

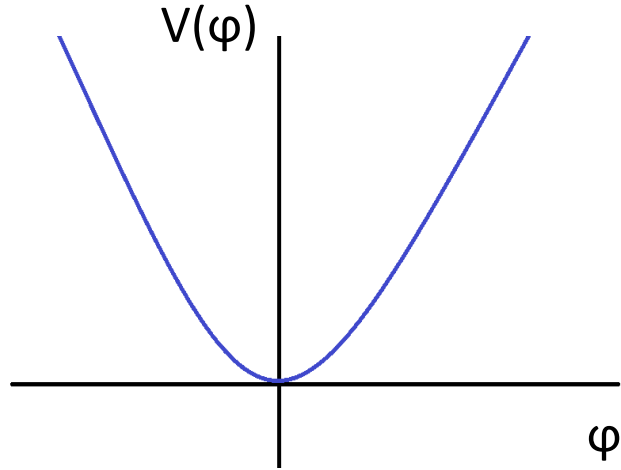
Towards the spontaneous symmetry
breaking:
First order phase transition in QFT

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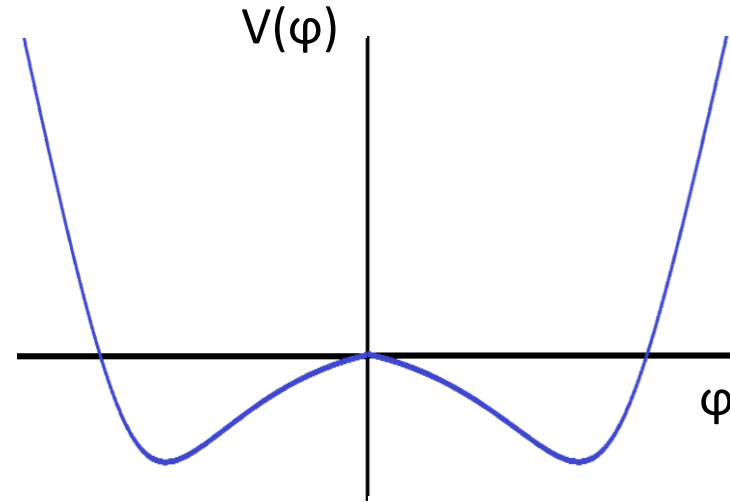
UNCE seminar

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After the Big Bang - high temperature
Only the symmetric vacuum in
unbroken symmetry



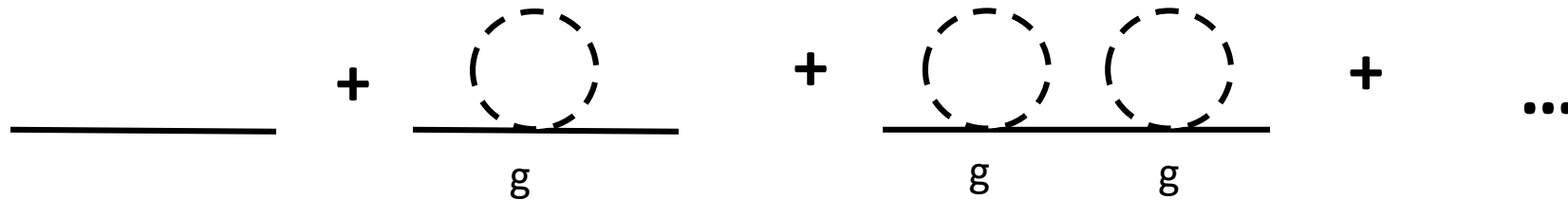
Now - low temperature
 $V(\varphi) \sim m^2\varphi^2 + \lambda\varphi^4$
The symmetry is broken



The goal: To get from high temperature to low temperature in a meaningful way

From high temperature to low temperature

- Sum up the diagrams to get the effective masses
- This effect might be so strong, it generates the asymmetric vacuum even for the high temperature case



