

Two-loop five-point amplitudes in massless QCD with finite fields

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Outline

Introduction

Processes

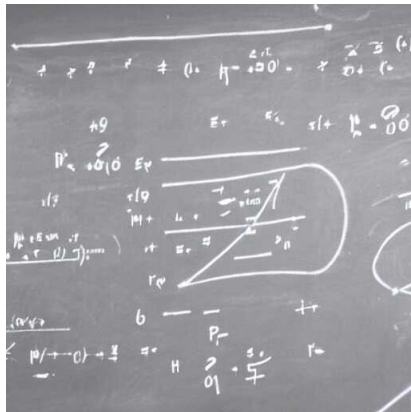
Computation

Finite fields

Reconstruction

Performance

Conclusion



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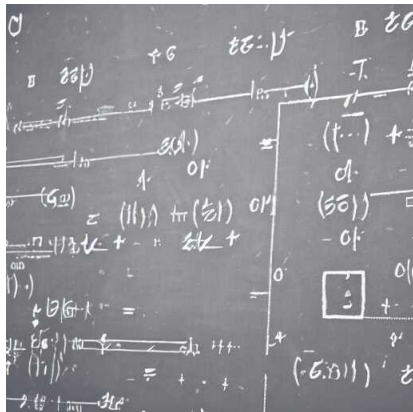
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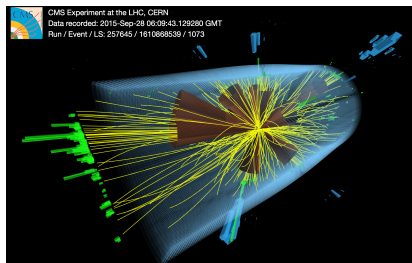


Precision frontier

- Better understand properties of SM
- Indirect probe of new physics by small deviations
- LHC demanding increasing precision
- Theory predictions at 1%
- Requires \geq NNLO QCD

$$\alpha_s \approx 0.1$$

$$d\sigma = d\sigma_{\text{LO}} + \alpha_s d\sigma_{\text{NLO}} + \alpha_s^2 d\sigma_{\text{NNLO}} + \dots$$



[McCauley et al. 2015]

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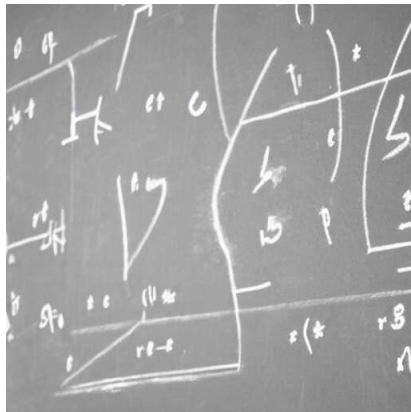
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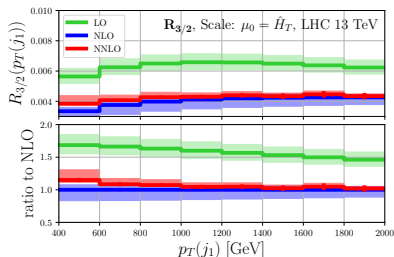
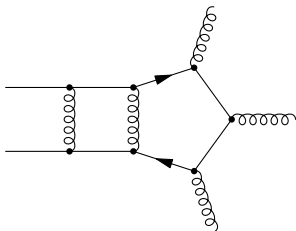
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Trijet production

- NNLO leading-colour $pp \rightarrow jjj$

[Badger, Brønnum-Hansen, Hartanto, RM, Peraro, Zoia n.d.]



