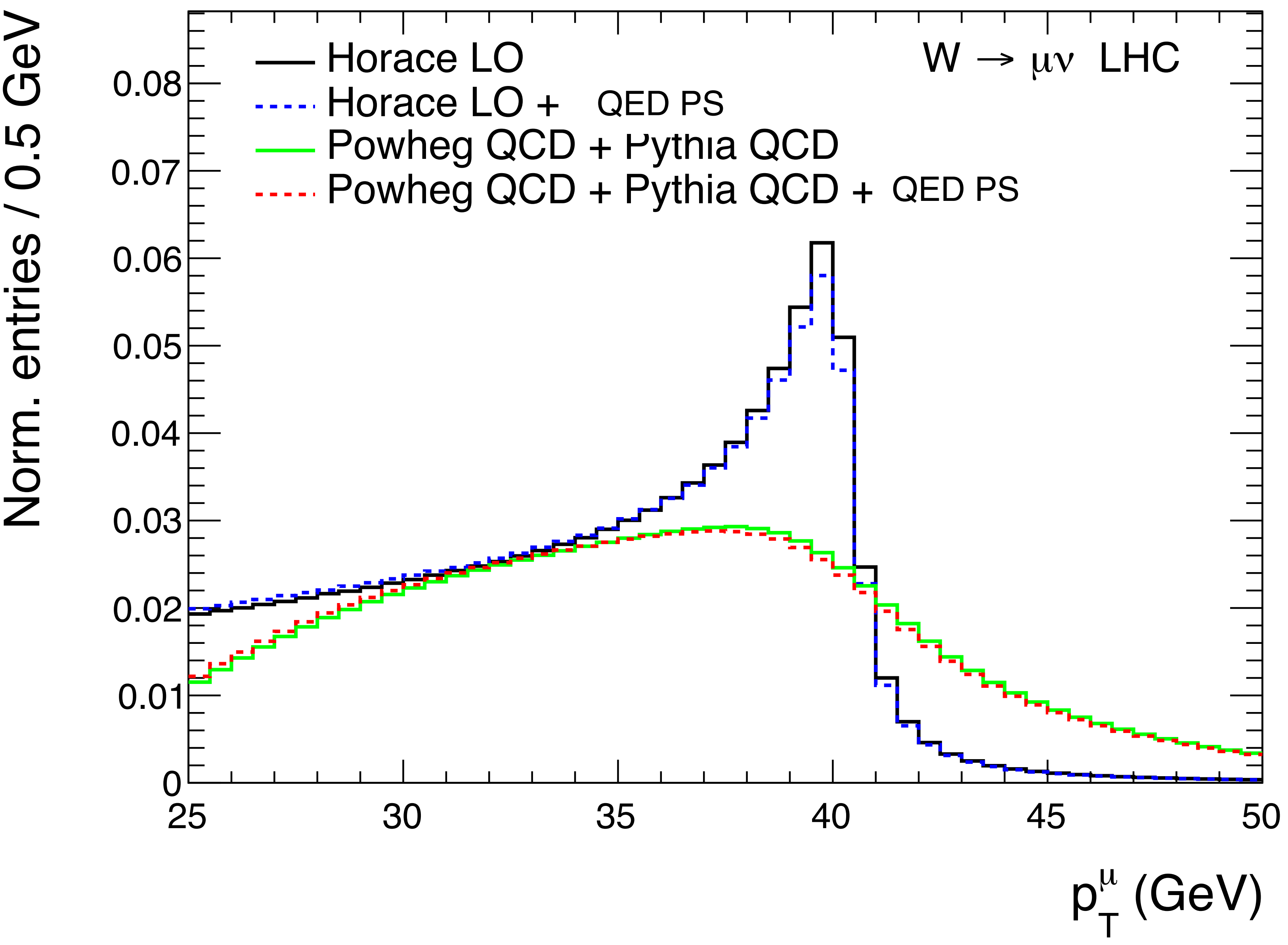
Given the very high precision goal

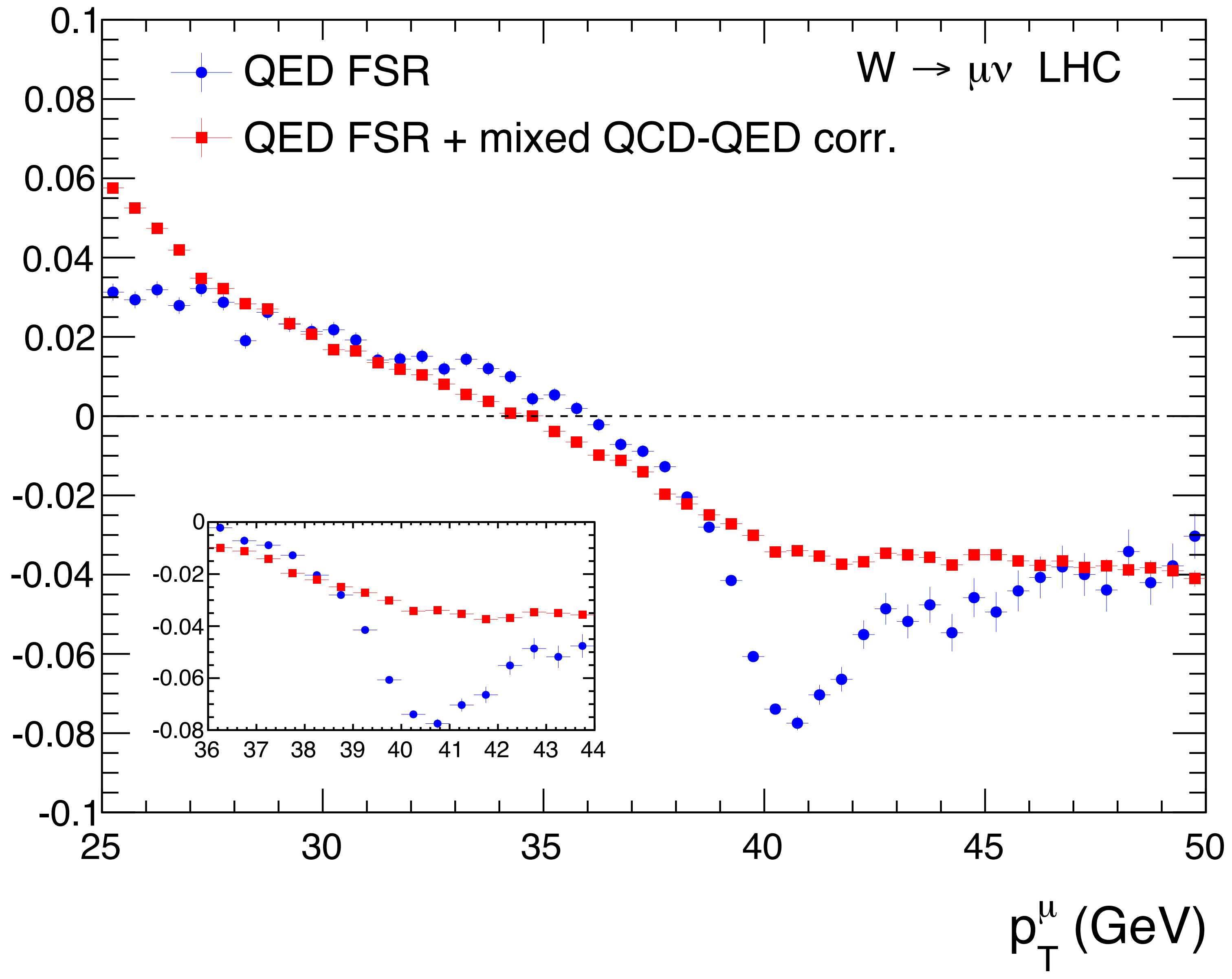
$$\delta m_W/m_W \sim 1\cdot 10^{-4}, \qquad \delta \sin^2\theta_{eff}/\sin^2\theta_{eff} \sim 1\cdot 10^{-3}$$

control on the shape of the distributions
at the sub-percent level is needed,
at a hadron collider…









The convolution of QED-FSR with QCD catches the bulk of the radiative effects.