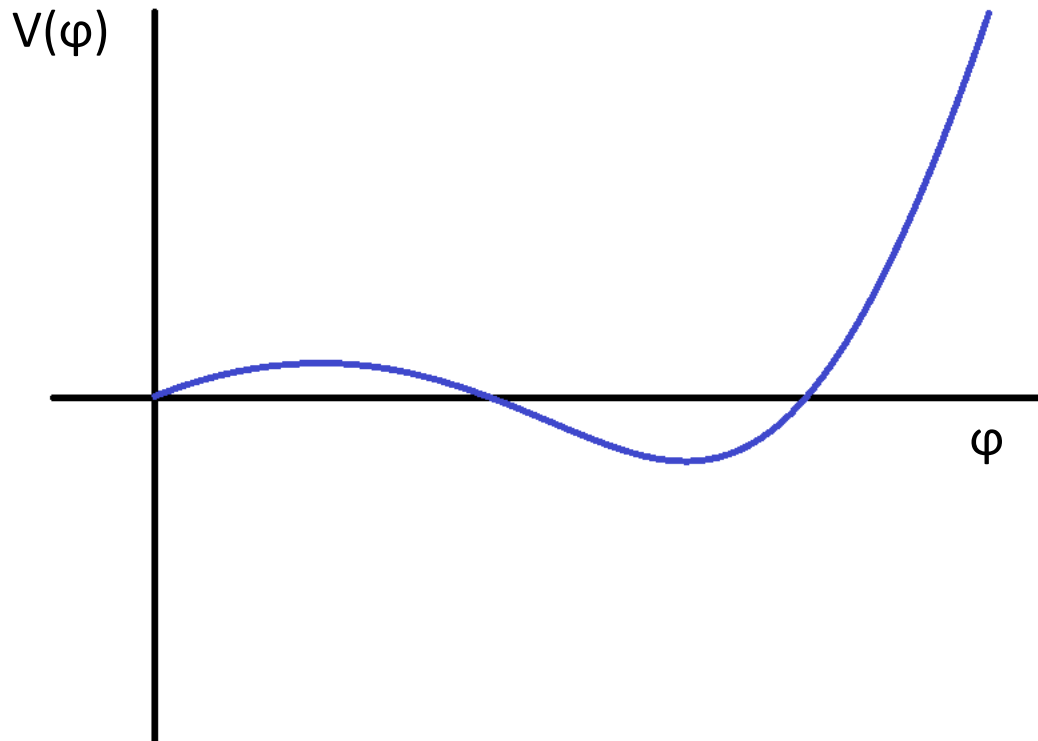
E.g. in theory with one vector boson and one SU(2) singlet with $L \ni g^2 AA\varphi\varphi$, we get in the lowest order (in 3 dim)

$$V_{eff,0}(\varphi) = \frac{\varphi^4 4 \pi \lambda - g^3 \varphi^3 + 8 \pi m^2 \varphi^2}{16 \pi}$$

This $g^3 \varphi^3$ term can even create a barrier!

