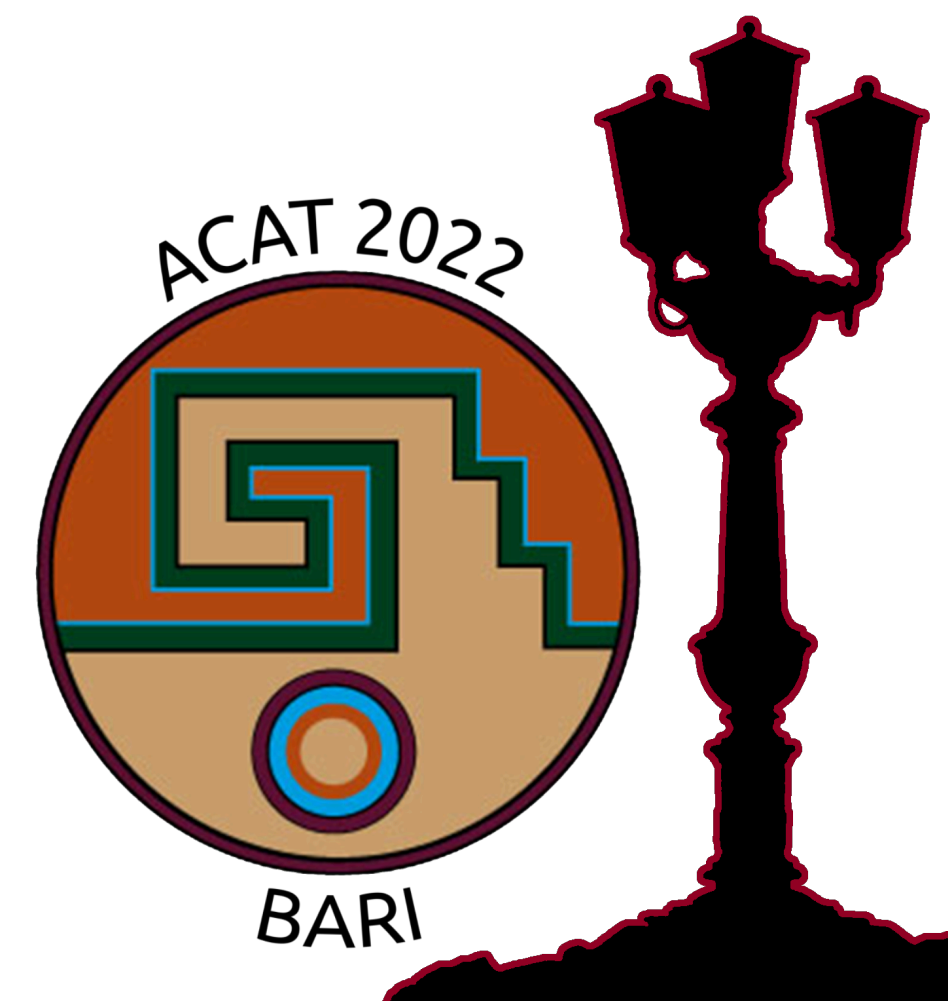


Track 3 Summary

Computations in Theoretical Physics: Techniques and Methods

Anke Biekötter - Leonardo Cosmai -
Joshua Davies - Latifa Elouadrhiri

ACAT 2022 -
24-28 October 2022
Villa Romanazzi Carducci, Bari, Italy



ACAT 2022



A promotional poster for Eurospin. The top section is a red banner with white text: "ALTROCONSUMO APPROVATO", "Eurospin approvato da Altroconsumo come discount salvaprezzo in Italia, per la categoria prodotti più economici, classifica unica iper, super e discount.", "Pubblicato il 09/2022", "PRODOTTI PIÙ ECONOMICI", "Bilievazione prezzi svolta dal 01/03/22 al 01/04/22 in 107 punti vendita, in 67 città, su 126 categorie di prodotti alimentari confezionati e freschi, cura casa e persona, pet food.", "DISCOUNT SALVAPREZZO IN ITALIA". Below the banner is a photograph of a man with white hair and a mustache, wearing a red sweater, standing with his arms crossed in a supermarket aisle. A shopping cart filled with groceries is in the foreground. The text "E=mc²" is overlaid on the image, with "Eurospin" and "Massima Convenienza" below it. At the bottom, a blue banner reads: "Prodotti di qualità alla massima convenienza tutti i giorni: questa è la Spesa intelligente." and the Eurospin logo.

ACAT 2022

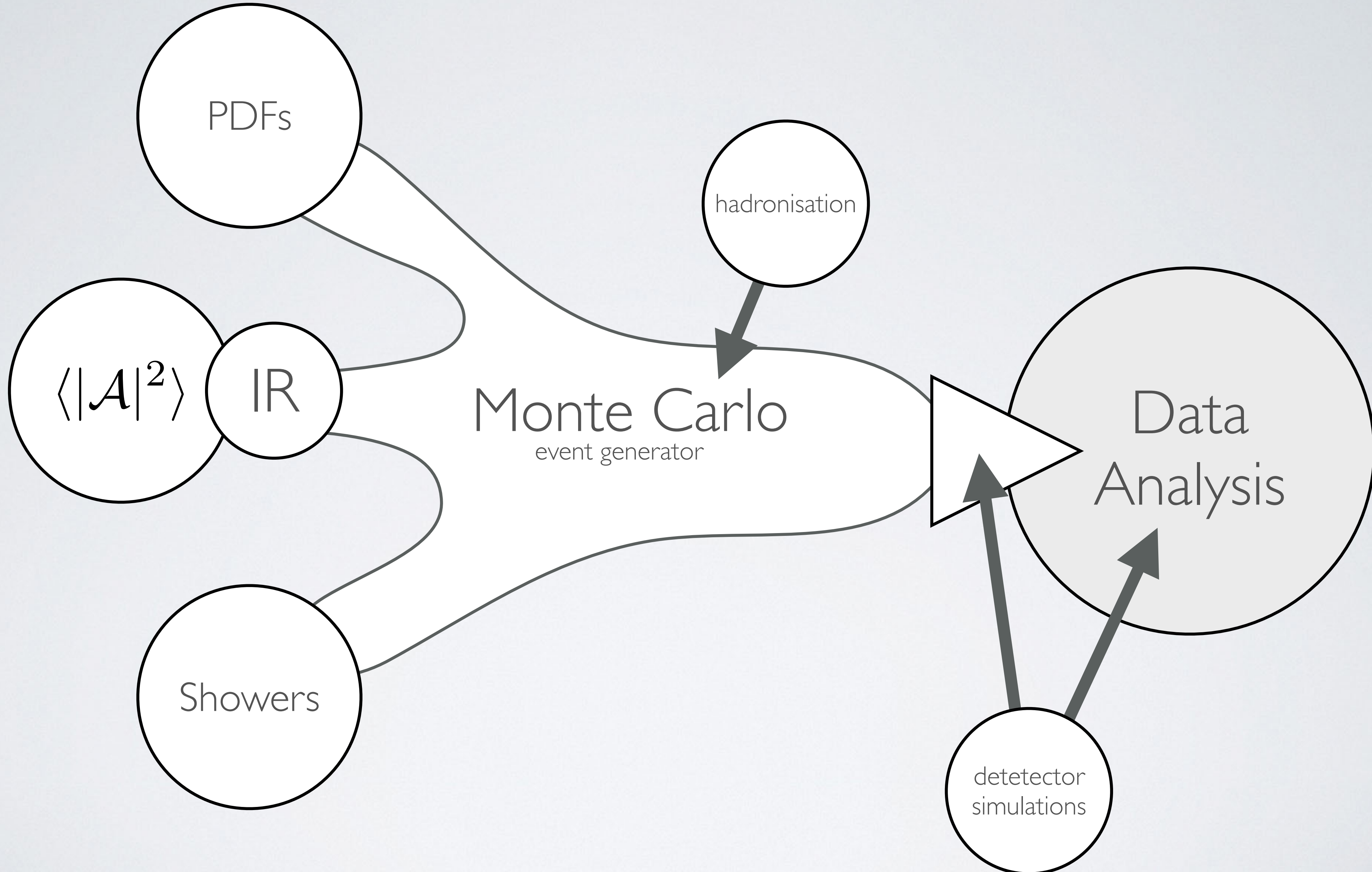


Thank you for your contributions!

