

Extracting the nucleon axial form factor from Lattice QCD using chiral perturbation theory

Fernando Alvarado (falvar@ific.uv.es)

Luis Alvarez-Ruso

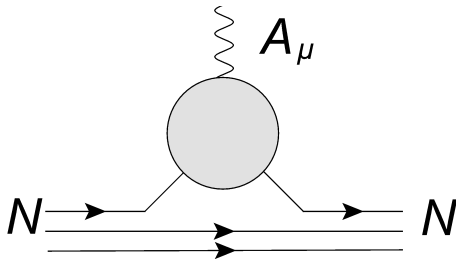


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Nucleon Axial Form Factor

- ▶ Nucleon Axial Form Factor, $F_A(q^2)$
 - ▶ Electroweak interactions open a doorway to fundamental properties of strong interacting matter: spins distribution
 - ▶ $A_\mu^i(x) = \bar{q}(x)\gamma_\mu\gamma_5\frac{\tau^i}{2}q(x)$
 - ▶ $\langle N(p')|A_\mu^i|N(p)\rangle = \bar{u}\left\{\gamma_\mu F_A(q^2) + \frac{q_\mu}{2m_N}G_P(q^2)\right\}\gamma_5\frac{\tau^i}{2}u(p)$
- ▶
$$F_A(q^2) = g_A\left[1 + \frac{1}{6}\langle r_A^2\rangle q^2 + \mathcal{O}(q^4)\right]$$
- ▶ g_A and F_A dependence in q^2 are necessary in ν oscillations experiments
- ▶ μ capture, β -decay
- ▶ Chiral Perturbation Theory calculation of F_A
 \Rightarrow extract $\langle r_A^2\rangle$ from lattice QCD without ad-hoc parametrization



Axial form factor, F_A

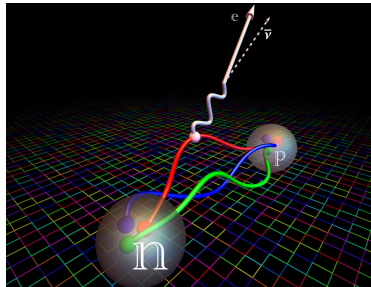
► Empirical determinations

- Rely on neutrino-induced charged-current quasielastic scattering on deuteron targets, muon capture in muonic hydrogen and pion electro-production.

► LQCD

- Several studies on $F_A(q^2) \rightarrow$ technical difficulties
 \Rightarrow significantly improved control of the systematic error
- Tension between LQCD and empirical determinations
- Experimental and lattice q^2 parametrisation:

$$\left. \begin{array}{l} \text{- dipole ansatz} \\ \text{- z-expansion} \end{array} \right\} \Rightarrow \text{different } \langle r_A^2 \rangle$$



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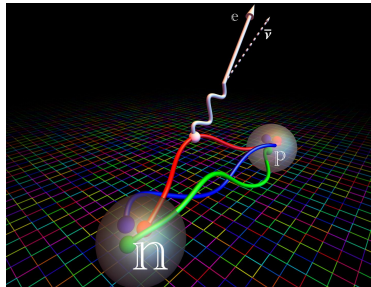
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► Chiral Perturbation Theory (χ PT)

- EFT for QCD at low energy
- QCD based parametrization of q^2 and M_π dependencies
 \Rightarrow extrapolate lattice results to the phys. point and extract $\langle r_A^2 \rangle$ from the lattice simulations
- Account for finite volume, lattice spacing and excited states
- Determining χ PT LECs from the lattice
 \Rightarrow predicting other observables



F_A Calculation

- ▶ NNLO $\mathcal{O}(p^4)$ in relativistic Baryon χ PT
- ▶ Baryon χ PT
 - ▶ Problem: $\underbrace{m_B \rightarrow 0}_{\chi \text{ limit}} \Rightarrow$ Power Counting Breaking (PCB)
 - ▶ \Rightarrow additional finite renormalisation: extended on mass-shell (EOMS)
 - ▶ PCB terms absorbed by LECs
 - ▶ Covariance and analytic properties of loops preserved \Rightarrow appropriate for chiral extrapolations

