

New Horizon of Particle Physics Explored by Higgs Physics

Kei Yagyu (Osaka U.)



2021, 14th Dec., U. of Toyama

Toyama Station: Before → After

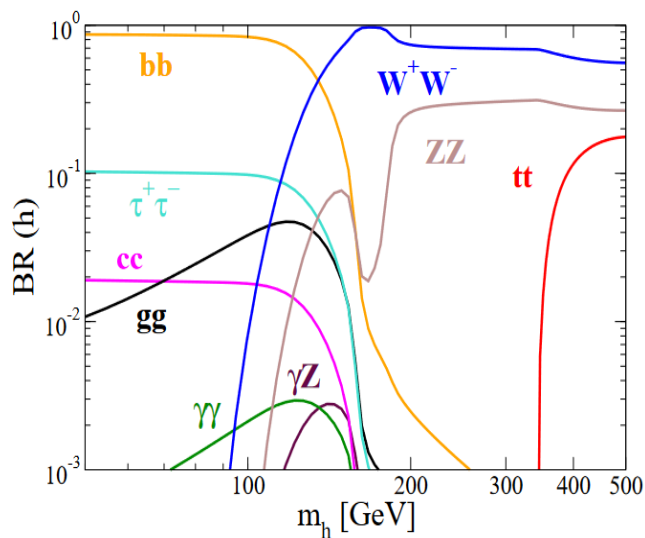


Particle Physics: Before → After

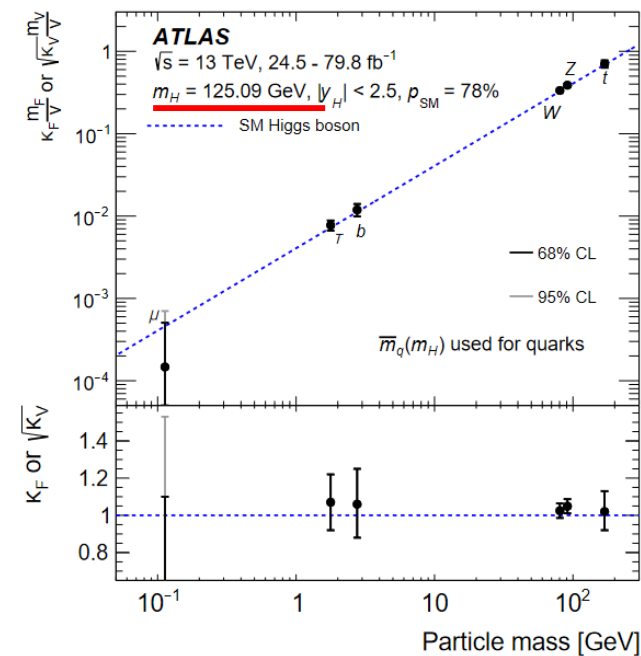
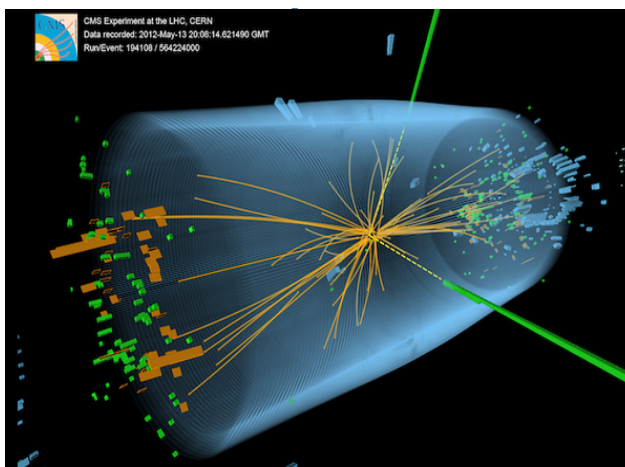
It was a hypothetical particle.

It was discovered, measured precisely.

2012, July



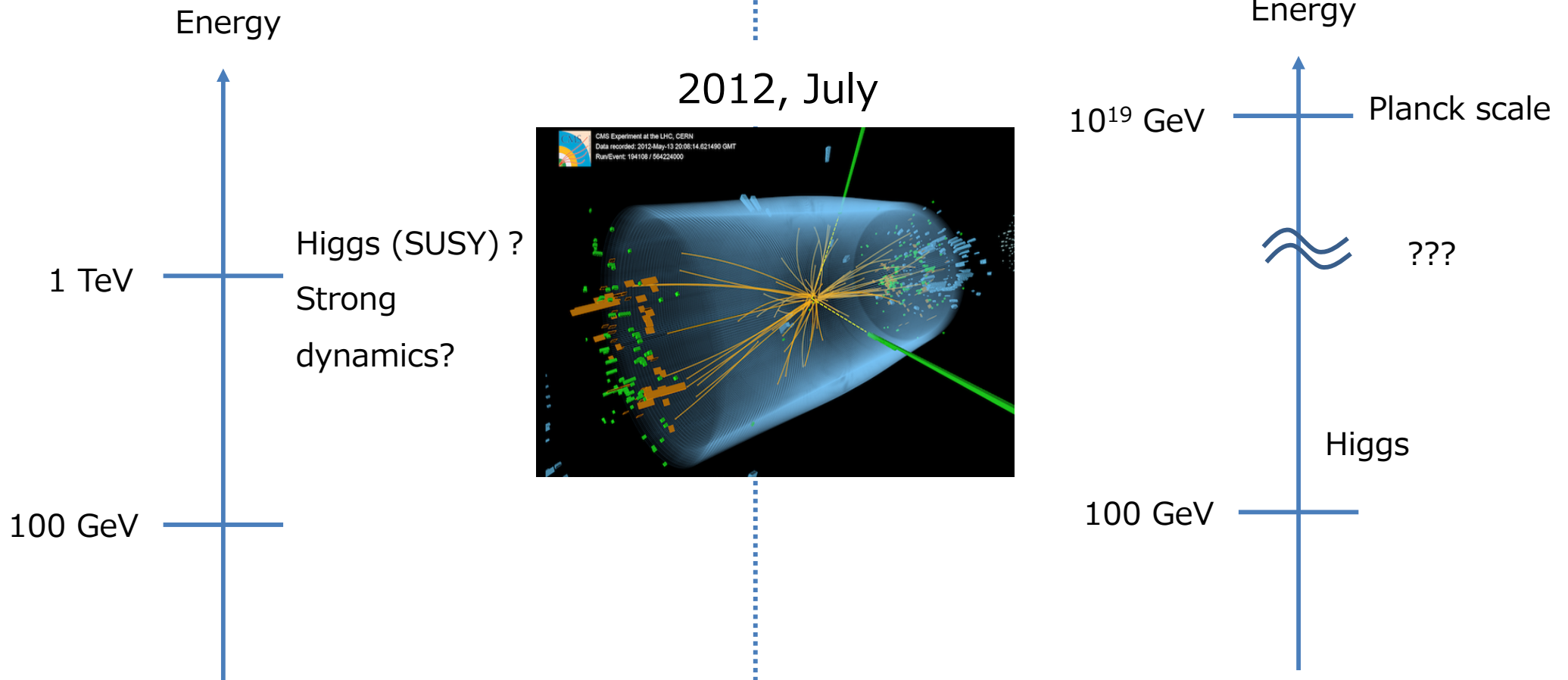
KY, PhD thesis (2012)



Particle Physics: Before → After

There was a **no-loose theorem**.

Previous concepts are broken, and now is a kind of **Sengoku-era**.



Contents

I. What is the Higgs boson?

II. Why is Higgs Physics important?

III. How can we determine the Higgs sector?

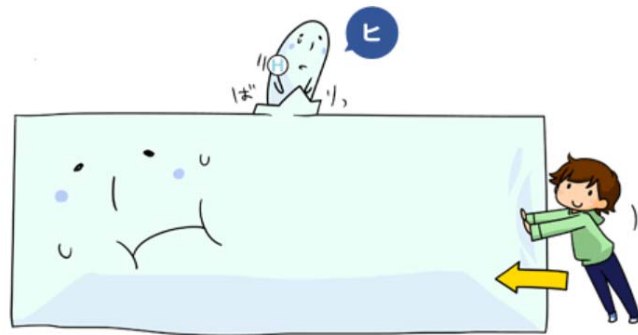
- Bottom-up approach
- Top-down approach

IV. Summary

The Higgs boson

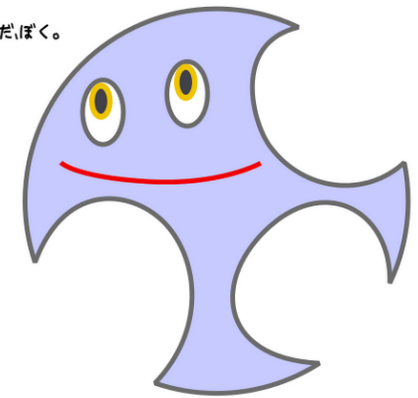


"Higgs-Tan"



"Higgs-Kun"

それでも強く
生きていくんだぼく。



"Higgs-Doll"

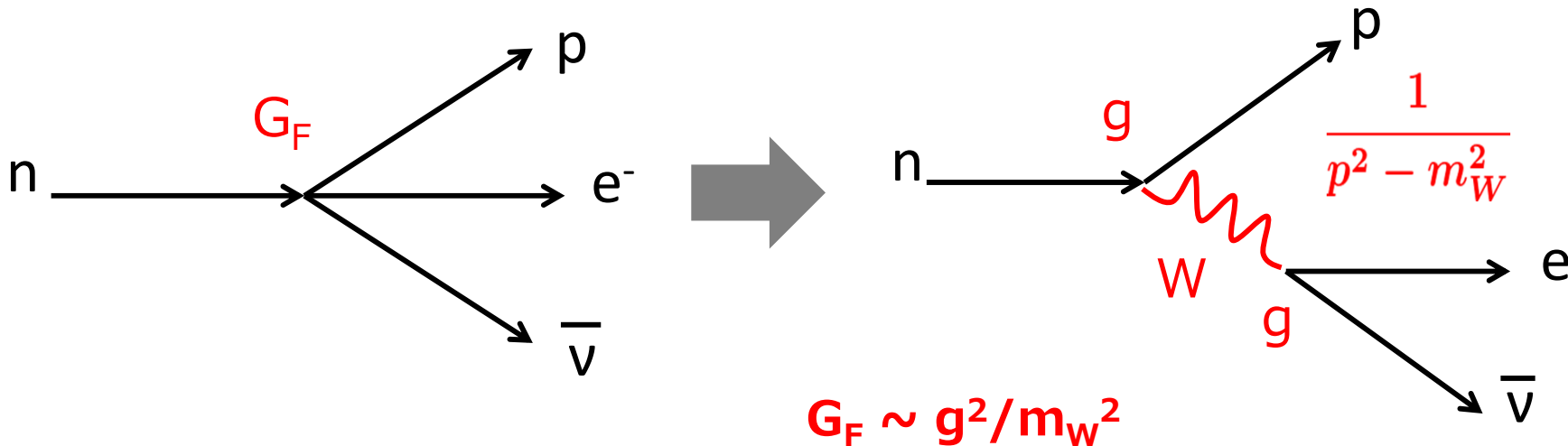
Weak force

1933: Fermi's theory for beta decays

G_F : Fermi const. $\sim 10^{-5} \text{ GeV}^{-2}$



E. Fermi



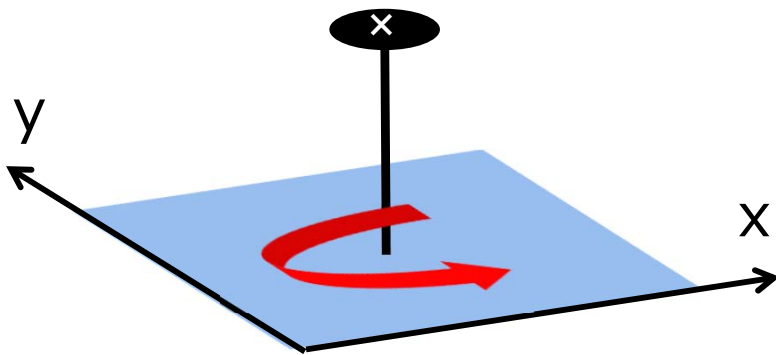
$$m_W = \mathcal{O}(100) \text{ GeV}$$

Gauge boson masses are forbidden by the gauge principle.

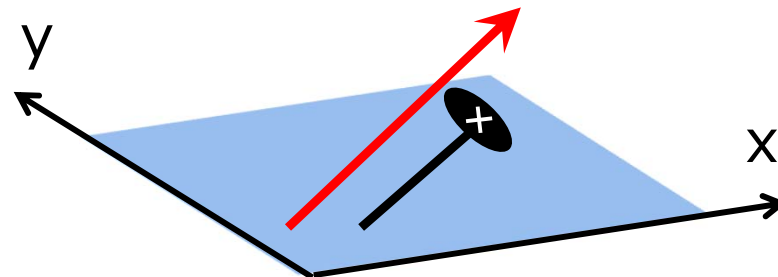
Spontaneous symmetry breaking



Y. Nambu



Unstable, but symmetric under rotation



Stable, but not symmetric under rotation

Gauge boson masses are generated by spontaneous symmetry breaking.

Scalar QED (Higgs model)

Particles : Comp. scalar $\Phi(x)$, gauge field $A_\mu(x)$

Symmetry : Local U(1)

$$\begin{aligned}\phi &\rightarrow e^{i\alpha(x)}\phi \\ A_\mu &\rightarrow A_\mu + \frac{1}{e}\partial_\mu\alpha(x)\end{aligned}$$



P. Higgs R. Brout F. Englert

Lagrangian :

$$\mathcal{L} = (D^\mu\phi)^*(D_\mu\phi) - V(|\phi|^2) - \frac{1}{4}F^{\mu\nu}F_{\mu\nu}$$

$$D_\mu\phi = (\partial_\mu - ieA_\mu)\phi \quad F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$$

$$m_A^2 A_\mu A^\mu \rightarrow m_A^2 (A_\mu + \partial_\mu\alpha/e)(A^\mu + \partial^\mu\alpha/e)$$

Potential :

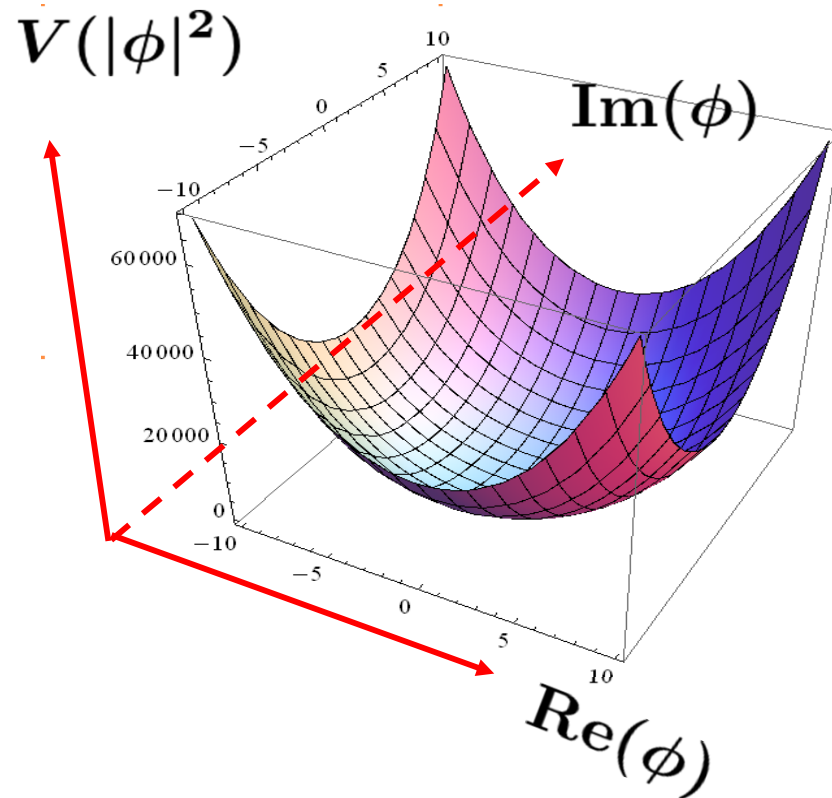
$$V = \mu^2|\phi|^2 + \lambda|\phi|^4$$

Scalar QED (Higgs model)

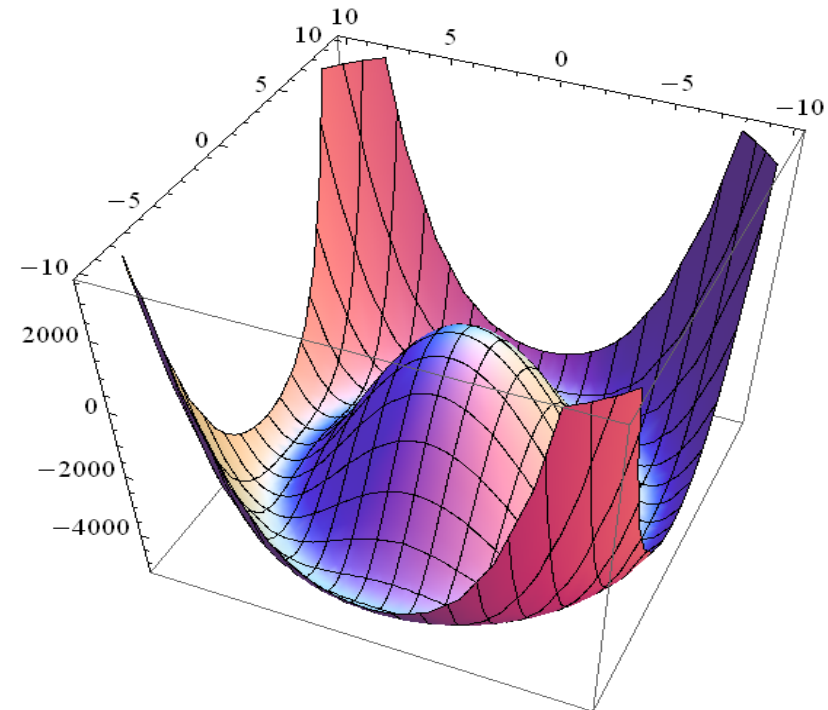
$$V = \mu^2 |\phi|^2 + \lambda |\phi|^4$$



$\mu^2 > 0$



$\mu^2 < 0$



Scalar QED (Higgs model)

□ Field expansion at around the minimum :

$$\phi(x) = [h(x) + v]e^{i\xi(x)/v}$$

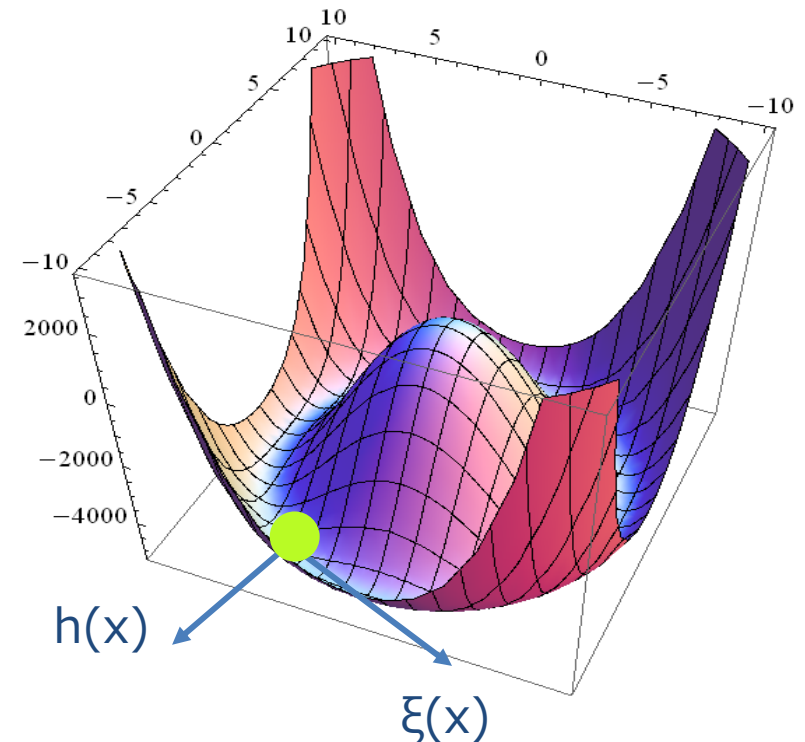
Nambu-Goldstone boson (NGB)

$$\phi \rightarrow e^{i\alpha(x)} \phi$$

$$A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha(x)$$

$$\phi \rightarrow e^{-\frac{i\xi(x)}{v}} \phi$$

$$A^\mu \rightarrow A^\mu - \frac{1}{ev} \partial^\mu \xi(x)$$



$$D_\mu \phi = (\partial_\mu - ieA_\mu) \phi$$

$$(D_\mu \phi)^* (D^\mu \phi) \rightarrow (ev)^2 A_\mu A^\mu + \dots$$

Gauge boson mass

NGB is “eaten” by the longitudinal comp. of the gauge boson,
and h (= Higgs particle) remains as a physical d.o.f.

