

# Massless Preheating and Electroweak Vacuum Metastability

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[arXiv:2105.06939]

collaborators on this work:  
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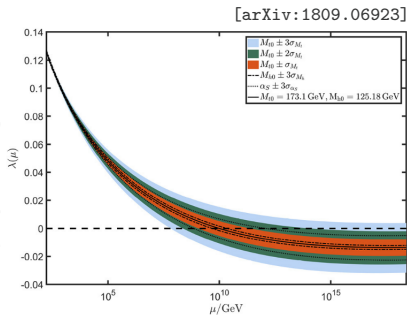
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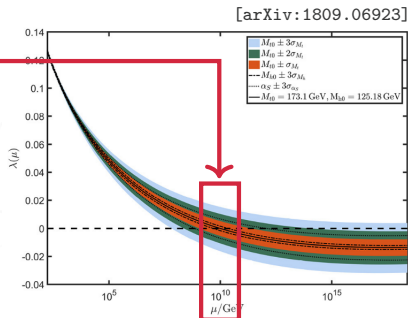
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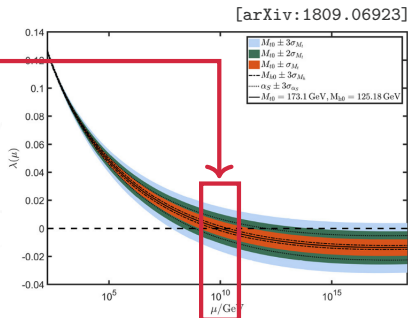


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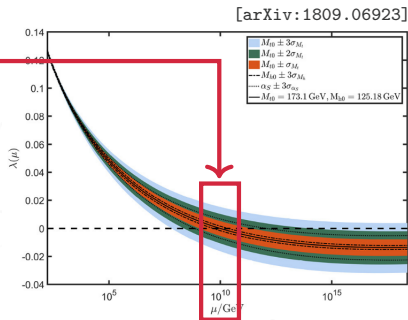
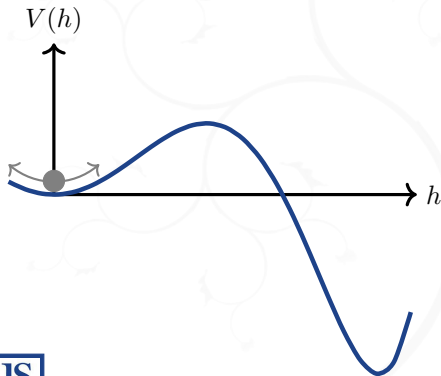


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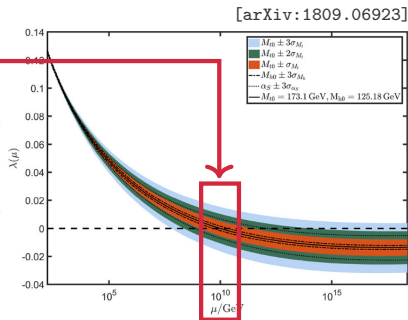
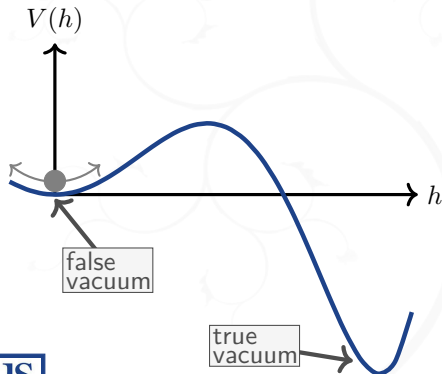


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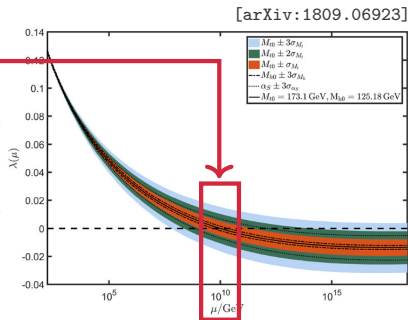
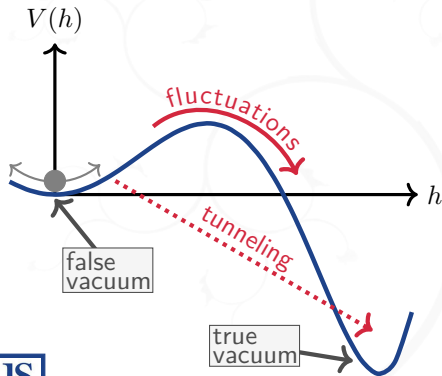


