



UNIVERSITY OF  
TEXAS  
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ICECUBE  
South Pole Neutrino Observatory



# Beyond-Standard-Model Neutrino Oscillations Analyses in IceCube

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PARALLEL SESSION 2: ASTROPHYSICAL NEUTRINOS



# Overview

## 1. A quick introduction to IceCube

- Detector design
- Signal principles
- Atmospheric neutrinos

## 2. BSM oscillations studies

- Heavy neutral leptons
- Neutrino decoherence
- Neutrino-matter non-standard interactions
  - High-Energy
  - Low-Energy

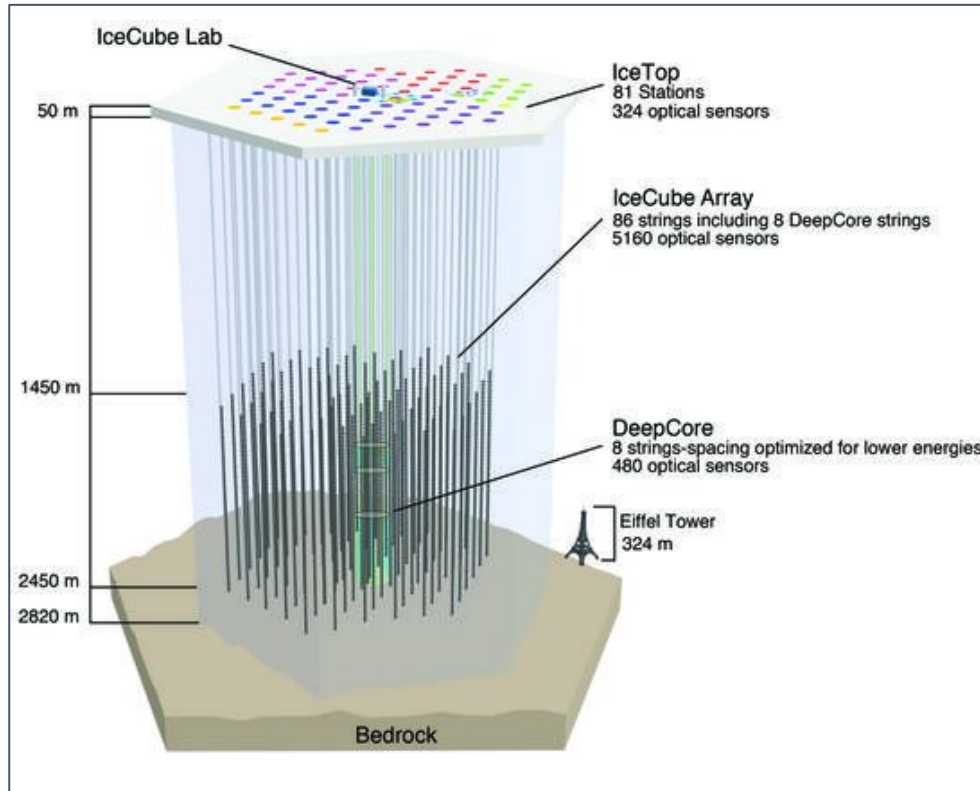
## 3. Concluding Remarks



# Detector Principles



# The IceCube Neutrino Observatory

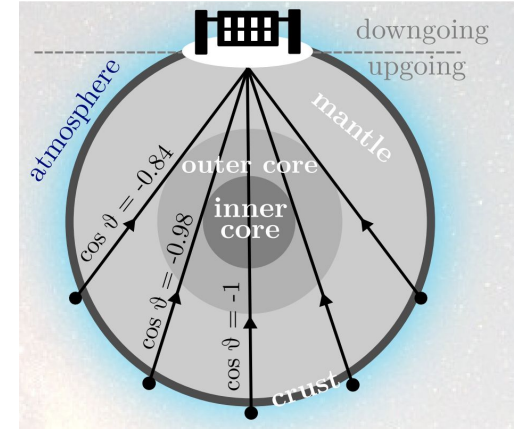


**Figure:** The IceCube neutrino observatory. DeepCore is an additional collection of center strings that allow for signals as low as 5 GeV. Credit: IceCube Collaboration gallery.