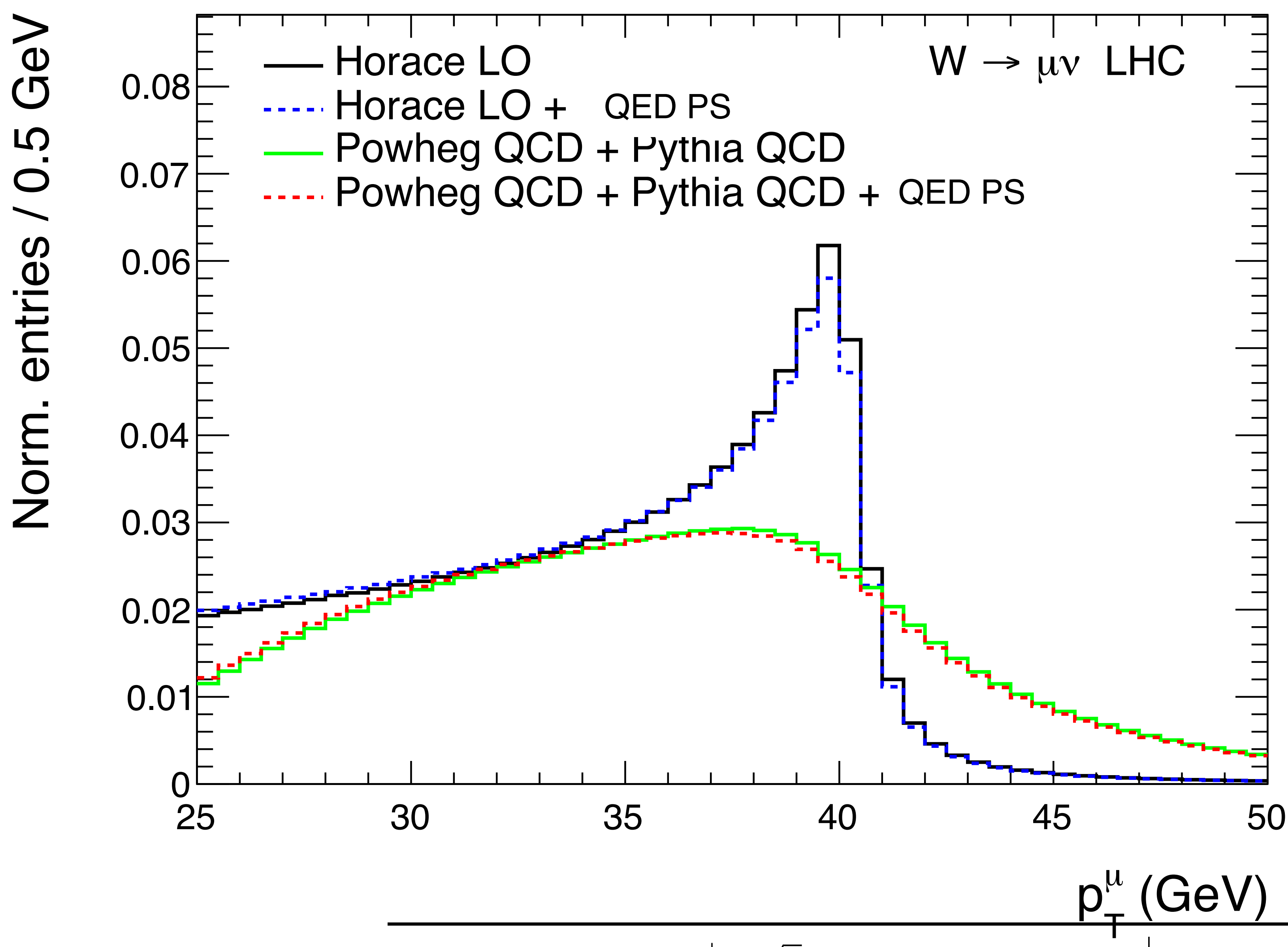
Mixed effects are still very large

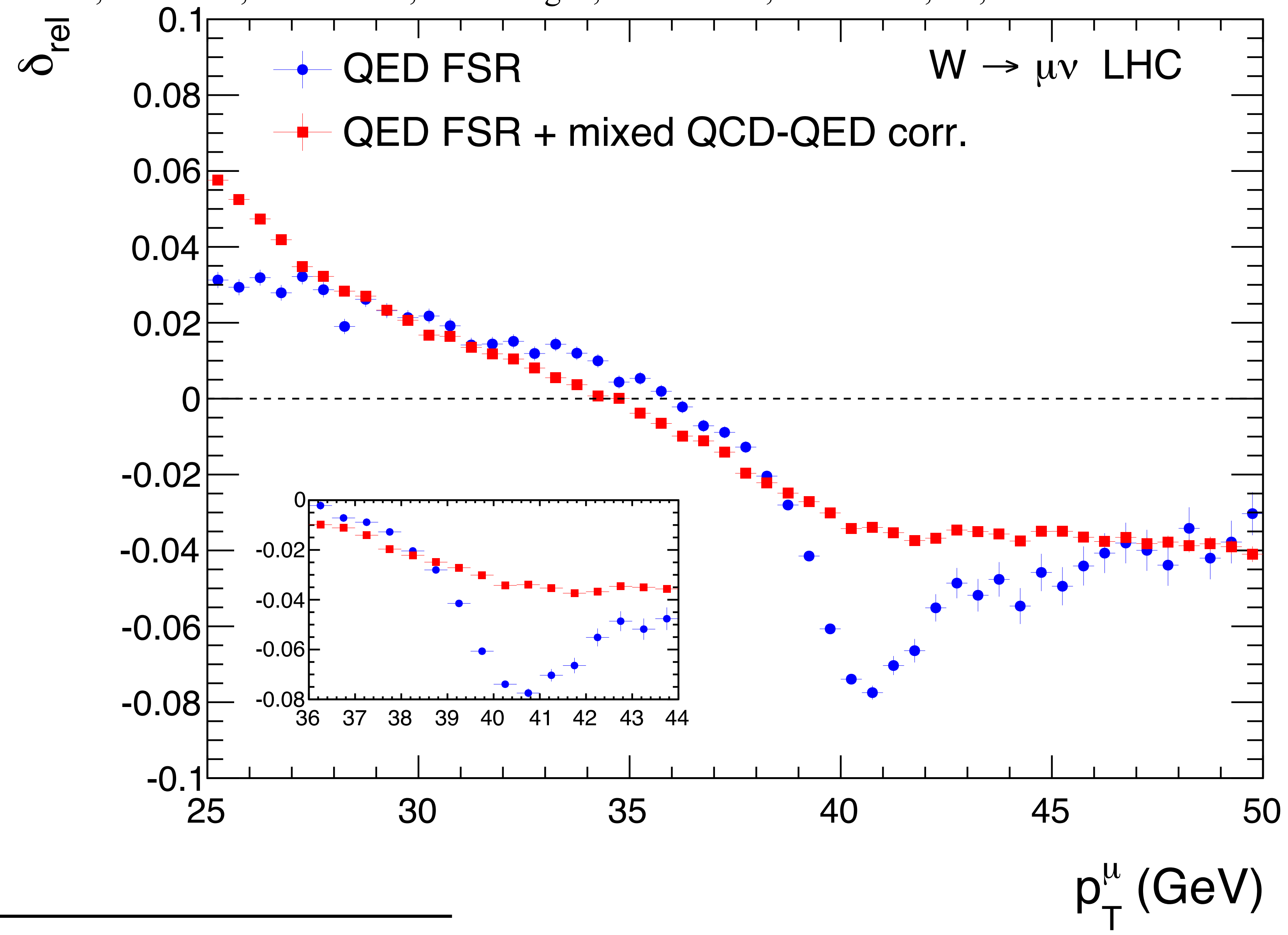
No uncertainty is assigned to this combination

→ need of full NNLO QCD-EW results

Combined QCD-EW simulation tools: impact of QED-FSR on MW

C.Carloni Calame, M.Chiesa, H.Martinez, G.Montagna, O.Nicrosini, F.Piccinini, AV, arXiv:1612.02841





$pp \to W^+$, $\sqrt{s} = 14$ TeV; Templates accuracy: LO; Pseudo–data accuracy		M_W shifts (MeV) $W^+ \to \mu^+\nu$ M_T	$W^+ \to \mu^+\nu$ p_T^ℓ	$W^+ \to e^+\nu$ M_T	$W^+ \to e^+\nu$ p_T^ℓ
1	HORACE only FSR-LL at $\mathcal{O}(\alpha)$	-94±1	-104±1	-204±1	-230±2
2	HORACE FSR-LL	-89±1	-97±1	-179±1	-195±1
3	HORACE NLO-EW with QED shower	-90±1	-94±1	-177±1	-190±2
4	HORACE FSR-LL + Pairs	-94±1	-102±1	-182±2	-199±1
5	PHOTOS FSR-LL	-92±1	-100±2	-182±1	-199±2

$pp \to W^+$, $\sqrt{s} = 14$ TeV; Templates accuracy: NLO-QCD+QCD$_{PS}$; Pseudodata accuracy		QED FSR	M_W shifts (MeV) $W^+ \to \mu^+\nu$ M_T	$W^+ \to \mu^+\nu$ p_T^ℓ	$W^+ \to e^+\nu$(dres) M_T	$W^+ \to e^+\nu$(dres) p_T^ℓ
1	NLO-QCD+(QCD+QED)$_{PS}$	PYTHIA	-95.2±0.6	-400±3	-38.0±0.6	-149±2
2	NLO-QCD+(QCD+QED)$_{PS}$	PHOTOS	-88.0±0.6	-368±2	-38.4±0.6	-150±3
3	NLO-(QCD+EW)+(QCD+QED)$_{PS}$two-rad	PYTHIA	-89.0±0.6	-371±3	-38.8±0.6	-157±3
4	NLO-(QCD+EW)+(QCD+QED)$_{PS}$two-rad	PHOTOS	-88.6±0.6	-370±3	-39.2±0.6	-159±2

the impact on MW of the mixed QCD QED-FSR corrections strongly depends on the underlying QCD shape/model

given that the bulk of the corrections is included in the analyses
- what is the associated uncertainty ?
- what happens if we change the underlying QCD model ?

can we constrain the formulation, for the $\alpha\alpha_s$ contribution ?

The Neutral Current Drell-Yan cross section in the SM: perturbative expansion

$$\begin{aligned}\sigma(h_1 h_2 \to \ell\bar{\ell} + X) = \sigma^{(0,0)} +& \\ &\alpha_s\,\sigma^{(1,0)} + \alpha\,\sigma^{(0,1)} + \\ &\alpha_s^2\,\sigma^{(2,0)} + \alpha\,\alpha_s\,\sigma^{(1,1)} + \alpha^2\,\sigma^{(0,2)} + \\ &\alpha_s^3\,\sigma^{(3,0)} + \ldots\end{aligned}$$

$$\sigma(h_1 h_2 \to l\bar{l} + X) = \sum_{i,j=q\bar{q},g,\gamma} \int dx_1\, dx_2\, f_i^{h_1}(x_1,\mu_F) f_j^{h_2}(x_2,\mu_F)\, \hat{\sigma}(ij \to l\bar{l} + X)$$

$\sigma^{(1,1)}$ requires the evaluation of the xsecs of the following processes, including photon-induced

0 additional partons	$q\bar{q} \to l\bar{l},\ \gamma\gamma \to l\bar{l}$	including virtual corrections of $\mathcal{O}(\alpha_s), \mathcal{O}(\alpha), \mathcal{O}(\alpha\alpha_s)$
1 additional parton	$q\bar{q} \to l\bar{l}g,\ qg \to l\bar{l}q$	including virtual corrections of $\mathcal{O}(\alpha)$
	$q\bar{q} \to l\bar{l}\gamma,\ q\gamma \to l\bar{l}q$	including virtual corrections of $\mathcal{O}(\alpha_s)$
2 additional partons	$q\bar{q} \to l\bar{l}g\gamma, qg \to l\bar{l}q\gamma, q\gamma \to l\bar{l}qg, g\gamma \to l\bar{l}q\bar{q}$	
	$q\bar{q} \to l\bar{l}q\bar{q},\ q\bar{q} \to l\bar{l}q'\bar{q}',\ qq' \to l\bar{l}qq',\ q\bar{q}' \to l\bar{l}q\bar{q}',\ qq \to l\bar{l}qq$	at tree level

Computational framework

The complete calculation has been included in the Munich/Matrix framework

- fully automatic generation and bookkeeping of all the double-real and real-virtual contributions based on an interface with OpenLoops and Recola/Collier
- the 2-loop virtual corrections are separately computed and provided in fast-evaluation format

In this specific framework, main compatibility requirement to include the double-virtual corrections:

the q_T-subtraction formalism to handle the IR singularities (Catani, Grazzini, 2007)

Upon inclusion of the appropriate scheme-dependent subtraction term,

the double virtual corrections can be used with any other simulation code

General structure of the inclusive cross section and the q_T-subtraction formalism

$$d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathscr{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right]_{q_T/Q > r_{cut}}$$

IR structure associated to the QCD-QED part derived from NNLO-QCD results via abelianisation
(de Florian, Rodrigo, Sborlini, 2016, de Florian, Der , Fabre, 2018)

the q_T-subtraction formalism has been extended to the case of final-state emitters (heavy quarks in QCD, leptons in EW)
(Catani, Torre, Grazzini, 2014, Buonocore, Grazzini, Tramontano 2019.)

General structure of the inclusive cross section and the q_T-subtraction formalism

$$\boxed{d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathscr{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)}\right]_{q_T/Q>r_{cut}}}$$

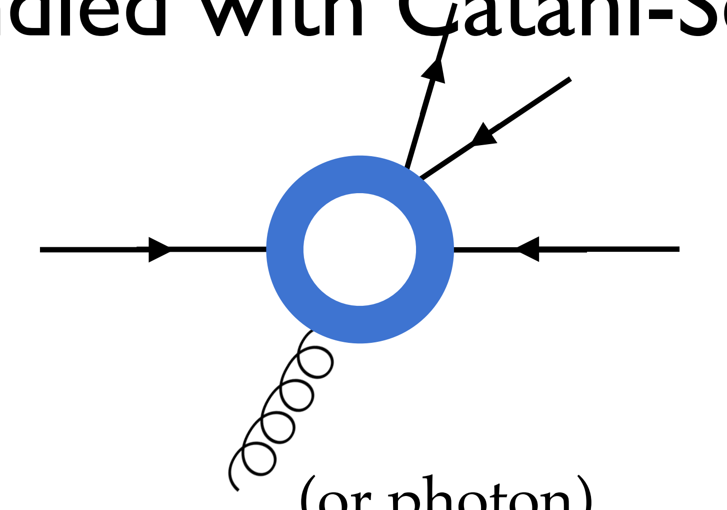
IR structure associated to the QCD-QED part derived from NNLO-QCD results via abelianisation
(de Florian, Rodrigo, Sborlini, 2016, de Florian, Der , Fabre, 2018)

the q_T-subtraction formalism has been extended to the case of final-state emitters (heavy quarks in QCD, leptons in EW)
(Catani, Torre, Grazzini, 2014, Buonocore, Grazzini, Tramontano 2019.)