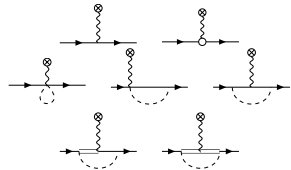


Figure: $\mathcal{O}(p)$ and $\mathcal{O}(p^3)$ (w. f. renormalisation not shown)

- ▶ Explicit $\Delta(1232)$
 - ▶ SSE: $\delta = m_\Delta - m_N \sim \mathcal{O}(p)$
- ▶ $F_A = \hat{g}_A + 4d_{16}M_\pi^2 + d_{22}t + \text{loops}(M_\pi, t)$
 - ▶ $\mathcal{L}_{\pi N}^{(1)} \Rightarrow \hat{g}_A, \quad \mathcal{L}_{\pi N}^{(3)} \Rightarrow d_{16},$
 $\mathcal{L}_{\pi N}^{(2)} \Rightarrow c_1, c_2, c_3, c_4$
 - ▶ $\mathcal{L}_{\pi N\Delta}^{(1)} \Rightarrow h_A, g_1,$
 $\mathcal{L}_{\pi N\Delta}^{(2)} \Rightarrow a_1, \quad \mathcal{L}_{\pi N\Delta}^{(2)} \Rightarrow b_4, b_5$

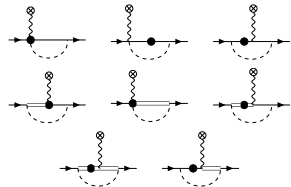
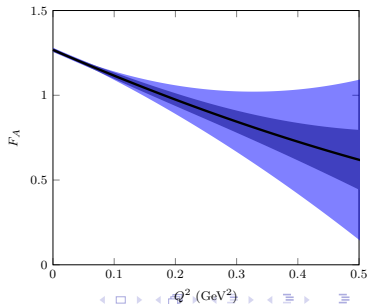
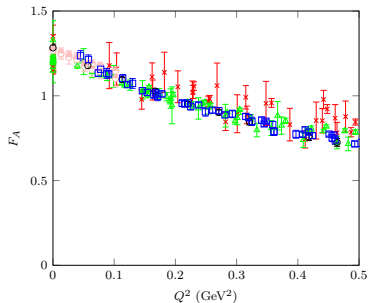


Figure: $\mathcal{O}(p^4)$

Combined fit to lattice data

► Lattice data

- Many recent works \Rightarrow substantial improvements
- RQCD^[1] + PNDME^[2] + "Mainz"^[3] + PACS^[4] + ETMC^[5]
- data without q^2 , finite volume, lattice spacing or M_π extrapolation
- large vol. only, $M_\pi L \geq 3.5$
- we correct lattice spacing a :
 $F_A(a) = F_A + \sum_i (x_i + t y_i) a^{n_i}$



[1] Bali et al. JHEP 05 (2020)

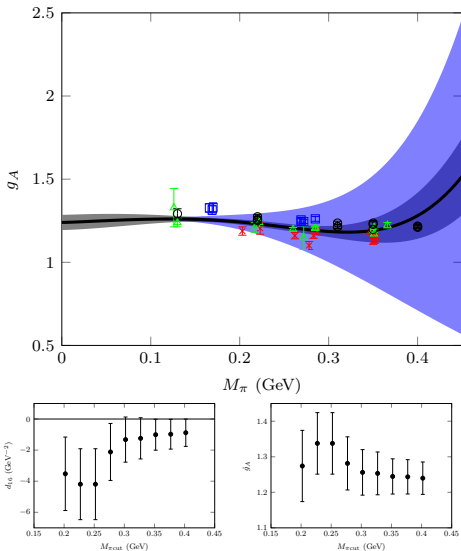
[2] Park et al. 2103.05599

[3] Meyer et al. Modern Phys. A 34 (2019)