Scalar QED (Higgs model)

Field expansion at around the minimum :

$$\phi(x) = [h(x) + v]e^{i\xi(x)/v}$$

Nambu-Goldstone boson (NGB)

$$\phi \rightarrow e^{i\alpha(x)} \phi$$

$$A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha(x)$$

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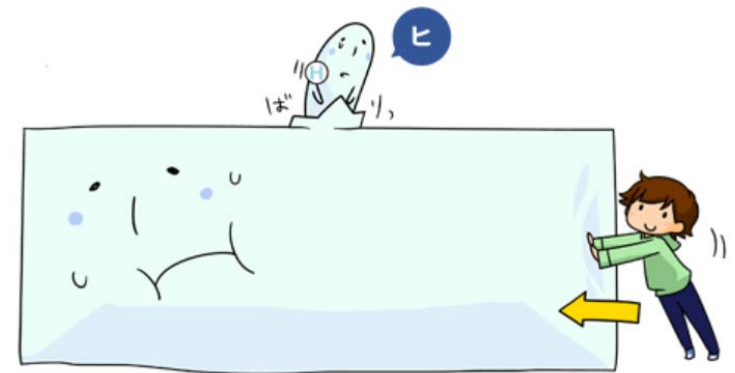
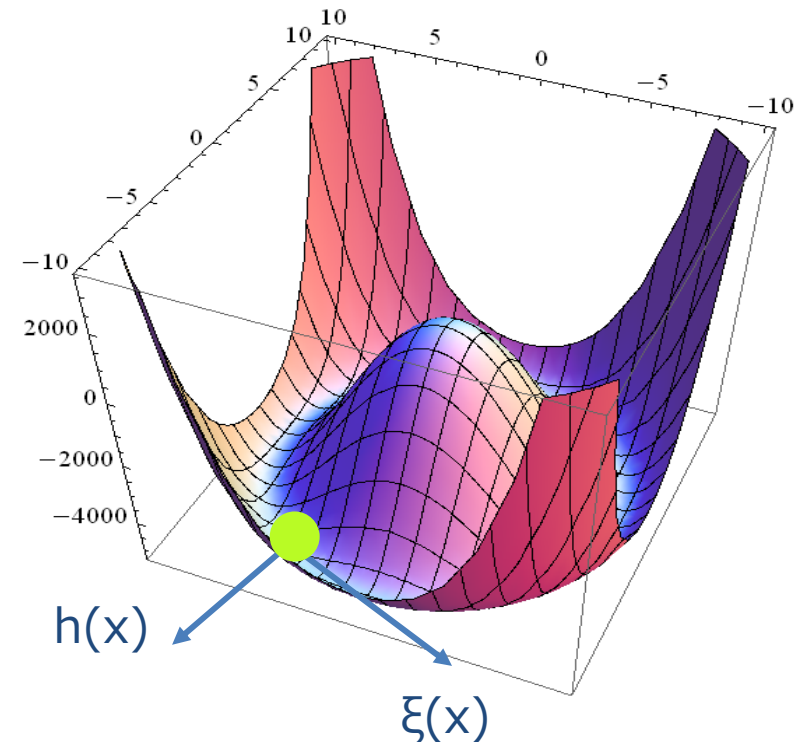
$$A^\mu \rightarrow A^\mu - \frac{1}{ev} \partial^\mu \xi(x)$$

$$D_\mu \phi = (\partial_\mu - ieA_\mu) \phi$$

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Gauge boson mass

NGB is "eaten" by the longitudinal comp.
and h (= Higgs particle) remains as a ph.



Spontaneous EWSB

Higgs field $\Phi(x)$: SU(2) doublet $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$, Gauge fields (W^+, W^0, W^-), B^0

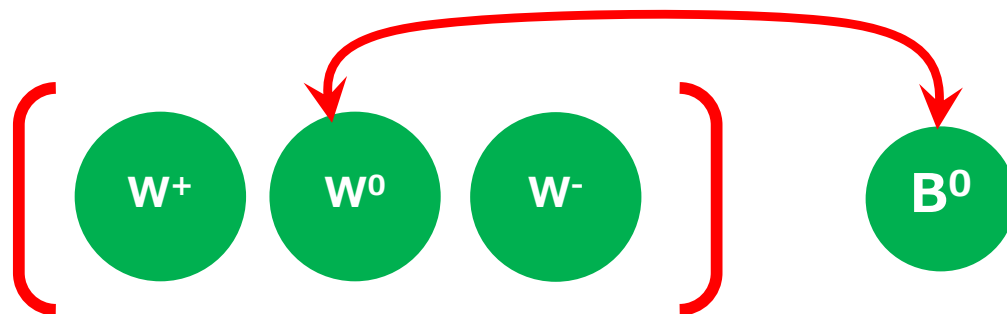
Symmetry: Local $SU(2)_I \times U(1)_Y$

$$V(\Phi) = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

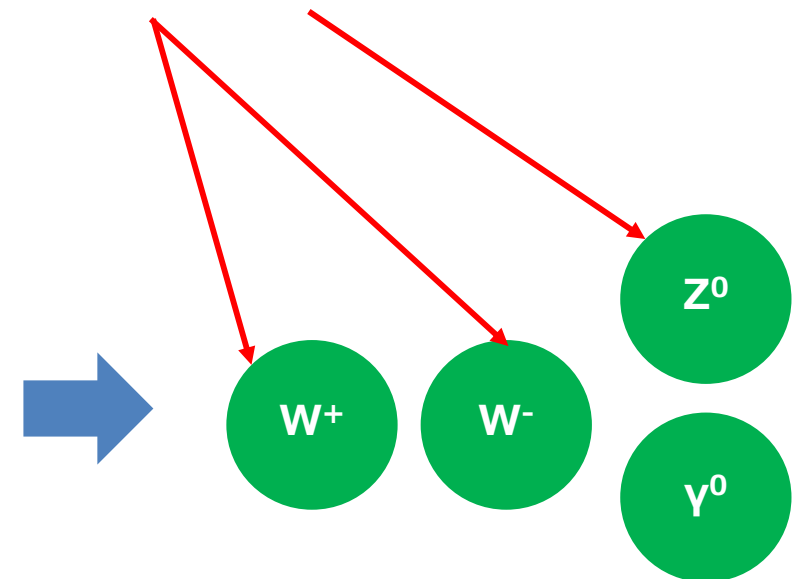
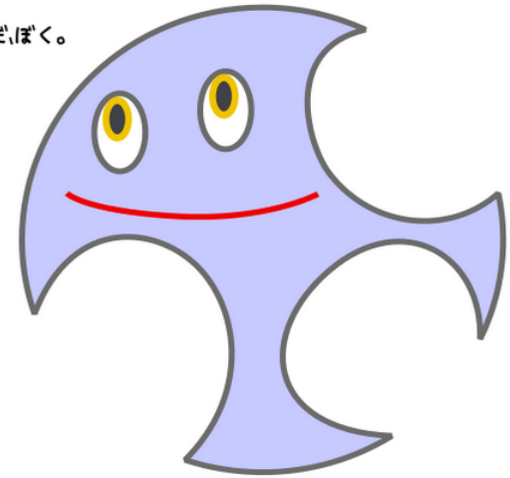
$$\langle \phi^0 \rangle = v$$

➔ $SU(2)_L \times U(1)_Y \rightarrow U(1)_{em}$

mixing

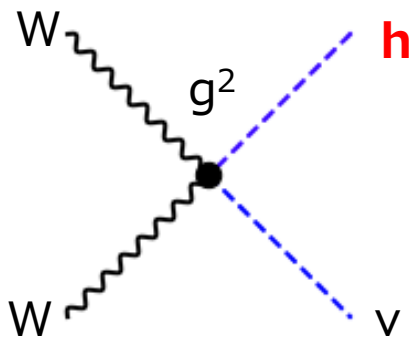
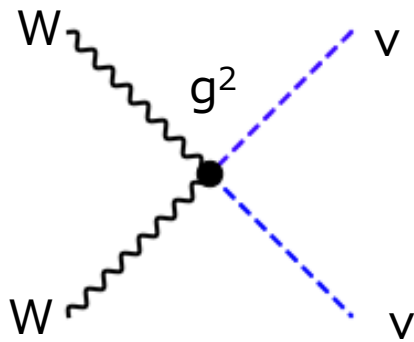


それでも強く
生きていくんだ、ぼく。



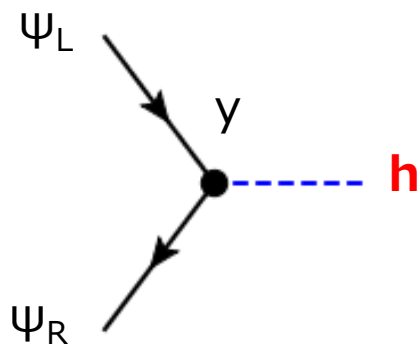
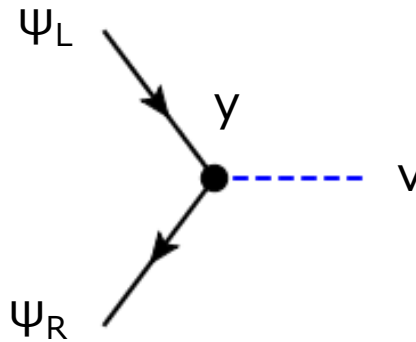
Origin of Mass

Kin. term $|D_\mu \Phi|^2$



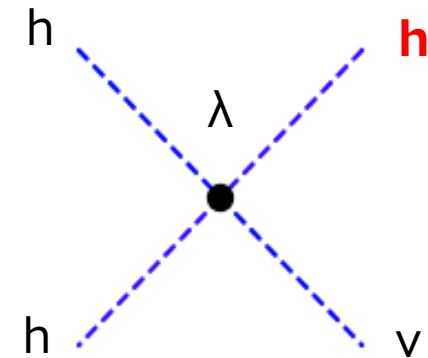
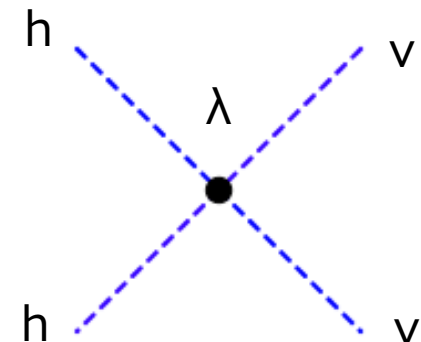
$$g_{hWW} \propto \frac{m_W^2}{v}$$

Yukawa $y \bar{\Psi}_L \Phi \Psi_R$



$$g_{hff} \propto \frac{m_f}{v}$$

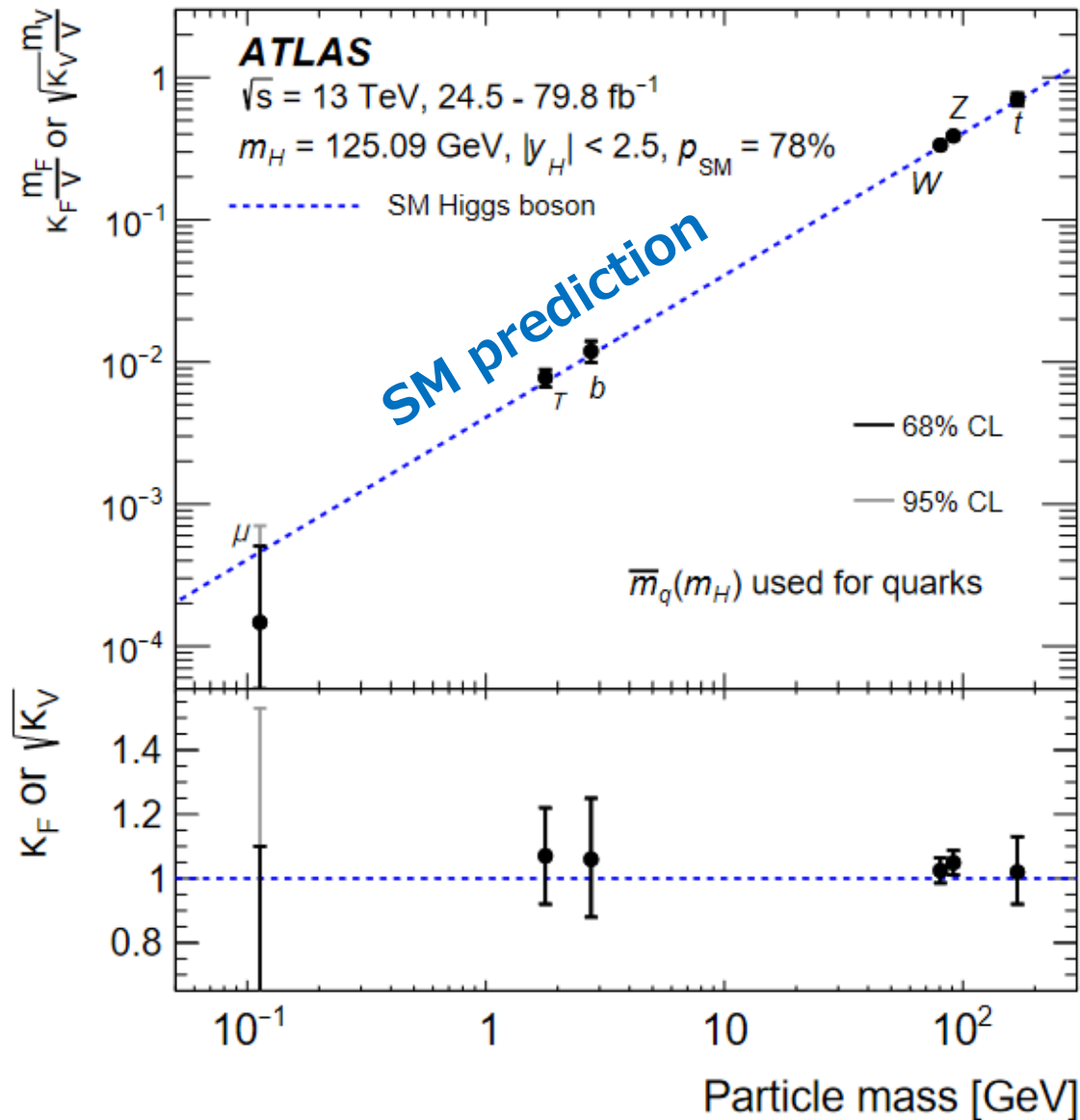
Potential $V(\Phi)$



$$g_{hhh} \propto \frac{m_h^2}{v}$$

$$\kappa_X \equiv \frac{g_{hXX}^{\text{exp}}}{g_{hXX}^{\text{SM}}}$$

ATLAS-CONF-2021-053



Parameter	(a) $B_i = B_u = 0$
κ_Z	0.99 ± 0.06
κ_W	1.06 ± 0.06
κ_b	0.87 ± 0.11
κ_t	0.92 ± 0.10
κ_μ	$1.07^{+0.25}_{-0.30}$
κ_τ	0.92 ± 0.07
κ_γ	1.04 ± 0.06
$\kappa_{Z\gamma}$	$1.37^{+0.31}_{-0.37}$
κ_g	$0.92^{+0.07}_{-0.06}$
B_i	-
B_u	-

Great success of the standard model.

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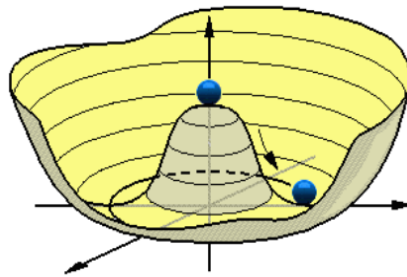
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Higgs Problems

- So far, the SM Higgs sector can successfully describe current experimental data.

$$V(\Phi) = -\mu^2|\Phi|^2 + \lambda|\Phi|^4$$

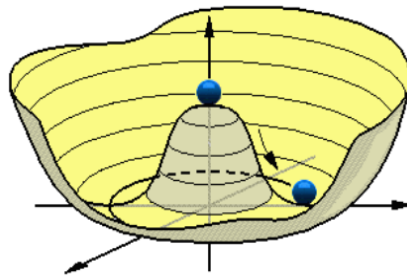


Higgs Problems

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Origin of the negative mass term

$$V(\Phi) = -\mu^2|\Phi|^2 + \lambda|\Phi|^4$$



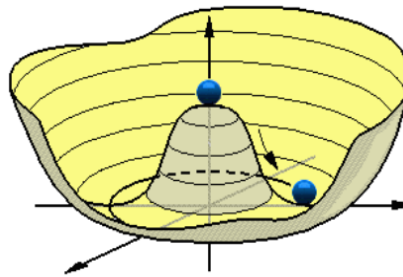
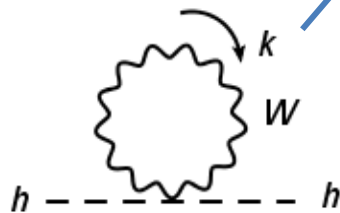
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$$V(\Phi) = -\mu^2|\Phi|^2 + \lambda|\Phi|^4$$

Quadratic divergence



$$m_h^2 \sim \frac{\Lambda^2}{16\pi^2} \gg (125 \text{ GeV})^2$$

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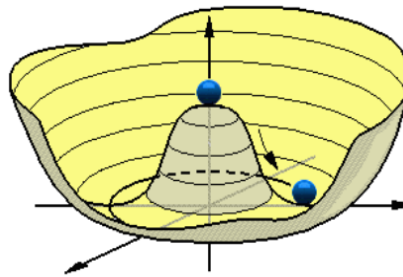
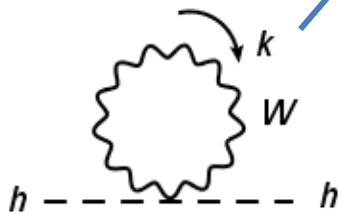
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One doublet?

Any other representations?

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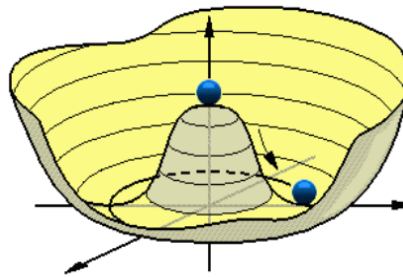
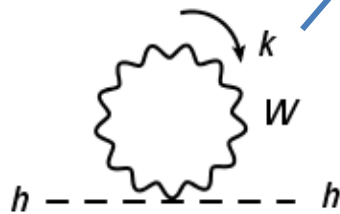
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Yukawa coupling?

$$\frac{y_e}{y_t} \sim 10^{-5}$$

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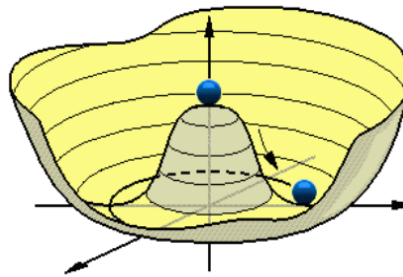
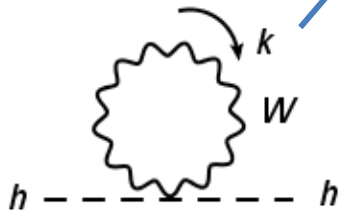
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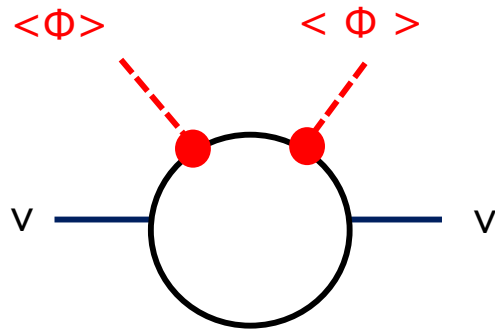
$$m_h^2 \sim \frac{\Lambda^2}{16\pi^2} \gg (125 \text{ GeV})^2$$

The Higgs sector is the center of the problem in the SM.
There should be new dynamics behind the Higgs sector.

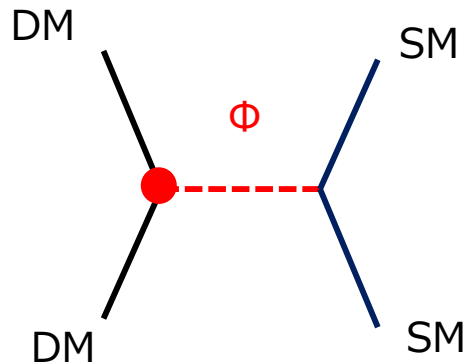
Higgs sector and BSM phenomena

- Phenomena cannot be explained in the SM.

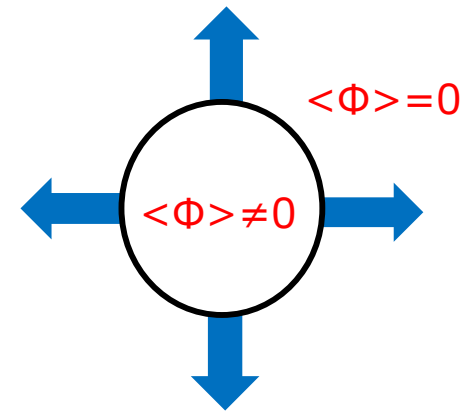
Neutrino mass



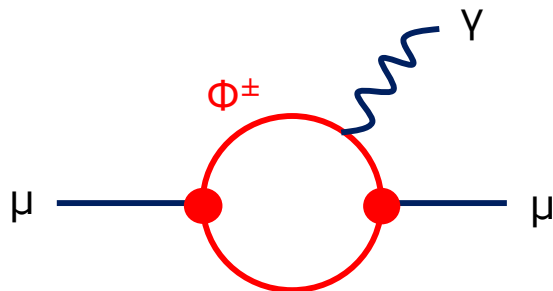
Dark matter



Baryon asymmetry of Universe

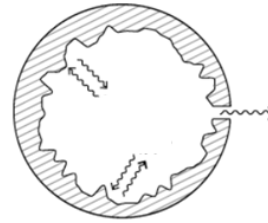


Flavor anomaly $(g-2)_\mu$



Higgs physics can strongly be related to BSMs.