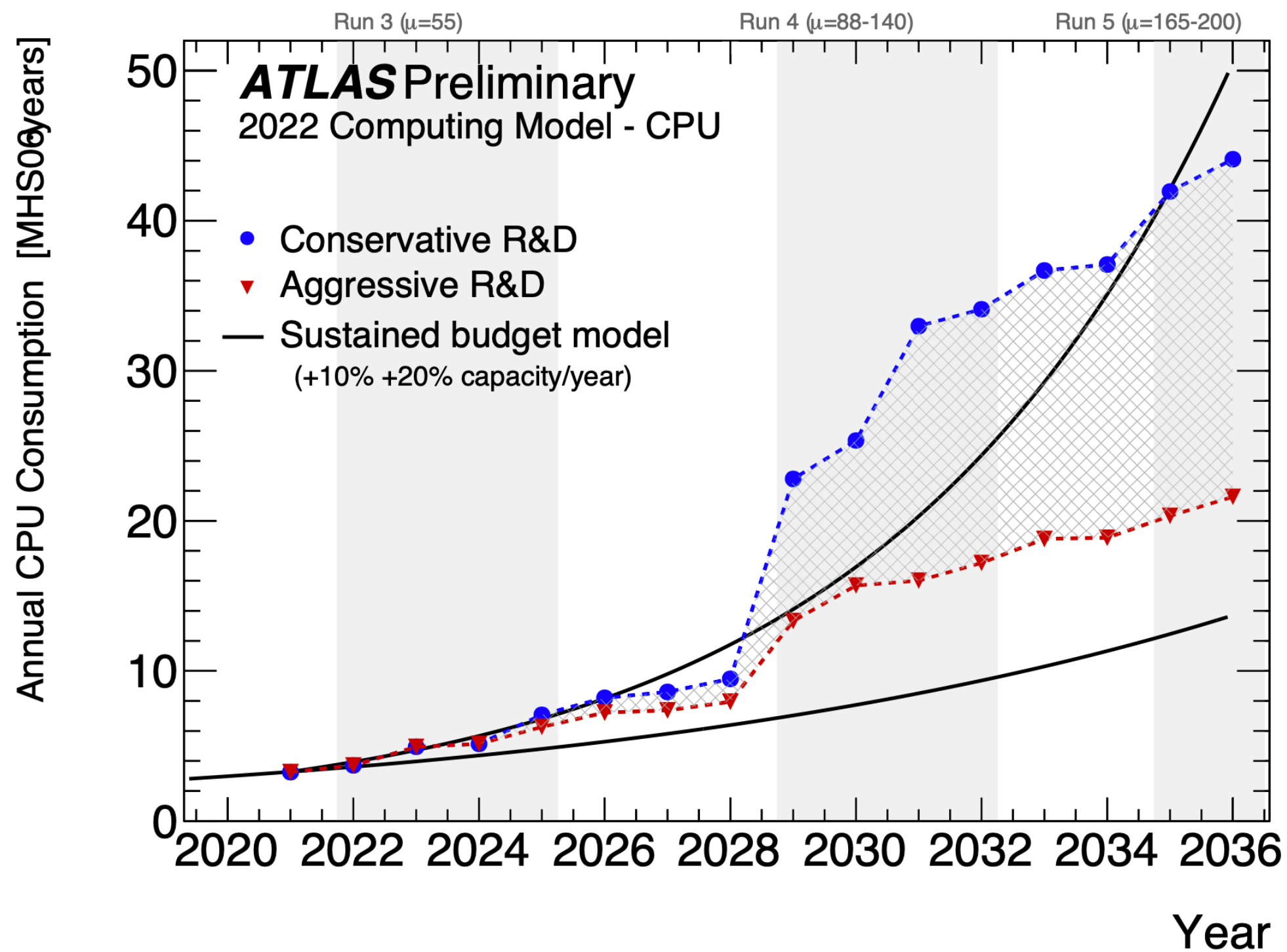
Track 3 Highlights

- Monte Carlo generation
- Precision frontier
- Beyond Standard Model physics
- Towards Quantum Computing

This is a biased selection



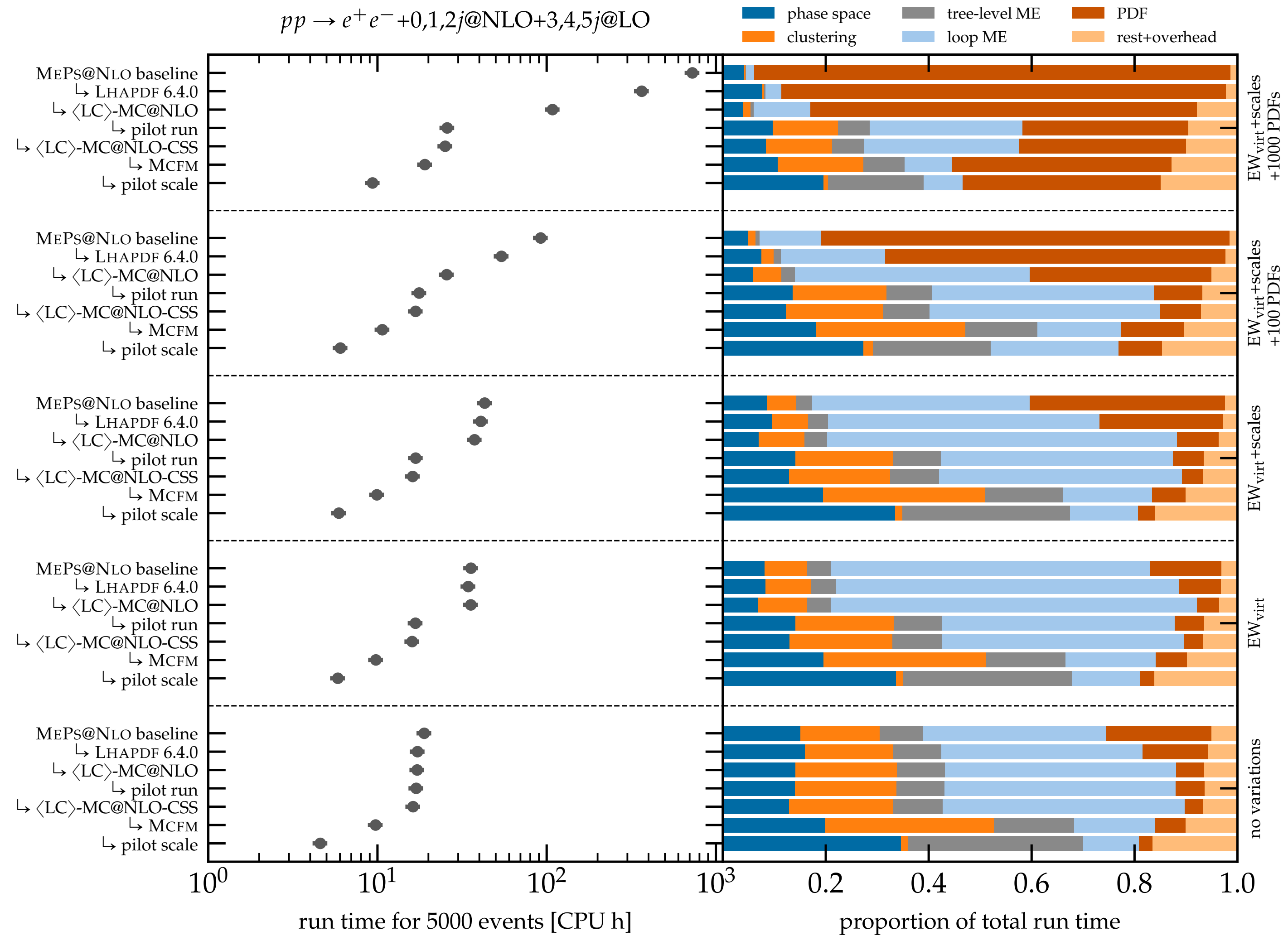
Speeding up Monte Carlo event generators



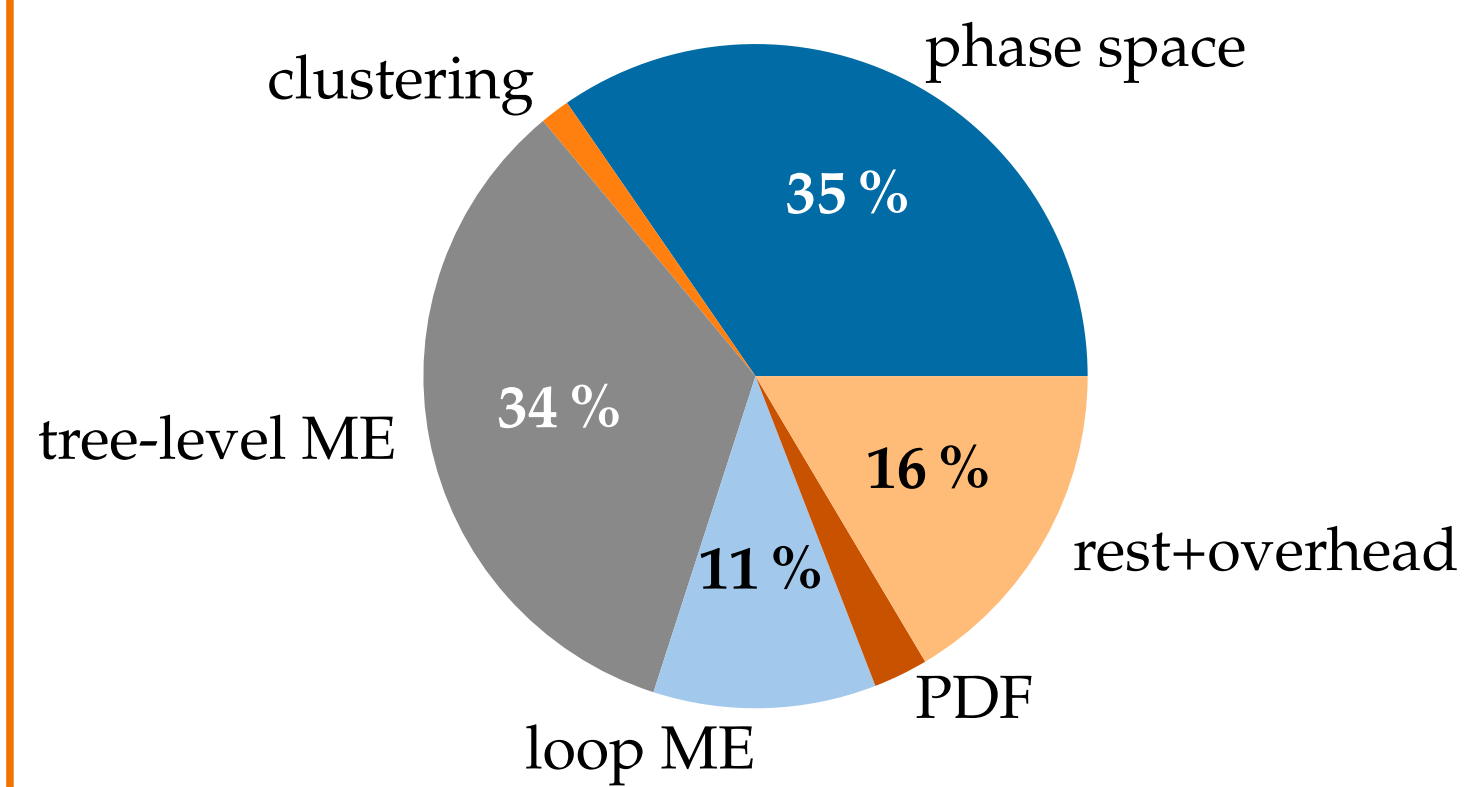
- Performance analysis
- Pilot runs (what do we need when?)
- New architectures - GPUs, vector CPUs
- Portability (Kokkos, Alpaka, ...)
- Physics ideas and analytic results