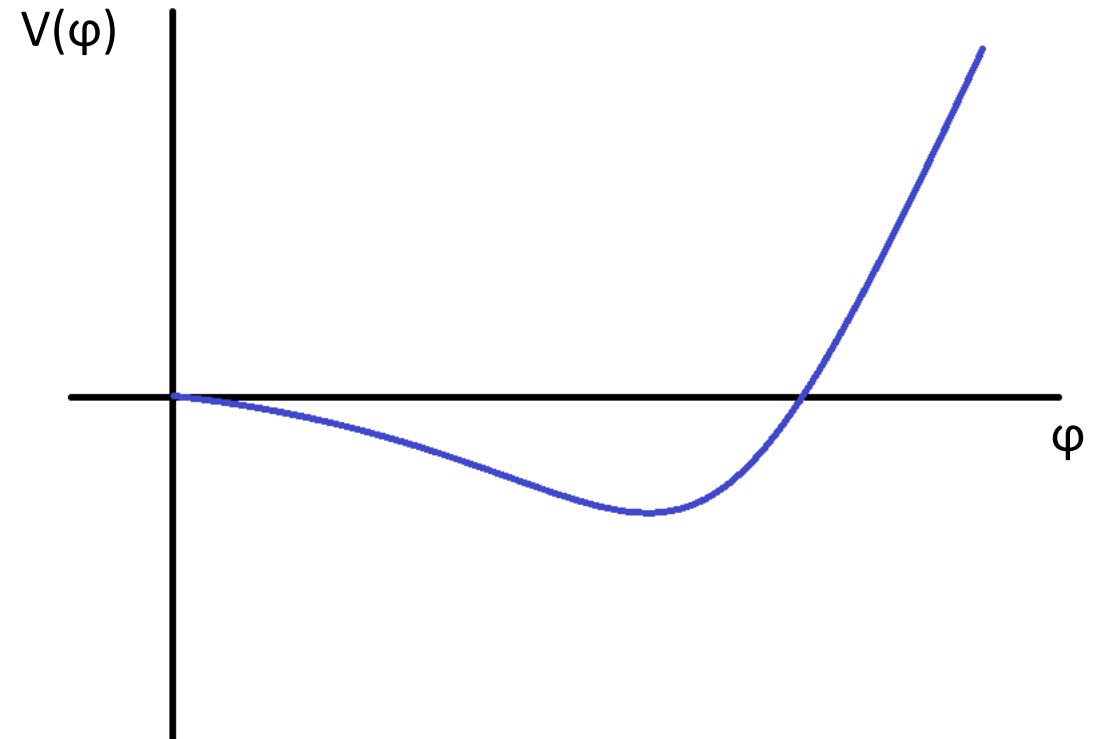
First order phase transition

- Tunneling creates bubbles
- The critical temperature is where the tunnelling is possible for the first time



Second order phase transition

- Smooth process



Why first order phase transitions?

- To generate any matter/antimatter asymmetry, three conditions are to be met
 - Baryon and lepton number violating processes – e.g. if the bubble refracts left-handed and right-handed particles differently (weak sphalerons), this is already in the SM
 - CP violation, not enough in the SM
 - Loss of thermal equilibrium – any asymmetry is washed out if there is equilibrium, there is no way we can lose thermal equilibrium in the SM
- The accelerating bubbles would have left traces in the universe – mainly the gravitational wave background as it is a turbulent event.

Model: Scalar σ , doublet H and boson A, dim 3

$$L \ni \mu_1^2 HH^* + \mu_2^2 \sigma^2 + \lambda_1 (HH^*)^2 + \lambda_2 \sigma^4 + \xi (HH^*) \sigma^2 + g^2 HH^* A^2$$

$$\text{Shift } H = \frac{1}{\sqrt{2}} \begin{pmatrix} \Psi_1 + i \Psi_2 \\ \varphi_1 + \Psi_3 + i \Psi_4 \end{pmatrix}, \quad \sigma = \varphi + \varphi_2$$

With background fields φ_1, φ_2

Some effective masses after diagonalization to the first order:

$$M_1^2 = \mu_2^2/2 + 1/2 \lambda_1 \varphi_1^2 + 1/2 \xi \varphi_2^2$$

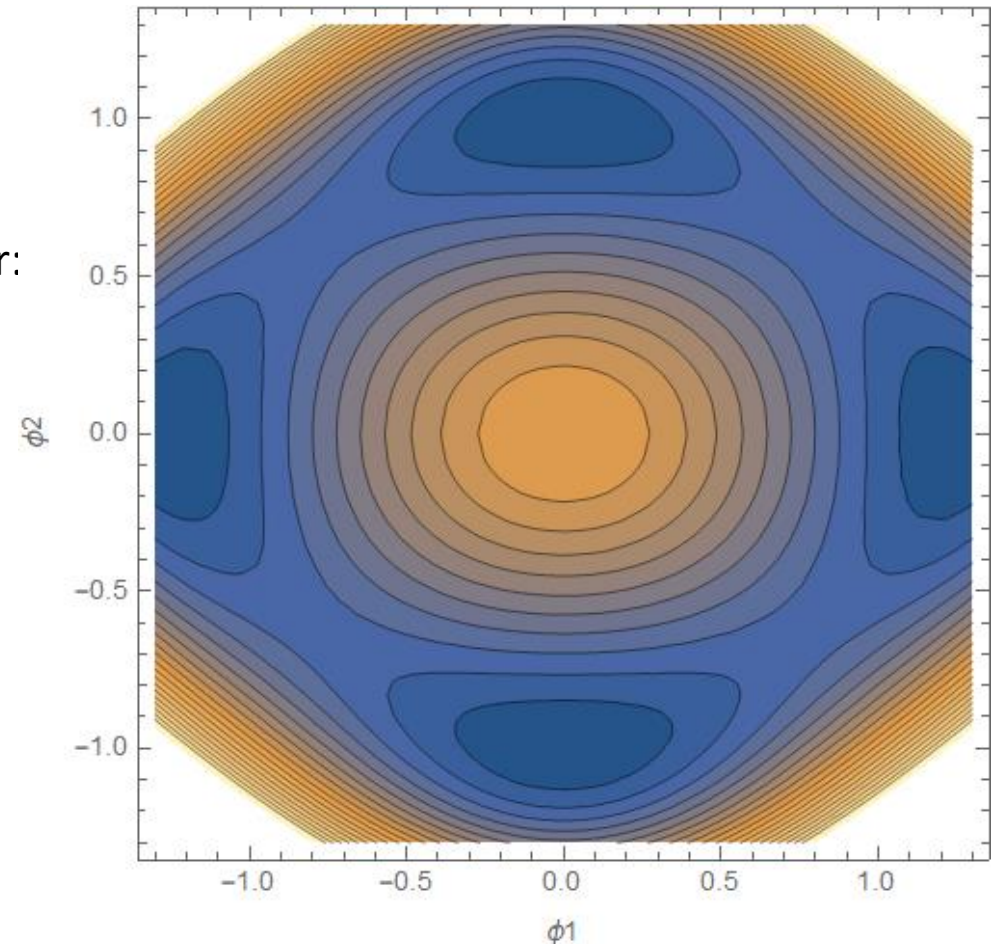
$$M_\varphi^2 = \mu_2^2 + 1/2 \xi \varphi_1^2 + 6 \lambda_2 \varphi_2^2$$

$$M_A^2 = 3/2 g^2 \varphi_1^2$$

The potential:

$$V_0(\varphi_1, \varphi_2) = \frac{\mu_1^2 \varphi_1^2}{2} + \frac{\lambda_1 \varphi_1^4}{4} + \mu_2^2 \varphi_2^2 + \frac{1}{2} \xi \varphi_1^2 \varphi_2^2 + \lambda_2 \varphi_2^4$$

