

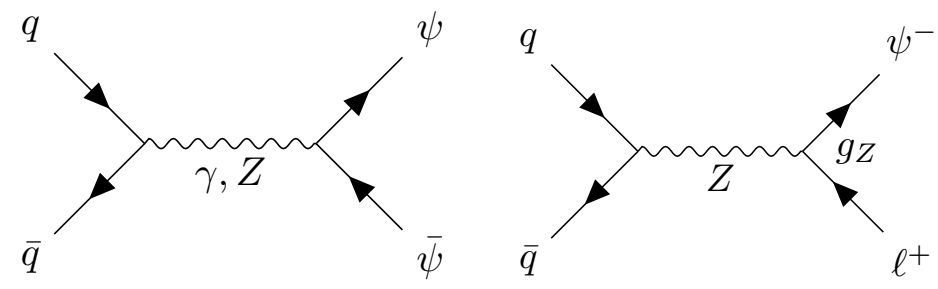
SM backgrounds for VLLs

Jonas M. Lindert



UK Research
and Innovation

Asymptotic Safety meets Particle Physics & Friends
17th December 2021



$$\psi_i \rightarrow S_{ij}^* l_j^- \rightarrow l_i^- l_j^+ l_j^-$$

$$\psi_i^0 \rightarrow S_{ji} \nu_j \rightarrow l_i^- l_j^+ \nu_j$$

....

Signatures for VLL's

	Processes	Signature	SM background
Singlet-model	$pp \rightarrow \psi_i \bar{\psi}_i \rightarrow l_i^- l_i^+ l_j^+ l_j^- l_k^+ l_k^-$ for $i, j, k = 1, 2, 3$	6L : 3 OS-SF leptons pairs	ZZZ Triboson
	$pp \rightarrow \psi_i \bar{\psi}_i \rightarrow l_i^- l_i^+ q_j \bar{q}_j l_k^+ l_k^-$ for $i, k = 1, 2$		
	$pp \rightarrow \psi_i \bar{\psi}_i \rightarrow l_i^- l_i^+ l_j^+ l_j^- \nu_k \bar{\nu}_k$ for $i, j = 1, 2$		
	$pp \rightarrow \psi_i \bar{\psi}_i \rightarrow \nu_i l_i^+ l_j^+ l_j^- l_k^- \bar{\nu}_k$ for $i, j, k = 1, 2$		
	$pp \rightarrow \psi_i \bar{\psi}_i \rightarrow l_i^- \bar{\nu}_i l_j^+ l_j^- l_k^+ \nu_k$ for $i, j, k = 1, 2$		
	$pp \rightarrow \psi_i l_i^+ \rightarrow l_i^- l_j^+ l_j^- l_i^+$ for $i, j = 1, 2$		
	$pp \rightarrow \bar{\psi}_i l_i^- \rightarrow l_i^+ l_j^+ l_j^- l_i^-$ for $i, j = 1, 2$		
Doublet-model	$pp \rightarrow \psi_i^0 \bar{\psi}_i^0 \rightarrow \nu_j \bar{\nu}_k l_j^+ l_i^- l_i^+ l_k^-$ for $i, j, k = 1, 2$	4L + jets / 4L+MET	ZZ Diboson
	$pp \rightarrow \psi_i^- \bar{\psi}_i^0 \rightarrow l_i^- l_i^+ \bar{\nu}_j l_j^- l_k^+ l_k^-$ for $i, j, k = 1, 2$		
	$pp \rightarrow \psi_i^0 \psi_i^+ \rightarrow l_i^- l_i^+ l_j^+ \nu_j l_k^+ l_k^-$ for $i, j, k = 1, 2$		
	$pp \rightarrow \psi_i^- \bar{\psi}_i^0 \rightarrow l_i^- l_i^+ \bar{q}_j q_j l_k^+ l_k^-$ for $i, k = 1, 2$		
	$pp \rightarrow \psi_i^0 \psi_i^+ \rightarrow l_i^- l_i^+ \bar{q}_j q_j l_k^+ l_k^-$ for $i, k = 1, 2$		
		5L +MET	ZZW Triboson

6L/5L very clean. However, small rates!

-> Look at 4L/3L signatures

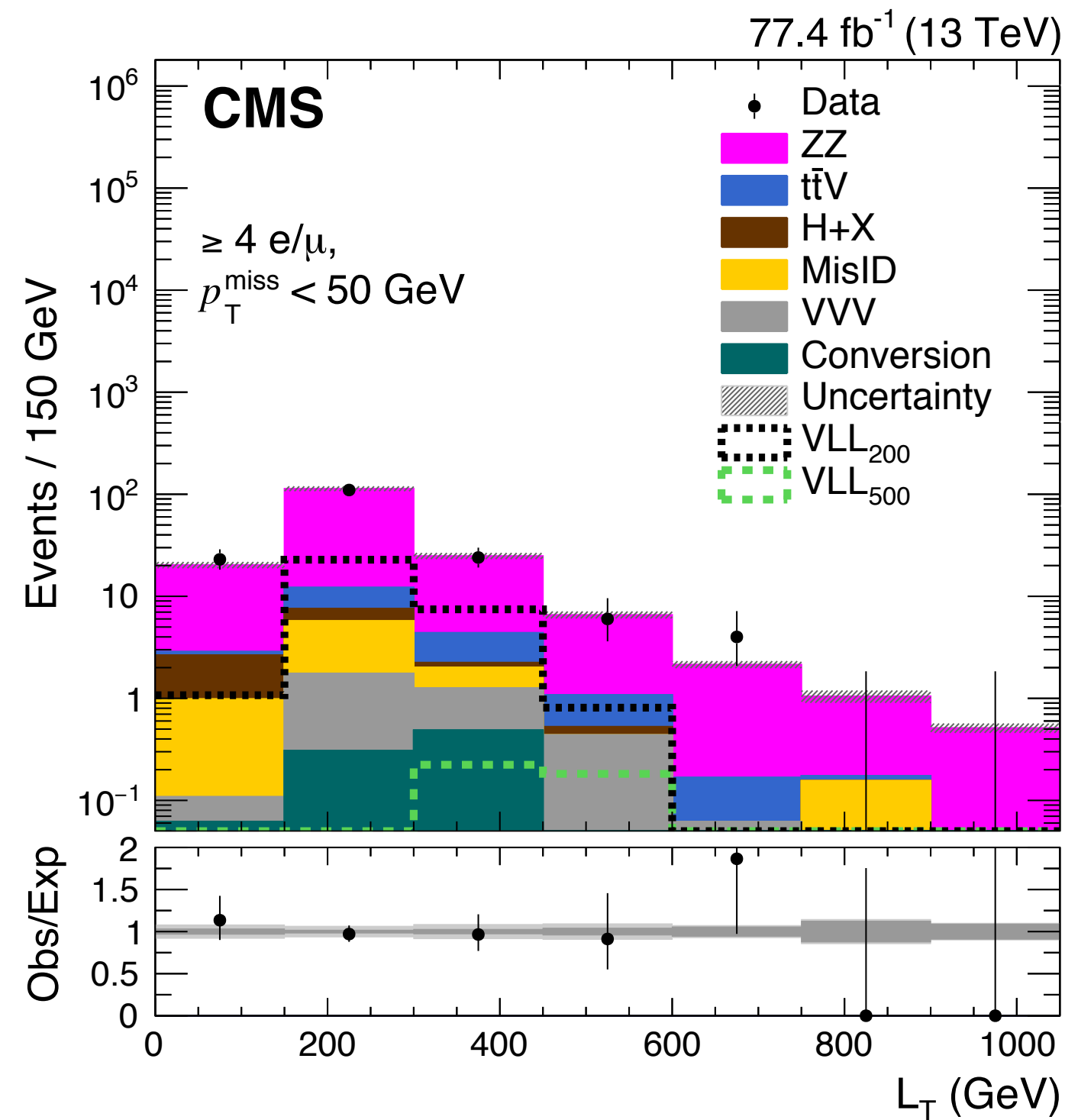
[Multi-lepton signatures of vector-like leptons with flavor, Bißmann, Hiller, Hormigos-Feliu, Litim, 20]

Available searches: CMS-EXO-18-005

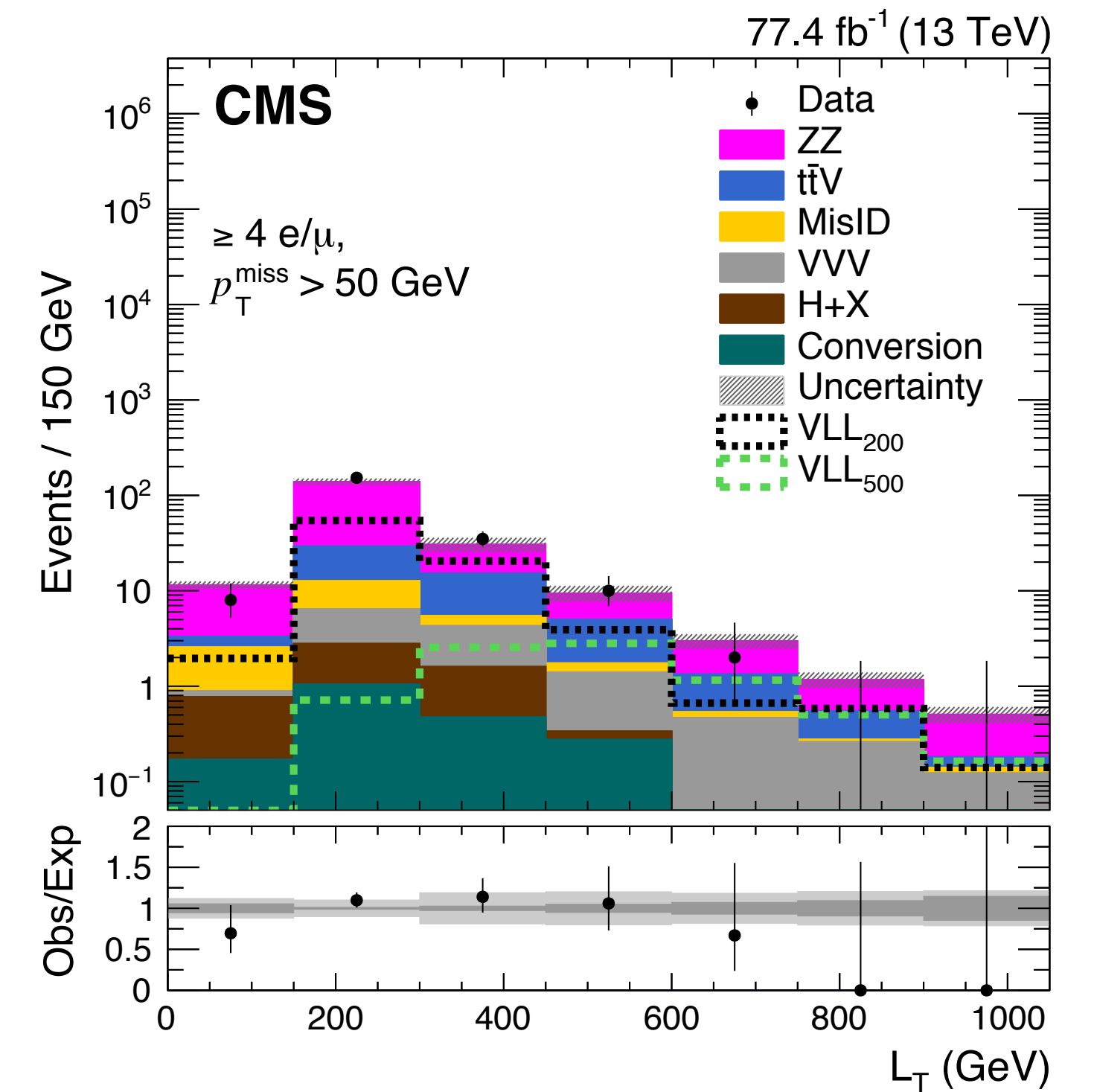
N_{leptons}	p_T^{miss} (GeV)
$\geq 4e/\mu$	< 50 > 50
$3e/\mu$	< 150 > 150
$2e/\mu$ OS (or SS) + $\geq 1\tau_h$	< 150 > 150

Dominant background

$$L_T = \sum_{i \in \ell} p_{T,i}$$



ZZ

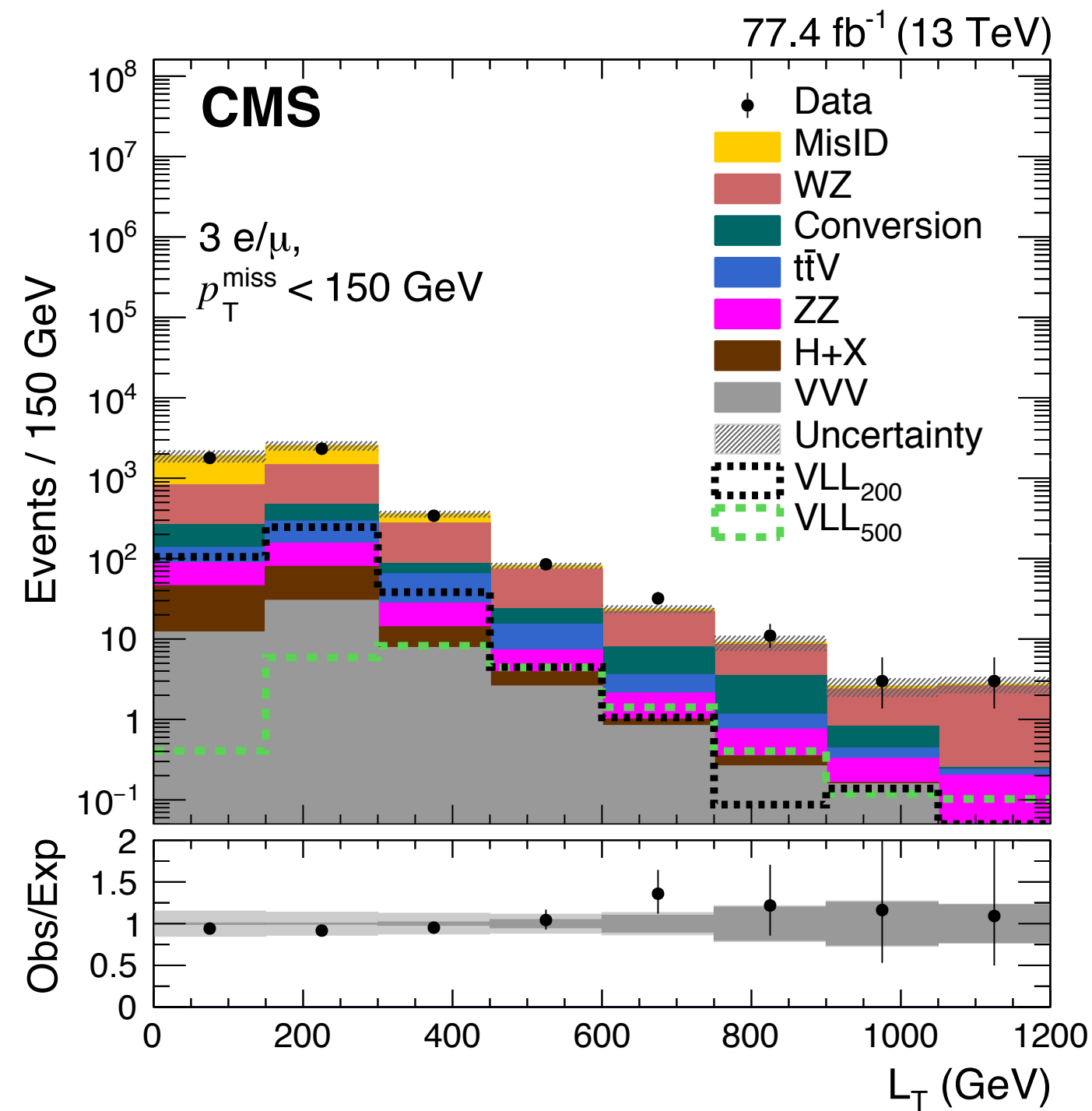


ZZ / (ttV)

also: misID -> Z+jets / ttbar+jets

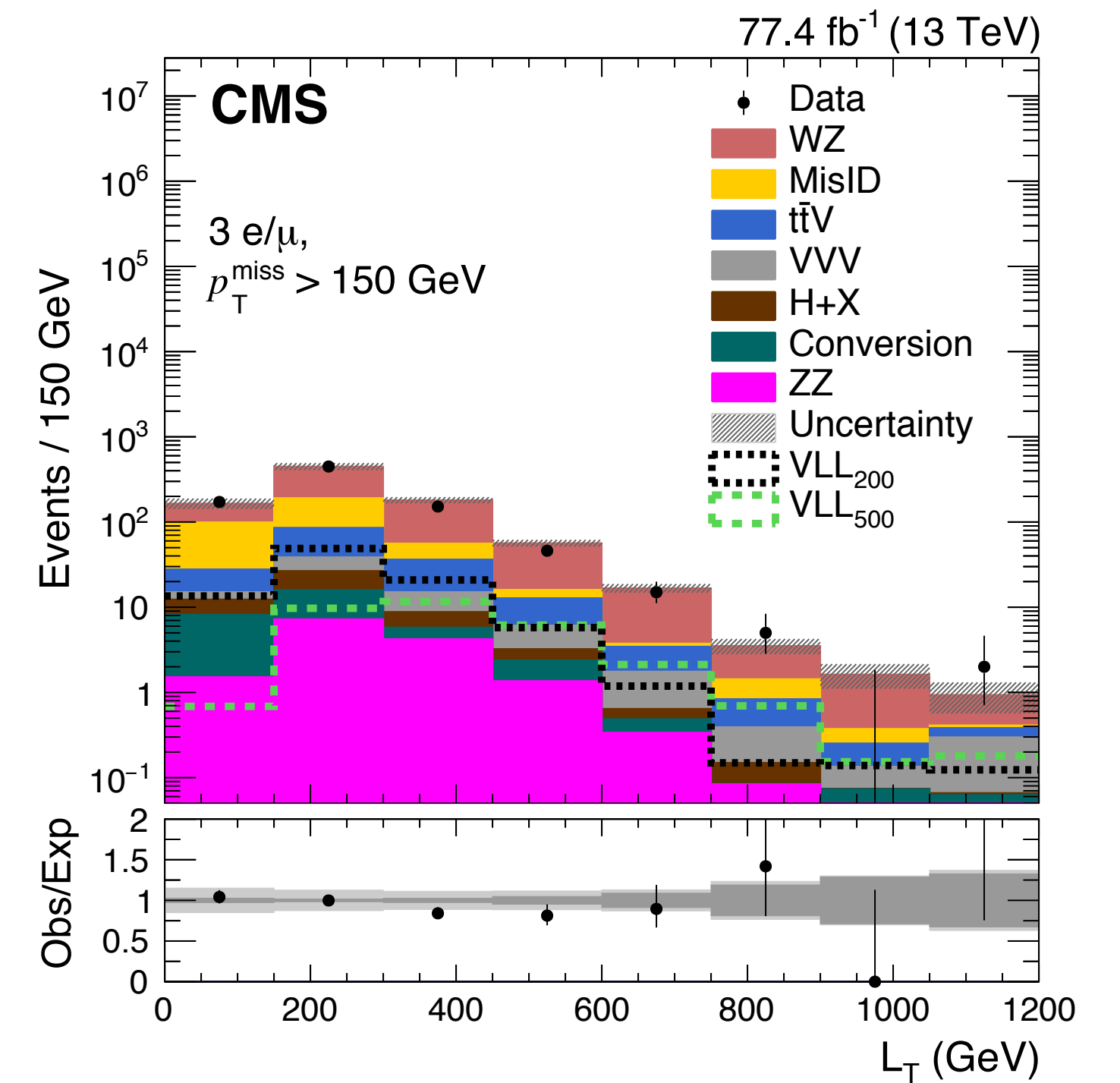
Available searches: CMS-EXO-18-005

N_{leptons}	p_T^{miss} (GeV)
$\geq 4e/\mu$	< 50 > 50
$3e/\mu$	< 150 > 150
$2e/\mu$ OS (or SS) + $\geq 1\tau_h$	< 150 > 150



Dominant background

WZ



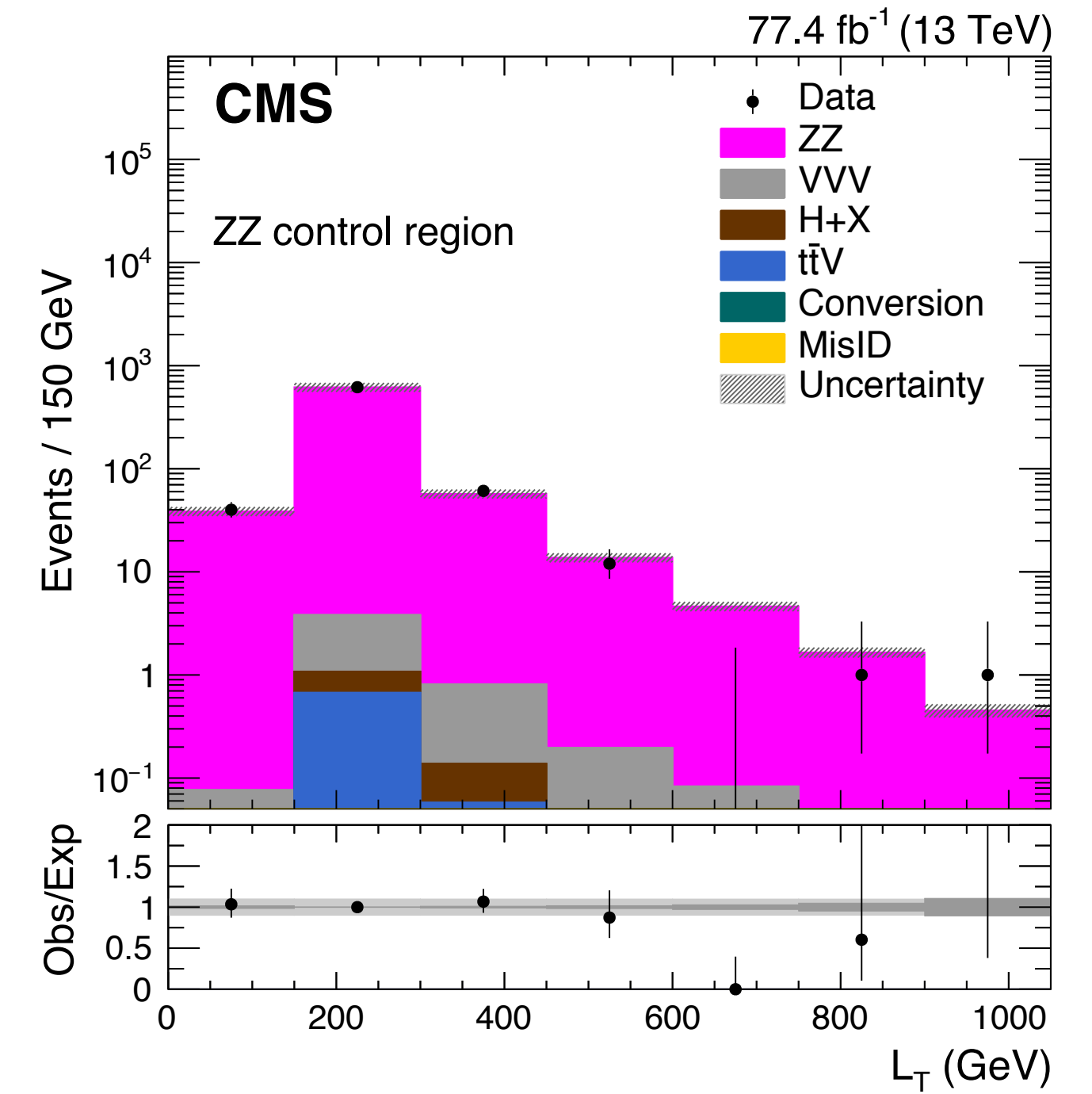
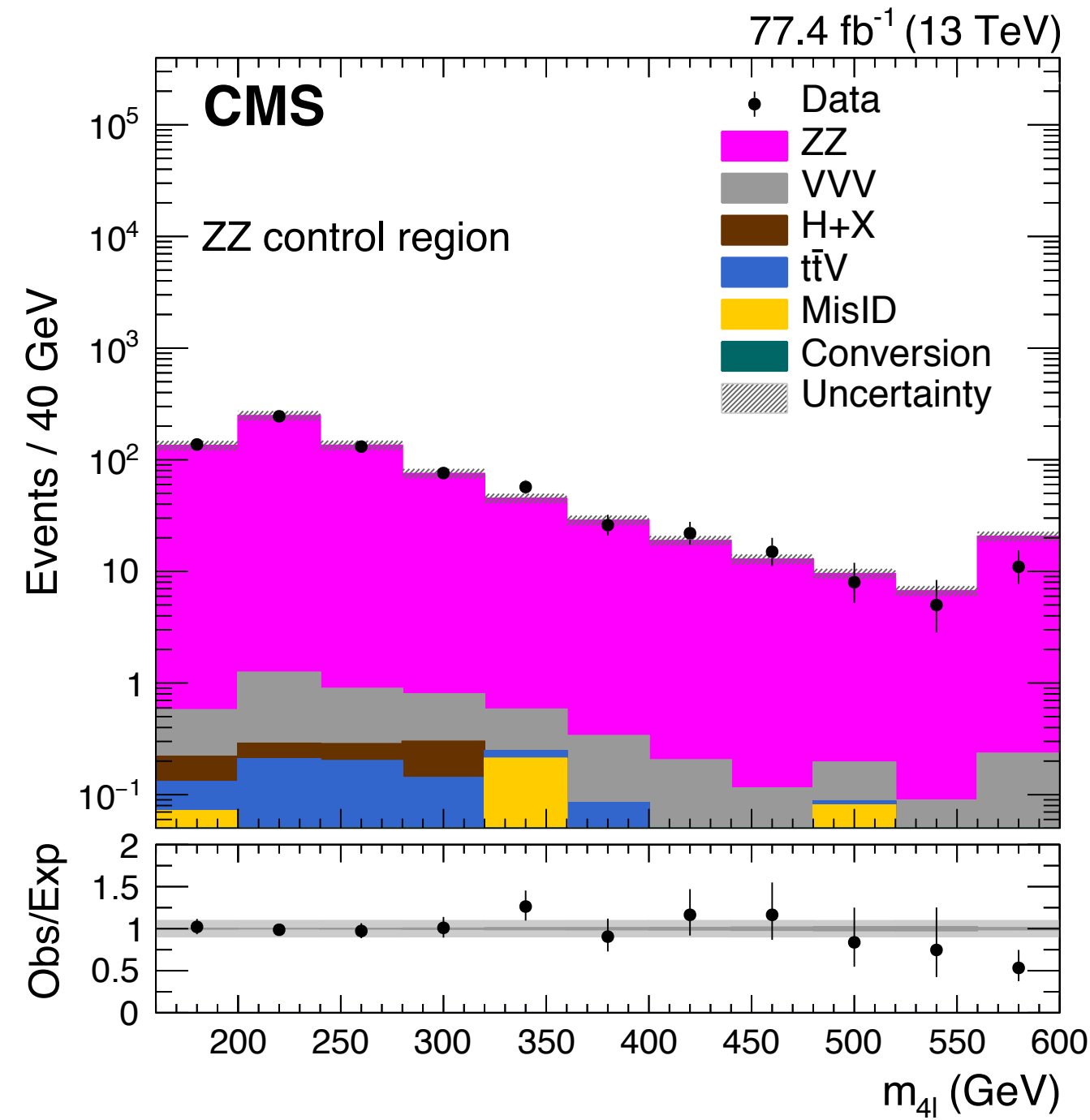
WZ

Background determination: VV@NLO+PS & overall normalization from CR

Available searches: CMS-EXO-18-005

N_{leptons}	p_T^{miss} (GeV)
$\geq 4e/\mu$	< 50 > 50
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$2e/\mu$ OS (or SS) + $\geq 1\tau_h$	< 150 > 150

