



# Higgs Early Dark Energy by non-thermal trapping effect

**Physics in LHC and Beyond**  
**May.15, 2022**

Shota Nakagawa (Tohoku U.)

Based on arXiv: 2205(6?).XXXXX with F. Takahashi and W. Yin  
cf) PRD 105 (2022) 10. [arXiv:2111.06696] with N. Kitajima and F. Takahashi

# 1. Introduction

The standard cosmology ( $\Lambda$ CDM) predicts the cosmic evolution which is consistent with observations. However, there may be deviations from  $\Lambda$ CDM model.

## Hubble ( $H_0$ ) tension

Direct measurement (CMB+ $\Lambda$ CDM)

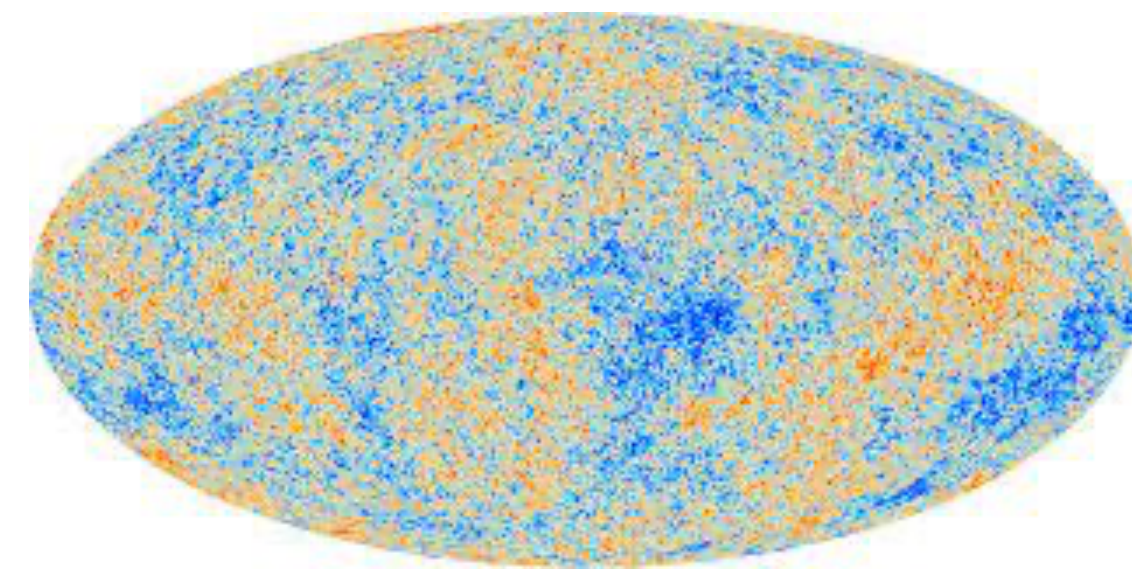
$$H_0^{(\text{direct})} = 67.27 \pm 0.60 \text{ km/s/Mpc}$$

Planck2018, 1807.06209

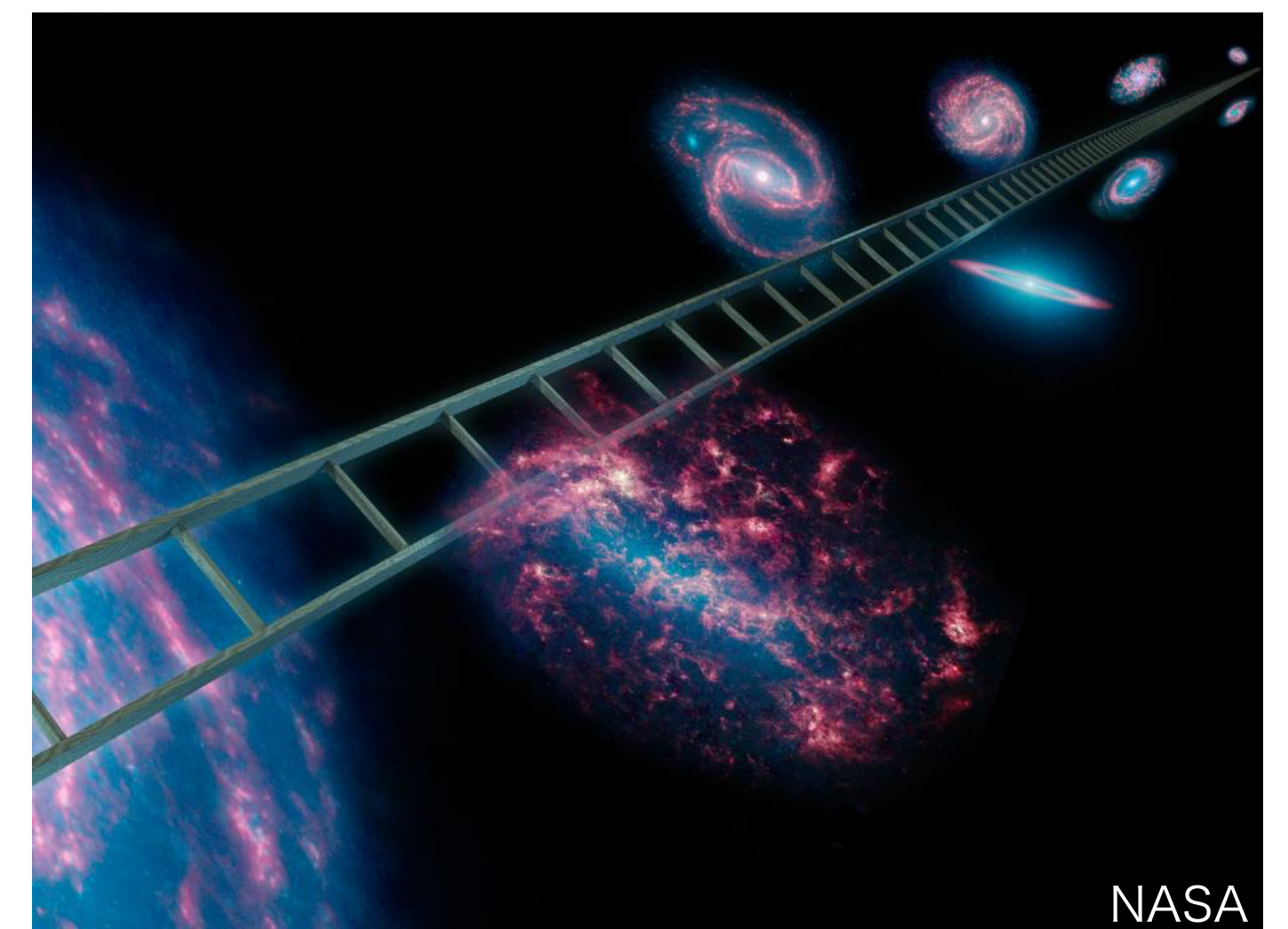
Indirect measurement w/ distance ladder

$$H_0^{(\text{indirect})} = 73.04 \pm 1.04 \text{ km/s/Mpc}$$

A. Riess, et al. 2112.04510 (SNIa+Cepheid stars)



Planck collaboration



# 1. Introduction

The standard cosmology ( $\Lambda$ CDM) predicts the cosmic evolution which is consistent with observations. However, there may be deviations from  $\Lambda$ CDM model.

## Hubble ( $H_0$ ) tension

Direct measurement (CMB+ $\Lambda$ CDM)

$$H_0^{(\text{direct})} = 67.27 \pm 0.60 \text{ km/s/Mpc}$$

Planck2018, 1807.06209

Indirect measurement w/ distance ladder

$$H_0^{(\text{indirect})} = 73.04 \pm 1.04 \text{ km/s/Mpc}$$

A. Riess, et al. 2112.04510 (SNIa+Cepheid stars)



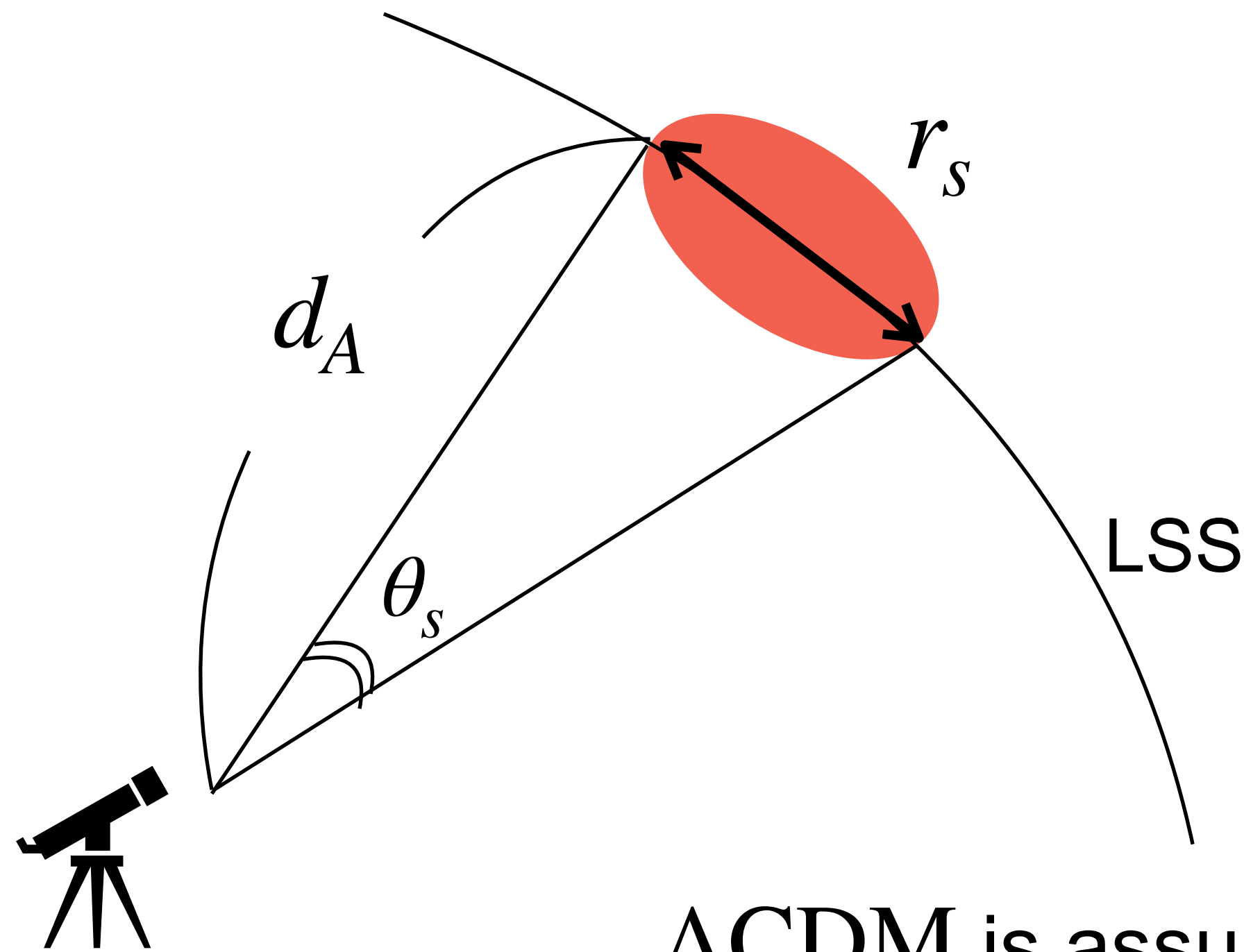
Modified so that  $H_0$  increases

Planck collaboration

$\sim 5\sigma!$

New Physics??