[Czakon, Mitov, Poncelet 2021]

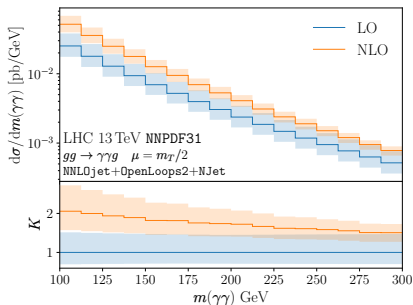
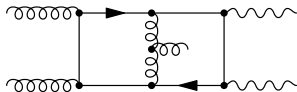
- Probe fundamental parameters of SM, $R_{3/2} \rightarrow \alpha_s$
- Significant theory uncertainty reduction

[Abreu, Febres Cordero, Ita, Page, Sotnikov 2021] [Czakon, Mitov, Poncelet 2021]

Diphoton-plus-jet production via gluon fusion

- NLO full-colour $gg \rightarrow g\gamma\gamma$ (N³LO $pp \rightarrow g\gamma\gamma$)

[Badger, Brønnum-Hansen, Chicherin, Gehrmann, Hartanto, Henn, Marcoli, RM, Peraro, Zoia 2021]



[Badger, Gehrmann, Marcoli, RM 2022]

- Important background for measuring Higgs properties
- Significant NLO corrections, enhances NNLO $pp \rightarrow g\gamma\gamma$

[Chawdhry, Czakon, Mitov, Poncelet 2021]

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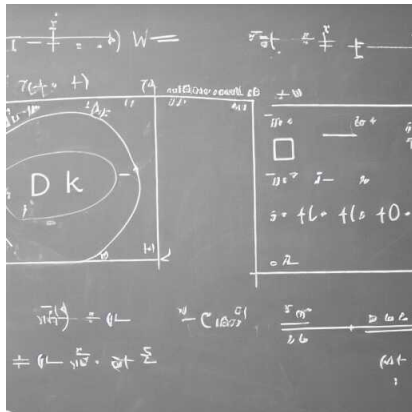
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Five-point two-loop computation

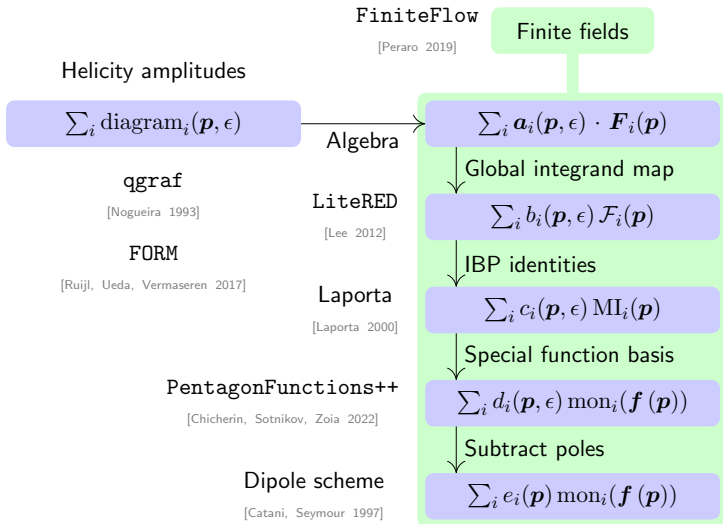
- Major theoretical challenge
 - New methods
- Reconstruct over finite fields

[Peraro 2019]
- Pentagon function basis

[Gehrmann, Henn, Lo Presti 2018]
[Chicherin, Sotnikov 2020]
[Chicherin, Sotnikov, Zoia 2022]
- Fast and stable C++ implementation in NJet

[Badger, Biedermann, RM, Uwer, Yundin 2021]

Method overview



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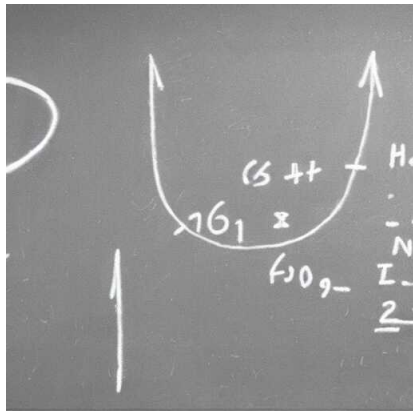
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Finite field arithmetic