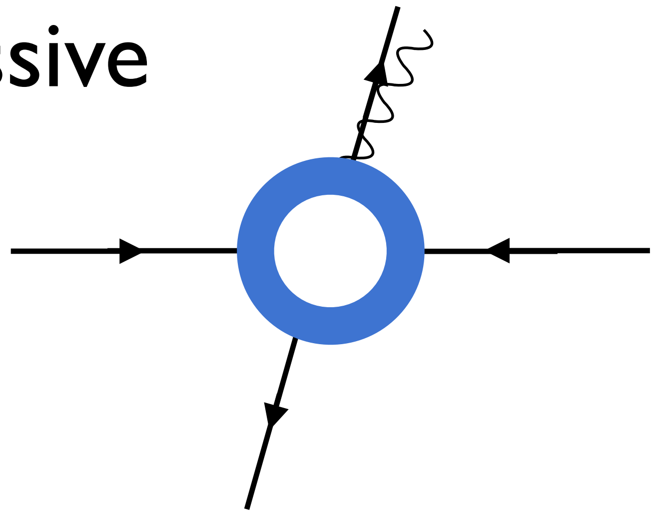
the gauge-boson phase space is split into $q_T = 0$ and $q_T > 0$ regions

$r_{cut} = q_T^{cut}/Q$

for ISR, if $q_T > 0$ the emitted parton is always resolved
and the process under study receives only NLO corrections
which can be handled with Catani-Seymour dipoles



in the FSR case, with $q_T > 0$,
the emitted parton is always resolved
only if the emitter is massive



the final state consists of a pair of massive leptons (treated as bare) to regulate the collinear (mass) singularities

The q_T-subtraction and the residual cut-off dependency

$$d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathscr{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right]_{q_T/Q > r_{cut}}$$

When $q_T/Q > r_{cut}$ the double-real and the real-virtual contributions, subtracted with CS dipoles, are finite

$d\sigma_{CT}^{(1,1)}$ is obtained by expanding to fixed order the q_T resummation formula

The q_T-subtraction and the residual cut-off dependency

$$d\sigma = \sum_{m,n=0}^{\infty} d\sigma^{(m,n)} \qquad d\sigma^{(1,1)} = \mathscr{H}^{(1,1)} \otimes d\sigma_{LO} + \left[d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right]_{q_T/Q > r_{cut}}$$

When $q_T/Q > r_{cut}$ the double-real and the real-virtual contributions, subtracted with CS dipoles, are finite

$d\sigma_{CT}^{(1,1)}$ is obtained by expanding to fixed order the q_T resummation formula

Logarithmic sensitivity on r_{cut} in the double unresolved limit $\int d\sigma_R^{(1,1)} \sim \sum_{i=1}^{4} c_i \ln^i r_{cut} + c_0 + \mathcal{O}(r_{cut}^m)$

The counterterm removes the IR sensitivity to the cutoff variable $\int \left(d\sigma_R^{(1,1)} - d\sigma_{CT}^{(1,1)} \right) \sim c_0 + \mathcal{O}(r_{cut}^m)$

→ we need small values of the cutoff

→ explicit numerical tests to quantify the bias induced by the cutoff choice (cfr. Buonocore, Kallweit, Rottoli, Wiesemann, arXiv:2111.13661 Camarda, Cieri, Ferrera, arXiv:2111.14509)

we can fit the r_{cut} dependence and extrapolate in the $r_{cut} \to 0$ limit

Subtraction of the IR divergences from the 2-loop amplitude

we identify QCD-QED (poles up to $1/\varepsilon^4$) and QCD-weak (poles up to $1/\varepsilon^2$ with cumbersome coefficients) diagrams

$$|\mathcal{M}^{(1,0),fin}\rangle = |\mathcal{M}^{(1,0)}\rangle - \mathcal{I}^{(1,0)}|\mathcal{M}^{(0)}\rangle\,,$$

standard NLO-QCD subtraction

$$|\mathcal{M}^{(0,1),fin}\rangle = |\mathcal{M}^{(0,1)}\rangle - \mathcal{I}^{(0,1)}|\mathcal{M}^{(0)}\rangle\,.$$

NLO-EW subtraction, with massive leptons

$$|\mathcal{M}^{(1,1),fin}\rangle = |\mathcal{M}^{(1,1)}\rangle - \mathcal{I}^{(1,1)}|\mathcal{M}^{(0)}\rangle - \tilde{\mathcal{I}}^{(0,1)}|\mathcal{M}^{(1,0),fin}\rangle - \tilde{\mathcal{I}}^{(1,0)}|\mathcal{M}^{(0,1),fin}\rangle\,.$$

$$\mathcal{I}^{(1,0)} = \left(\frac{\alpha_s}{4\pi}\right)\left(\frac{s}{\mu^2}\right)^{-\epsilon} C_F\left(-\frac{2}{\epsilon^2} - \frac{1}{\epsilon}(3+2i\pi) + \zeta_2\right)\,,$$

$$\mathcal{I}^{(0,1)} = \left(\frac{\alpha}{4\pi}\right)\left(\frac{s}{\mu^2}\right)^{-\epsilon}\left[Q_u^2\left(-\frac{2}{\epsilon^2} - \frac{1}{\epsilon}(3+2i\pi) + \zeta_2\right) + \frac{4}{\epsilon}\Gamma_l^{(0,1)}\right],$$

$$\Gamma_l^{(0,1)} = Q_u Q_l \log\left(\frac{2p_1.p_3}{2p_2.p_3}\right) + \frac{Q_l^2}{2}\left(-1 - \frac{1+x_l^2}{1-x_l^2}\log(x_l)\right).$$

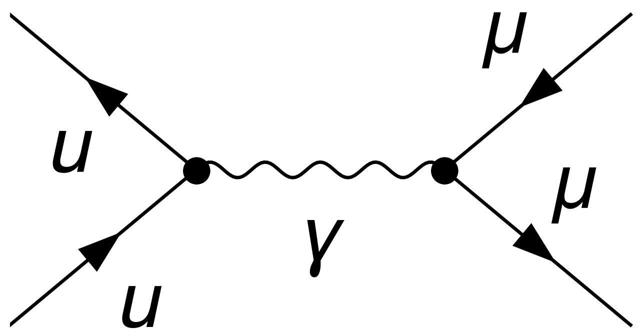
$$\mathcal{I}^{(1,1)} = \left(\frac{\alpha_s}{4\pi}\right)\left(\frac{\alpha}{4\pi}\right)\left(\frac{s}{\mu^2}\right)^{-2\epsilon} C_F Q_u^2\left(\frac{4}{\epsilon^4} + \frac{1}{\epsilon^3}(12+8i\pi) + \frac{1}{\epsilon^2}(9 - 28\zeta_2 + 12i\pi)\right.$$
$$\left. + \frac{1}{\epsilon}\left(-\frac{3}{2} + 6\zeta_2 - 24\zeta_3 - 4i\pi\zeta_2\right)\right).$$

$2\mathrm{Re}\langle \mathscr{M}^{(0,0)} | \mathscr{M}^{(1,1),fin}\rangle$ is free of any singularity

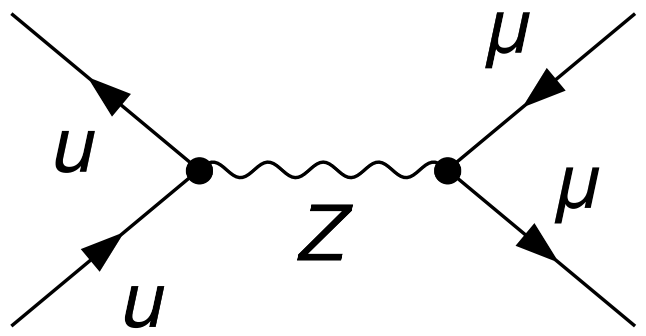
the analytical check of the cancellation of the IR poles in the QCD-weak sector is one very demanding test of the calculation