Paradigm Shift in Early 20th Century



Black-body Radiation

Classical Theory

- Newton Dynamics
- Maxwell Electromagnetism

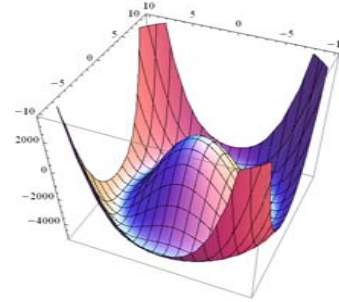
Quantum Theory

- Nuclear Physics
- Particle Physics, ...

Planck's Law

$$u(\nu, T) = \frac{8\pi\nu^2}{c^3} \frac{h\nu}{\exp(h\nu/k_B T) - 1}$$

Paradigm Shift in Early 21st Century



Higgs Physics

Standard Model

- Gauge Principle
- Higgs Mechanism

New Physics

- New dynamics
- New symmetries
- Unifications, ...

$$V(\Phi) = -\mu^2 |\Phi|^2 + \lambda |\Phi|^4$$

Higgs as a Probe of New Physics!!

Higgs as a window to New Physics

New Physics

Top-down

(Extended) Higgs sector

Bottom-up

Experimental data

High energy exp. (LHC, ...)
High intensity exp. (KEKB, Super K, ...)
Space based exp. (Planck, ...)

Theory requirements

Gauge principle
Unitarity, Vacuum stability
(Renormalizability)

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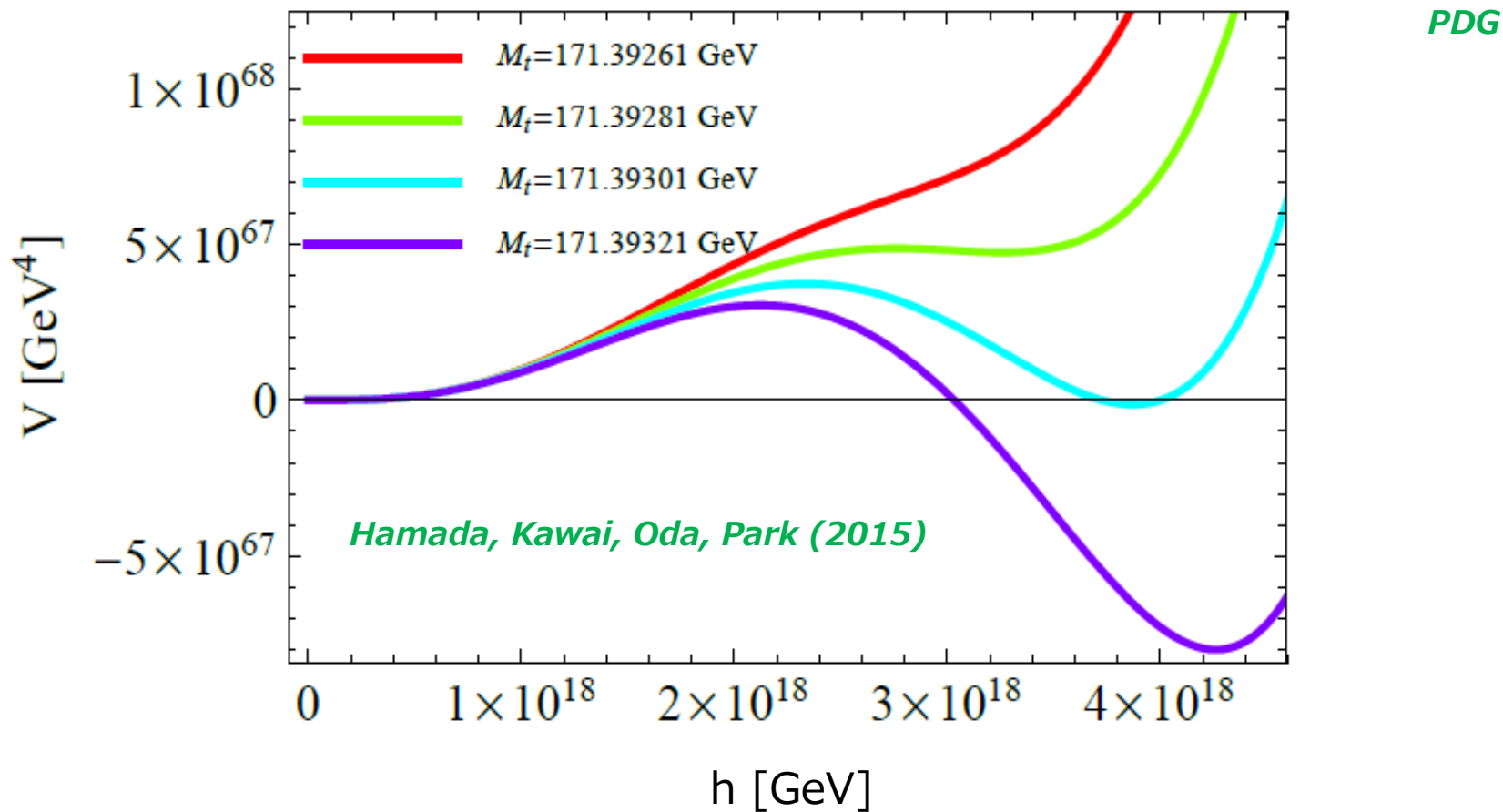
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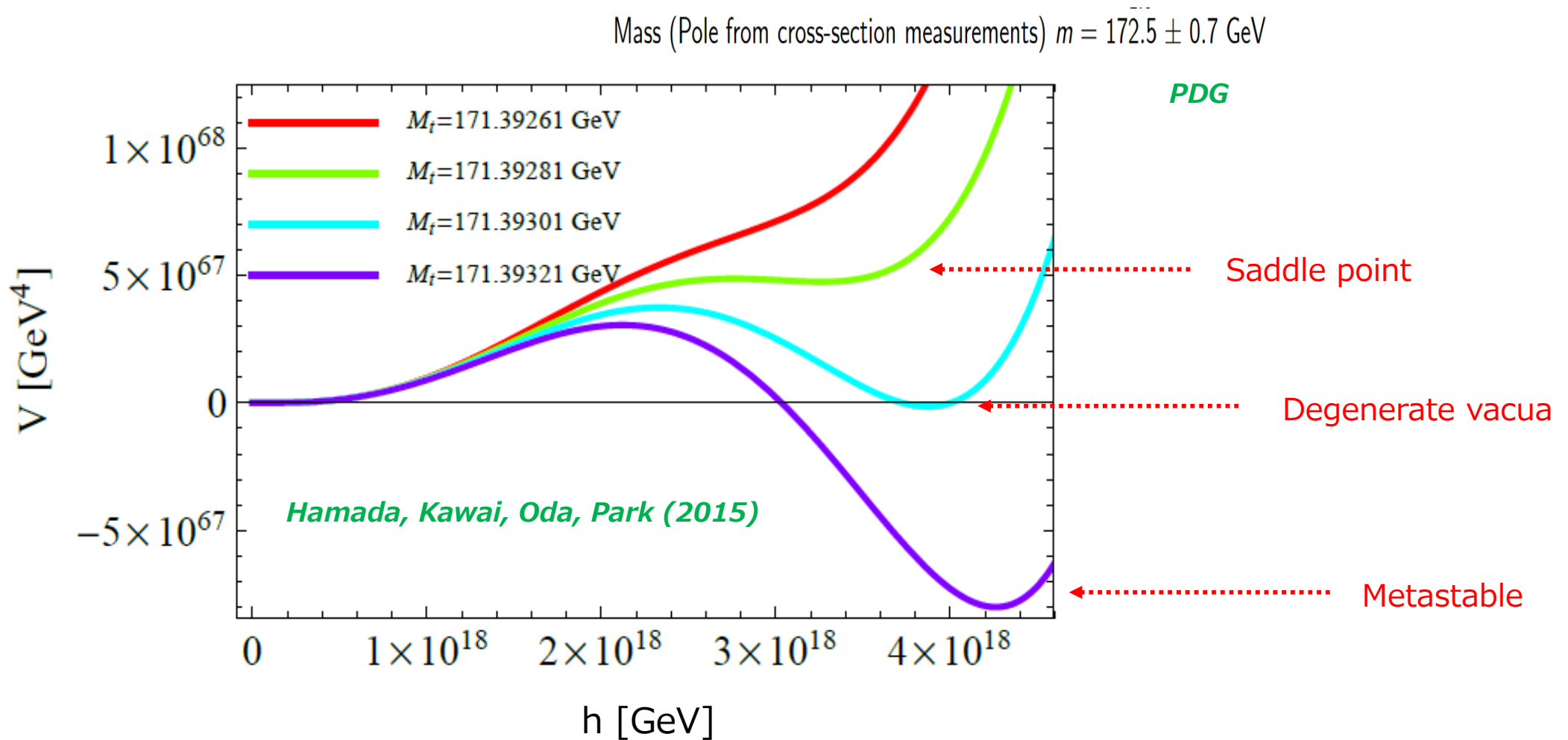
*Yuta Hamada (Harvard U.), Hikaru Kawai (National Taiwan U.),
Kiyoharu Kawana (Seoul National U.), Kin-ya Oda (Tokyo Woman's Christian U.), KY
2008.08700 [hep-ph], 2102.04617 [hep-ph] + paper in progress,*

Implication of the 125 GeV Higgs mass

Mass (Pole from cross-section measurements) $m = 172.5 \pm 0.7$ GeV



Implication of the 125 GeV Higgs mass



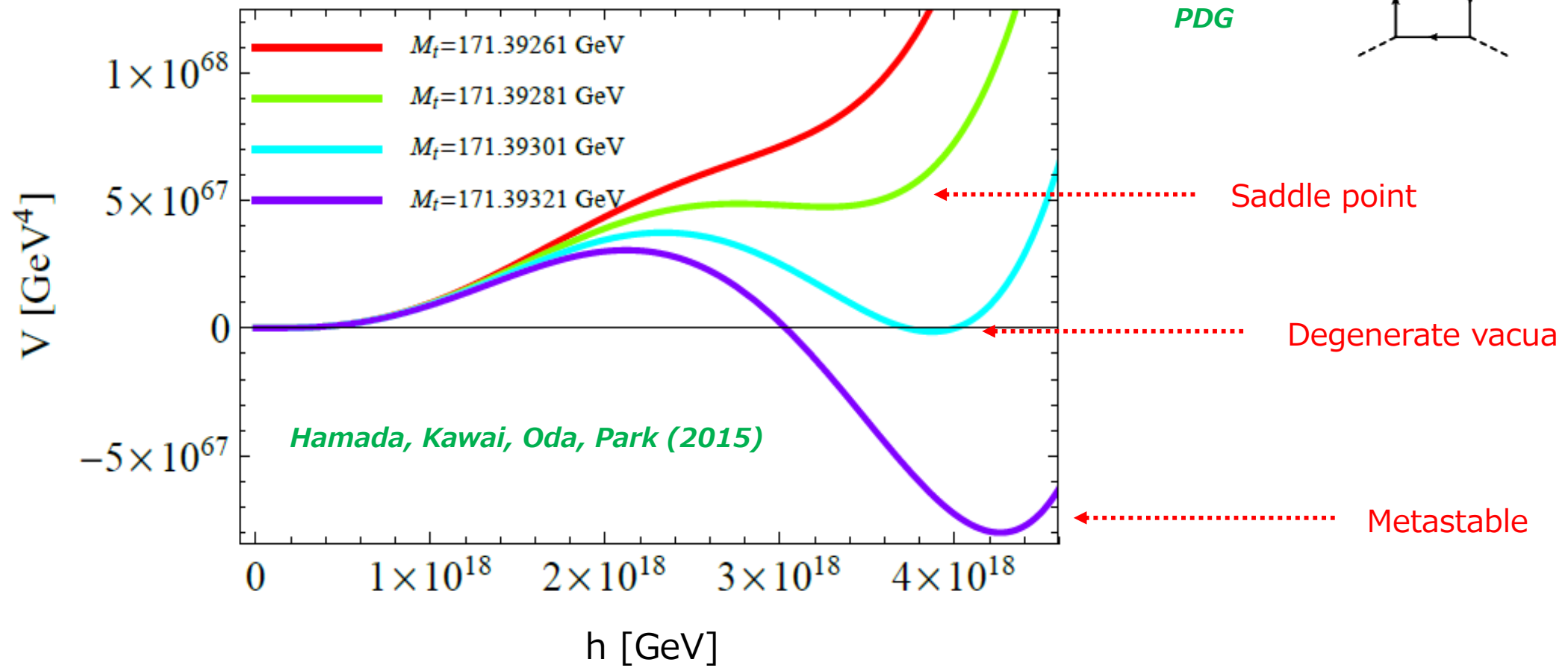
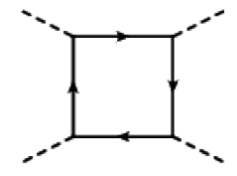
Implication of the 125 GeV Higgs mass

$$V_{\text{eff}} = \frac{\lambda_{\text{eff}}(h)}{4} h^4 \quad @ h \gg v, \mu = h$$

$$\mu \frac{d\lambda}{d\mu} = \frac{1}{16\pi^2} (24\lambda^2 - 6y_t^4 + \dots)$$



Mass (Pole from cross-section measurements) $m = 172.5 \pm 0.7 \text{ GeV}$



Froggatt & Nielsen (1995)



ELSEVIER

Physics Letters B 368 (1996) 96–102

PHYSICS LETTERS B

Standard model criticality prediction top mass 173 ± 5 GeV
and Higgs mass 135 ± 9 GeV

C.D. Froggatt^a, H.B. Nielsen^b

^a *Department of Physics and Astronomy, Glasgow University, Glasgow G12 8QQ, Scotland, UK*

^b *The Niels Bohr Institute, Blegdamsvej 17, DK-2100 Copenhagen Ø, Denmark*

Received 4 November 1995

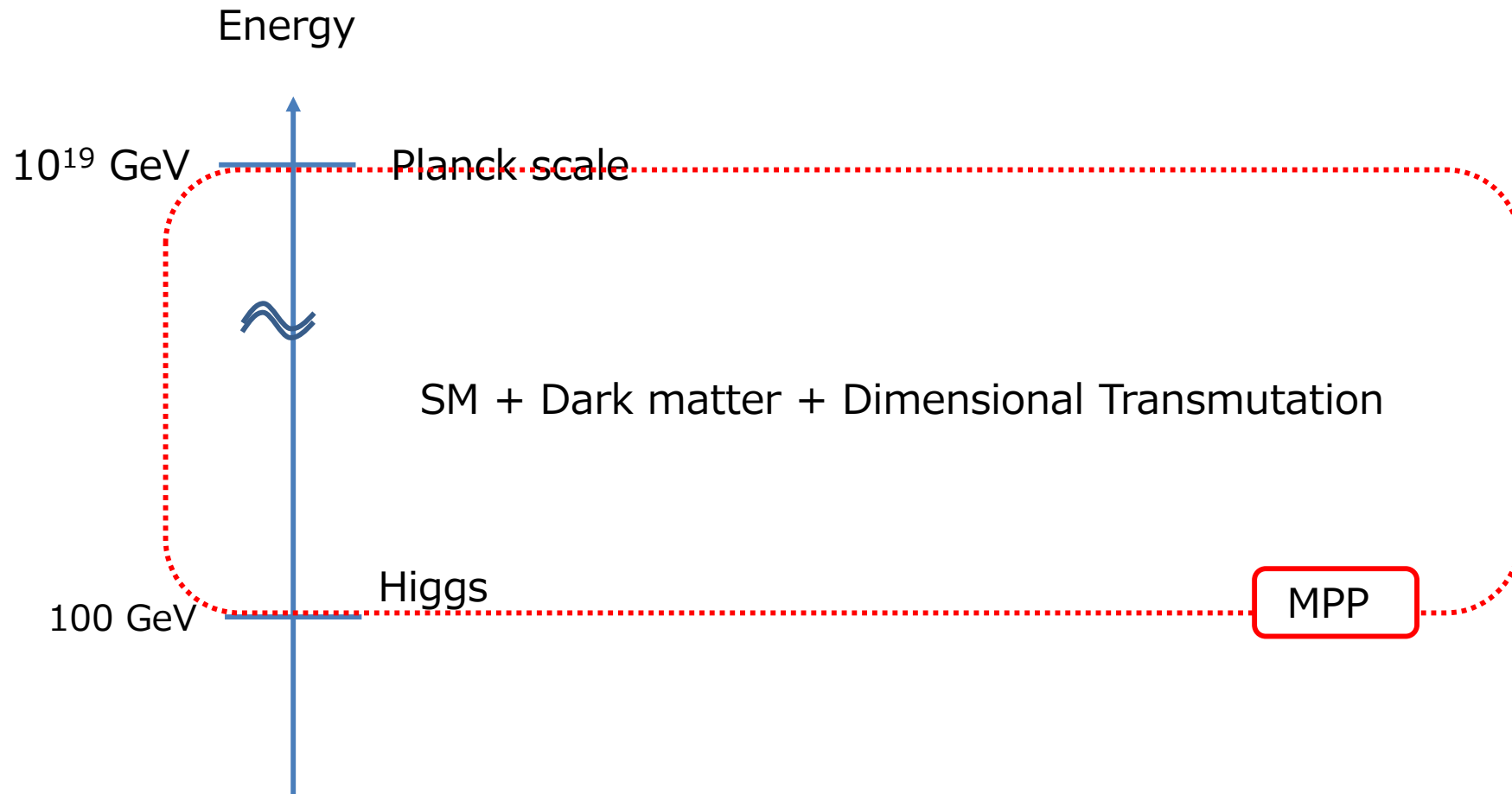
Editor: P.V. Landshoff

They predicted the Higgs mass in 1995!!

The Higgs potential is realized at a critical point?

→ Multicritical point principle (MPP)

Our Scenario



We proposed a minimal model to realize the scenario based on the MPP framework.

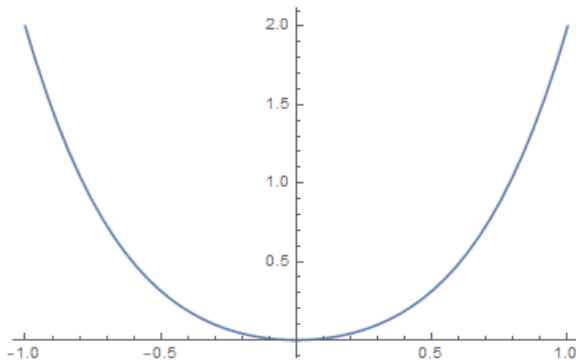
Multicritical point principle (MPP)

Froggatt, Nielsen (1995)

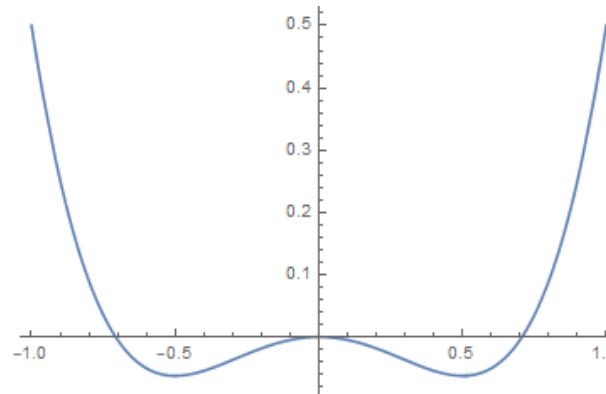
MPP: More parameters in the effective potential are tuned to a set of critical values, which is more likely to be realized by nature.

Examples: $V = \mu^2 |H|^2 + \lambda |H|^4$

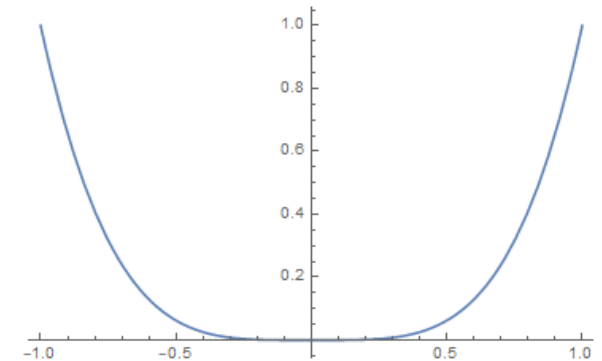
$\mu^2 > 0$: Stable



$\mu^2 < 0$: SSB



$\mu^2 = 0$: Critical point



Classical Scale Invariance (CSI)

CSI can be understood as a special realization of the MPP.

