

B anomaly hunting at the LHC: $\tau b + \text{missing}$ search for leptoquarks

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with

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Physics in LHC and Beyond
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名古屋大学

高等研究院



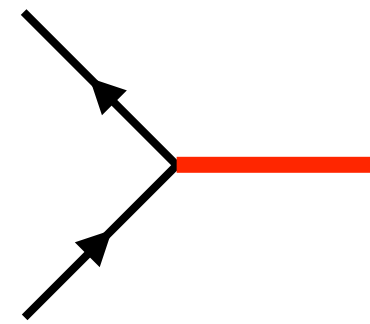
KMI

Nagoya University

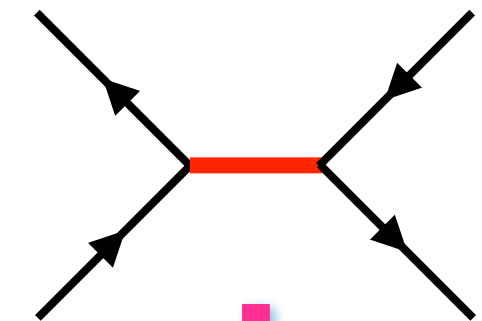
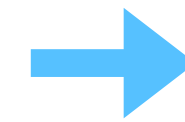


hunt a new interaction

SM particle



New particle

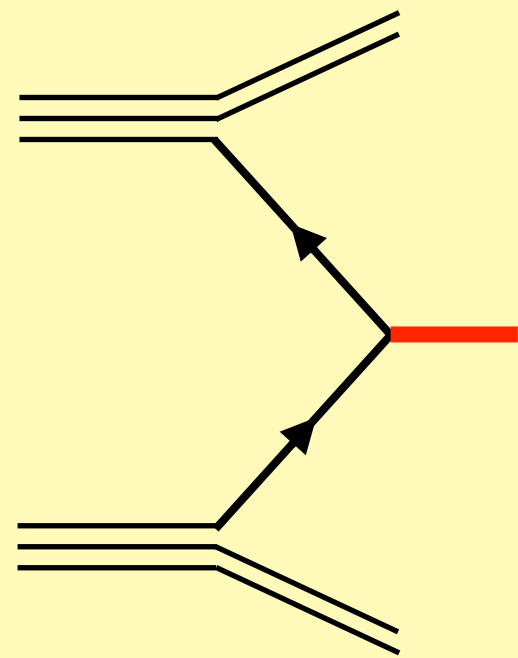


Effective field theory

Collider physics

High-energy frontier

$$\sqrt{s} > m_{NP}$$



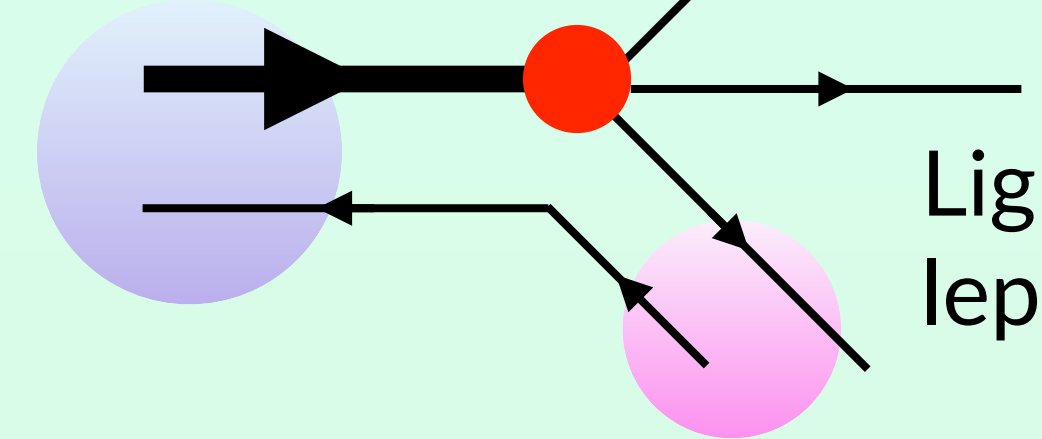
complementary

Precision measurement

High-intensity frontier

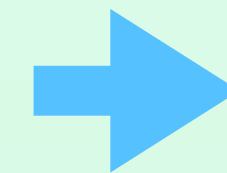


Heavy quark



Light meson, lepton

Small theoretical uncertainty



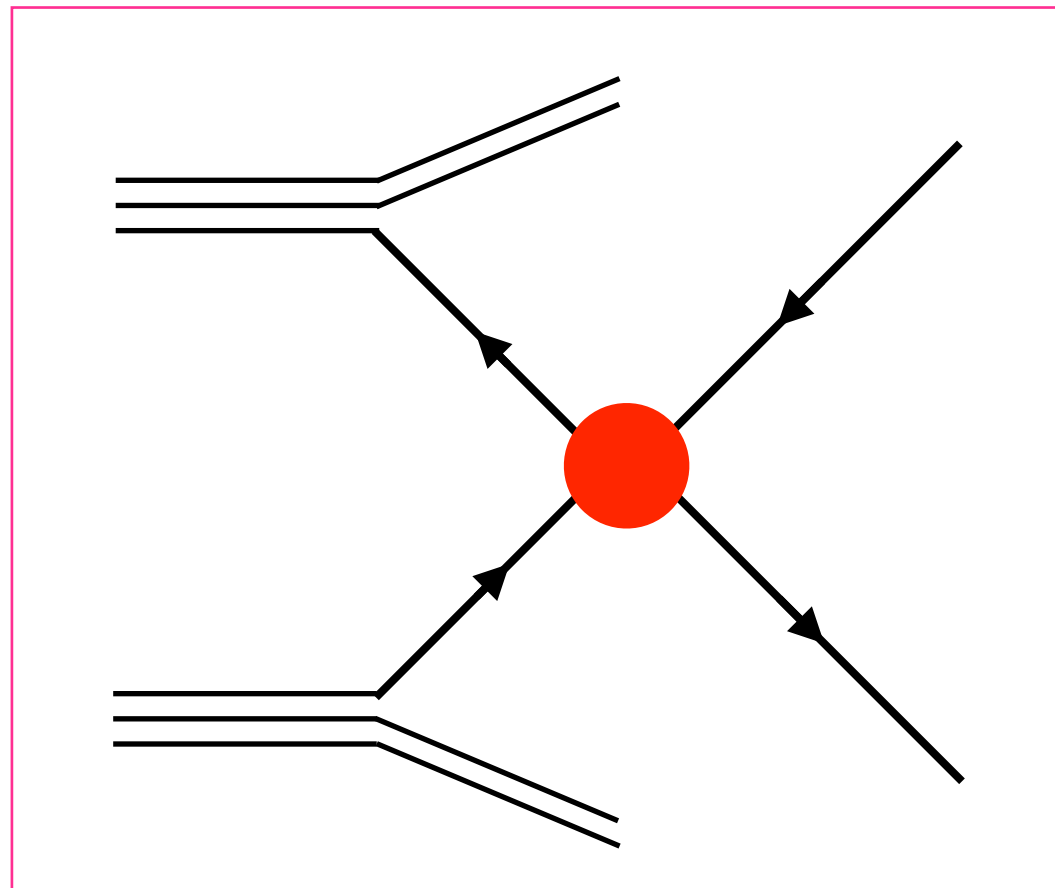
Good new physics sensitivity

no new resonance

non-resonance search
→ This talk

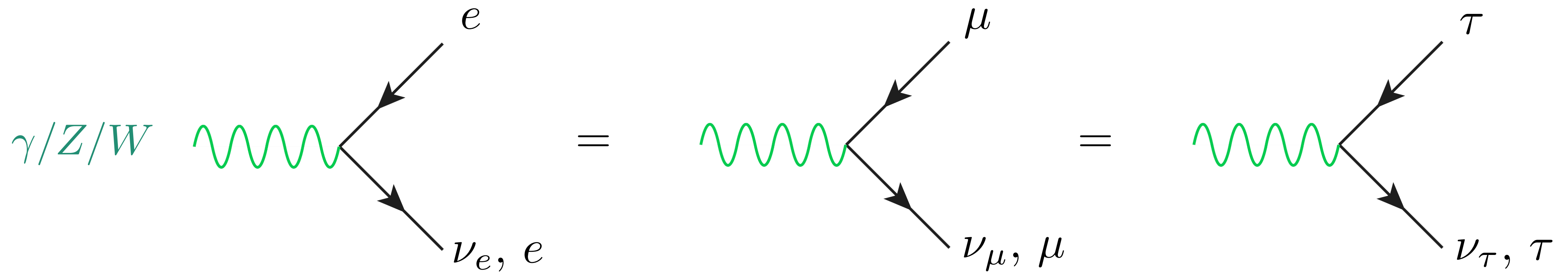
B anomaly!

jet flavor-tagging is very important



Test of Lepton Flavor Universality (LFU)

- ◆ Gauge symmetry predicts lepton flavor universal (LFU) phenomena:



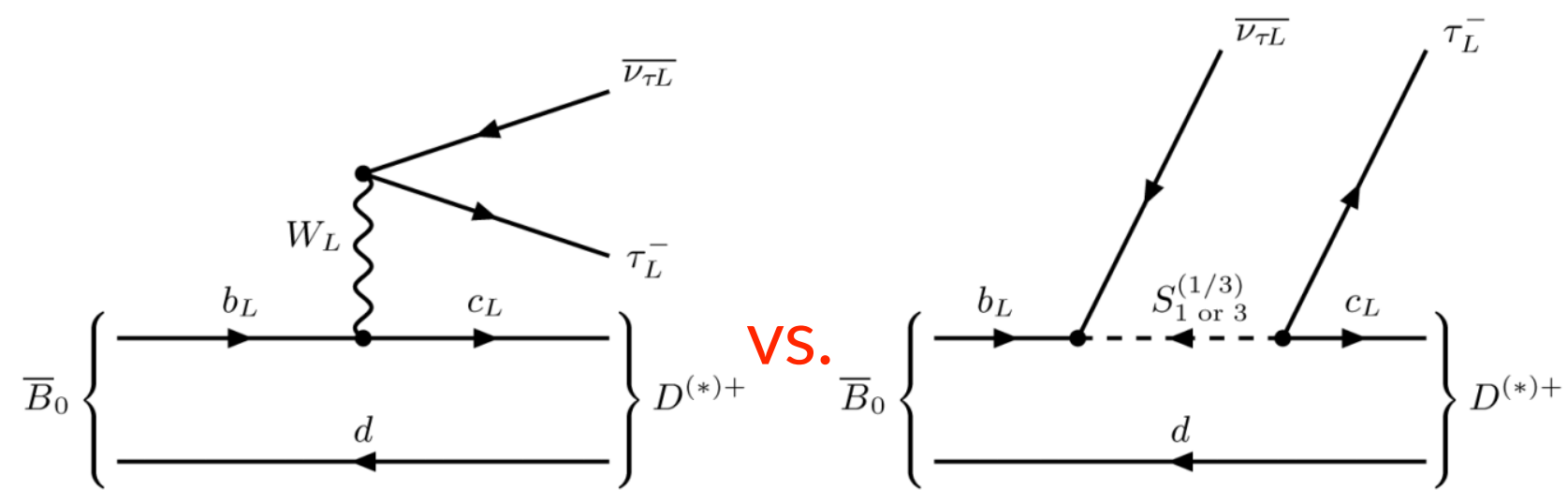
- ◆ Only charged-lepton mass violates the LFU within the SM

$$m_e = 0.5 \text{ MeV}, \quad m_\mu = 105 \text{ MeV}, \quad m_\tau = 1776 \text{ MeV}$$

LFU observables $R(D)$ and $R(D^*)$

[HFLAV 2019, New SM= Iguro, Watanabe, [2004.10208](https://arxiv.org/abs/2004.10208)]

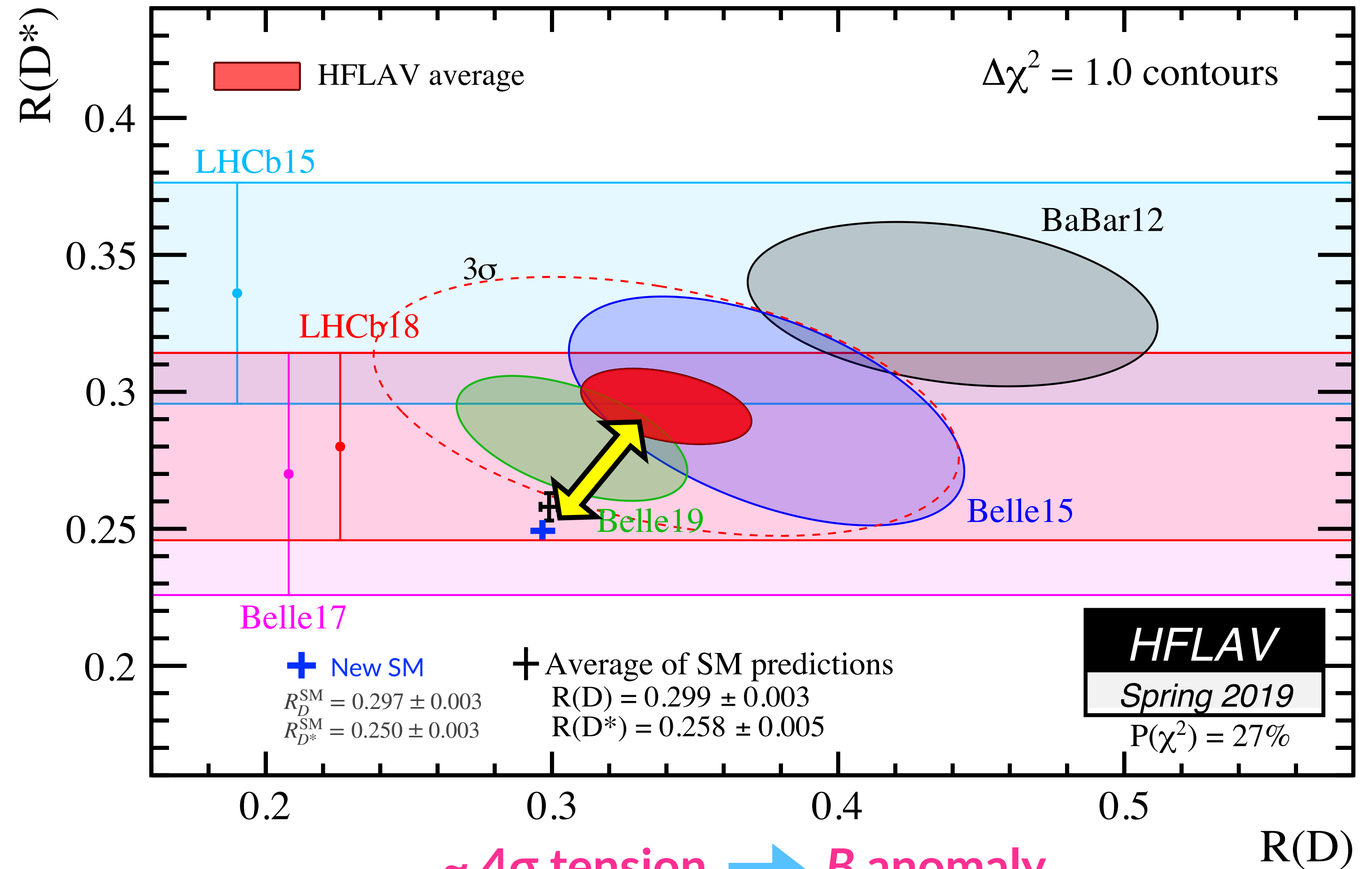
$$R(D^{(*)}) = \frac{\text{BR}(B \rightarrow D^{(*)} \bar{\tau} \nu_{\tau})}{\text{BR}(B \rightarrow D^{(*)} \bar{\ell} \nu_{\ell})} \quad (\ell = e, \mu)$$



SM

New physics, e.g.,
Leptoquark

cf. $\mathcal{B}(B \rightarrow D \ell \nu) = 2\%$, $\mathcal{B}(B \rightarrow D^* \ell \nu) = 5\%$,