The double virtual amplitude: generation of the amplitude

$$\mathcal{M}^{(0,0)}(q\bar{q} \to l\bar{l}) =$$



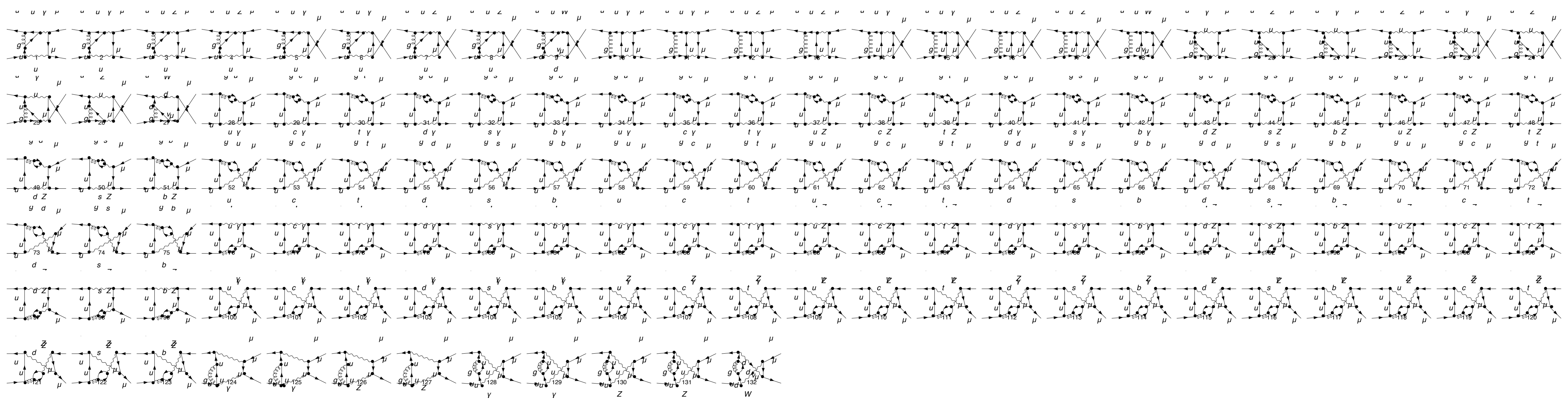
1

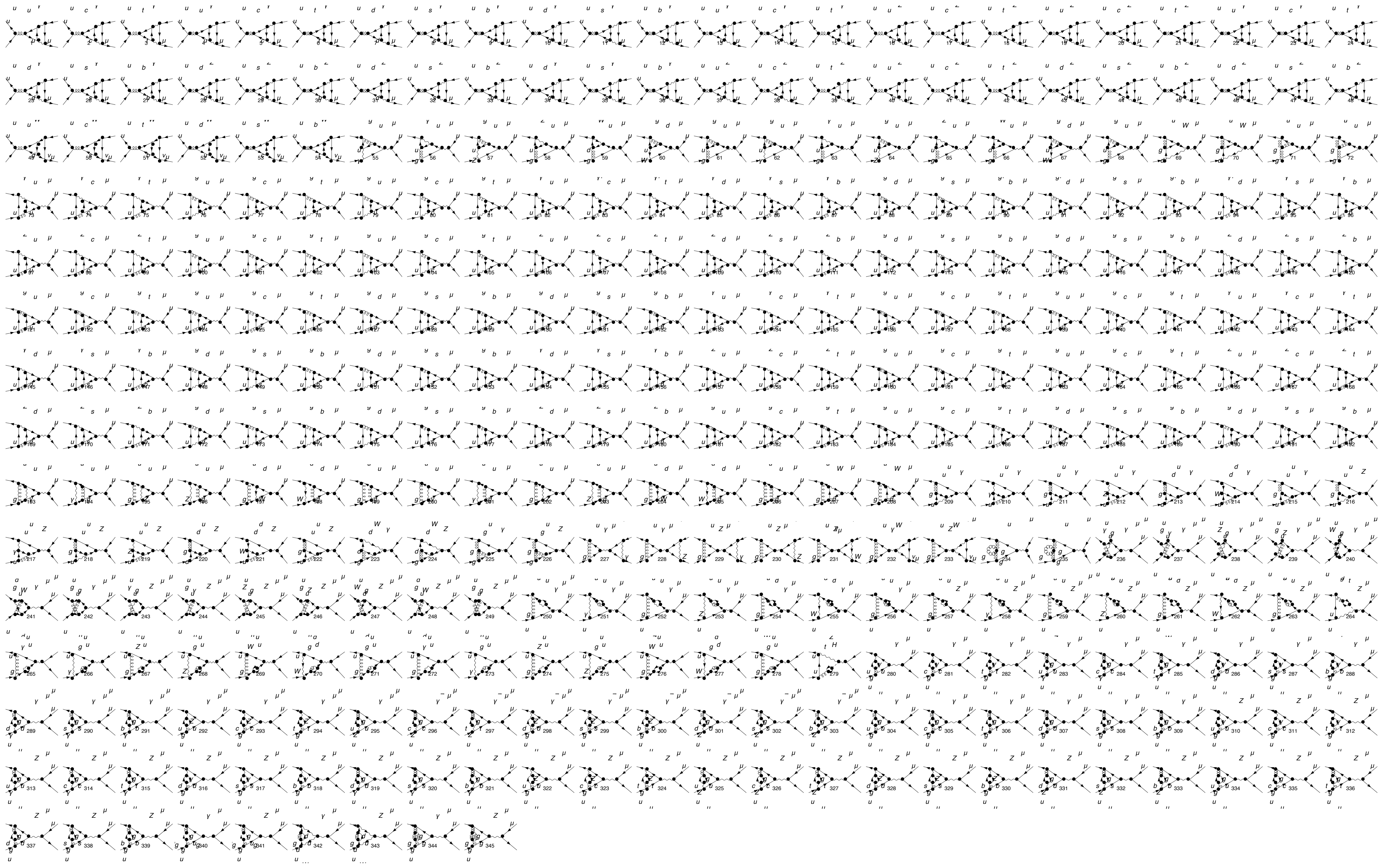


2

$$\mathcal{M}^{(1,1)}(q\bar{q} \to l\bar{l}) =$$

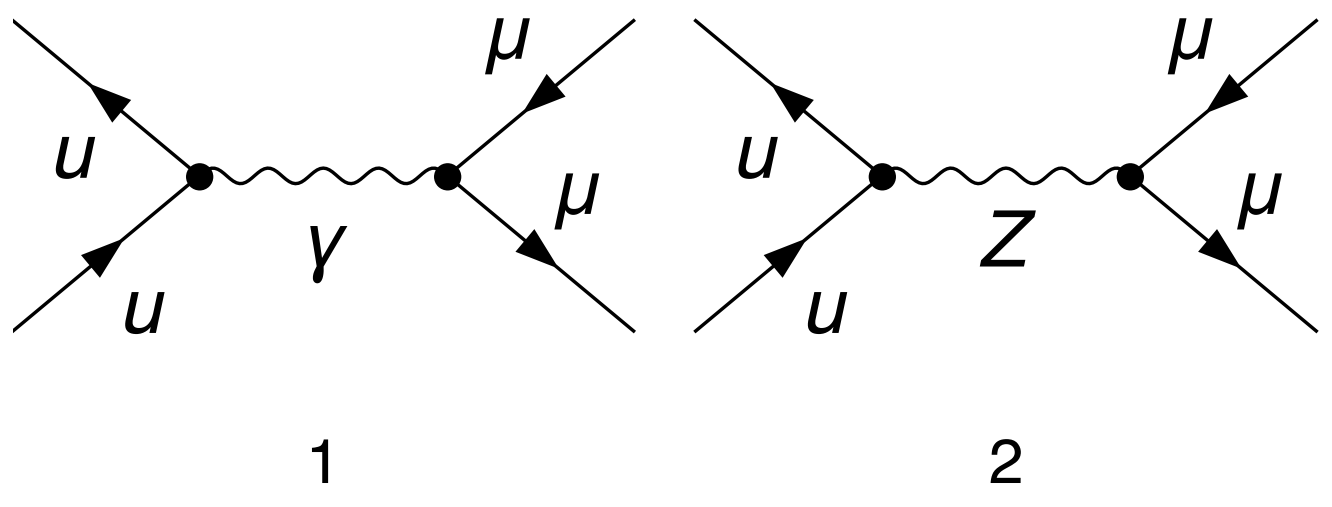
O(1000) self-energies + O(300) vertex corrections +O(130) box corrections + 1loop x 1loop
(before discarding all those vanishing for colour conservation, e.g. no fermonic triangles)





The double virtual amplitude: generation of the amplitude

$$\mathscr{M}^{(0,0)}(q\bar{q} \to l\bar{l}) =$$



$$\mathscr{M}^{(1,1)}(q\bar{q} \to l\bar{l}) =$$

O(1000) self-energies + O(300) vertex corrections +O(130) box corrections + 1loop x 1loop
(before discarding all those vanishing for colour conservation, e.g. no fermonic triangles)

Two independent calculations based on QGraf and FeynArts in the EW Background Field Gauge

The BFG choice guarantees the validity of EW Ward identities for the initial state vertex → additional technical checks

- UV finiteness when combining 2-loop vertex and quark WF in the full EW SM → that combination has only IR poles
- UV renormalisation is confined to the gauge-boson propagators sector, where IR divergences are absent

The 1-loop check of the gauge-parameter independence identifies those subsets of diagrams yielding the cancellation.

The 2-loop calculation is organised splitting the total amplitude in the combination of different subsets, according to their EW charges (# of Ws, Zs, γs)

The double virtual amplitude: UV renormalization

G.Degrassi, AV, hep-ph/0307122, S.Dittmaier,T.Schmidt,J.Schwarz, arXiv:2009.02229 S.Dittmaier, arXiv:2101.05154

Complex mass scheme

$$\mu_{W0}^2 = \mu_W^2 + \delta\mu_W^2, \quad \mu_{Z0}^2 = \mu_Z^2 + \delta\mu_Z^2, \quad e_0 = e + \delta e$$

$$\frac{\delta s^2}{s^2} = \frac{c^2}{s^2}\left(\frac{\delta\mu_Z^2}{\mu_Z^2} - \frac{\delta\mu_W^2}{\mu_W^2}\right)$$

the mass counterterms are defined
at the complex pole of the propagator

the weak mixing angle is complex valued $c^2 \equiv \mu_W^2/\mu_Z^2$

BFG EW Ward identity → cancellation of the UV divergences combining vertex and fermion WF corrections

The bare couplings of Z and photon to fermions in the (G_μ, μ_W, μ_Z) input scheme are given by

$$\frac{g_0}{c_0} = \sqrt{4\sqrt{2}G_\mu\mu_Z^2}\left[1 - \frac{1}{2}\Delta r + \frac{1}{2}\left(2\frac{\delta e}{e} + \frac{s^2-c^2}{c^2}\frac{\delta s^2}{s^2}\right)\right] \equiv \sqrt{4\sqrt{2}G_\mu\mu_Z^2}\left(1 + \delta g_Z^{G_\mu}\right)$$

$$g_0 s_0 = \sqrt{4\sqrt{2}G_\mu\mu_W^2 s^2}\left[1 + \tfrac{1}{2}\left(-\Delta r + 2\tfrac{\delta e}{e}\right)\right] \equiv e_{ren}^{G_\mu}\left(1 + \delta g_A^{G_\mu}\right)$$

Gauge boson renormalised propagators

$$\Sigma_{R,T}^{AA}(q^2) = \Sigma_T^{AA}(q^2) + 2\,q^2\,\delta g_A$$
$$\Sigma_{R,T}^{ZZ}(q^2) = \Sigma_T^{ZZ}(q^2) - \delta\mu_Z^2 + 2\,(q^2 - \mu_Z^2)\,\delta g_Z$$

$$\Sigma_{R,T}^{AZ}(q^2) = \Sigma_T^{AZ}(q^2) - q^2\,\frac{\delta s^2}{sc}$$
$$\Sigma_{R,T}^{ZA}(q^2) = \Sigma_T^{ZA}(q^2) - q^2\,\frac{\delta s^2}{sc},$$

After the UV renormalisation, the singular structure is entirely due to IR soft and/or collinear singularities

The double virtual amplitude: γ_5 treatment

The absence of a consistent definition of γ_5 in $n = 4 - 2\varepsilon$ dimensions yields a practical problem

The trace of Dirac matrices and γ_5 is a polynomial in ε
The UV or IR divergences of Feynman integrals appear as poles $1/\varepsilon$

$$\mathrm{Tr}(\gamma_\alpha \ldots \gamma_\mu \gamma_5) \times \int d^n k \frac{1}{[k^2 - m_0^2][(k+q_1)^2 - m_1^2][(k+q_2)^2 - m_2^2]} \sim (a_0 + a_1 \varepsilon + \ldots) \times \left(\frac{c_{-2}}{\varepsilon^2} + \frac{c_{-1}}{\varepsilon} + c_0 + \ldots \right)$$

If a_1 is evaluated in a non-consistent way,
then poles might not cancel and the finite part of the xsec might have a spurious contribution

The double virtual amplitude: γ_5 treatment

The absence of a consistent definition of γ_5 in $n = 4 - 2\varepsilon$ dimensions yields a practical problem

The trace of Dirac matrices and γ_5 is a polynomial in ε
The UV or IR divergences of Feynman integrals appear as poles $1/\varepsilon$

$$Tr(\gamma_\alpha \dots \gamma_\mu \gamma_5) \times \int d^n k \frac{1}{[k^2 - m_0^2][(k+q_1)^2 - m_1^2][(k+q_2)^2 - m_2^2]} \sim (a_0 + a_1 \varepsilon + \dots) \times \left(\frac{c_{-2}}{\varepsilon^2} + \frac{c_{-1}}{\varepsilon} + c_0 + \dots \right)$$

If a_1 is evaluated in a non-consistent way,
then poles might not cancel and the finite part of the xsec might have a spurious contribution

- 't Hooft-Veltman treat γ_5 (anti)commuting in (4) $n-4$ dimensions preserving the cyclicity of the traces (one counterterm is needed)
- Kreimer treats γ_5 anticommuting in n dimensions, abandoning the cyclicity of the traces (→ need of a starting point)

- Heller, von Manteuffel, Schabinger verified that the IR-subtracted squared matrix element are identical in the two approaches

- we adopted the naive anticommuting prescription (Kreimer); we use $\gamma_5 = \frac{i}{4!}\epsilon_{\mu\nu\rho\sigma}\gamma^\mu\gamma^\nu\gamma^\rho\gamma^\sigma$ to compute traces with one γ_5
  - we computed the 2-loop amplitude and, independently, the IR subtraction term; both depend on the prescription chosen
  - the cancellation of all the lowest order poles is checked (and non trivial)
  - absence of fermionic triangles because of colour conservation

The double virtual amplitude: solution and evaluation of the Master Integrals

The system of first-order linear differential equations satisfied by the 36 QCD-weak Master Integrals chosen by Bonciani et al. can be written in dlog form

$$d\mathbf{I} = \epsilon\, d\mathbb{A}\, \mathbf{I}\,, \qquad d\mathbb{A} = \sum_{i=1}^{n} \mathbb{M}_i \; \mathrm{dlog}\, \eta_i\,.$$

The letters η_i provide the complete information about the singular structure of the amplitude

Boundary Conditions

The BCs have been evaluated outside the physical phase space and are expressed in exact form

Master Integrals 1-31

When the letters have a rational (linear) expression, it is possible to integrate the system in terms of GPLs

The appearance, for kinematical reasons, of four square roots among the letters is handled with a change of variables that makes all the new letters linear, leading to a GPL solution in the new variables

Master Integrals 32-36

The appearance of another distinct square root among the letters, makes it impossibile to linearise the weights → the equations are formally solved with a Chen-Goncharov iterated representation

- the poles of these MIs contain Chen-Goncharov functions, but the latter cancel in the physical amplitude
- in the finite part, the Chen-Goncharov functions remain → problems to evaluate the amplitude in the physical region

Total cross section in the fiducial region

$G_\mu = 1.1663787 \times 10^{-5}\,\mathrm{GeV}^{-2}$, $M_W = 80.358$ GeV, $\Gamma_W = 2.084$ GeV, $M_Z = 91.1535$ GeV, $\Gamma_Z = 2.4943$ GeV

$M_H = 125.09$ GeV, $m_t = 173.07$ GeV   NNPDF31_nnlo_as_0118_luxqed

$p_T^{\mu^\pm} > 25\,\mathrm{GeV}$, $|\eta^{\mu^\pm}| < 2.5$, $m_{\mu^+\mu^-} > 50\,\mathrm{GeV}$, $\mu_R = \mu_F = M_Z$

σ [pb]	σ_{LO}	$\sigma^{(1,0)}$	$\sigma^{(0,1)}$	$\sigma^{(2,0)}$	$\sigma^{(1,1)}$
$q\bar{q}$	809.56(1)	191.85(1)	−33.76(1)	49.9(7)	−4.8(3)
qg	—	−158.08(2)	—	−74.8(5)	8.6(1)
$q(g)\gamma$	—	—	−0.839(2)	—	0.084(3)
$q(\bar{q})q'$	—	—	—	6.3(1)	0.19(0)
gg	—	—	—	18.1(2)	—
$\gamma\gamma$	1.42(0)	—	−0.0117(4)	—	—
tot	810.98(1)	33.77(2)	−34.61(1)	−0.5(9)	4.0(3)
$\sigma^{(m,n)}/\sigma_{\mathrm{LO}}$		+4.2 %	−4.3 %	∼ 0 %	+0.5 %