Breakdown of CPU budget in $V+jets$

Christian Gütschow

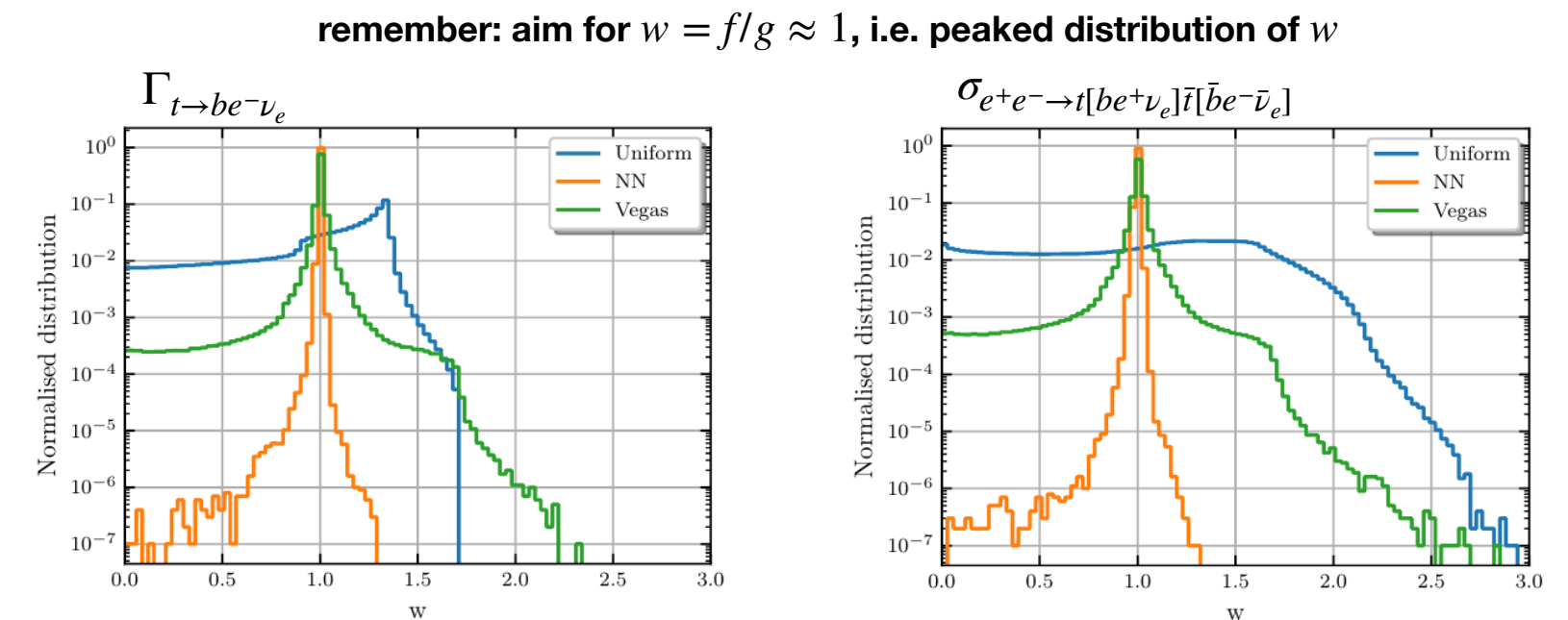


$$pp \rightarrow e^+e^- + 0,1,2j @ \text{NLO} + 3,4,5j @ \text{LO}$$



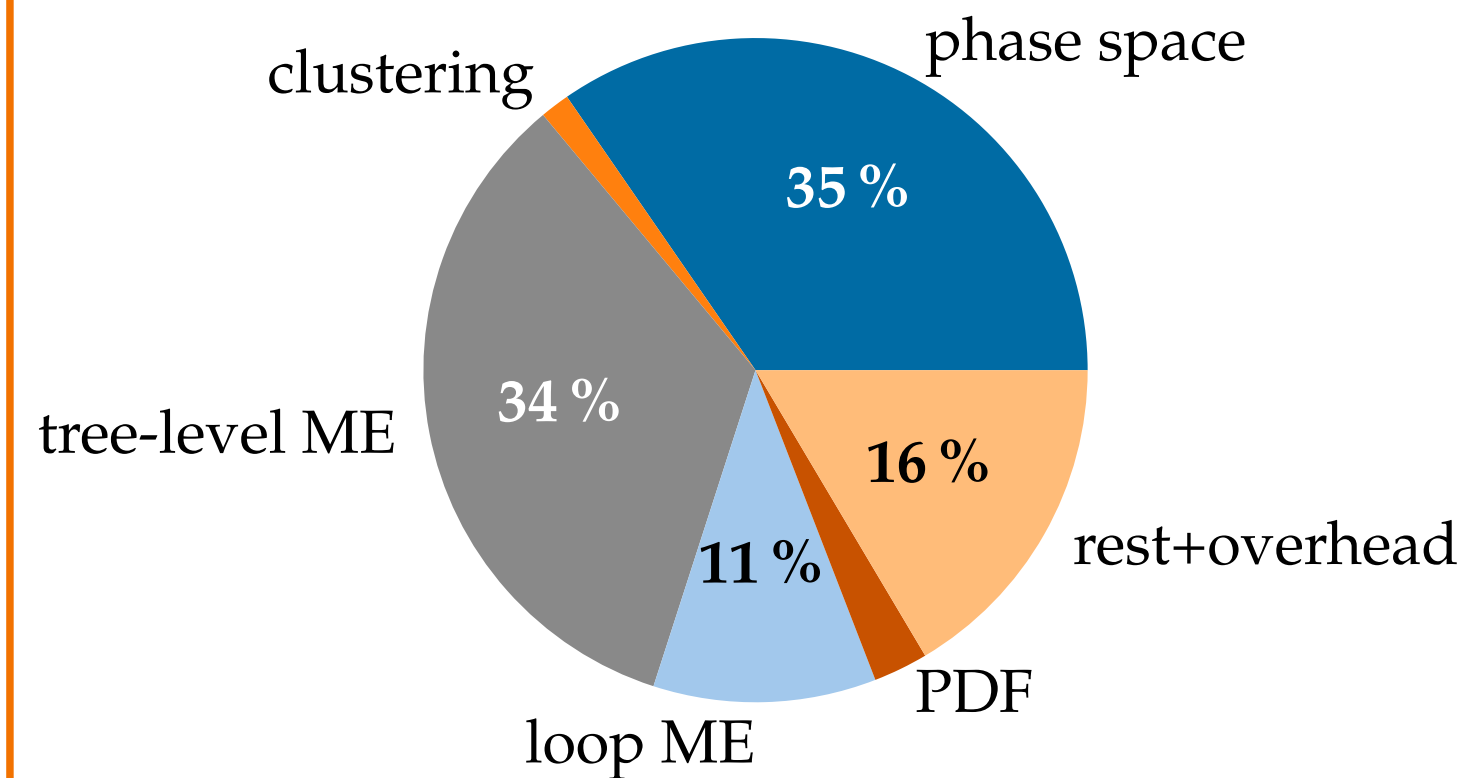
Neural Importance Sampling – Results

Enrico
Bothmann



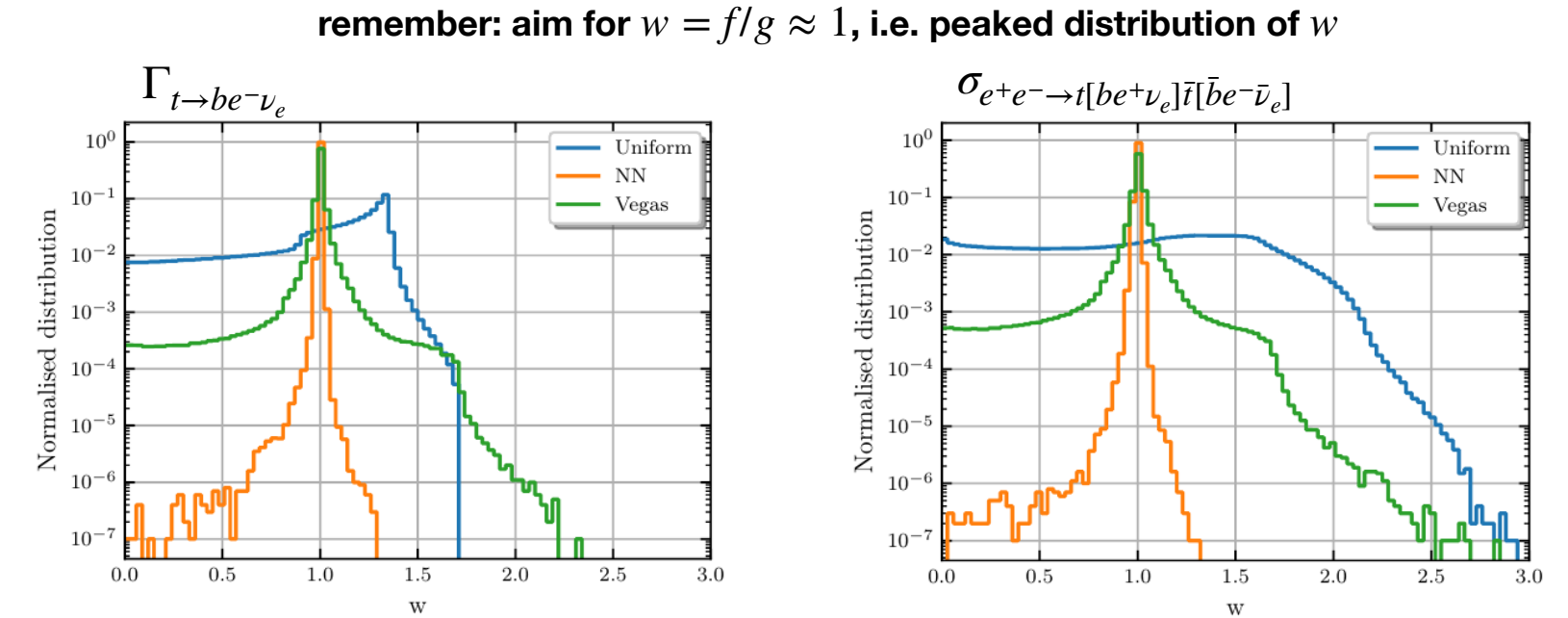
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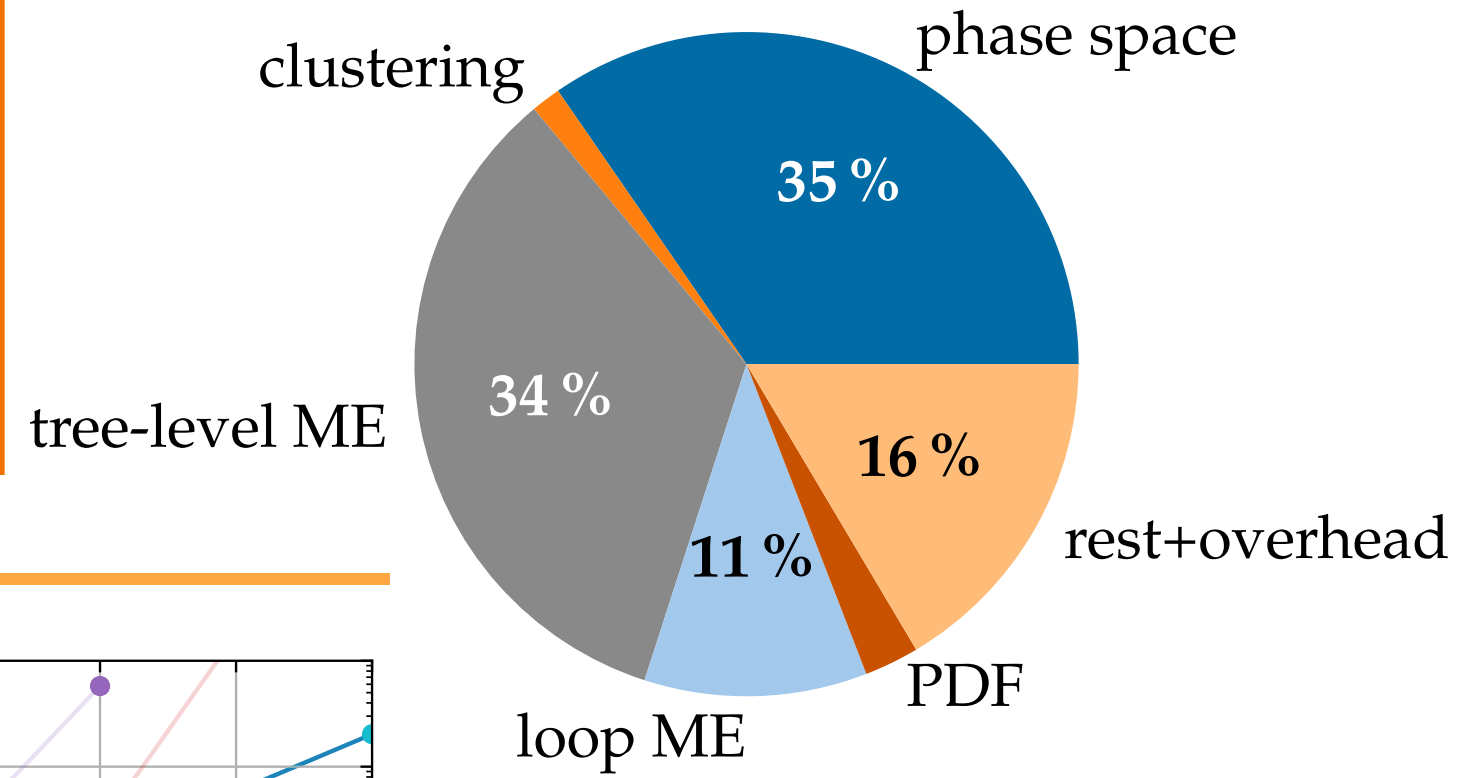
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Enrico
Bothmann



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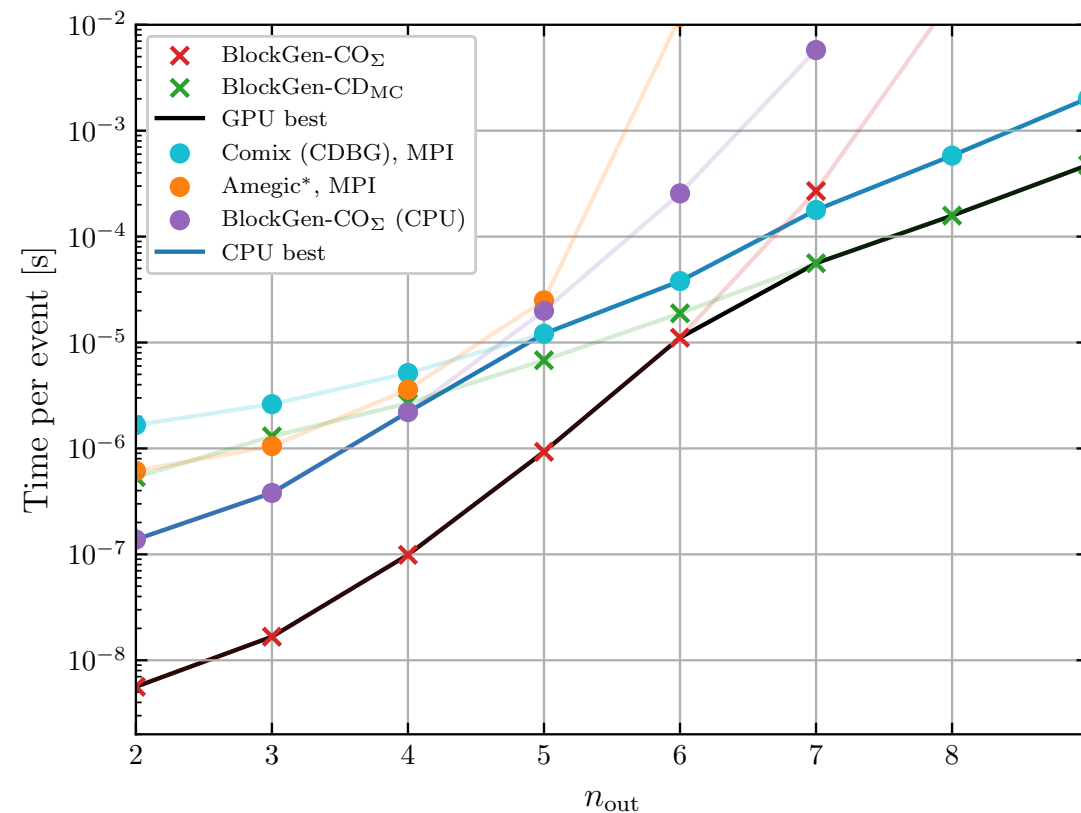
10

Max Knobbe

3rd Component: The Color Sum [Bothmann, Giele, Höche, Isaacson, MK, 2106.06507]

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- Caveat: Color-sampling comes with penalty factor from slower convergence

⇒ Algorithmic choice: Sum colors



MadEvent/CUDA for $gg \rightarrow t\bar{t}gg$ (improved at ACAT2022)

CUDA grid size		ICHEP2022			madevent		standalone	
$gg \rightarrow t\bar{t}gg$	MES precision	$t_{TOT} = t_{Mad} + t_{MEs}$ [sec]	N_{events}/t_{TOT} [events/sec]	N_{events}/t_{MEs} [MEs/sec]	8192	524288		
Fortran	double	58.3 = 5.2 + 53.1	1.55E3 (=1.0)	1.70E3 (=1.0)				
CUDA	double	6.1 = 5.7 + 0.36	1.49E4 (x9.6)	2.54E5 (x149)	2.51E5	4.20E5 (x247)		
CUDA	float	5.7 = 5.4 + 0.24	1.59E4 (x10.3)	3.82E5 (x224)	3.98E5	8.75E5 (x515)		

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CUDA	float	2.8 = 2.6 + 0.24	3.24E4 (x19.9)	3.83E5 (x225)	3.96E5	8.77E5 (x)		