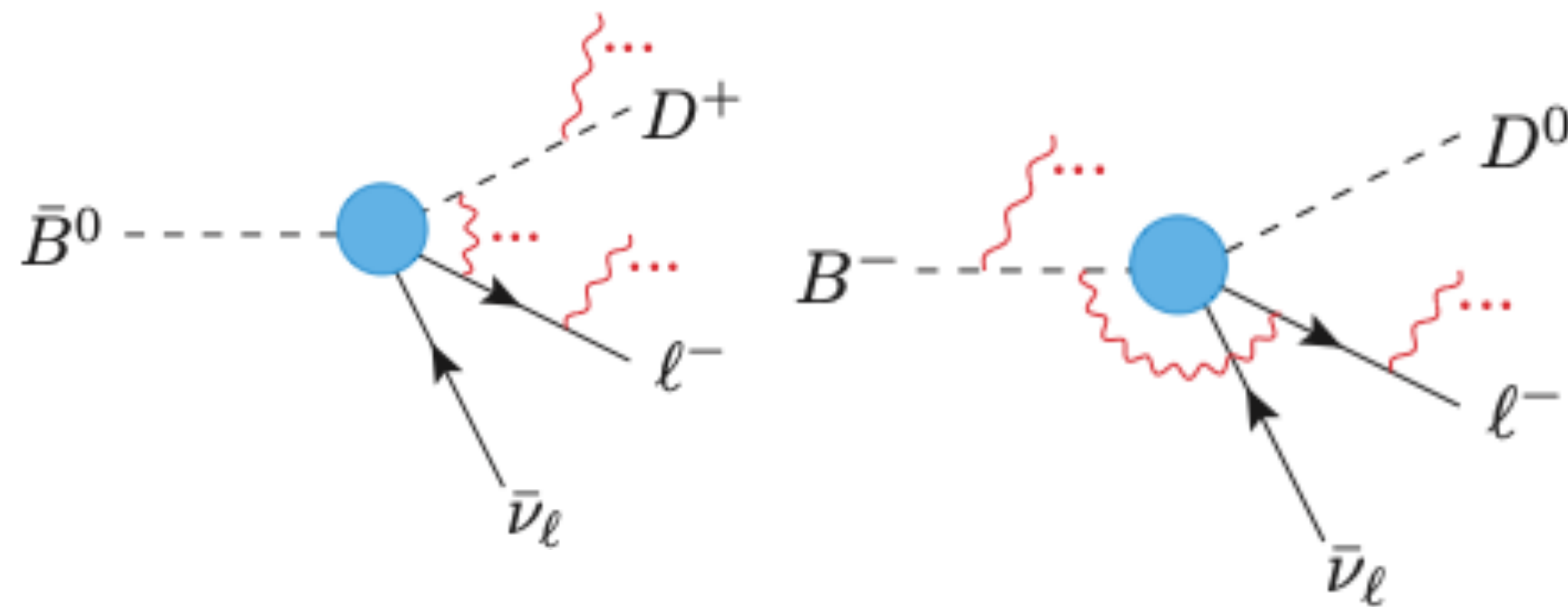


$\sim 4\sigma$ tension \rightarrow B anomaly

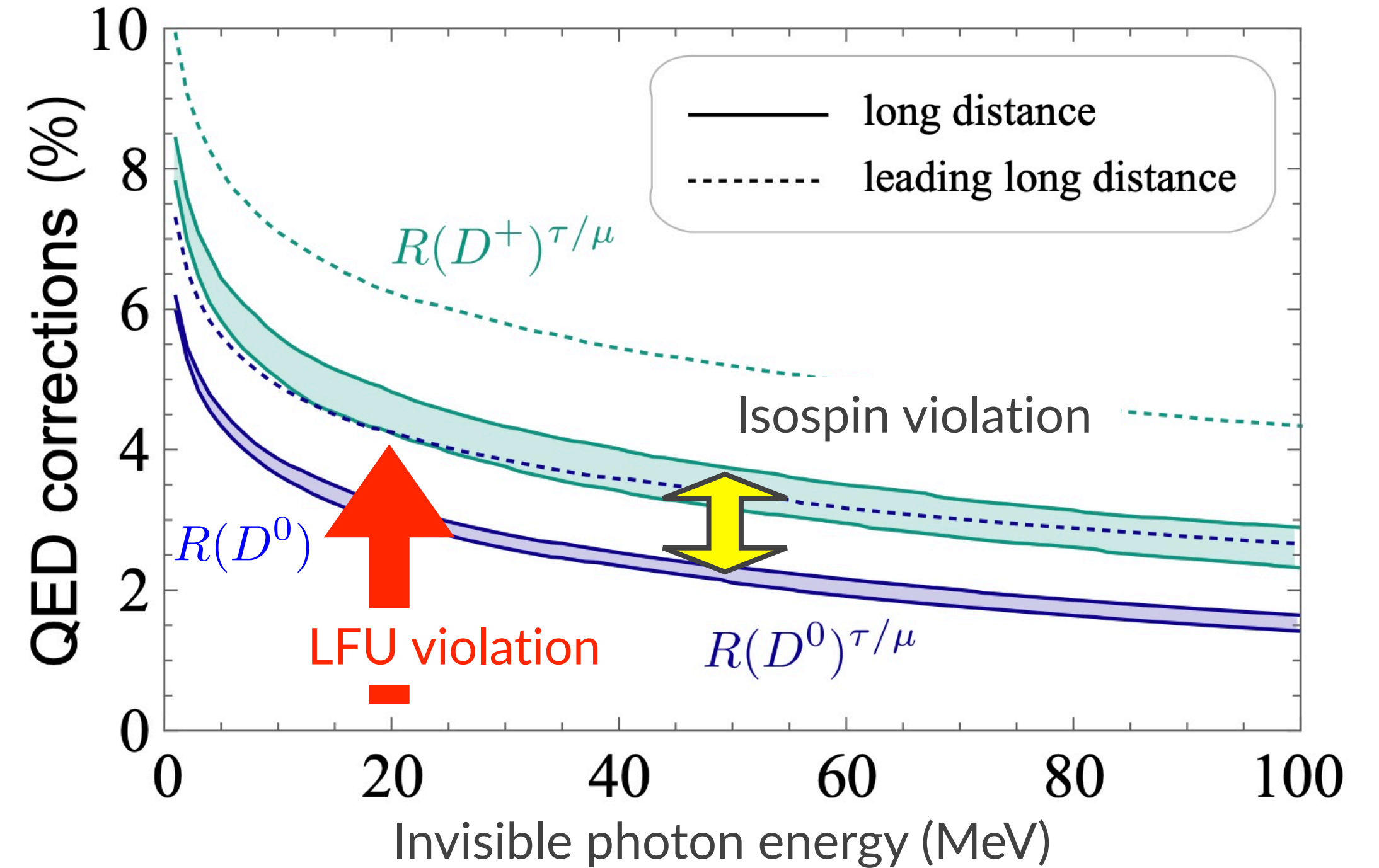
Comment on QED correction within the SM

- ◆ Long-distance QED correction could violate the lepton flavor universality

[de Boer, TK, Nisandzic, [1803.05881](#)] + [Calí, Klaver, Rotondo, Sciascia, [1905.02702](#); Isidori, Nabeebaccus, Zwicky, [2009.00929](#)]



It is found that the QED corrections depend on the lepton velocity (=the lepton masses) with opposite sign (light lepton vs heavy lepton)



[de Boer, TK, Nisandzic, [1803.05881](#)]

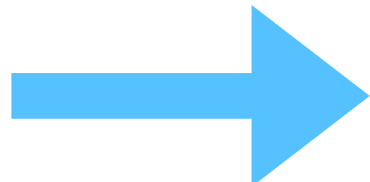
General operator basis

- ◆ General NP for R(D^(*)) anomaly can be model-independently described by

$$\mathcal{H}_{\text{eff}} = 2\sqrt{2}G_F V_{cb} \left[\overset{\text{SM}}{\downarrow} (1 + C_{V_1}) (\bar{c}\gamma^\mu P_L b) (\bar{\tau}\gamma_\mu P_L \nu_\tau) + C_{V_2} (\bar{c}\gamma^\mu P_R b) (\bar{\tau}\gamma_\mu P_L \nu_\tau) \right. \\ \left. + C_{S_1} (\bar{c}P_R b) (\bar{\tau}P_L \nu_\tau) + C_{S_2} (\bar{c}P_L b) (\bar{\tau}P_L \nu_\tau) \right. \\ \left. + C_T (\bar{c}\sigma^{\mu\nu} P_L b) (\bar{\tau}\sigma_{\mu\nu} P_L \nu_\tau) \right] + \text{h.c.},$$

- ◆ C_{V_2} appears dimension-eight at the SMEFT, while others are dimension-six
- ◆ Light right-handed neutrinos ($m_{\nu_R} < m_B$) can also be included, but are constrained from the collider bounds

Single new-particle interpretations

- ◆ W' : C_{V_1}
 - ◆ **Severely constrained** from ΔM_s , $Z' \rightarrow \tau\tau$ search [Faroughy, Greljo, Kamenik, [1609.07138](#)], and $W' \rightarrow \tau\nu$ search [Abdullah, Calle, Dutta, Flores, Restrepo, [1805.01869](#)]
 - ◆ Charged Higgs with generic flavor structure : C_{S_1} , C_{S_2}
 - ◆ **Severely constrained** from $B_c \rightarrow \tau\nu$ and $H^\pm \rightarrow \tau\nu$ search [Iguro, Tobe, [1708.06176](#); Iguro, [2201.06565](#)]
 - ◆ Leptoquark : C_{V_1} , C_{V_2} , C_{S_1} , C_{S_2} , C_T
 - ◆ Collider bound comes from $gg \rightarrow LQ LQ^*$, and broad parameter regions are still allowed
-  We will focus on **heavy LQ scenarios**, then there is no direct bound from LHC

Leptoquark catalogue

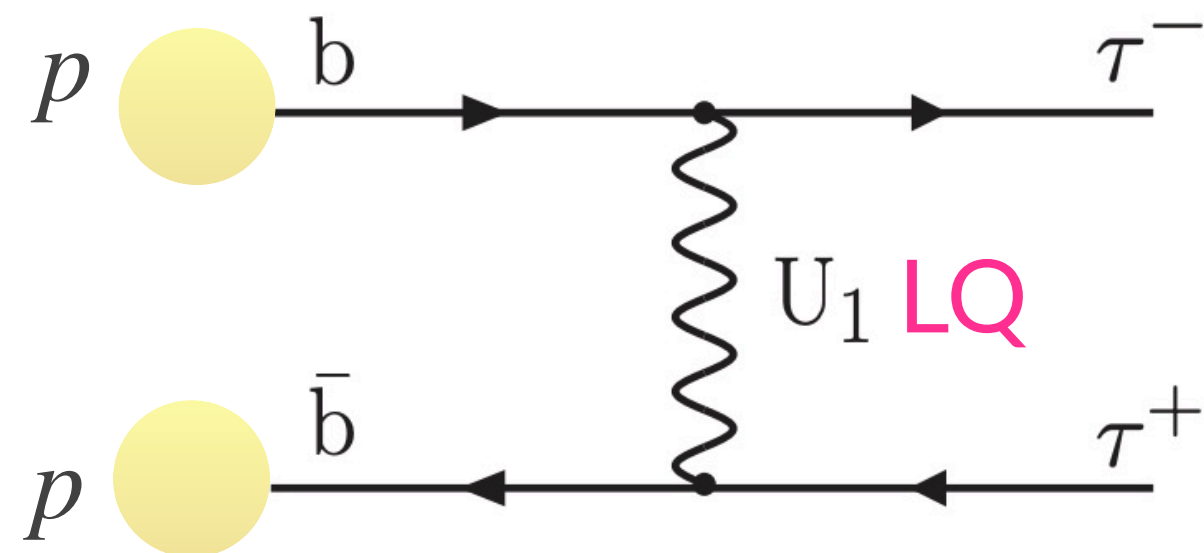
[cf. Angelescu, Bečirević, Faroughy, Jaffredo, Sumensari, [2103.12504](#); Athron, Balazs, Jacob, Kotlarski, Stockinger, Stockinger-Kim, [2104.03691](#)]

- ◆ Leptoquarks that do not lead to proton decay and can contribute precision measurements