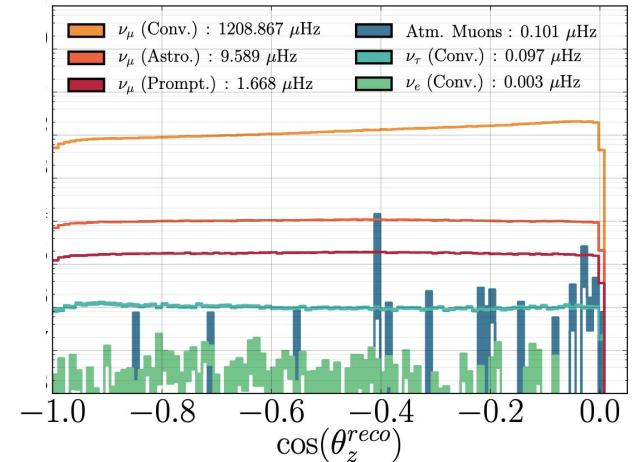
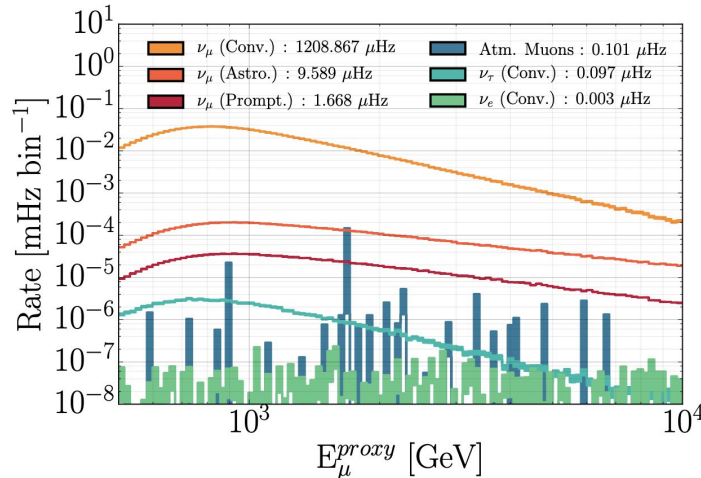
- Gigaton-scale, 1 km<sup>3</sup> volume Cherenkov detector and telescope, located within the South Pole glacial ice.
  - 78 main detector strings (125 m string spacing) with 60 digital optical modules (DOMs) each (17 m DOM spacing)
  - 8 additional center strings of higher efficiency DOMs (~40% improvement), 7 m DOM spacing
- Detection principle:
  - Neutrino-matter interaction products with sufficient energy produce Cherenkov (photon) radiation. Light collected by DOMs allows for neutrino direction and energy reconstruction.
- Oscillation analyses study energy and zenith dependence of both tracks (muons, linear trajectories) and cascades (EM/hadronic, blob-like signal).

# Atmospheric Neutrinos as Data

- Oscillations studies often require large baselines and high statistics. IceCube meets both of these requirements by detecting atmospheric neutrinos generated across Earth's surface.
- BSM signals often rely on matter effects, in which case the Earth provides a suitably large target medium. IceCube atmospheric neutrino samples frequently consist of “upgoing” signals.
- Additional IceCube advantages are the sensitivity to high event energies and possible access to tau appearance.



**Figure:** (Bottom) Signal and background distributions at IceCube binned in muon energy (left) and zenith (right) [Aartsen et al 2020]. (Top Right) Diagram demonstrating how atmospheric neutrinos can travel to IceCube. Credit: E. Lohfink for the IceCube collaboration.