- Set of $n \in \mathbb{P}^p$ non-negative integers

$$\mathbb{F}_n = \{0, \dots, n-1\}$$

- Arithmetic operations modulo n

$$+ \quad - \quad \times \quad \div$$

- Modular multiplicative inverse $x = a^{-1} \pmod{n}$

$$ax = 1 \pmod{n} \qquad a \neq 0$$

- One-to-many map $\mathbb{Q} \rightarrow \mathbb{F}_n$

$$\frac{a}{b} \rightarrow ab^{-1} \pmod{n}$$

Numeric representation for computation

`float` Catastrophic cancellation

`int` Overflow

\mathbb{Q} Slow

- \mathbb{F}_n
 - No precision loss (n large)
 - Fast: fixed size integer
 - Recover $\mathbb{F}_n \rightarrow \mathbb{Q}$
 - Chinese Remainder Theorem: $\{\mathbb{F}_n\} \rightarrow \mathbb{F}_{\prod n}$

- Amplitudes

- Large intermediate expressions
- Bypass with numerical evaluation over \mathbb{F}_n
- Reconstruct analytic expression

Rational on-shell parametrisation

- Momentum twistor variables

$$\langle ij \rangle, [ij] \rightarrow x_i$$

$$Z = \begin{pmatrix} \lambda_1 & \cdots & \lambda_n \\ \mu_1(\tilde{\lambda}) & \cdots & \mu_n(\tilde{\lambda}) \end{pmatrix}$$

- Rational functions
 - $p_i(\mathbf{x}), s_{ij}(\mathbf{x}), \text{tr}_5(\mathbf{x}), \dots$
- \mathbf{x} unconstrained
 - On-shell
 - Momentum conserving

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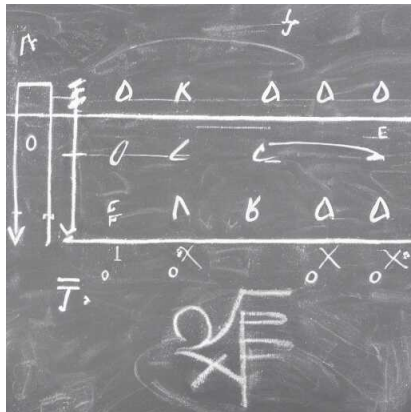
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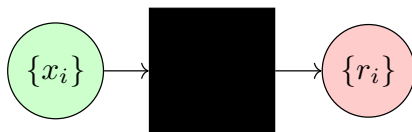


Reconstruction

- Amplitude components

$$F(\mathbf{x}) = \sum_i r_i(\mathbf{x}) \text{mon}_i(f)$$

- Have numerical algorithm for r_i
- Reconstruct expression from sufficient evaluations
- `FiniteFlow` [Peraro 2019]



- Strategies to optimise and compactify

Linear relations in the coefficients

- Linearise the r_i

$$\{r_i\}_{i \in S} \rightarrow \{r_i\}_{i \in T \subseteq S}$$

- Numerically solve:

$$\sum_i a_i r_i(\mathbf{x}) = 0$$

and choose lowest degrees

Matching factors

- Coefficient ansatz

$$r(\mathbf{x}) = \frac{n(\mathbf{x})}{\prod_k \ell_k^{e_k(\mathbf{x})}} \quad \ell_k \in \text{pentagon alphabet}$$

- Determine e_k by reconstructing r on univariate slice

$$\mathbf{x} = \mathbf{c}_0 + \mathbf{c}_1 t \quad \rightsquigarrow \quad r(t)$$

and matching RHS

- Fix denominators

$$\{r_i\} \rightarrow \{n_i\}$$

[Abreu, Dormans, Febres Cordero, Ita, Page 2019]

Univariate partial fraction decomposition

- Example in y

$$\frac{n(x, y)}{x^2 y^2 (x^2 + y^2)} = q_0(x) y^{d_n - d_d} + \frac{q_1(x)}{y} + \frac{q_2(x) + q_3(x)y}{y^2} + \frac{q_4(x) + q_5(x)y}{x^2 + y^2}$$