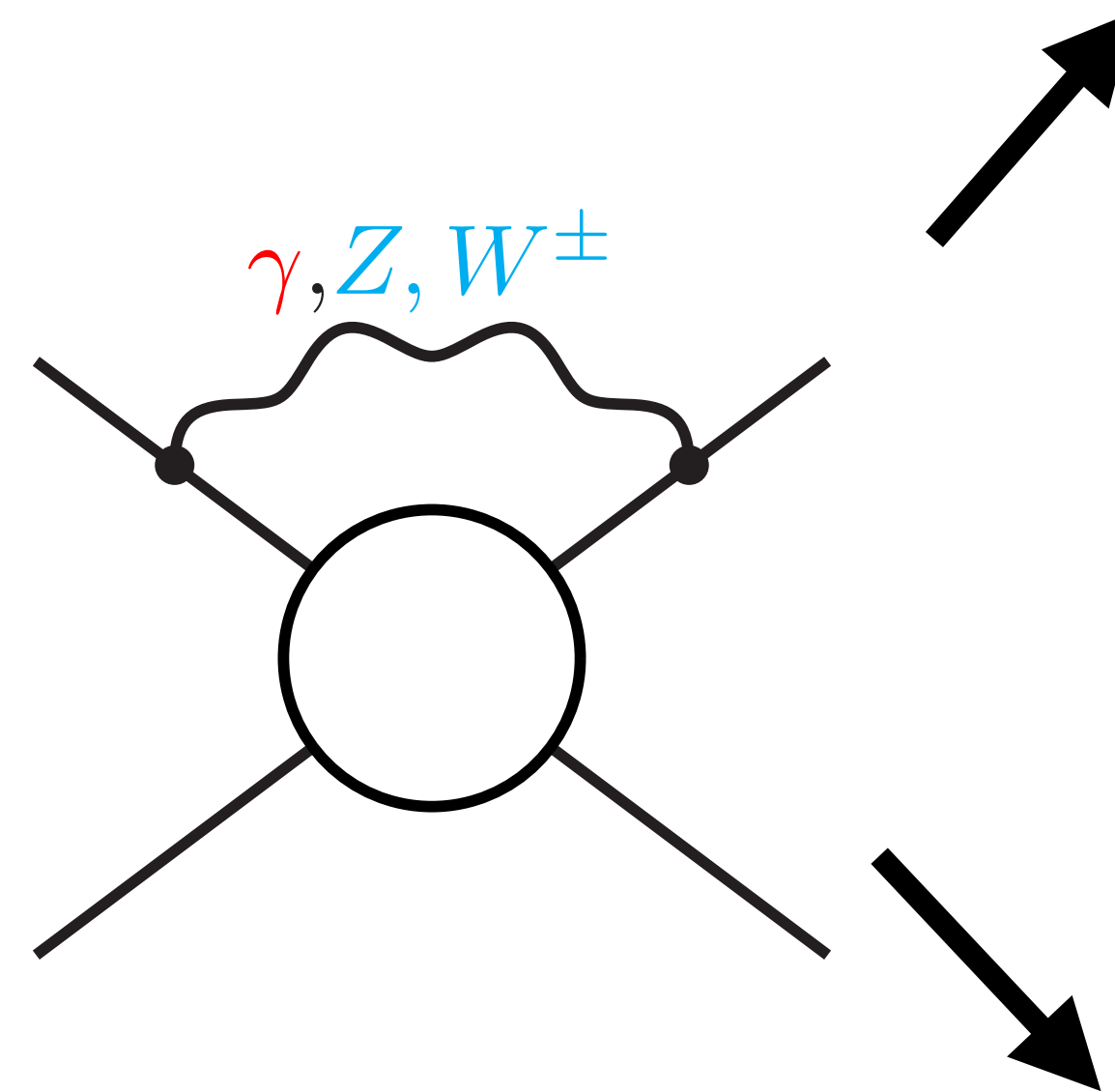
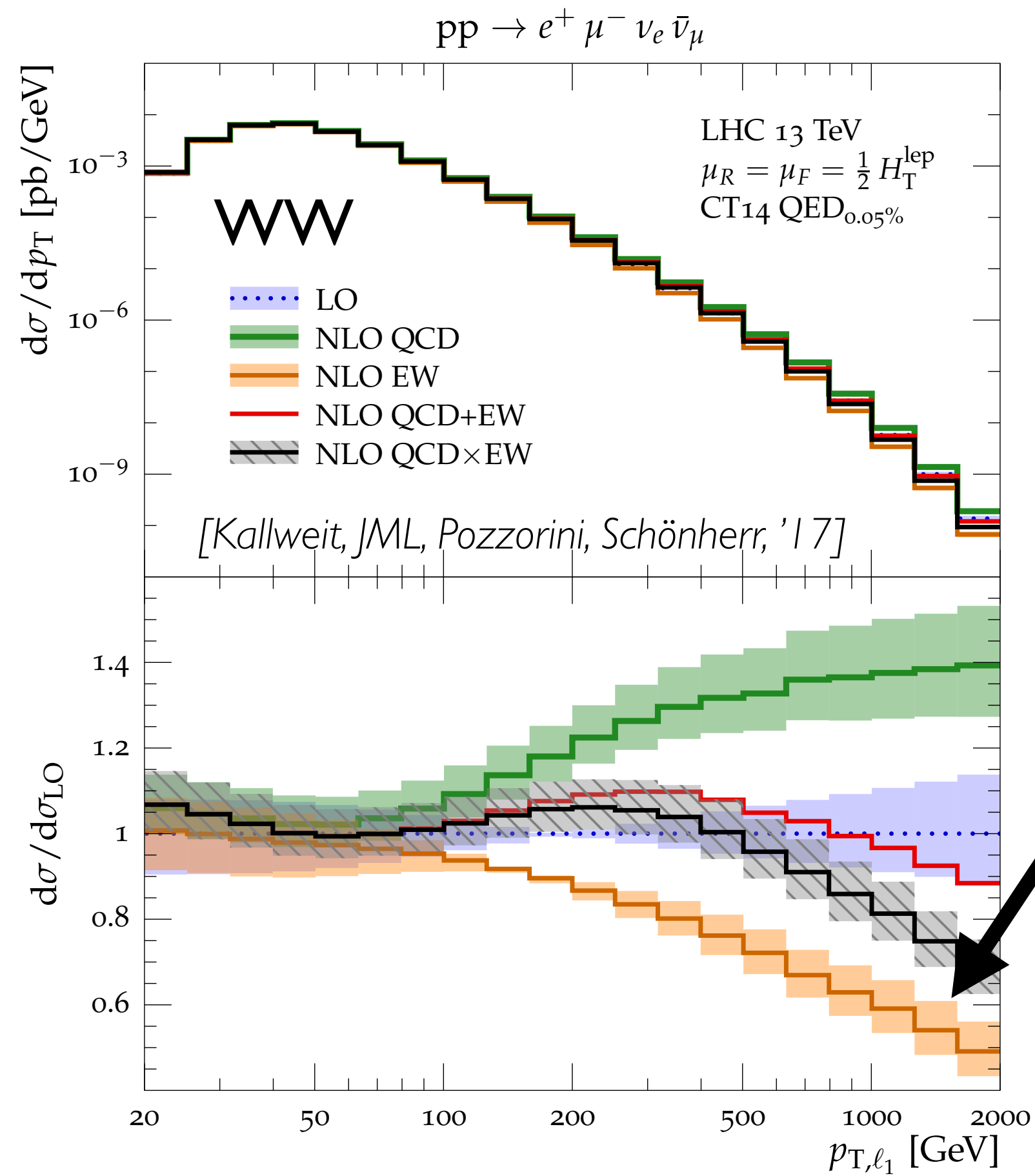
Control region



This will not be good enough with HL-LHC data samples!

Relevance of EW higher-order corrections: Sudakov logs in the tails

I. Possible large (negative) enhancement due to soft/collinear **logs** from virtual EW gauge bosons:



[Ciafaloni, Comelli, '98;
Lipatov, Fadin, Martin, Melles, '99;
Kuehen, Penin, Smirnov, '99;
Denner, Pozzorini, '00]

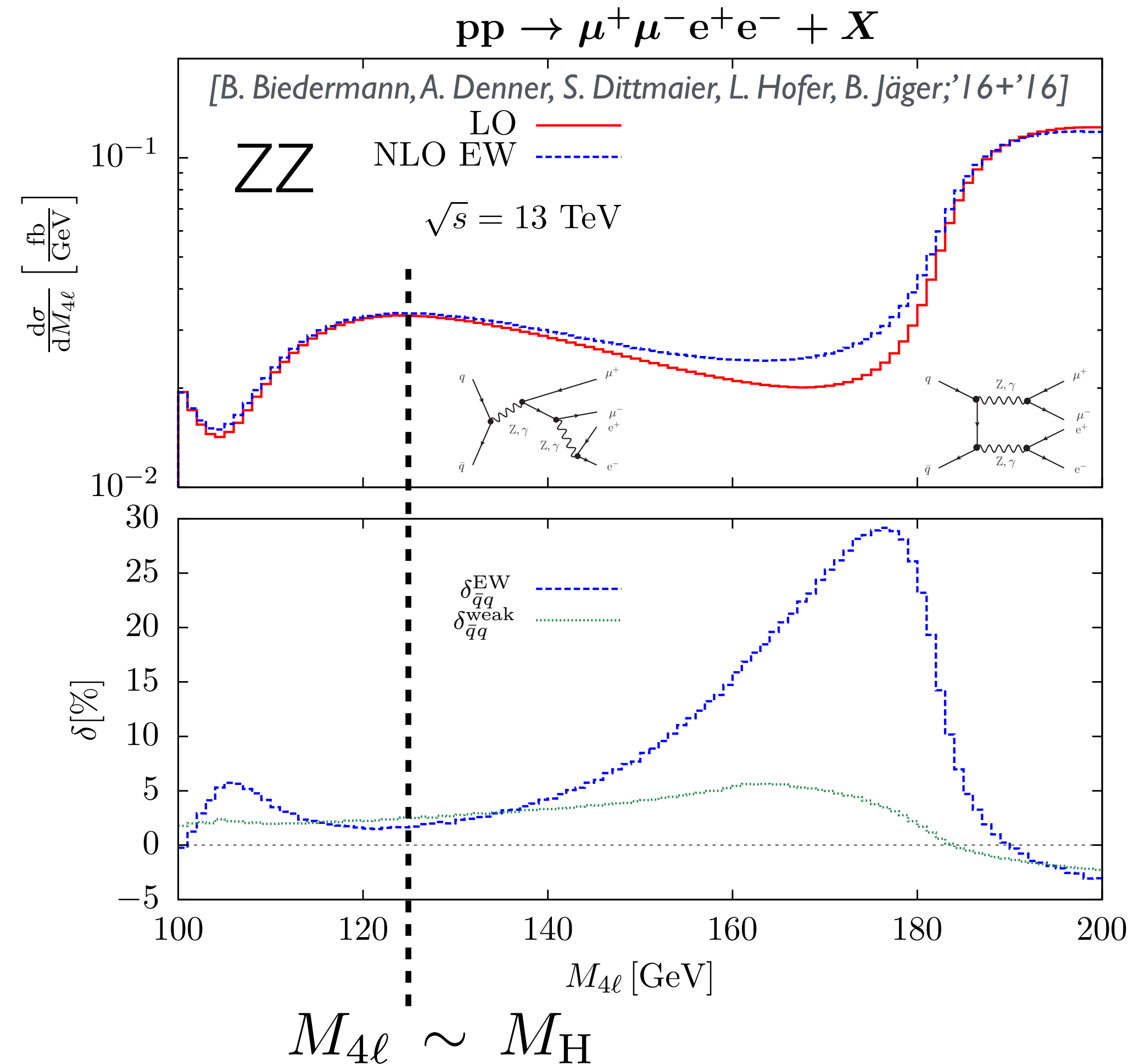
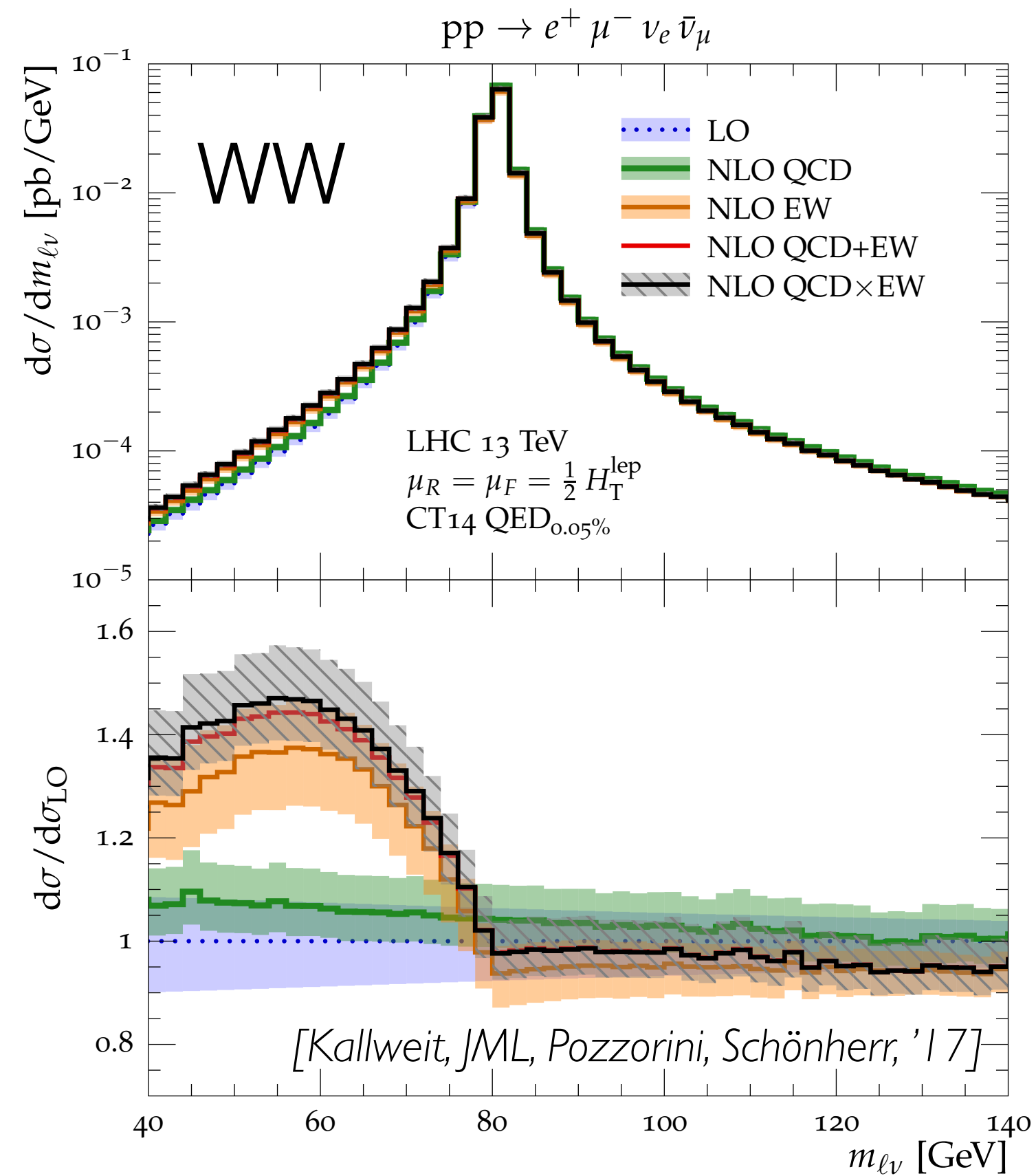
Universality and factorisation: [Denner, Pozzorini, '01]

$$\delta\mathcal{M}_{\text{LL+NLL}}^{1\text{-loop}} = \frac{\alpha}{4\pi} \sum_{k=1}^n \left\{ \frac{1}{2} \sum_{l \neq k} \sum_{a=\gamma, Z, W^\pm} I^a(k) I^{\bar{a}}(l) \ln^2 \frac{\hat{s}_{kl}}{M^2} + \gamma^{\text{ew}}(k) \ln \frac{\hat{s}}{M^2} \right\} \mathcal{M}_0$$

→ overall large (negative) effect in the tails of distributions:
p_T, m_{inv}, H_T, ... (relevant for BSM searches!)

Relevance of EW higher-order corrections: collinear QED radiation

- II. Possible large enhancement due to soft/collinear **logs** from photon radiation $\sim \alpha \log\left(\frac{m_f^2}{Q^2}\right)$ in sufficiently exclusive observables.



→ important for radiative tails, Higgs backgrounds etc.

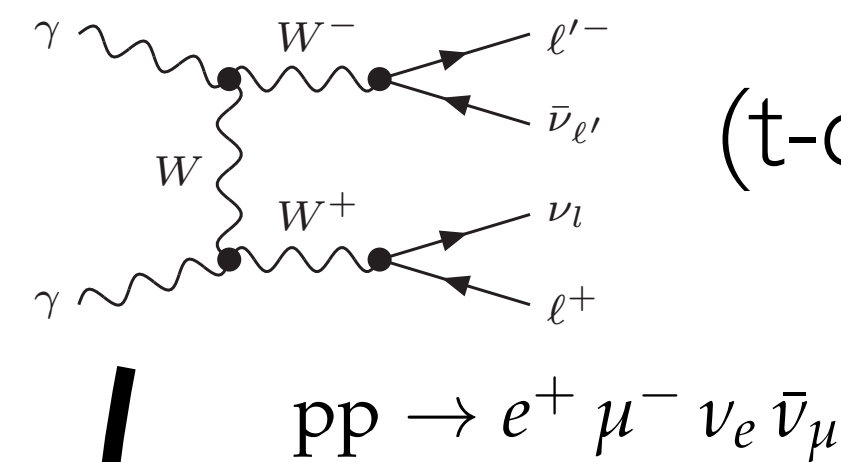
→ typically considered via QED PS (PHOTOS / YFS)

Relevance of EW higher-order corrections: photon-induced channels

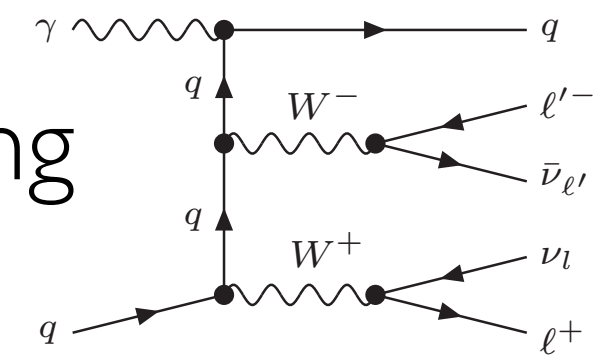
III. QED factorisation and thus photon luminosities needed to absorb IS photon singularities.

→ Possible large enhancement due to photon-induced channels in the tails of kinematic distributions,

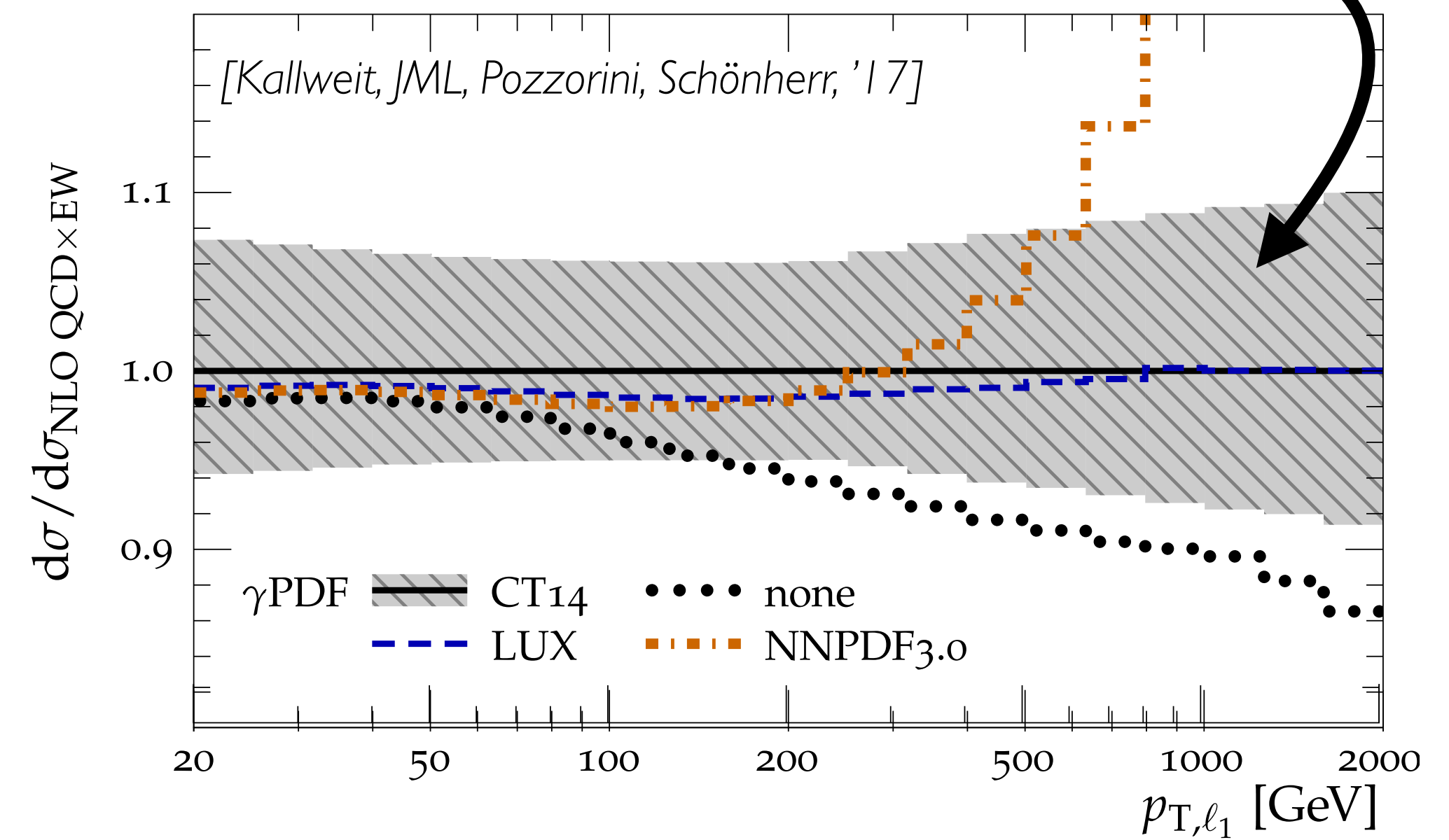
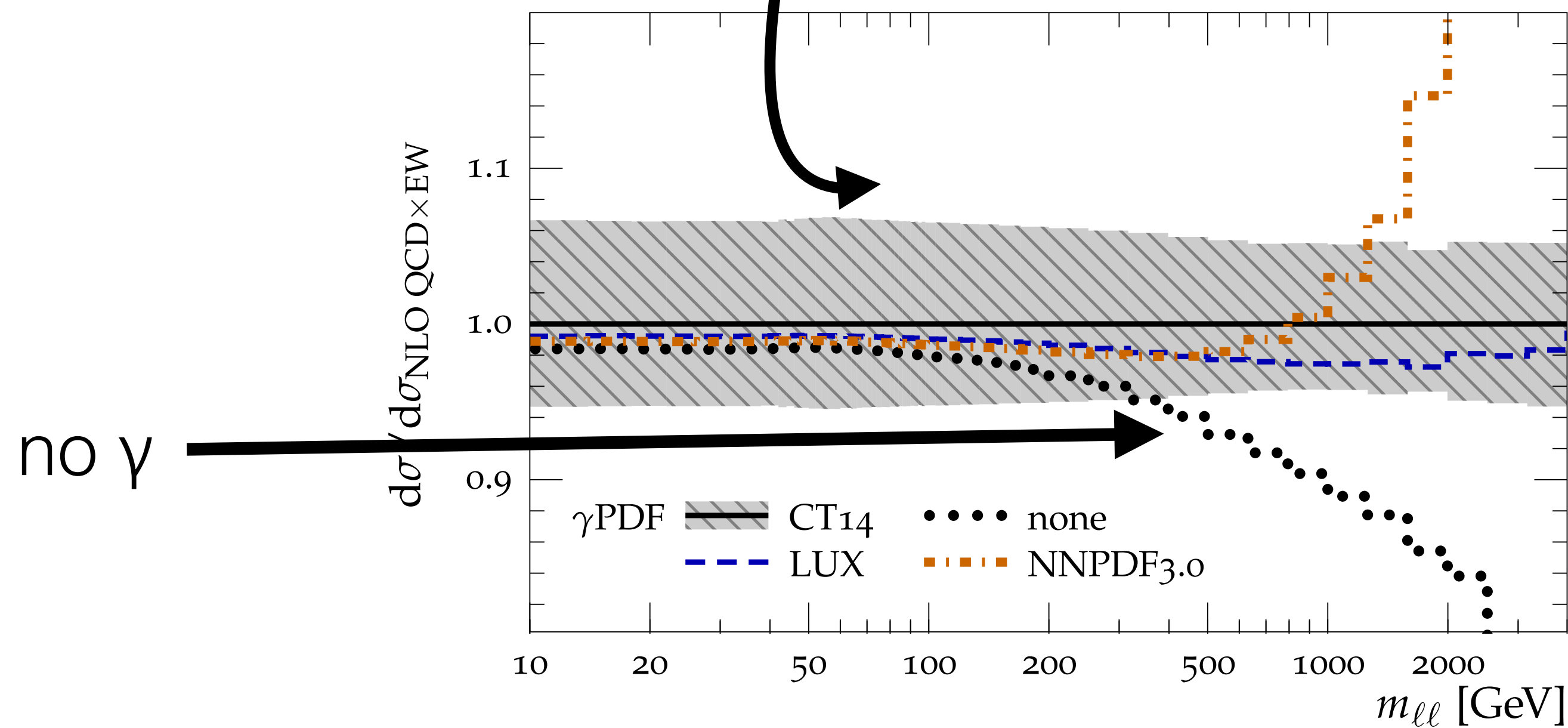
in particular in WW:



(t-channel enhancement), but also in Bremsstrahlung



$$d\sigma_{\text{LO}} = d\sigma_{\text{LO}}^{q\bar{q}} + d\sigma_{\text{LO}}^{\gamma\gamma}$$



→ large differences between different photon descriptions. Now settled: LUXqed superior

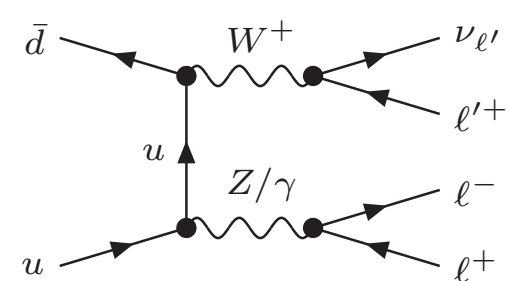
→ O(10%) contributions from photon-induced channels

Perturbative expansion for VV

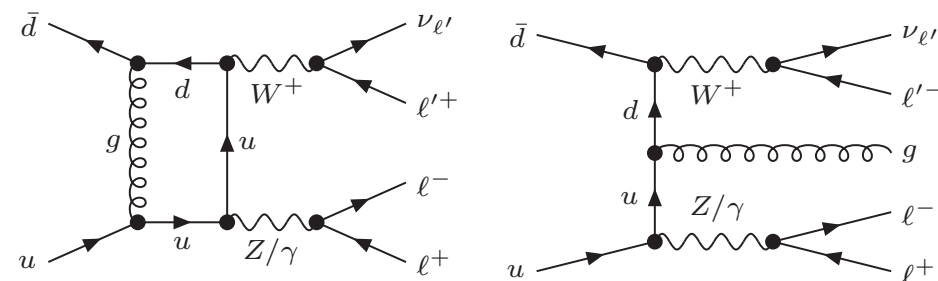
$$\begin{aligned}
 d\sigma = & d\sigma_{\text{LO}} + \alpha_S d\sigma_{\text{NLO QCD}} + \alpha_{\text{EW}} d\sigma_{\text{NLO EW}} \\
 & + \alpha_S^2 d\sigma_{\text{NNLO QCD}} + \alpha_{\text{EW}}^2 d\sigma_{\text{NNLO EW}} + \alpha_S \alpha_{\text{EW}} d\sigma_{\text{NNLO QCD} \times \text{EW}} \\
 & + \alpha_S^3 d\sigma_{\text{N3LO QCD}} + \dots
 \end{aligned}$$

?
?
?