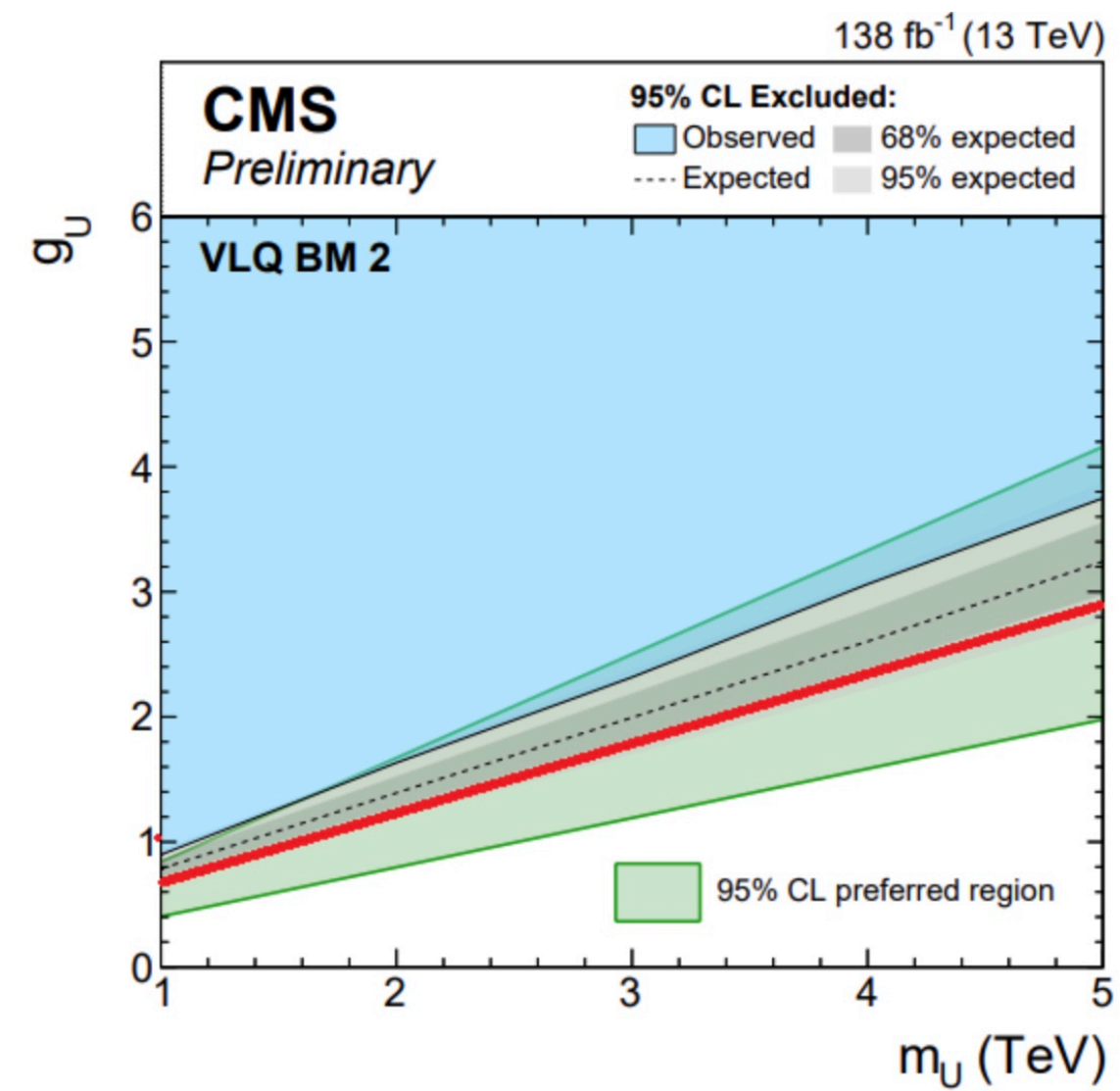
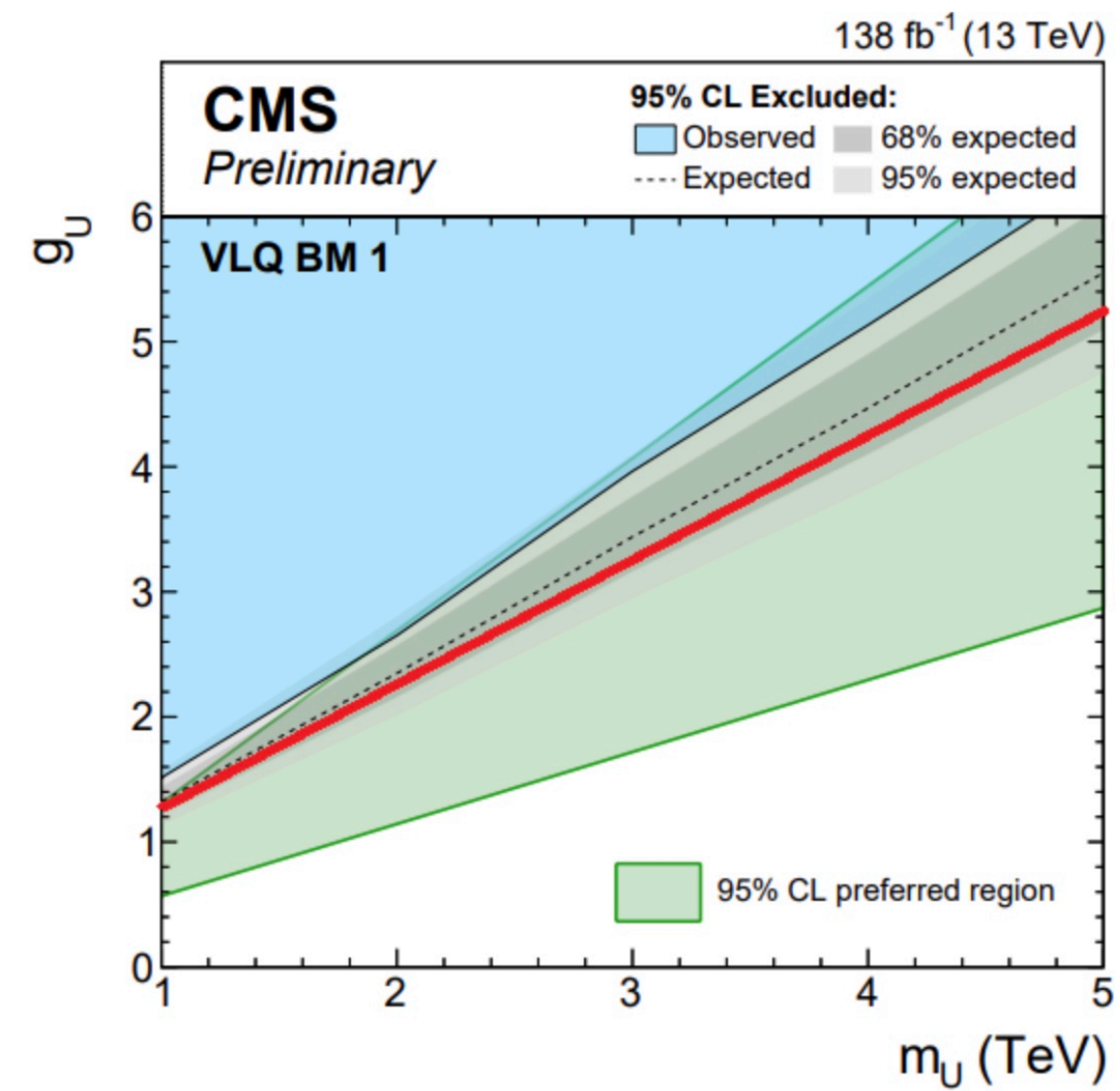
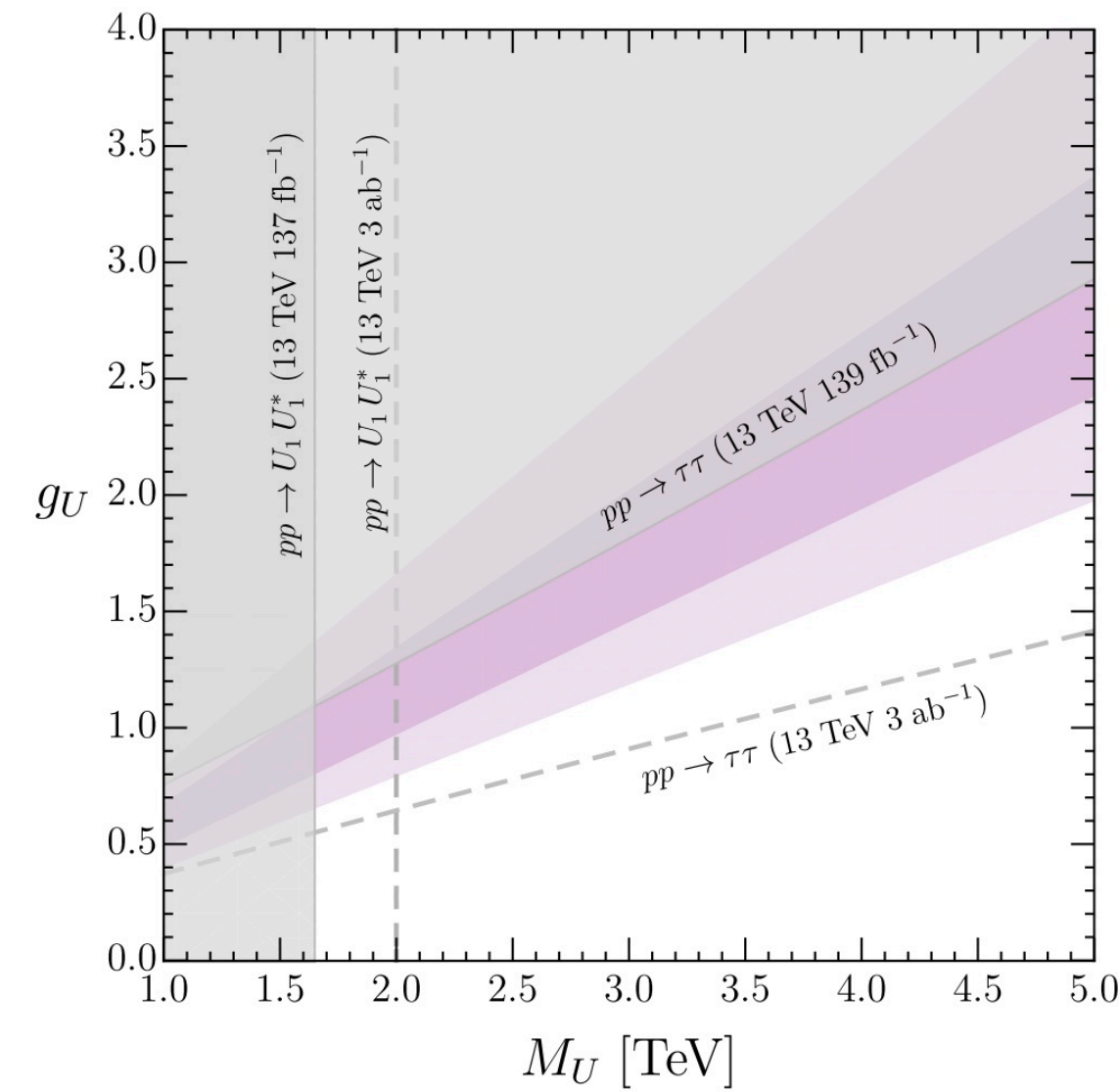
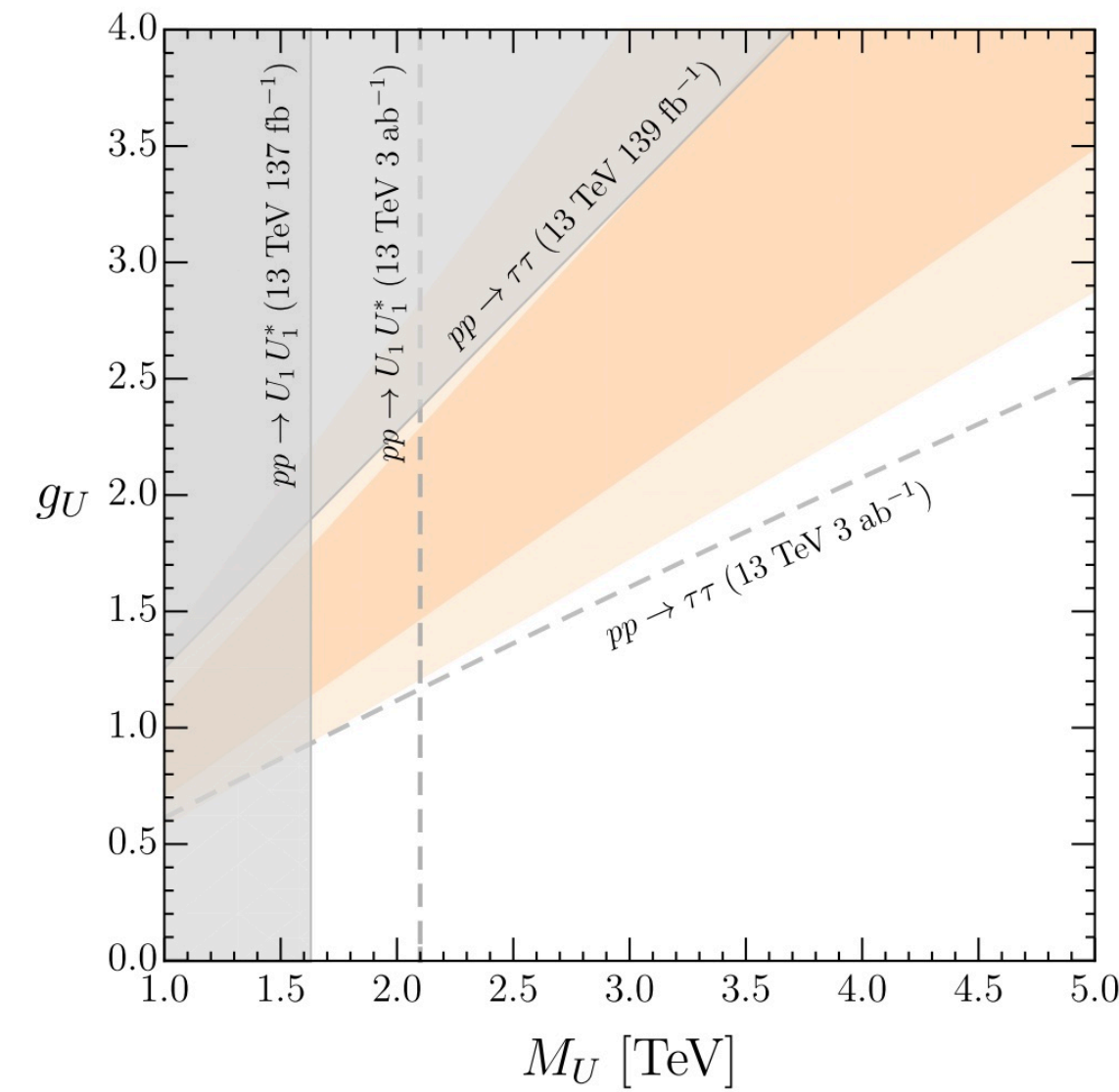
Label	Spin	Charge	R(D ^(*))	R(K ^(*))	muon g-2	M _W
S ₁ LQ	0	$(\bar{3}, 1, 1/3)$	✓	Loop	✓	With S ₃
U ₁ LQ	1	$(3, 1, 2/3)$	✓	✓	✗	✗
R ₂ LQ	0	$(3, 2, 7/6 [1/6])$	✓	Loop	✓	Small
V ₂ LQ	1	$(\bar{3}, 2, 5/6)$	Small	Electron	Small	✓
S ₃ LQ	0	$(\bar{3}, 3, 1/3)$	✗	✓	✗	With S ₁
U ₃ LQ	1	$(3, 3, 2/3)$	✗	✓	✗	?