[4] Shintani et al. PRD 102 (2020)

[5] Alexandrou et al. PRD 103 (2021)

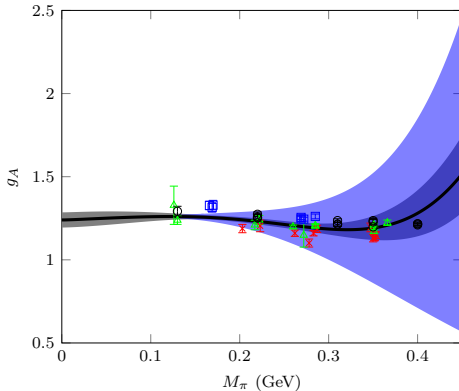
$$F_A(q^2 = 0) = g_A$$



- ▶ $g_A(M_\pi)$: interesting puzzle by itself
 - ▶ we saw that Δ LECs from πN elastic and inelastic scattering fail to describe its M_π dependence [Alvarado & Alvarez-Ruso PRD 105 \(2021\)](#)
- ▶ Differences between $\mathcal{O}(p^3)$ and $\mathcal{O}(p^4)$ are considerable (at larger M_π) and provide a measure of the systematic error [6] arising from the truncation of the perturbative expansion:

$$\Delta g_{A\chi}^{(4)} = \max \left\{ \left(\frac{M_\pi}{\Lambda} \right)^4 |g_A^{(4)}|, \left(\frac{M_\pi}{\Lambda} \right)^2 |g_A^{(3)}|, \frac{M_\pi}{\Lambda} |g_A^{(4)}| \right\}$$
- ▶ $\Delta F_{A\chi}$ is added to LQCD errors in the χ^2
- ▶ LECs have naturalness priors
- ▶ χ^2 plateau $\Rightarrow M_\pi^{\text{cut}} \simeq 400$ MeV, $Q_{\text{cut}}^2 = 0.36$ GeV²
- ▶ Δ baryon is a necessary d.o.f.
- ▶ **Good fit:**
very accurate description at the physical point
- ▶ $\mathcal{O}(p^5)$ still needed for full convergence

$$F_A(q^2 = 0) = g_A$$



► Axial charge results from $F_A(q^2)$ fit

- $g_A(M_{\pi\text{phys}}) = 1.273 \pm 0.014$
vs $g_A^{\text{exp}} = 1.2754(13)_{\text{exp}(2)_{\text{RC}}} \Rightarrow$ **excellent agreement with exp.**
vs $g_A^{\text{FLAG}} = 1.246 \pm 0.028$

► $g_A(M_\pi) = \dot{g}_A + 4d_{16}M_\pi^2 + \text{loop}(M_\pi)$

► $d_{16} = -1.46 \pm 1.00 \text{ GeV}^{-2}$

→ M_π dependence of long range nuclear forces

- Can not be extracted from πN elastic scattering
- In line with $d_{16} = -1.0 \pm 1.0 \text{ GeV}^{-2}$ from $\pi N \rightarrow \pi\pi N$ [6]

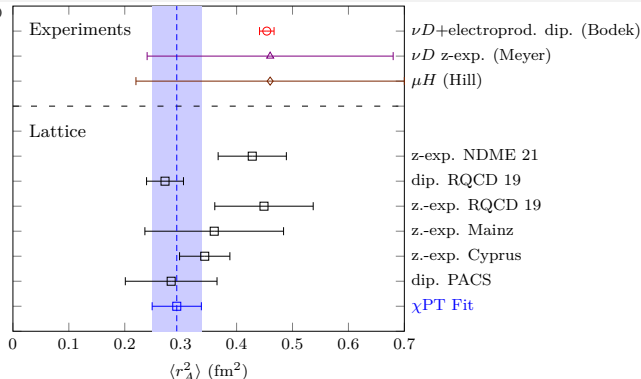
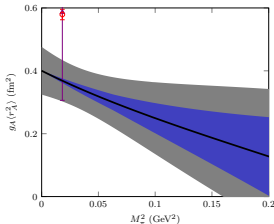
[6] Siemens et al. PRC 96 (2017)
(value converted to standard EOMS)