# Beyond-Standard-Model Analyses

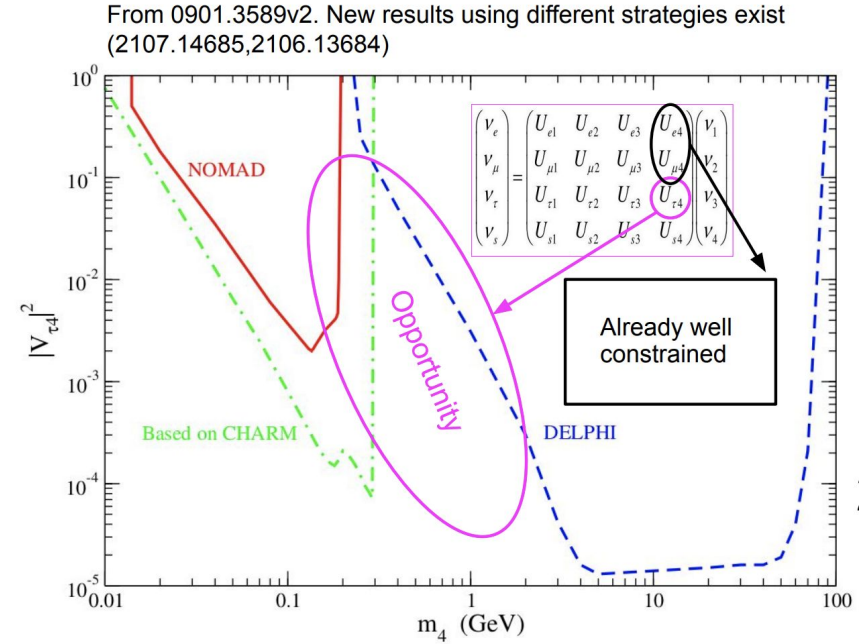


# Heavy Neutral Leptons



# Heavy Neutral Leptons

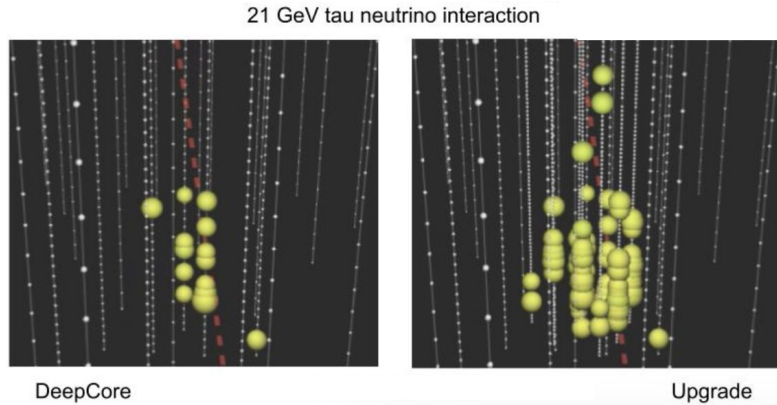
- Previous IceCube sterile searches have assumed light-masses, leading to strong constraints in  $U_{e4}$  and  $U_{\mu 4}$ .
- Lacking constraints on tau-sterile mixing ( $U_{\tau 4}$ ) allows for a larger sterile (heavy neutral lepton, HNL)
- A possible signature of a heavy neutral lepton decay would appear similar to the double-bang cascade signal of  $\nu_\tau$ -CC interactions.
- Depending on the HNL mass, the HNL lifetime may be longer than that of the tau lepton, allowing for topological discrimination of signals.



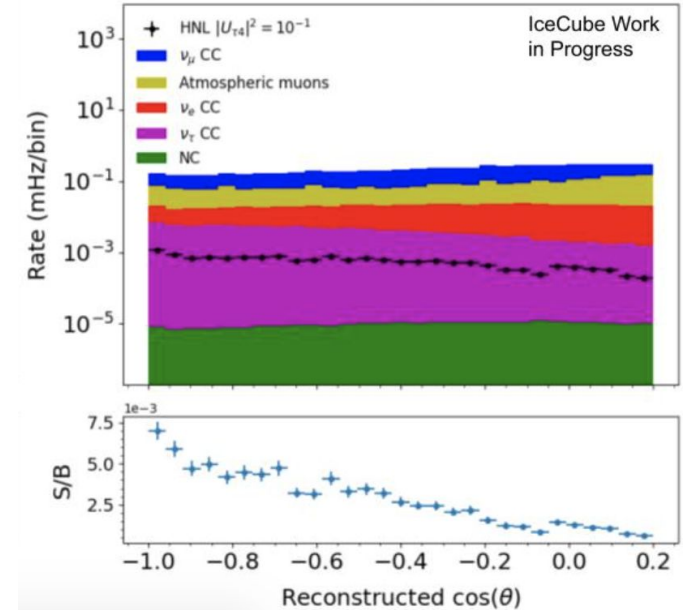


# Heavy Neutral Leptons

- Recently, a low-energy event selection was paired with a boosted decision tree (BDT) discriminator to determine the a sensitivity.
- Simulation found that current IC array is not sensitive to HNL decay due to heavy backgrounds and poor low-energy cascade reconstruction (right figure).
- IceCube Upgrade (additional 7 strings of specialized DOMs) will push us into HNL sensitivity from improved low-energy cascade reconstruction (below figure):



**Figure:** Comparison of a tau neutrino event in DeepCore (left) vs DeepCore+Upgrade (right).  
Credit: Summer Blot for the IceCube Collaboration



**Figure:** Comparison of simulated HNL decay event rate to IceCube backgrounds. Credit: David Vannerom for the IceCube Collaboration.