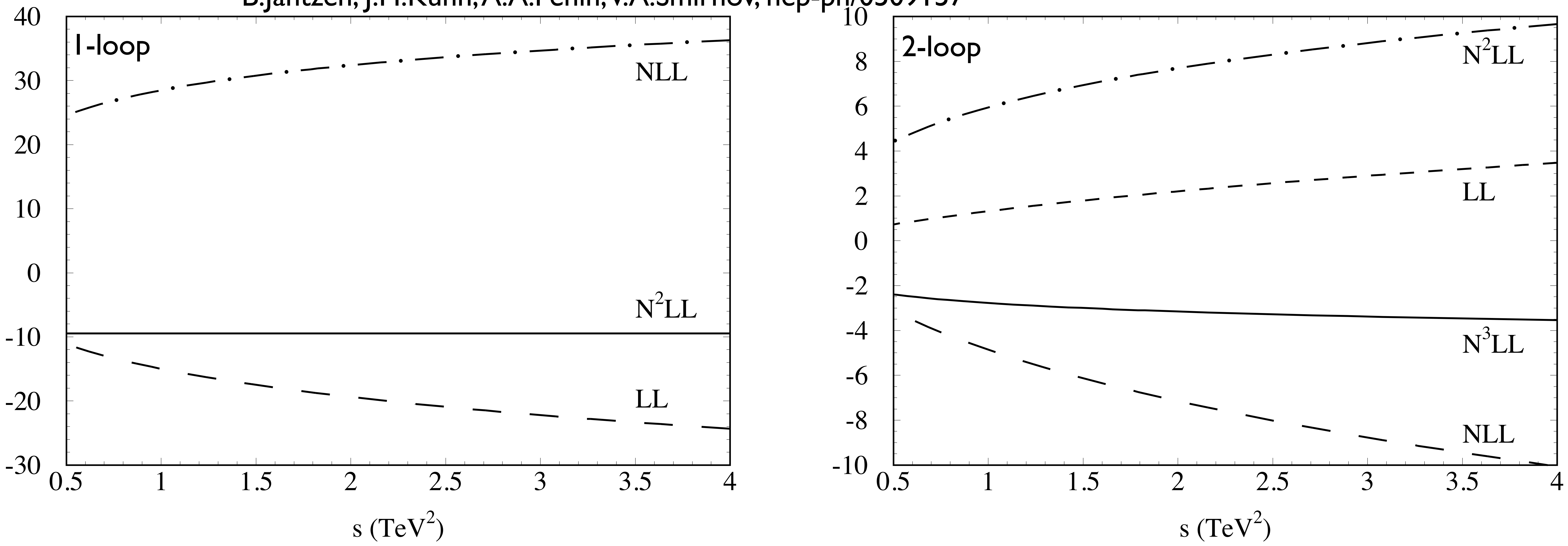
Accidental cancellation of NLO-QCD and NLO-EW, small contribution from NNLO-QCD

→ the NNLO QCD-EW is comparable (or larger) in size than the combination of the previous orders

Towards the NNLO-EW calculation ?

The NNLO-EW corrections could modify in a non-trivial way the large-mass/momentum tails of the distributions
Large logarithmic corrections (EW Sudakov logs) appear in the virtual corrections
At two-loop level, we have up to the fourth power of $\log(s/m_V^2)$,
the different corrections are comparable in size and with alternate signs
→ how can we estimate the constant term ?

B.Jantzen, J.H.Kühn, A.A.Penin, V.A.Smirnov, hep-ph/0509157

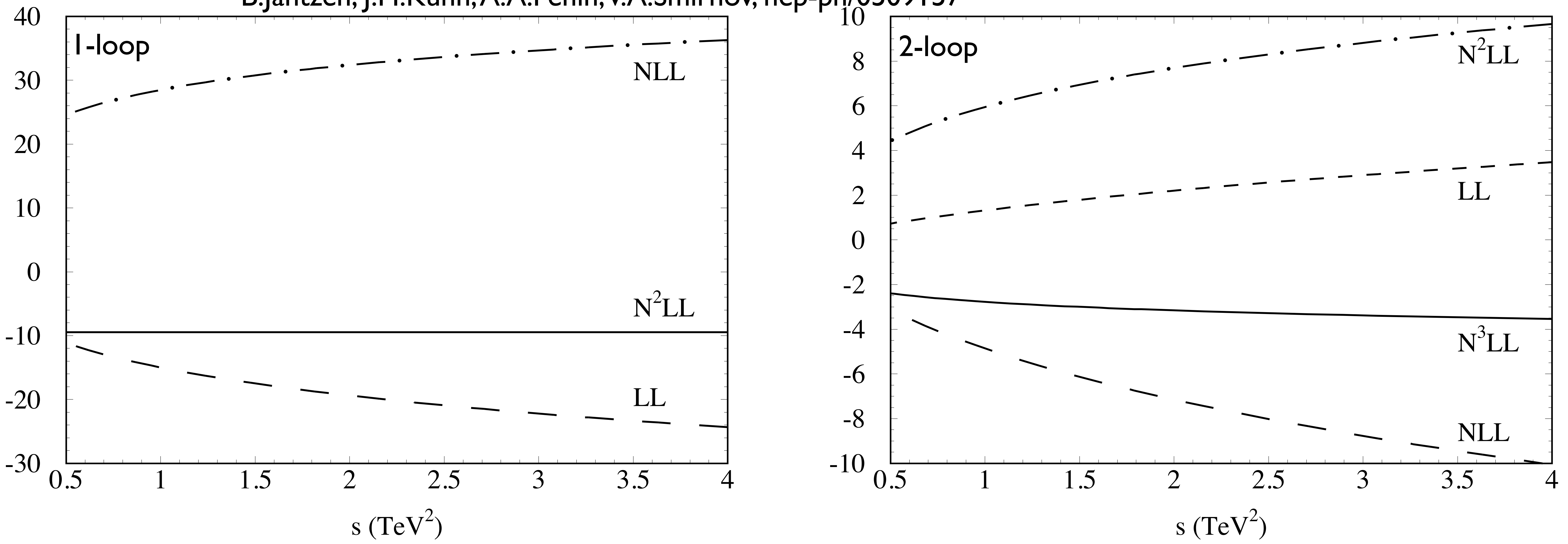


corrections to $e^+e^- \to q\bar{q}$
due to EW Sudakov logs

Towards the NNLO-EW calculation ?

The NNLO-EW corrections could modify in a non-trivial way the large-mass/momentum tails of the distributions
Large logarithmic corrections (EW Sudakov logs) appear in the virtual corrections
At two-loop level, we have up to the fourth power of $\log(s/m_V^2)$,
the different corrections are comparable in size and with alternate signs
→ how can we estimate the constant term ?

B.Jantzen, J.H.Kühn, A.A.Penin, V.A.Smirnov, hep-ph/0509157



corrections to $e^+e^- \to q\bar{q}$
due to EW Sudakov logs

The NNLO-EW corrections will require an extra step compared to the mixed QCD-EW case
- for the number of additional Master Integrals (→ automation)
- for the complexity of the amplitudes (size problems? large cancellations?)
- for the conceptual problems (γ_5?, complex-mass scheme at two-loop?)

but the discussion has started

Differential distributions: exact vs approximated predictions

The exact $\mathcal{O}(\alpha\alpha_s)$ corrections allow to test the validity of different recipes based on NLO-QCD and NLO-EW results

factorised Ansatz

$$\frac{d\sigma_{fact}}{dX} = \frac{d\sigma^{(0,0)}}{dX}\left[1+\frac{d\sigma^{(1,0)}}{dX}\left(\frac{d\sigma^{(0,0)}}{dX}\right)^{-1}\right]\times\left[1+\frac{d\sigma^{(0,1)}}{dX}\left(\frac{d\sigma^{(0,0)}}{dX}\right)^{-1}\right]$$

$$\simeq \frac{d\sigma^{(0,0)}}{dX}+\frac{d\sigma^{(1,0)}}{dX}+\frac{d\sigma^{(0,1)}}{dX}+\frac{d\sigma^{(0,1)}}{dX}\frac{d\sigma^{(1,0)}}{dX}\left(\frac{d\sigma^{(0,0)}}{dX}\right)^{-1}$$

- the last term is absent in a purely additive formulation
- Factorisation is expected to work when both QCD and EW corrections factorise w.r.t. the gauge boson production
- the giant K-factors (qg and qγ processes) should not be applied to photon-induced channels

pole approximation

in the hard coefficient, the 2-loop virtual is approximated by $H_{PA}^{(1,1)} = \frac{2\mathrm{Re}(\mathcal{M}^{(1,1)}\mathcal{M}^{(0,0)*})_{PA}}{|\mathcal{M}_{PA}^{(0,0)}|^2}$

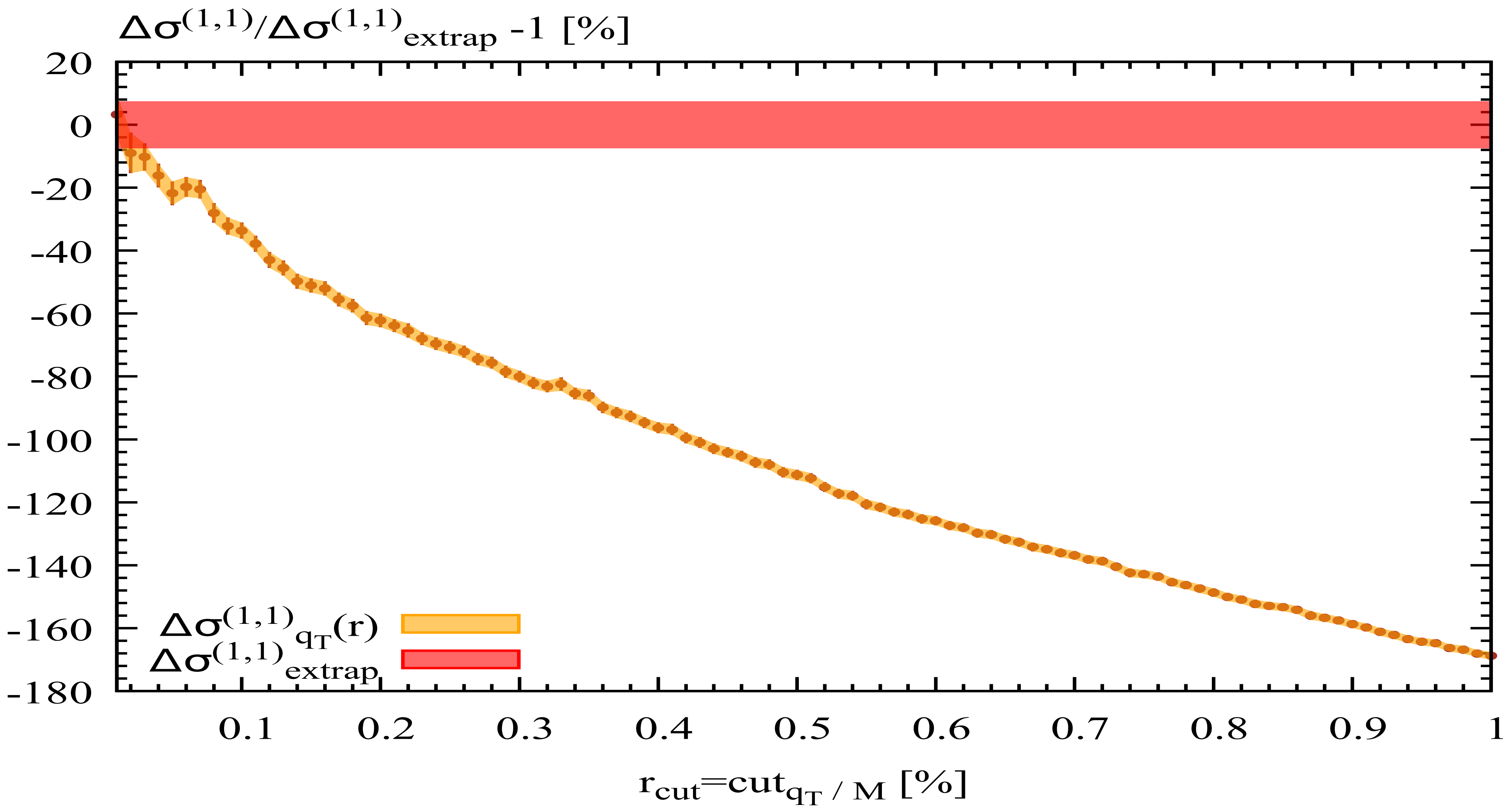
the 2-loop virtual corrections are evaluated in pole approximation
- on-shell Z boson form factor
- the resonant contributions of the γZ box diagrams cancel