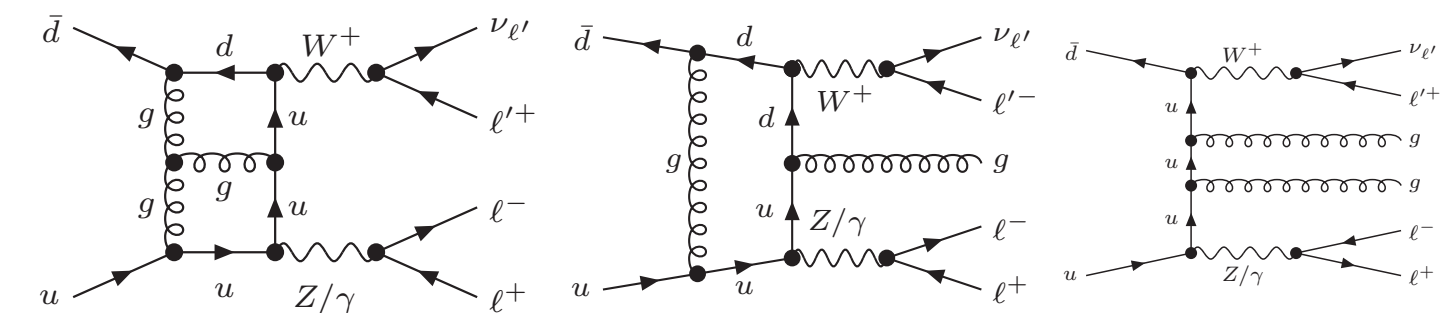
LO



NLO QCD



NNLO QCD



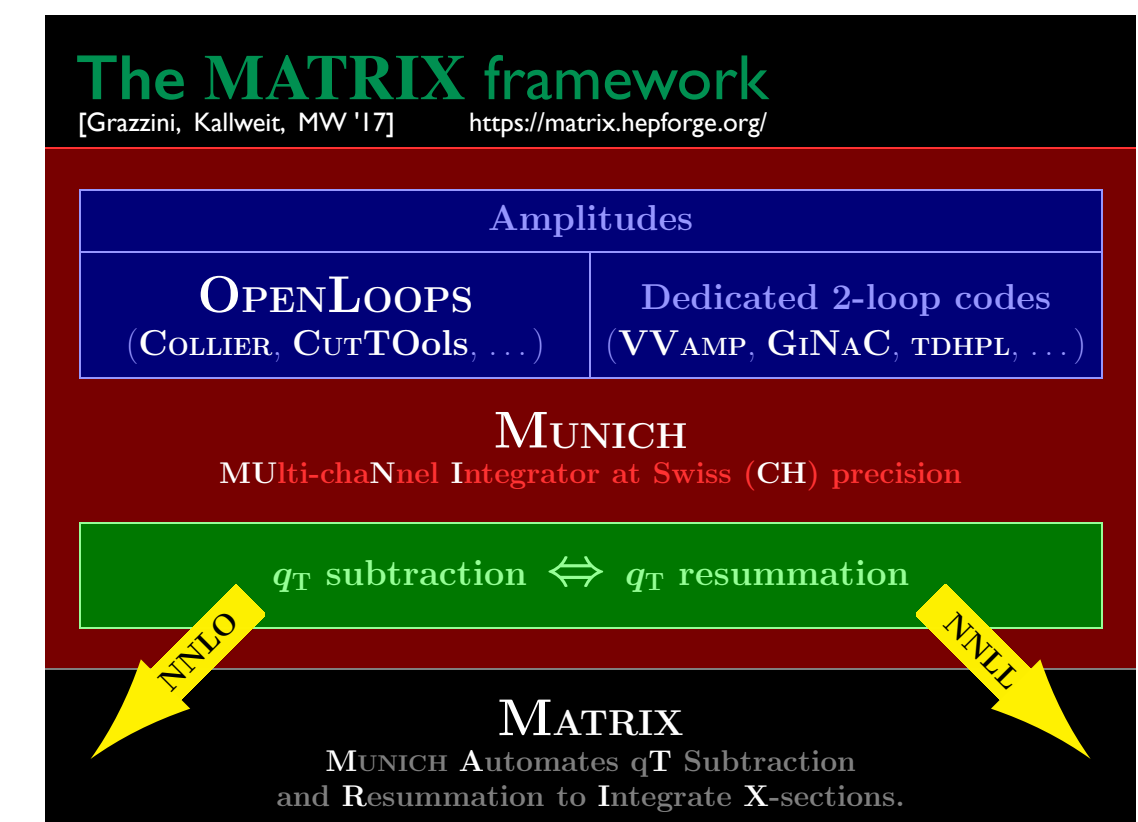
Perturbative expansion for VV

$$\begin{aligned}
 d\sigma = & \underbrace{d\sigma_{\text{LO}}}_{\text{LO}} + \underbrace{\alpha_S d\sigma_{\text{NLO}}}_{\text{NLO QCD}} + \underbrace{\alpha_{\text{EW}} d\sigma_{\text{NLO EW}}}_{\text{NLO EW}} \\
 & + \underbrace{\alpha_S^2 d\sigma_{\text{NNLO}}}_{\text{NNLO QCD}} + \underbrace{\alpha_{\text{EW}}^2 d\sigma_{\text{NNLO EW}}}_{\text{NNLO EW}} + \underbrace{\alpha_S \alpha_{\text{EW}} d\sigma_{\text{NNLO QCD} \times \text{EW}}}_{\text{NNLO QCD-EW}} \\
 & + \underbrace{\alpha_S^3 d\sigma_{\text{NNLO}}}_{\text{N3LO QCD}} + \dots
 \end{aligned}$$

?
?
?

NNLO QCD

In MATRIX [Grazzini, Kallweit, Wiesemann '17] all on-shell & off-shell diboson processes $pp \rightarrow VV$ are available



Perturbative expansion for VV

$$\begin{aligned}
 d\sigma = & \underbrace{d\sigma_{\text{LO}}}_{\text{LO}} + \underbrace{\alpha_S d\sigma_{\text{NLO}}}_{\text{NLO QCD}} + \underbrace{\alpha_{\text{EW}} d\sigma_{\text{NLO EW}}}_{\text{NLO EW}} \\
 & + \underbrace{\alpha_S^2 d\sigma_{\text{NNLO}}}_{\text{NNLO QCD}} + \underbrace{\alpha_{\text{EW}}^2 d\sigma_{\text{NNLO EW}}}_{\text{NNLO EW}} + \underbrace{\alpha_S \alpha_{\text{EW}} d\sigma_{\text{NNLO QCD} \times \text{EW}}}_{\text{NNLO QCD-EW}} \\
 & + \underbrace{\alpha_S^3 d\sigma_{\text{NNLO}}}_{\text{N3LO QCD}} + \dots
 \end{aligned}$$

?
?

NLO EW

- 4I-DF-ZZ Biedermann, Denner, Dittmaier, Hofer, Jäger; '16, '16
- 2I-DF-WWW Biedermann, Billoni, Denner, Dittmaier, Hofer, Jäger, Salfelder; '16
- 2I-SF-ZZ & 2I-SF-ZZWW & 2I-DF-WWW Kallweit, JML, Pozzorini, Schönherr; '17
- 3I-DF-WZ & 3I-DF-WZ Biedermann, Denner, Hofer; '17

Perturbative expansion for VV

$$\begin{aligned}
 d\sigma = & \underbrace{d\sigma_{\text{LO}}}_{\text{LO}} + \underbrace{\alpha_S d\sigma_{\text{NLO}}}_{\text{NLO QCD}} + \underbrace{\alpha_{\text{EW}} d\sigma_{\text{NLO EW}}}_{\text{NLO EW}} \\
 & + \underbrace{\alpha_S^2 d\sigma_{\text{NNLO}}}_{\text{NNLO QCD}} + \underbrace{\alpha_{\text{EW}}^2 d\sigma_{\text{NNLO EW}}}_{\text{NNLO EW}} + \underbrace{\alpha_S \alpha_{\text{EW}} d\sigma_{\text{NNLO QCD} \times \text{EW}}}_{\text{NNLO QCD-EW}} \\
 & + \underbrace{\alpha_S^3 d\sigma_{\text{NNLO}}}_{\text{N3LO QCD}} + \dots
 \end{aligned}$$

?
?
?

NNLO QCD + NLO EW

4l-SF-ZZ	$pp \rightarrow l^+ l^- l^+ l^-$	ZZ
4l-DF-ZZ	$pp \rightarrow l^+ l^- l'^+ l'^-$	ZZ
3l-SF-WZ	$pp \rightarrow l^+ l^- l \nu_\ell$	WZ
3l-DF-WZ	$pp \rightarrow l^+ l^- l' \nu_{\ell'}$	WZ
2l-SF-ZZ	$pp \rightarrow l^+ l^- \nu_\ell \bar{\nu}_{\ell'}$	ZZ
2l-SF-ZZWW	$pp \rightarrow l^+ l^- \nu_\ell \bar{\nu}_\ell$	ZZ, WW
2l-DF-WW	$pp \rightarrow l^+ l'^- \nu_\ell \bar{\nu}_{\ell'}$	WW

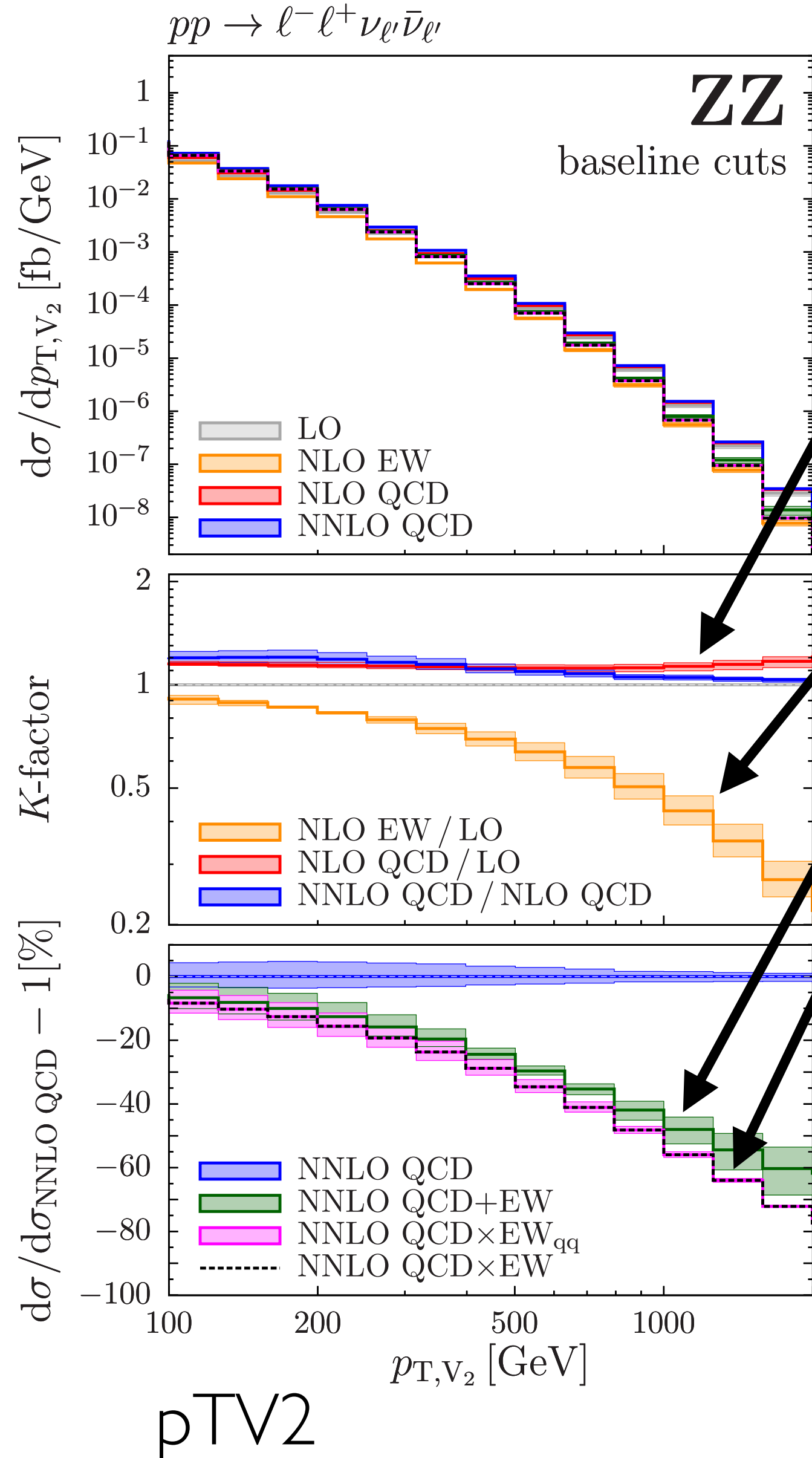
In Matrix+OpenLoops all (massive) diboson processes are now available at **NNLO QCD + NLO EW**

[M. Grazzini, S. Kallweit, JML, S. Pozzorini, M. Wiesemann; 1912.00068]

(code available upon request)

NNLO QCD + NLO EW for dibosons: pTV2

[M. Grazzini, S. Kallweit, JML, S. Pozzorini, M. Wiesemann; 1912.00068]



- moderate QCD corrections

- ▶ NNLO/NLO QCD very small at large pTV2

- ▶ NNLO QCD uncertainty: few percent

- NLO EW/LO = -(50-60)% @ 1 TeV

$$d\sigma_{\text{NNLO QCD+EW}} = d\sigma_{\text{LO}} (1 + \delta_{\text{QCD}} + \delta_{\text{EW}}) + d\sigma_{\text{LO}}^{gg}$$

$$\begin{aligned} d\sigma_{\text{NNLO QCD}\times\text{EW}} &= d\sigma_{\text{LO}} (1 + \delta_{\text{QCD}}) (1 + \delta_{\text{EW}}) + d\sigma_{\text{LO}}^{gg} \\ &= d\sigma_{\text{NNLO QCD+EW}} + d\sigma_{\text{LO}} \delta_{\text{QCD}} \delta_{\text{EW}} \end{aligned}$$

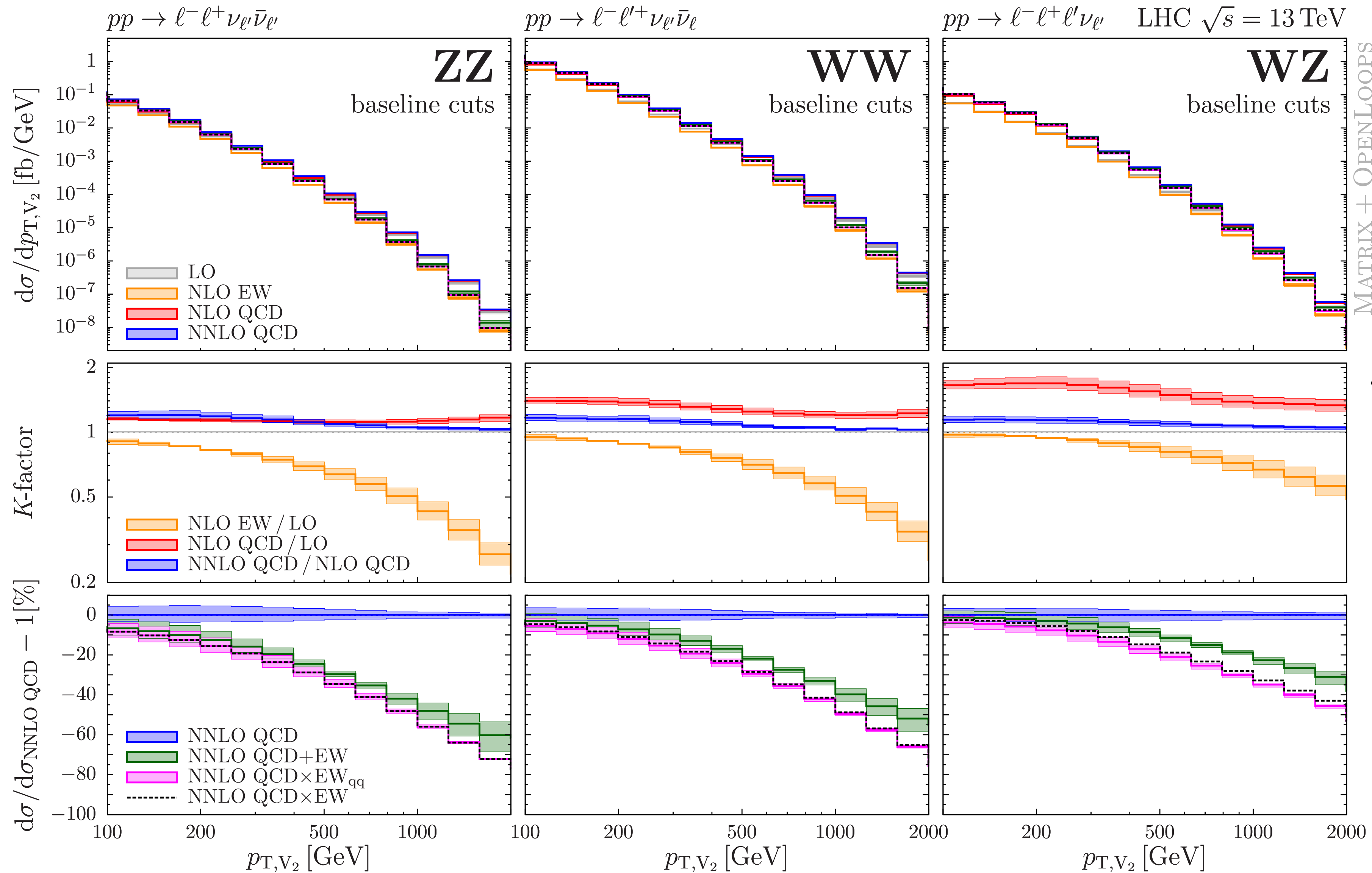
- difference very conservative upper bound on $\mathcal{O}(\alpha_S \alpha)$

- multiplicative/factorised combination clearly superior (EW Sudakov logs x soft QCD)

- dominant uncertainty at large pTV2: $\mathcal{O}(\alpha^2) \sim \alpha_w^2 \log^4(Q^2/M_W^2)$

Estimate: $\frac{1}{2} \delta_{\text{EW}}^2$

NNLO QCD + NLO EW for dibosons: pTV2



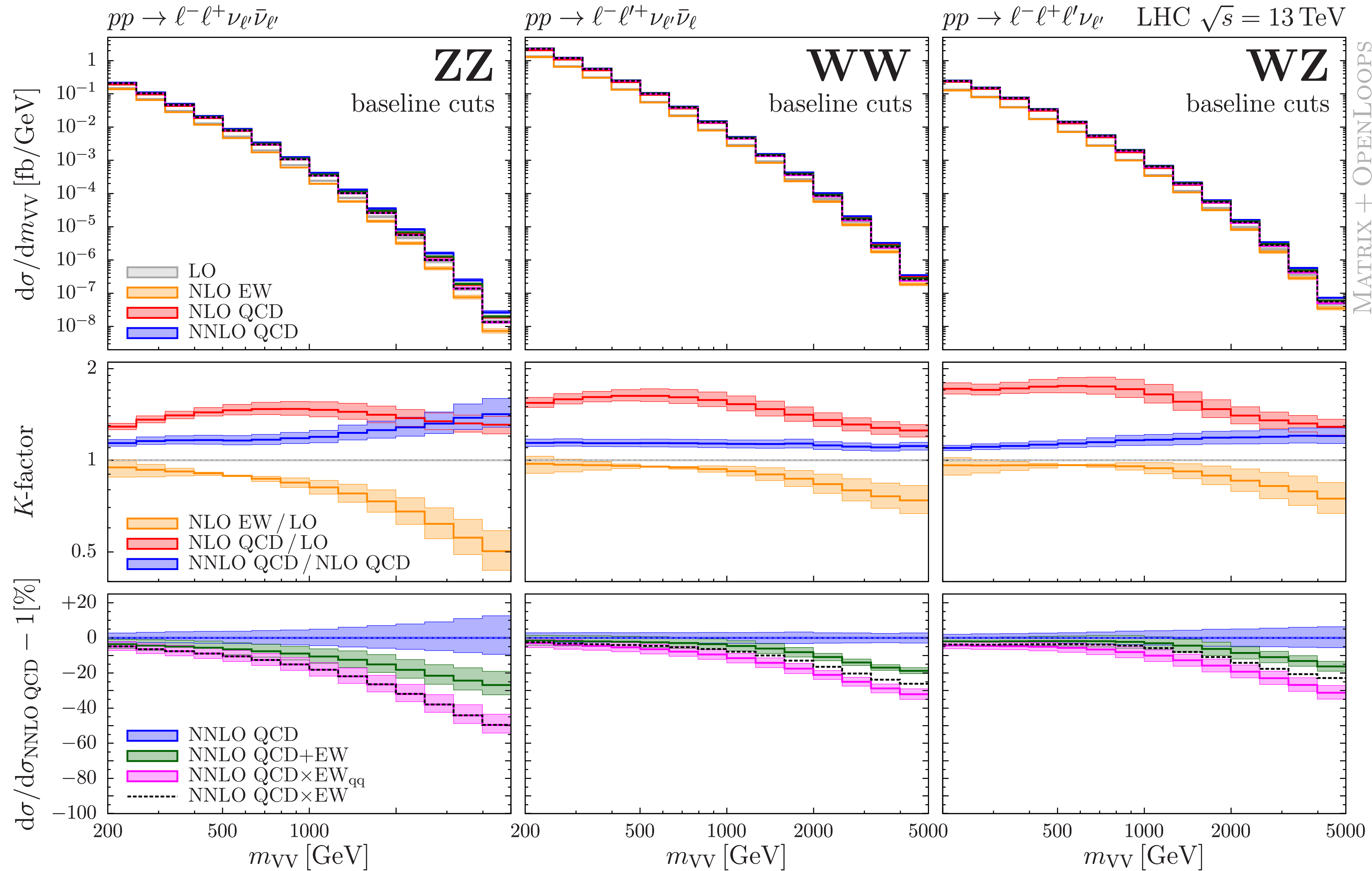
- consistent picture amongst all processes

- Largest QCD corrections in WZ (radiation zero at LO)

- Largest EW corrections in ZZ

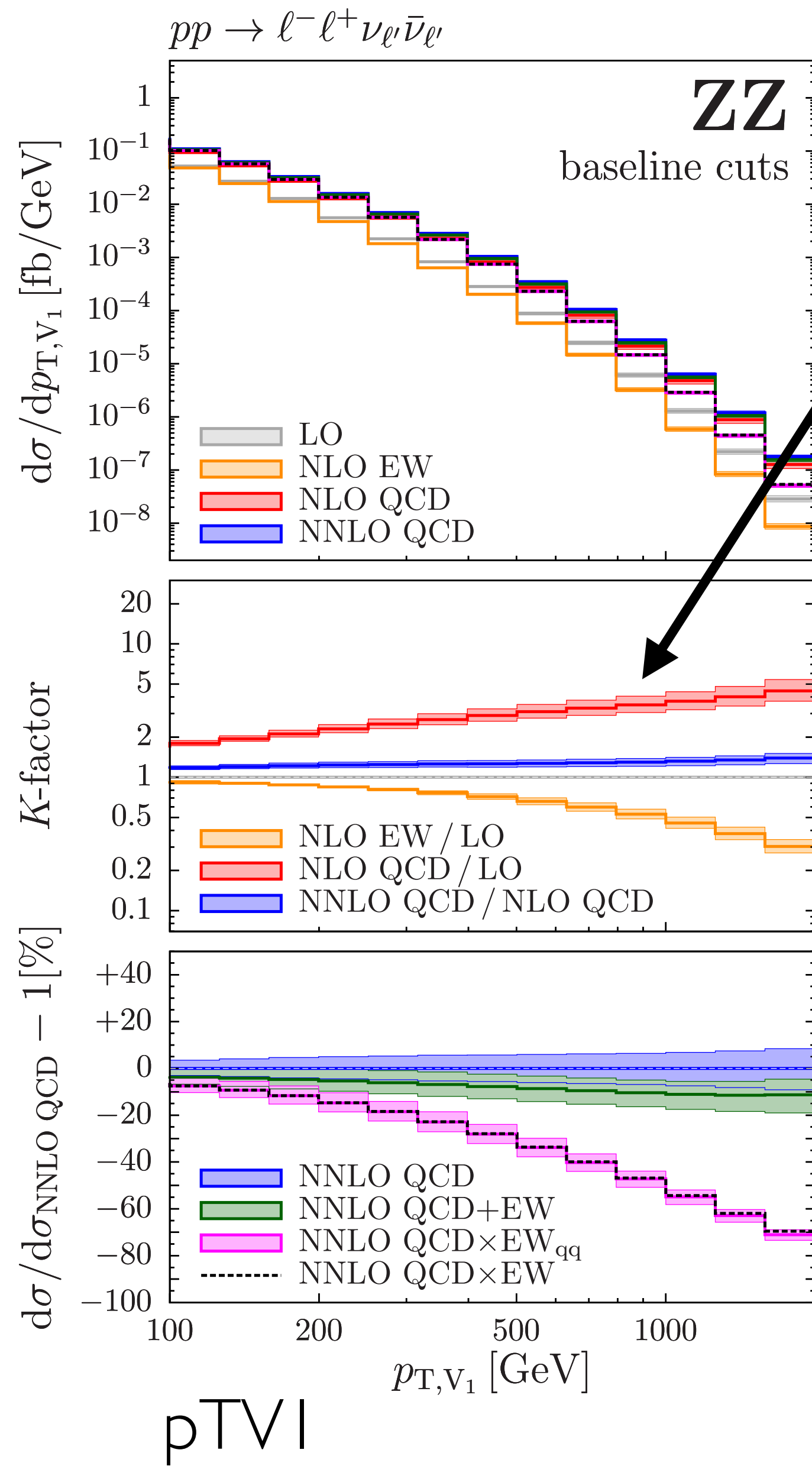
pTV2

NNLO QCD + NLO EW for dibosons: m_{VV}



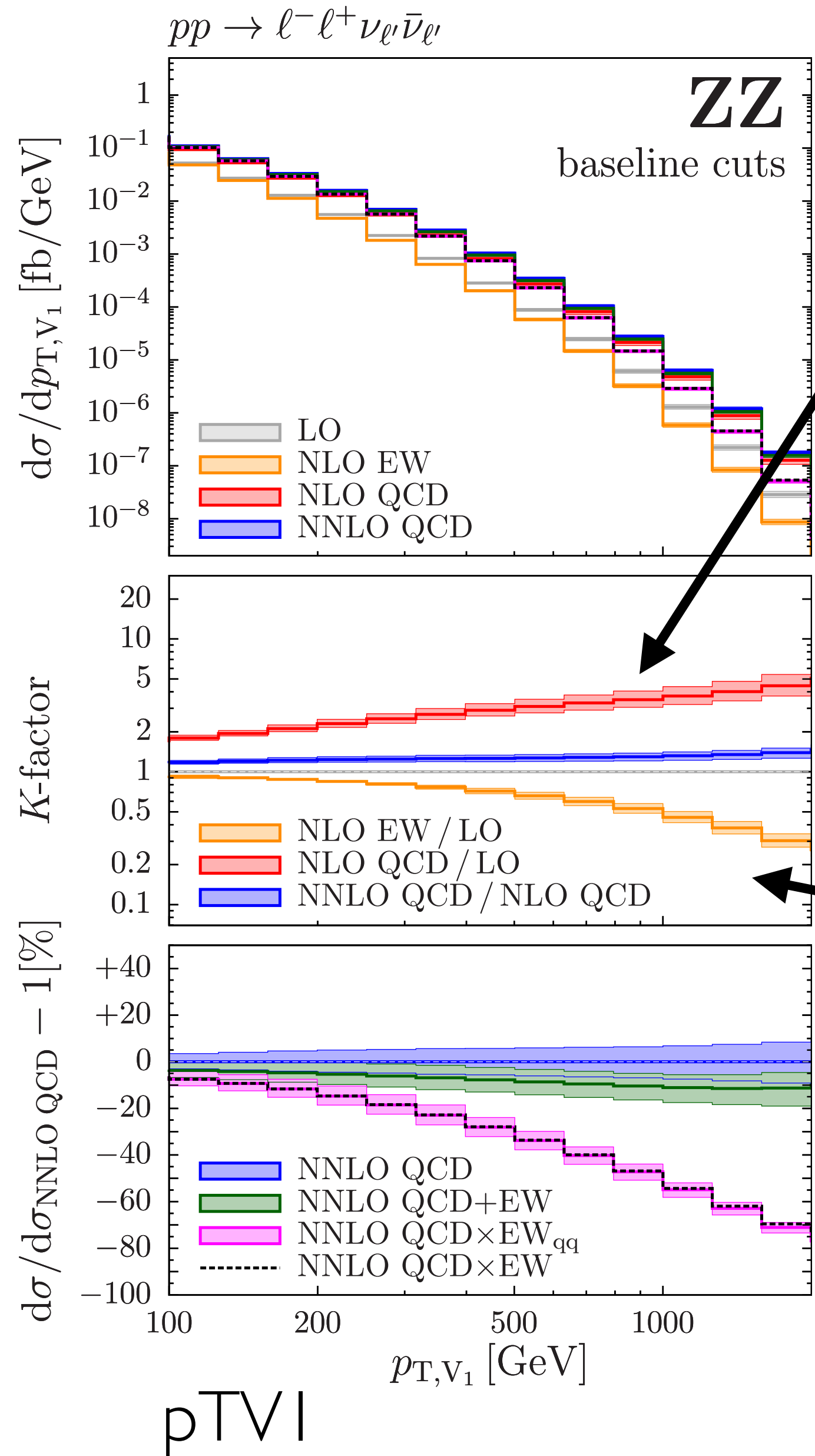
- NLO QCD/LO = 30-70%
- NNLO QCD/NLO = 10-20%
- NLO EW = -30/-20/-20% at 2 TeV