$$F_A = g_A \left(1 + \frac{1}{6} \langle r_A^2 \rangle q^2 \right) \text{ axial radius}$$

(Bodek)=Bodek, Eur. Phys. J. C 53, 349 (2008)

(Meyer)=Meyer, PRD 93, 113015 (2016)

(Hill)=Hill, Rept. Prog. Phys. 81 (2018)



► Our $\mathcal{O}(p^4)$ χ PT extraction:

- M_π slope driven by loops with Δ
- $d_{22} = 0.29 \pm 1.69 \text{ GeV}^{-2}$ (no assumptions on $\Delta\Delta\pi$ coupling enlarges error)
 - d_{22} compatible with $\mathcal{O}(p^3)$ π electroprod.

$$\langle r_A^2 \rangle(M_{\text{phys}}) = 0.293 \pm 0.044 \text{ fm}^2$$

- Empirical determinations (model dependent) are in **tension** with ours and with most of LQCD extractions
- Typically the extracted $\langle r_A^2 \rangle^{\text{phys}}$ value varies depending on the parametrisation
- Our QCD based parametrisation leads to a value in line with most of the individual LQCD extractions

Conclusions

- ▶ $F_A(q^2)$ essential in ν oscillations
- ▶ We extract $F_A(q^2)$ from LQCD using $\mathcal{O}(p^4)$ relativistic χ PT
- ▶ Our combined fit $\mathcal{O}(p^4)$ with Δ successfully describes the lattice data
 - ▶ Δ is a necessary d.o.f.
 - ▶ $g_A(M_{\pi\text{phys}}) = 1.273 \pm 0.014$ vs $g_A^{\text{exp}} = 1.2754(13)_{\text{exp}}(2)_{\text{RC}} \Rightarrow$ excellent agreement with exp.
 - ▶ There is tension between the experimental and lattice extraction of $\langle r_A^2 \rangle$
 - ▶ We extract $\langle r_A^2 \rangle^{\text{phys}} = 0.291 \pm 0.052 \text{ fm}^2$ without ad hoc parametrisations
- ▶ $d_{16} = -1.46 \pm 1.00 \text{ GeV}^{-2}$, $d_{22} = 0.29 \pm 1.69 \text{ GeV}^{-2}$ and other LECs have been extracted \Rightarrow agreement with different determinations at the physical point

Thanks!
Any questions?

Nucleon Axial Form Factor: Extra

- ▶ Dipole ansatz: $F_A(q^2) = g_A(1 - \frac{q^2}{M_A^2})^{-2}$
- ▶ z-exp.: $F_A(q^2) = \sum_k a_k z^k(q^2)$, with $z(q^2, t_{\text{cut}}, t_0)$

Extra

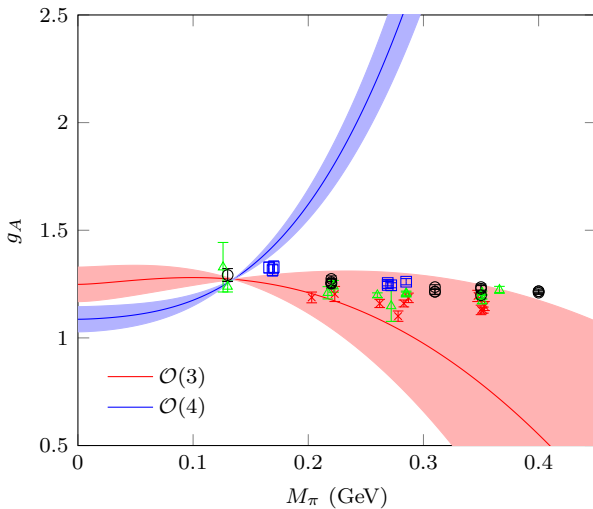


Figure: Pion-mass dependence of g_A at $\mathcal{O}(p^3)$ (red) and $\mathcal{O}(p^4)$ (blue) using phenomenological input from Ref. ? and 1σ error bands.

