# Core-valence absorption in breakup and stripping reactions and its isospin dependence

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- Absorption in nuclear breakup reactions
- Results
- 3 Absorption in stripping reactions
- Conclusions and outlook

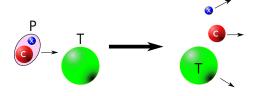


- Absorption in nuclear breakup reactions
- - ${}^{12}\text{C}({}^{11}\text{Be}, n^{10}\text{Be}){}^{12}\text{C}$  at 70 MeV/A
  - <sup>12</sup>C(<sup>41</sup>Ca,n<sup>40</sup>Ca)<sup>12</sup>C at 70 MeV/A



### Nuclear breakup reactions

- $\bullet$   $P(c+x)+T \rightarrow c+x+T$
- All particles are detected

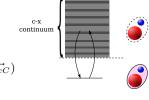


• Extensively used to analyze single-particle properties of nuclei, specially exotic nuclei



### Continuum-discretized coupled-channels

- State-of-the-art description of low- and mid-energy breakup reactions
- 3-body approximation of the full wavefunction, expanded in discretized C - x continuum



$$\Psi \simeq \Psi^{3b}(\vec{R_{PT}}, \vec{r_{xC}}) \simeq \sum_{nJ\pi} \chi_{nJ\pi}(\vec{R_{PT}}) \phi_{nJ\pi}(\vec{r_{xC}})$$
$$\phi_{nJ\pi}(k_n, \vec{r_{xC}}) = \sqrt{\frac{2}{\pi N}} \int_{k_n}^{k_n} \phi_{nJ\pi}(k, \vec{r_{xC}}) dk$$

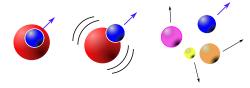
• Orthogonality of states used to solve equations:  $V_{xC}$  must be real

$$\sum_{nJ\pi} [(E - T - \epsilon_m) \underbrace{\langle \phi_{nJ\pi'} | \phi_{mJ\pi} \rangle}_{\delta_{nm}\delta_{J\pi J\pi'}} - \langle \phi_{nJ\pi'} | U_{xT} + U_{CT} | \phi_{mJ\pi} \rangle] \chi_{mJ\pi} = 0$$



### Absorption effects

- Imaginary parts of  $U_{xT}$  and  $U_{CT}$  describe absorption between x-T and C-T
- In the continuum, the interaction between x and C can excite C or x, which can then break up, removing flux



- $U_{xC}$  should be complex at positive energies, but then its eigenstates  $\varphi$  are no longer orthogonal!!!
- Binormal basis  $\tilde{\varphi}$  is orthogonal to a set of non-orthogonal states  $\varphi$

$$\sum_{nJ\pi} [(E - T - \epsilon_m) \underbrace{\langle \tilde{\varphi}_{nJ\pi'} | \varphi_{mJ\pi} \rangle}_{\delta_{nm} \delta_{J\pi J\pi'}} - \langle \tilde{\varphi}_{nJ\pi'} | U_{xT} + U_{CT} | \varphi_{mJ\pi} \rangle] \chi_{mJ\pi} = 0$$

$$\tilde{\varphi}_n^{(-)} \sim \varphi_n^{(+)*}$$

$$\tilde{\varphi}_i^{(-)} = \sum_i \mathcal{A}_{ij}^{-1} \varphi_j^{(+)}$$

$$\tilde{\varphi}_n^{(-)} \sim \varphi_n^{(+)*} \qquad \quad \tilde{\varphi}_i^{(-)} = \sum_j \mathcal{A}_{ij}^{-1} \varphi_j^{(+)*} \qquad \quad \mathcal{A}_{ji} = \left\langle \varphi_j^{(+)*} | \varphi_i^{(-)} \right\rangle$$



- Results



- Absorption in nuclear breakup reactions
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### Interaction potential

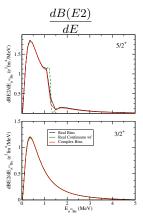
- Real part: Potential from Capel et al (PRC 70, 064605 (2004)), reproduces bound states and low-energy resonance
- Imag part: Adjusted to reproduce reaction cross sections for n-9Be (compilation by A. Bonaccorso and R.J. Charity PRC 89,024619 (2014)), rescaled through  $A^{2/3}$