Giant QCD K-factors and EW corrections: pTVI

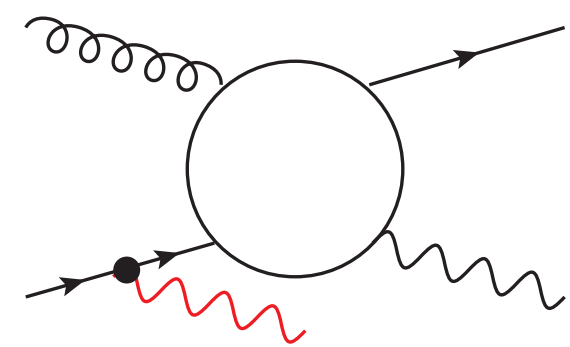


• NLO QCD/LO=2-5! (“giant K-factor”)

Giant QCD K-factors and EW corrections: pTVI



- NLO QCD/LO=2-5! (“giant K-factor”)
- at large pTVI: VV phase-space is dominated by V+jet (w/ soft V radiation)



$$\frac{d\sigma^{V(V)j}}{d\sigma_{VV}^{\text{LO}}} \propto \alpha_S \log^2 \left(\frac{Q^2}{M_W^2} \right) \simeq 3 \quad \text{at } Q = 1 \text{ TeV}$$

- NNLO / NLO QCD moderate and NNLO uncert. 5-10%
- NLO EW/LO=-(40-50)%

• Very large difference $d\sigma_{\text{NNLO QCD+EW}}$ vs. $d\sigma_{\text{NNLO QCD} \times \text{EW}}$

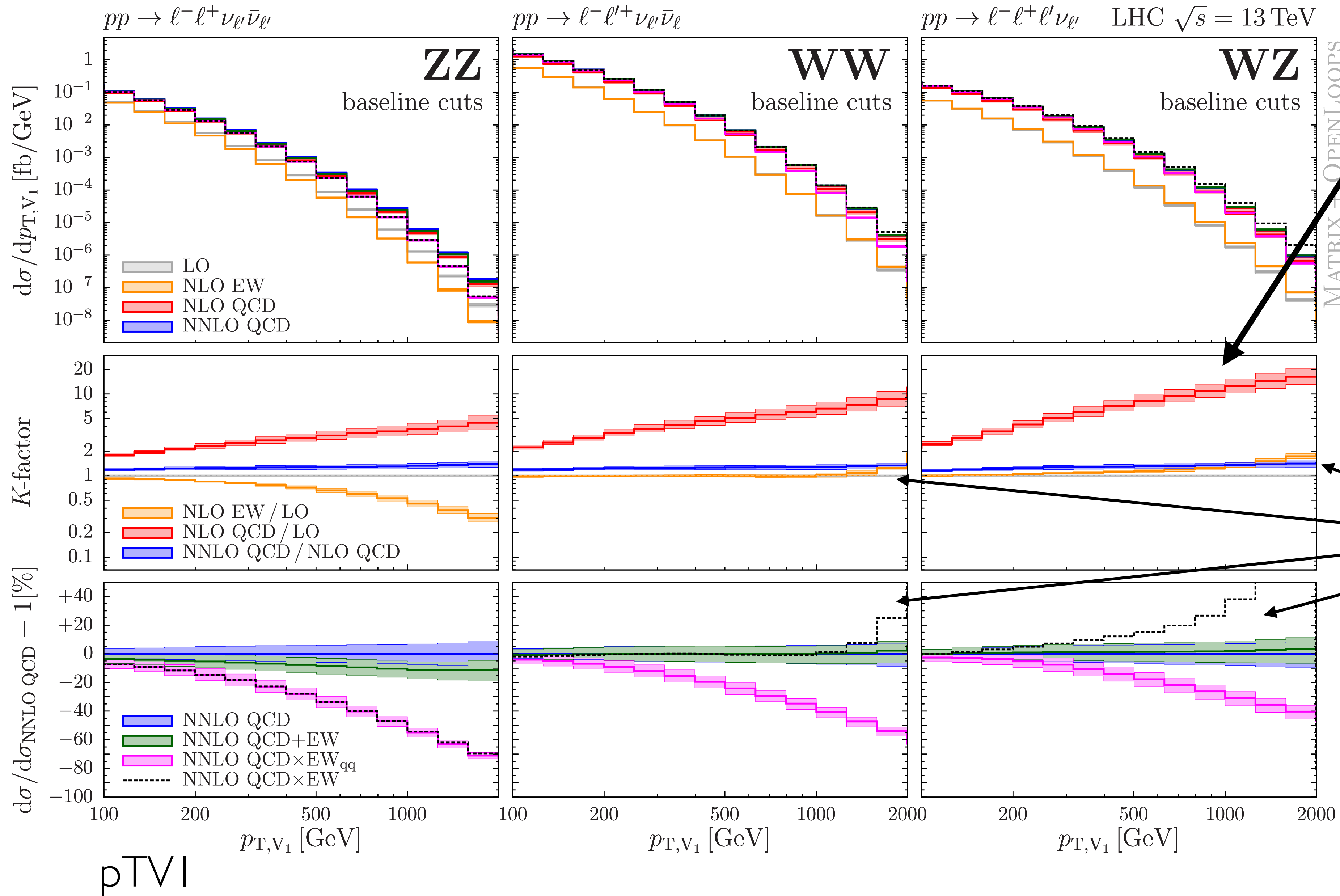
• Problems:

1. In additive combination dominant Vj topology does not receive any EW corrections
2. In multiplicative combination EW correction for VV is applied to Vj hard process

• Pragmatic solution: **take average as nominal and spread as uncertainty**

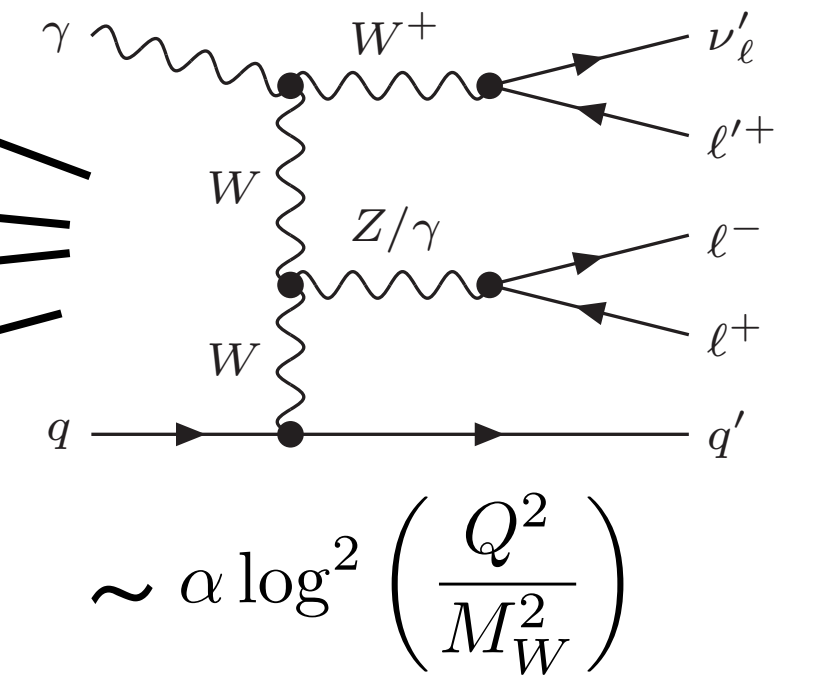
• Rigorous solution: merge VVj incl. EW corrections with VV retaining NNLO QCD + EW

Giant QCD K-factors and EW corrections: pTVI

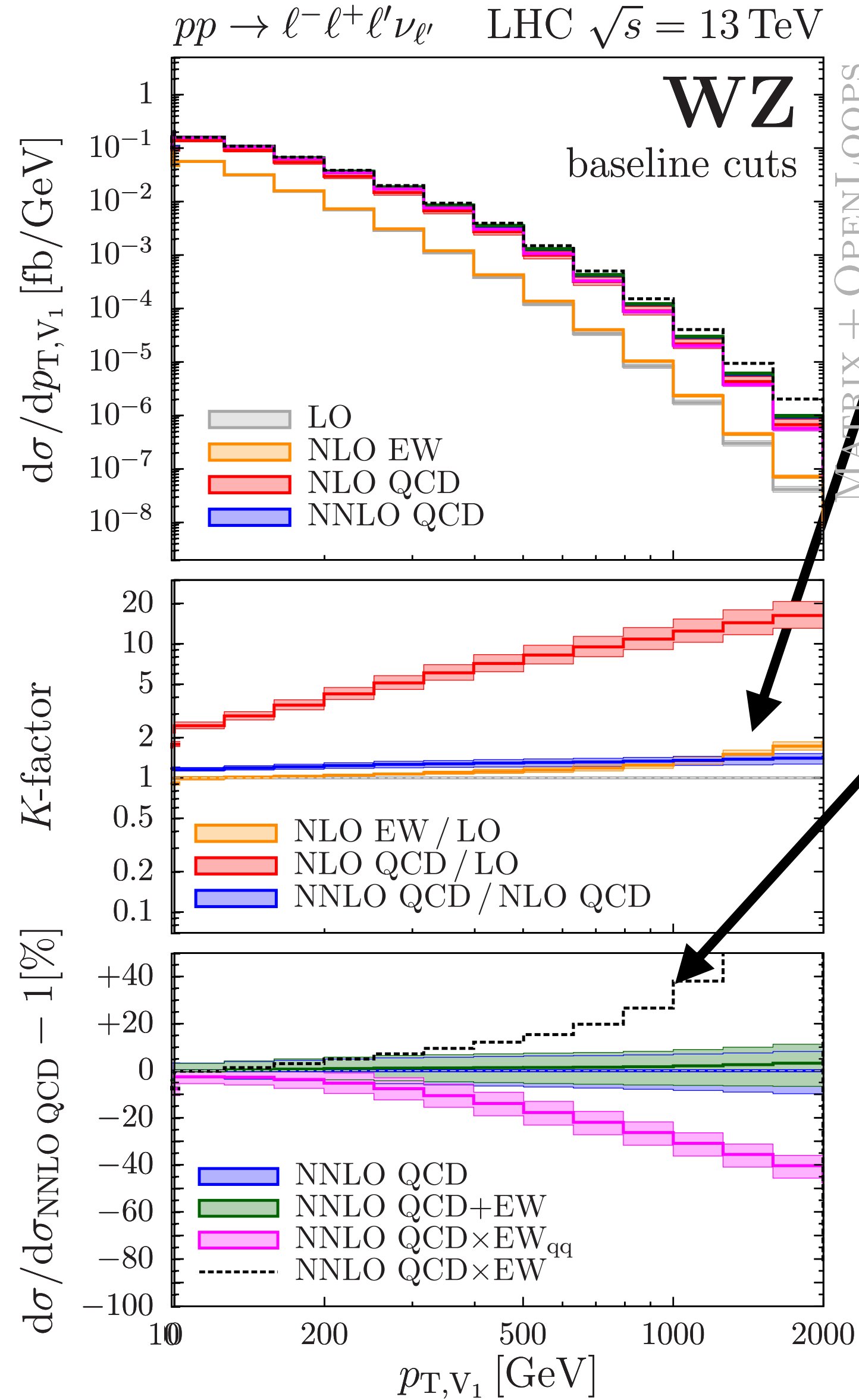


• NLO QCD/LO=5-10!

• Similar giant K-factor mechanism also in gamma-induced:



Giant QCD K-factors and EW corrections: pTVI



- Giant EW K-factor

$$\bullet d\sigma_{\text{NNLO QCD} \times \text{EW}} = \left[d\sigma_{\text{LO}}^{q\bar{q}} \left(1 + \delta_{\text{QCD}}^{q\bar{q}} \right) + d\sigma_{\text{LO}}^{\gamma\gamma} \right] (1 + \delta_{\text{EW}}) + d\sigma_{\text{LO}}^{gg}$$

$$\text{with } \delta_{\text{EW}} = \frac{\delta_{\text{EW}}^{q\bar{q}} d\sigma_{\text{LO}}^{q\bar{q}} + \delta_{\text{EW}}^{\gamma\gamma/\gamma q} d\sigma_{\text{LO}}^{\gamma\gamma}}{d\sigma_{\text{LO}}^{q\bar{q}} + d\sigma_{\text{LO}}^{\gamma\gamma}} \quad \text{averaged EW corr. factor}$$

yields pathological behaviour when δ_{EW} is dominated by giant EW K-factors

- alternative/modified multiplicative ansatz:

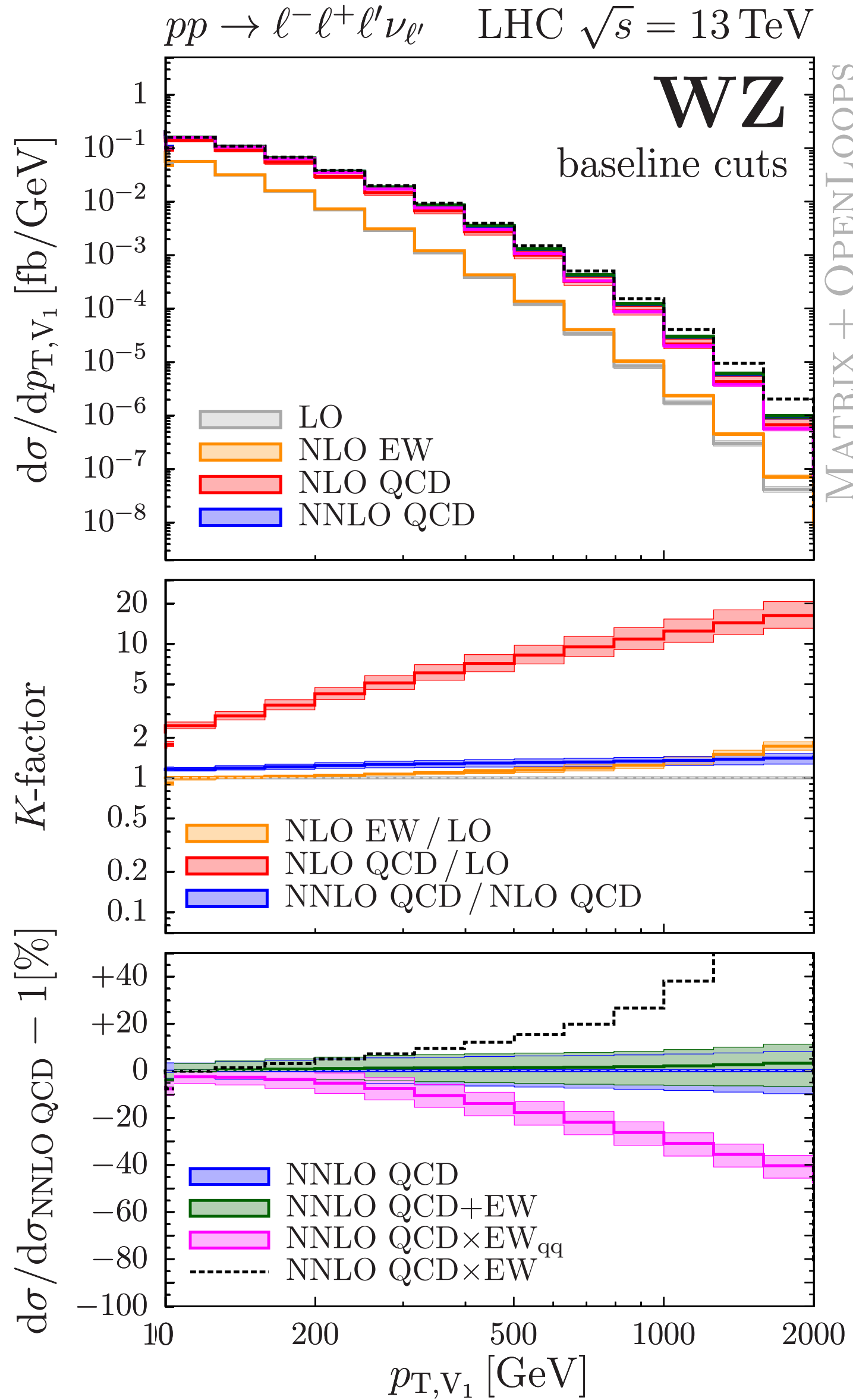
$$\begin{aligned} d\sigma_{\text{NNLO QCD} \times \text{EW}_{\text{qq}}} &= d\sigma_{\text{LO}}^{q\bar{q}} \left(1 + \delta_{\text{QCD}}^{q\bar{q}} \right) \left(1 + \delta_{\text{EW}}^{q\bar{q}} \right) + d\sigma_{\text{LO}}^{\gamma\gamma} \left(1 + \delta_{\text{EW}}^{\gamma\gamma/q\gamma} \right) + d\sigma_{\text{LO}}^{gg} \\ &= d\sigma_{\text{NNLO QCD} + \text{EW}} + d\sigma_{\text{LO}} \delta_{\text{QCD}} \delta_{\text{EW}}^{q\bar{q}} \end{aligned}$$

yields behaviour consistent with EW Sudakov logs

- Thus: discard $d\sigma_{\text{NNLO QCD} \times \text{EW}}$

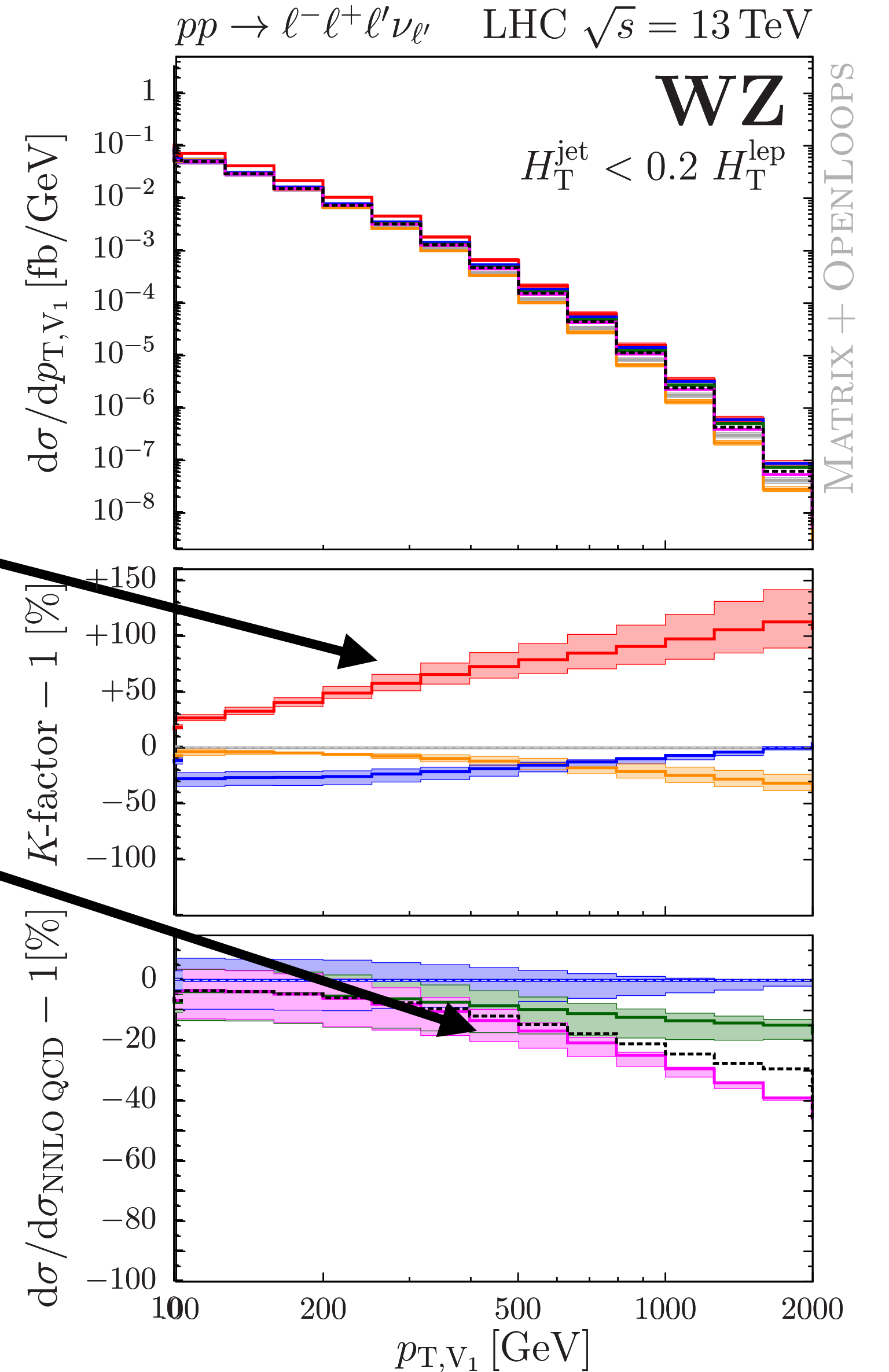
- Caveat: splitting in $q\bar{q}$ and $\gamma\gamma/\gamma q$ channels is ad-hoc/scheme dependent

Giant QCD K-factors and EW corrections: pTVI

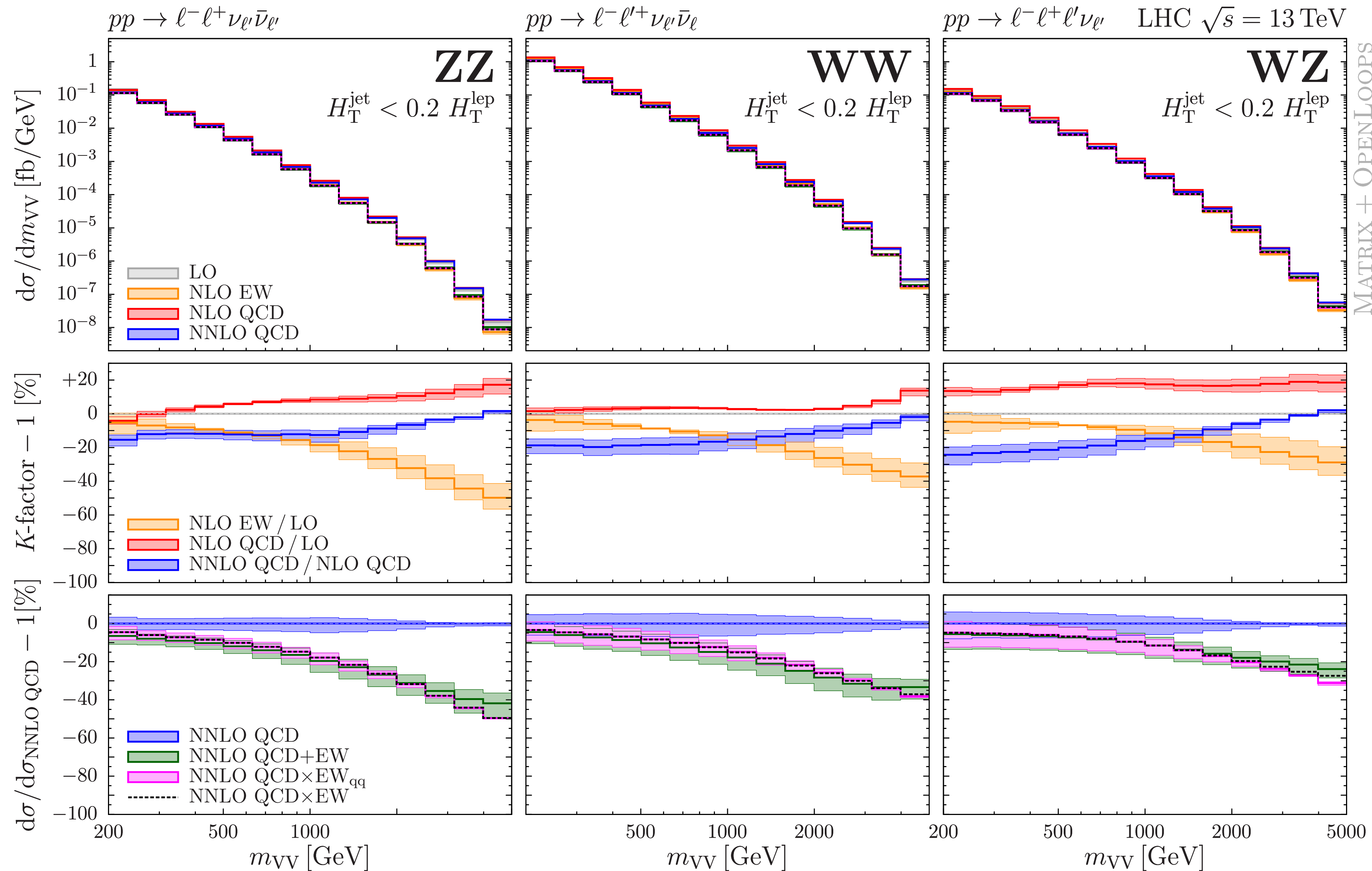


jet veto
 $H_T^{\text{jet}} < 0.2 H_T^{\text{lep}}$

- NLO QCD/LO = $\sim < 1.5$ (“normal K-factor”)
- small differences between the two multiplicative combinations
 ➔ check!
- consistent results for all processes
- In case we are now really dominated by VV topologies: multiplicative dominations should be seen as superior
 ➔ check!