# $H_0$ from CMB + cosmological model



Useful object : Sound horizon

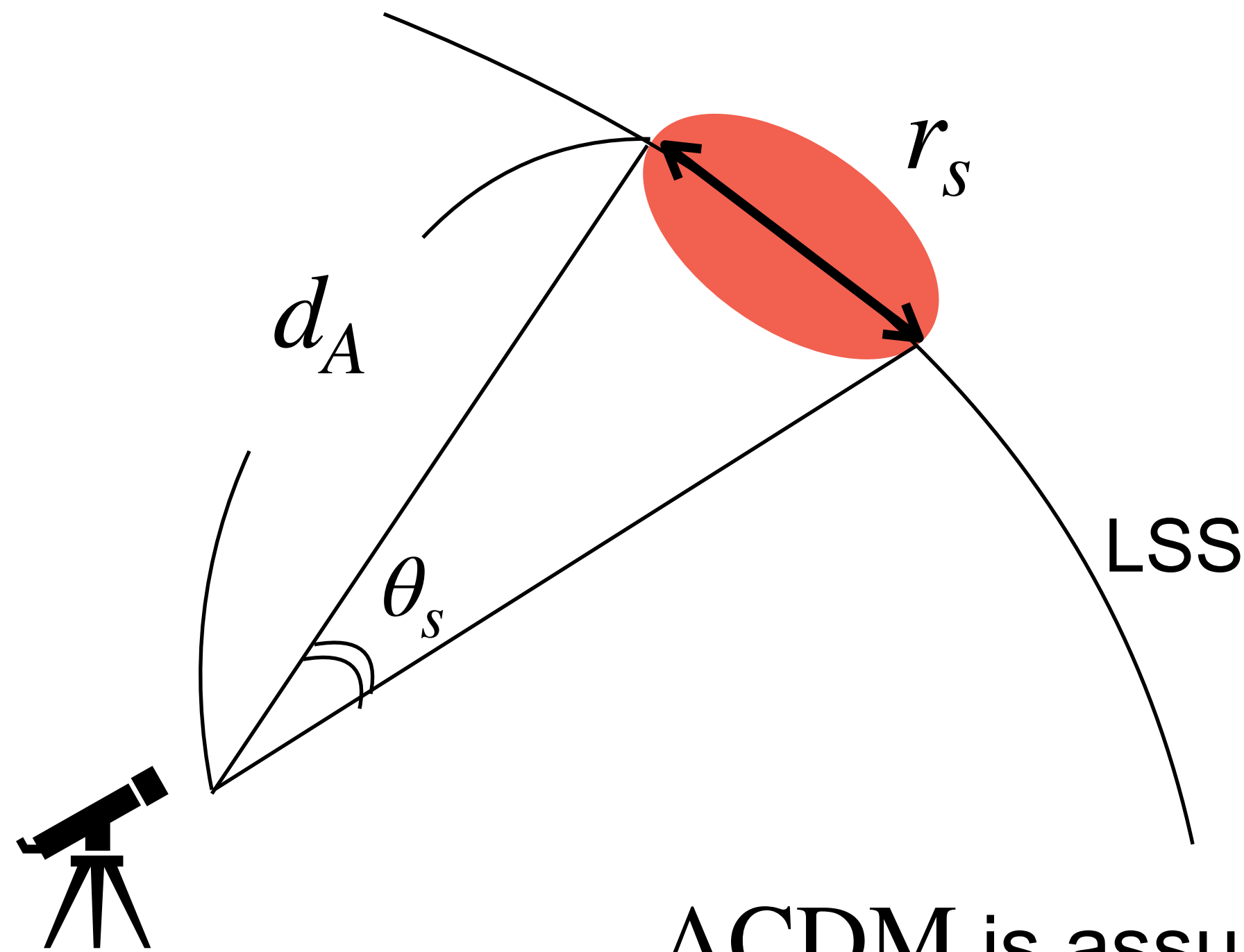
$$r_s \equiv \int_{z_{\text{rec}}}^{\infty} dz \frac{c_s(z)}{H(z)}$$

$\Lambda$ CDM is assumed.

$$\theta_s = \frac{r_s}{d_A} \stackrel{\downarrow}{\approx} \frac{H_0 r_s}{\int_0^{z_{\text{rec}}} dz [(1+z)^3 \Omega_m + (1-\Omega_m)]^{-1/2}}$$

CMB data gives  $\theta_s$ ,  $\Omega_m$ ,  $\Omega_b$ ,  $\Omega_r \dots$  and  $r_s$  is derived.

# $H_0$ from CMB + cosmological model



Useful object : Sound horizon

$$r_s \equiv \int_{z_{\text{rec}}}^{\infty} dz \frac{c_s(z)}{H(z)}$$

$\Lambda$ CDM is assumed.

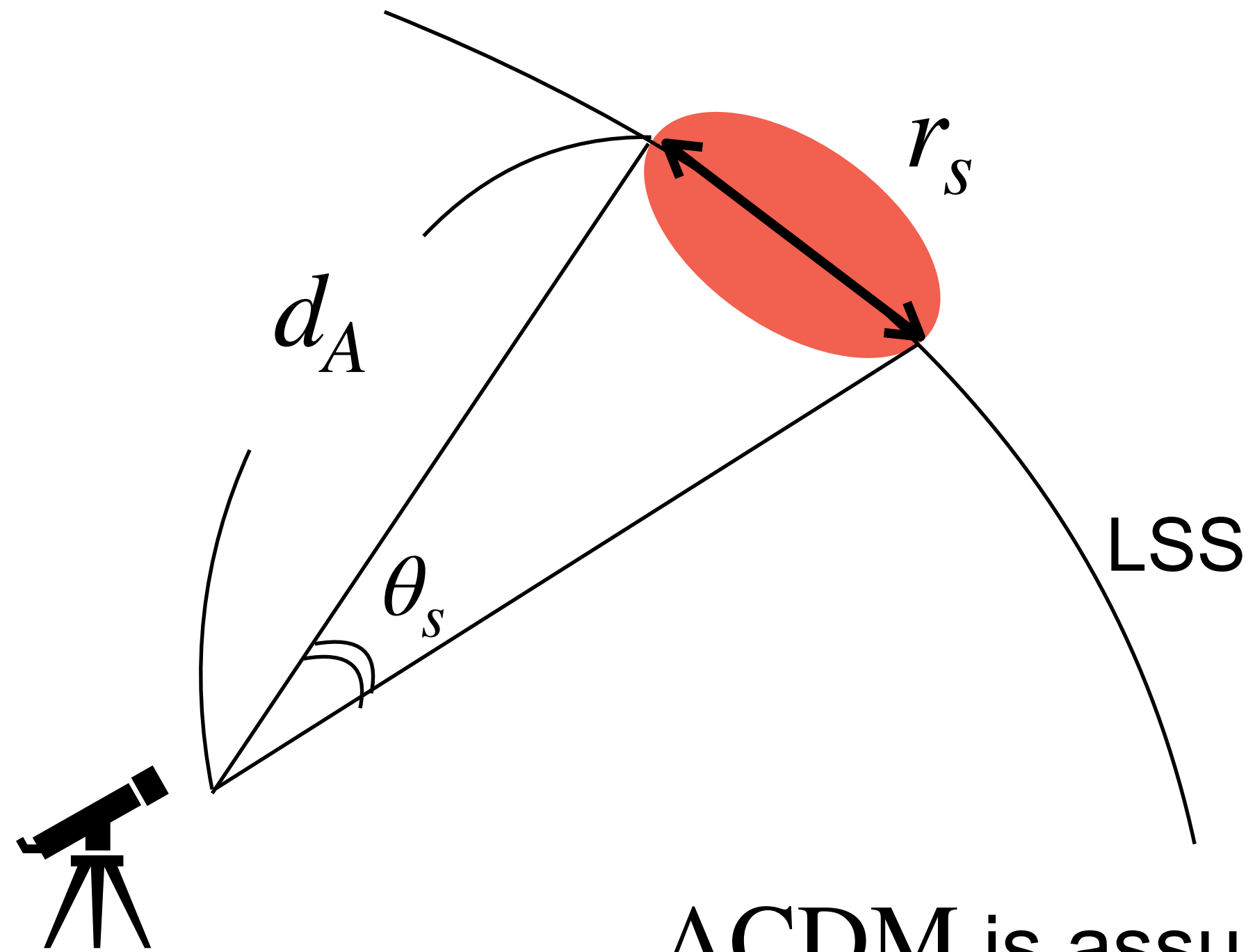
$$\underline{\underline{\theta_s}} = \frac{r_s}{d_A} \underset{\downarrow}{\simeq} \frac{H_0 r_s}{\int_0^{z_{\text{rec}}} dz [(1+z)^3 \Omega_m + (1-\Omega_m)]^{-1/2}}$$

Observable

With Late-time evolution left unchanged

CMB data gives  $\theta_s, \Omega_m, \Omega_b, \Omega_r \dots$  and  $r_s$  is derived.