Speeding up Madgraph5_aMC@NLO through CPU vectorization and GPUs

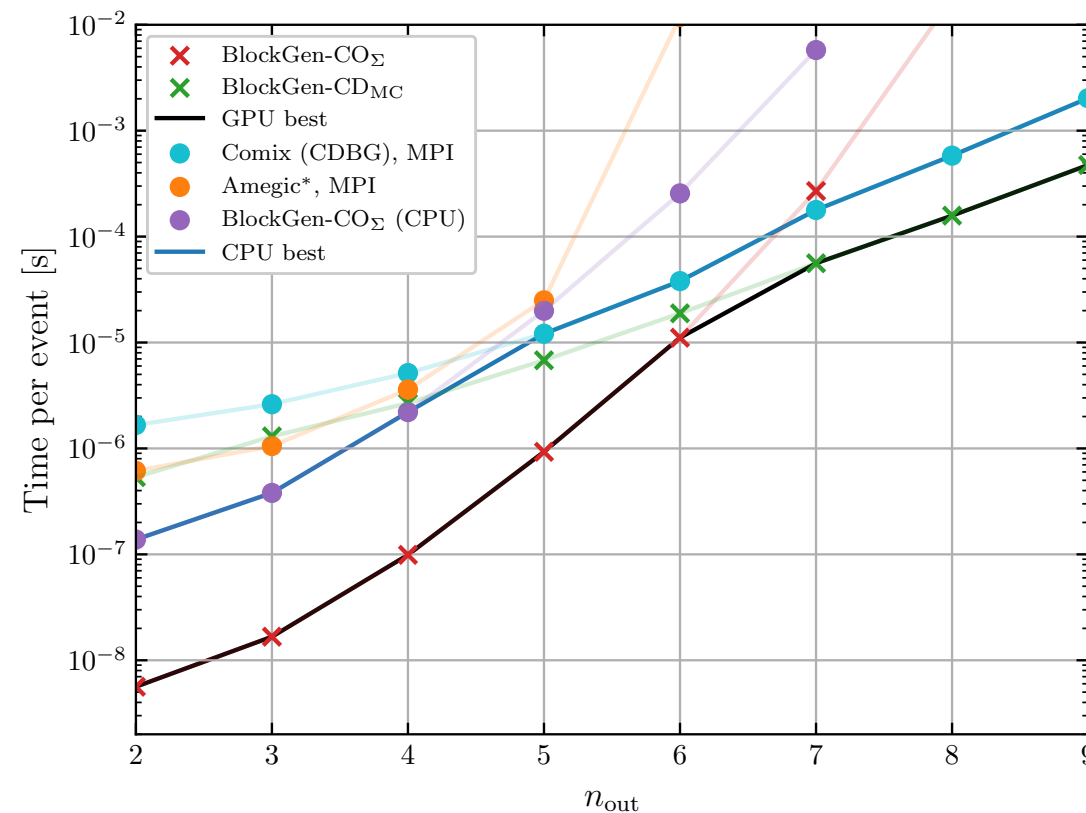
A. Valassi - ACAT, Bari, 24 October 2022

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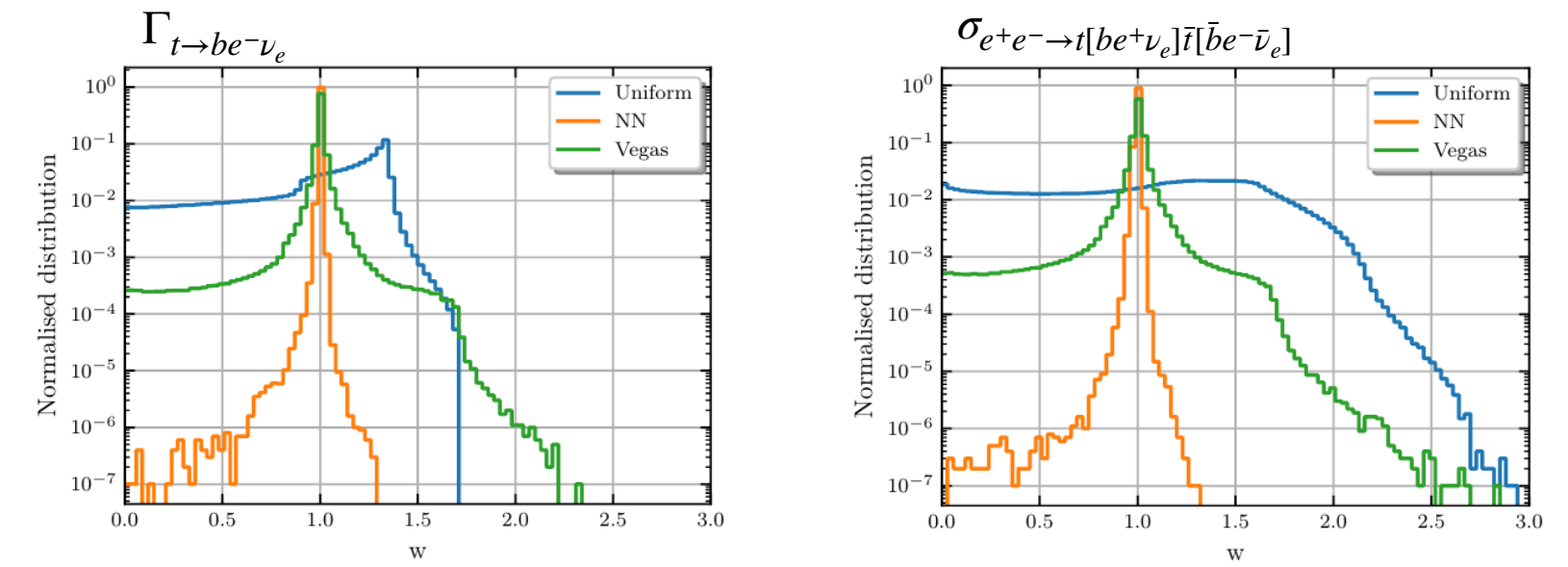
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Enrico Bothmann