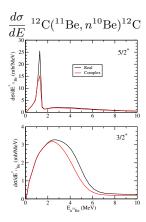
$$W(E,r) = \frac{W_0(E)}{1 + \exp{(r - R)/a_0}} \quad W_0(E) = \frac{(a(E - E_b) + b)E^4}{E^4 + E_b^4},$$

$$\frac{\frac{800}{200}}{\frac{2}{200}} \quad \frac{\frac{\alpha}{2}}{\frac{\alpha}{200}} \quad \frac{\frac{\alpha}{2}}{\frac{\alpha}{2}} \quad \frac{\frac{\alpha}{2}}{\frac{\alpha}$$



# $5/2^+$ and $3/2^+$ states

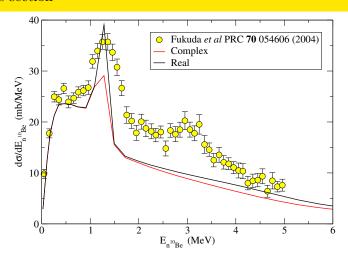




- Coulomb breakup barely affected by absorption (larger x-C distance)
- Resonances severely affected



#### Full cross section

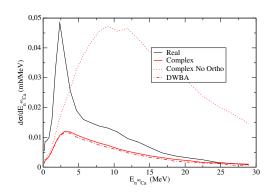


- Small effect of absorption:  $\sim 10\%$
- Resonance too severely affected (absorption threshold possibly too low)
- Core-excitation effects have been predicted for these data (A.M. M. and J.A. Lay PRL 109 232502 (2012)) but are not included here

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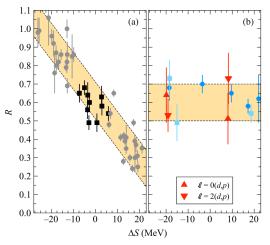


# $^{12}\text{C}(^{41}\text{Ca}, n^{40}\text{Ca})^{12}\text{C}$ at 70 MeV/A



- Much stronger effect:  $\sim 50\%$  reduction
- Breakup of more tightly bound nucleon explores higher energies with larger absorption, and there are more open channels
- Large effect of non-orthogonality with the ground state, use of  $\varphi^{(+)*}$  is not enough

## Implications?



B.P. Kay et al PRL 129 152501 (2022) J. A. Tostevin and A. Gade PRC 103 054610 (2021)

• Answer to open question?



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# Absorption in stripping

• In usual eikonal treatment uses closure to obtain a density of the bound nucleon

$$\rho(\mathbf{r_1},\mathbf{r_2}) = \phi_b^*(\mathbf{r_1})\phi_b(\mathbf{r_2}) \int \mathrm{d}\mathbf{k} \phi_{xC}^{\left(+\right)}(\mathbf{k},\mathbf{r_1})\phi_{xC}^{*\left(+\right)}(\mathbf{k},\mathbf{r_2}) = \delta(\mathbf{r_1} - \mathbf{r_2})\phi_b^*(\mathbf{r_1})\phi_b(\mathbf{r_2}) = \left|\phi_b(\mathbf{r_1})\right|^2$$

- This is only true for real, energy-independent  $V_{xC}$ .
- For absorptive potentials we can define an effective density for an average position

$$\rho^{\mathrm{eff}}(x,y) = \int \mathrm{d}\mathbf{r_1} \mathrm{d}\mathbf{r_2} \delta(x - \frac{x_1 + x_2}{2}) \delta(y - \sqrt{\frac{y_1^2 + y_2^2}{2}}) \phi_b^*(\mathbf{r_1}) \phi_b(\mathbf{r_2}) \int \mathrm{d}\mathbf{k} \phi_{xC}^{(+)}(\mathbf{k},\mathbf{r_1}) \phi_{xC}^{*(+)}(\mathbf{k},\mathbf{r_2})$$