# $H_0$ from CMB + cosmological model



Useful object : Sound horizon

$$r_s \equiv \int_{z_{\text{rec}}}^{\infty} dz \frac{c_s(z)}{H(z)}$$

Early dark energy (EDE)

$\Lambda$ CDM is assumed.

$$\theta_s = \frac{r_s}{d_A} \approx \frac{H_0 r_s}{\int_0^{z_{\text{rec}}} dz [(1+z)^3 \Omega_m + (1 - \Omega_m)]^{-1/2}}$$

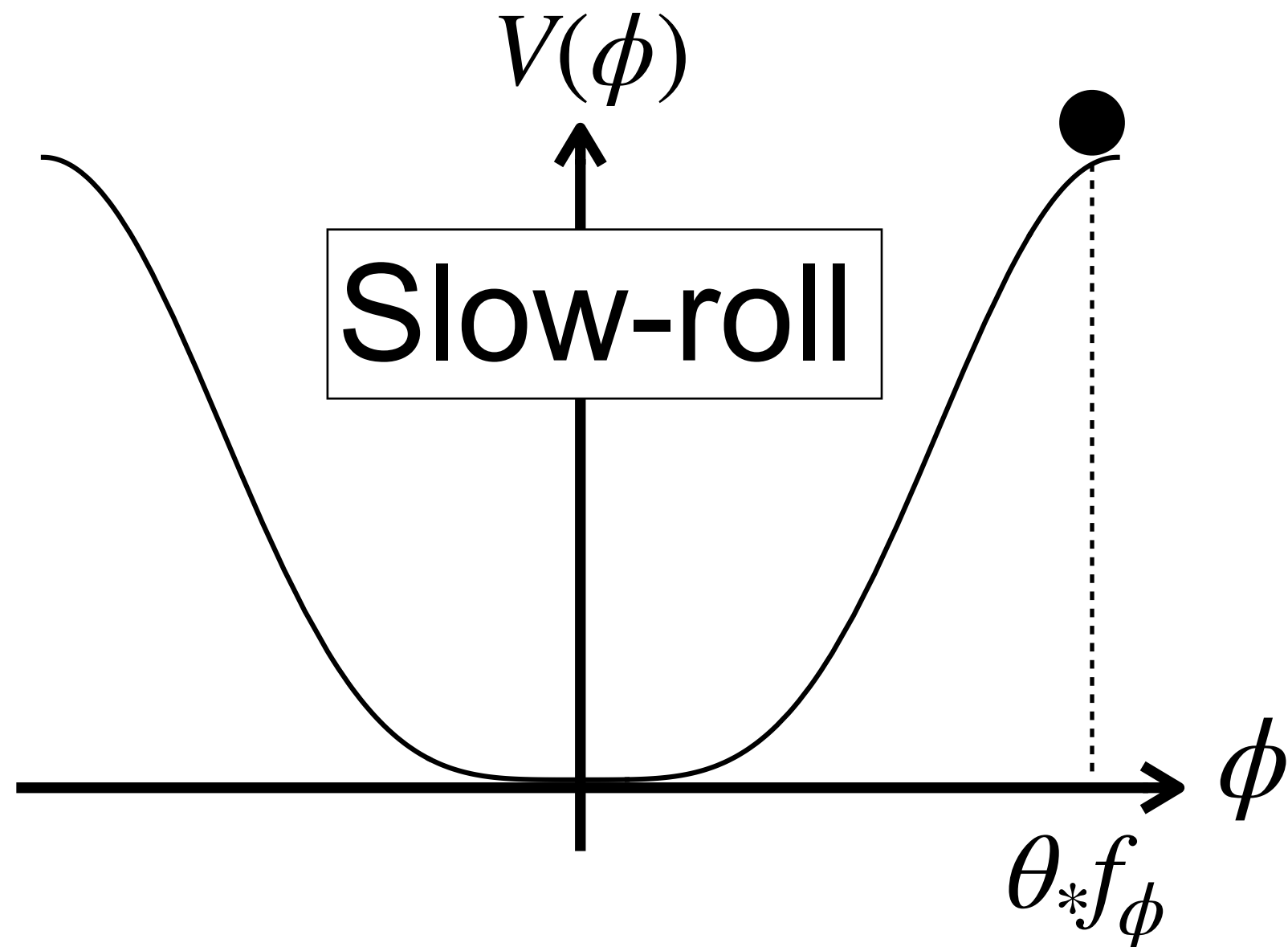
Observable

With Late-time evolution left unchanged

CMB data gives  $\theta_s, \Omega_m, \Omega_b, \Omega_r \dots$  and  $r_s$  is derived.

# Example for EDE

V. Poulin, et al. (2019)

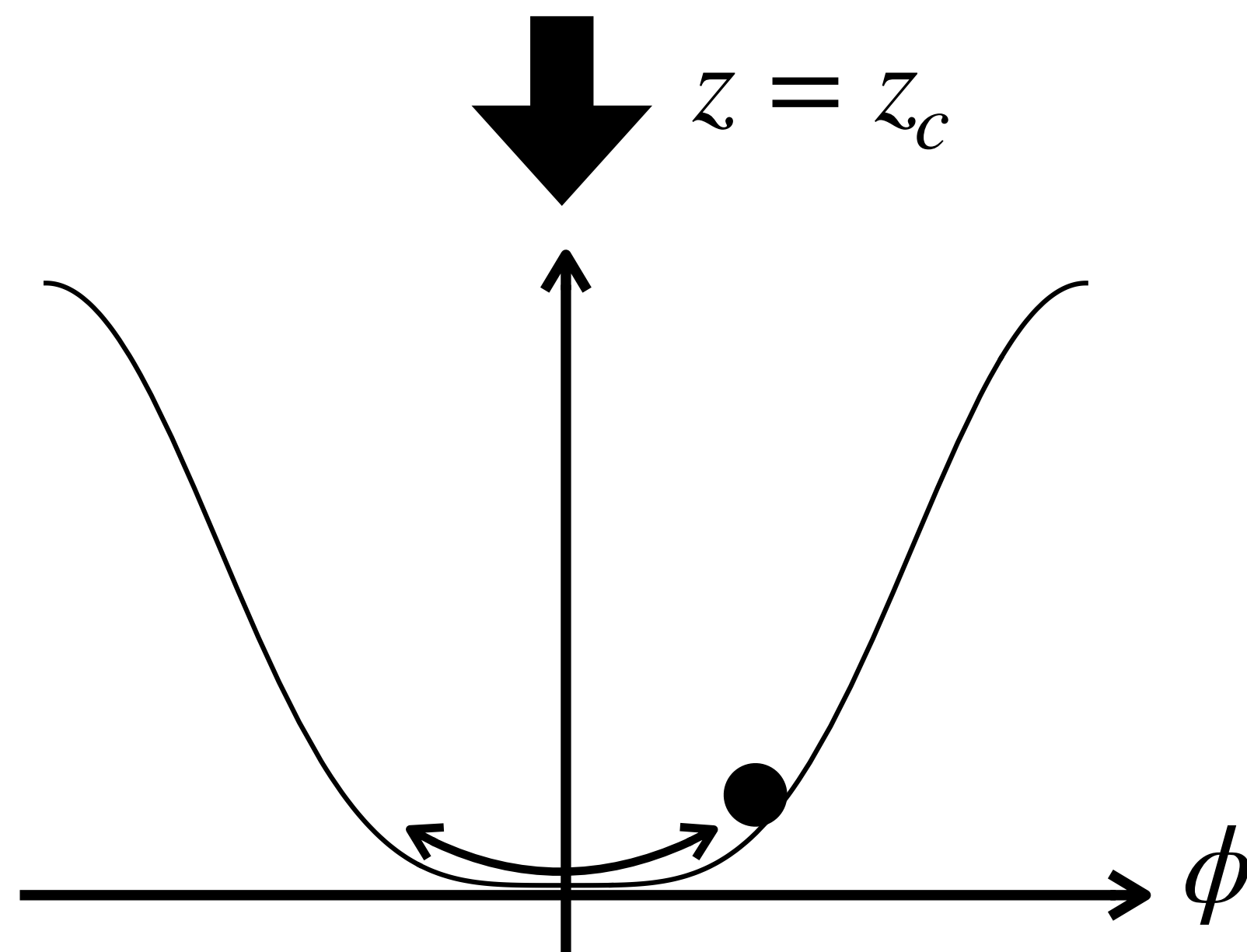


$$V(\phi) = m_\phi^2 f_\phi^2 \left( 1 - \cos \frac{\phi}{f_\phi} \right)^n$$

$\phi$  : axion  
 $f_\phi$  : decay constant

M. Kamionkowski, et al. 1409.0549

$n \geq 2$  Diluted equal to or faster than radiation **Late-time evolution is not affected.**