- Only need to know y degree of n , d_n
- Choose y by studying one-loop
- Reconstruct r_i directly in decomposed form

$$\{r_i(\bar{\mathbf{x}}, y)\} \rightarrow \{q_i(\bar{\mathbf{x}})\}$$

- Reduce variables by one
- Lower degrees

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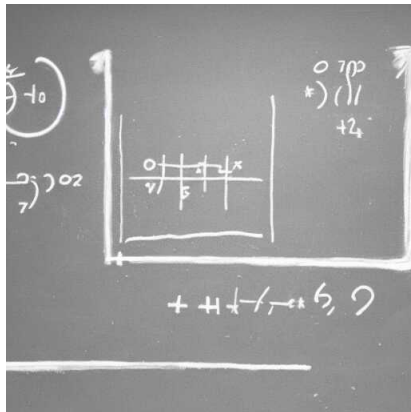
Computation

Finite fields

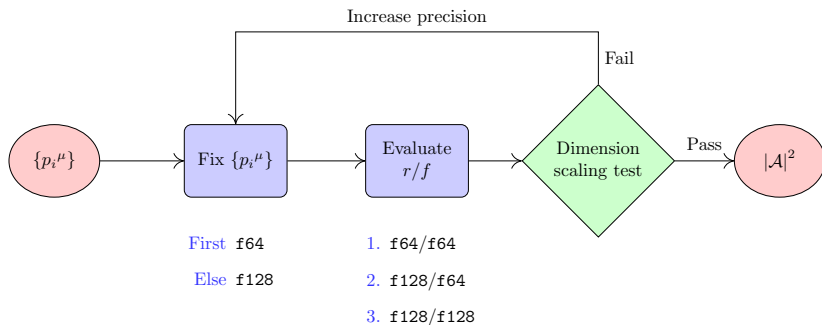
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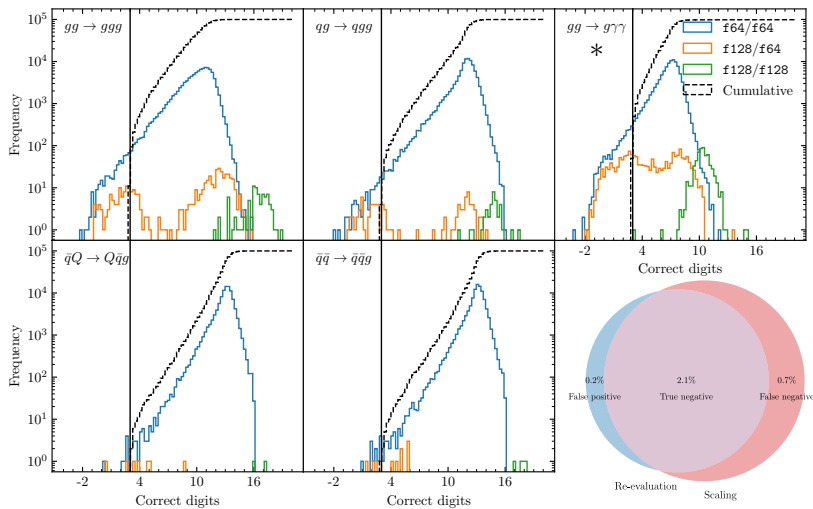
Evaluation strategy