Dependence on r_{cut} of the NNLO QCD-EW corrections to NC DY

courtesy of S.Kallweit

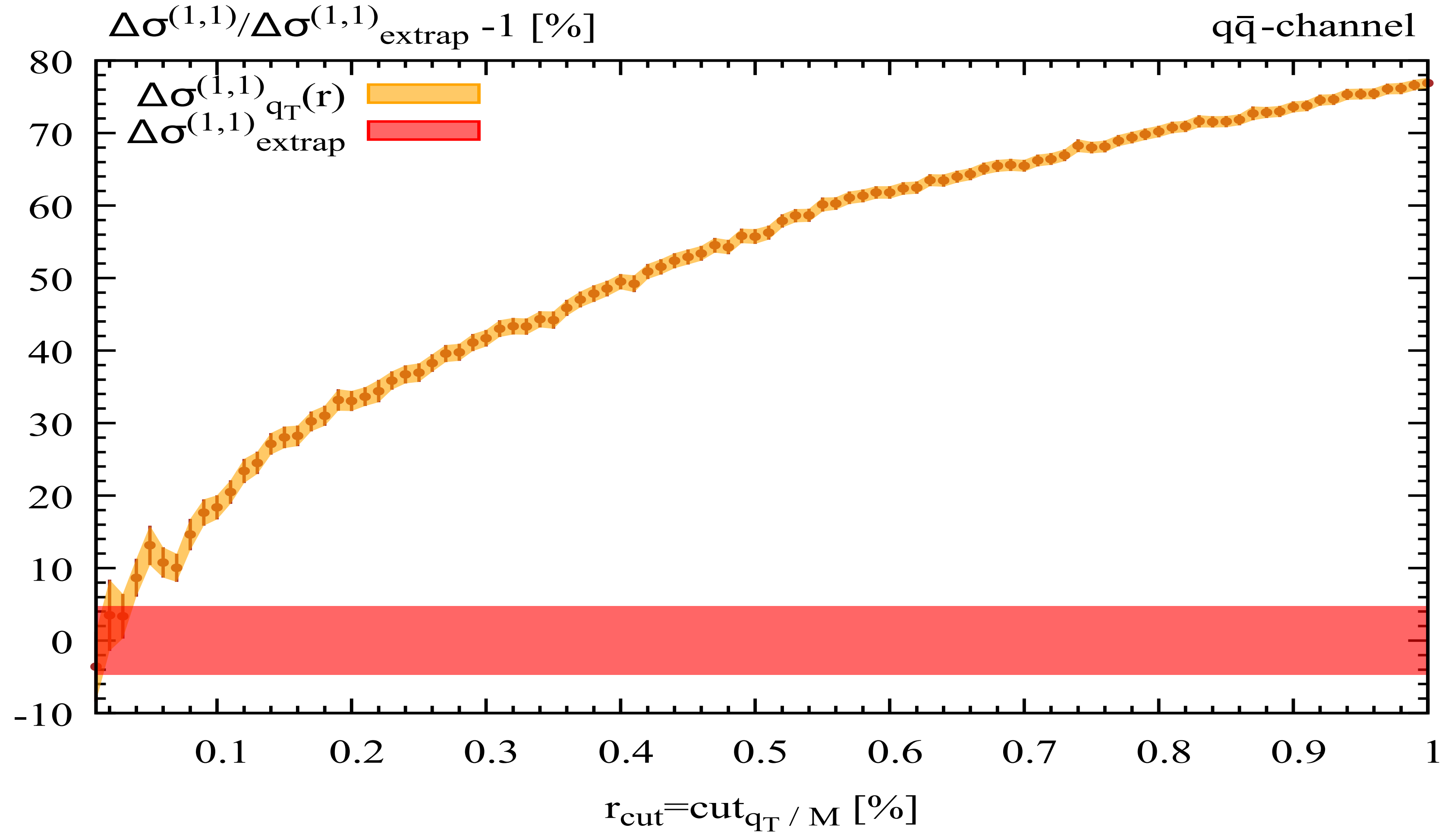
Symmetric-cut scenario

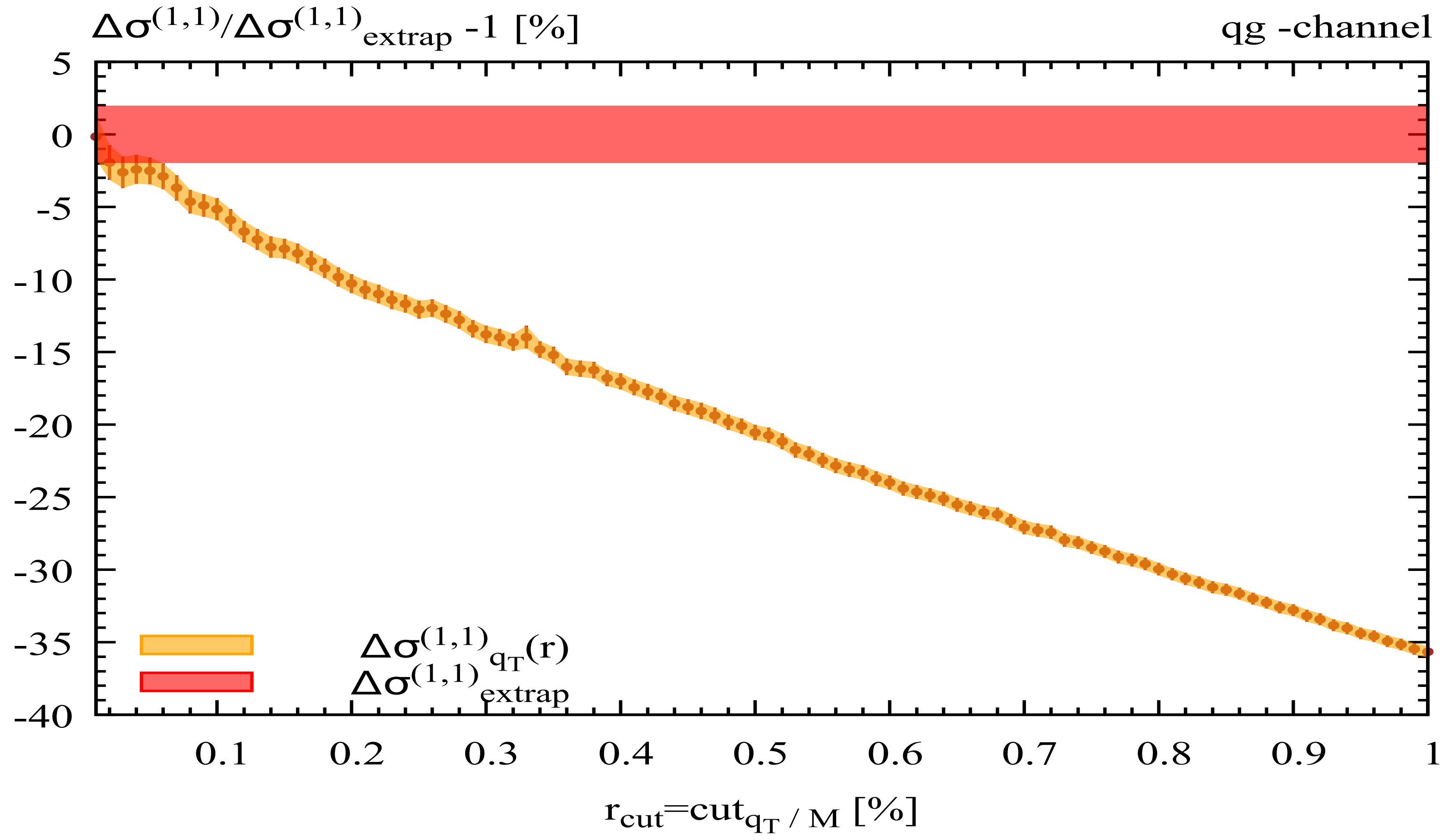
$p_{T,\ell^\pm} > 25$ GeV $\quad y_{\ell^\pm} < 2.5 \quad m_{\ell\ell} > 50$ GeV



- **large power corrections in r_{cut} for mixed corrections**
  - explained by overall small size of corrections, and in parts also by cancellation between partonic channels
- **by far less dramatic dependence at level of cross sections**
  - better than permille precision at inclusive level

Splitting into partonic channels





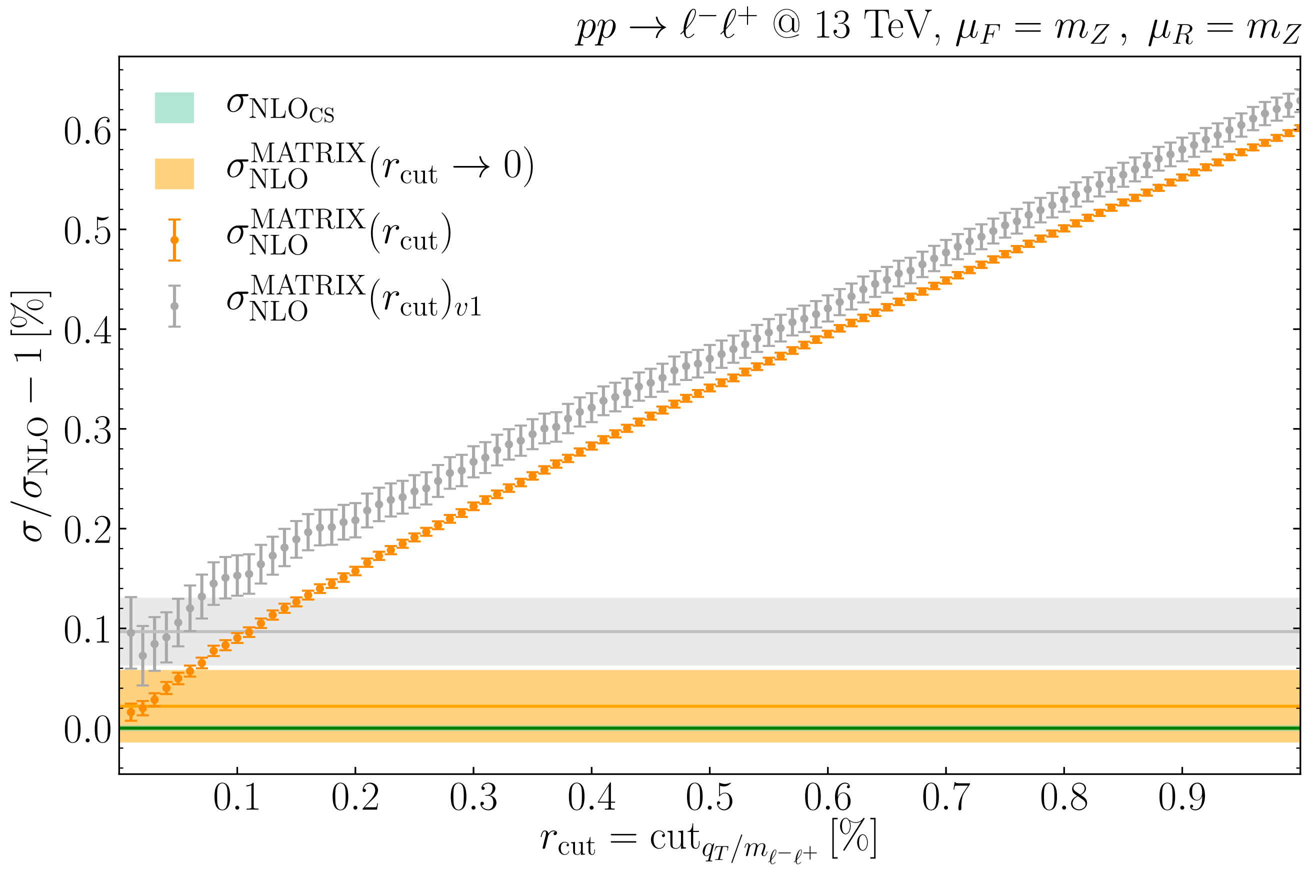
The q_T-subtraction and the residual cut-off dependency in different acceptance setups