← different model points →

Thanks Hirose-san



Unfortunately, the LHC sensitivity depends on the flavor structure of the leptoquark, significantly



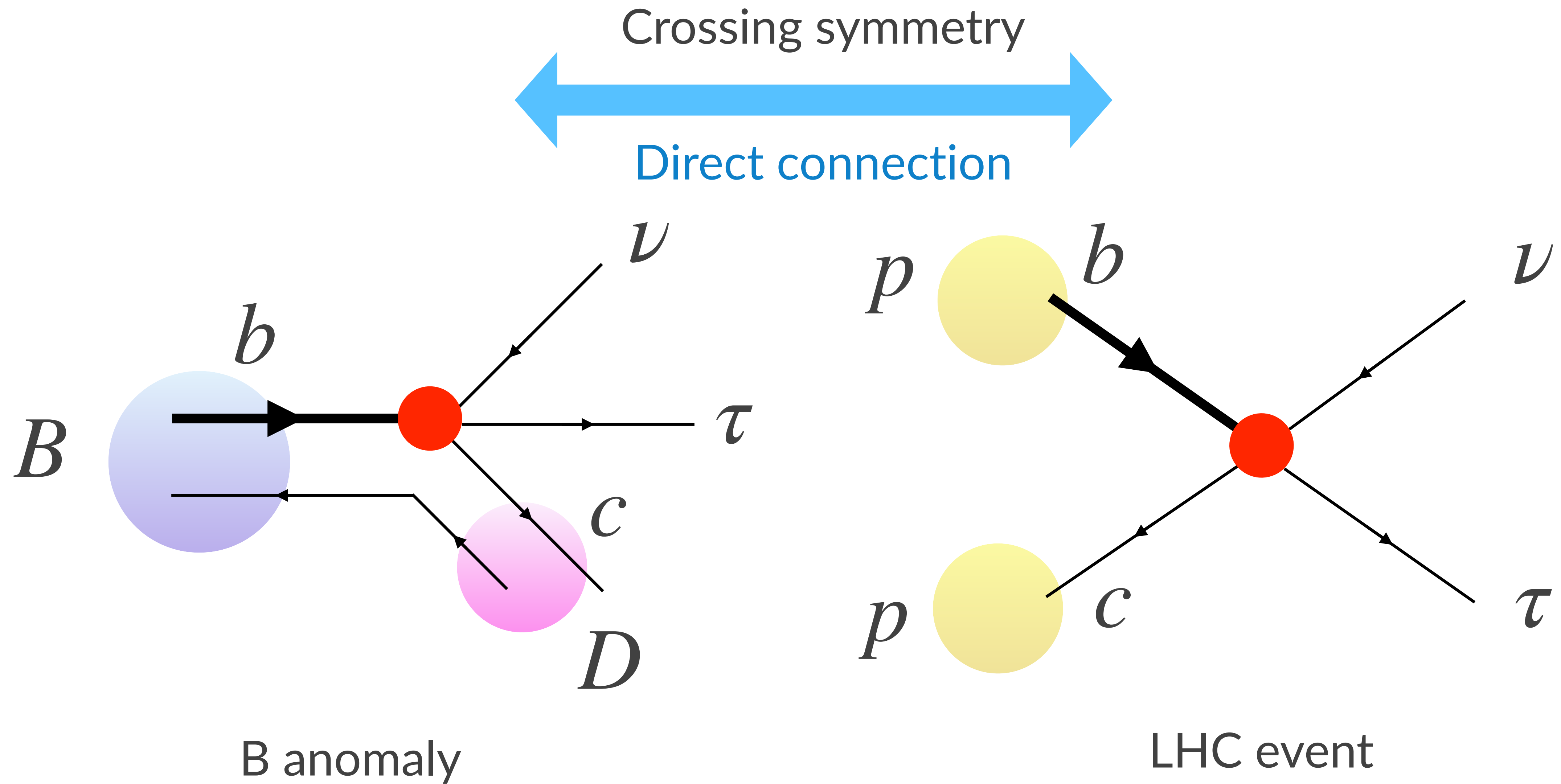
Recast the ATLAS data of $\tau^+\tau^-$ resonant search

[Cornella, Farouhy, Fuentes-Martin, Isidori, Neubert, [2103.16558](#); ATLAS, [2002.12223](#)]

The CMS data of $\tau^+\tau^-$ non-resonant search

[CMS, CMS PAS HIG-21-001]

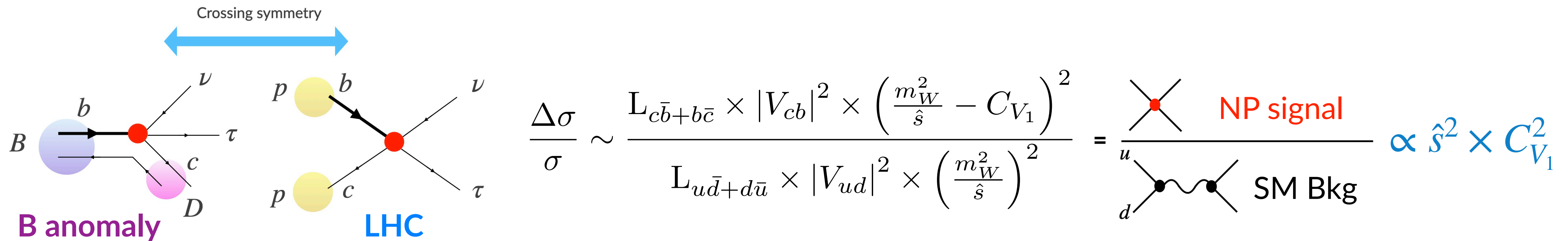
ATLAS
B anomaly



Collider search for $pp \rightarrow \tau\nu$

- ◆ The direct collider bound comes from **high- p_T tails in mono- τ searches**

[Greljo, Camalich, Ruiz-Alvarez, [1811.07920](#); Marzocca, Min, Son, [2008.07541](#); Iguro, Takeuchi, Watanabe, [2011.02486](#)]



LQ signal over Bkg is amplified at large \hat{s} : **hard τ with large missing transverse energy (= ν_τ)**

2 σ upper bound from data (36fb⁻¹)