NNLO QCD + NLO EW for dibosons: m_{VV}

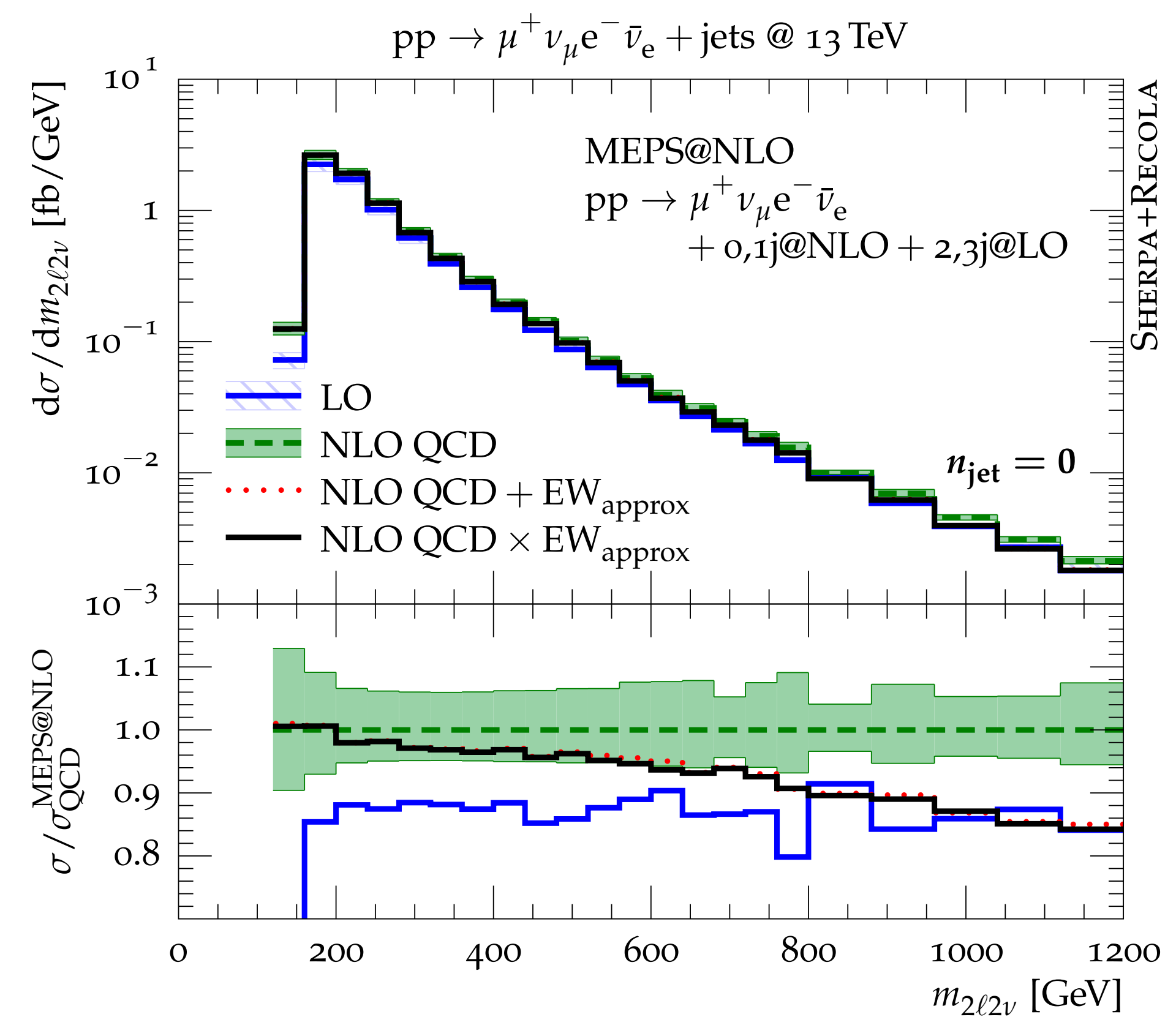
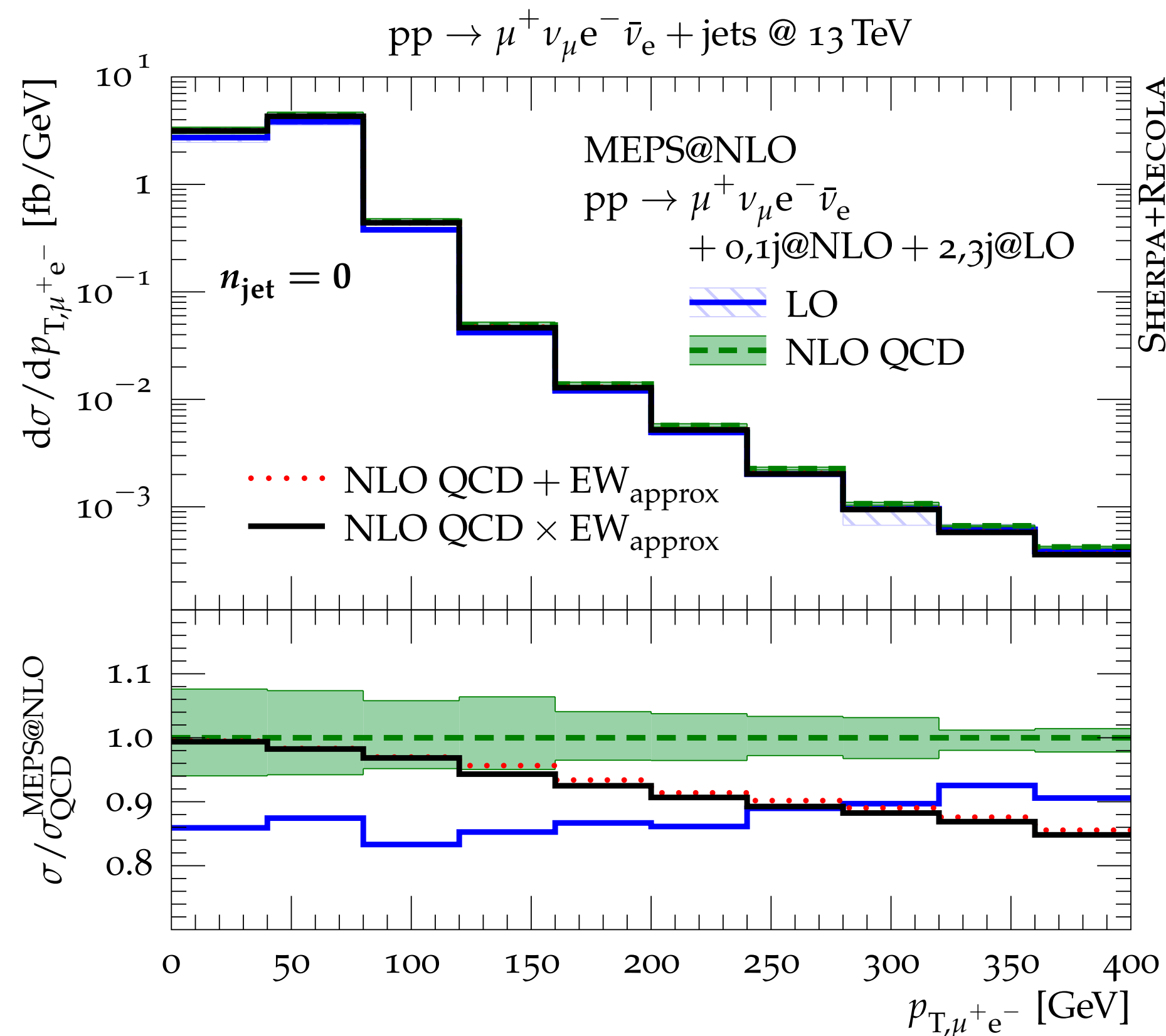


- NLO QCD/LO reduced to 10-20%
- Very small difference between additive and multiplicative QCD-EW combinations

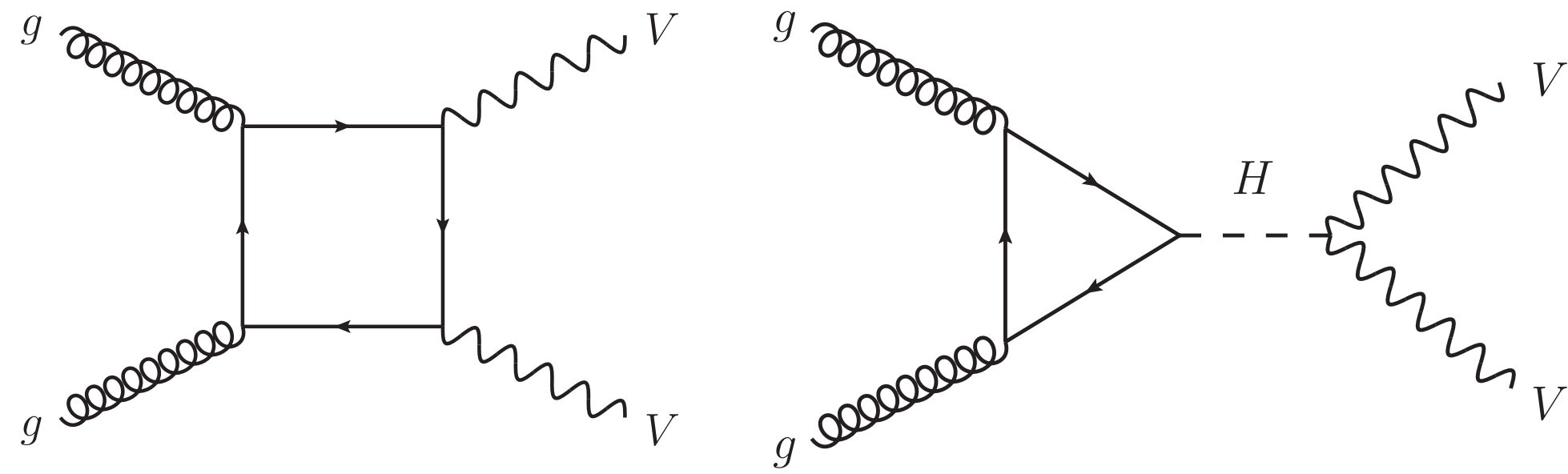
MEPS @ NLO QCD + EW

[Bräuer, Denner, Pellen, Schönherr, Schumann; '20]
 [Bothmann, Napoletano, Schönherr, Schumann, Villani; '21]

- “Rigorous approximate solution”: merge VVj incl. approx. EW corrections with VV with Sherpa’s MEPS@NLO

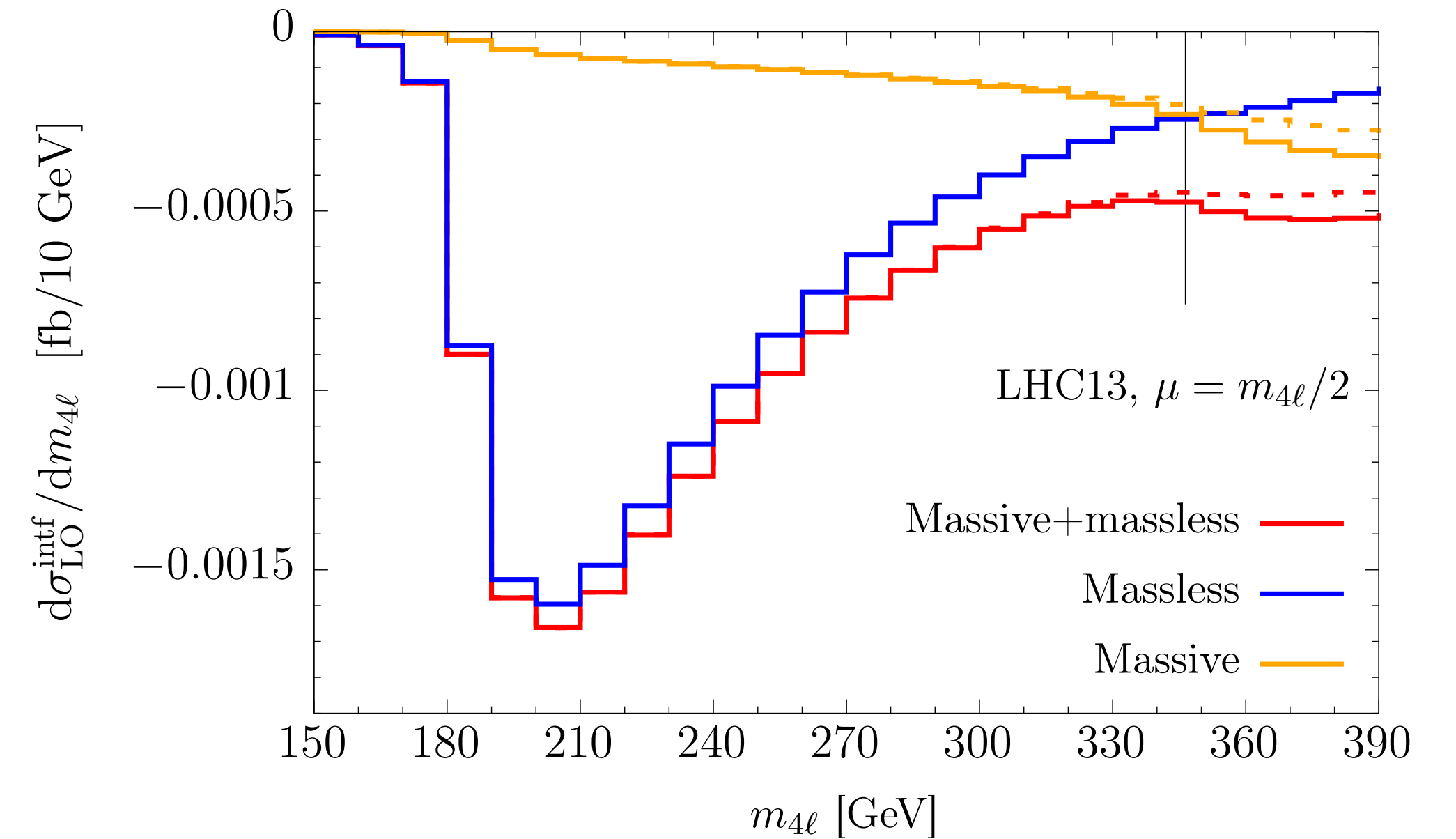
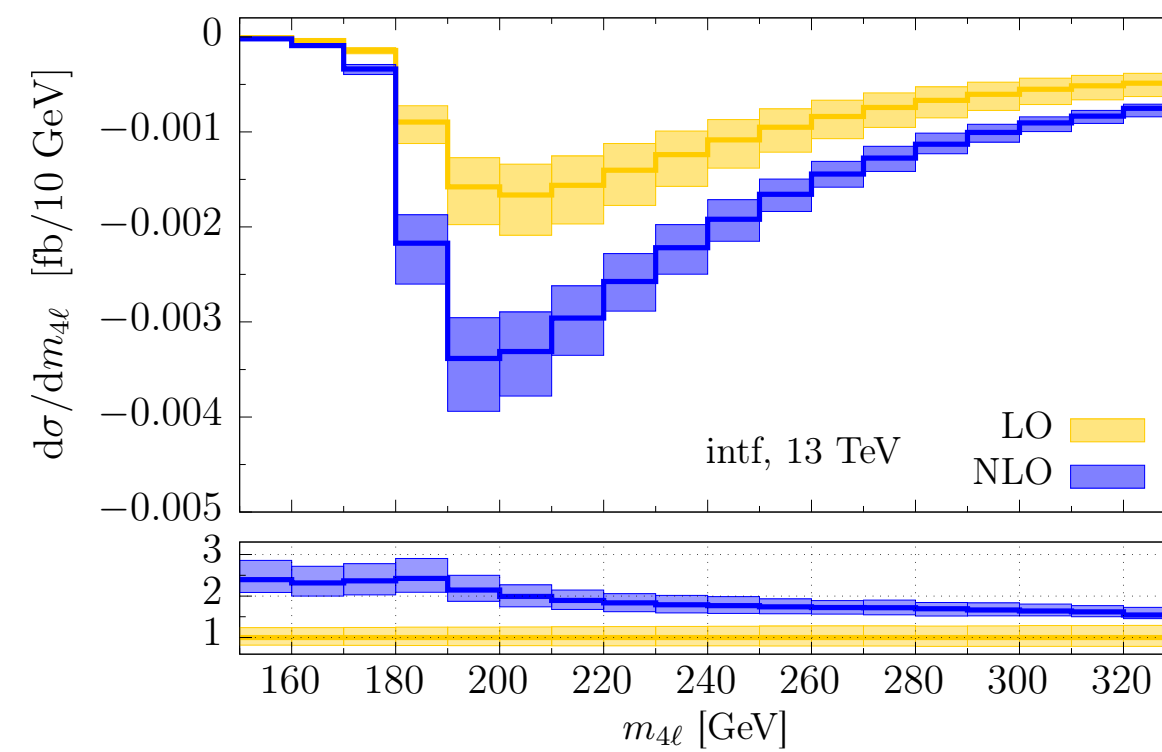
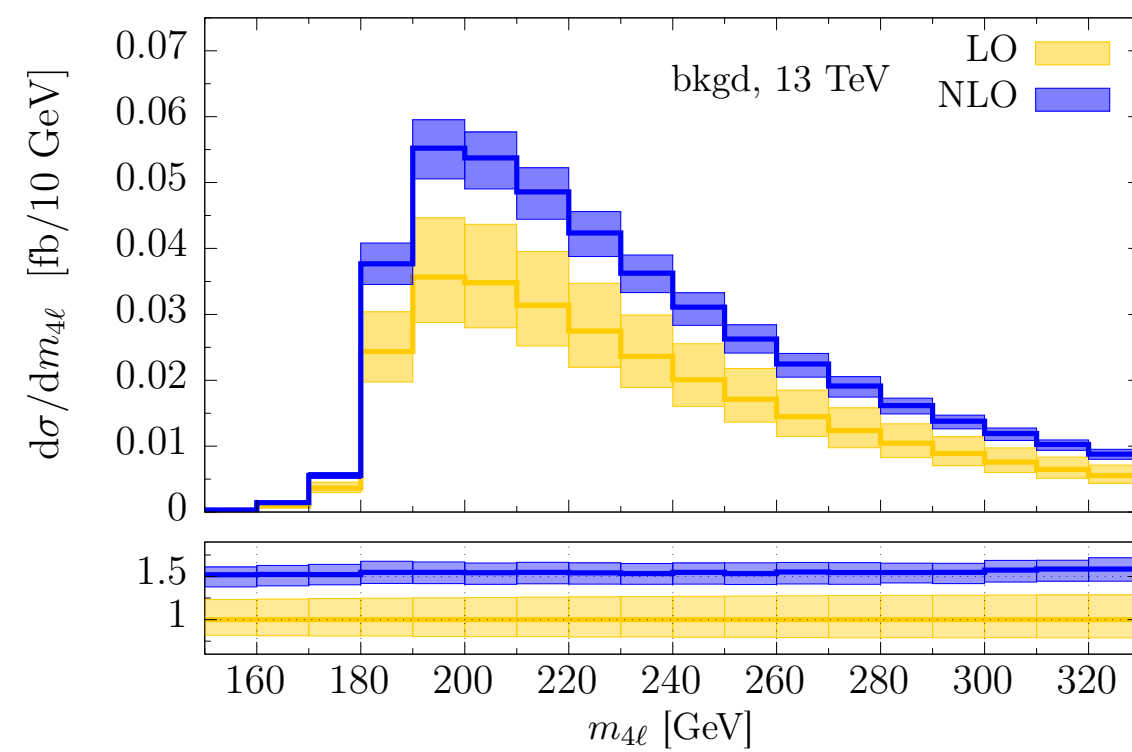
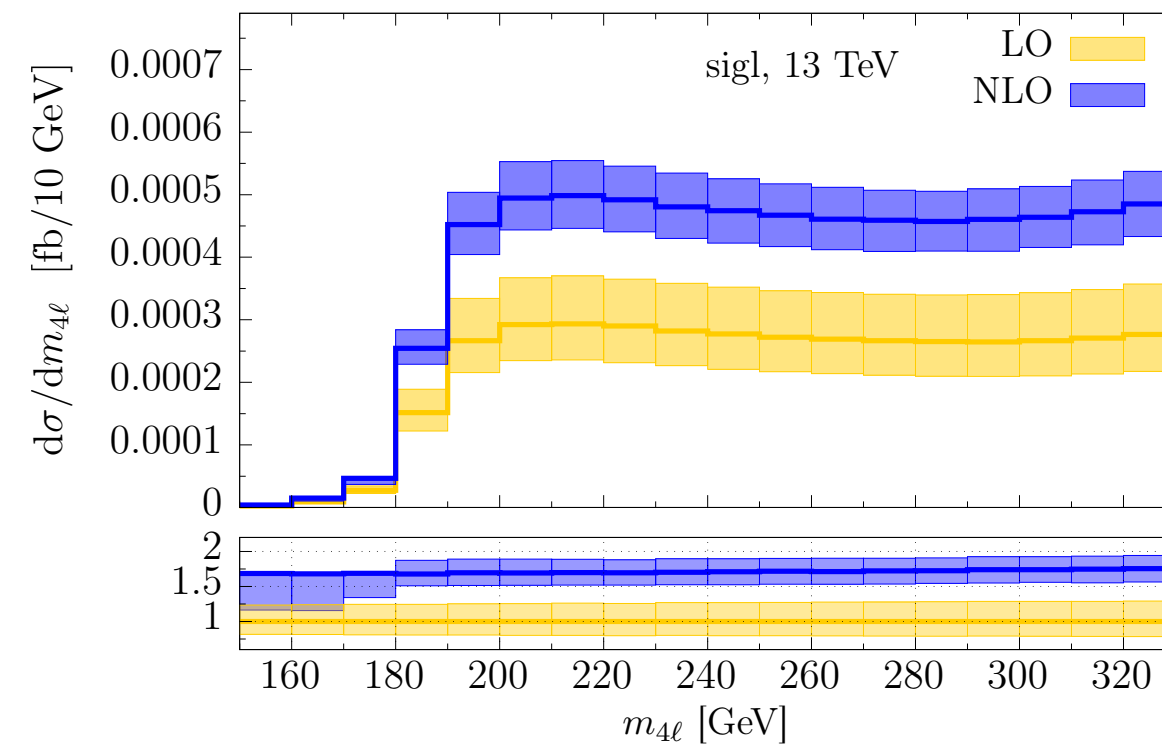
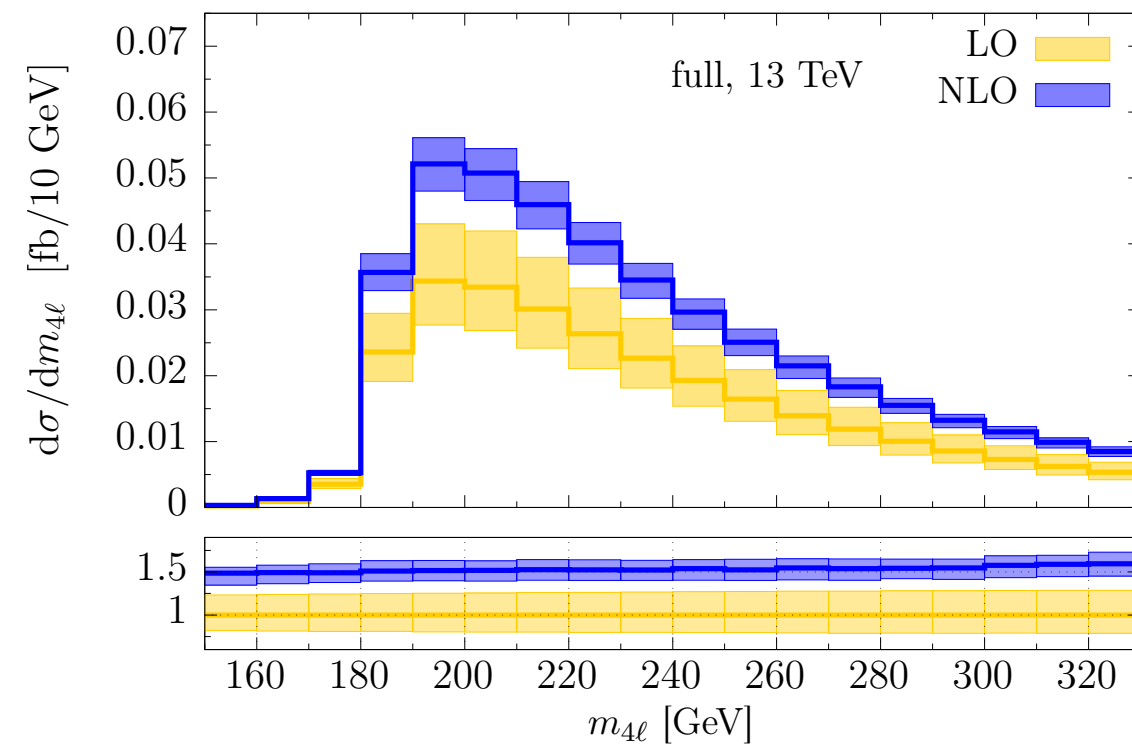