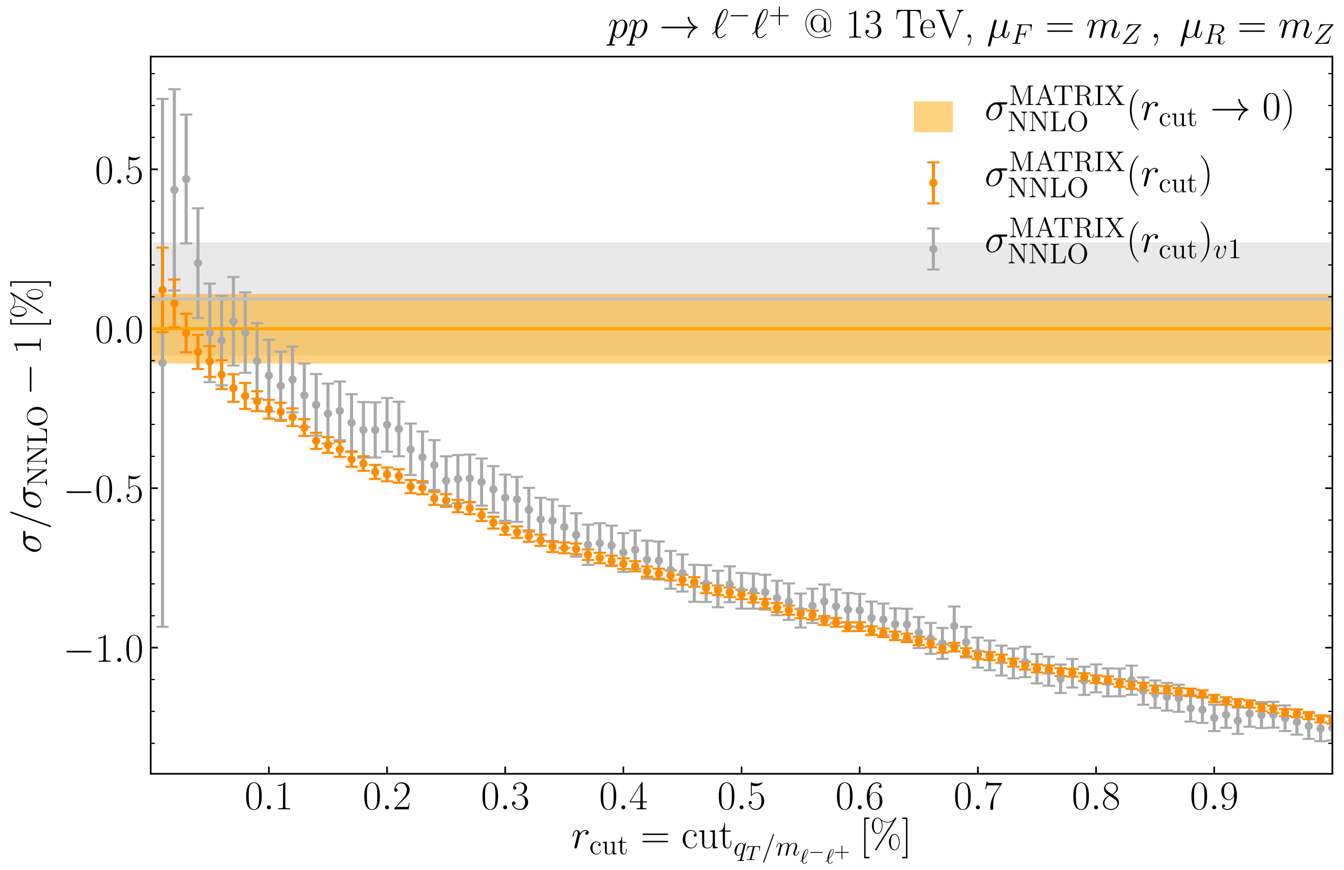
courtesy of S.Kallweit

(cfr. Buonocore, Kallweit, Rottoli, Wiesemann, 2111.13661)