We require

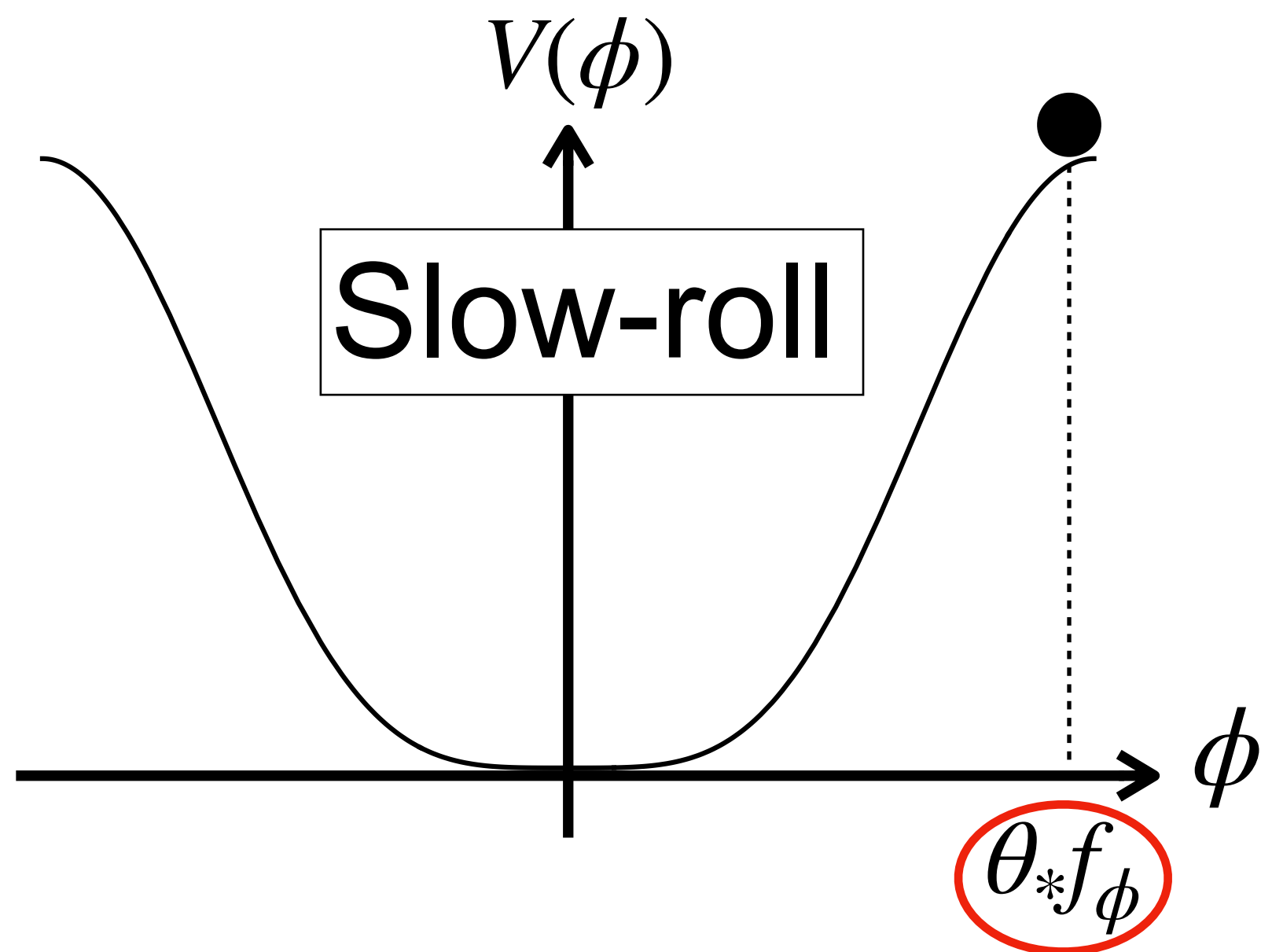
$$f_{\text{EDE}} \equiv \frac{\rho_\phi}{\rho_{\text{tot}}}(z_c) \simeq 0.05$$

$$z = z_c \simeq 5000$$

# Example for EDE

V. Poulin, et al. (2019)

Can the potential shape be simpler?



$$V(\phi) = m_\phi^2 f_\phi^2 \left( 1 - \cos \frac{\phi}{f_\phi} \right)^n$$

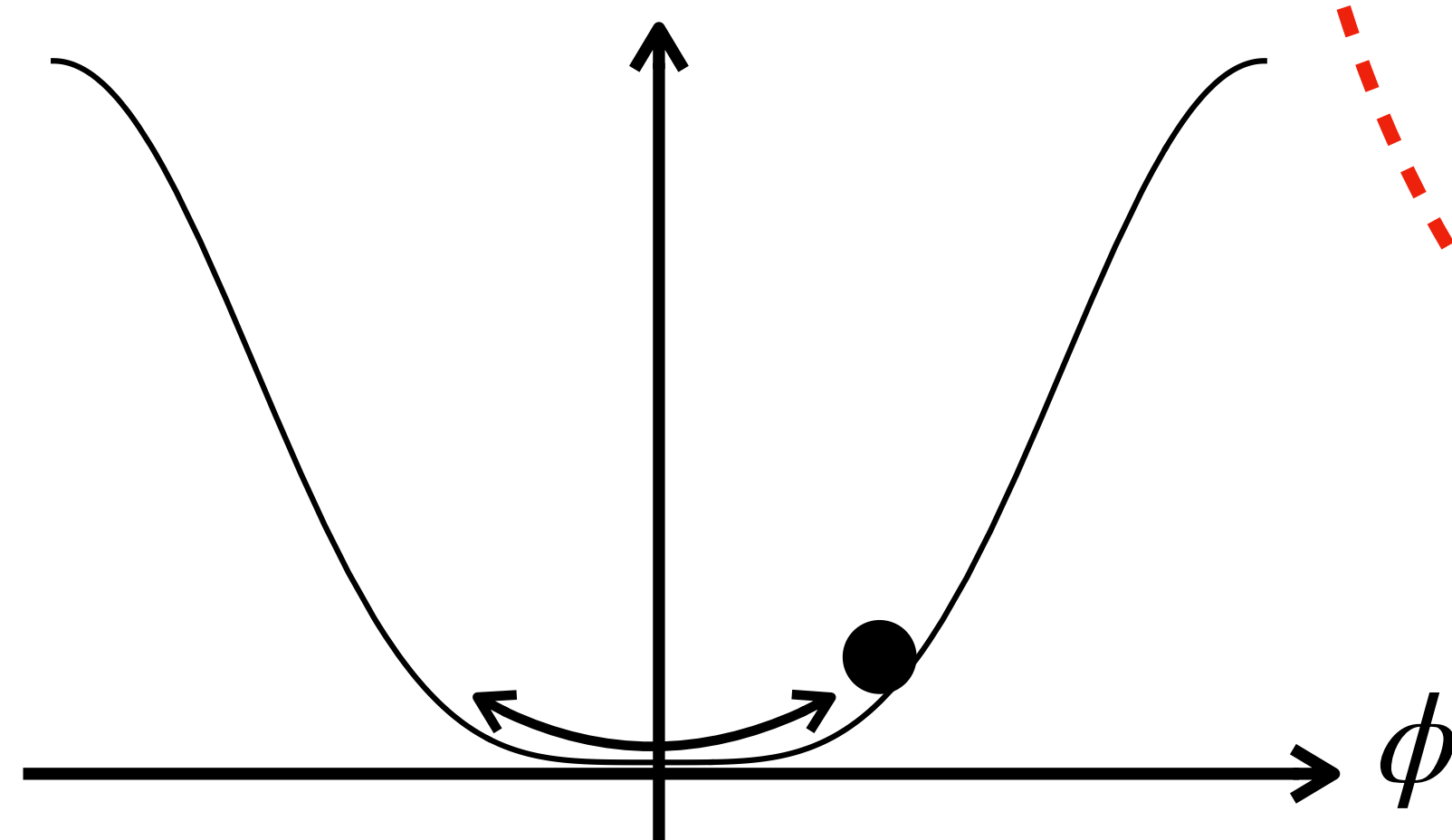
$\phi$  : axion  
 $f_\phi$  : decay constant

M. Kamionkowski, et al. 1409.0549

$n \geq 2$  Diluted equal to or faster than radiation

Late-time evolution is not affected.

$z = z_c$



We require

$$f_{\text{EDE}} \equiv \frac{\rho_\phi}{\rho_{\text{tot}}}(z_c) \simeq 0.05$$

$$z = z_c \simeq 5000$$

$|\pi - \theta_*| \ll 1$



# What we did

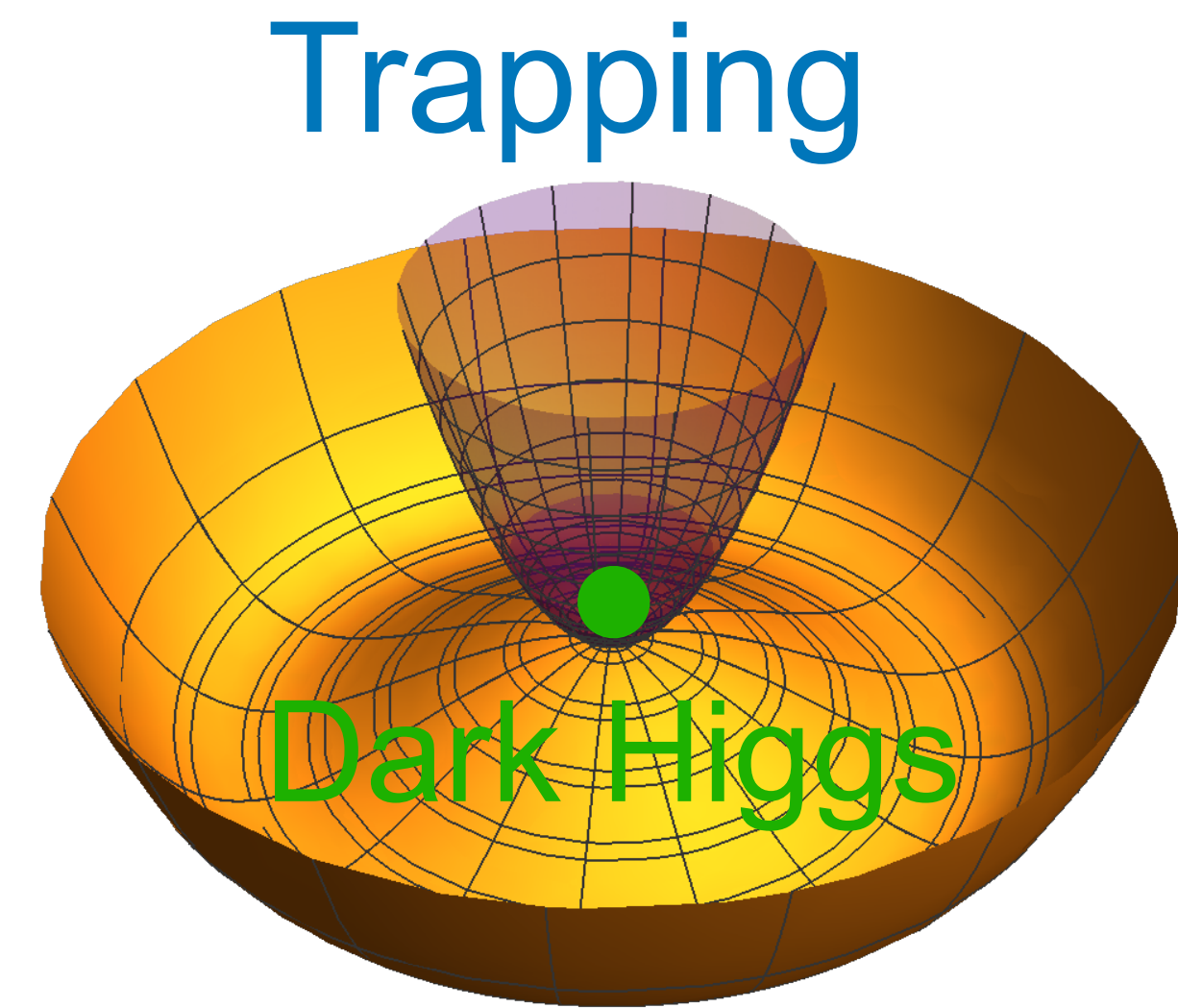
## Our suggestions for simple EDE potential

- Trapped scalar, not slow-rolling

The origin of this problem comes from slow-roll.  
So, we consider trapped scalar field as EDE.