$$V_1(\varphi_1, \varphi_2) = -\frac{1}{12\pi} \sum \text{eigenvalues}(M^2)$$



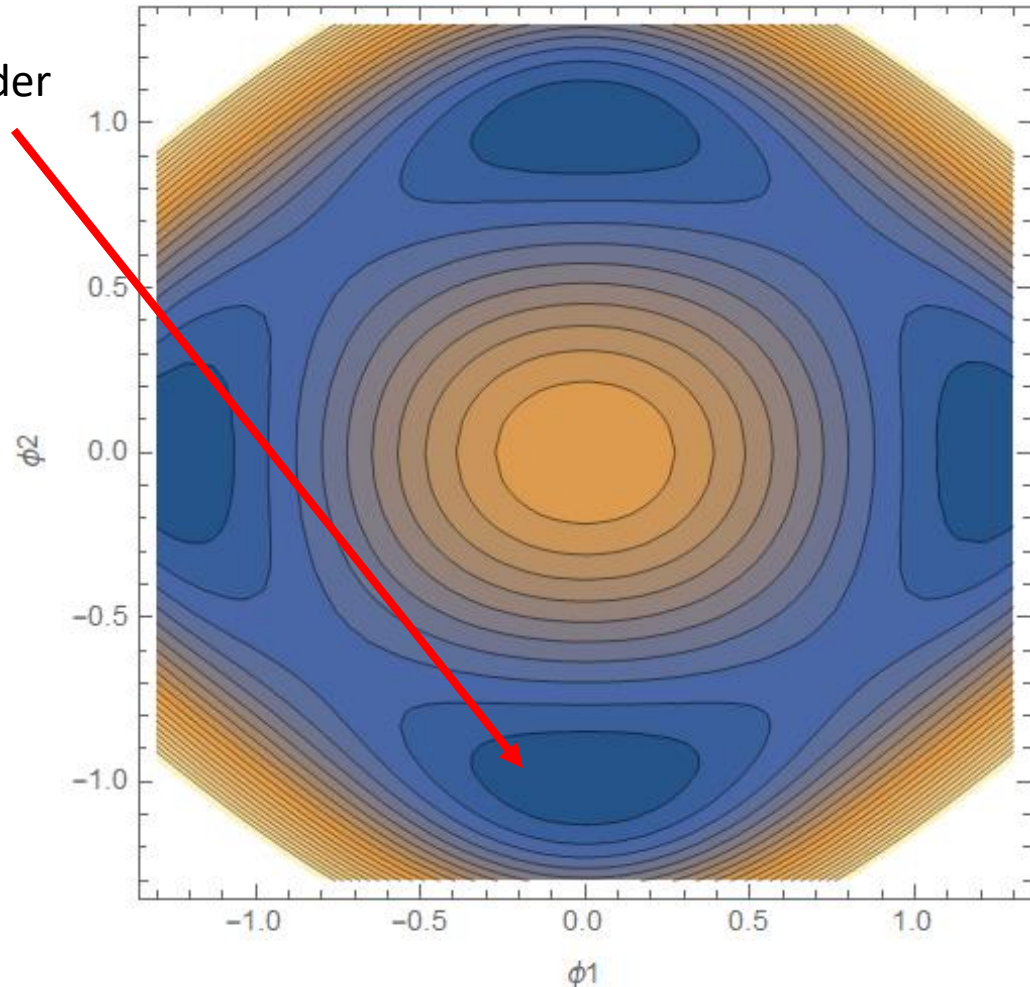
Model: Scalar σ , doublet H and boson A, dim 3

Some VEVs at the critical temperature at the lowest order

$$\langle H \rangle = \frac{\mu_1^2}{2\sqrt{\lambda_1}}$$

$$\langle H^2 \rangle = -\frac{\mu_1^4}{4\lambda_1}$$

$$\langle A \rangle = 0$$



Some VEVs at the critical temperature at the first order

$$\langle H \rangle =$$

$$\frac{8 \sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}} - \frac{\sqrt{2} g^3 \sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}}}{\lambda_1^{3/2}} + \sqrt{2} \sqrt{\mu_{22} - \frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}} + \frac{\varepsilon \sqrt{\mu_{22} \left(4 - \frac{2\varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}\right)}}{\lambda_1}}{32 \pi} - \frac{1}{1536 \pi^2 \sqrt{\lambda_2}} \left(-\frac{4}{\sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}}} + \frac{g^3}{\sqrt{2} \lambda_1^{3/2} \sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}}} - \frac{1}{\sqrt{2} \sqrt{\mu_{22} - \frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}}} - \frac{\varepsilon^2}{\lambda_1^2 \sqrt{\mu_{22} \left(4 - \frac{2\varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}\right)}} \right) \left(2 \sqrt{2} \sqrt{\lambda_2} \left(g^2 \sqrt{-\frac{g^2 \mu_{22}}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} + \varepsilon \sqrt{-\frac{\mu_{22} \varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} \right) - \sqrt{\lambda_1} \left(2 \lambda_2 \left(8 \sqrt{2} \sqrt{-\mu_{22}} + \sqrt{2} \sqrt{\mu_{22}} + 3 \sqrt{2} \sqrt{-\frac{\mu_{22} \varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} - \sqrt{\mu_{22} \left(2 - \frac{\varepsilon}{\lambda_2}\right)} \right) + \varepsilon \left(3 \sqrt{-\frac{\mu_{22} \varepsilon}{\lambda_2}} + \sqrt{\mu_{22} \left(2 - \frac{\varepsilon}{\lambda_2}\right)} \right) \right) \right)$$