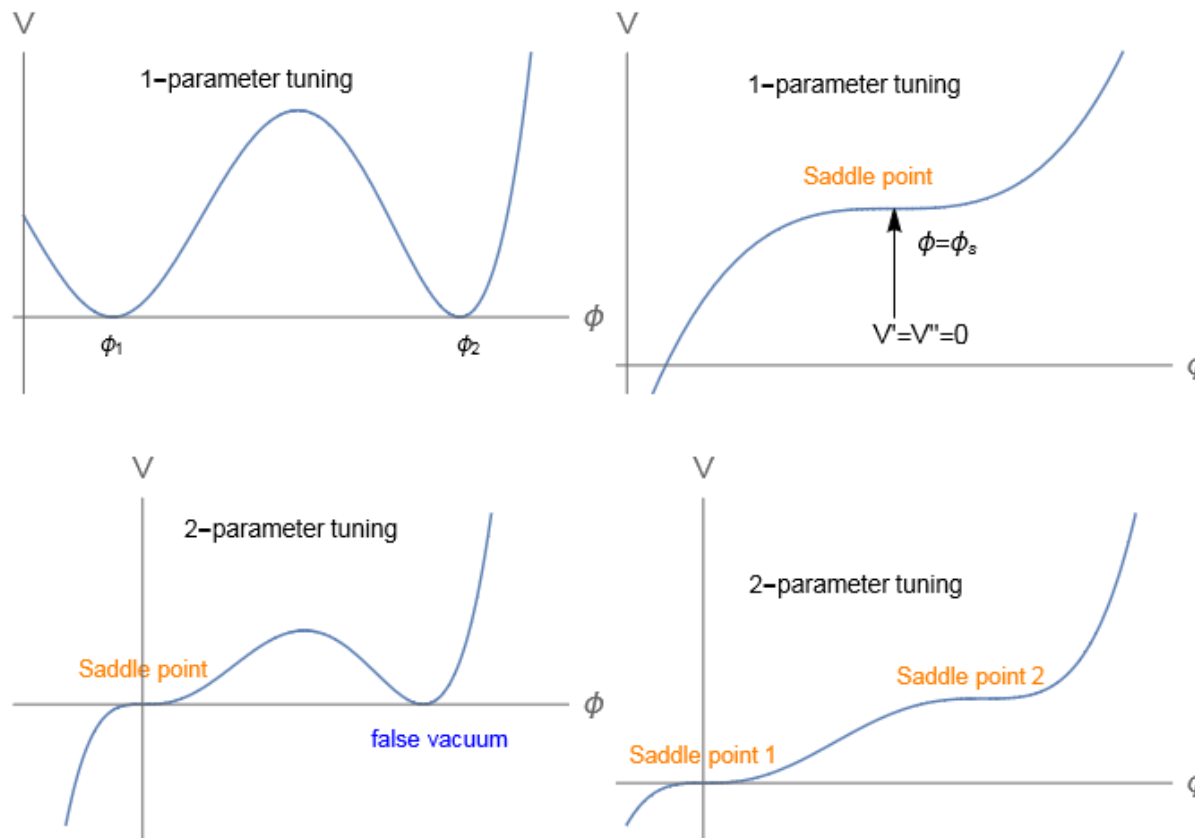
Multicritical point principle (MPP)

Froggatt, Nielsen (1995)

MPP: More parameters in the effective potential are tuned to a set of critical values, which is more likely to be realized by nature.

Examples:

Kawai, Kawana, 2107.10720 [hep-th]

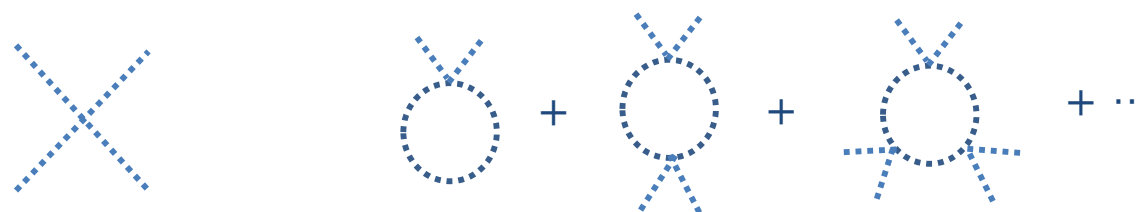


Dimensional transmutation

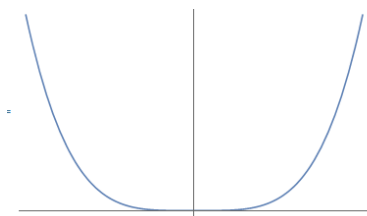
Coleman, Weinberg (1972)

Massless (CSI) Φ^4 theory: $V_{\text{tree}} = \frac{\lambda}{4!} \phi^4$

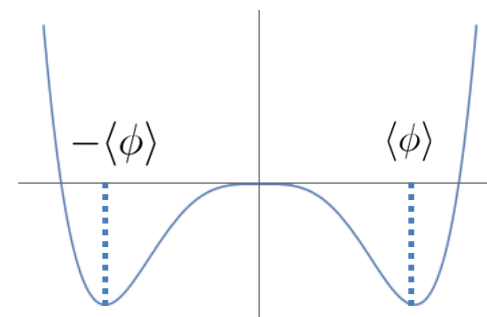
Eff. potential at 1-loop:



$$V_{\text{eff}} = \underbrace{\frac{\lambda}{4!} \phi_c^4}_{\text{tree}} + \underbrace{\frac{m_\phi^4}{64\pi^2} \left[\ln \frac{m_\phi^2}{\mu^2} - \frac{3}{2} \right]}_{\text{1-loop}} \quad m_\phi^2 = \frac{\lambda}{2} \phi_c^2$$



Dimensional
transmutation



$$\langle \phi \rangle \propto M_{\text{pl}} \exp \left(-\frac{16\pi^2}{3\lambda} + \frac{1}{2} \right) \quad (\mu = M_{\text{pl}})$$

However, it does not work...

2-scalar model

Bornholdt, Tetradis, Wetterich (1994)

Jennifer, Adams, Tetradis (1995)

Φ and S with a Z_2 symmetry $S \rightarrow -S$ and $\langle S \rangle = 0$

$$V_{\text{tree}} = \frac{\lambda}{4!} \phi^4 + \frac{\lambda_S}{4!} S^4 + \frac{\lambda_{\phi S}}{4} \phi^2 S^2$$

$$m_\phi^2 = \frac{\lambda}{2} \phi_c^2$$

$$m_S^2 = \frac{\lambda_{\phi S}}{2} \phi_c^2$$

Negligible if $\lambda \ll \lambda_{\phi S}$

$$V_{\text{eff}} = \frac{\lambda}{4!} \phi_c^4 + \frac{m_\phi^4}{64\pi^2} \left[\ln \frac{m_\phi^2}{\mu^2} - \frac{3}{2} \right] + \frac{m_S^4}{64\pi^2} \left[\ln \frac{m_S^2}{\mu^2} - \frac{3}{2} \right]$$

Dimensional transmutation: $\langle \phi \rangle \propto M_{\text{pl}} \exp \left(-\frac{16\pi^2}{3} \frac{\lambda}{\lambda_{\phi S}^2} + \frac{1}{2} \right)$

Taking $\lambda/\lambda_{\phi S}^2 = O(0.1)$, the **EW scale is generated** from the Planck scale!

In addition, S can be identified as a **dark matter** candidate!

Maximal Criticality

Kawai, Kawana, 2107.10720 [hep-th]

Generally, the effective potential for Φ is written as

$$V_{\text{eff}} = \mu_1^3 \phi_c + \frac{\mu_2^2}{2} \phi_c^2 + \frac{\mu_3}{3!} \phi_c^3 + \frac{\lambda_\phi}{4!} \phi_c^4 + \frac{m_S^4}{64\pi^2} \left[\ln \frac{m_S^2}{\mu^2} - \frac{3}{2} \right]$$

We can take $\mu = M$ at which λ_ϕ term vanishes.

$$V_{\text{eff}} = \mu_1^3 \phi_c + \frac{\mu_2^2}{2} \phi_c^2 + \frac{\mu_3}{3!} \phi_c^3 + \frac{\lambda_{\phi S}^2}{64\pi^2} \ln \frac{\phi_c^2}{M^2}$$

- The shape of the potential is determined by 3 parameters e.g., μ_1 , μ_2 and μ_3 .
- This potential generally has 5 extrema.

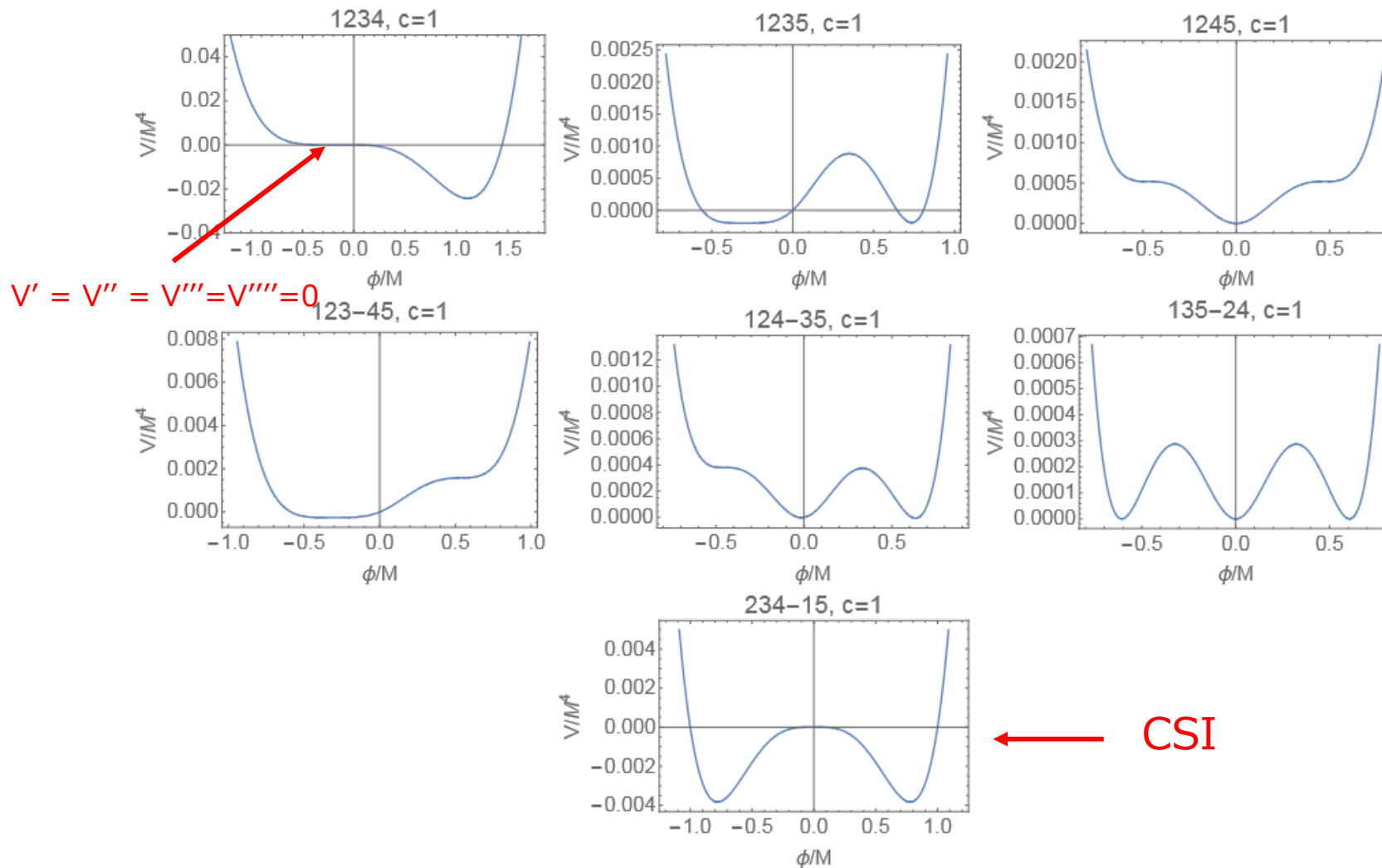


We can take **triple criticality** as the maximal critical point.

Maximal Criticality

Kawai, Kawana, 2107.10720 [hep-th]

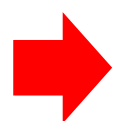
There are 7 independent triply critical points.



We focus on CP 1234, which may realize the strongly 1st order PT → Gravitational wave?

CP 1234

CP 1234: $\left. \frac{dV_{\text{eff}}}{d\phi} \right|_{\phi \rightarrow \phi_S} = \left. \frac{d^2V_{\text{eff}}}{d\phi^2} \right|_{\phi \rightarrow \phi_S} = \left. \frac{d^3V_{\text{eff}}}{d\phi^3} \right|_{\phi \rightarrow \phi_S} = \left. \frac{d^4V_{\text{eff}}}{d\phi^4} \right|_{\phi \rightarrow \phi_S} = 0.$

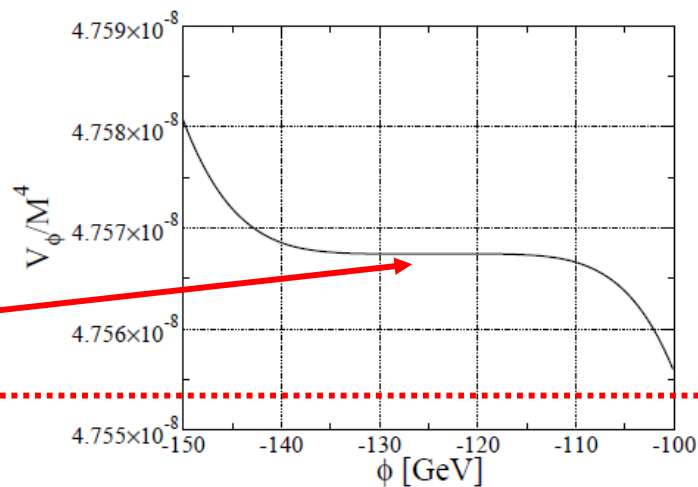
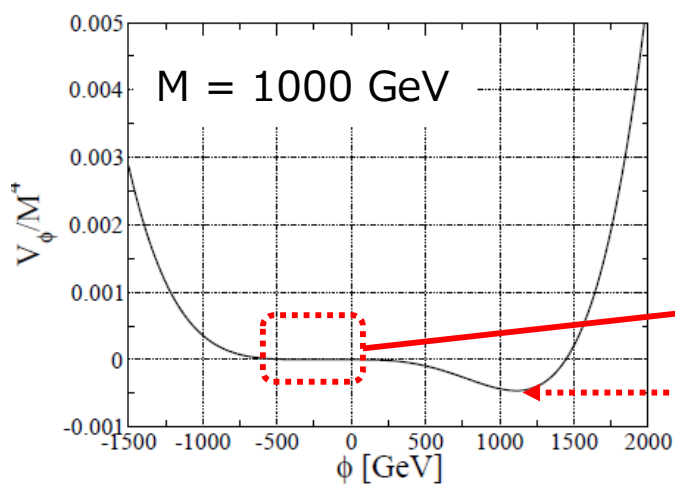


$$\mu_1^3 = -\frac{\kappa M^3}{18e^{25/4}}, \quad \mu_2^2 = -\frac{\kappa M^2}{4e^{25/6}}, \quad \mu_3 = -\frac{\kappa M}{e^{25/12}}, \quad \phi_S = -\frac{M}{e^{25/12}}$$

$$\kappa = \frac{3}{16\pi^2} \lambda_{\phi_S}^2$$

$$\bar{\phi} = \frac{\phi}{M}$$

$$V_{\text{eff}} = \kappa M^4 \left[-\frac{\bar{\phi}_c}{18e^{25/4}} - \frac{\bar{\phi}_c^2}{8e^{25/6}} - \frac{\bar{\phi}_c^3}{6e^{25/12}} + \frac{\bar{\phi}_c^4}{48} \ln \bar{\phi}_c^2 \right]$$



True vacuum

$$v_\phi \simeq 1.1M$$

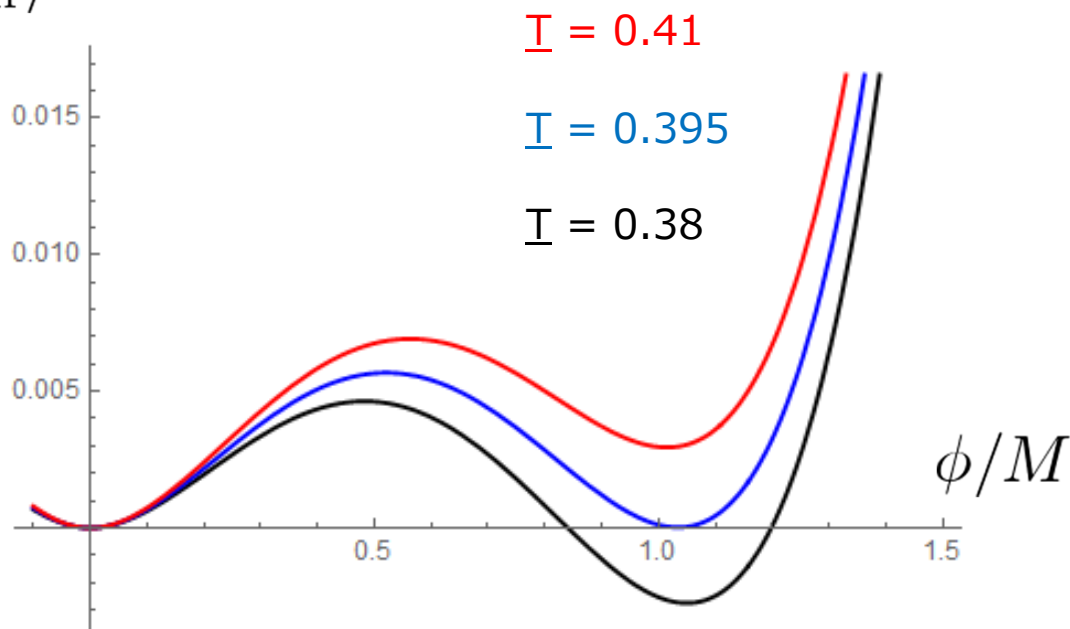
Potential at finite temperature

$$V_{\text{eff}}(\bar{\phi}, T) = V_{\text{eff}}(\bar{\phi}, T = 0) + \Delta V_{\text{eff}}(\bar{\phi}, T) \quad \bar{\phi} = \frac{\phi}{M} \quad \bar{T} = \frac{T}{M} \frac{1}{\sqrt{2\pi}} \left(\frac{3}{\kappa}\right)^{1/4} = \frac{T}{M} \sqrt{\frac{2}{\lambda_{\phi S}}}$$

$$= \kappa M^4 \left[-\frac{\bar{\phi}}{18e^{25/4}} - \frac{\bar{\phi}^2}{8e^{25/6}} - \frac{\bar{\phi}^3}{6e^{25/12}} + \frac{\bar{\phi}^4}{48} \ln \bar{\phi}^2 + \frac{2\bar{T}^4}{3} \int_0^\infty dx x^2 \ln \left(1 - e^{-\sqrt{x^2 + \bar{\phi}/\bar{T}^2}} \right) \right]$$

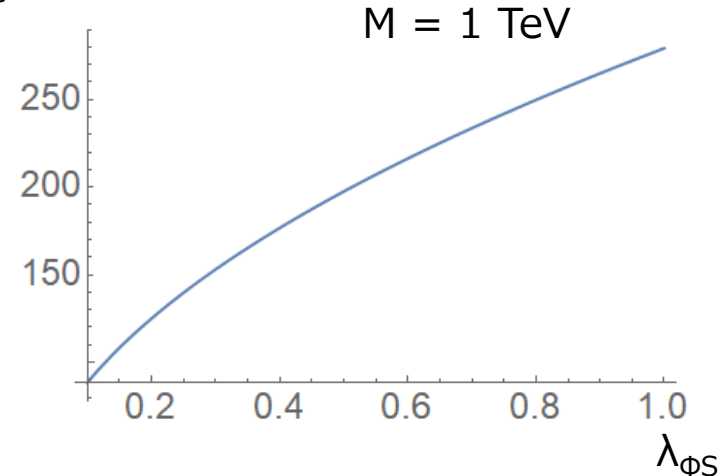
The shape of the potential is determined only by the $\bar{\mathbb{I}}$ parameter!

$V_{\text{eff}}/M^4 \kappa$



Critical temperature $\bar{\mathbb{T}}_c \sim 0.395!$

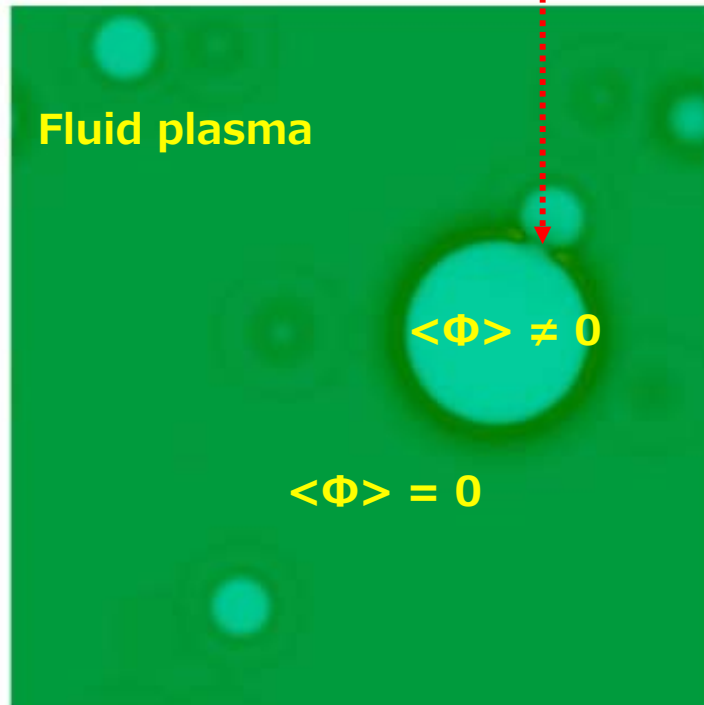
T_c [GeV]



1st order phase transition can be realized!