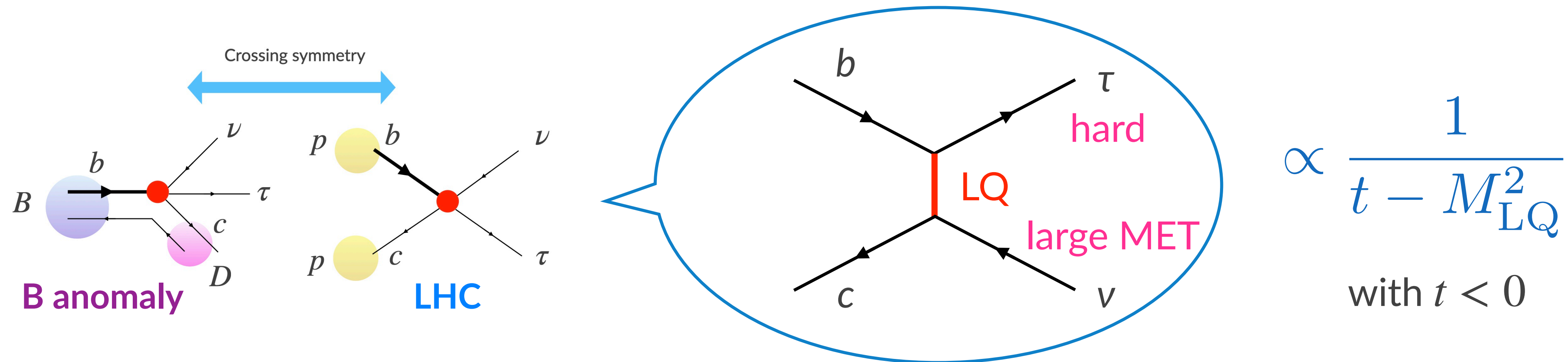
$$|C_{V_{1/2}}| < 0.32, \quad |C_{S_{1/2}}| < 0.57, \quad |C_T| < 0.16$$

[Greljo, Camalich, Ruiz-Alvarez, [1811.07920](#)]

Assuming EFT limit, namely $M_{LQ} \rightarrow \infty$

Leptoquark mass dependence

- ◆ We investigate LQ mass effects, which relax the collider bound [Iguro, Takeuchi, Watanabe, [2011.02486](#)]



To amplify the LQ signal over Bkg ratio, $t \sim -\mathcal{O}(1)\text{TeV}^2$ is expected.

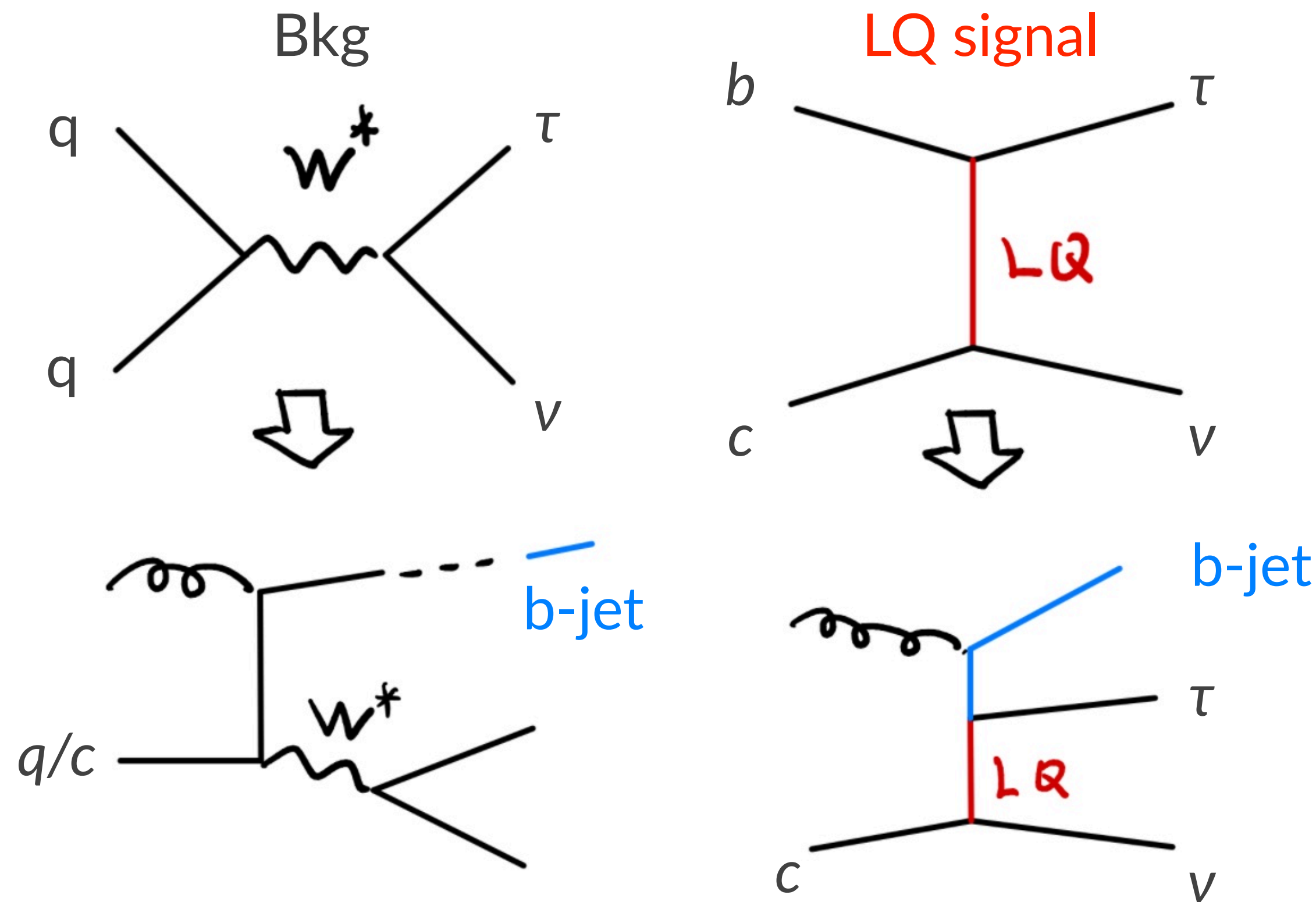
Now, the LQ mass receives additional effective mass via the large negative t

$$\mathcal{M}_{\text{LQ}} \approx \frac{g_{\text{LQ}}^2}{|t - M_{\text{LQ}}^2|} < \frac{g_{\text{LQ}}^2}{M_{\text{LQ}}^2} \approx \mathcal{M}_{\text{EFT}}$$

Collider bound should be relaxed when LQ mass is light

Additional b -jet in the final state

- ◆ Furthermore, a requirement of an additional b -jet is powerful [Marzocca, Min, Son, [2008.07541](#)]



Additional b -jet tagging suppresses the Bkg by several of powers

LQ signal suppresses only by a several factors

→ Signal over Bkg ratio can be amplified by the additional b -jet tagging

MC simulation (MadGraph5 + Pythia8 + Delphes) of both SM Bkg and leptoquark signal

SM Bkg (that mimics $\tau b + \text{missing}$)

BG (cut \mathbf{b})	Wjj	$Zjj (Z \rightarrow \nu\bar{\nu})$	$t\bar{t}$	$Z, \gamma DY$	VV	single t
number of jets	6693.4	235099	346.7	1813.2	125.8	151.8
number of τ	3173.5	5617.1	73.9	894.9	59.7	34.0
number of b	90.6	305.5	35.9	163.9	5.28	18.8
isolated lepton	90.5	305.5	29.7	10.4	1.38	17.0
τ kinematics	78.8	20.8	23.6	9.19	1.13	14.0
MET cut	71.2	4.62	20.9	2.52	0.98	12.7
back-to-back	7.84	3.61	1.67	0.57	0.18	0.54
$0.7 < m_T < 1 \text{ TeV}$	0.58	0.37	0.056	0.28	0.018	0.029
$1 \text{ TeV} < m_T$	0.16	0.06	0.01	0.007	0.005	0.005
$1 \text{ TeV} < m_T$ [34]	0.18(5)	0.21(12)	0.29(3)	$4.2(4) \times 10^{-5}$	0.35(5)	0.067(7)

Examples of LQs [U_1 LQ scenarios]