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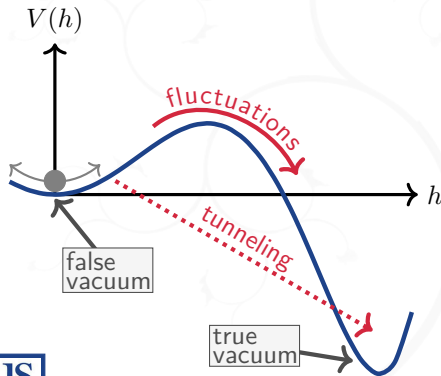


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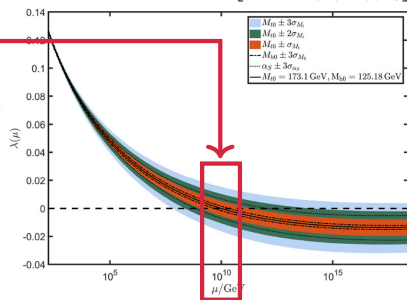
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[arXiv:1809.06923]



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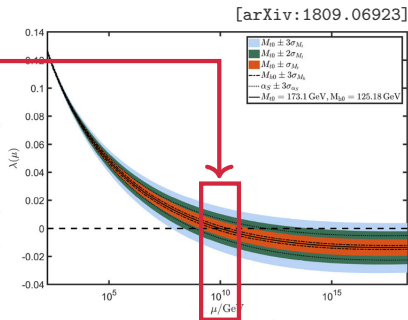
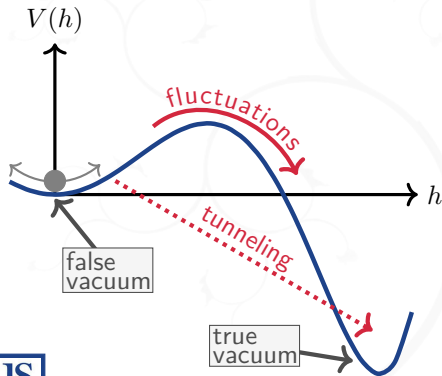


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The fact that false vacuum has *persisted* may provide *window* into *early-universe* dynamics involving the Higgs field.

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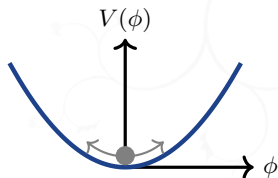
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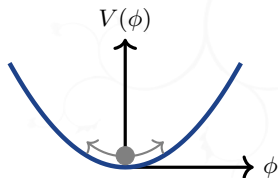
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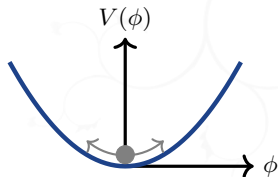
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⇒ leads to **amplification** of fluctuations  $h_k$

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$$\frac{d^2\varphi}{d\eta^2} + \lambda_\phi\varphi^3 = 0$$

for conformal time  $\eta$  and  $\varphi \equiv a\phi$

$$\varphi(x) = \bar{\varphi} \operatorname{cn}\left(x - x_0, \frac{1}{\sqrt{2}}\right),$$

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Similarly, Higgs fluctuations grow *steadily* and **uninterrupted**—appears catastrophic for EW metastability.

on closer inspection,  
is there a **regime of viability**  
for massless preheating?

# Comments On Inflationary Regime

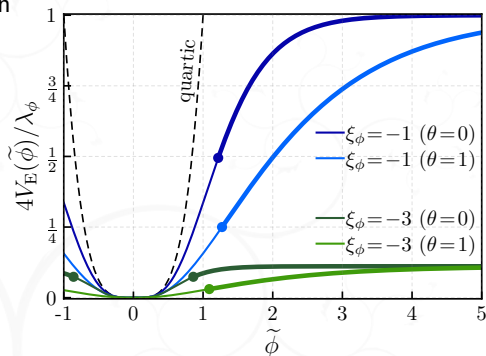
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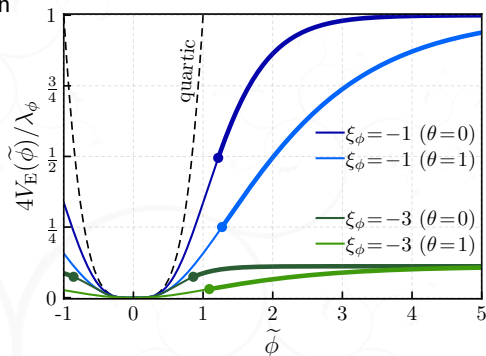