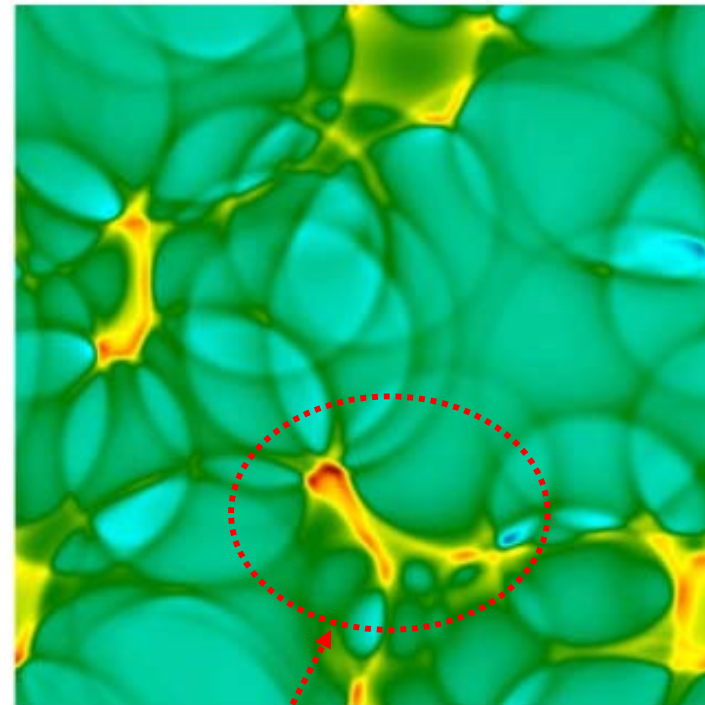
Gravitational Waves from 1st OPT

Energy density of fluid



Bubble collision

Hindmarsh, Huber, Rummukainen, Weir (2013)



Grojean, Servant (2007)

Leitao, Megevand, Sanchez (2012)

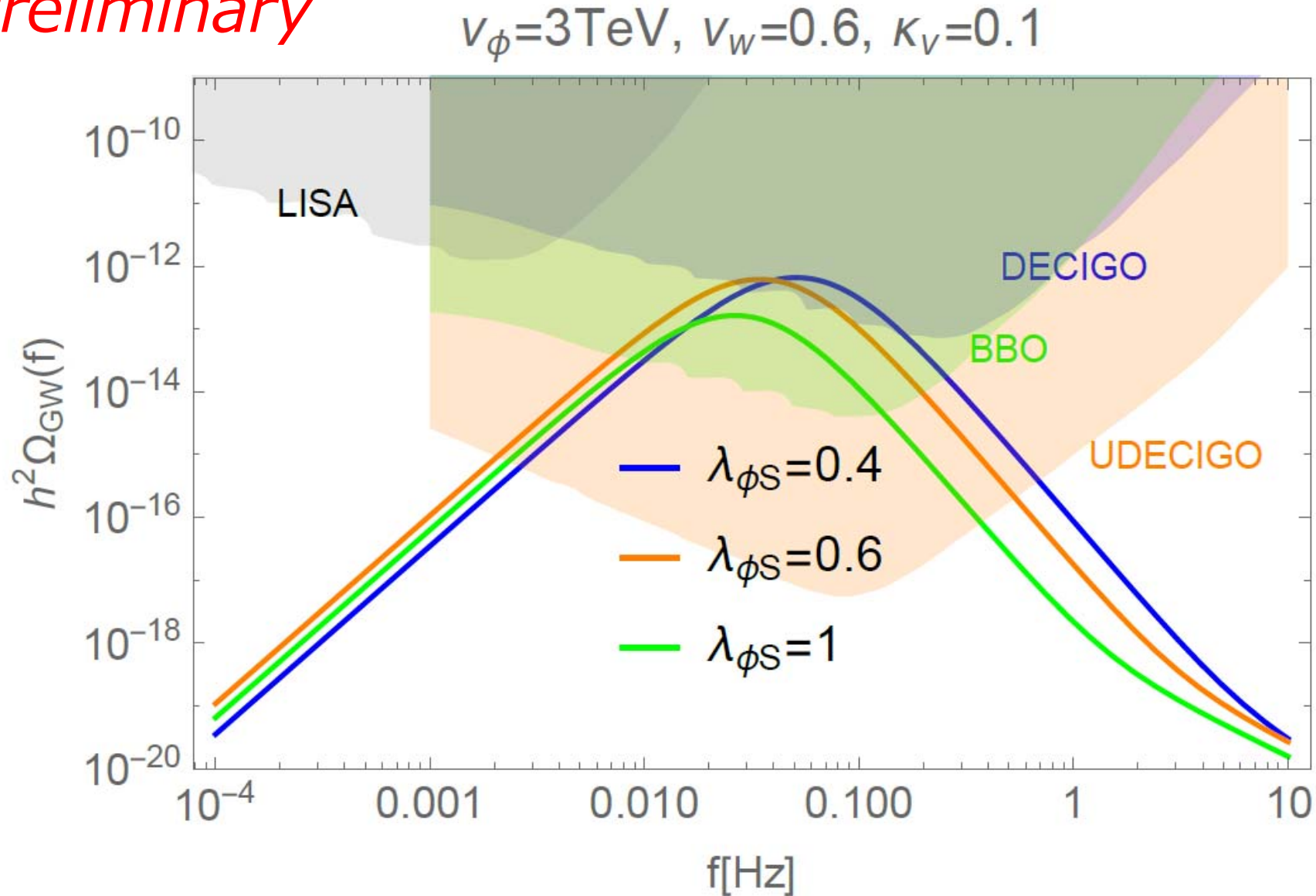
Jinno, Nakayama, Takimoto (2016)

Hashino, Kakizaki, Kanemura, Matsui (2016), ...

Gravitational Waves

Hamada, Kawai, Kawana, Oda, KY, in progress

Preliminary



Implementation to the SM

$$\mathcal{L} = \mathcal{L}_{\text{SM}}^{\text{w/o pot.}} - V_{\text{MPP}}(H, \phi, S) \quad V_{\text{MPP}} = V_{\text{CSI}} + V_{\phi}^{\text{CP-1234}}$$

Trigger the EWSB after Φ gets the VEV.

$$V_{\text{CSI}} = \frac{\lambda_H}{2} (H^\dagger H)^2 + \frac{\lambda_S}{4!} S^4 - \frac{\lambda_{\phi H}}{2} \phi^2 (H^\dagger H) + \frac{\lambda_{\phi S}}{4} \phi^2 S^2 + \frac{\lambda_{SH}}{2} S^2 (H^\dagger H)$$

$$\rightarrow v_H = v_\phi \sqrt{\frac{\lambda_{\phi H}}{\lambda_H}} \sim 246 \text{ GeV}$$

$$m_h^2 \simeq v_H^2 \lambda_H \sim (125 \text{ GeV})^2$$

$$m_H^2 \simeq 0.23 \kappa v_\phi^2 \simeq \left(130 \text{ GeV} \times \frac{\lambda_{\phi S}}{1} \times \frac{v_\phi}{2 \text{ TeV}} \right)^2$$

Only 3 free parameters! $(m_S, \lambda_{HS}, v_\phi)$

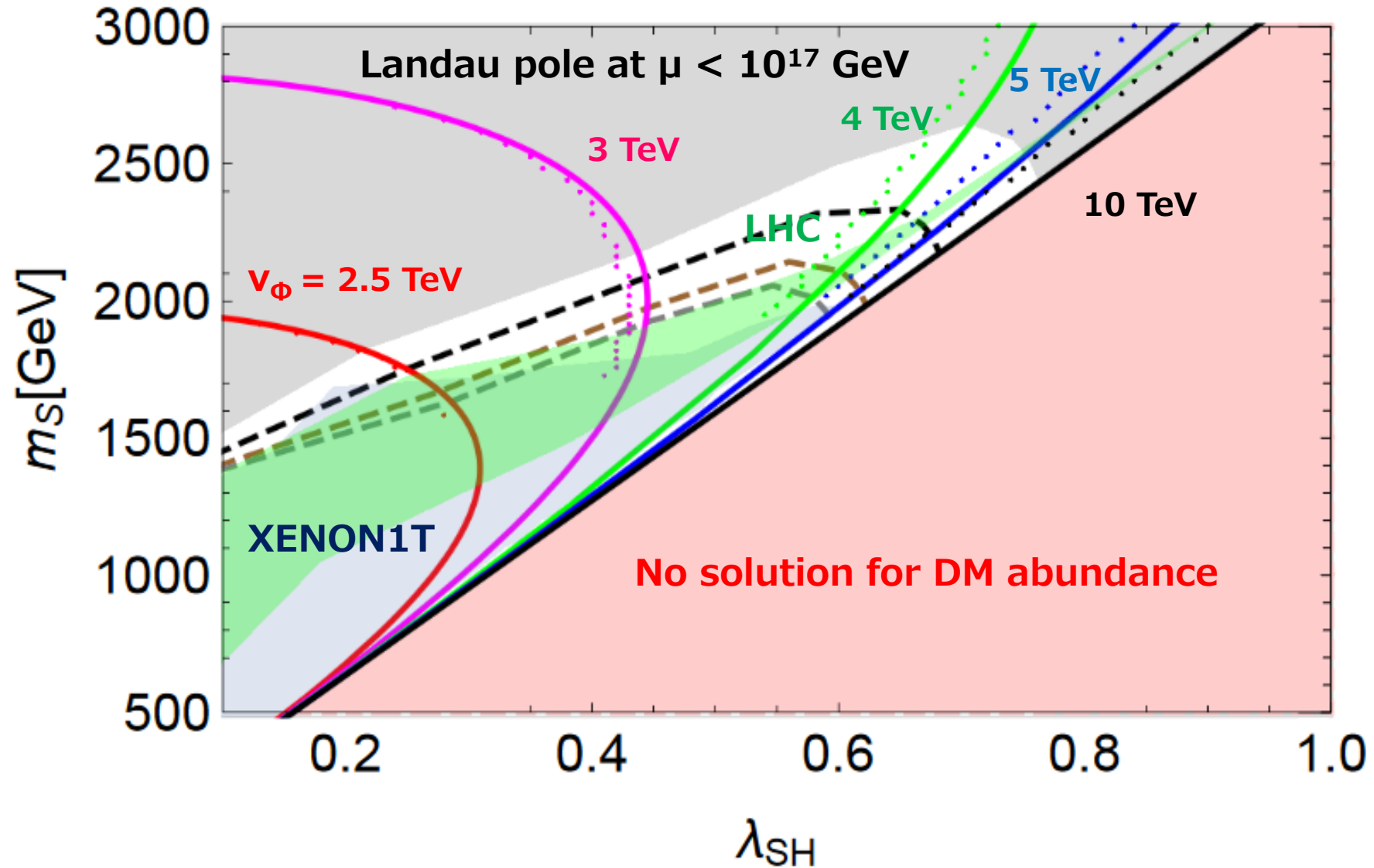
DM mass

DM-Higgs coupling

Phenomenology

Hamada, Kawai, Oda, KY, 2008.08700 [hep-ph]

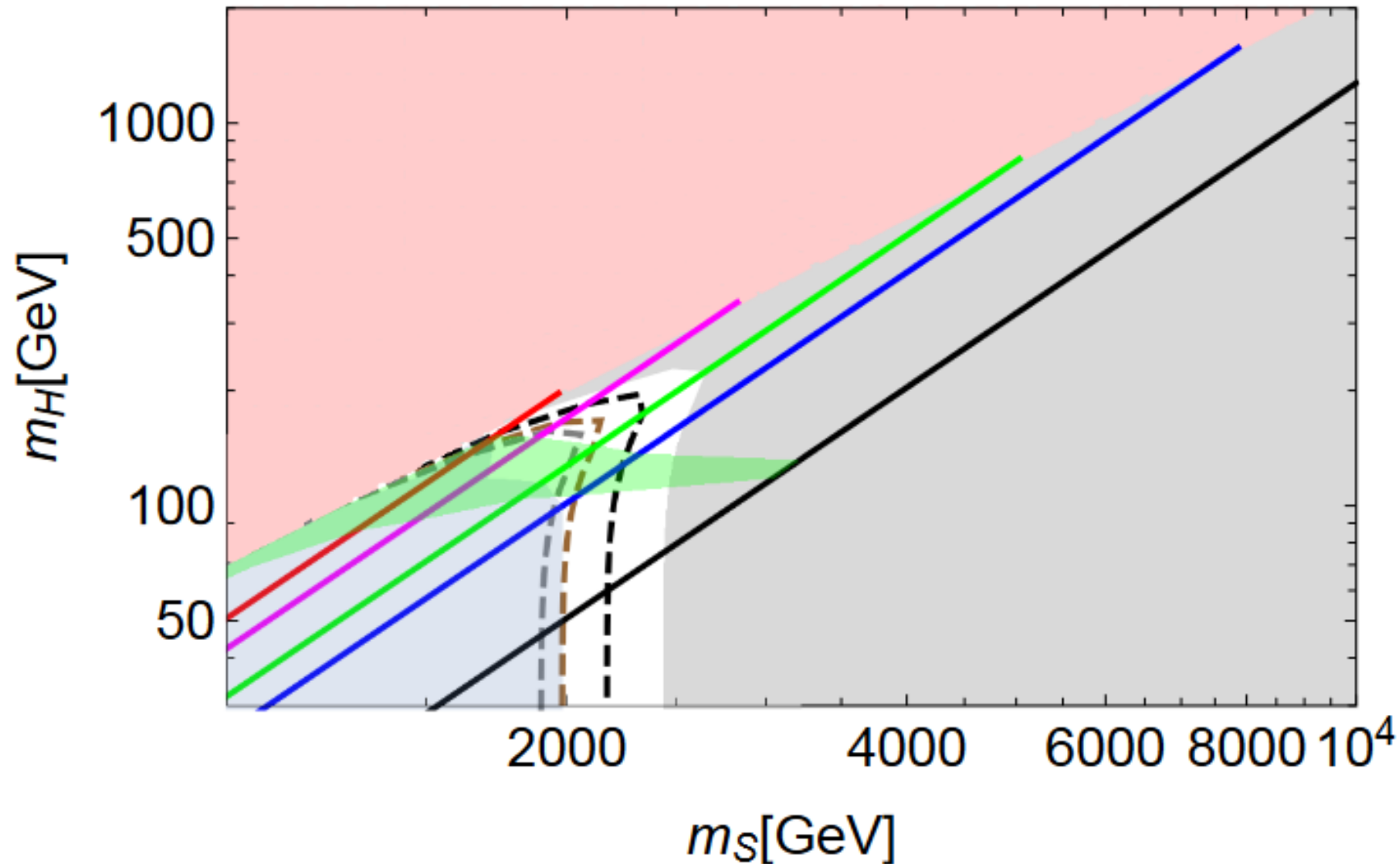
*This result is based on the other CP, but the result is almost the same.



Phenomenology

Hamada, Kawai, Oda, KY, 2008.08700 [hep-ph]

*This result is based on the other CP, but the result is almost the same.



Additional Higgs boson mass is predicted to be at around 150-200 GeV.

Contents

I. What is the Higgs boson?

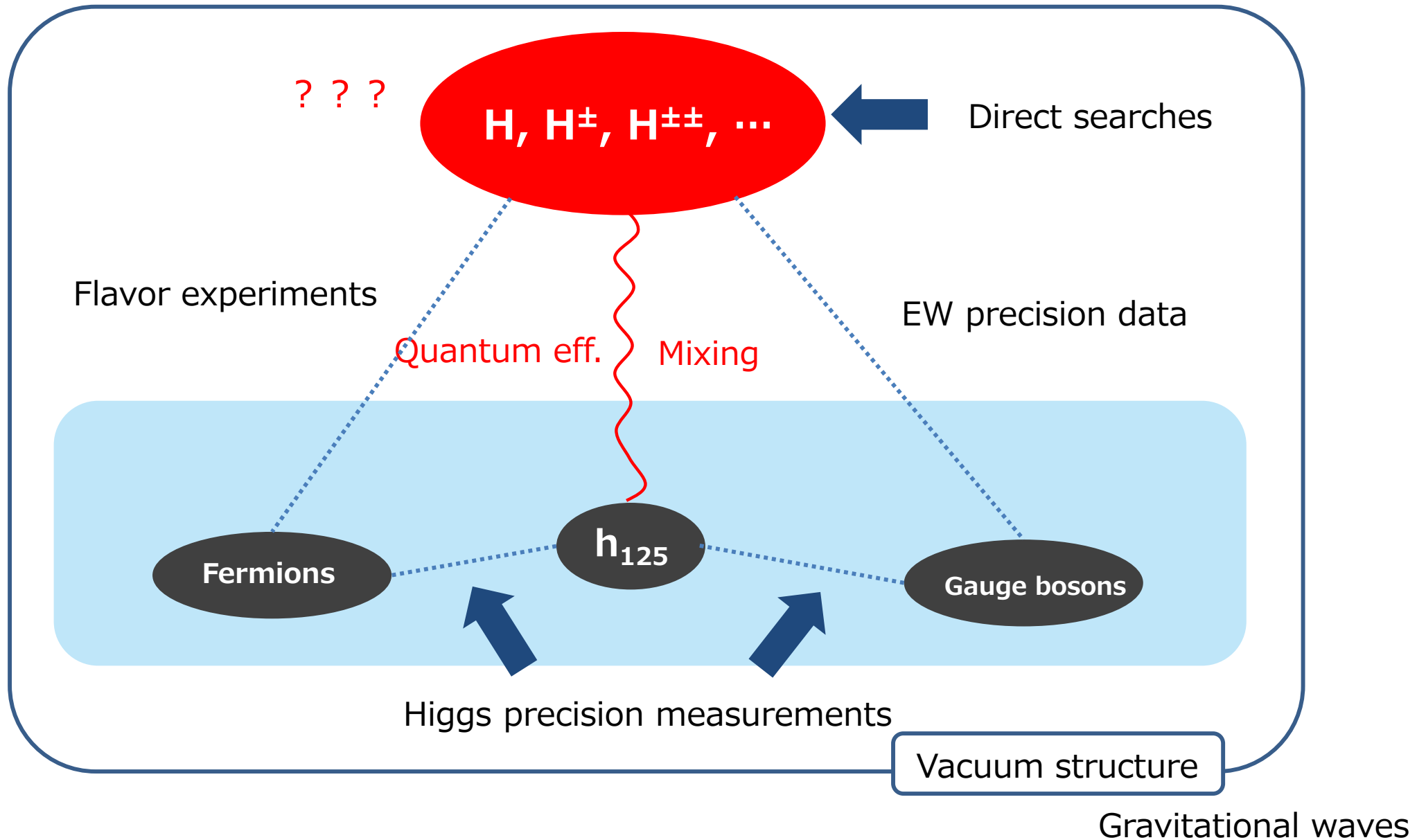
II. Why is Higgs Physics important?

III. How can we determine the Higgs sector?

- Bottom-up approach
- Top-down approach

IV. Summary

Exploring the Higgs sector



Constraints from EWPO

$$\rho_{\text{tree}} = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = \frac{\sum_b v_b^2 [T_b(T_b + 1) - Y_b^2]}{2 \sum_a v_a^2 Y_a^2}$$

$$\rho_{\text{exp}} = 1.00038 \pm 0.0002$$

$$\rho_{\text{tree}} \neq 1$$

Arbitrary fields

$$\rho_{\text{tree}} = 1$$

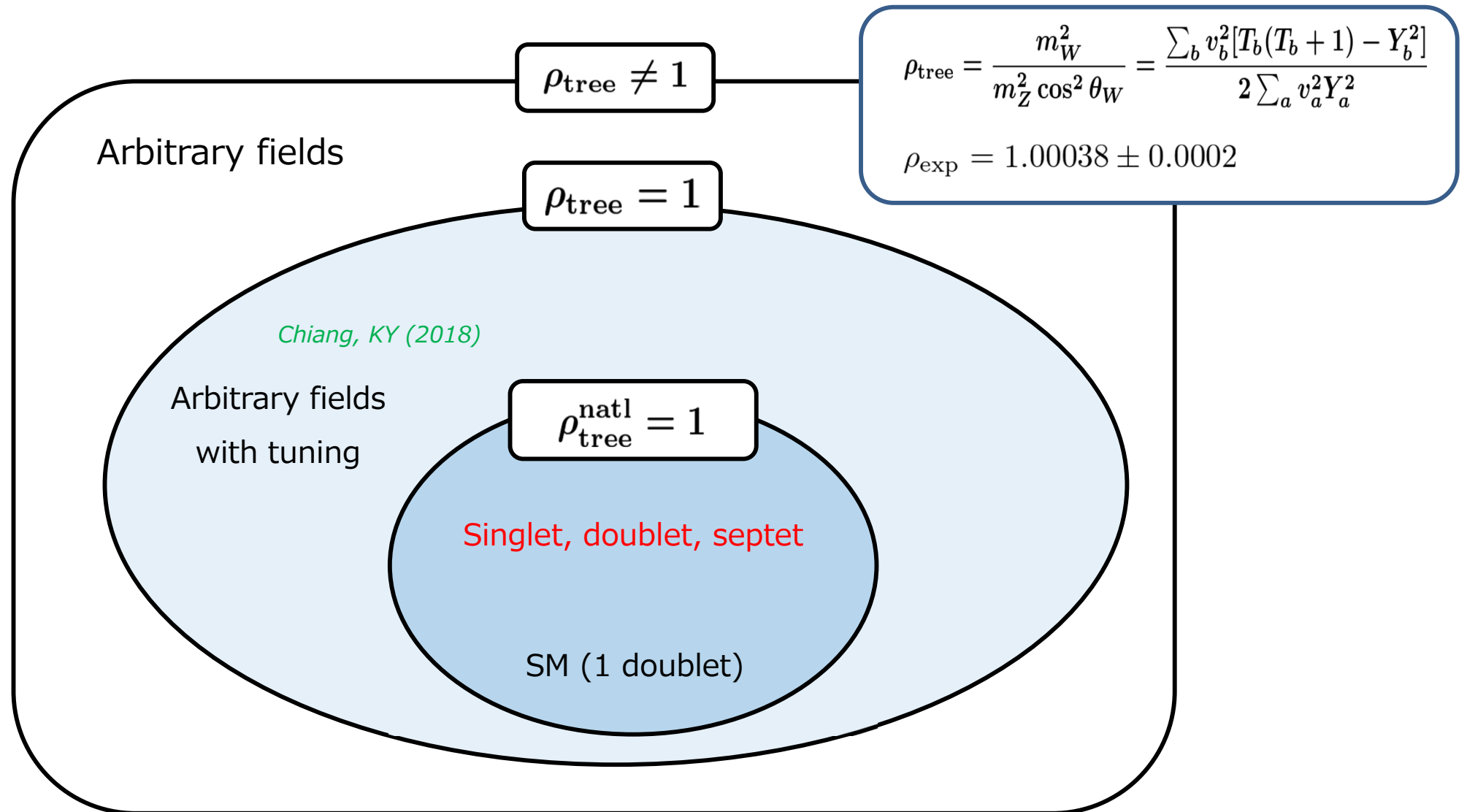
Chiang, KY (2018)

Arbitrary fields
with tuning

$$\rho_{\text{tree}}^{\text{natl}} = 1$$

Singlet, doublet, septet

SM (1 doublet)



Constraints from Flavor Experiments

Multi-doublet structures introduces FCNCs

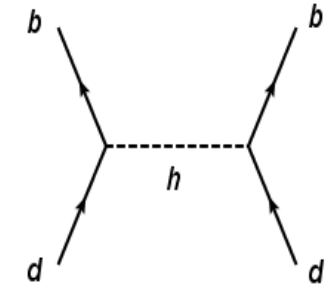
$$\mathcal{L}_Y = \bar{Q}_L (Y_1 \Phi_1 + Y_2 \Phi_2) d_R + \text{h.c.}$$

➔ $M = v_1 Y_1 + v_2 Y_2$

➔ $Y = \langle h_1 | h \rangle Y_1 + \langle h_2 | h \rangle Y_2$

Mass matrix

Interaction matrix



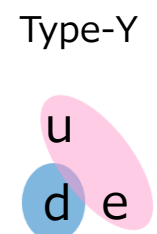
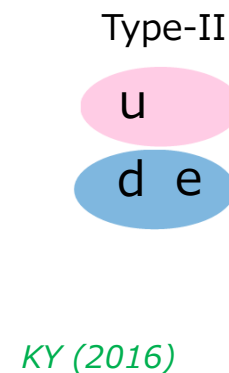
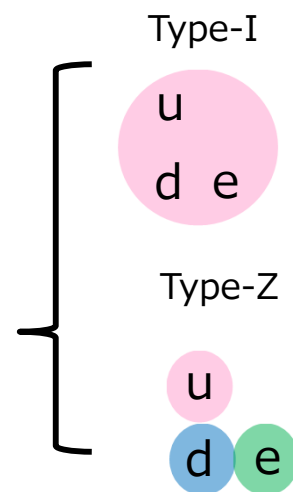
$M \gtrsim \mathcal{O}(100)$ TeV
for $\mathcal{O}(0.1)$ coupling

Generally $M \not\propto Y$ ➔ FCNC!