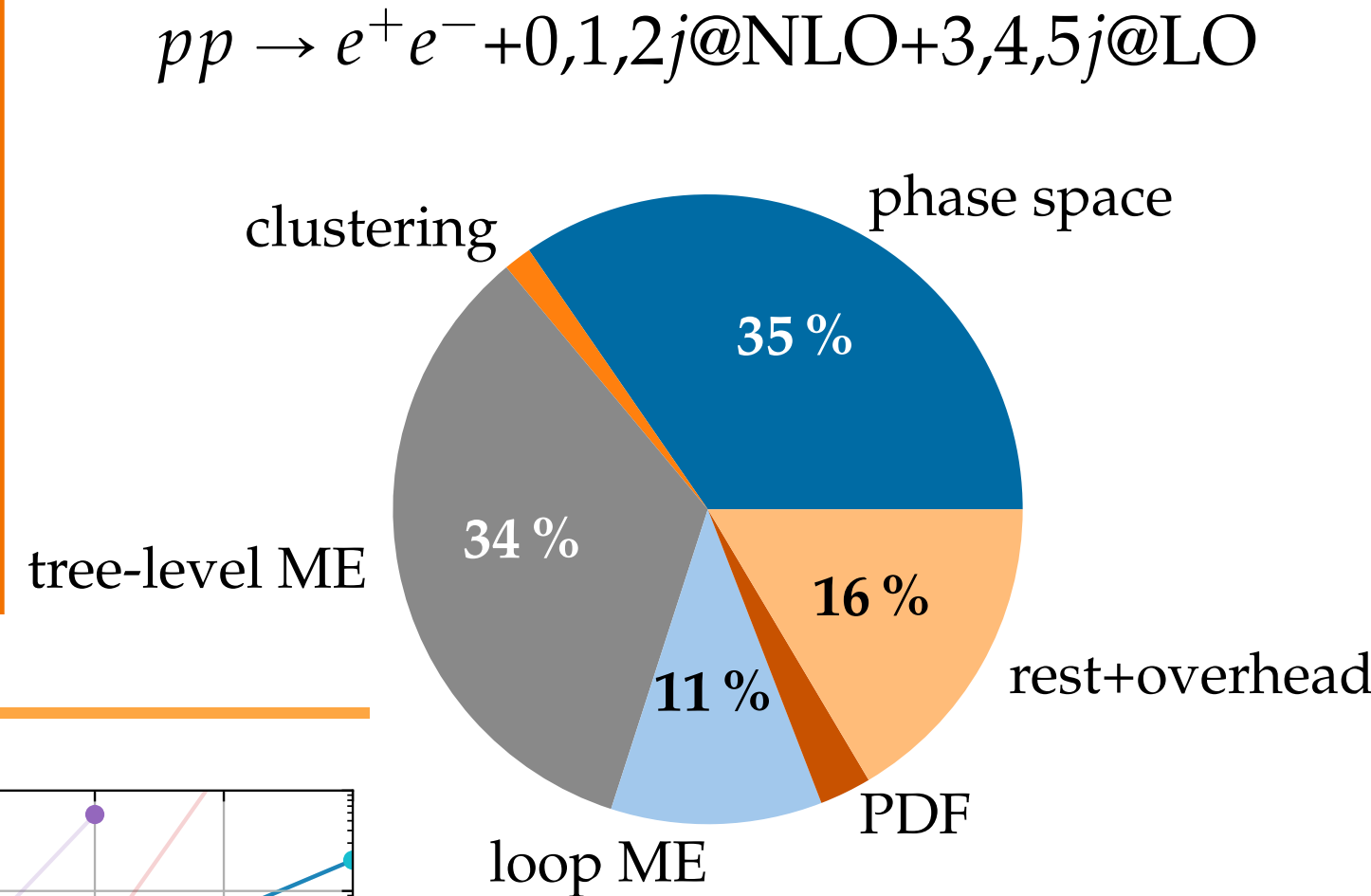
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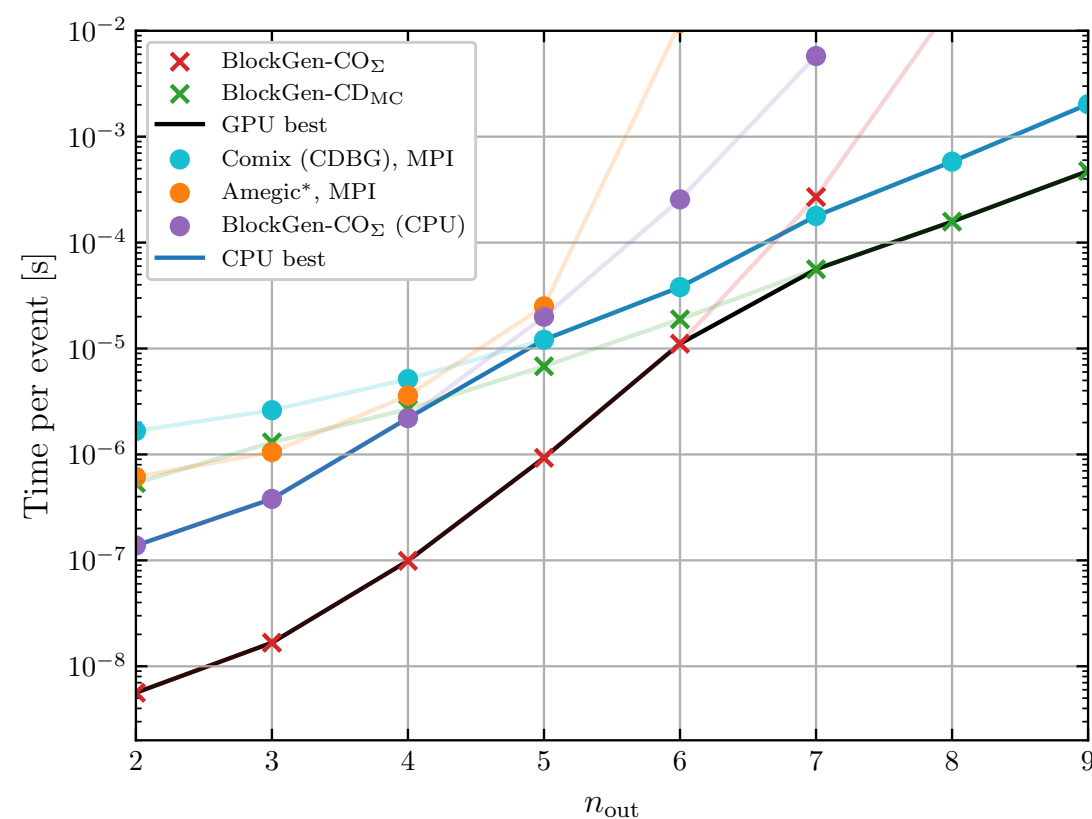
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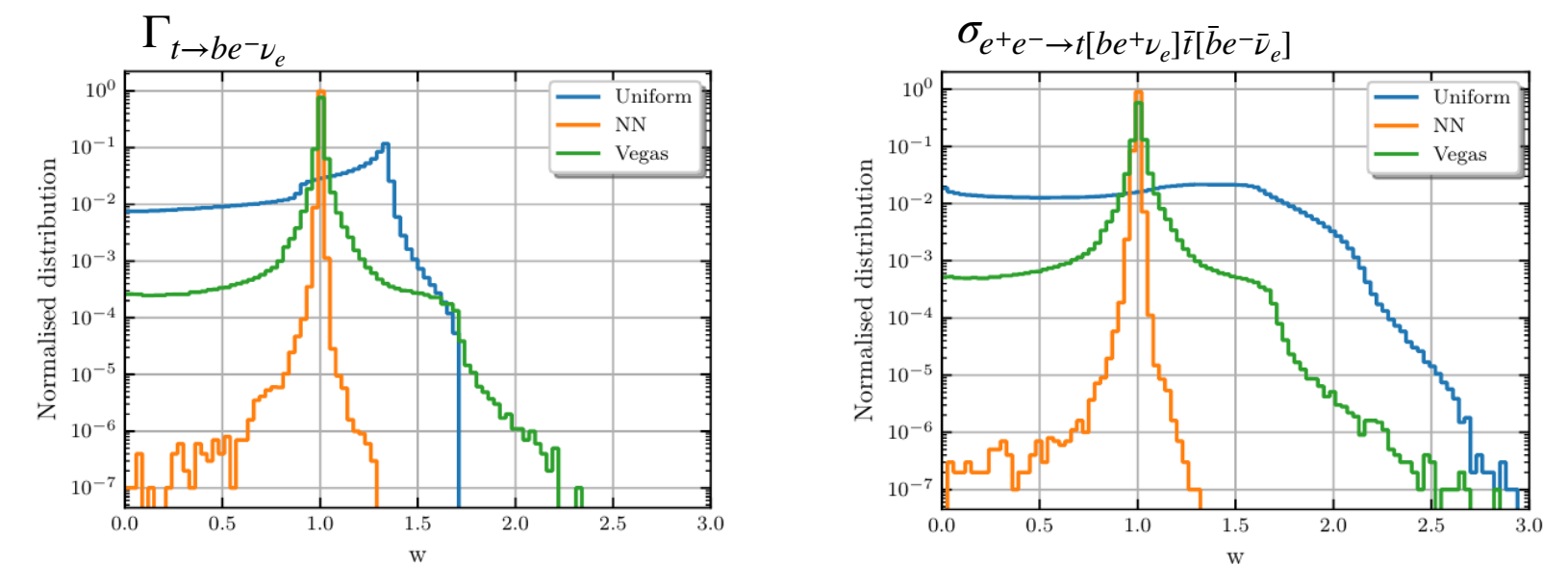
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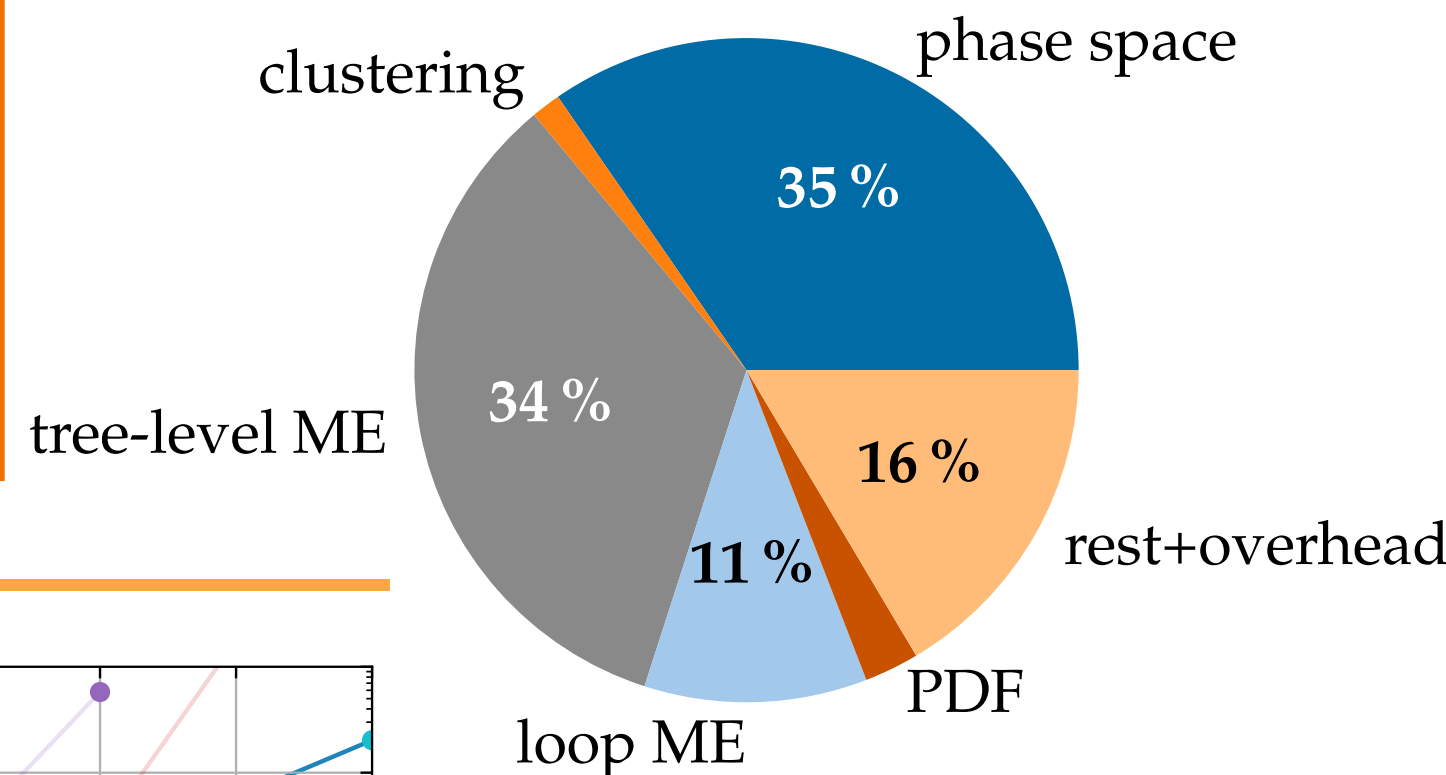