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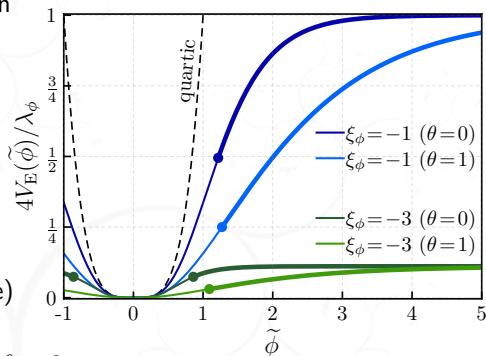
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- Overall, we have the (Einstein-frame) potential in the inflationary regime

$$V_E(\tilde{\phi}) = \frac{\lambda_\phi}{4\xi_\phi^2} \begin{cases} \tanh^4(\sqrt{-\xi_\phi}\tilde{\phi}) & \text{for } \theta = 0 \\ (1 - e^{-\sqrt{\frac{2}{3}}\tilde{\phi}})^2 & \text{for } \theta = 1 \end{cases}$$

where  $\tilde{\phi}$  is the canonical inflaton field.



# Production of Higgs Particles

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- tachyonic production** driven by curvature interactions [relatively short lived since terms dissipate as  $1/a^2$ ]

$$n_h \simeq \frac{1}{8} \left( \sqrt{\lambda_\phi \bar{\varphi}} \right)^3 \left( \frac{3^{9/4} \sqrt{\xi}}{2\pi x^2} \right)^{3/2} \left( \frac{x}{x_0} \right)^4 \sqrt{\frac{2\xi}{3\sqrt{3}}} \quad \text{for } x \lesssim \sqrt{6\xi} \text{ (and } \xi > 0)$$



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[terms do not dissipate due to scale invariance ⇒ unimpeded Higgs production]

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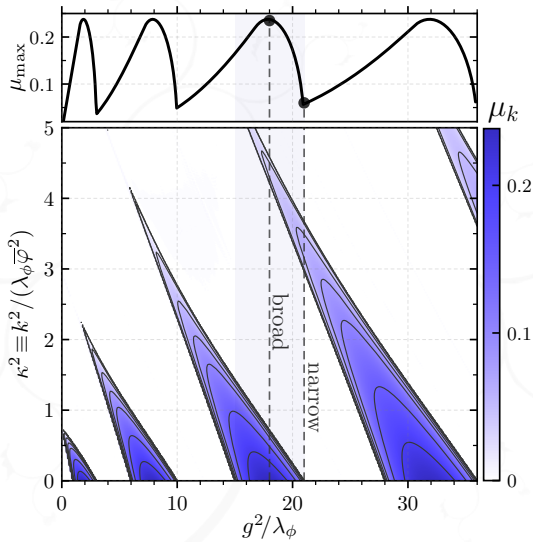
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growth rate  $\mu_{\max}$  has non-trivial dependence on coupling  $g^2/\lambda_\phi$

Higgs modes grow as  $n_{h_k} \propto e^{2\mu_k x}$   
 [for parametric resonance]

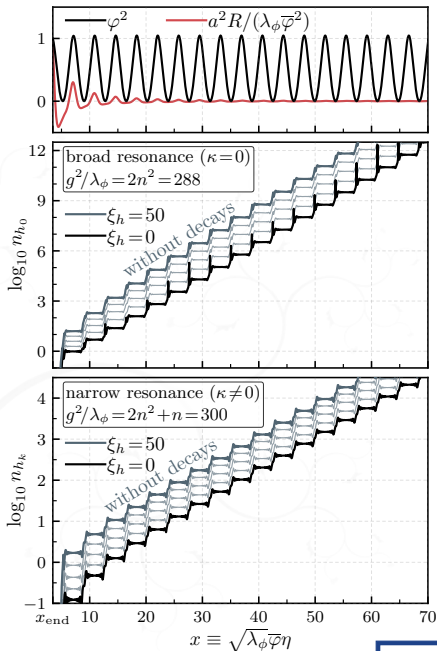


Two important effects neglected thus far:

- ① **perturbative decays** of produced Higgs particles
- ② **backreaction** of particle production on the system