$$\langle H^2 \rangle =$$

$$\frac{\sqrt{2} g^3 \left(-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}\right)^{3/2} + \frac{\lambda_1 \mu_{22} \left(12 \lambda_1 \sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}} + \varepsilon \sqrt{\mu_{22} \left(4 - \frac{2\varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}\right)}\right)}{\sqrt{\lambda_2}}}{32 \pi \lambda_1^{5/2}} + \frac{1}{1536 \pi^2 \lambda_1^{5/2} \sqrt{\lambda_2}} \left(-\frac{3 g^3 \sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}}}{\sqrt{2}} + \frac{1}{2} \sqrt{\lambda_1} \left(36 \lambda_1 \sqrt{-\frac{\sqrt{\lambda_1} \mu_{22}}{\sqrt{\lambda_2}}} + \frac{(4 \sqrt{\lambda_1} \sqrt{\lambda_2} - 3 \varepsilon) \varepsilon \sqrt{\mu_{22} \left(4 - \frac{2\varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}\right)}}{2 \sqrt{\lambda_1} \sqrt{\lambda_2} - \varepsilon} \right) \right) \left(2 \sqrt{2} \sqrt{\lambda_2} \left(g^2 \sqrt{-\frac{g^2 \mu_{22}}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} + \varepsilon \sqrt{-\frac{\mu_{22} \varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} \right) - \sqrt{\lambda_1} \left(2 \lambda_2 \left(8 \sqrt{2} \sqrt{-\mu_{22}} + \sqrt{2} \sqrt{\mu_{22}} + 3 \sqrt{2} \sqrt{-\frac{\mu_{22} \varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} - \sqrt{\mu_{22} \left(2 - \frac{\varepsilon}{\lambda_2}\right)} \right) + \varepsilon \left(3 \sqrt{-\frac{\mu_{22} \varepsilon}{\lambda_2}} + \sqrt{\mu_{22} \left(2 - \frac{\varepsilon}{\lambda_2}\right)} \right) \right) \right)$$

$$\langle A \rangle =$$

$$\frac{g^2 \left(-\frac{\mu_{22}}{\sqrt{\lambda_1} \sqrt{\lambda_2}}\right)^{3/2}}{8 \sqrt{2} \pi} + \frac{\mu_{22} (4 \sqrt{\lambda_1} \sqrt{\lambda_2} - 3 \varepsilon) \left(-2 \sqrt{2} \sqrt{\lambda_2} \left(g^2 \sqrt{-\frac{g^2 \mu_{22}}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} + \varepsilon \sqrt{-\frac{\mu_{22} \varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} \right) + \sqrt{\lambda_1} \left(2 \lambda_2 \left(8 \sqrt{2} \sqrt{-\mu_{22}} + \sqrt{2} \sqrt{\mu_{22}} + 3 \sqrt{2} \sqrt{-\frac{\mu_{22} \varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}} - \sqrt{\mu_{22} \left(2 - \frac{\varepsilon}{\lambda_2}\right)} \right) + \varepsilon \left(3 \sqrt{-\frac{\mu_{22} \varepsilon}{\lambda_2}} + \sqrt{\mu_{22} \left(2 - \frac{\varepsilon}{\lambda_2}\right)} \right) \right) \right)}{1536 \pi^2 \lambda_1^{3/2} \lambda_2 \sqrt{\mu_{22} \left(4 - \frac{2\varepsilon}{\sqrt{\lambda_1} \sqrt{\lambda_2}}\right)}}$$

What next?

- Rewrite the couplings constants from 3D to 4D. This has to be done for every model separately.
- Insert numbers
- Calculate higher order correction - still not clear how to incorporate them