# Oscillation Decoherence



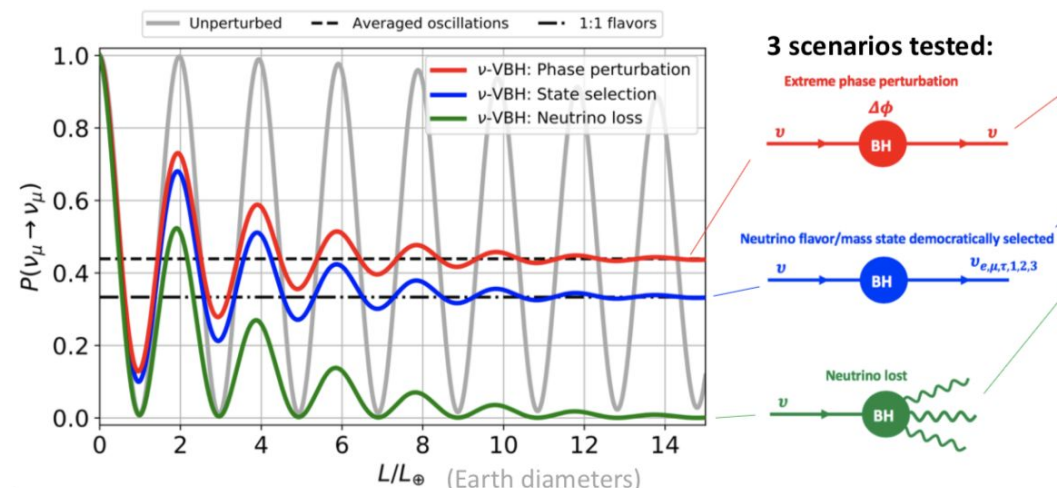


# Oscillation Decoherence Analyses

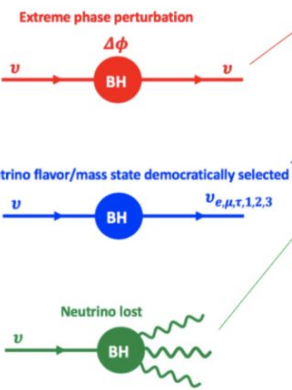
- Spacetime fluctuations at the quantum scale may allow for virtual black hole (VBH) pair production, a form of quantum foam.
- We call the subsequent loss of coherent SM oscillations **decoherence**.

- Decoherence can be introduced as an operator in the quantum system formalism below:

$$\dot{\rho} = -i[H, \rho] - \mathcal{D}[\rho] \implies \mathcal{D}[\rho] = (D_{\mu\nu} \rho^\nu) b^\mu$$



## 3 scenarios tested:



$$D_{\text{phase perturbation}} = \text{diag} (0 \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma)$$

$$D_{\text{state selected}} = \text{diag} (0 \ \Gamma \ \Gamma \ 0 \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ 0)$$

$$D_{\text{neutrino lost}} = \text{diag} (\Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma \ \Gamma)$$

**Figure:** (Left) Modified neutrino survival probabilities as a function of coherence length for the 3 tested decoherence scenarios. (Center) Diagrammed neutrino-VBH interactions for the 3 scenarios. (Right) The operator matrix representation for each scenario. Credit: T. Stuttard for the IceCube collaboration [arxiv:2007.00068].

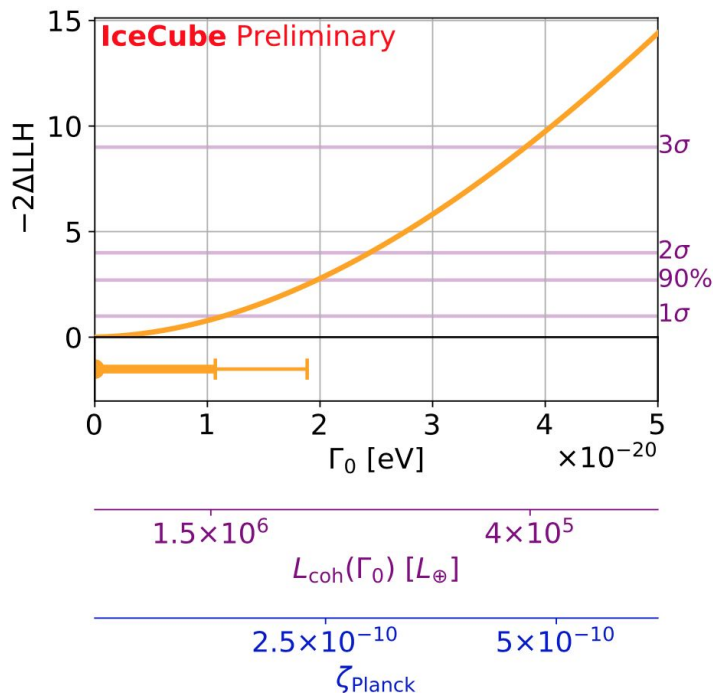
# Oscillation Decoherence Analyses

- Coherence length  $L_{coh}$  is the mean free path, defined as  $1/\Gamma$ , where  $\Gamma$  is the decoherence damping factor.

$$\Gamma(E) = \zeta_{\text{Planck}} \frac{E^n}{M_{\text{Planck}}^{n-1}}$$

with  $\zeta_{\text{Planck}}$  characterizing the Planck-scale coherence length.