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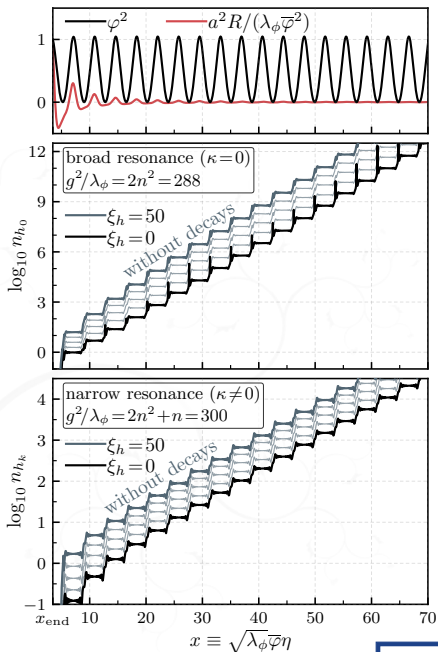


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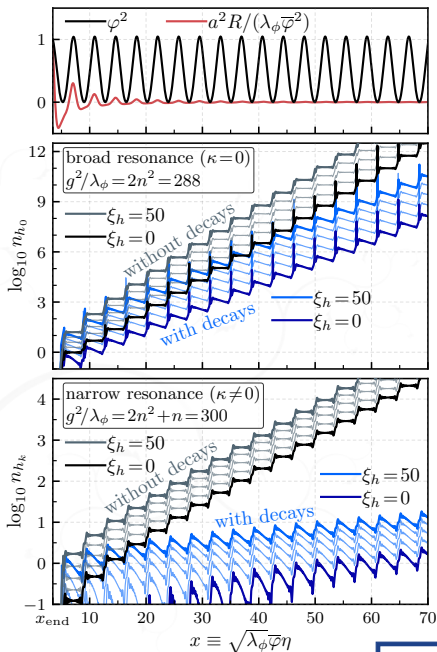
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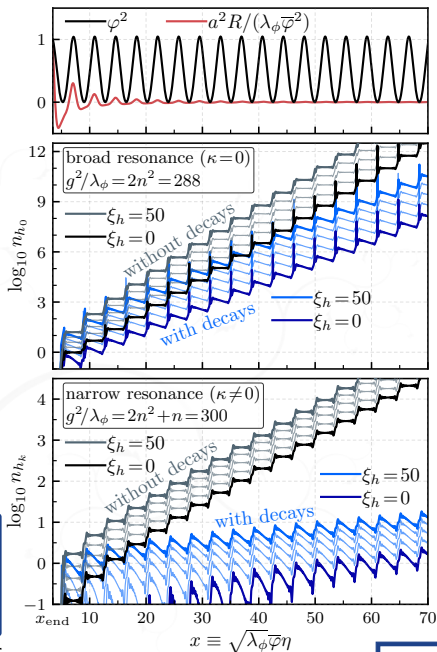
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The decay exponent depends **linearly** on time, same as growth exponents  $2\mu_k x$ .  
 $\Rightarrow$  decays could **entirely suppress production** of Higgs particles.

**does not occur** in massive preheating





## ② Backreaction

- Although fluctuations grow unimpeded, eventually their energy density will be comparable to inflaton background

$$\ddot{\phi} + 3H\dot{\phi} + \lambda_{\phi}\phi^3 + \underbrace{(3\lambda_{\phi}\langle\phi^2\rangle + g^2\langle h^2\rangle)}_{\text{backreaction on inflaton}}\phi = 0$$

variance of fluctuations

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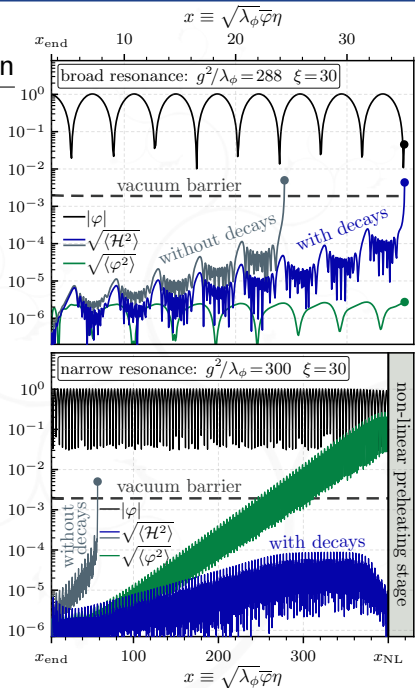
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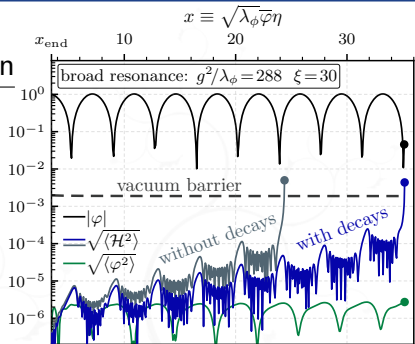
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⇒ tachyonic contribution ( $3\lambda_h\langle h^2\rangle < 0$ ) can **destabilize Higgs** prior to  $x_{\text{NL}}$ .