gg-induced WW and ZZ production



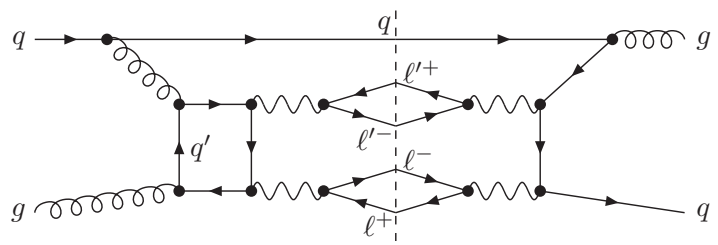
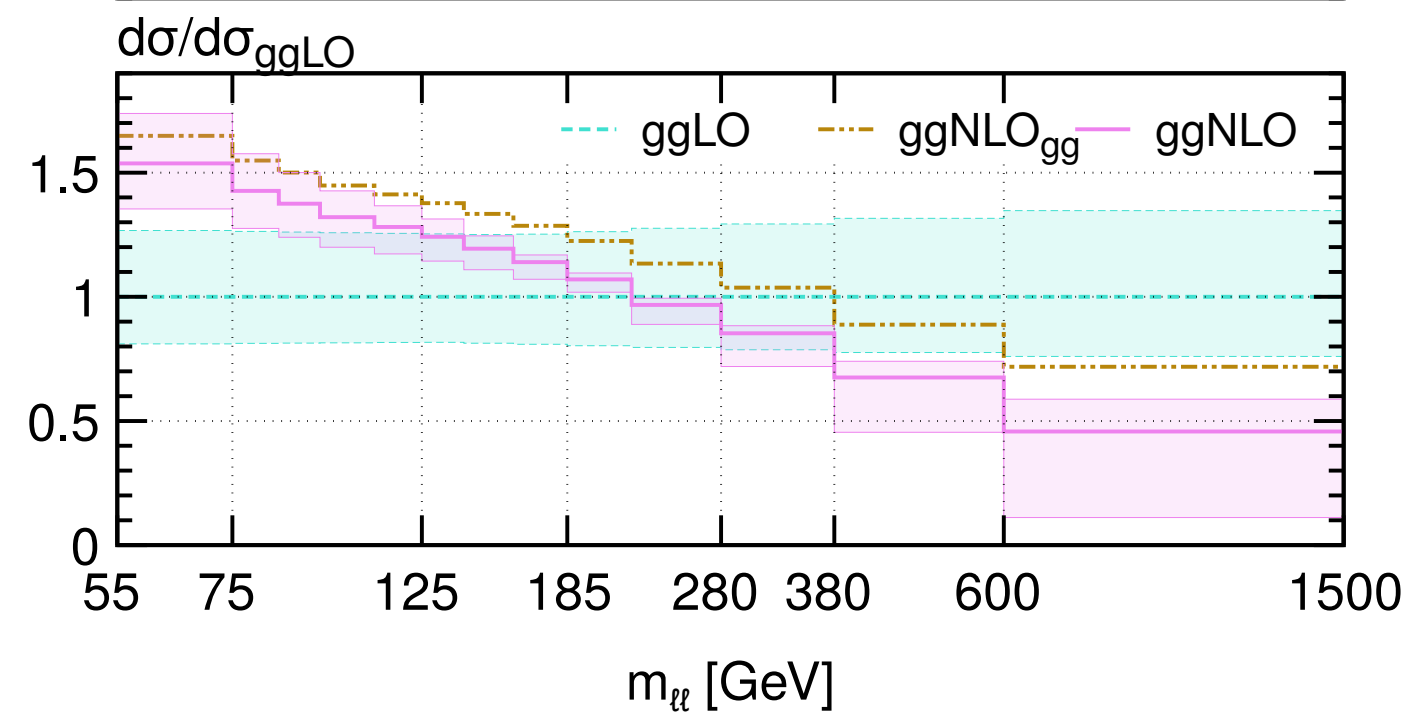
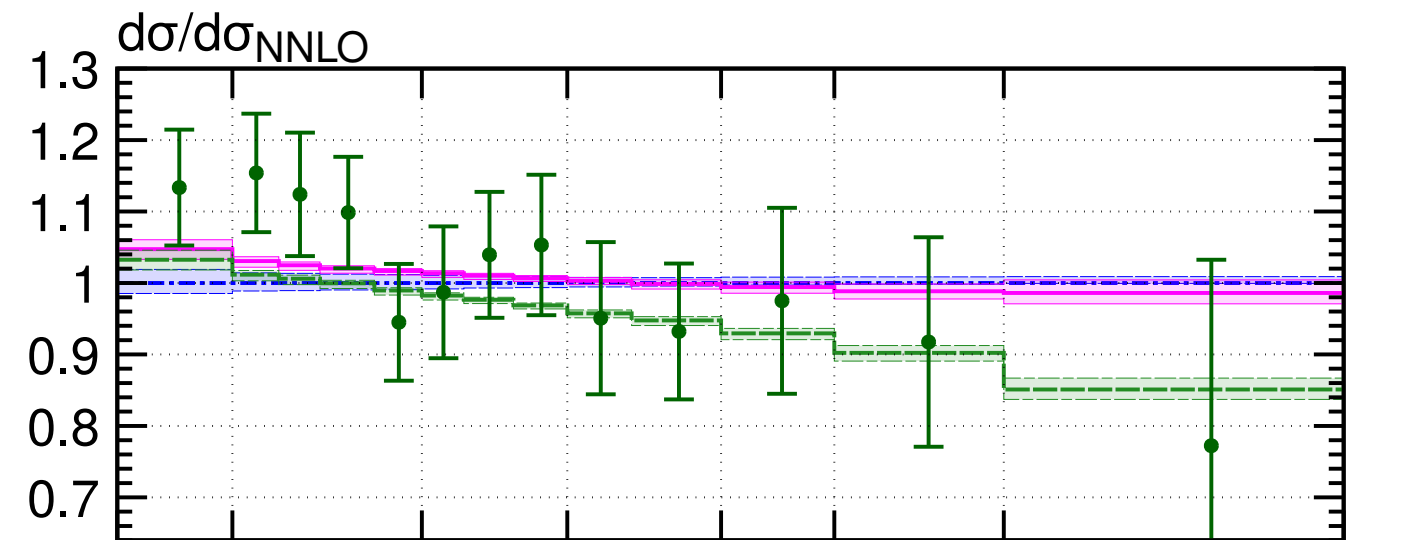
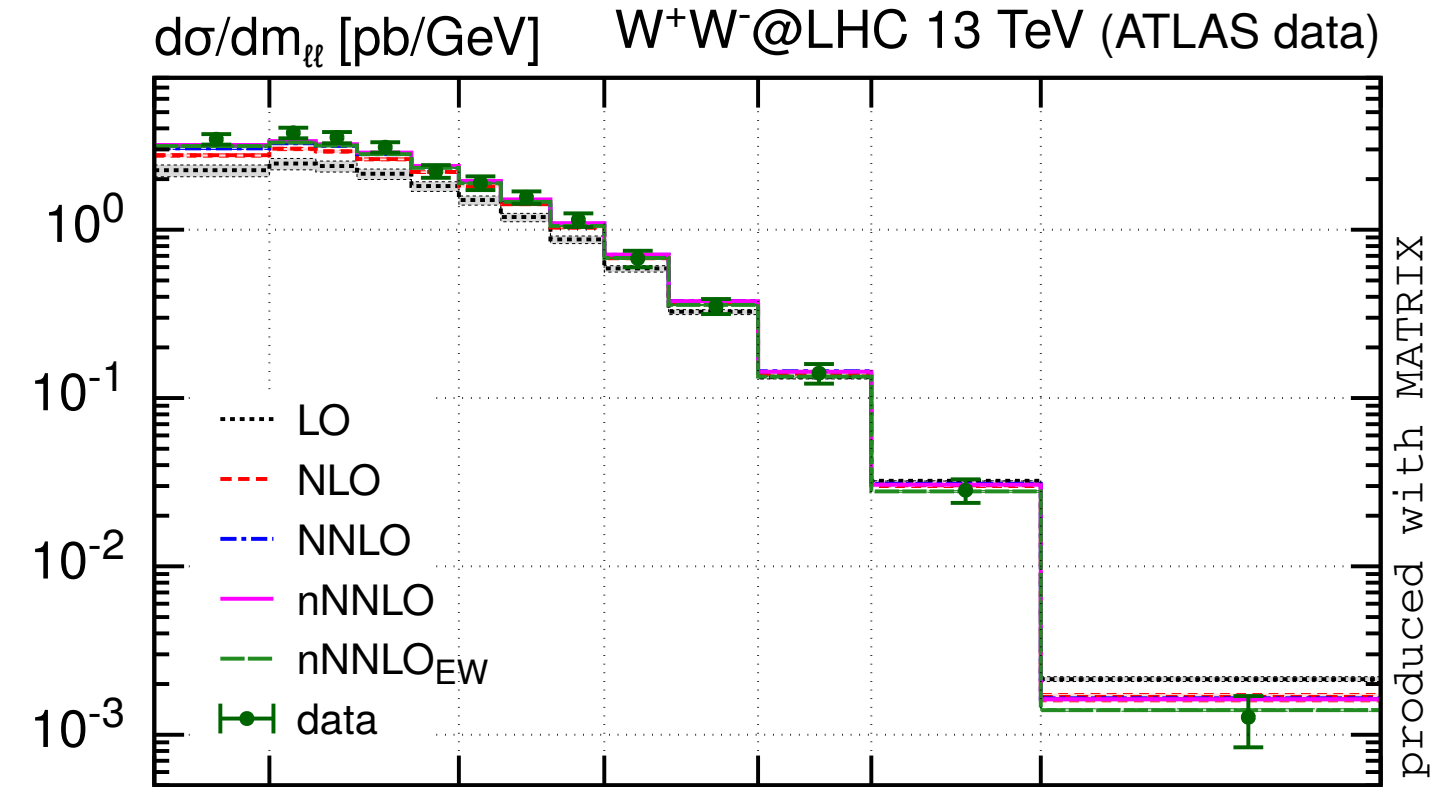
- Formally same order as NNLO QCD
- Enhanced due to gg flux
- Interference with H->VV

- Sizeable QCD corrections (formally N3LO QCD)
- For $m_{4\ell} < 340$ GeV $1/M_t$ expansion reliable



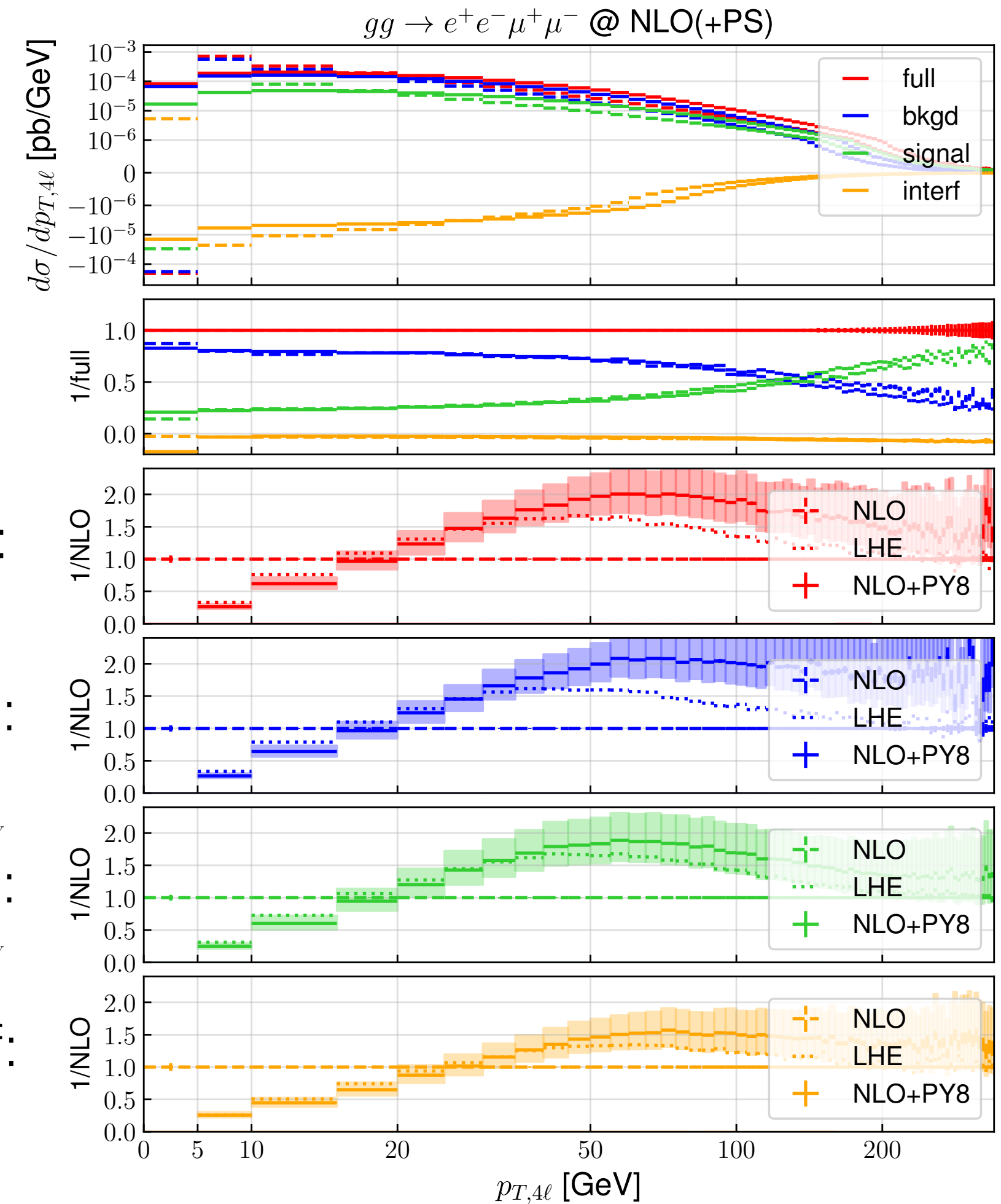
NLO QCD_{gg}

[M. Grazzini, S. Kallweit, J.Y. Yook, M. Wiesemann; WW: '20, ZZ: '21]

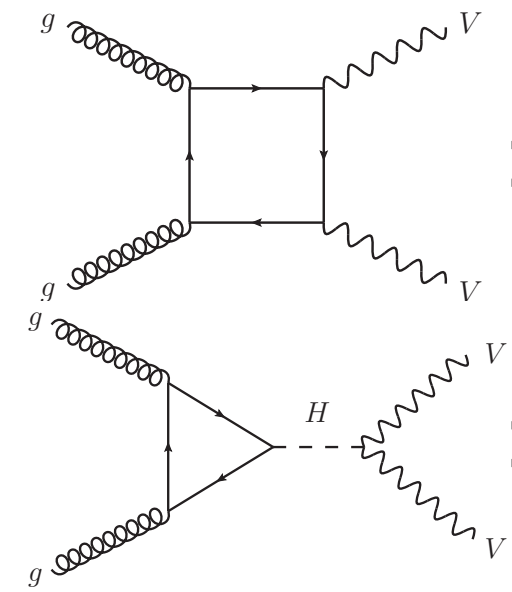


- Very good data agreement with NNLO QCD + NLO QCD_{gg} + NLO EW

[Alioli, Ferrario Ravasio, JML, Röntschi, '21]



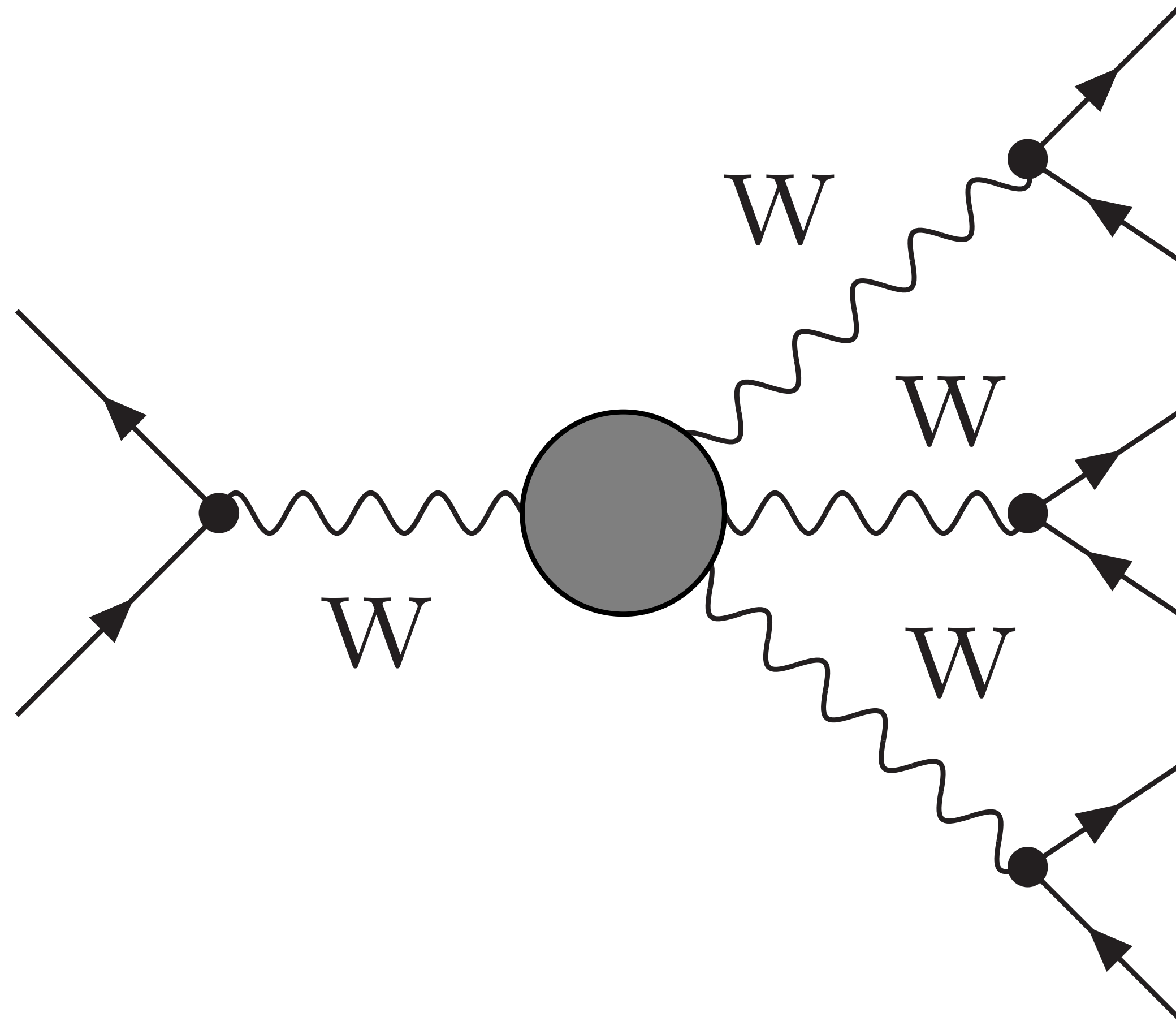
sum:



interf:

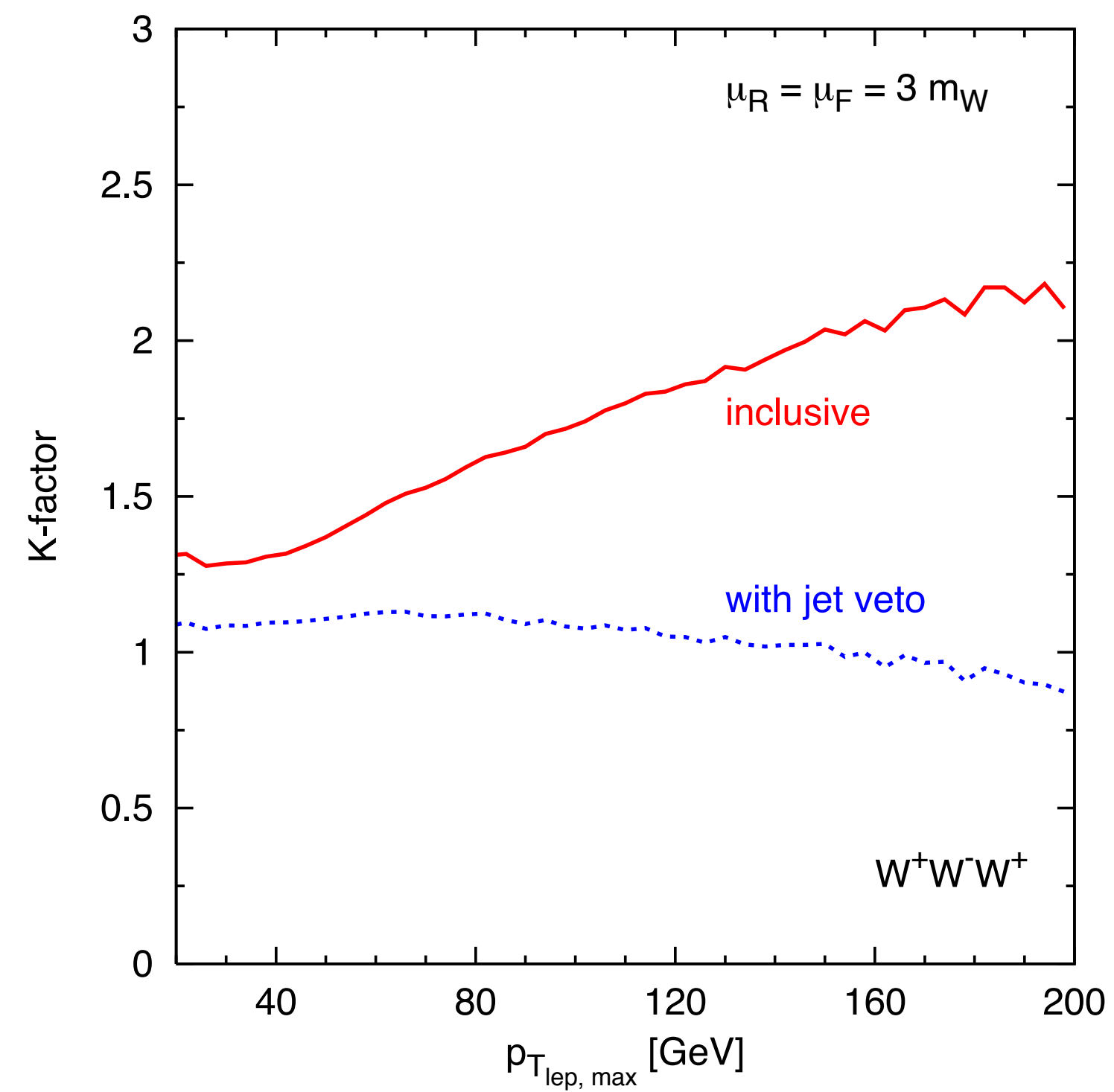
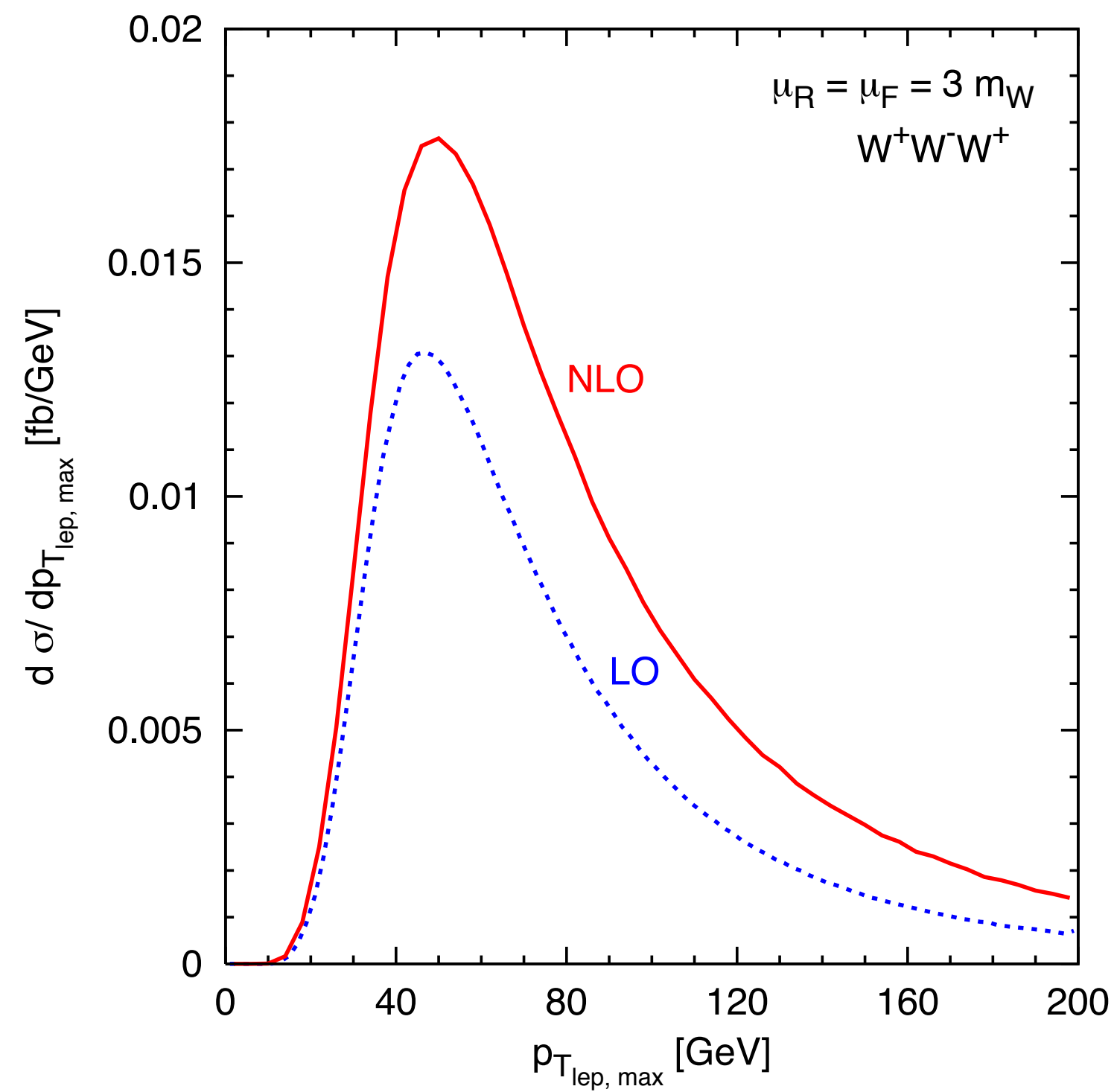
- ggWW/ggZZ @ NLO QCD + PS available! (VV-cont., H → VV & interference)