Thermal potential is difficult because the universe is already cold enough around  $z \sim z_c$ .

→ **Non-thermal trapping of dark Higgs** N. Kitajima, S.N., F. Takahashi (2022)



- How to dilute so as not to affect late-time universe

The dark Higgs decays into dark photons promptly at  $z \sim z_c$ .

# 2. Non-thermal trapping effect

Abelian Higgs model with an axion coupled to dark photons

$$\mathcal{L} = (D_\mu \Psi)^\dagger D^\mu \Psi - V_\Psi(\Psi, \Psi^\dagger) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V_\phi(\phi) - \frac{\beta}{4f_\phi} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

where  $D_\mu = \partial_\mu - ieA_\mu$

$\Psi$  : dark Higgs

$A_\mu$  : dark photon

$\phi$  : axion

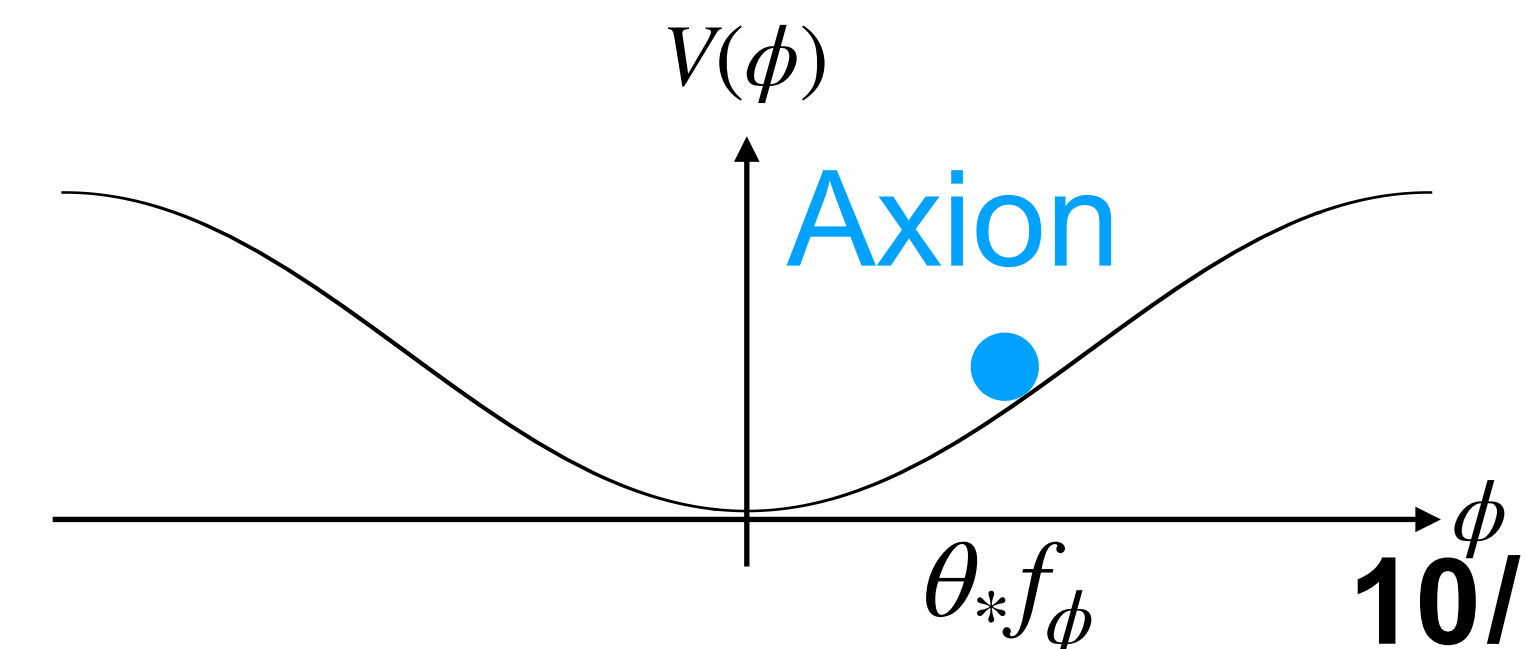
Higgs potential

$$V_\Psi(\Psi, \Psi^\dagger) = \frac{\lambda}{4} (|\Psi|^2 - v^2)^2$$

where  $f_{\text{EDE}} \equiv V_\Psi / \rho_{\text{tot}}(z_c) \simeq 0.05$

Axion potential

$$V_\phi(\phi) = m_\phi^2 f_\phi^2 \left[ 1 - \cos \left( \frac{\phi}{f_\phi} \right) \right]$$



# 2. Non-thermal trapping effect

Abelian Higgs model with an axion coupled to dark photons

$$\mathcal{L} = (D_\mu \Psi)^\dagger D^\mu \Psi - V_\Psi(\Psi, \Psi^\dagger) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V_\phi(\phi) - \frac{\beta}{4f_\phi} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

where

$$(D_\mu \Psi)^\dagger D^\mu \Psi = |\partial_\mu \Psi|^2 + \underline{e^2 A_\mu A^\mu |\Psi|^2}$$

Dark Higgs acquires an effective mass from dark photon production.

$A_\mu$  : dark photon

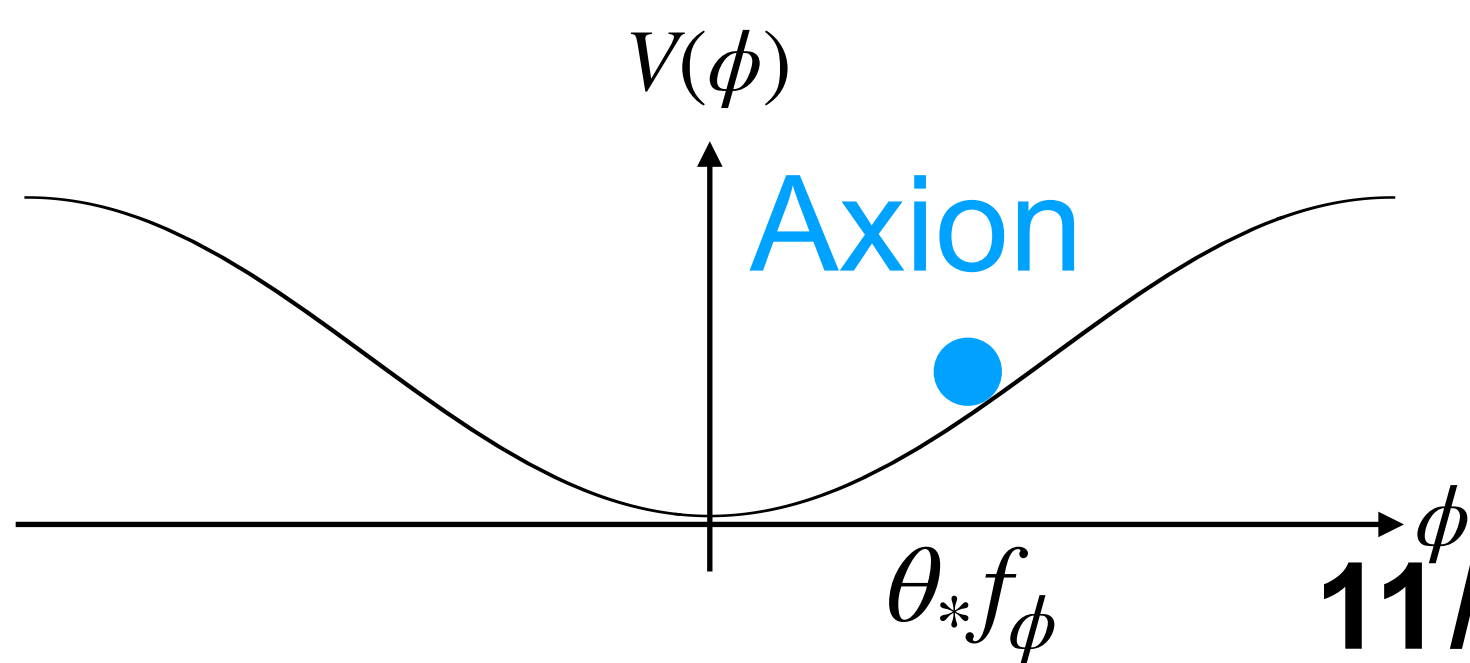


Higgs p

where  $f_{\text{EDE}} \equiv V_\Psi / \rho_{\text{tot}}(z_c) \simeq 0.05$

Axion potential

$$V_\phi(\phi) = m_\phi^2 f_\phi^2 \left[ 1 - \cos \left( \frac{\phi}{f_\phi} \right) \right]$$