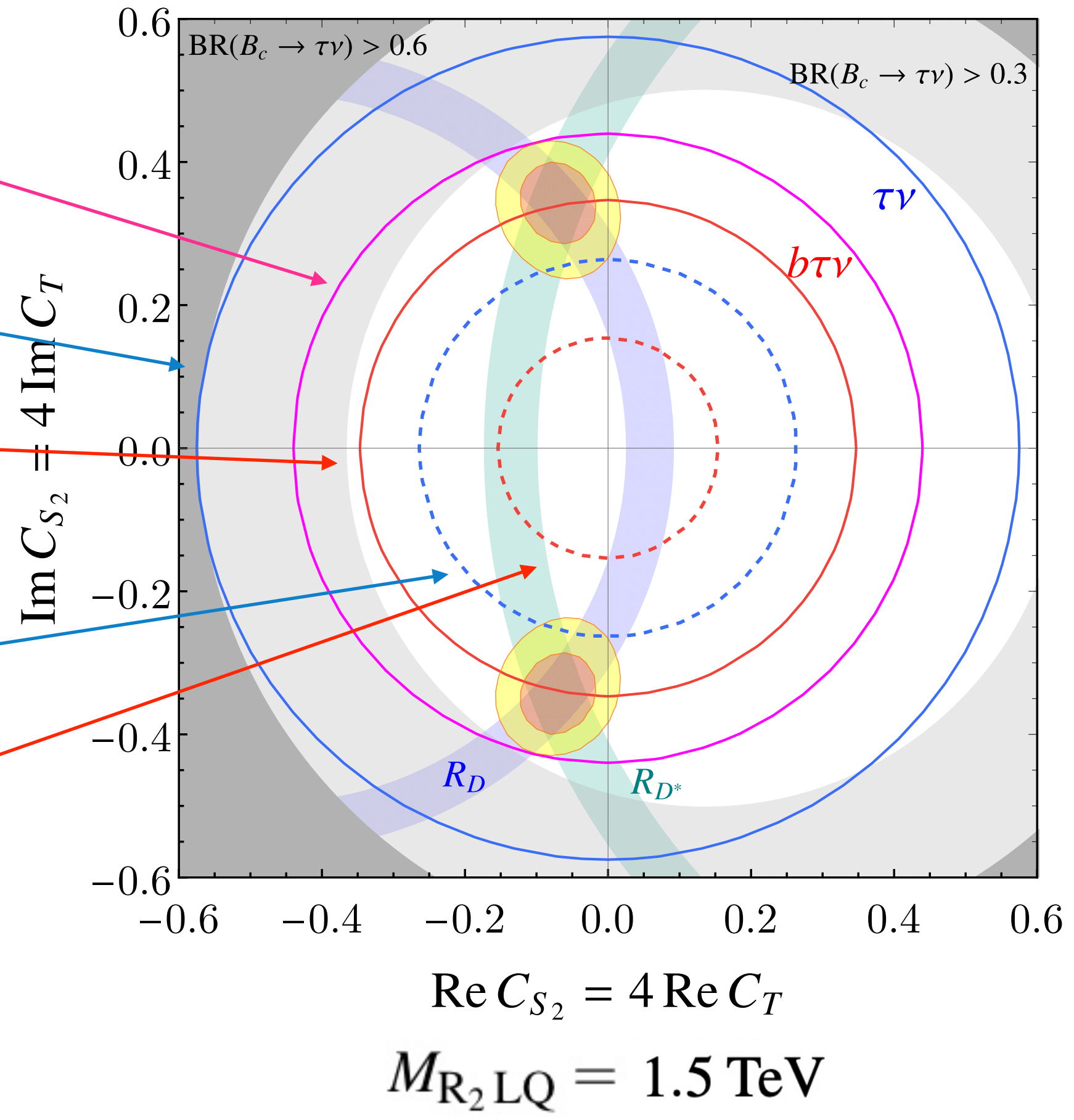
	$U_{1,1.5 \text{ TeV}}^{R_{s/b}=0}$	$U_{1,EFT}^{R_{s/b}=0}$	$U_{1,1.5 \text{ TeV}}^{R_{s/b}=1}$	$U_{1,EFT}^{R_{s/b}=1}$	BG
number of jets	4245	6085	5966	8376	244230
number of τ	2024	2941	2898	4168	9853
number of b	460	692	535	754	620.0
isolated lepton	454	685	485	747	454.5
τ kinematics	424	637	451	692	147.5
MET cut	350	540	371	590	112.9
back-to-back	258	402	263	443	14.4
$0.7 < m_T < 1 \text{ TeV}$	53.9	86.4	55.8	92.0	1.33
$1 \text{ TeV} < m_T$	26.0	71.6	30.7	101	0.25

A transverse mass defined as:
$$m_T = \sqrt{2p_T^\tau E_T^{\text{miss}} [1 - \cos \Delta\phi(\vec{p}_T^\tau, \vec{p}_T^{\text{miss}})]}$$

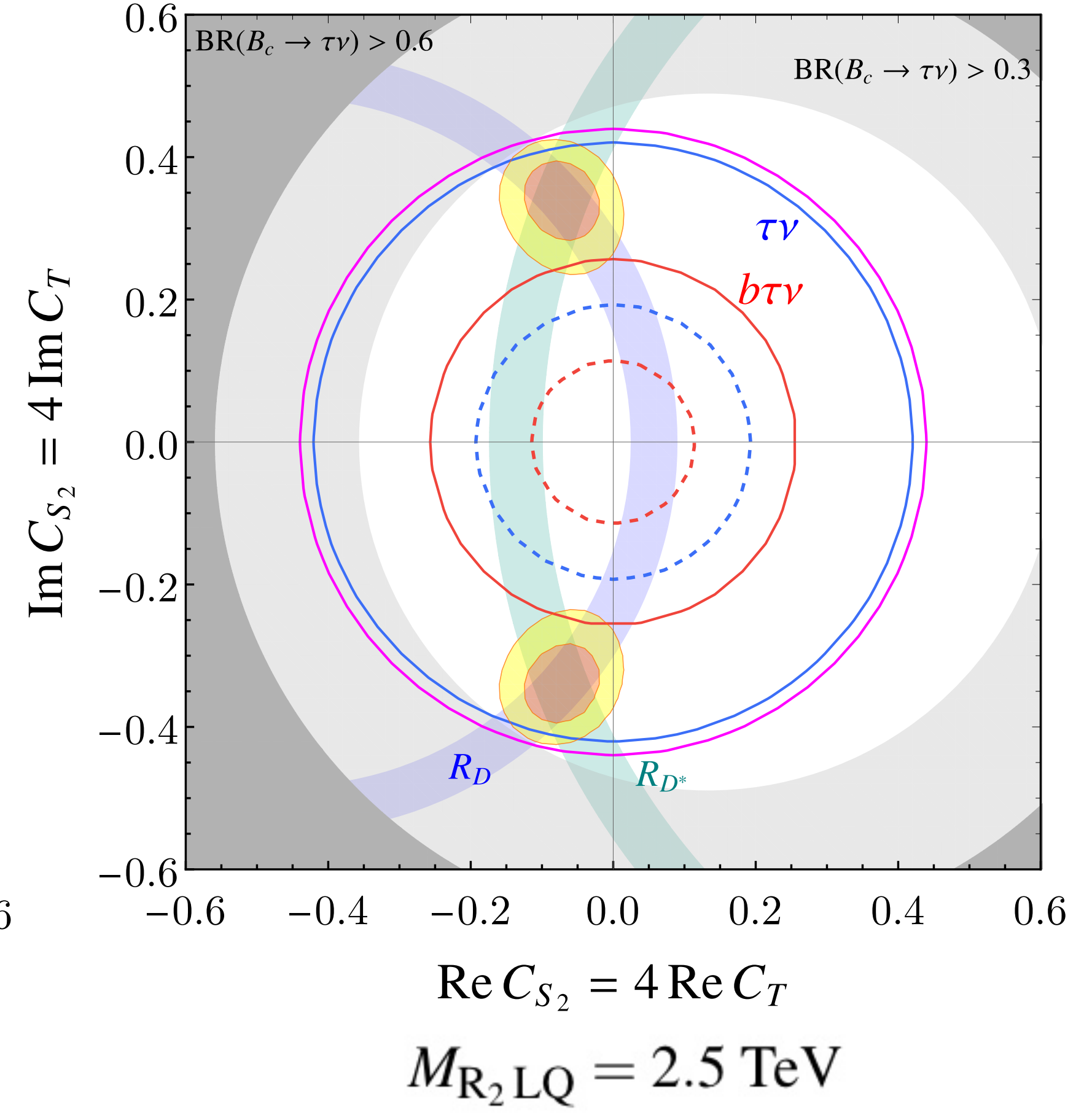
R₂ Leptoquark scenario

[Endo, Iguro, TK, Takeuchi, Watanabe, [2111.04748](#)]

- τ +MET search
36fb⁻¹ exclusion
- τ +MET search
139fb⁻¹ sensitivity
- τ +MET+b search
139fb⁻¹ sensitivity
- τ +MET search
3000fb⁻¹ sensitivity
- τ +MET+b search
3000fb⁻¹ sensitivity



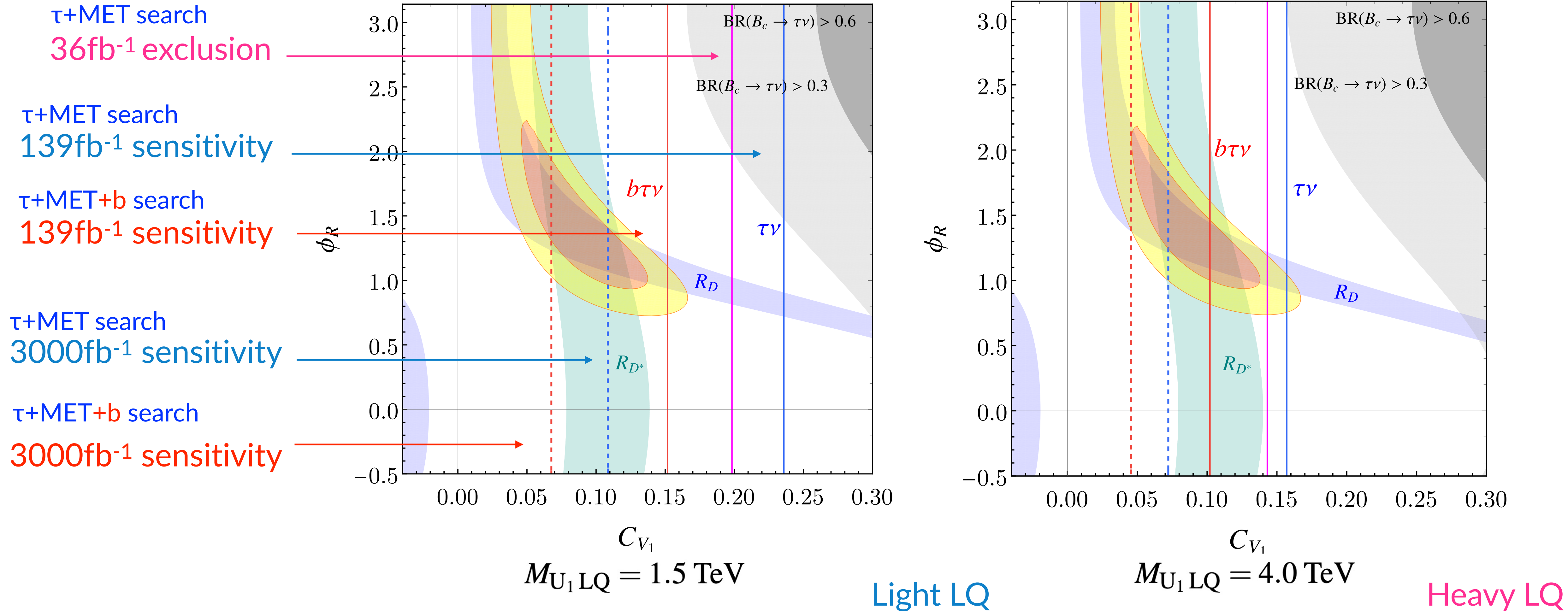
Light LQ



Heavy LQ

U₁ Leptoquark scenario

[Endo, Iguro, TK, Takeuchi, Watanabe, [2111.04748](#)]