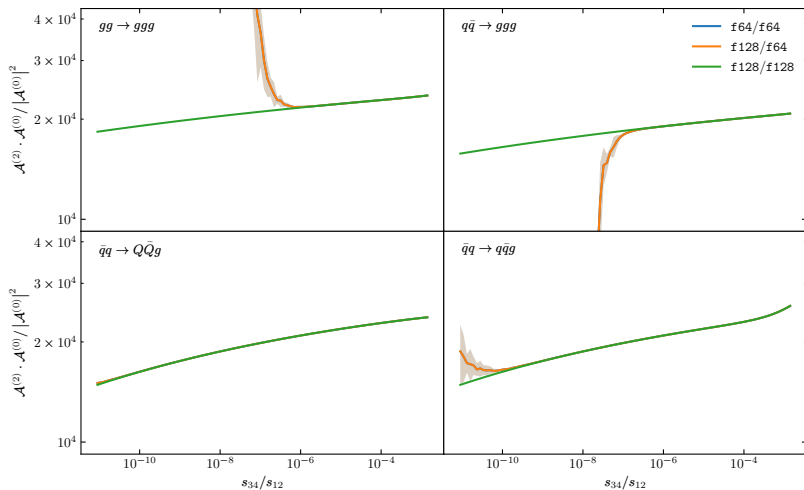
Timing

Channel	f64/f64		Evaluation strategy	
	Time (s)	f (%)	Time (s)	f (%)
$gg \rightarrow ggg$	1.39	69	1.89	77
$gg \rightarrow \bar{q}qg$	1.35	91	1.37	91
$qg \rightarrow qgg$	1.34	92	1.57	93
$q\bar{q} \rightarrow ggg$	1.34	93	1.38	93
$\bar{q}Q \rightarrow Q\bar{q}g$	1.14	99	1.16	99
$\bar{q}\bar{Q} \rightarrow \bar{q}\bar{Q}g$	1.36	99	1.39	99
$\bar{q}g \rightarrow \bar{q}Q\bar{Q}$	1.36	99	1.39	99
$\bar{q}q \rightarrow Q\bar{Q}g$	1.14	99	1.14	99
$\bar{q}g \rightarrow \bar{q}q\bar{q}$	1.84	99	1.90	99
$\bar{q}\bar{q} \rightarrow \bar{q}\bar{q}g$	1.82	99	1.94	99
$\bar{q}q \rightarrow q\bar{q}g$	1.71	99	1.77	99
$gg \rightarrow \gamma\gamma g$ *	9	99	26	99

Stability

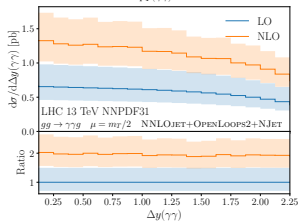
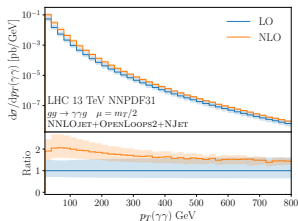


IR performance



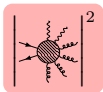
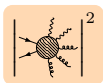
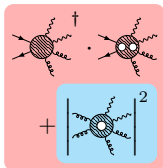
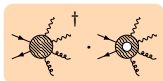
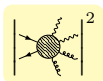
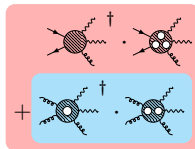
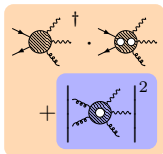
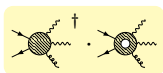
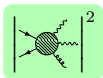
Conclusion

- Huge success of technical advancements in precision QCD
 - Numerical evaluation over finite fields
 - Bypass intermediate complexity
 - Reconstruct compact analytic forms
- Efficient and stable evaluation in NJet
 - $pp \rightarrow jjj$ (LC)
 - $gg \rightarrow g\gamma\gamma$
- Driving pheno predictions towards 1% precision



[Badger, Gehrmann, Marcoli, RM 2022]

$$pp \rightarrow g\gamma\gamma$$



LO
NLO
NNLO
N³LO

Loop induced:
 LO
NLO

Implementation

- Finite remainders coded up in C++ as analytic expressions
- Construct partial amplitudes as

$$F^h = r_i^h(\boldsymbol{x}) M_{ij}^h f_j^h$$

f_j^h	special function monomials	global
M_{ij}^h	rational sparse matrices	partials
r_i^h	independent rational coefficients	helicities
\boldsymbol{x}	momentum twistor variables	global

- Independent helicities permuted to all mostly-plus
 - Mostly-minus: $r_i^* M_{ij} P(f_j)$

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