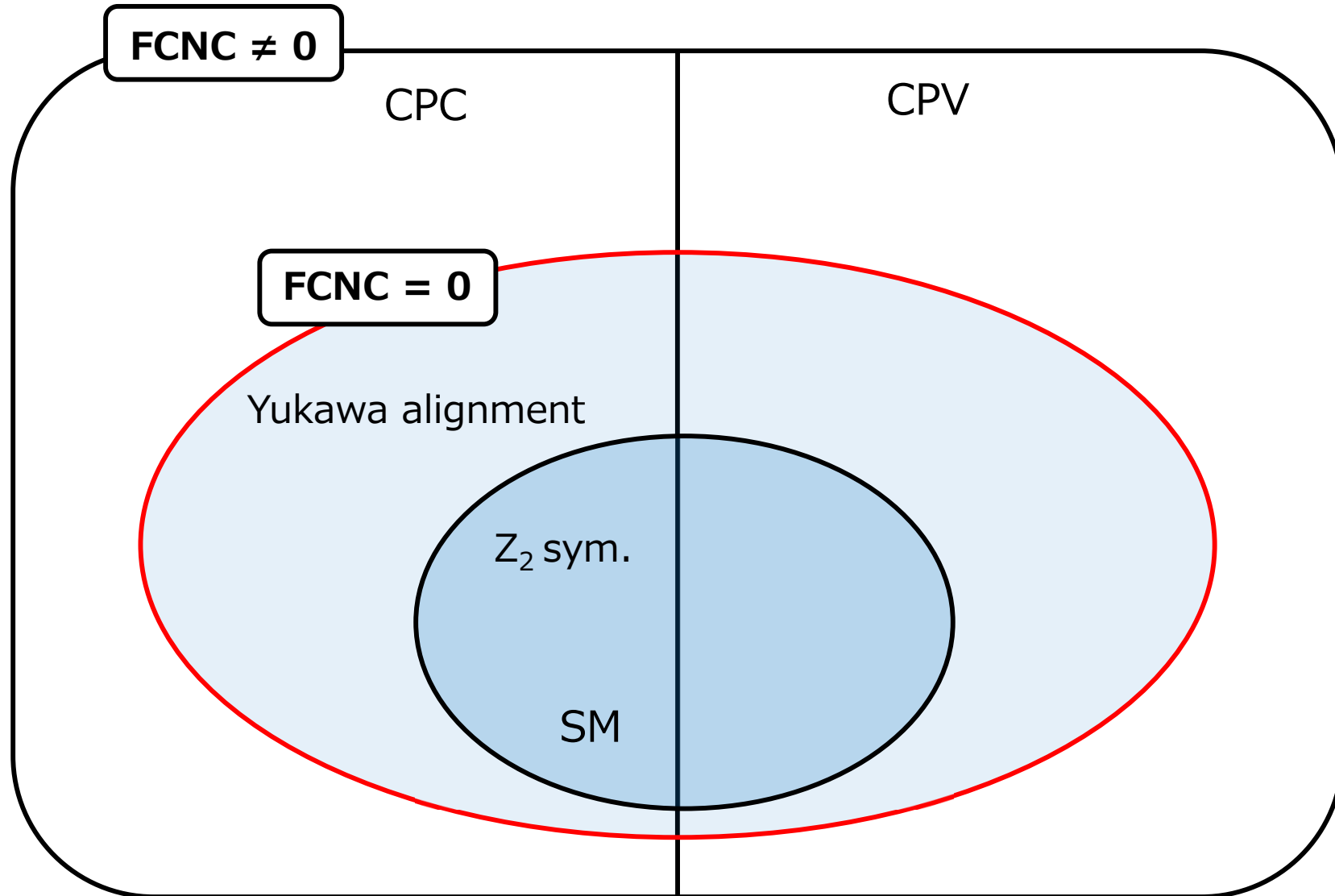
Barger, Hewett, Phillips (1990), Grossman (1994)
Aoki, Kanemura, Tsumura, KY (2009), Cite 300+

No FCNC {

- Pich, Tuzon (2009)*
Yukawa alignment
 $Y_1 \propto Y_2$
- Glashow, Weinberg (1977)*
 Z_2 symmetry
 $Y_2 \rightarrow 0$

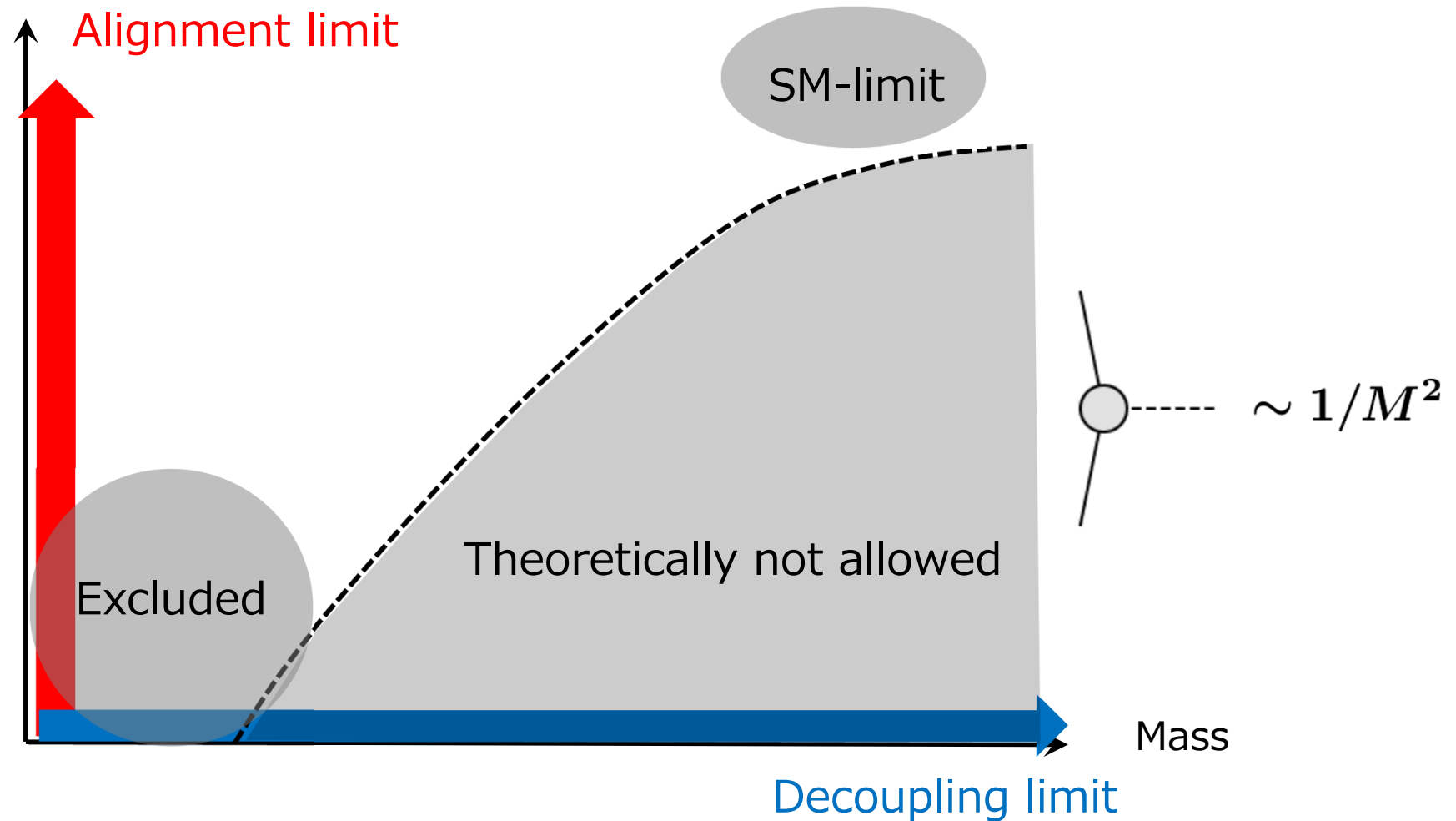


Classification of Multi-Doublet Models



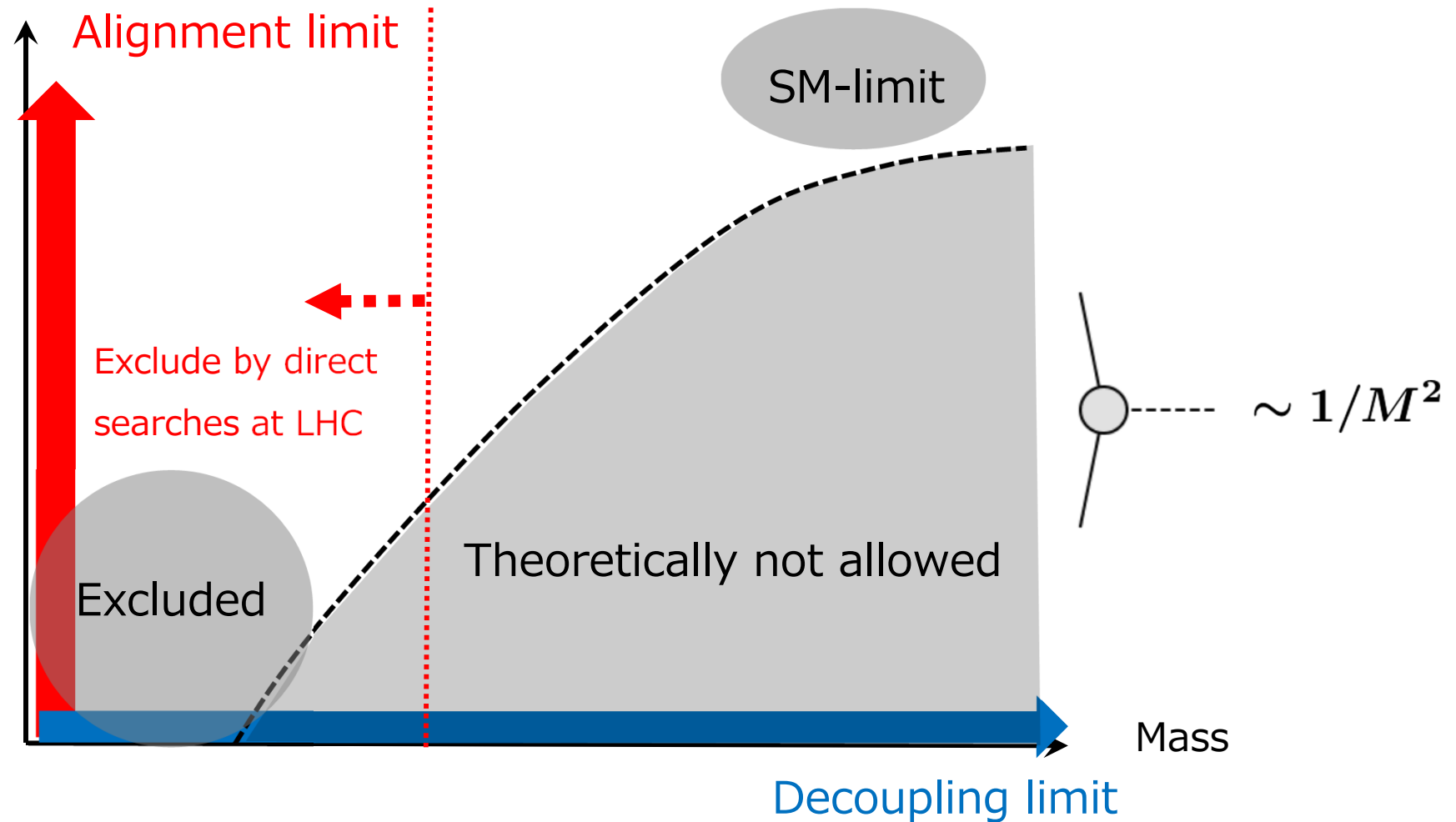
Alignment & Decoupling

SM-likeness of $h(125)$



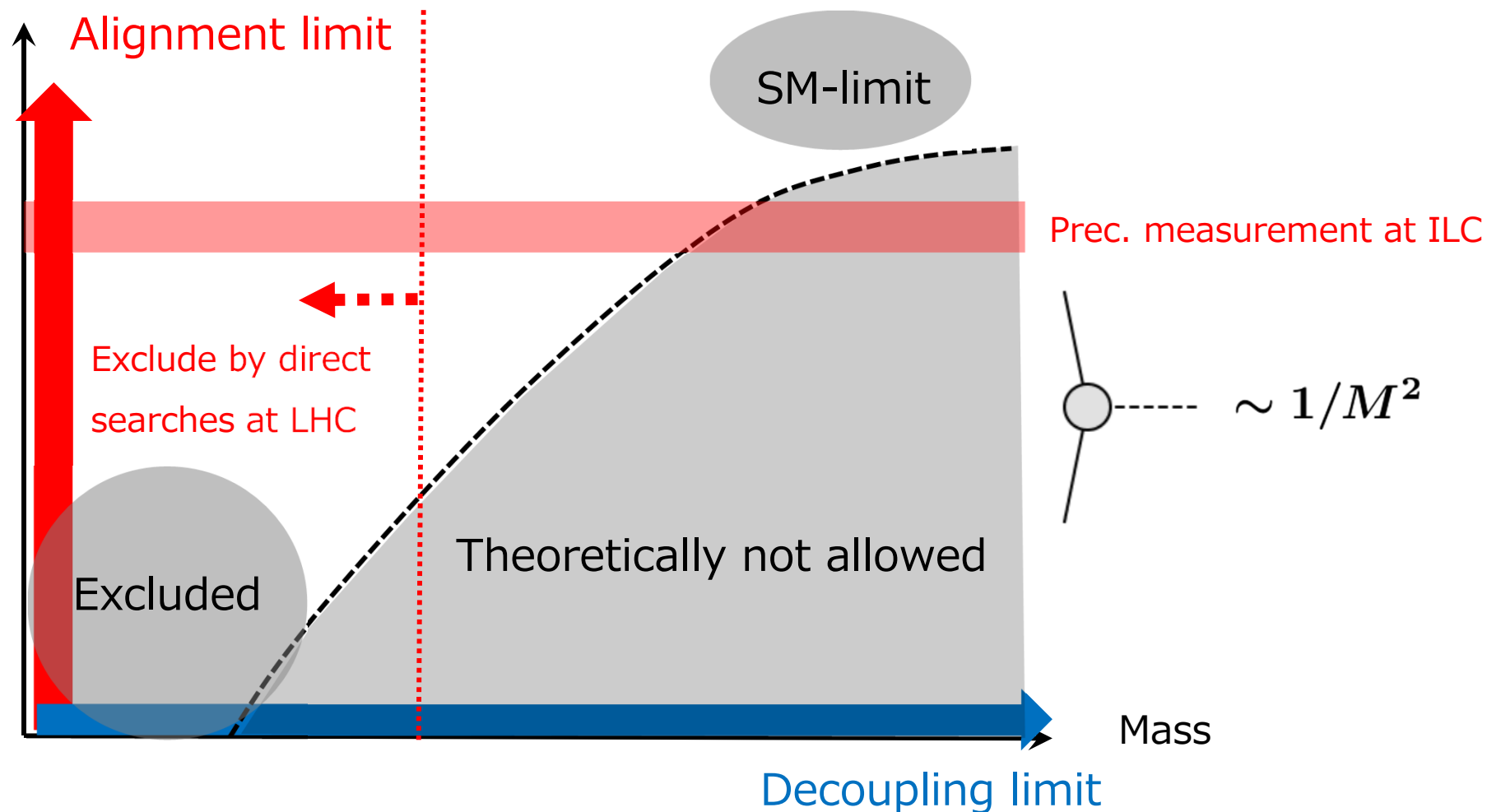
Alignment & Decoupling

SM-likeness of h(125)



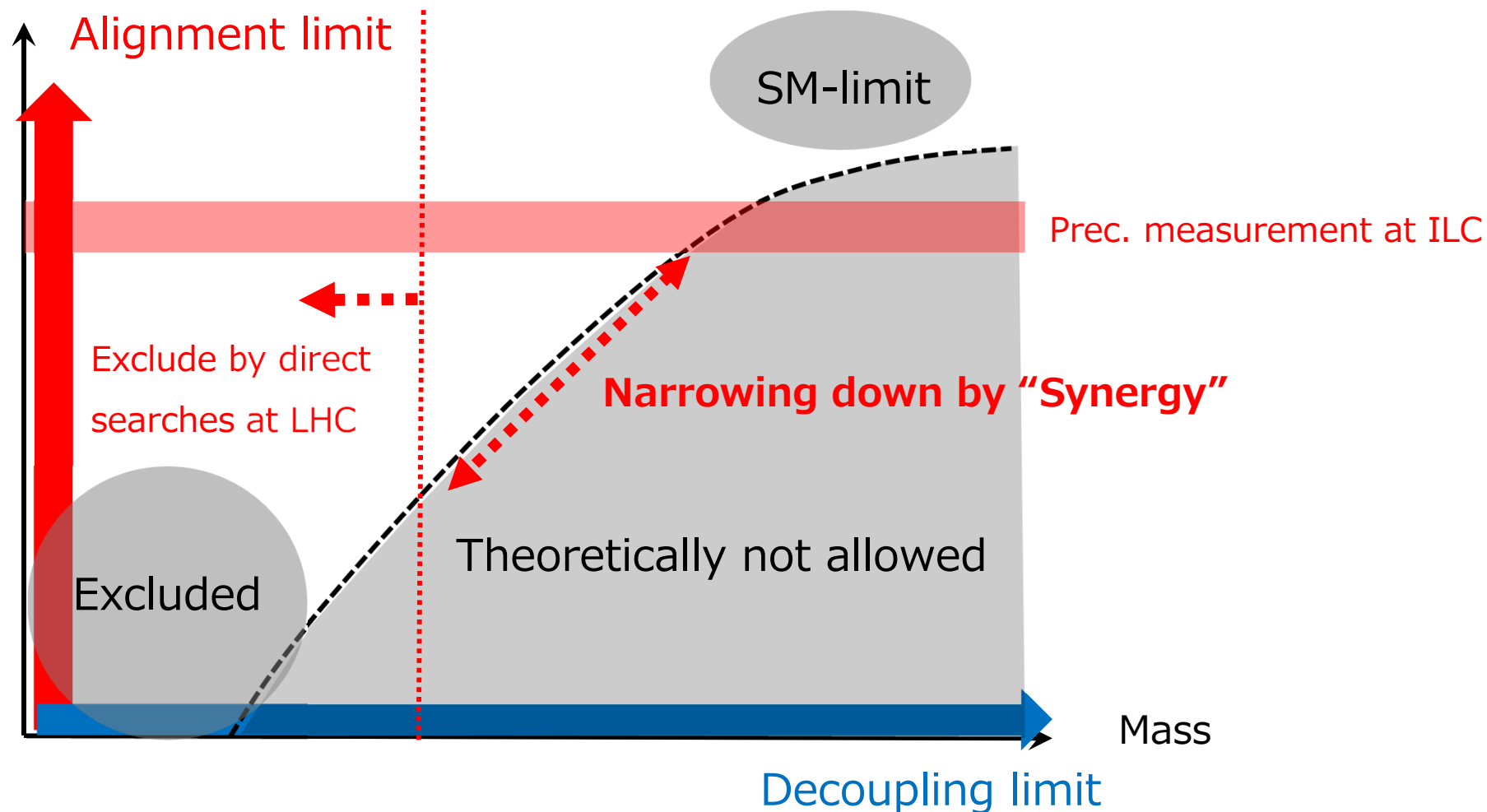
Alignment & Decoupling

SM-likeness of h(125)