Tribosons



Triboson production @ NLO QCD

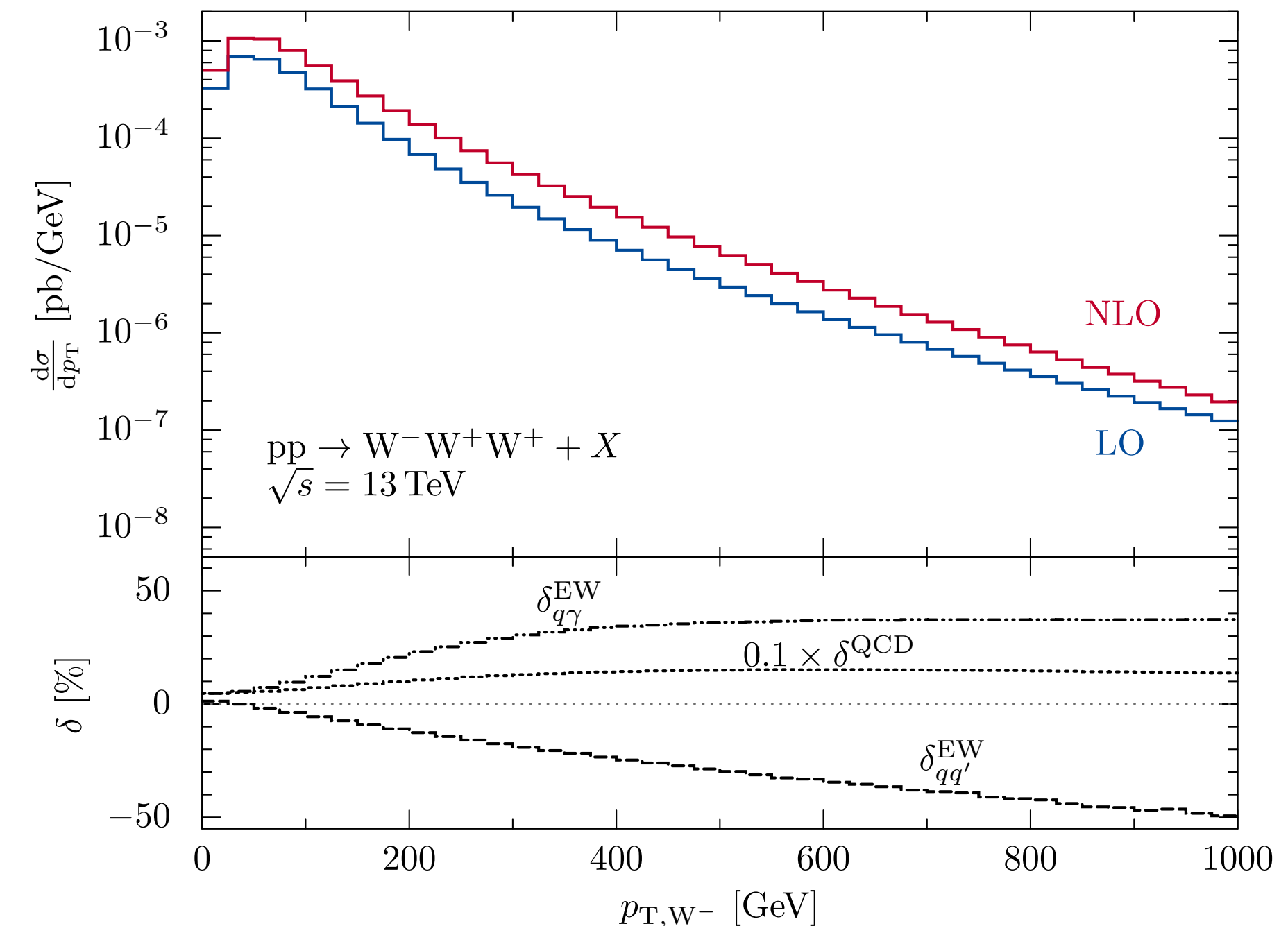
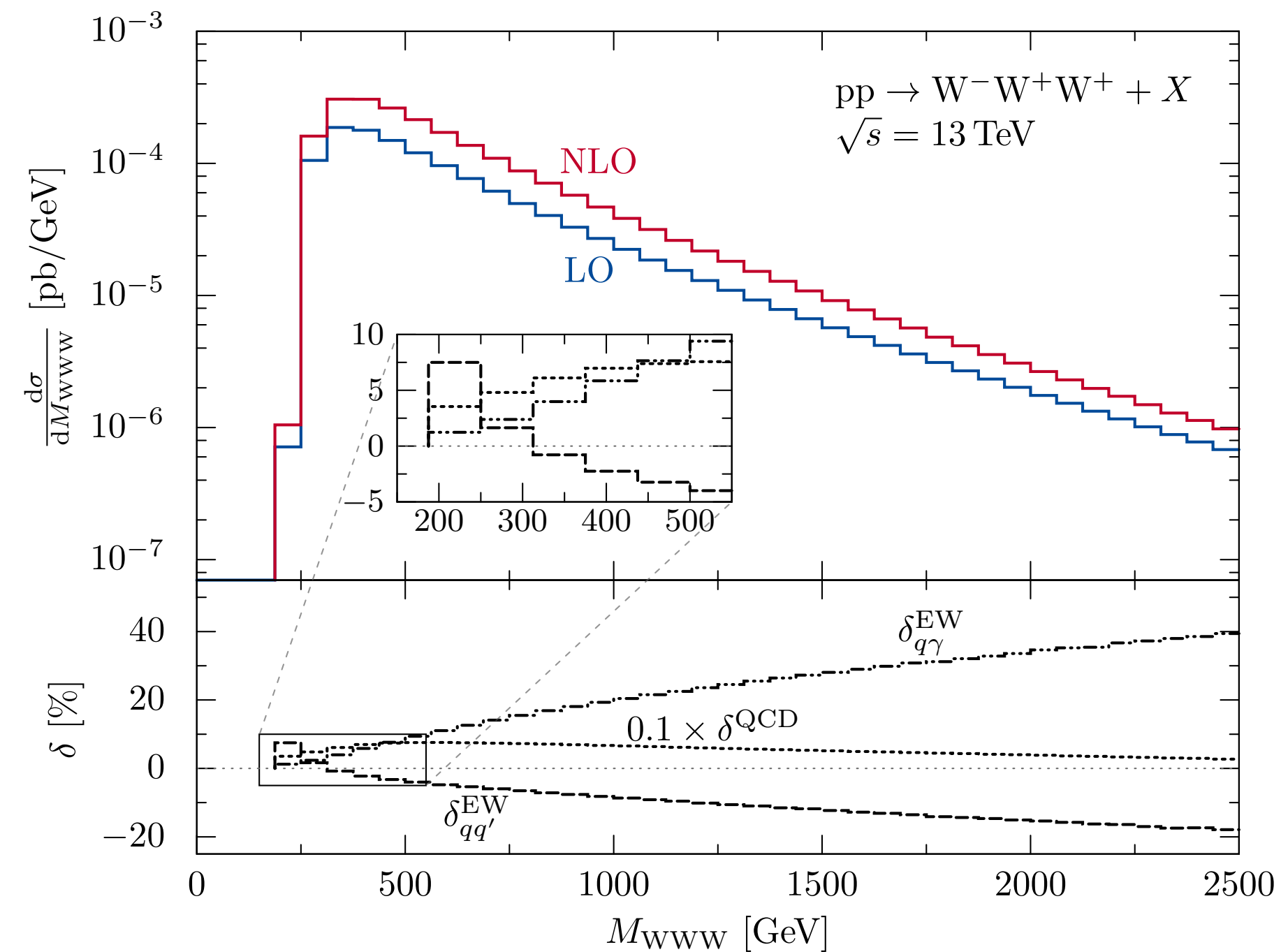
[Campanario et.al., '08]



- QCD correction driven by additional jet activity: VV +jet topologies with soft V
 - 'giant K-factors'
 - strong observable dependence
 - NLO mandatory
- jet veto ($p_{T_{cut}} = 50$ GeV) reduces size and phase space dependence
 - better: multi-jet merging

WWW @ NLO QCD+EW

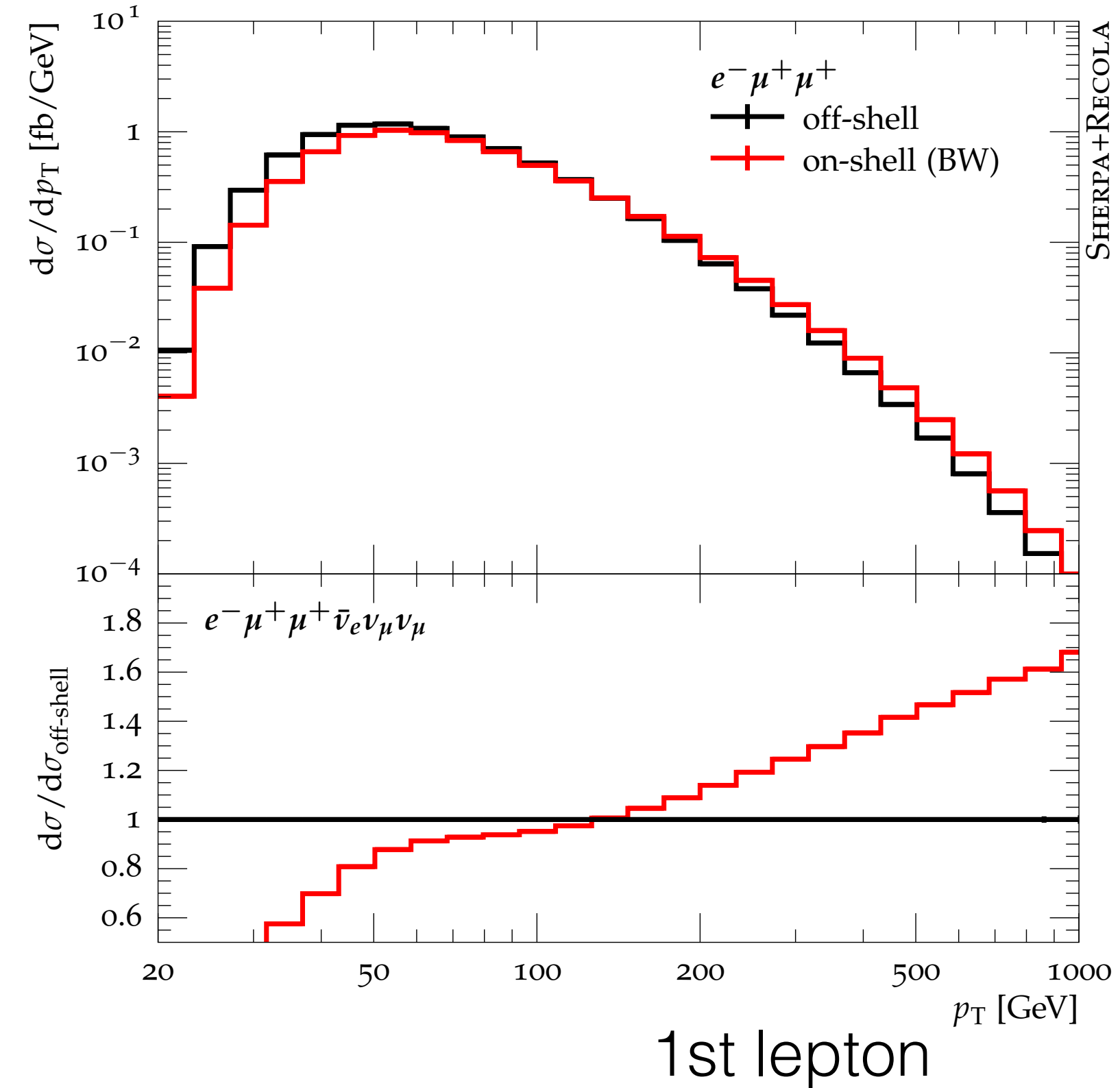
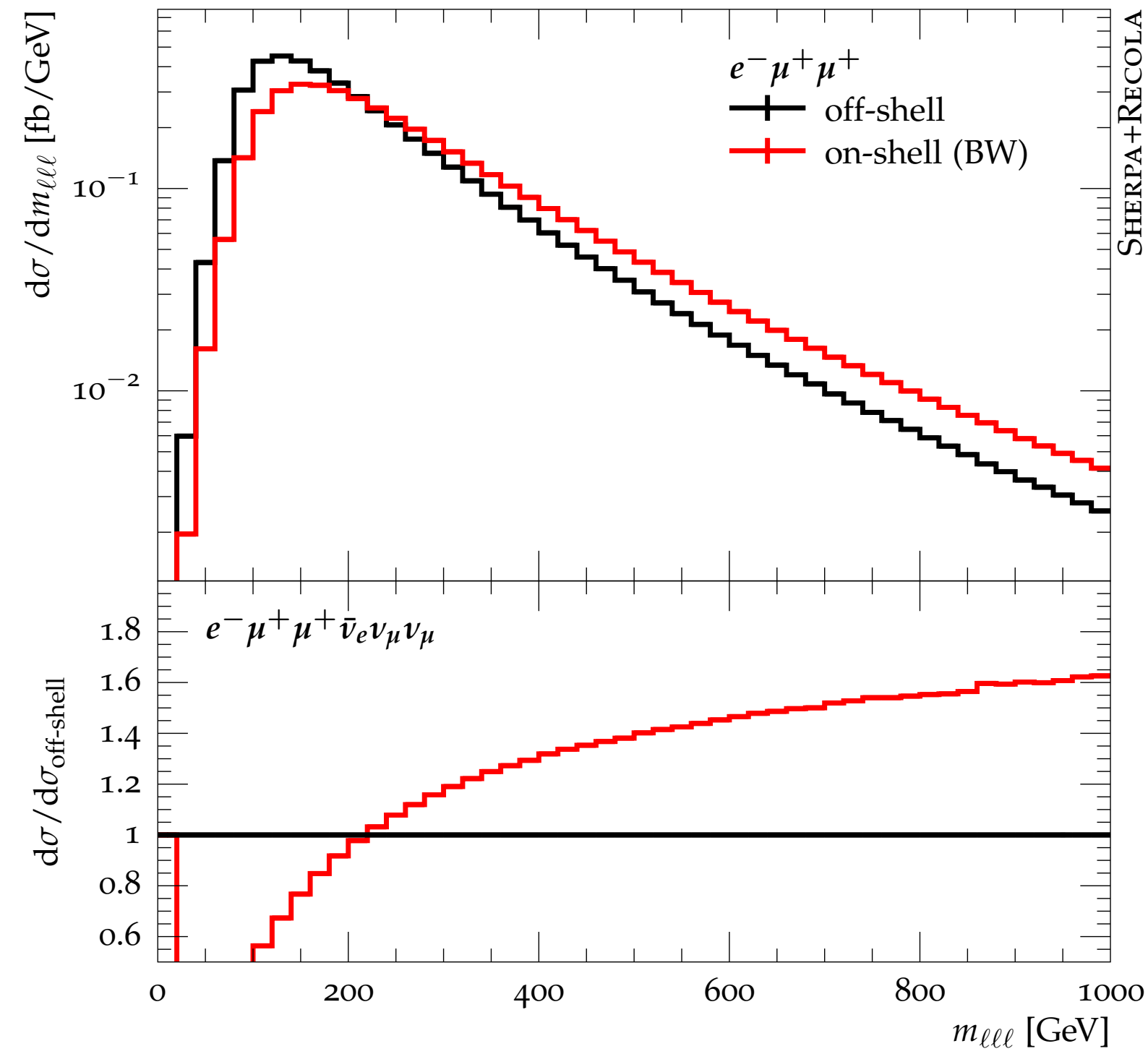
Dittmaier, Huss, Knippen arXiv:1705.03722



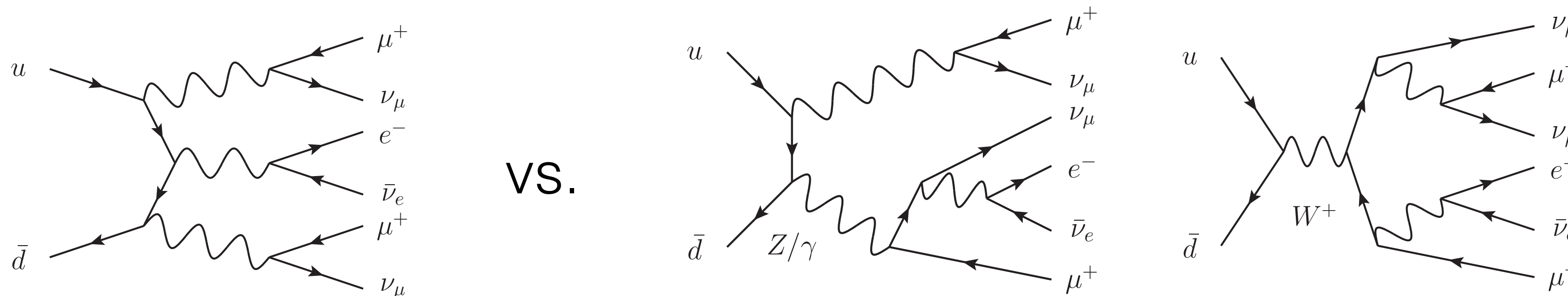
- corrections w/o jet veto
- QCD corrections $\approx 70\%$, slight observable dependence
- γ -induced EW corrections large and observable dependent
 \rightarrow large accidental cancellations with EW corrections in $q\bar{q}$ -channel

Triboson production: on-shell vs. off-shell

[M. Schönherr, '18]



- at large $m_{\ell\ell}$ and $p_{T,\ell\ell}$ large interference with other resonance structures



Off-shell VVV(3l+MET) production @ NLO EW

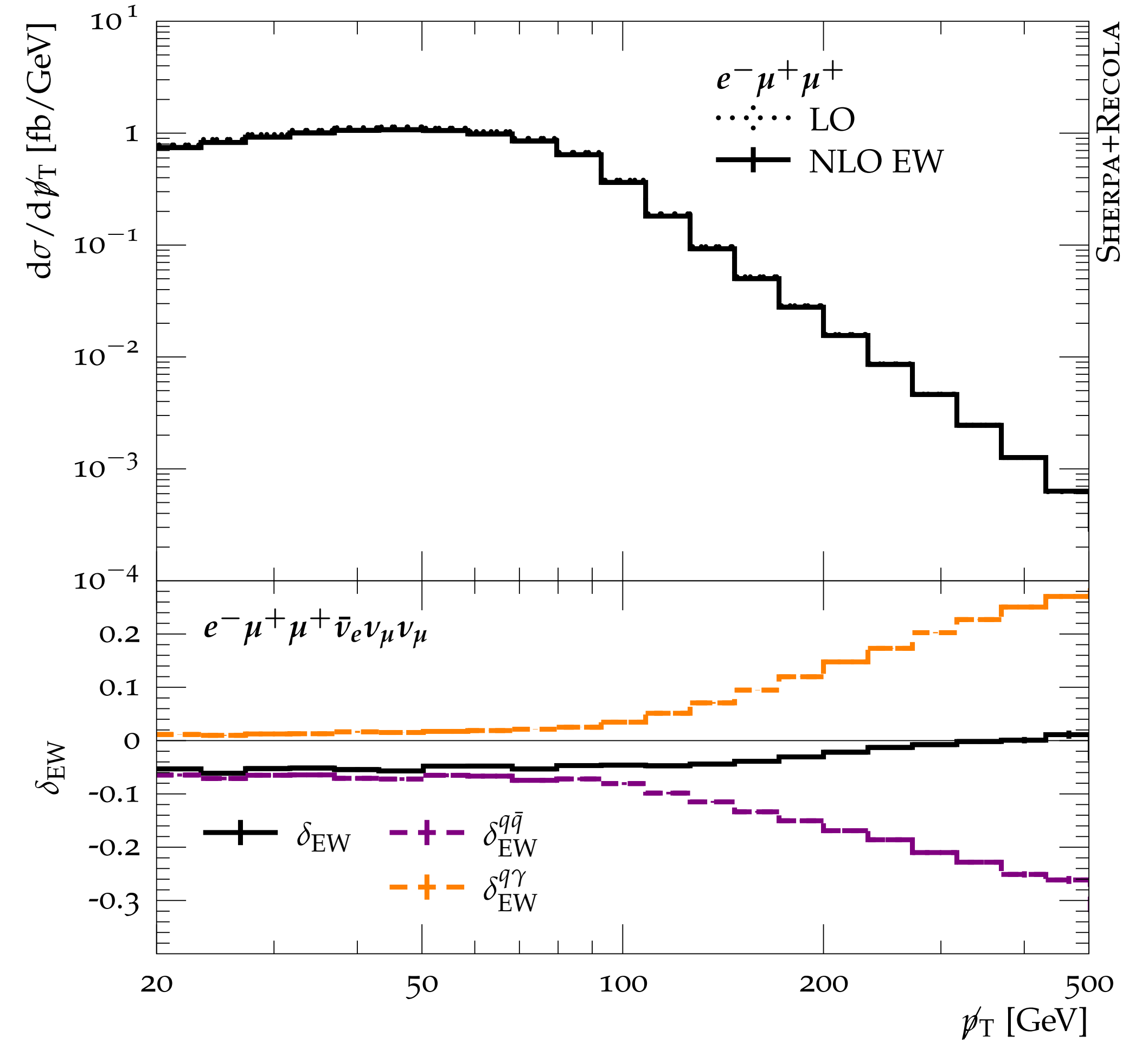
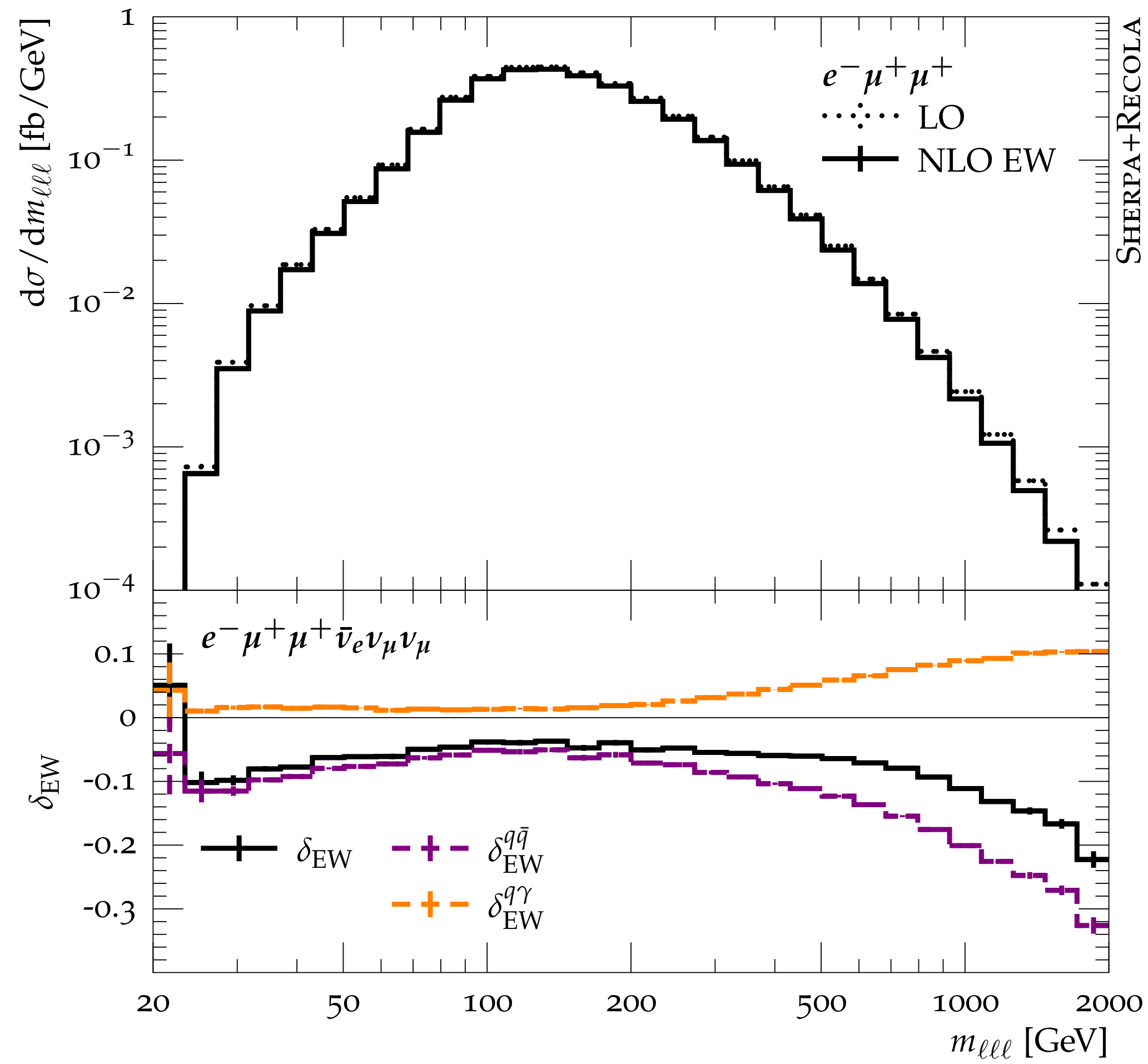
[M. Schönherr, '18]

	inclusive			
	LO [fb]	δ_{EW}	$\delta_{q\bar{q}}^{EW}$	$\delta_{q\gamma/\bar{q}\gamma}^{EW}$
$l^- l^+ l^+$	0.4209	-2.0 %	-5.2 %	3.2 %
$e^- e^+ e^+$	0.0212	-3.4 %	-7.1 %	3.6 %
$e^- e^+ e^+ \bar{\nu}_e \nu_e \nu_e$	0.0206	-3.4 %	-7.0 %	3.6 %
$e^- e^+ e^+ \bar{\nu}_{\mu/\tau} \nu_{\mu/\tau} \nu_e$	0.0006	-5.4 %	-9.5 %	4.1 %
$e^- e^+ \mu^+$	0.0938	-1.4 %	-5.4 %	4.1 %
$e^- e^+ \mu^+ \bar{\nu}_e \nu_e \nu_\mu$	0.0924	-1.4 %	-5.4 %	4.1 %
$e^- e^+ \mu^+ \bar{\nu}_\mu \nu_\mu \nu_\mu$	0.0007	-2.9 %	-6.1 %	3.2 %
$e^- e^+ \mu^+ \bar{\nu}_\tau \nu_\tau \nu_\mu$	0.0007	-2.7 %	-6.2 %	3.5 %
$e^- \mu^+ \mu^+$	0.0955	-2.2 %	-4.6 %	2.4 %
$e^- \mu^+ \mu^+ \bar{\nu}_e \nu_\mu \nu_\mu$	0.0955	-2.2 %	-4.6 %	2.4 %

- large accidental and cut dependent cancellations of Sudakov-type neg. EW corrections and γ -induced pos. contribs w/ extra jet activity
- **WWW** channels receive smaller corrections than pure **WZZ** channels

Off-shell VV production @ NLO EW

[M. Schönherr, '18]



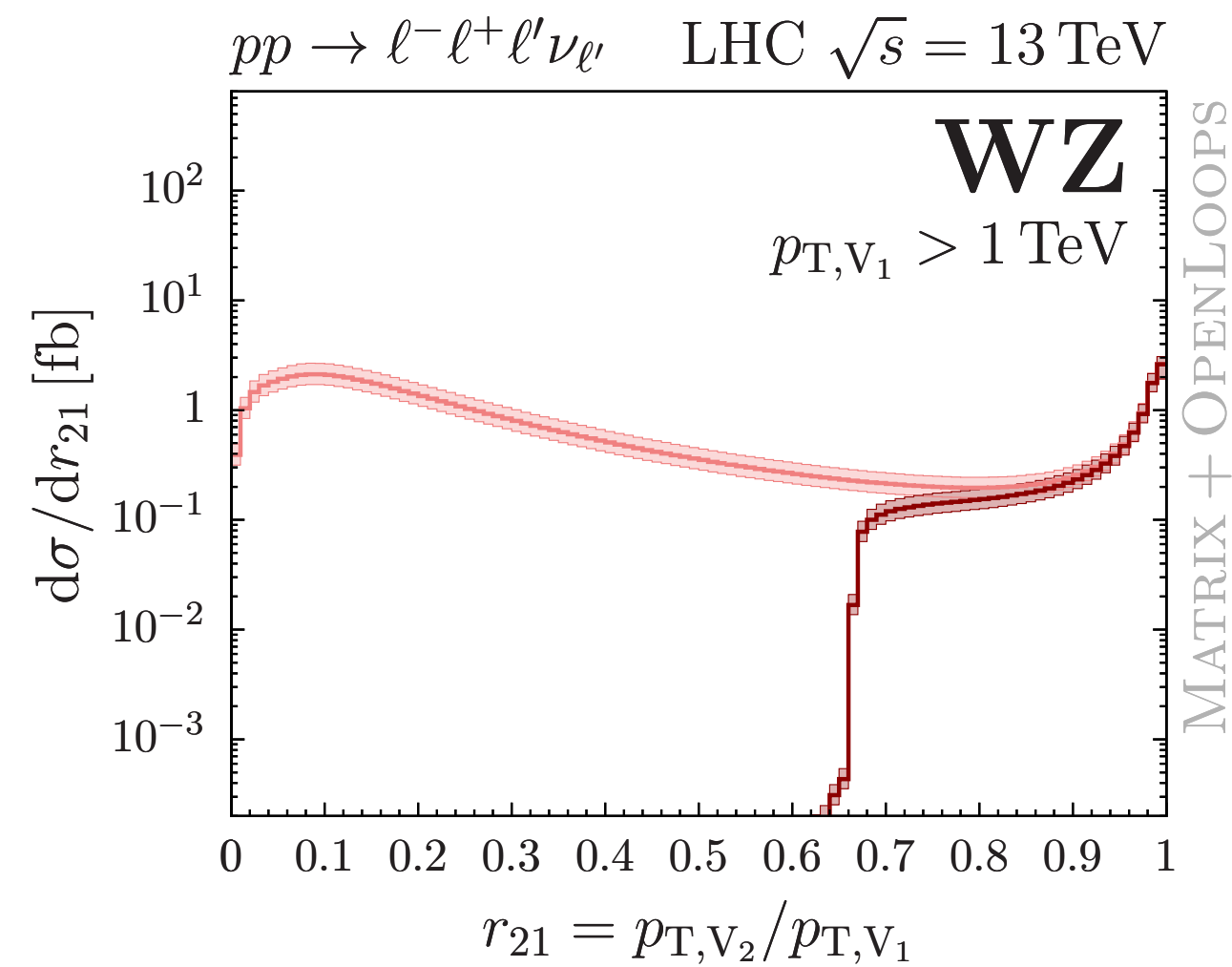
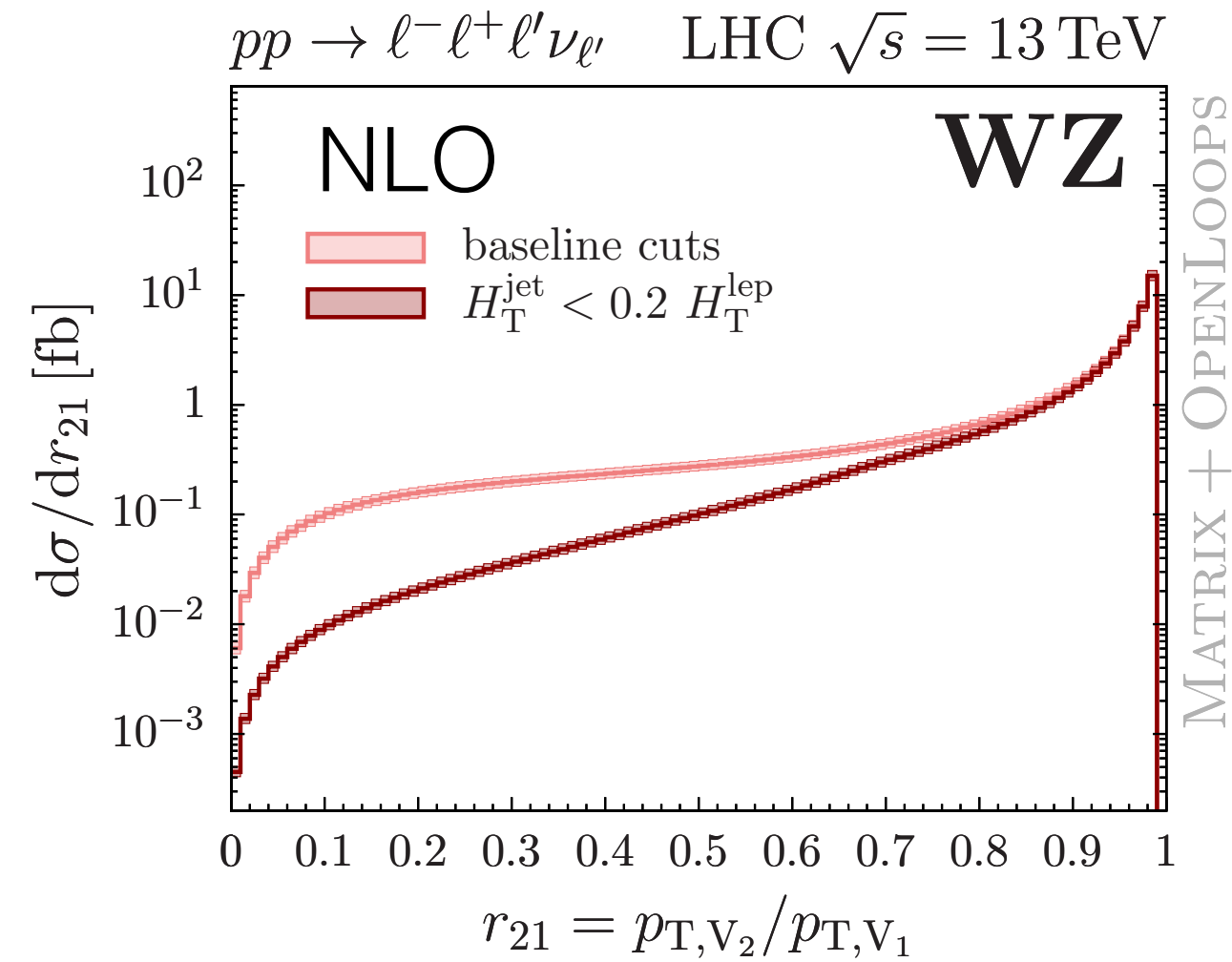
- cancellations of EW corr. in $q\bar{q}$ and $q\gamma$ channels highly observable dependent