Extra

| | $\mathcal{O}(p^3) \Delta$ | $\mathcal{O}(p^4) \Delta$ | $\mathcal{O}(p^3) \Delta$ | $\mathcal{O}(p^4) \Delta$ |
|--|---------------------------|---------------------------|---------------------------|--|
| \tilde{g}_A (free) | 1.1782 ± 0.0073 | | 1.2041 ± 0.0074 | 1.274 ± 0.041 |
| d_{16} (GeV^{-2}) (free) | -1.021 ± 0.048 | | 0.983 ± 0.062 | -1.46 ± 1.00 |
| d_{22} (GeV^{-2}) (free) | 1.275 ± 0.086 | | 3.77 ± 1.96 | 0.29 ± 1.69 large error (free g_1) |
| h_A | - | - | 1.35 | 1.35 |
| g_1 (free) | - | - | -0.69 ± 0.69 | 0.66 ± 0.56 |
| c_1 (GeV^{-1}) | - | -0.89 ± 0.06 | - | -1.15 ± 0.05 |
| c_2 (GeV^{-1}) | - | 3.38 ± 0.15 | - | 1.57 ± 0.10 |
| c_3 (GeV^{-1}) | - | -4.59 ± 0.09 | - | -2.54 ± 0.05 |
| c_4 (GeV^{-1}) | - | 3.31 ± 0.13 | - | 2.61 ± 0.10 |
| a_1 (GeV^{-1}) | - | - | - | 0.90 |
| b_1 (GeV^{-2}) (free) | - | - | - | -0.27 ± 4.96 |
| b_2 (GeV^{-2}) (free) | - | - | - | 2.27 ± 2.28 |
| \tilde{b}_4 (GeV^{-2}) (free) | - | - | - | -12.48 ± 1.28 |
| x_1 (fm^{-2}) (free) | -8.4 ± 5.8 | - | -5.6 ± 5.9 | -0.25 ± 16.5 (consistent) |
| x_2 (fm^{-2}) (free) | -8.6 ± 2.6 | - | -7.1 ± 2.6 | -6.36 ± 4.20 |
| x_3 (fm^{-1}) (free) | -0.25 ± 0.21 | - | -0.08 ± 0.22 | 0.36 ± 0.47 |
| y_1 ($\text{fm}^{-2} \text{ GeV}^{-2}$) (free) | -100 ± 40 | - | -76 ± 44 | -64 ± 121 |
| y_2 ($\text{fm}^{-2} \text{ GeV}^{-2}$) (free) | -31 ± 21 | - | -21 ± 22 | -15 ± 46 |
| y_3 ($\text{fm}^{-1} \text{ GeV}^{-2}$) (free) | -0.63 ± 1.49 | - | 0.36 ± 1.63 | 2.54 ± 3.98 |
| \tilde{m} (GeV) | 0.874 | 0.874 | 0.855 | 0.855 |
| \tilde{m}_Δ (GeV) | - | - | 1.166 | 1.166 |
| χ^2/dof | $46.13/(127-9) = 0.391$ | | $39.17/(127-10) = 0.326$ | $14.64/(127-13) = 0.129$ |
| χ_0^2/dof | $857.31/(127-9) = 7.27$ | | $533.87/(127-10) = 4.45$ | $196.58/(127-13) = 1.724$ |