- Two** analyses underway:
  - Low energy (~ 5 GeV - 300 GeV) with DeepCore 8-year sample of tracks + cascades.
  - High energy (~ 0.5 TeV - 10 TeV) using 2020 8-year sterile analysis track sample.
  - The complementarity of these analyses allows for a big-picture view of  $\Gamma$  and  $n$ -index values across a wide energy spectrum.



**Figure:** IceCube sensitivity to  $\Gamma_0$ ,  $L_{coh}$ , and  $\zeta_{\text{Planck}}$  for an  $n=0$  in the “democratic selection” scenario.  
[\[Stuttard T., Jensen M. 2020; arxiv:2007.00068\]](#)



# Nonstandard Interactions



# Brief NSI Review

- Neutrino masses are a BSM phenomenon:
  - Neutrino masses can be added through the dim-5 Wolfenstein operator.
  - Dim-6 operators may contribute new interactions scaled by a characteristic energy.
- IceCube is sensitive to *propagation* NSI (as opposed to *production* NSI).
- NSI are modeled by generalized coupling strengths corresponding to each flavor-violating (off-diagonal) and flavor-conserving (diagonal) instance (lower-right equation).

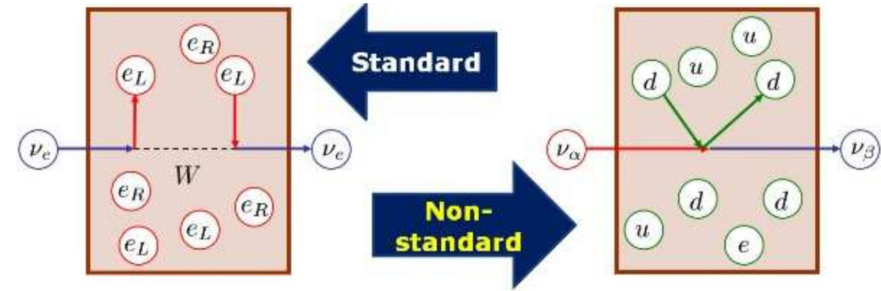


Figure: Diagramed neutrino-matter interactions for SM and NSI [Ohlson 2013 [arXiv:1209.2710](https://arxiv.org/abs/1209.2710)].

$$H_{\text{mat+NSI}} = V_{CC}(x) \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix}$$

Equation: Matter+NSI Hamiltonian. Each epsilon parameter is a weighted sum of the NSI contributions from the individual matter constituents.



# DeepCore Analyses

- Recently: All-parameter constraints using 3-years of 5 - 300 GeV events [arxiv:2106.07755], a first among NSI searches.

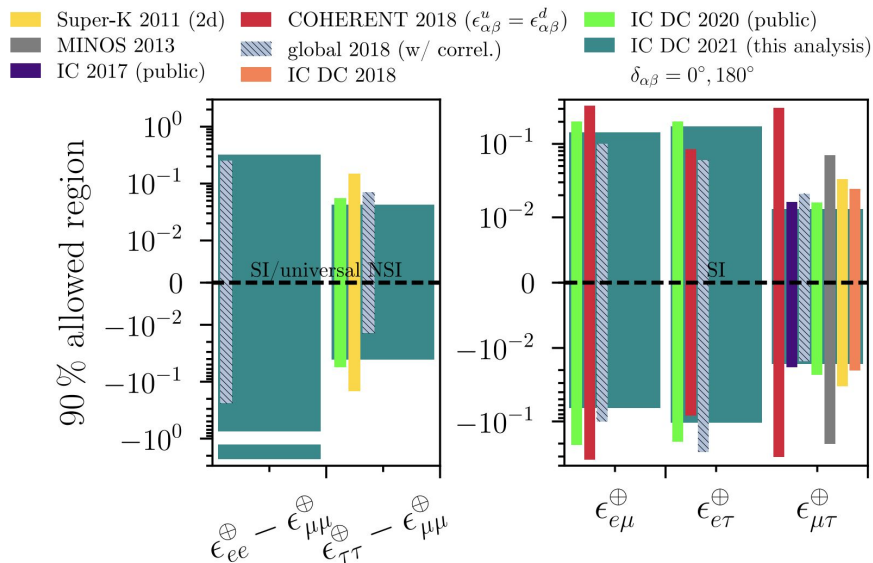


Figure: Results of the 3-year low-energy NSI analysis [arxiv:2106.07755].