Conclusions

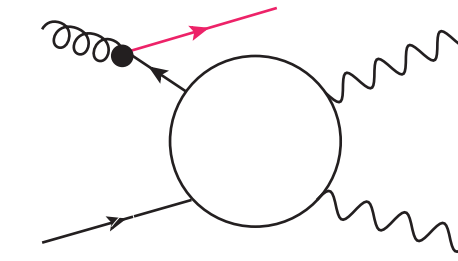
- Multilepton signatures dominated by $VV / VVV / ttV$ SM backgrounds
- Upcoming large data samples will require more careful background estimates
- It will become mandatory to include higher-order **QCD** & **EW**
- State-of-the-art for VV at fixed-order: NNLO **QCD** \times NLO **EW**
 - QCD \rightarrow giant K-factors, but NNLO often allows to reach few percent precision
 - EW \rightarrow large EW Sudakov corrections at large energies: several tens of percent
- Beyond fixed-order for VV :
 - NNLO QCD + PS
 - MEPS @ NLO QCD \times EW
 - NLO QCD + EW PS
- State-of-the-art for VVV at fixed-order: NLO **QCD** \times NLO **EW**
 - off-shell calculations mandatory
 - Very large cancellations between EW Sudakov and photon-induced corrections

Backup

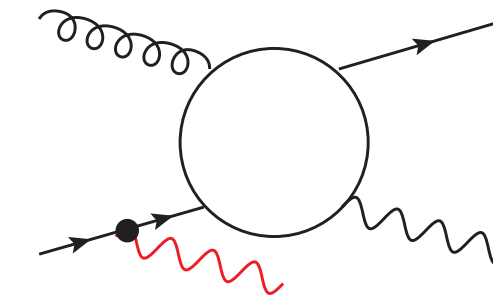
Giant K-factors and effect of jet veto



- at $r_{21} \rightarrow 1$: hard-VV topologies



- at $r_{21} \rightarrow 0$: hard-Vj topologies



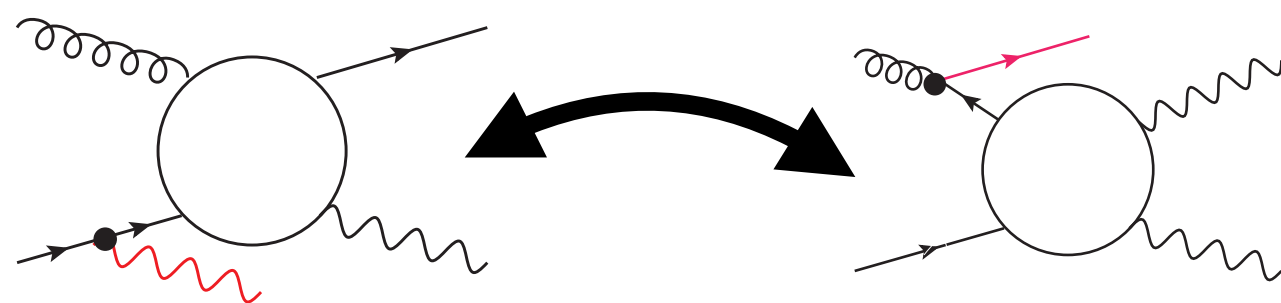
- for $p_{T,V_1} > 1$ TeV: hard-Vj topologies dominate over hard-VV

- Jet veto $H_T^{\text{jet}} < \xi_{\text{veto}} H_T^{\text{lep}}$ corresponds to

$$p_{T,V_2} \geq \frac{1 - \xi_{\text{veto}}}{1 + \xi_{\text{veto}}} p_{T,V_1} = \frac{2}{3} p_{T,V_1} \quad \text{for } \xi_{\text{veto}} = 0.2$$

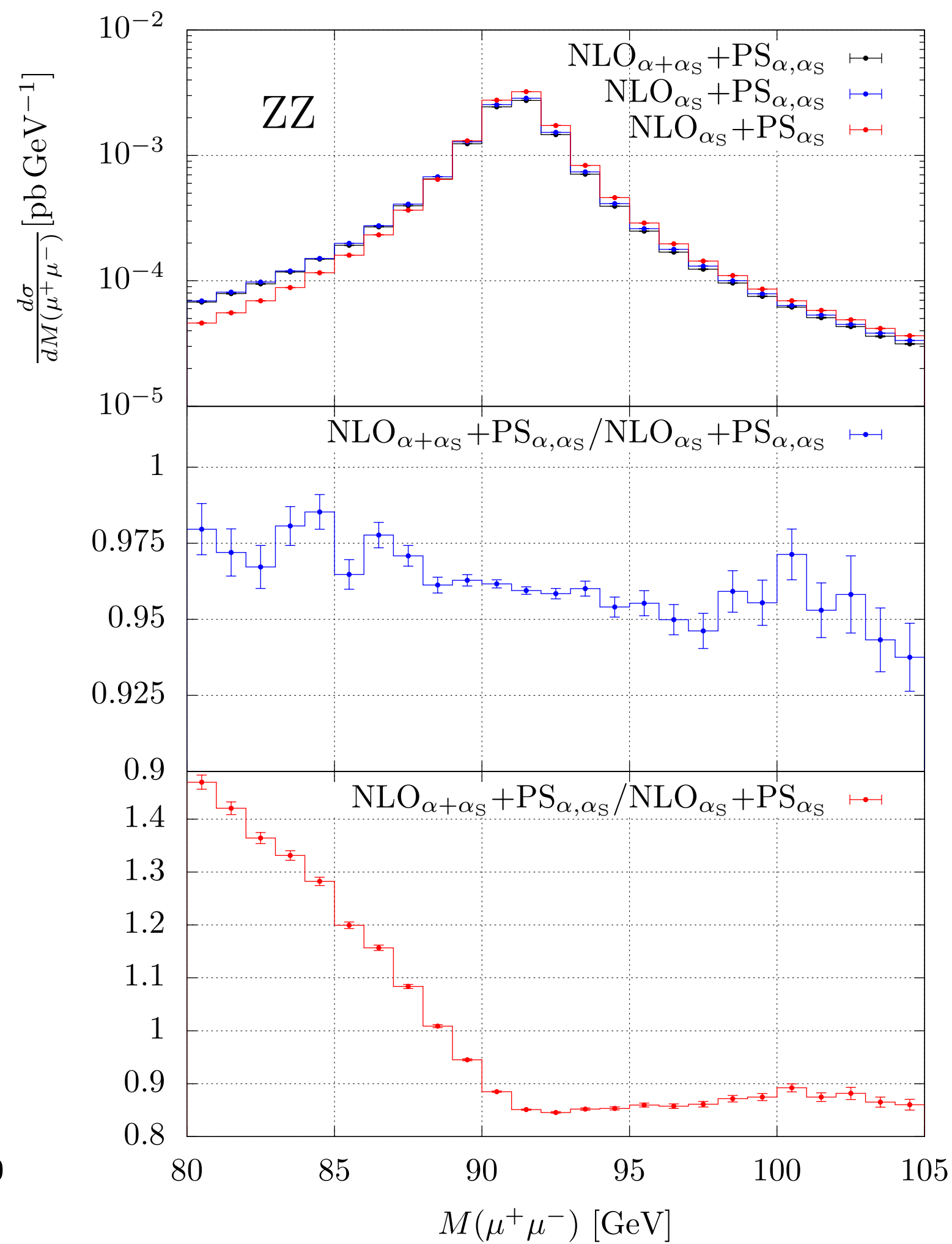
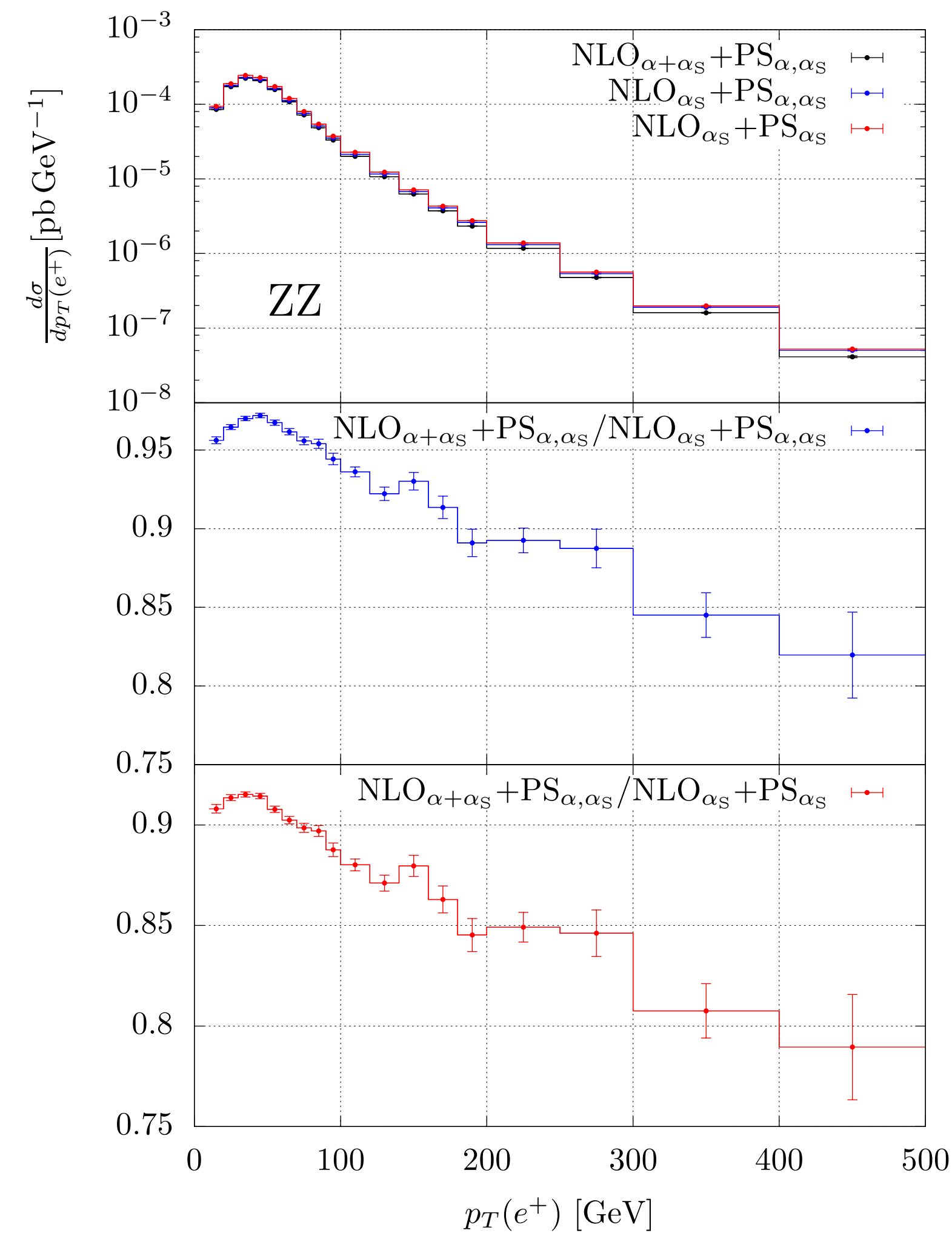
(violated by off-shell topologies)

- Jet veto results in phase-space dominated by hard-VV



PS MC: NLO QCD + EW PS

[Chiesa, Re, Oleari '20]



Available in POWHEG-BOX-RES
(Resonance aware matching)

NLO (QCD + EW) PS (QCD + QED) /
NLO QCD PS (QCD + QED)

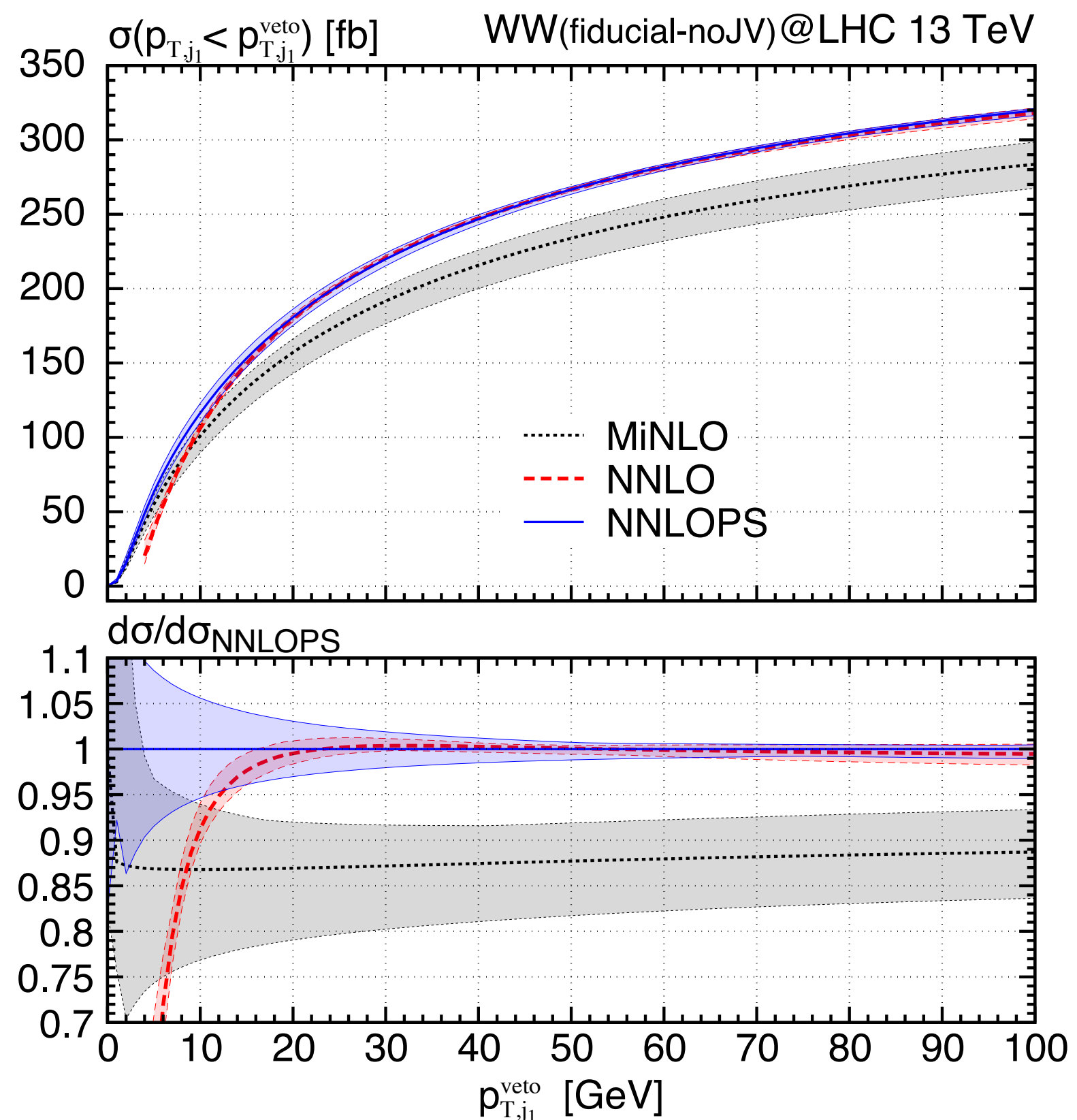
NLO (QCD + EW) PS (QCD + QED) /
NLO QCD PS QCD

- Missing: photon-induced channels
- Question: NLO (QCD + EW) PS (QCD + QED) / (NLO QCD PS QCD) \times NLO EW ?

Parton shower Monte Carlos: NNLO QCD + PS via Mi(N)NLO_{PS}

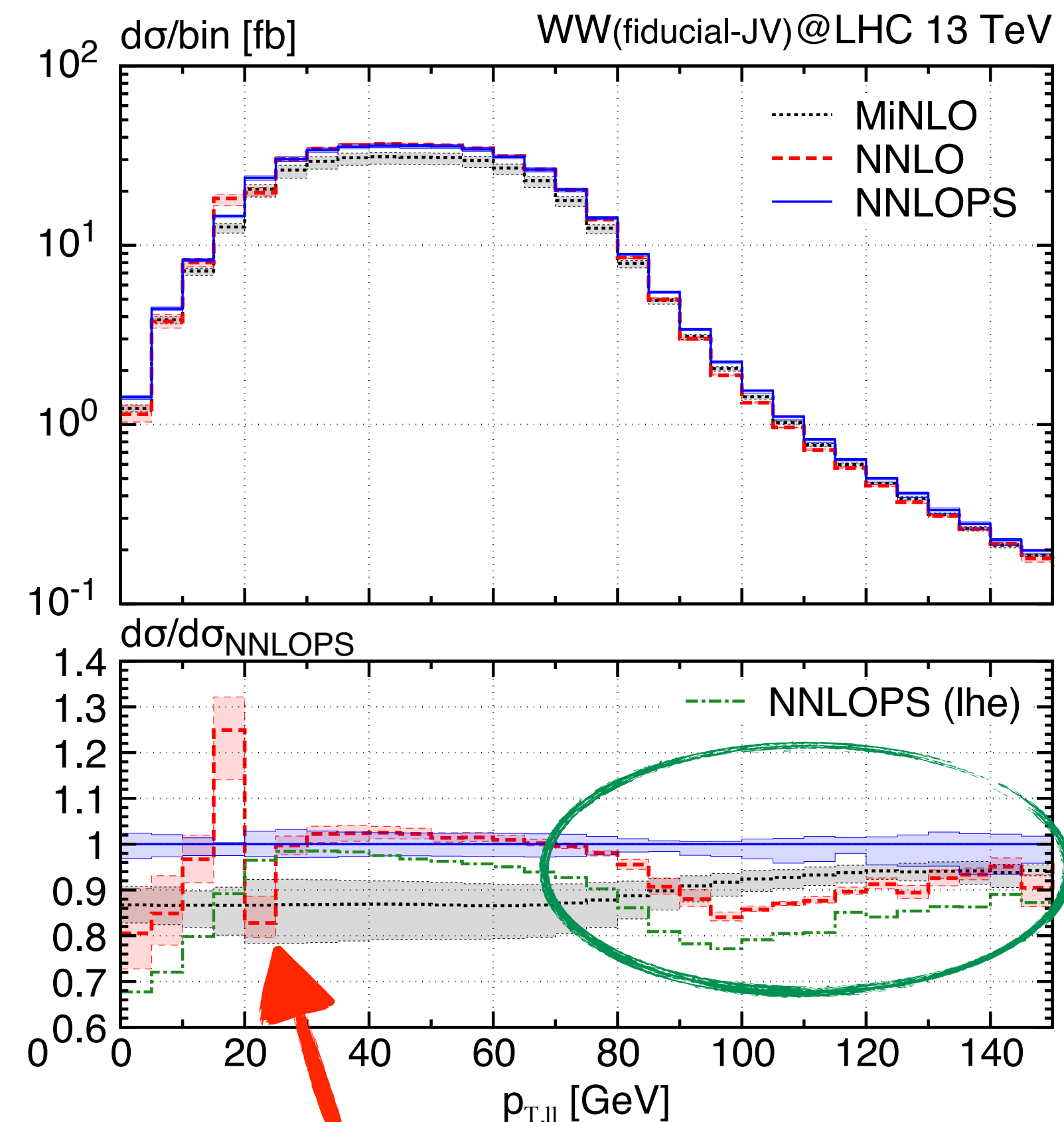
[Re, Wiesemann; Zanderighi '18]

Jet veto



→ **NNLOPS physical down to $p_T = 0$**

p_T of dilepton system



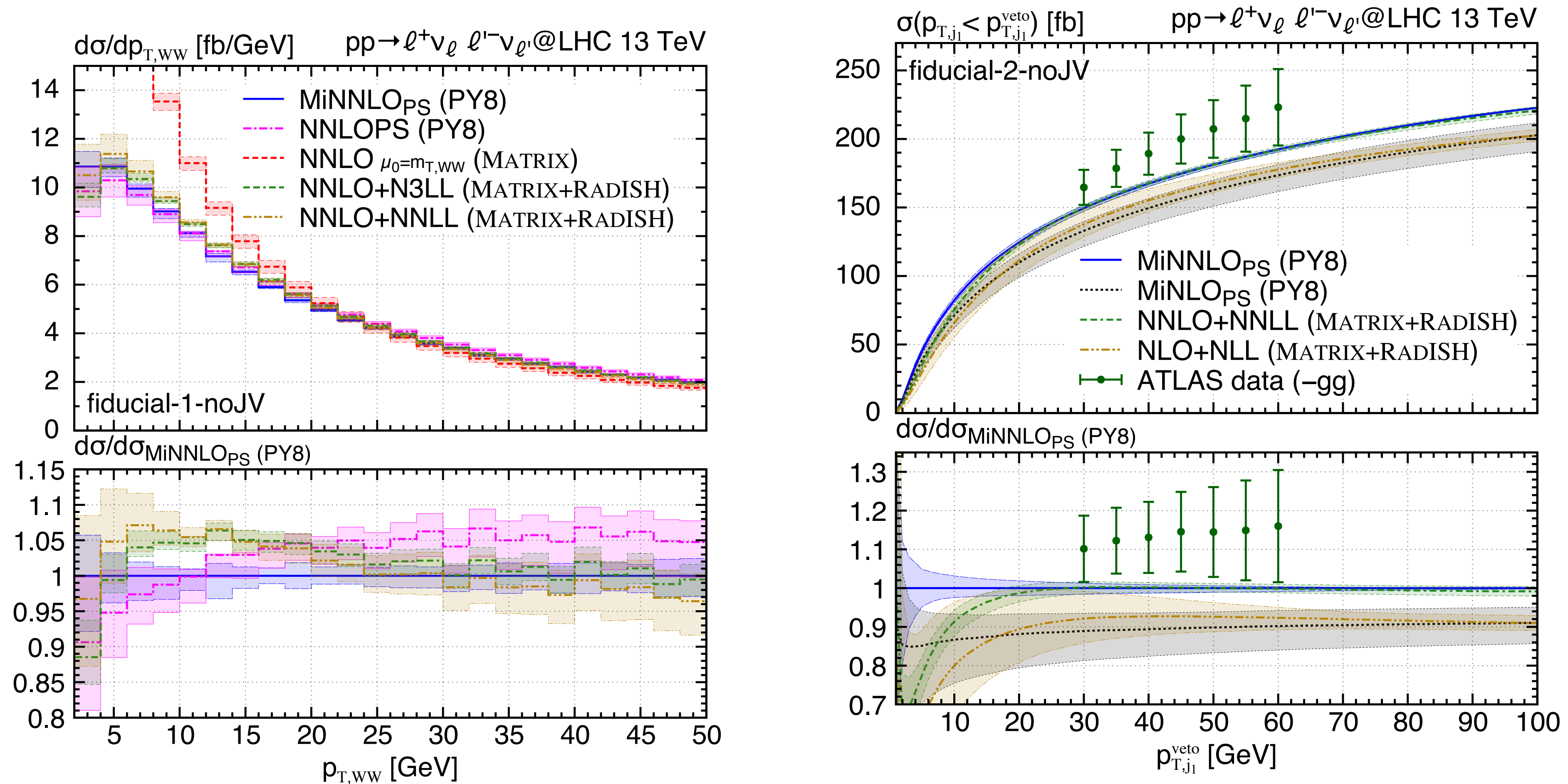
→ **NNLOPS cures perturbative instabilities (p_T^{miss} cut)**

→ **NNLOPS induces additional shape effects**

- Currently only available for WW via MiNLO_{PS} + reweighting
- Full MiNNLO_{PS} work in progress (recent results for $Z\gamma$ - [2010.10478](https://arxiv.org/abs/2010.10478))

PS MC: NNLO QCD + PS for VV via MiNNLO_{PS}

[Lombardi, Wiesemann; Zanderighi '21]



- MiNNLO_{PS} physical down to $p_{T,WW}=0$
- Latest implementation does not require computationally expensive reweighting required earlier
- Also available for Z γ : - [2010.10478](https://arxiv.org/abs/2010.10478) [Lombardi, Wiesemann; Zanderighi '20]
- Alternative NNLOPS approach available for ZZ in GENEVA [Alioli, Broggio, Gavardi, Kallweit, Lim, Nagar, Napoletano '21]

Triboson production calculations

NLO QCD

- on-shell fixed-order [Binoth et.al. arXiv:0804.0350](#)
off-shell fixed-order [Campanario et.al. arXiv:0809.0790](#)
→ available in aut. tools (MG5_aMC, SHERPA+OPENLOOPS/RECOLA)
- on-shell matched to parton showers in SHERPA,
multijet merged $WWW + 0, 1j$ [Höche et.al. arXiv:1403.7516](#)
- off-shell matched to parton showers should be available in
automated tools (MG5_aMC, SHERPA+OPENLOOPS/RECOLA)

NLO EW

- on-shell known for some time [Yong-Bai et.al. arXiv:1605.00554](#)
[Dittmaier, Huss, Knippen arXiv:1705.03722](#)
- off-shell recently computed [MS arXiv:1806.00307](#)
- no matching to parton showers available yet