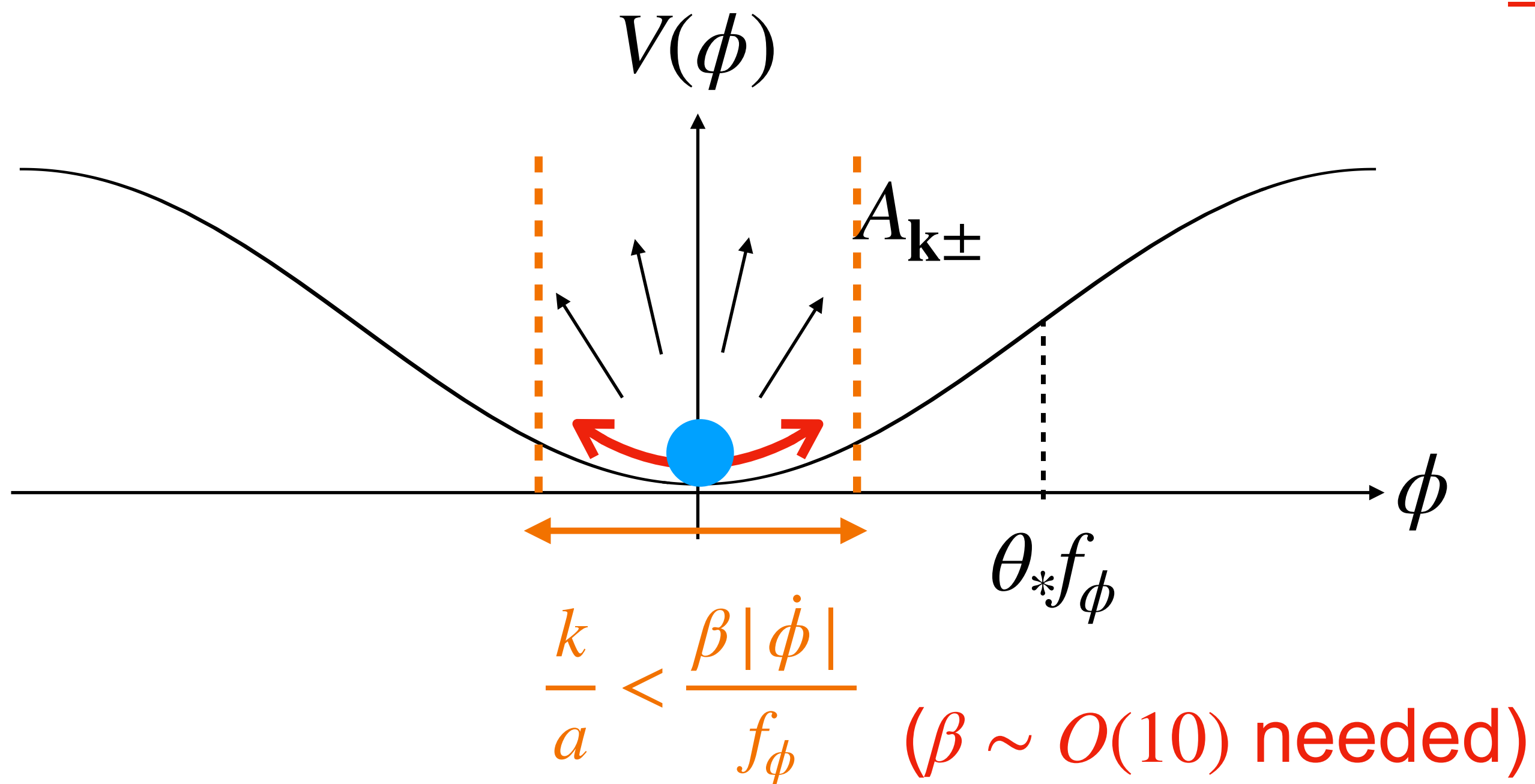


# Equation of motion for massless dark photon

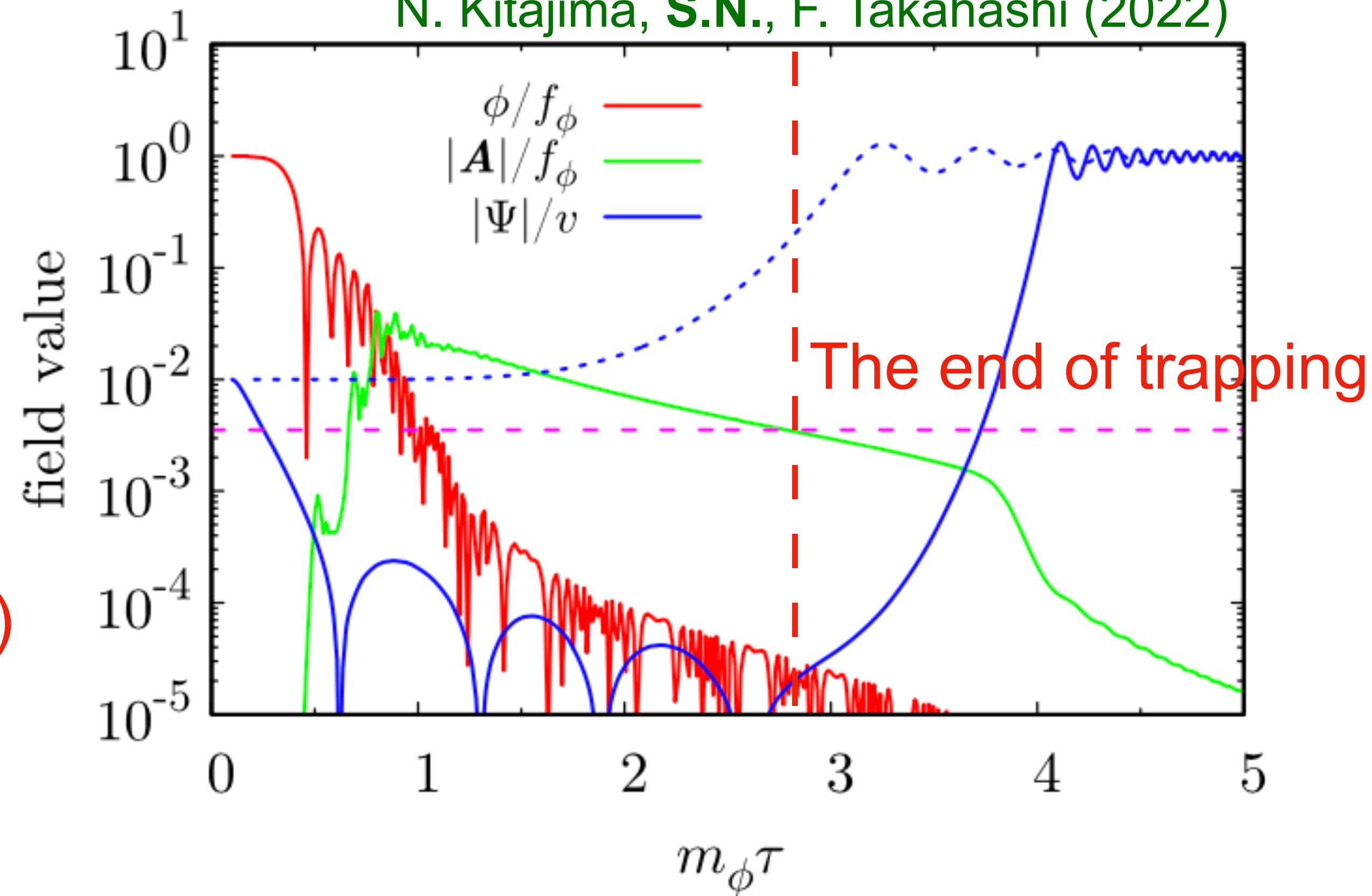
$$\ddot{A}_{\mathbf{k},\pm} + H\dot{A}_{\mathbf{k},\pm} + \frac{k}{a} \left( \frac{k}{a} \mp \frac{\beta\dot{\phi}}{f_\phi} \right) A_{\mathbf{k},\pm} = 0$$

$$\mathcal{L}_{\text{int}} = -\frac{\beta}{4f_\phi} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

N. Kitajima, S.N., F. Takahashi (2022)



Tachyonic production

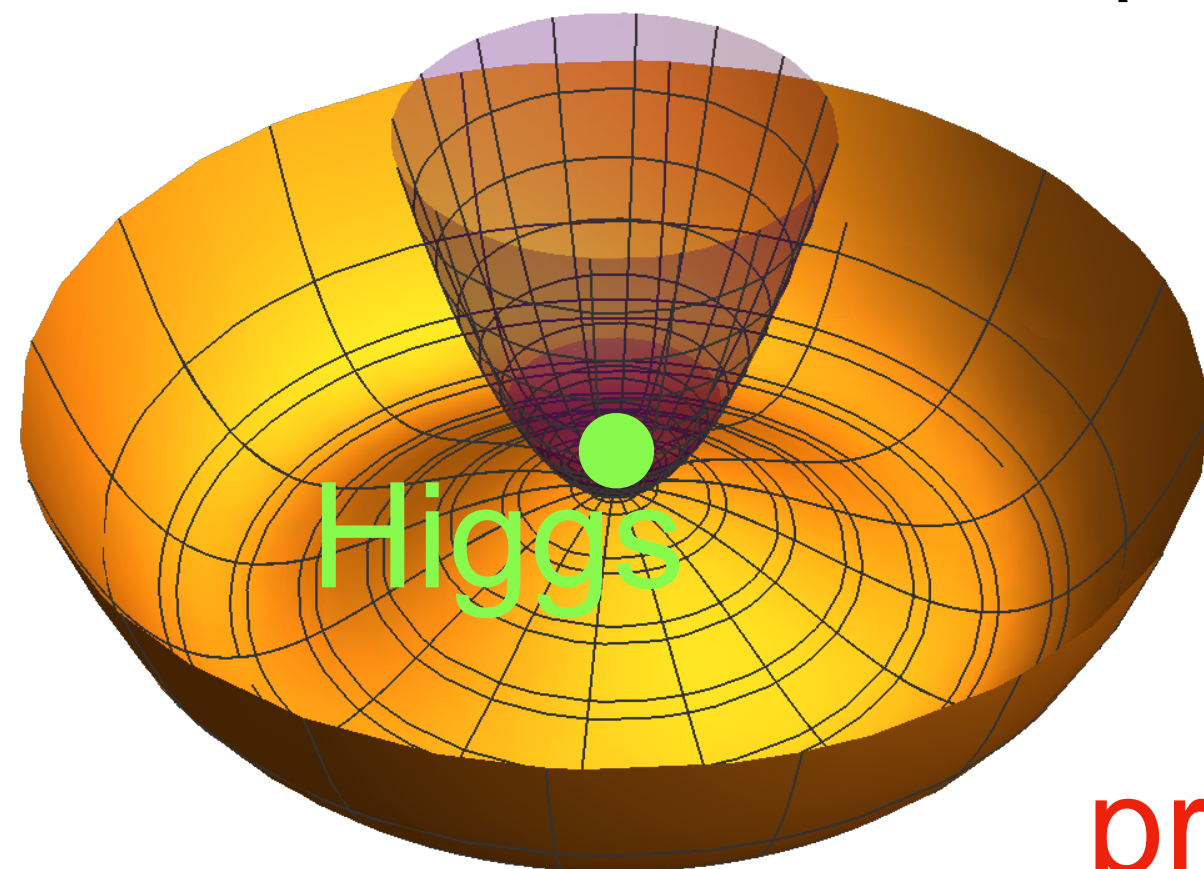


Non-thermal trapping 12/15

# Summary of our scenario

1. Dark photons are produced by tachyonic instability  
(Note: the axion becomes dark matter at the same time).
2. Dark Higgs is non-thermally trapped until  $z = z_c \simeq 5000$ .
3. It promptly decays into dark photons.