Alignment & Decoupling

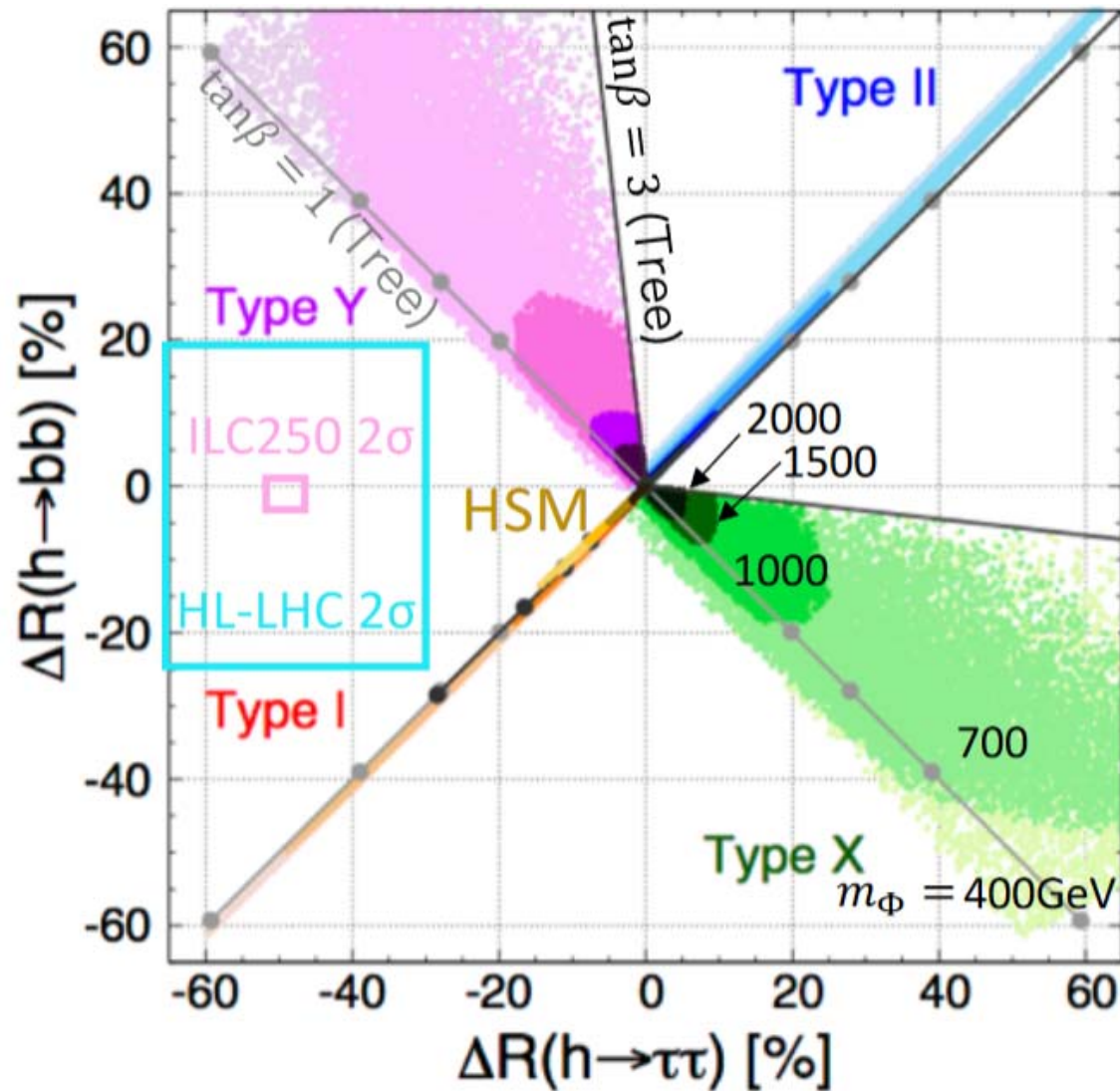
SM-likeness of $h(125)$



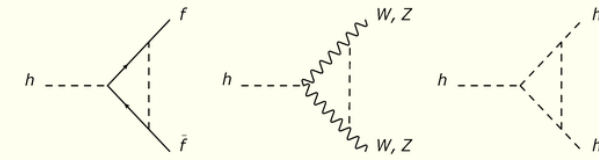
"New No Loose Theorem"

Fingerprinting

Kanemura, Kikuchi, Sakurai, Mawatari, KY, PLB783, 140 (2018)



H-COUP

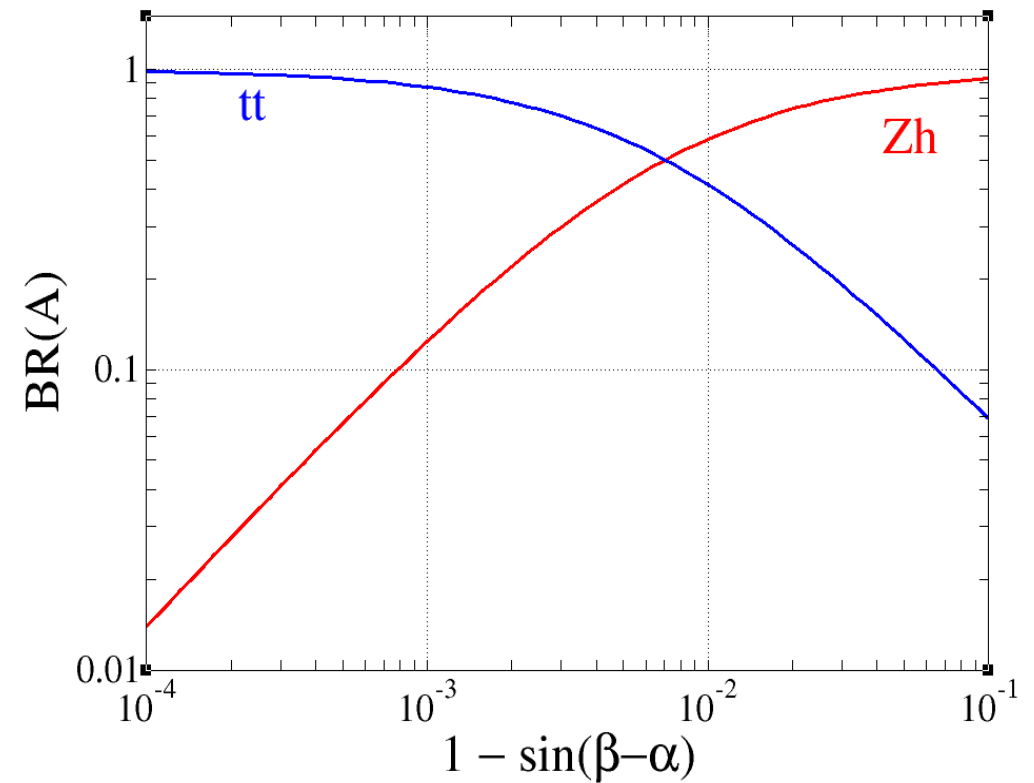
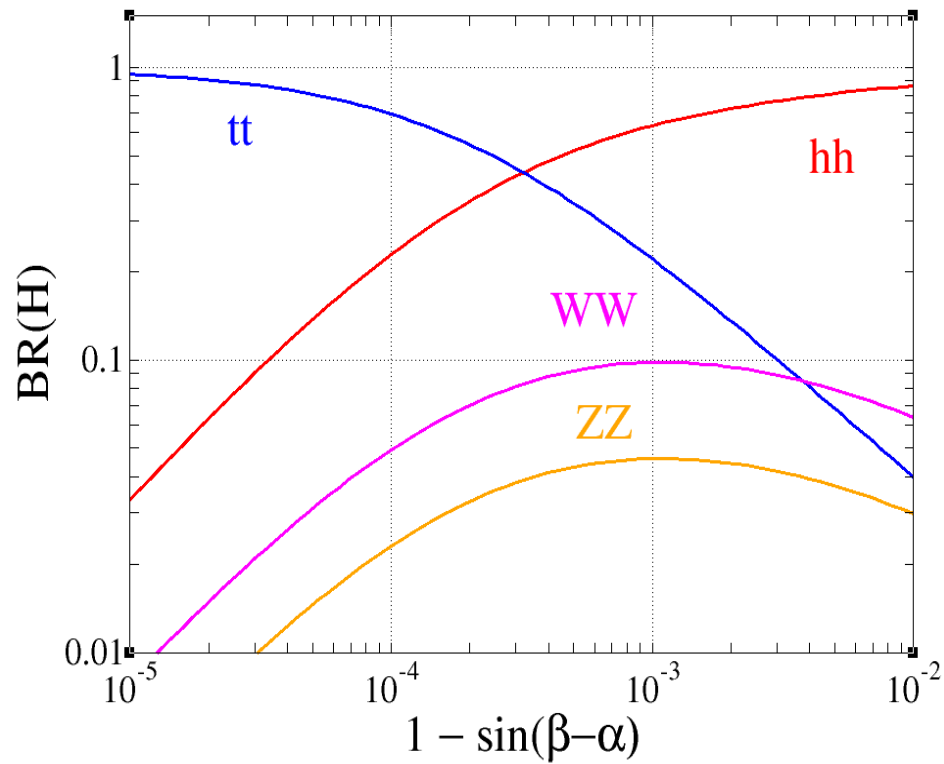


v1: Kanemura, Kikuchi, Sakurai, KY (2017)

v2: Kanemura, Kikuchi, Sakurai, Mawatari, KY (2019)

Higgs to Higgs decays @ near alignment

Type-I 2HDM with $m_H = m_A = m_{H^\pm} = M = 400$ GeV, $\tan\beta = 10$

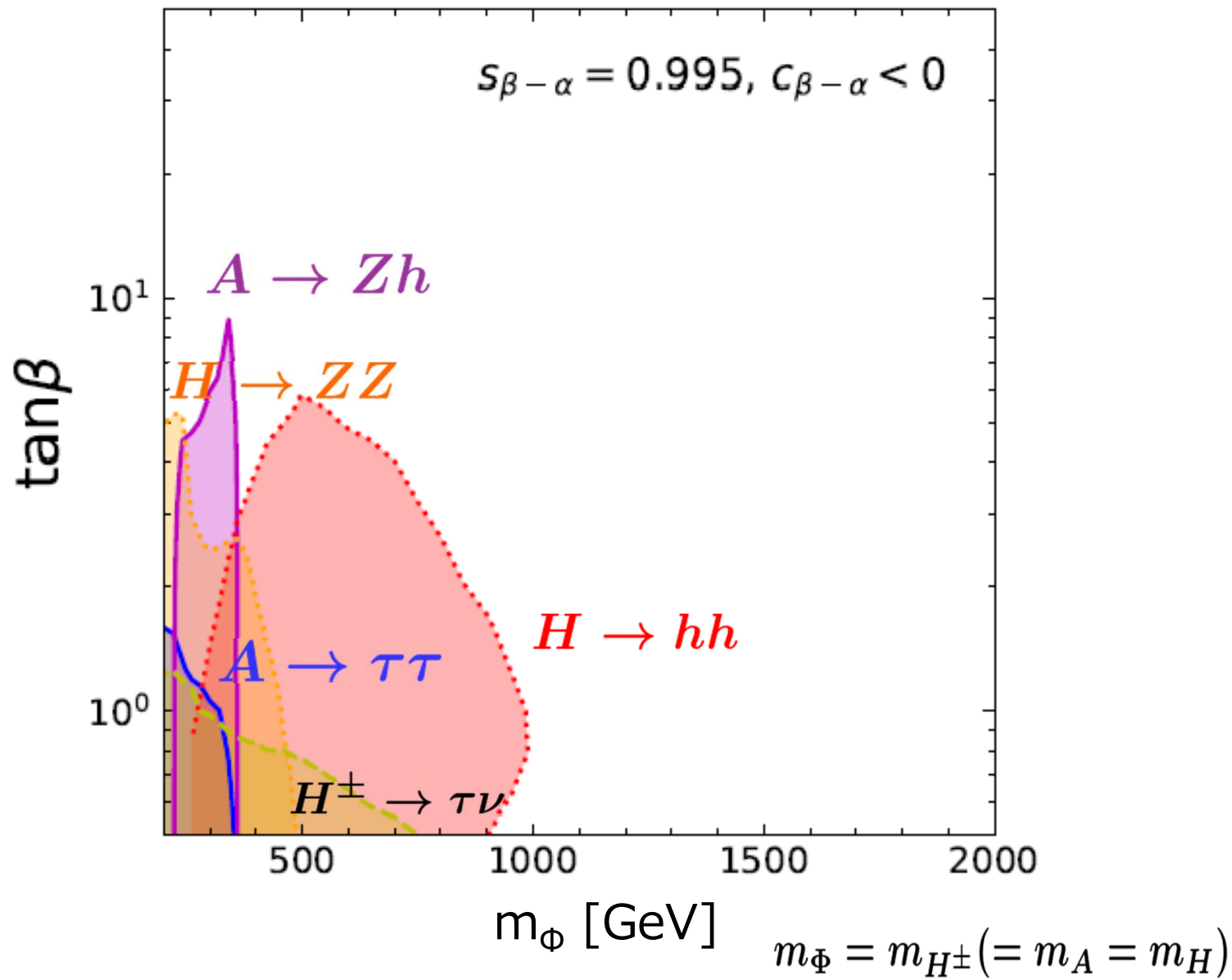


Alignment limit

Current LHC

Aiko, Kanemura, Kikuchi, Mawatari, Sakurai, KY (2020)

Type-I

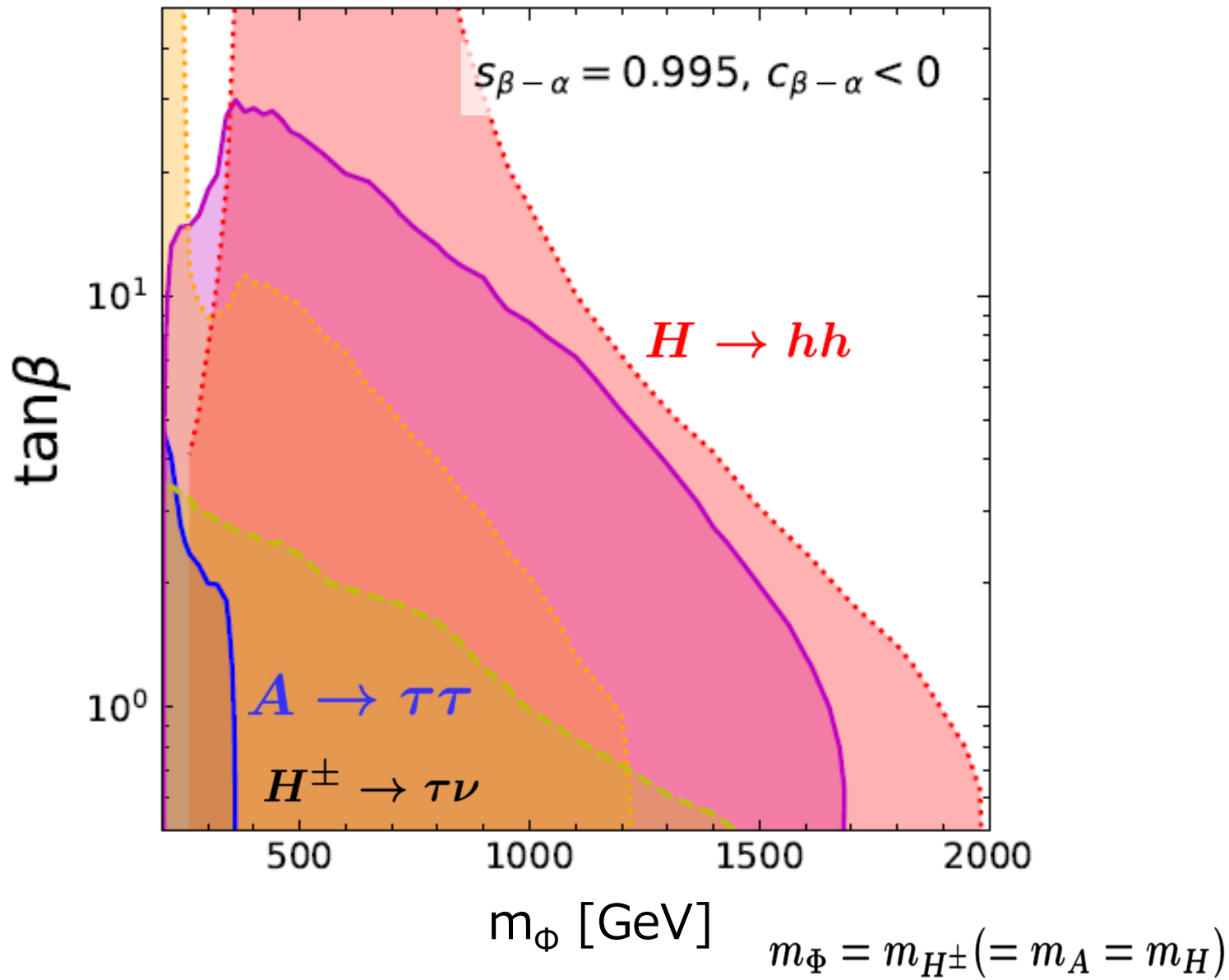


HL-LHC

Aiko, Kanemura, Kikuchi, Mawatari, Sakurai, KY (2020)

$H \rightarrow ZZ$

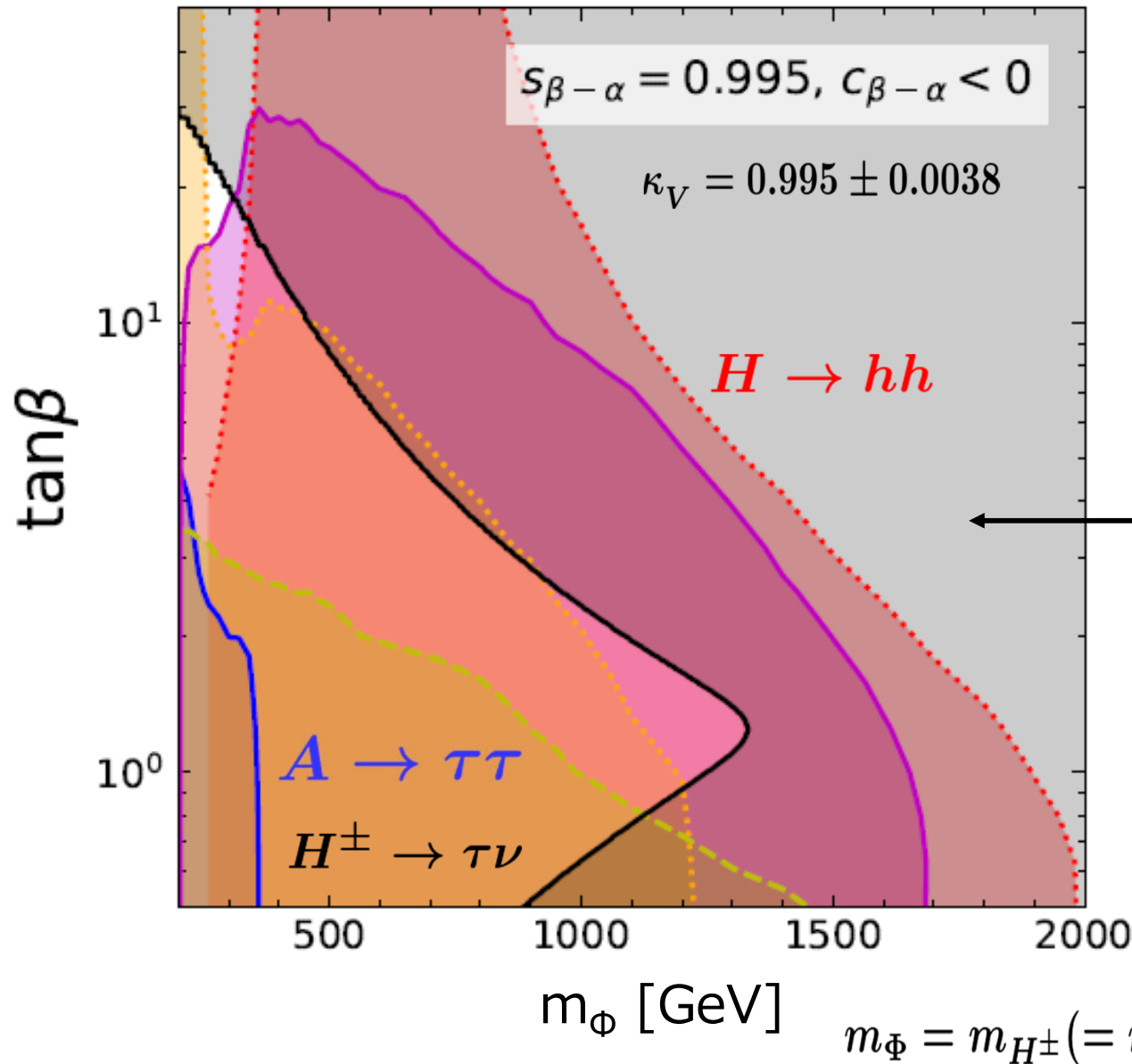
Type-I



HL-LHC & ILC

Aiko, Kanemura, Kikuchi, Mawatari, Sakurai, KY (2020)

$H \rightarrow ZZ$ Type-I



Excluded by unitarity and vacuum stab.