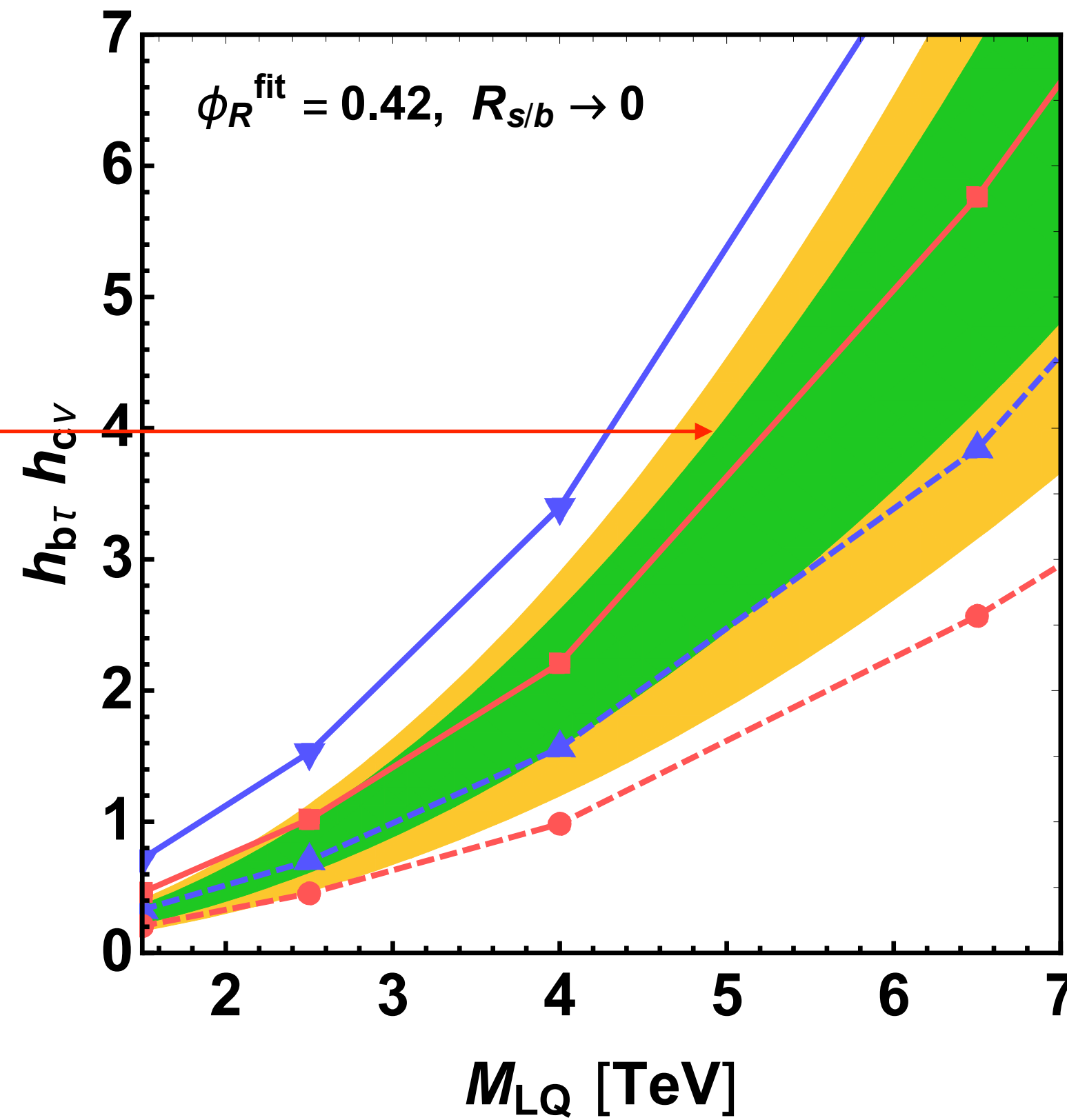
U₁ Leptoquark scenario: comparison

[Endo, Iguro, TK, Takeuchi, Watanabe, [2111.04748](#)]

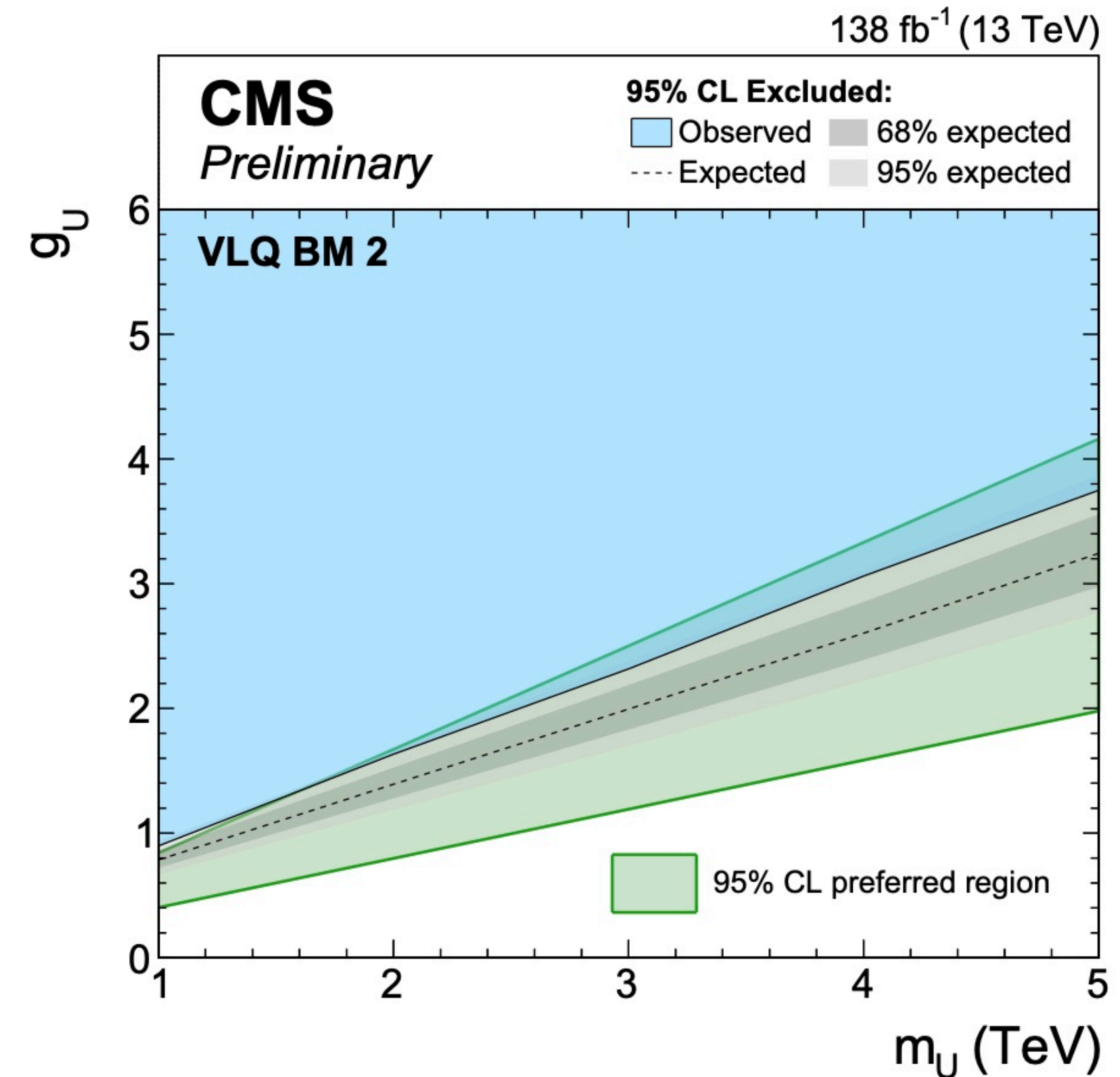
[CMS, CMS PAS HIG-21-001]

$$U_1 \text{ LQ: } C_{S_1} = -2 e^{i\phi_R} C_{V_1}$$

τ +MET+b search
139fb⁻¹ sensitivity



τb + missing sensitivity (139 fb⁻¹)



$\tau\tau$ + excluded region (138 fb⁻¹)

Conclusion

- ◆ Several leptoquark can easily explain both $R(D)$ and $R(D^*)$ 4σ levels anomalies
- ◆ We simulate the (HL-)LHC sensitivity of the leptoquark indirect search via $pp \rightarrow \tau \nu + b$ as well as $pp \rightarrow \tau \nu$
- ◆ We show that additional b-jet tagging significantly improves the LHC sensitivity by $\approx 40\%$
- ◆ We show the EFT breakdown behavior in light leptoquark region (see paper)
- ◆ We investigate the angular correlations to discriminate the LQ scenarios (see paper)

Backup

Latest SM predictions of $R(D)$ and $R(D^*)$

HFLAV theory average2019 $R(D)_{\text{SM}} = 0.299 \pm 0.003$ $R(D^*)_{\text{SM}} = 0.258 \pm 0.005$

New analyses

- ◆ + All lattice data, QCD sum rule, and the latest LCSR result [Gubernari, Kokulu, van Dyk '18]
 $\text{@}q^2 = q_{\text{max}}^2$ $\text{@}q^2 \leq 0$
- ◆ + All $\mathcal{O}(\Lambda^2/m_c^2)$ corrections in the heavy quark effective theory in all form factors [Jung, Straub '18]
- ◆ + Momentum distributions from Belle data [Bordone, Jung, van Dyk '19]
 $R(D)_{\text{SM}} = 0.297 \pm 0.003$ $R(D^*)_{\text{SM}} = 0.250 \pm 0.003$ [BJvD]
- ◆ + Angular distributions from Belle data [Iguro Watanabe '20] [Available Belle data
1510.03657;
1702.01521;
1809.03290]
 $R(D)_{\text{SM}} = 0.289 \pm 0.004$ $R(D^*)_{\text{SM}} = 0.248 \pm 0.001$ [IW]

$R(D)$: 1.4 [HFLAV2019] \rightarrow 1.4 [BJvD], 1.7 σ [IW]
 $R(D^*)$: 2.5 \rightarrow 3.2, 3.4 σ
combine: 3.1 \rightarrow 3.9, 4.2 σ