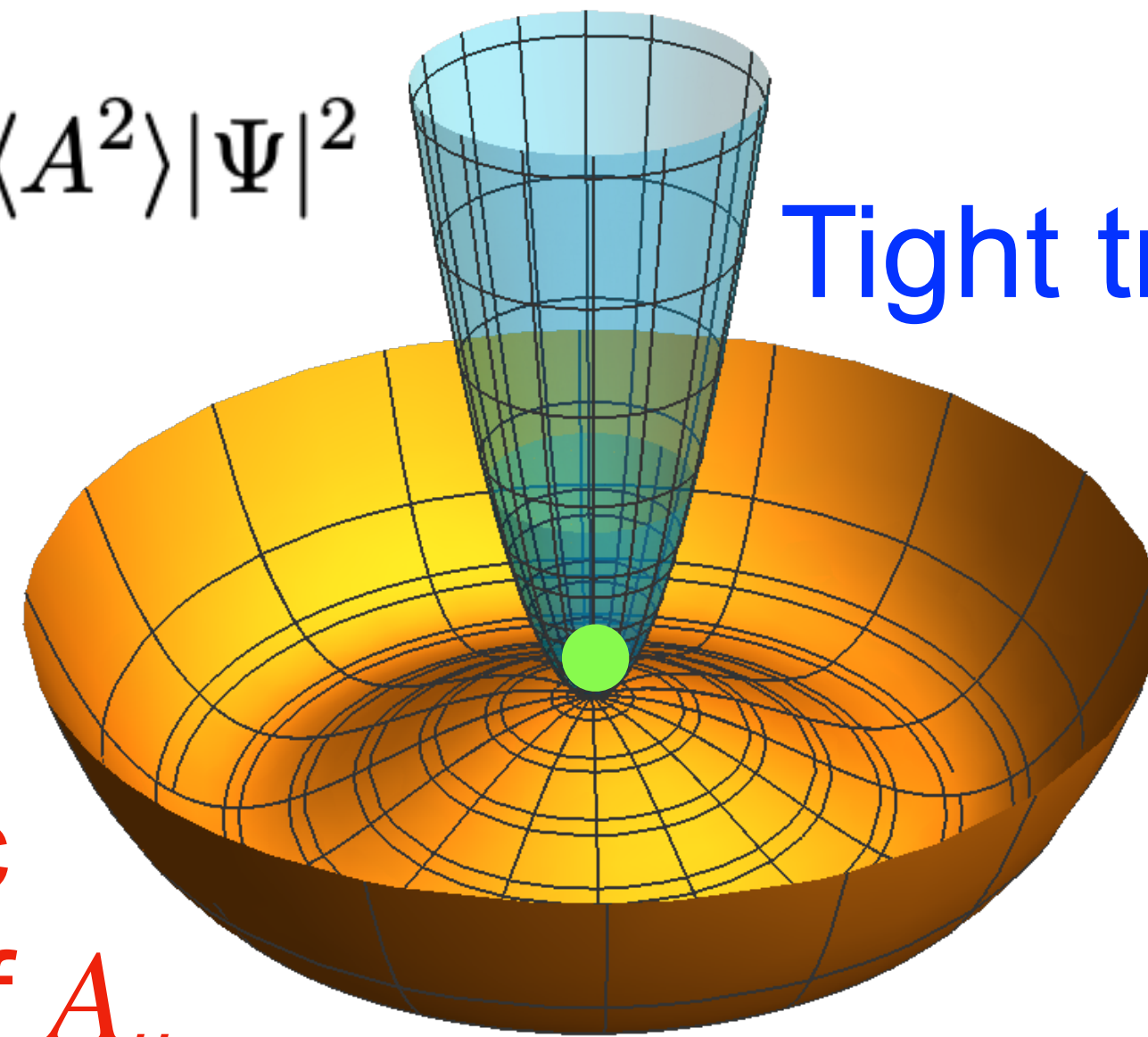
e.g. Higgs portal  
or non-minimal coupling