- Upcoming: All-parameter constraints with 8-years of DeepCore data (6-fold increase in sample size).

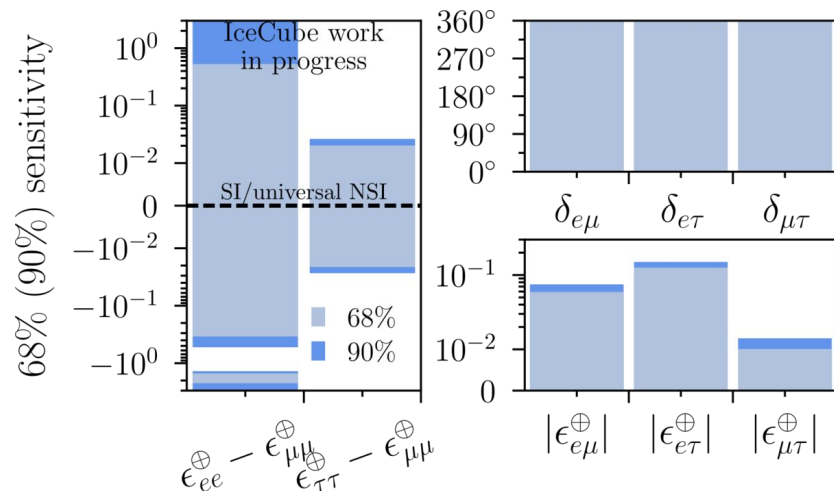


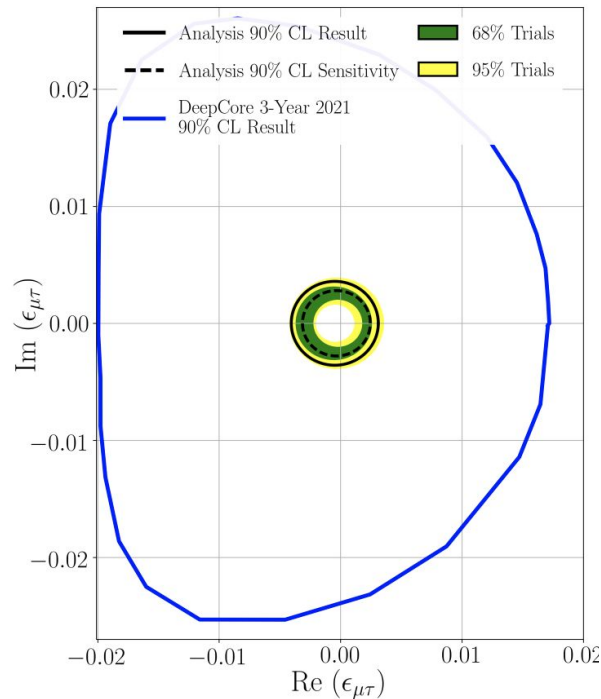
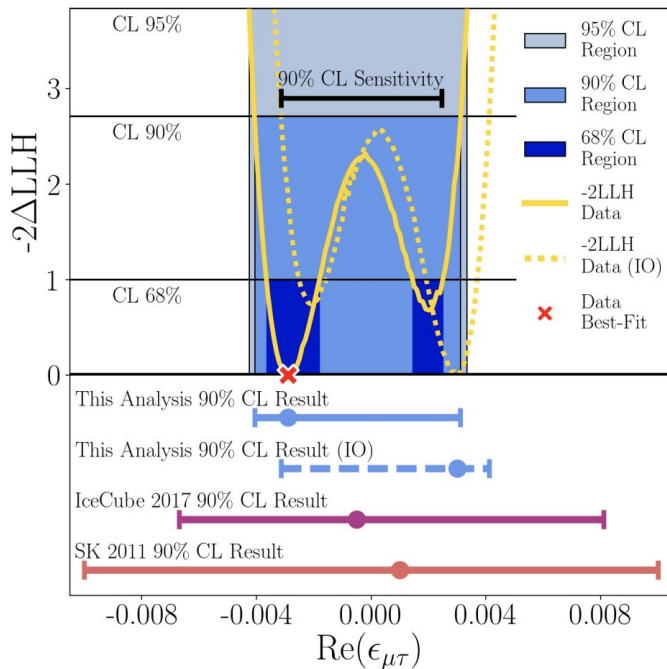
Figure: Sensitivities for the 8-year low-energy NSI analysis. Credit: E. Lohfink for the IceCube collaboration.