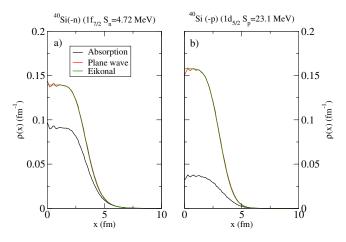
• This  $\rho^{\text{eff}}$  can be used in standard eikonal calculations

$$\sigma_{\text{str}} = \int d\mathbf{b} \int d\mathbf{b}_{\mathbf{V}\mathbf{C}} \rho^{\text{eff}}(x, y) |S_{CT}|^2 (1 - |S_{VT}|^2)$$



### Effective density

•  $U_{VC}$ : Imaginary part of Morillon potential (since we study absorption)

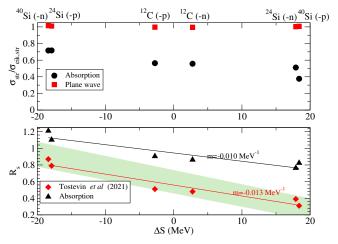


• Significant reduction, larger for deeply-bound nucleon



#### Effect on cross sections

Only computed for stripping, effect in diffraction assumed to be the same



- Small reduction in slope
- $R_s < 1$  for weakly-bound nucleons... problematic



### Elastic compound scattering

- Optical potential gives finite reaction cross section at low energies for weakly-bound nucleons (But there are no open channels!!!)
- This corresponds to compound nucleus which decays to elastic channel (This is not absorption) → Must be removed from potential
- Use compound-nucleus calculation (PACE4) to estimate and remove flux to elastic

$n+^{39}\mathrm{Si}$						$p+$ $^{\circ\circ}{ m Al}$ 1. Yields of residual nuclei							
	1	. Yiel	ds of r	esidual	nuclei	Z	N		Α	events	percent	x-section(mb	
Z	N	Α	events	percent	x-section(mb)	14	25	39	Si	13	1.3%	11.8	
14	25	39 Si	1000	100%	1.7e+03	14	24	38	Si	117	11.7%	106	
		55 51				14	23	37	SI	865	86.5%	786	
TOTAL			1000	100	1700.25	14	22	36	Si	5	0.5%	4.54	
						TOT	AI .			1000	100	908.47	

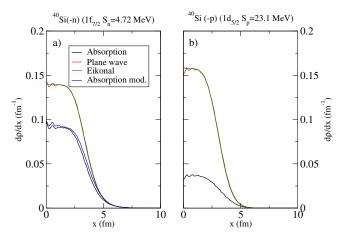
 Absorption unchanged for deeply-bound nucleons but severely reduced for weakly-bound at low energies



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### Effective density

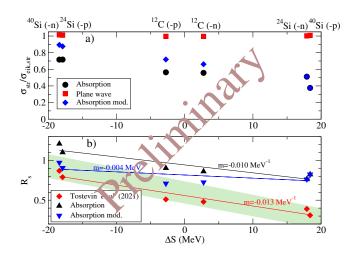
•  $U_{VC}$ : Imaginary part of Morillon potential (since we study absorption)



• Modification in tail (relevant for stripping)



#### Effect on cross sections



• Significant flattening, consistent with transfer



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#### Conclusions and outlook

- An extension of CDCC to include core-valence absorption in elastic breakup reactions has been developed
- Application to  ${}^{12}\text{C}({}^{11}\text{Be}, n^{10}\text{Be}){}^{12}\text{C}$  and  ${}^{12}\text{C}({}^{41}\text{Ca}, n^{40}\text{Ca}){}^{12}\text{C}$  at 70 MeV/A shows larger effect when removing more deeply-bound species
- Modification of eikonal formalism for stripping to include absorption
- When correcting for proper loss of flux significant reduction in isospin dependence (consistent with transfer)
- Lots of work to do!
  - Uncertainty in VC optical potentials, more reliable (ab initio, dispersive, measurements?) are required
  - Extension to diffractive dispersion
  - Inclusion of real part of VC interaction (bound states?)
  - Go beyond eikonal (Ichimura-Austern-Vincent?)
  - Complete Gade plot
  - Momentum distributions



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