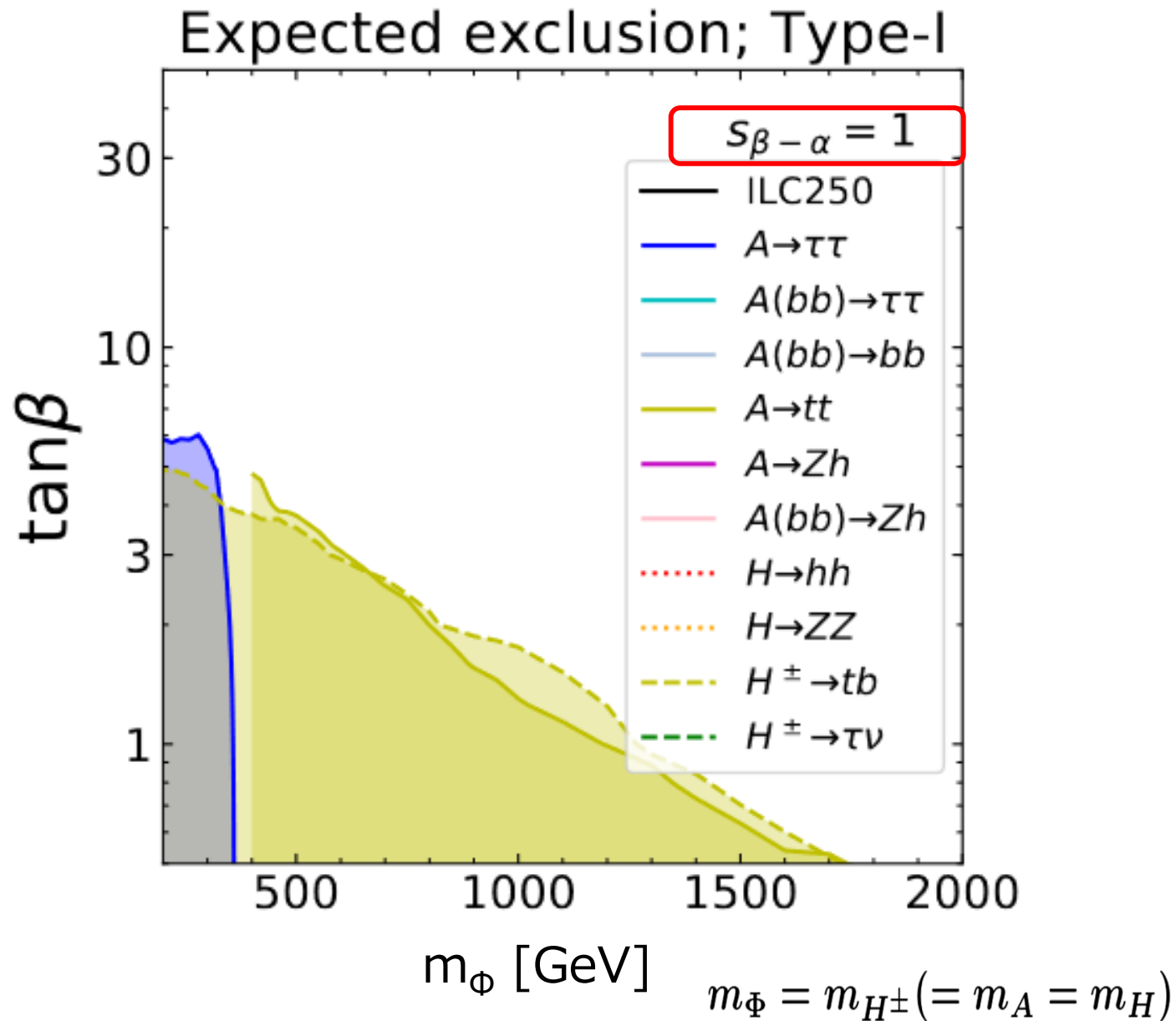
Despande, Ma (1978)

Kanemura, Kubota, Takasugi (1993)

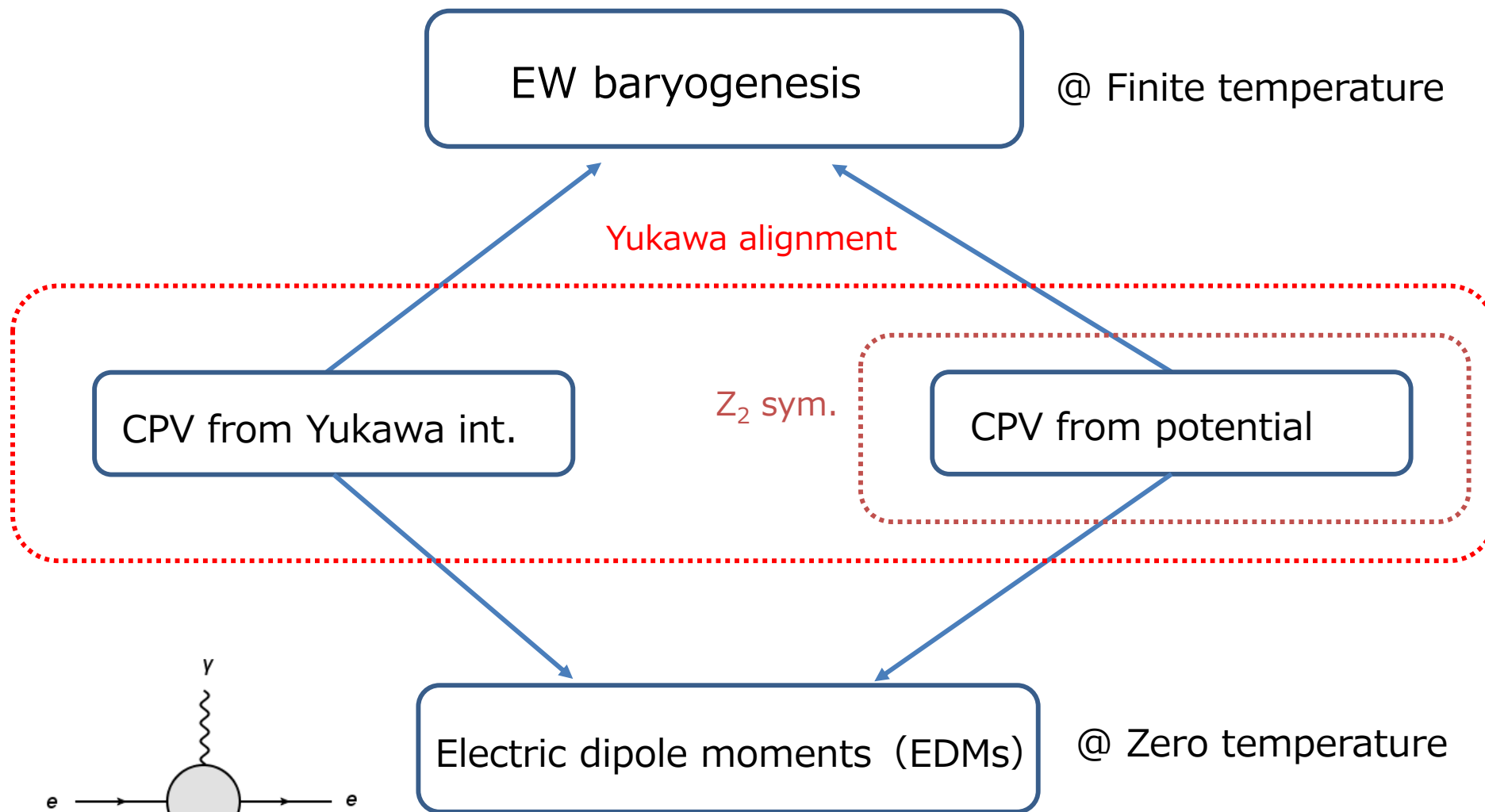
Kanemura, KY (2015)

Alignment limit

Aiko, Kanemura, Kikuchi, Mawatari, Sakurai, KY (2020)



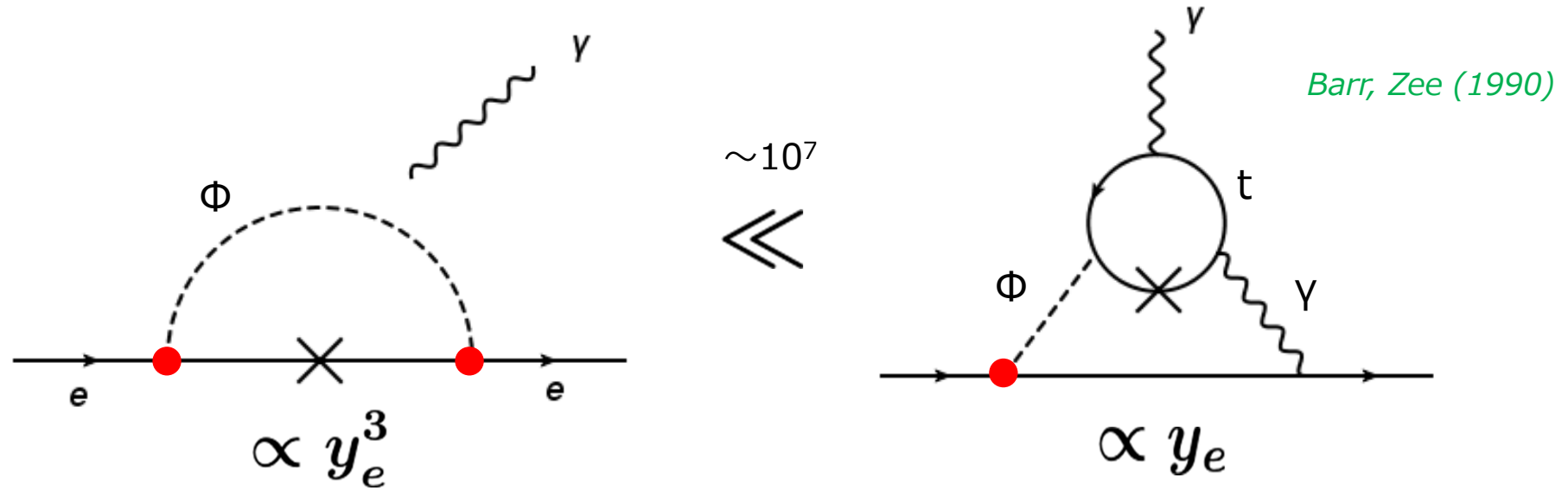
CP-violation in the alignment scenario



ACME Collaboration (2018)

$$|d_e| \leq 1.1 \times 10^{-29} e \text{ cm} \quad @90\% \text{ CL}$$

Contribution to EDMs

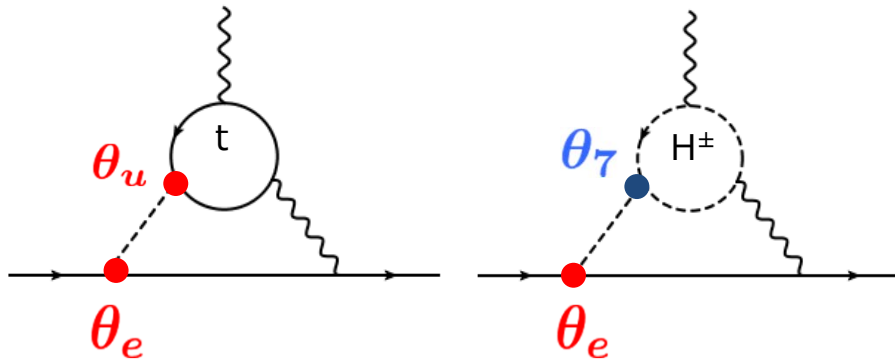


$$d_e \sim \left(\frac{1}{16\pi^2} \right)^2 \times e^3 G_F m_e \sim 10^{-27} \text{ e cm} \gg 10^{-29} \text{ e cm}$$

- Typically, new source of CPV is excluded by EDM experiments.
- We need a mechanism to get O(1) without conflict to EDMs.
- Destructive interference in EDMs can happen between Yukawa & potential CPV.

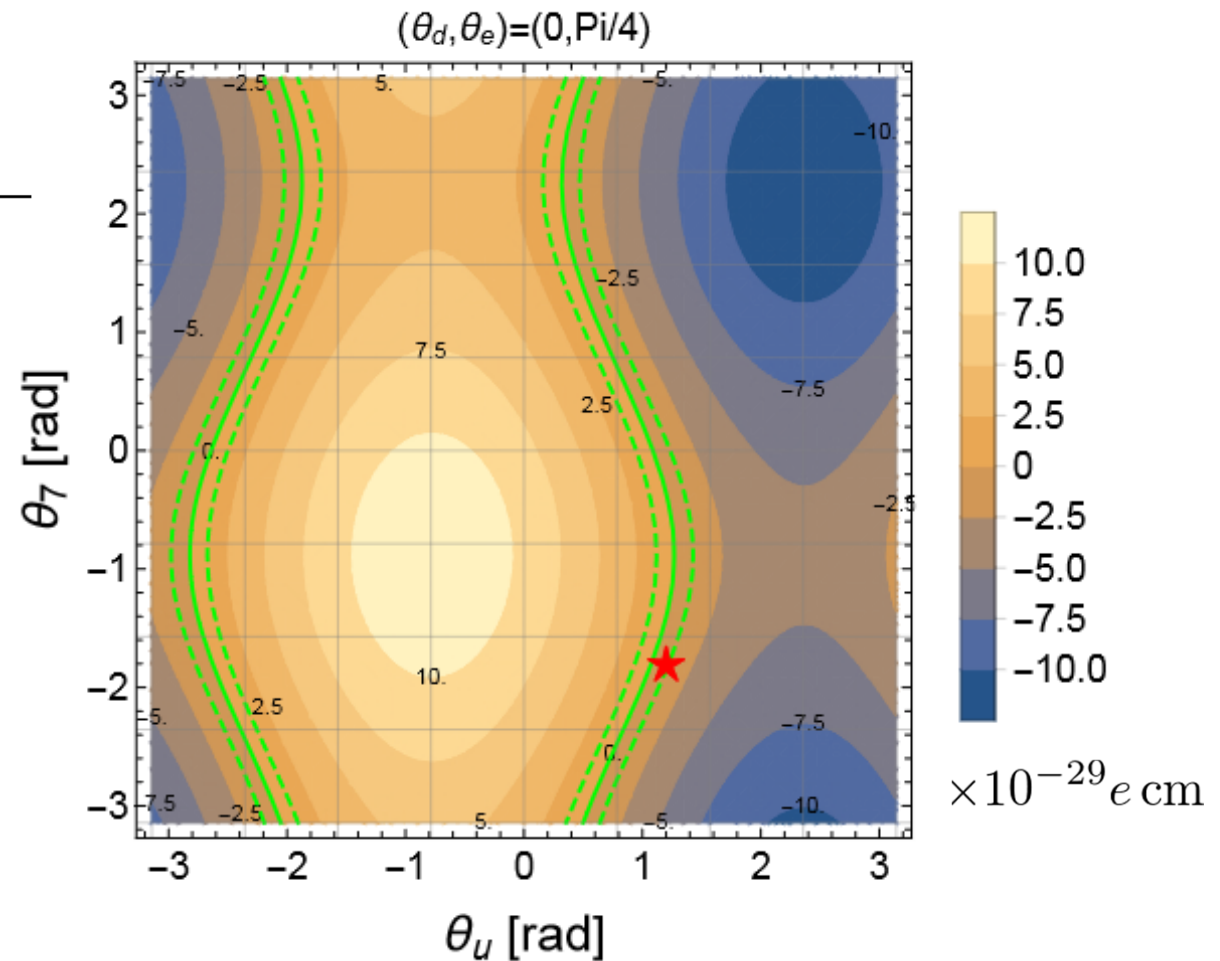
Destructive interference in EDM

Kanemura, Kubota, KY (2020)



• Benchmark point

$M = 240,$	$m_{H_2^0} = 280,$	$m_{H_3^0} = 230,$	$m_{H^\pm} = 230$	(in GeV).
$ \zeta_u = 0.01,$	$ \zeta_d = 0.1,$	$ \zeta_e = 0.5,$	$ \lambda_7 = 0.3,$	$\lambda_2 = 0.5.$
$\theta_u = 1.2,$	$\theta_d = 0,$	$\theta_e = \pi/4,$	$\theta_7 = -1.8$	(in rad).

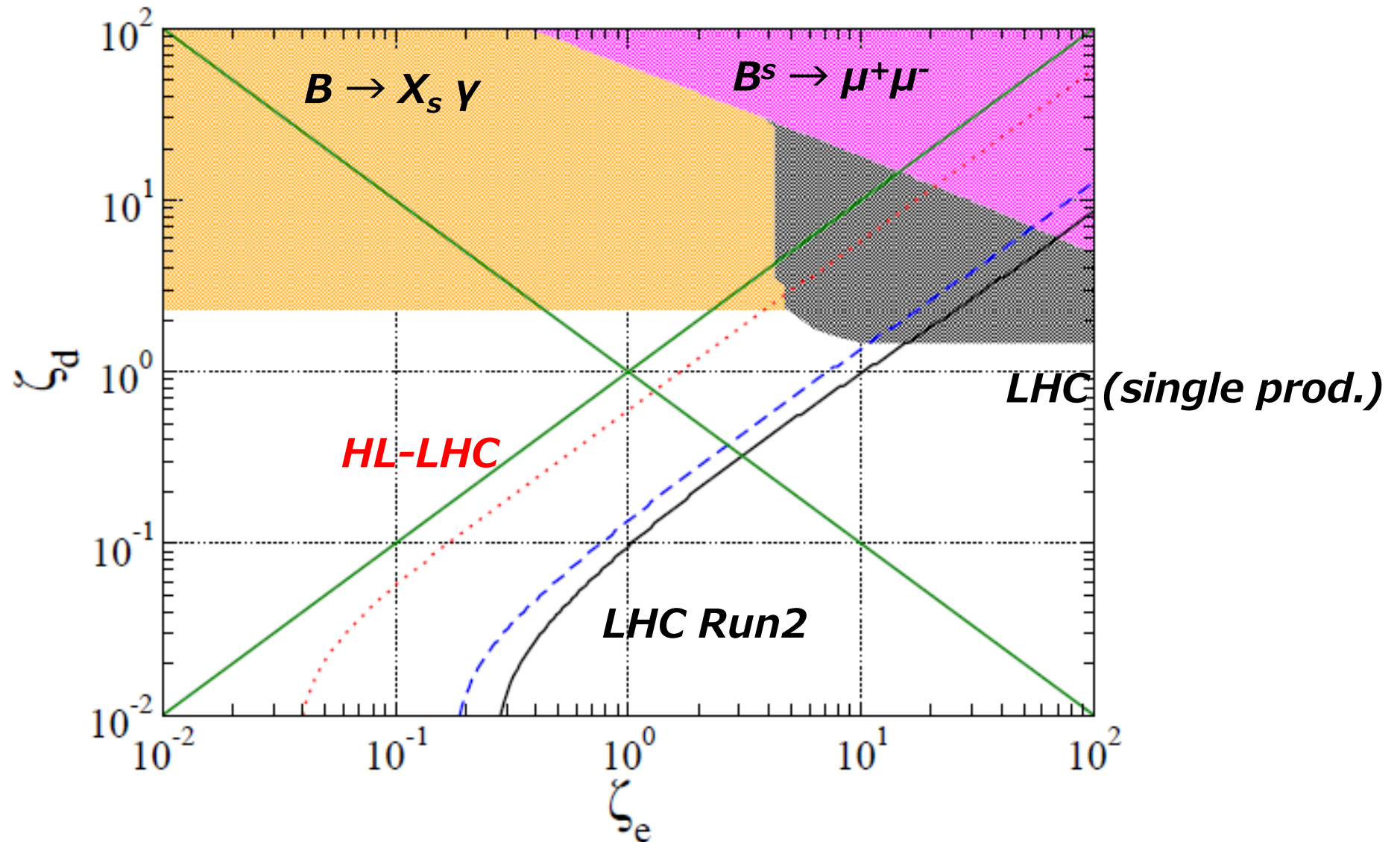


This opens a new possibility of the successful EW baryogenesis scenario.

Collider phenomenology

Kanemura, Takeuchi, KY, in preparation

$$m_{H_2} = 230 \text{ GeV}, m_{H_3} = 280 \text{ GeV}, m_{H^+} = 280 \text{ GeV}, \zeta_u = 0.1$$



Summary

- Higgs physics is a key to probe new physics BSM.
- The Higgs mass and top mass may indicate **criticality**.
- **EW scale** and **dark matter** can naturally be explained in the **MPP** framework.
This can be probed by **GW observations** at DECIGO and/or BBO, and
a light Higgs boson search at the **ILC**.
- Now, **“alignment”** is the important keyword.
- Near alignment : **Higgs to Higgs decay** & **Higgs precision** are important.
- Double alignment : **EW baryogenesis** is possible via the O(1) CPV phase.