# TeV-Scale Analysis

- Uses 2020 8-year sterile sample of atmospheric neutrinos, fits only  $\epsilon_{\mu\tau}$ .
- Set leading real-valued and complex constraints ([arxiv:2201.03566](https://arxiv.org/abs/2201.03566))

- Left figure displays the fit to real-valued NSI in the top panel, while the lower panel demonstrates the improvement in 90% CL limits.

- Right figure shows the complex 90% CL result compared to the DeepCore 3-year complex result.





# Concluding Remarks

- 10-years of data from the full IceCube array has opened significant constraint opportunities in BSM oscillations studies.
- With Earth as a variable matter baseline, IceCube has many opportunities to explore a vast range of BSM oscillation models and accomplish:
  - Leading measurements of sterile neutrino parameters after extensive work in data filtering, data stability, and nuisance parameter studies;
  - New sterile decay parameter constraints;
  - A new probe of quantum foam models;
  - World-leading NSI parameter constraints.







Thank You



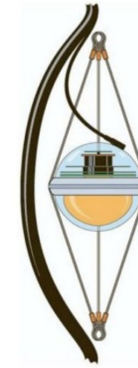
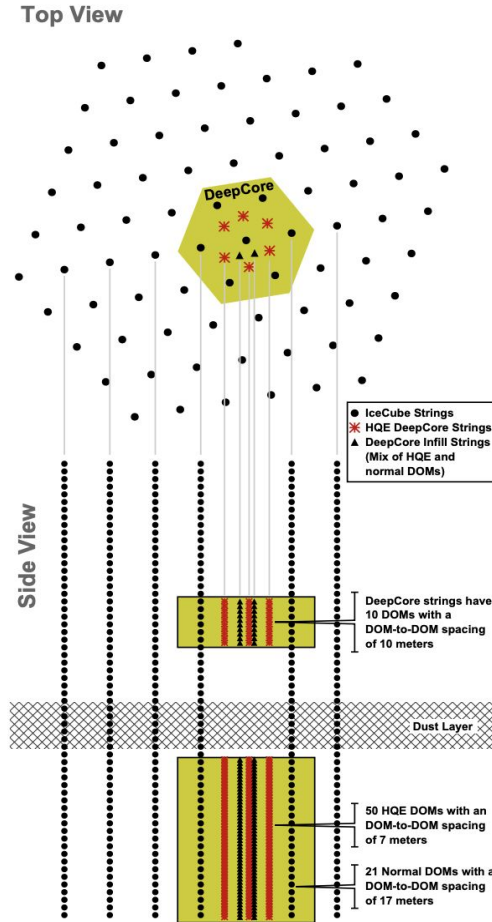
# Additional Material



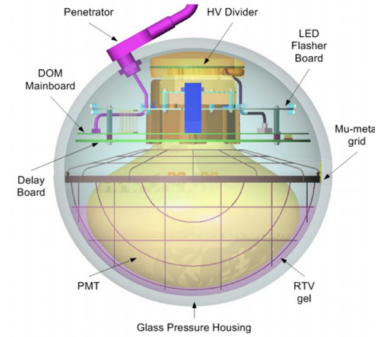
# DeepCore Design

DeepCore consists of 8 additional strings (red stars) of specialized DOMs placed within the center of the array.

The dust layer (2000 - 2100m) presents unknown absorption and scattering parameters for detected event light. This adds a systematic that is accounted for in IceCube analyses.



Cable: many twisted pairs, each pair carries power & communications for 2 DOMs



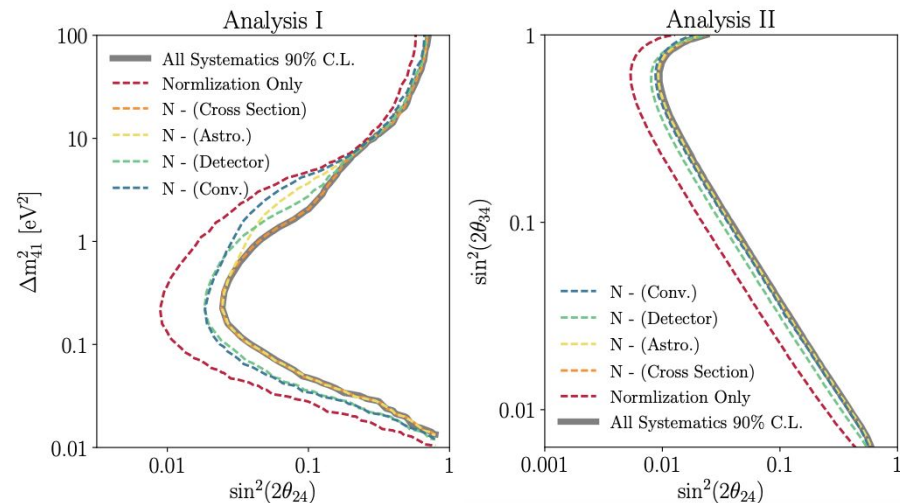
“High Quantum Efficiency” DOMs have ~40% improvement in optical sensitivity

DeepCore lowers the event energy threshold to ~10 GeV.