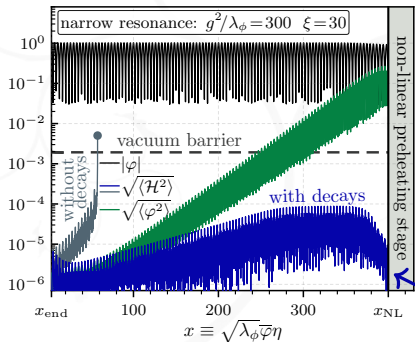
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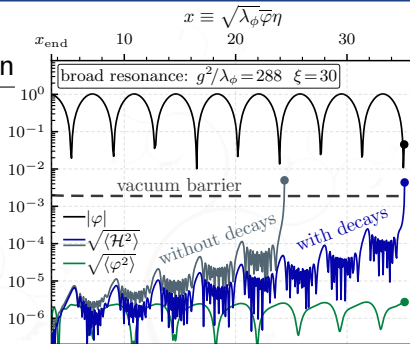
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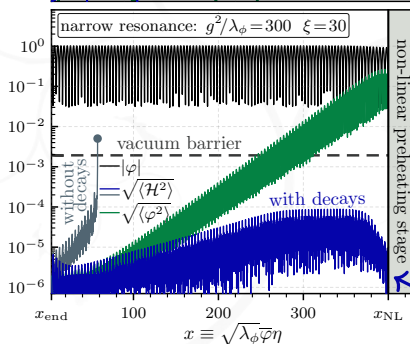
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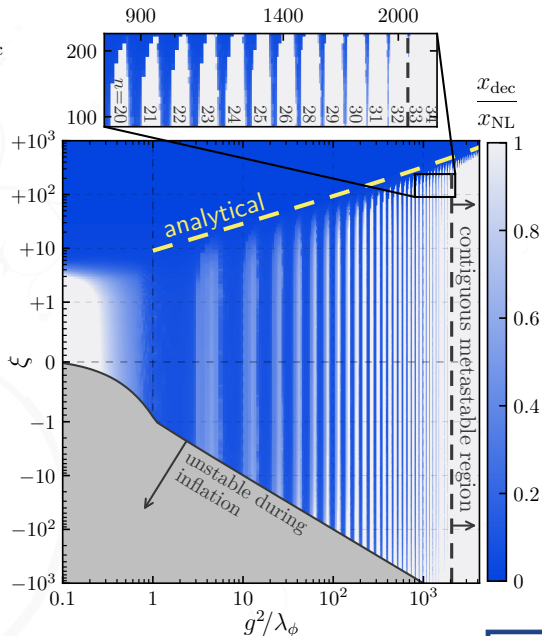