Initially trapped  
regime

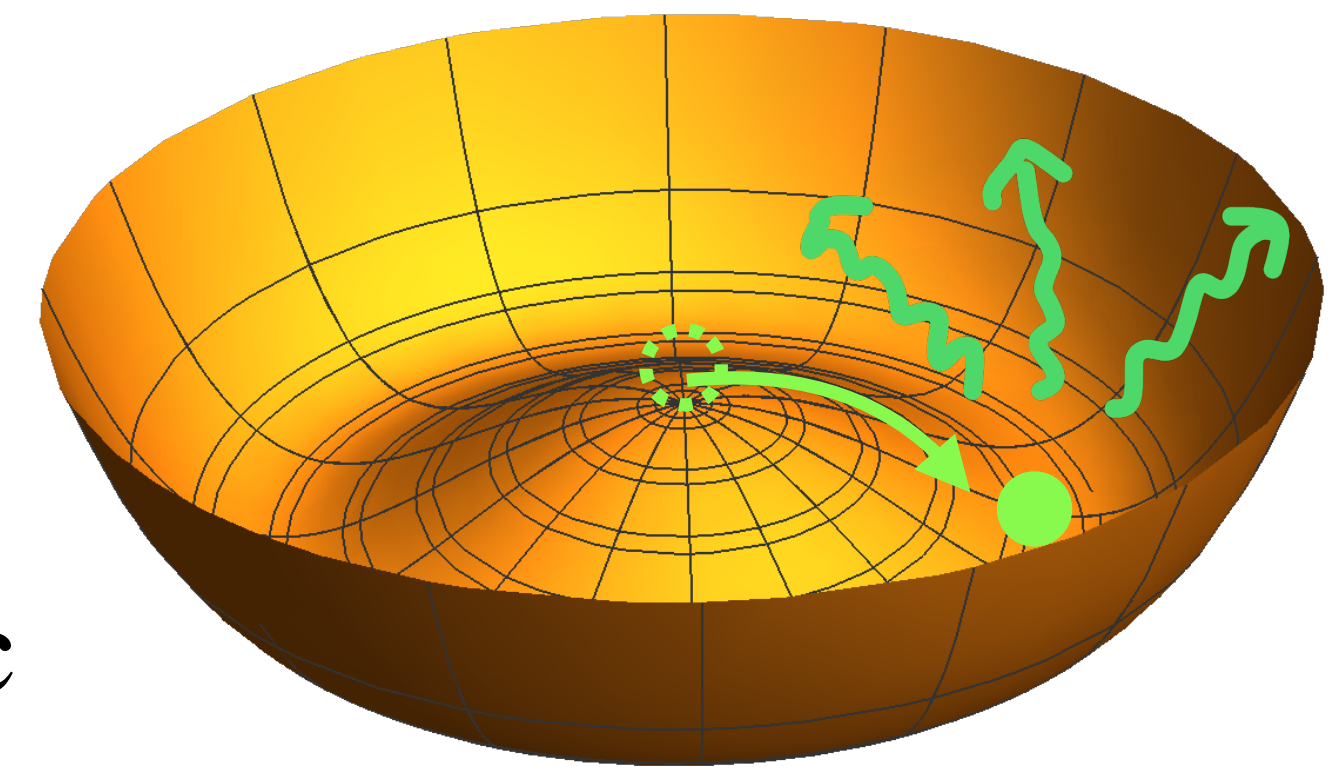
Tachyonic  
production of  $A_\mu$

$$e^2 \langle A^2 \rangle |\Psi|^2$$



Non-thermally  
trapped EDE

$$z = z_c$$



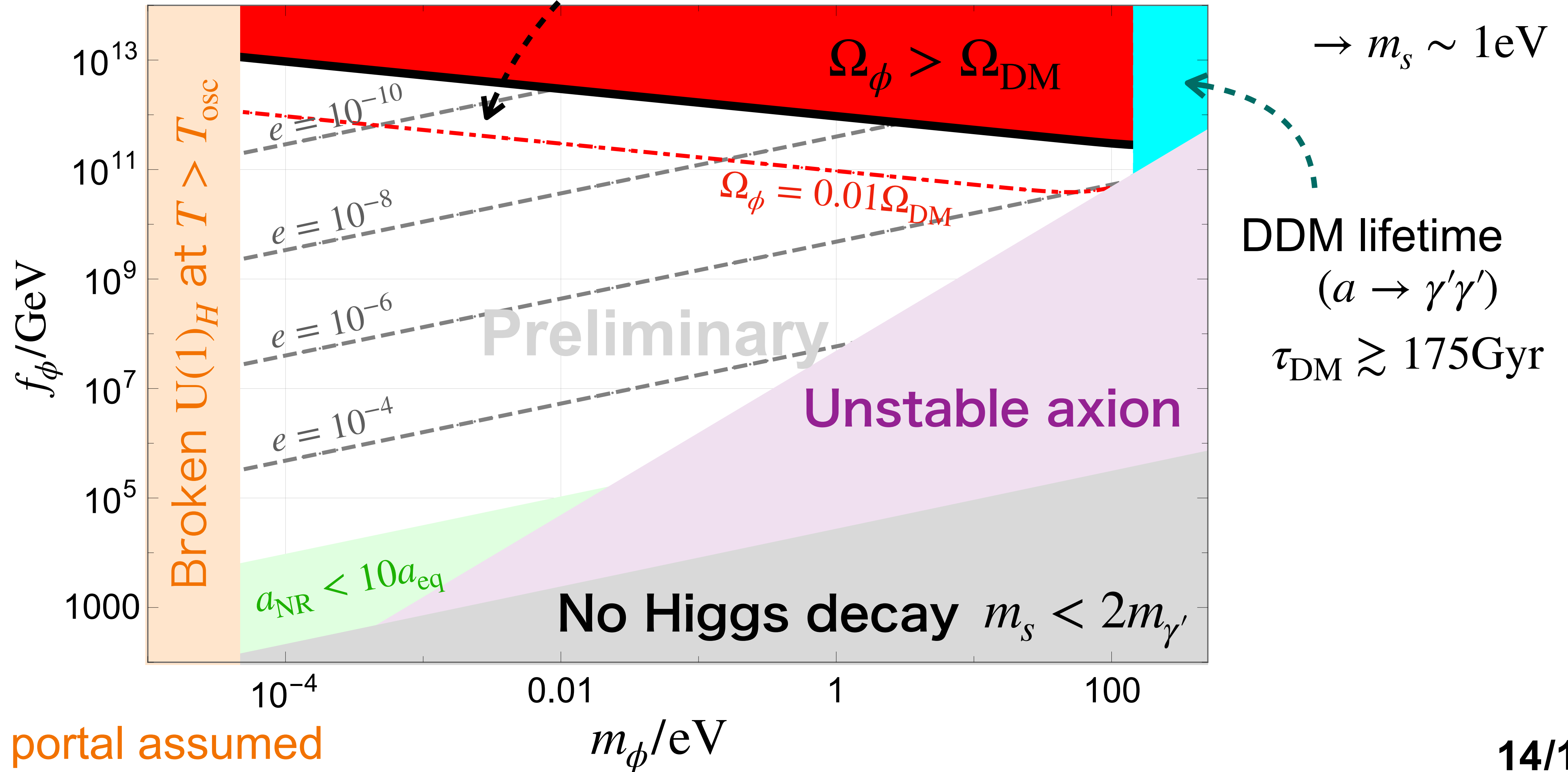
Decay into  
dark photons 13/15

# 3. Allowed region

Some tuning of  $e$  is needed for successful EDE and DM.

$$\theta_* = 1, \beta = 25, \lambda = 1, f_{\text{EDE}} = 0.05, z_c = 5000$$

$$\rightarrow m_s \sim 1\text{eV}$$



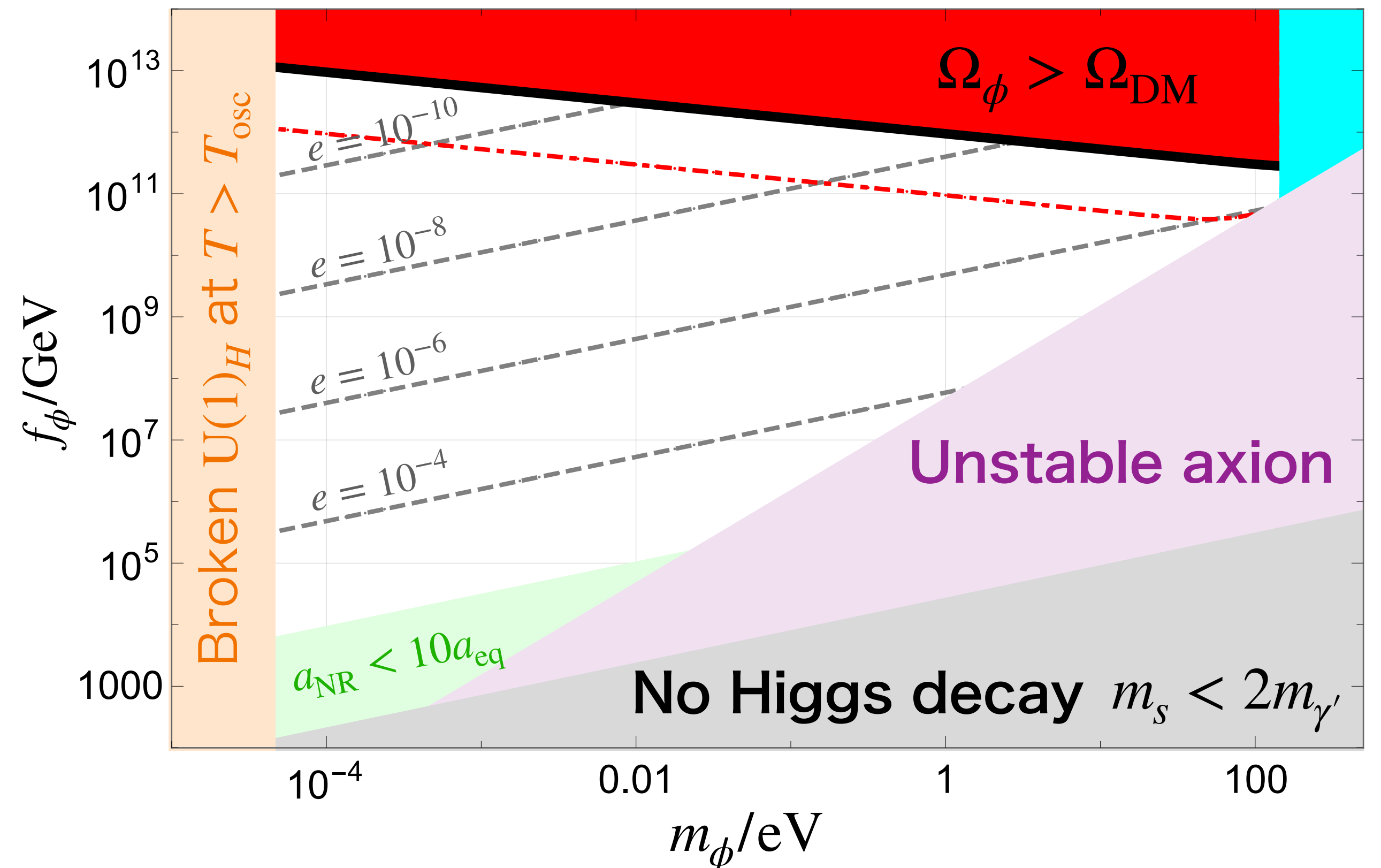
Higgs portal assumed

# Summary

- $H_0$  tension is being revealed by recent development of observation. We propose a new EDE model as a possible solution.

- The Higgs EDE can be triggered by axion dark matter which is strongly coupled to dark photons.

- $H_0$  tension can be alleviated in a vast parameter region (some tuning of  $e$  for successful EDE and DM).



**Back Up**



# Initial trapping regime

Enough tachyonic instability requires the absence of dark photon mass.

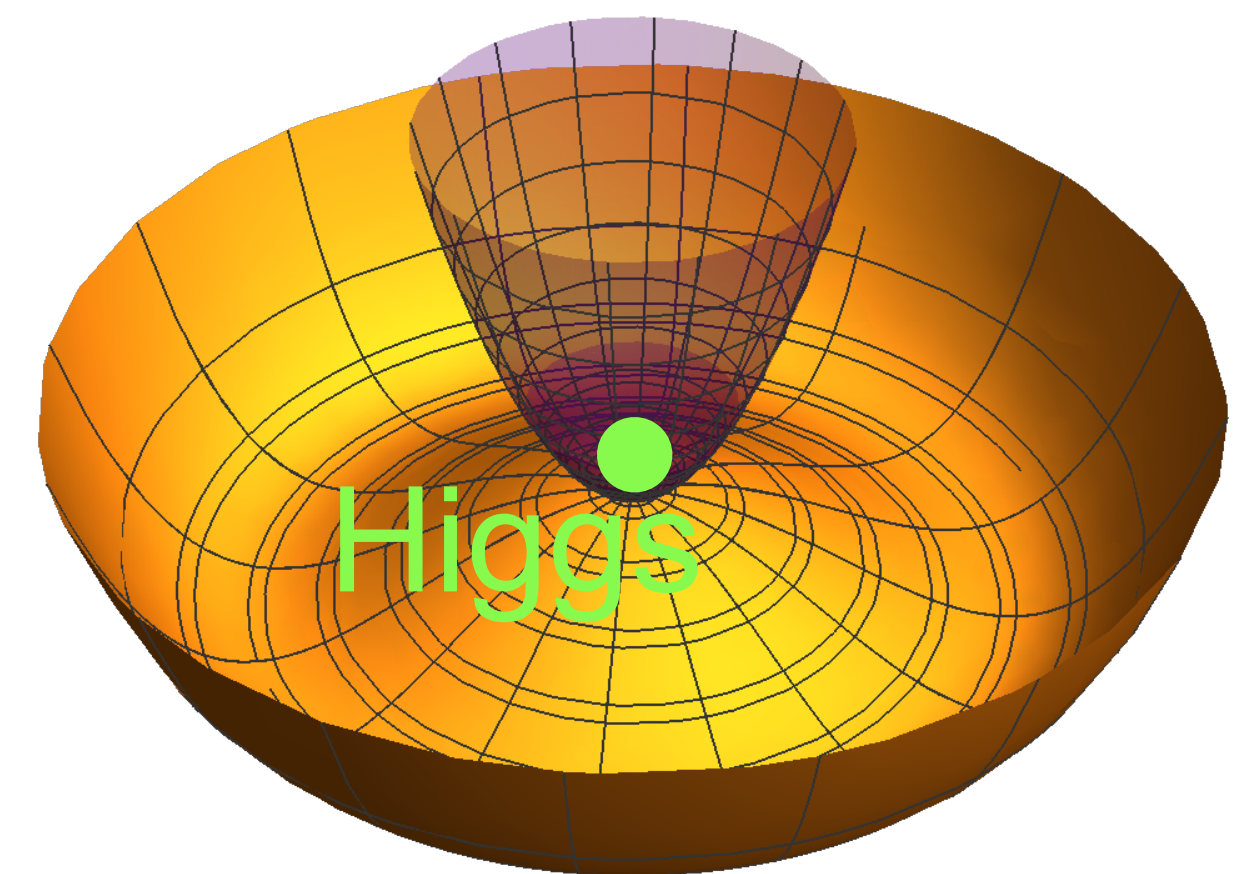
e.g. Higgs portal coupling

$$L = -\lambda_{H\Psi} H H^\dagger |\Psi|^2$$

The dark Higgs acquires thermal mass,  $m_{\text{th}} \sim \sqrt{\lambda_{H\Psi} T_{\text{osc}}}$  at the axion oscillation.

$$V_{\text{eff}} = \frac{\lambda}{4} v^4 + \underbrace{\left( m_{\text{th}}^2 - \frac{\lambda}{2} v^2 \right)}_{> 0} |\Psi|^2 + \dots$$

We can also consider non-minimal coupling to gravity, but it is somewhat weak.



# Sensitivity for axion search experiment

Assuming axion-photon coupling

For  $\lambda = 1, \beta = 25, \theta_* = 1$