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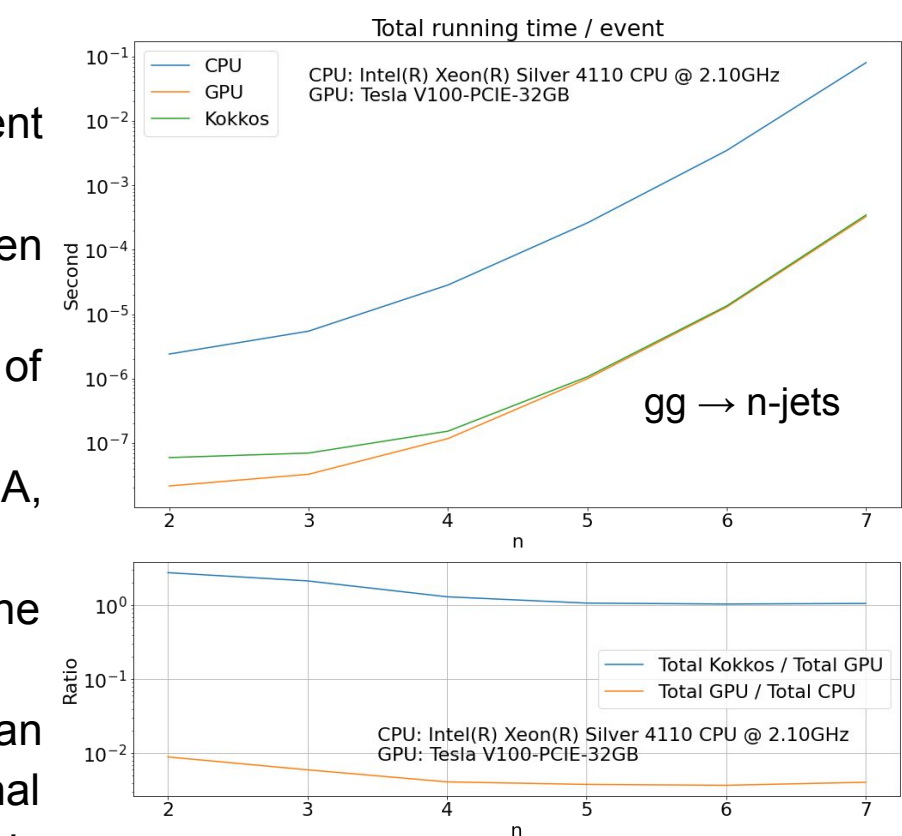
$$pp \rightarrow e^+e^- + 0,1,2j@NLO + 3,4,5j@LO$$



Performance of Kokkos

- So, does Kokkos provide equivalent performance?
- Plot shows early versions of BlockGen calculating the process: $gg \rightarrow n$ -jets
- Time per Event on y-axis, number of outgoing partons on x-axis
- Compare CPU with C++, GPU with CUDA, and GPU with Kokkos
- Can see the CUDA is 100x faster than the CPU for this example
- Kokkos is slightly less performant than CUDA at low multiplicity (low computational complexity), but reaches comparable performance as multiplicity increases.

Taylor Childers



Precision Frontier

Johann Usovitsch

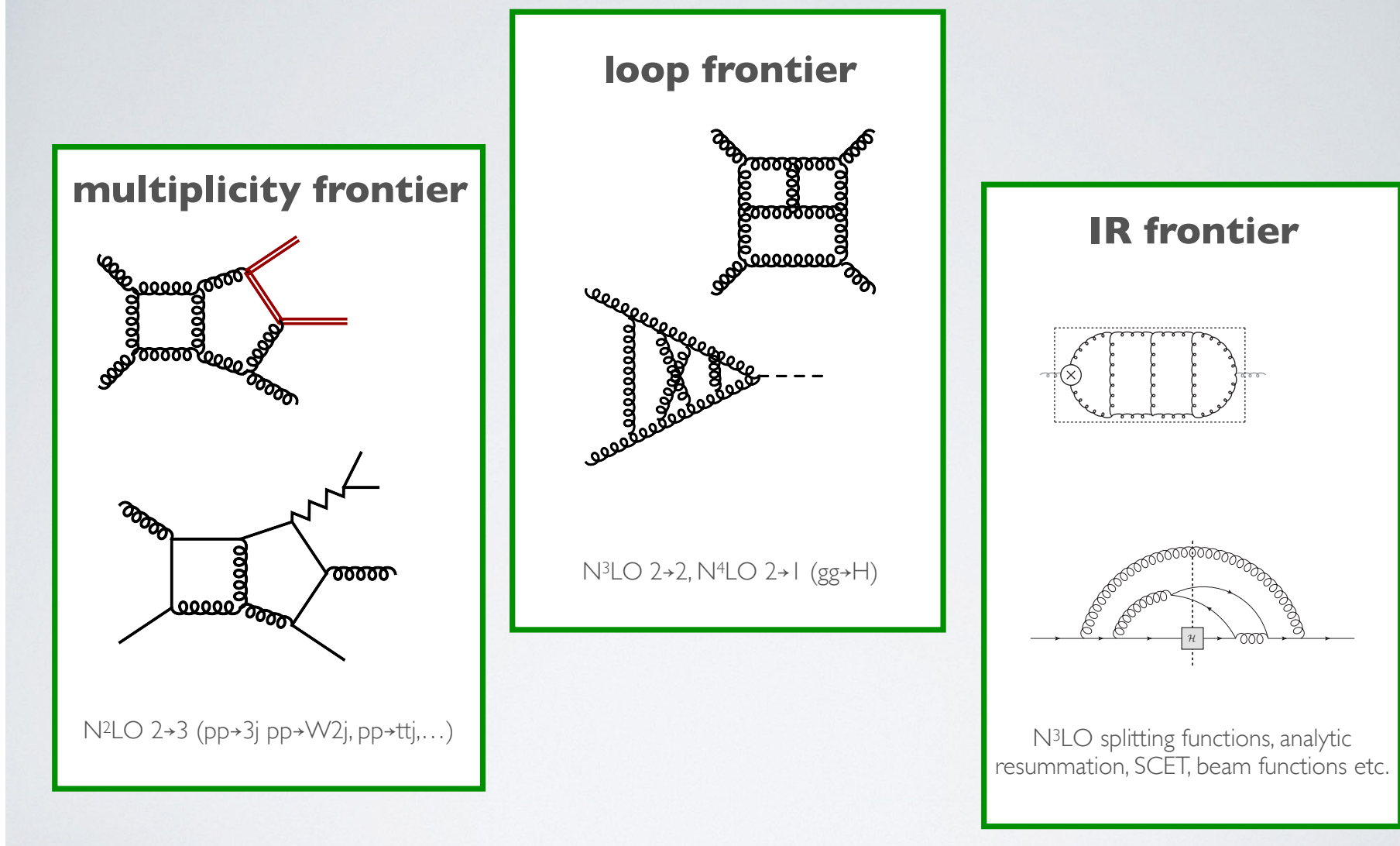
Precision test of the Standard Model Future prospects

Overview of future experiments as of 2022

	Experiment uncertainty			Theory uncertainty
	ILC	CEPC	FCC-ee	Current
M_W [MeV]	3-4	3	1 0.3	4
$\sin^2 \theta_{\text{eff}}^l$ [10^{-5}]	1	2.3	?0.6	4.5
Γ_Z [MeV]	0.8	0.5	0/10.025	0.4
R_f [10^{-5}]	14	17	ϕ1	15

- Recent update from [\[Alain Blondel, Patrick Janot, Eur.Phys.J.Plus 137 \(2022\) 1\]](#)
- To match the precision of the experiment we compute **3-loop** and **4-loop** Standard Model predictions

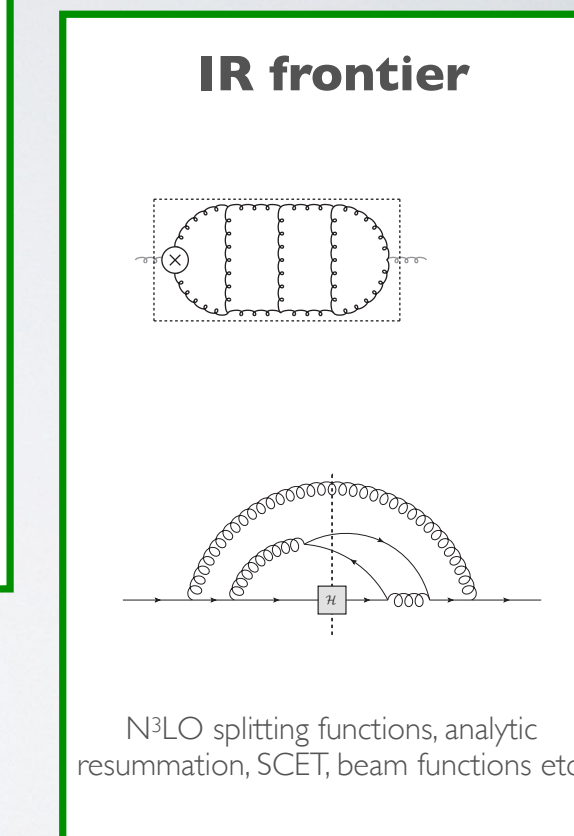
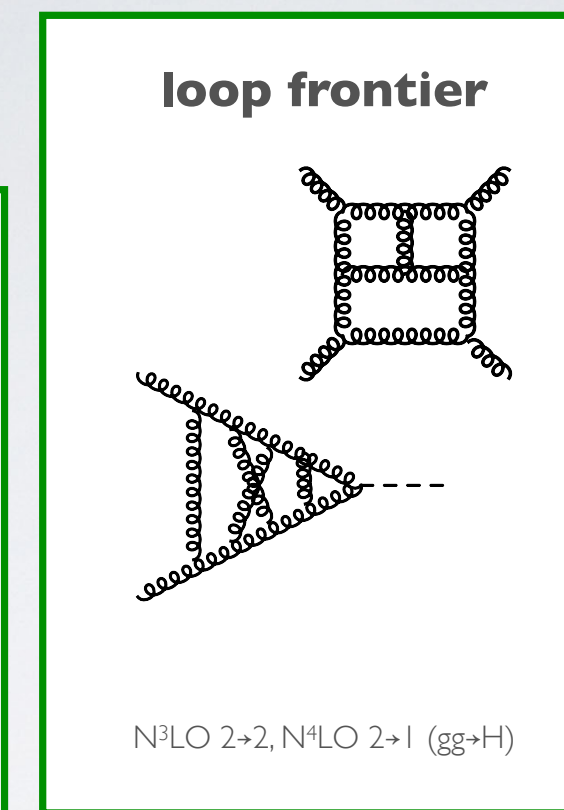
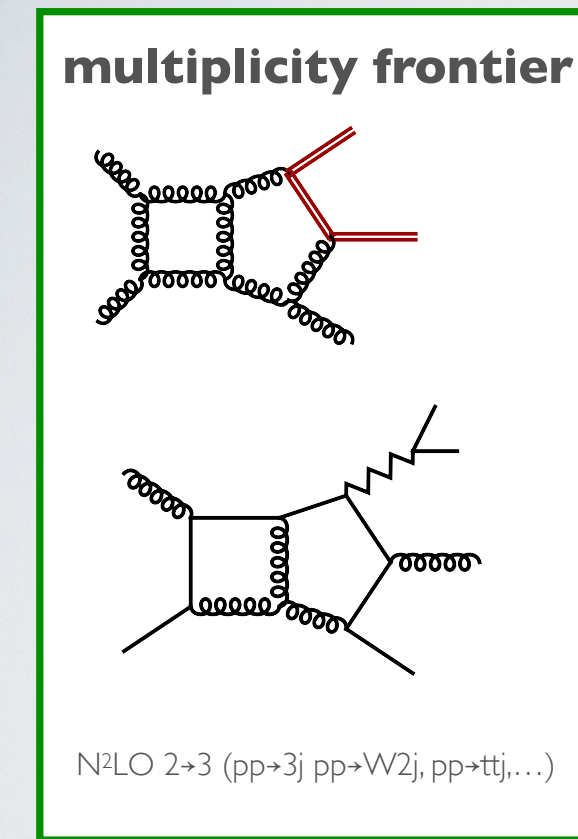
Simon
Badger