inflaton fluctuations terminate linear stage of preheating

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# Summary of Results

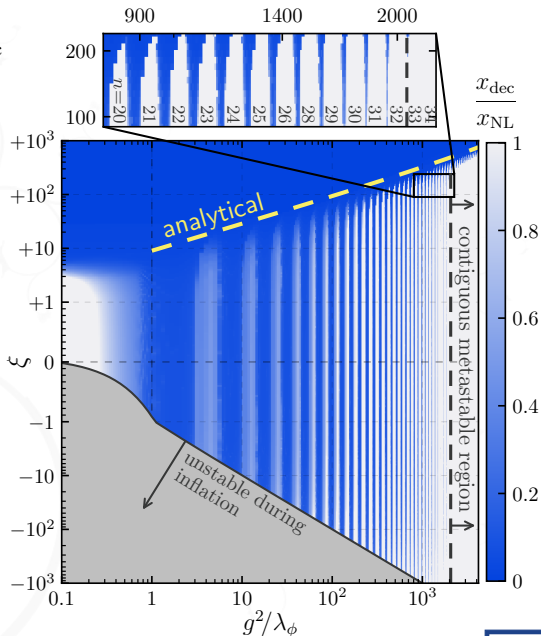
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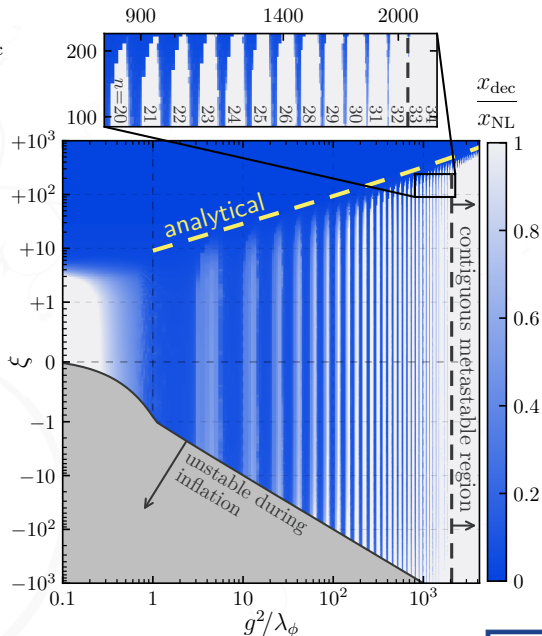


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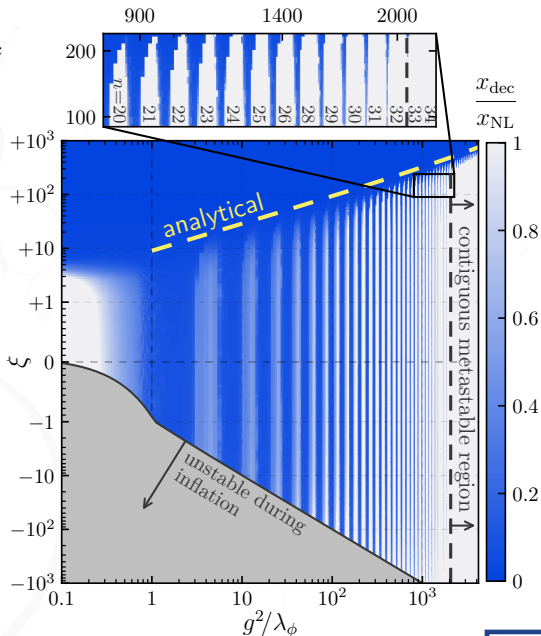


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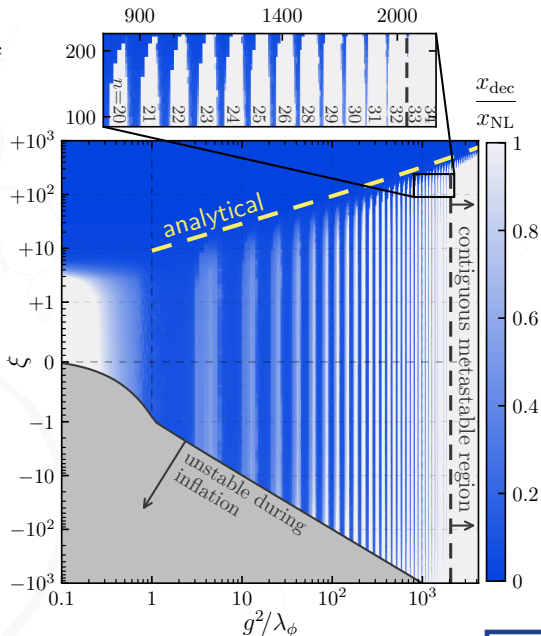


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- curvature coupling imposes envelope over metastable regions at  $\xi \lesssim g^2/\lambda_\phi$ —i.e., **large  $\xi$  viable** as long as  $g^2/\lambda_\phi$  is similarly large.



## TAKE-HOME MESSAGE:

- Although models that lead to massless preheating appear catastrophic for electroweak vacuum metastability, fully accounting for backreaction and perturbative decays reveals a **large number of disjoint (meta)stable regions**.
- In contrast to other (massive) preheating scenarios, the Higgs-inflaton coupling is ultimately **bounded from below** to ensure viability.

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## FUTURE WORK/DIRECTIONS:

- Inclusion of terms in inflaton potential, *e.g.*, small mass terms, that break scale invariance—new phases of evolution would be considered.
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**THANK YOU FOR YOUR ATTENTION!**