Symmetric cuts

- $p_{\mathrm{T},\ell^\pm} > 25$ GeV

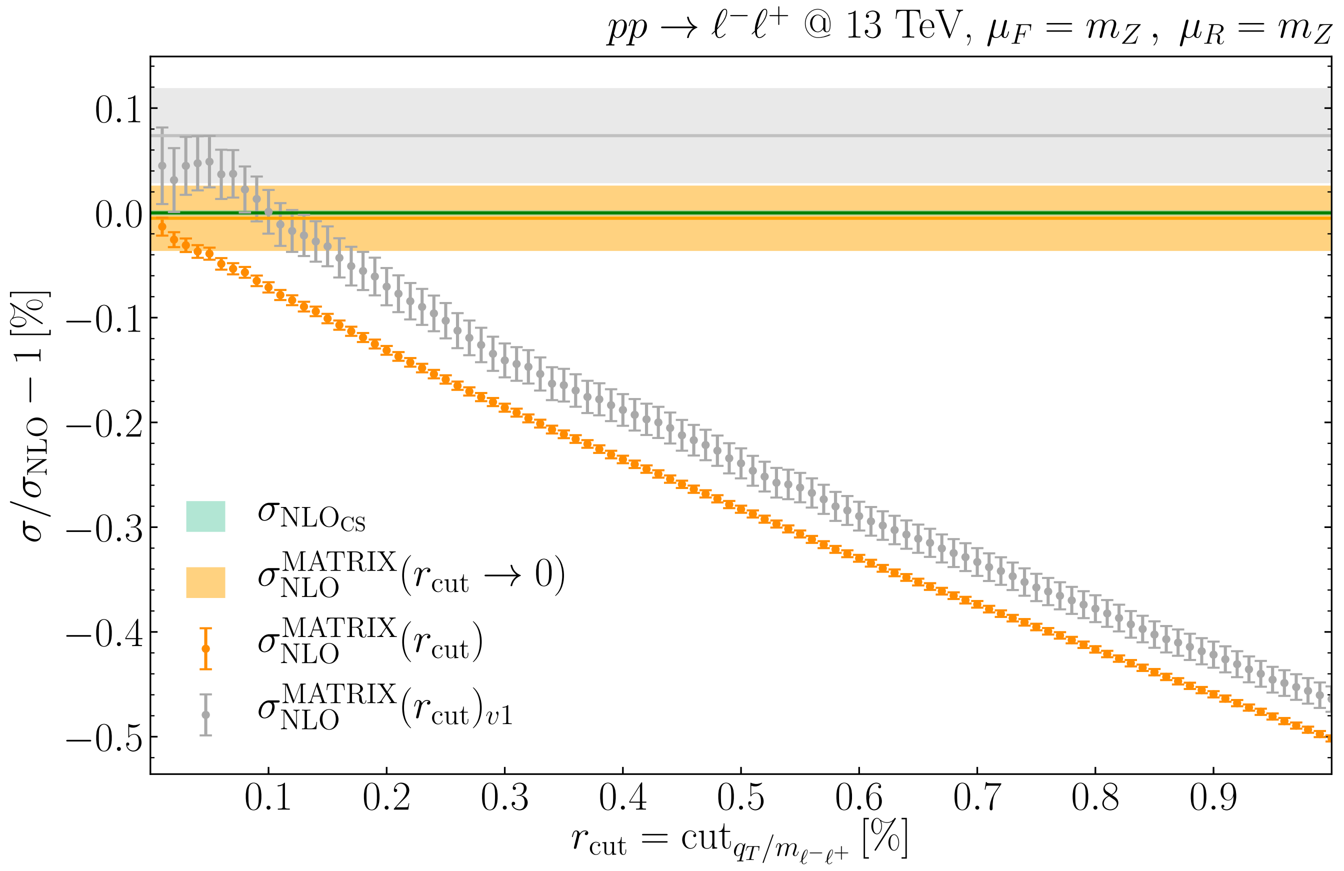


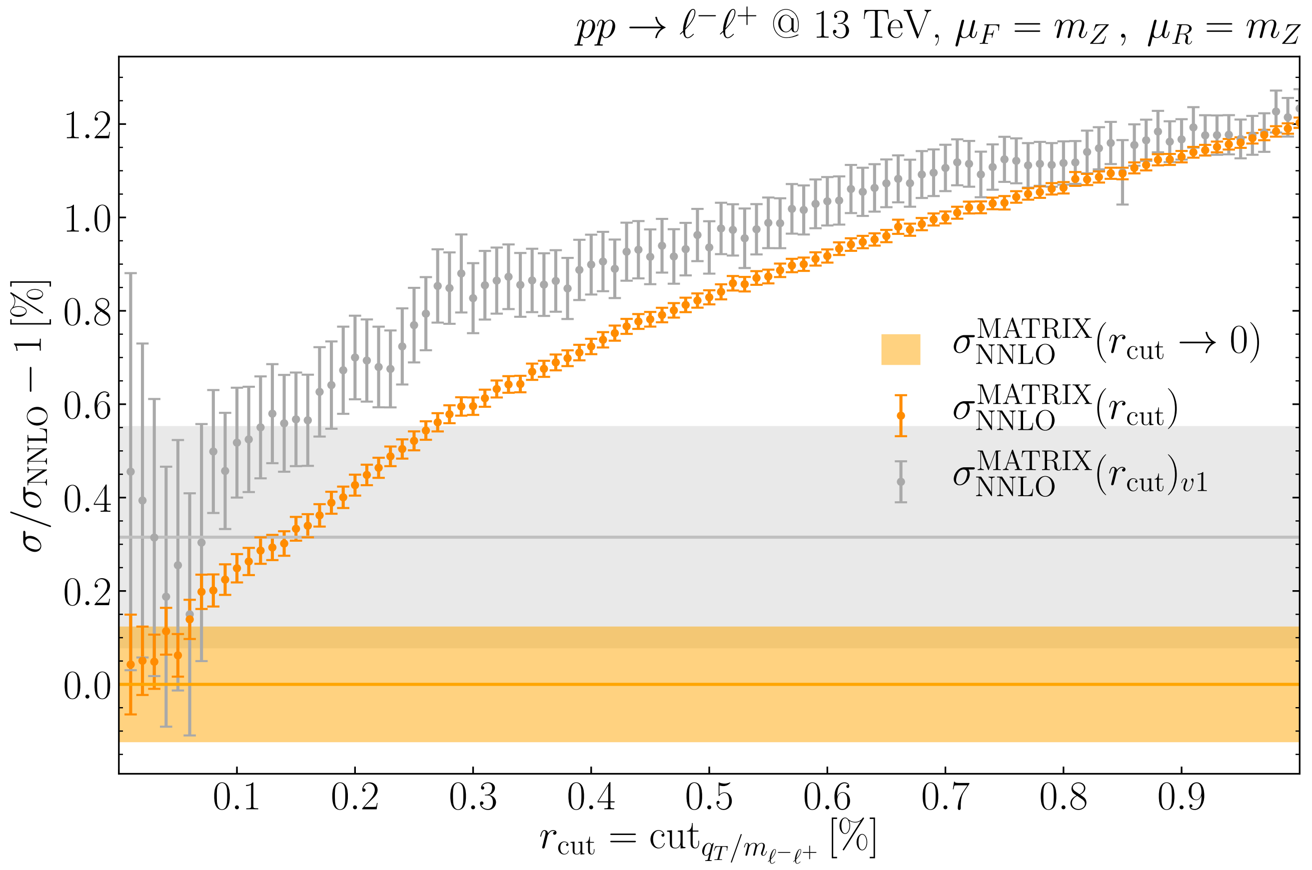


➥ large power corrections in r_{cut}

Asymmetric cuts on ℓ_1 and ℓ_2

- $p_{\mathrm{T},\ell_1} > 25$ GeV $p_{\mathrm{T},\ell_2} > 20$ GeV

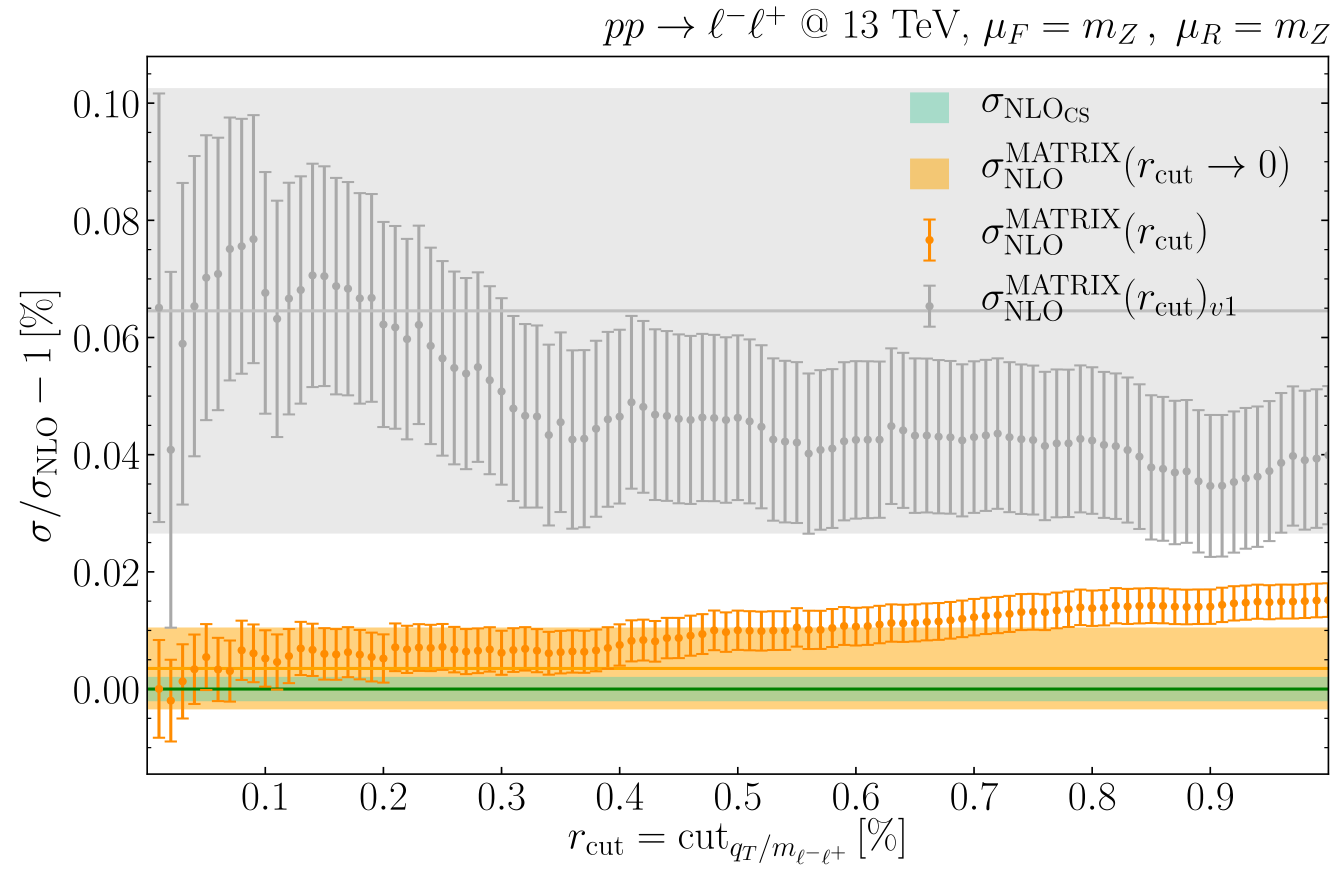


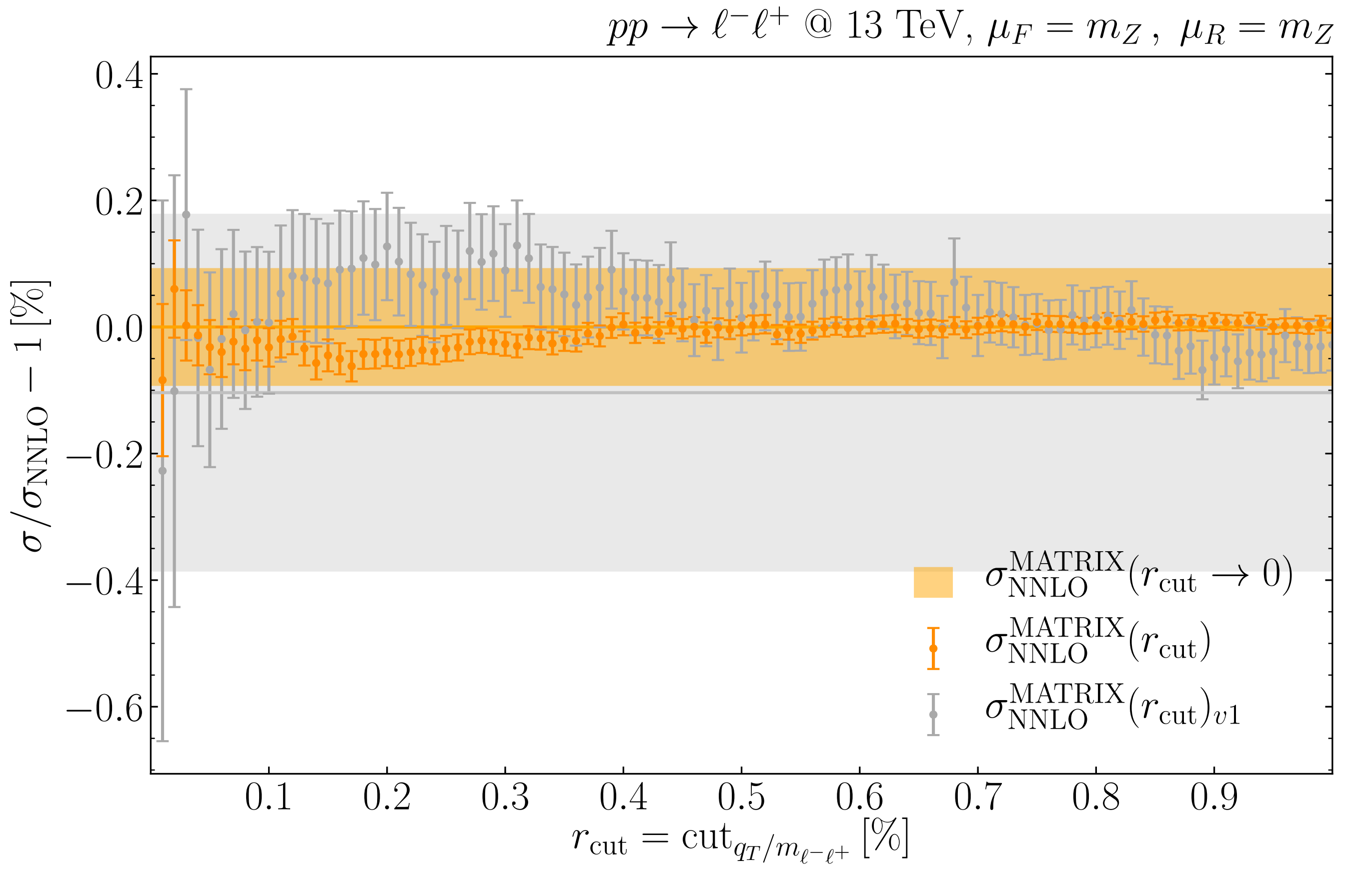


➥ large power corrections in r_{cut}

Asymmetric cuts on ℓ^+ and ℓ^-

- $p_{\mathrm{T},\ell^+} > 25$ GeV $p_{\mathrm{T},\ell^-} > 20$ GeV



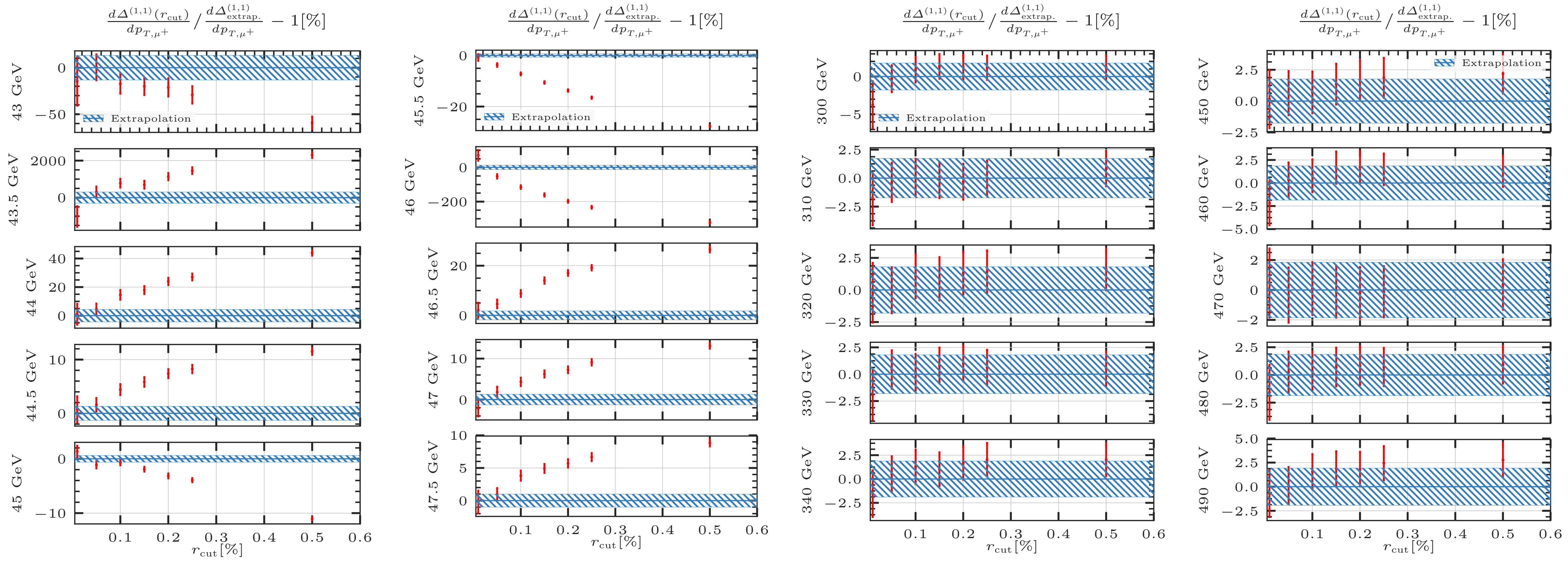


➥ no significant dependence on r_{cut}

Differential sensitivity to r_{cut}

Binwise r_{cut} dependence of the mixed NNLO QCD–EW corrections for NC Drell–Yan

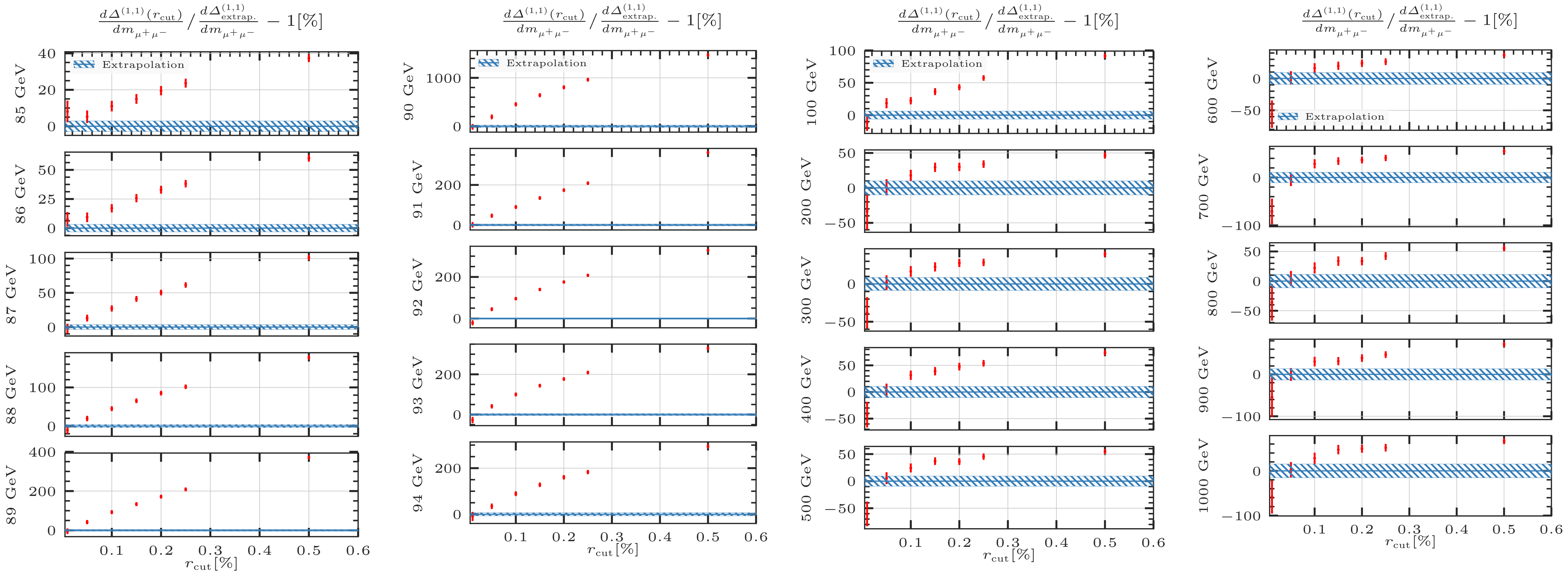
Differential distribution in p_{T,μ^+}: peak (left panels) and tail (right panels) regions



➥ large r_{cut} dependence in particular around the peak of the distribution, and typically precision of $\lesssim 3\%$ on the relative mixed QCD–EW corrections (artificially large where corrections are basically zero)

Binwise r_{cut} dependence of the mixed NNLO QCD–EW corrections for NC Drell–Yan

Differential distribution in $m_{\mu^+\mu^-}$: peak (left panels) and tail (right panels) regions



➥ quite large r_{cut} dependence throughout, and lower numerical precision of $\lesssim 10\%$ on the relative mixed QCD–EW corrections (but still permille-level precision at the level of cross sections)