# Studies of Nuisance Parameters

- **Bulk Ice:** Scattering and absorption of glacial ice at different depths has wide uncertainty.
  - New method developed to study correlation effects via perturbed central MC
- **Hole Ice:** Ice in sensor column has different scattering and absorption.
  - Broken into two parameters-- central MC selected after testing different parameter values on angular acceptance
- **DOM Efficiency:** Optimal module photon detection efficiency.
  - 5 distinct DOM efficiency sets were used to create a splined continuous parameterization fit to zenith and energy reconstructions
- **Atmospheric Neutrino Flux:** Neutrino production factors and uncertainties.
  - Switched from discrete CR + hadronic models to continuous parameterizations of physically-motivated variables (Barr hadronic parameterization, CR spectral slope, conventional normalization, kaon-nuclei x-section, atmospheric density)
- **Cosmic Ray and Neutrino Flux:** The cosmic ray spectral index and neutrino flux normalization have parameterized uncertainties.
  - Normalization and spectral index with correlated uncertainties to fit IceCube's measurements



**Figure:** 90% CL sensitivities for analyses I and II of the 8-year sterile analysis with individual nuisance parameter groups removed. Grey is the original sensitivity [[Aartsen et al 2020](#)]



# NSI Formalism

- NSI's come in two types:
  - Neutral-Current (NC)
  - Charged-Current (CC)
- In IceCube we only look for NC NSI:
  - NC affect propagation, so we use the Earth as a large matter filter to search for NSI signals from “upgoing” atmospheric neutrinos that reach IceCube.

Start with the NC NSI Lagrangian

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \sum_{f,P,\alpha,\beta} \varepsilon_{\alpha\beta}^{f,P} (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{f} \gamma_\mu P f)$$

Generalize the parameters


$$\varepsilon_{\alpha\beta} = \sum_{f,P} \varepsilon_{\alpha\beta}^{fP} \frac{N_f}{N_e}$$


Add the NSI Hamiltonian to the matter Hamiltonian.

$$H_{\text{mat}} = \sqrt{2}G_F N_e(x) \begin{pmatrix} 1 + \varepsilon_{ee}(x) & \varepsilon_{e\mu}(x) & \varepsilon_{e\tau}(x) \\ \varepsilon_{e\mu}^*(x) & \varepsilon_{\mu\mu}(x) & \varepsilon_{\mu\tau}(x) \\ \varepsilon_{e\tau}^*(x) & \varepsilon_{\mu\tau}^*(x) & \varepsilon_{\tau\tau}(x) \end{pmatrix}$$



# Symmetry Between +Im(NSI) and -Im(NSI)

- For a given Re, +Im and -Im are identical in the two-neutrino calculation (right)
- Therefore, we can:
  - 1: Scan only the Re axis to get all the information
  - 2: Calculate statistics object for 1 DoF

$$P(+) = \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 2(a + bi)V_d \right|^2 \left( \frac{L}{2} \right)^2$$

$$P(-) = \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 2(a - bi)V_d \right|^2 \left( \frac{L}{2} \right)^2$$

$$P(+)/P(-) = \frac{\left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 2(a + bi)V_d \right|^2}{\left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 2(a - bi)V_d \right|^2}$$

$$= \frac{\left( \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} \right)^2 + 4(a)V_d \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 4(a^2 + b^2)V_d^2}{\left( \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} \right)^2 + 4(a)V_d \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 4(a^2 + b^2)V_d^2}$$
$$= 1$$

# NSI LLH Offset

- At sample energies, the electron flavor state decouples from atmospheric oscillation.
- Below is the approximate calculation of the difference in probabilities for NSI values with equal imaginary components, opposite-sign real components.

$$\begin{aligned} P(-) - P(+) &= \\ \left[ \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 2(-a + bi)V_d \right|^2 - \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_\nu} + 2(a + bi)V_d \right|^2 \right] \left( \frac{L}{2} \right)^2 \\ &= -aV_d L^2 \sin 2\theta_{23} \frac{\Delta m_{31}^2}{E_\nu} \\ &= \text{negative value} \end{aligned}$$

This confirms what we see in the  $-2\Delta\text{LLH}$  distribution, as  $-2\Delta\text{LLH}(-a + bi) < -2\Delta\text{LLH}(a + bi)$ .