- Numerical methods
 - Avoiding algebraic complexity
- Physics informed
 - Exploiting known structures

Precision Frontier

Simon
Badger



Advanced
Pen and Paper Physics



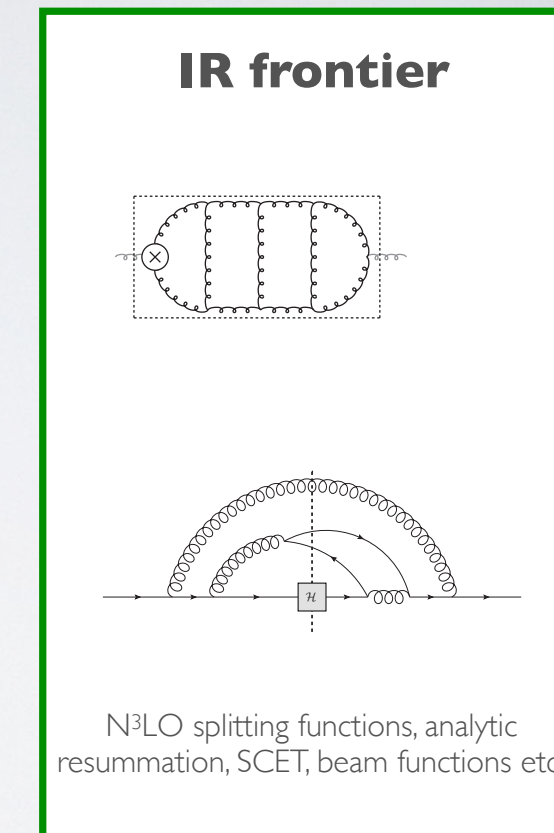
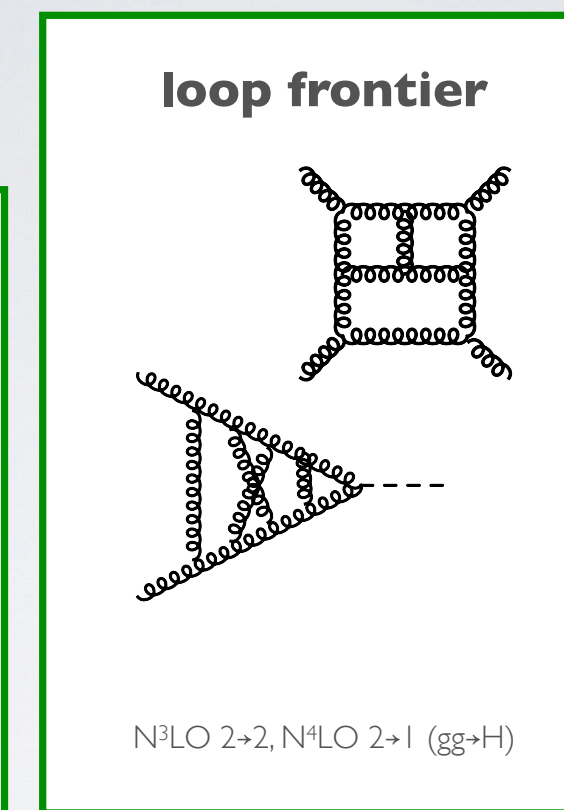
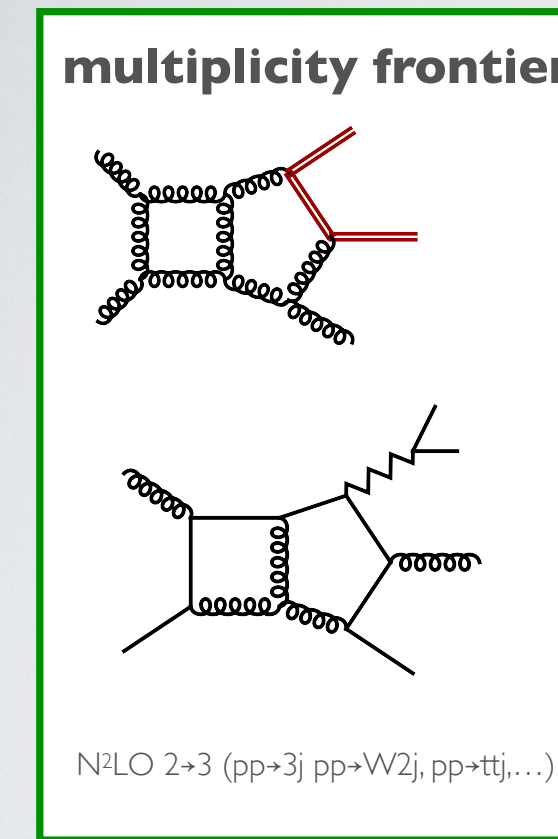
BARI

9 / 29

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Precision Frontier

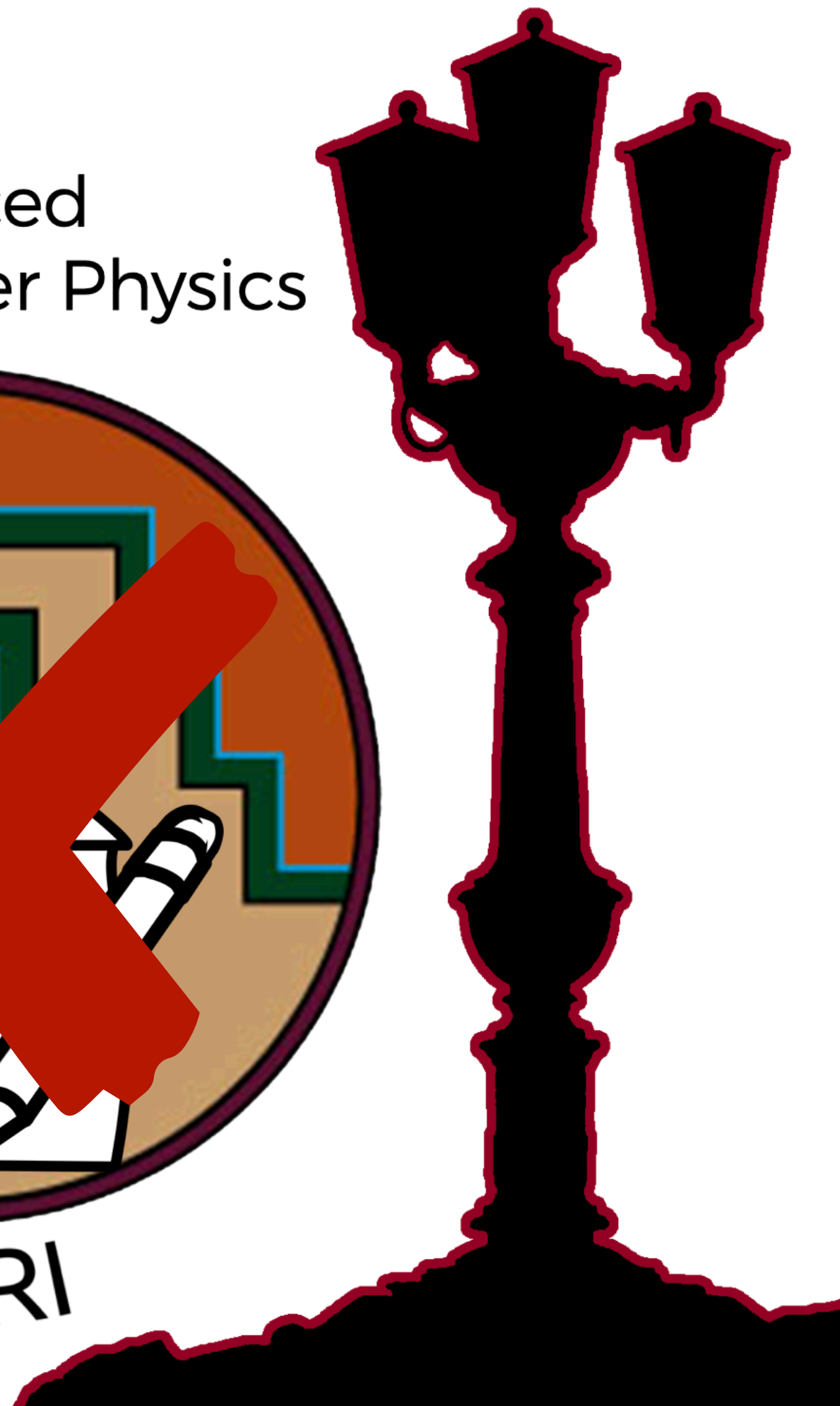
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BARI



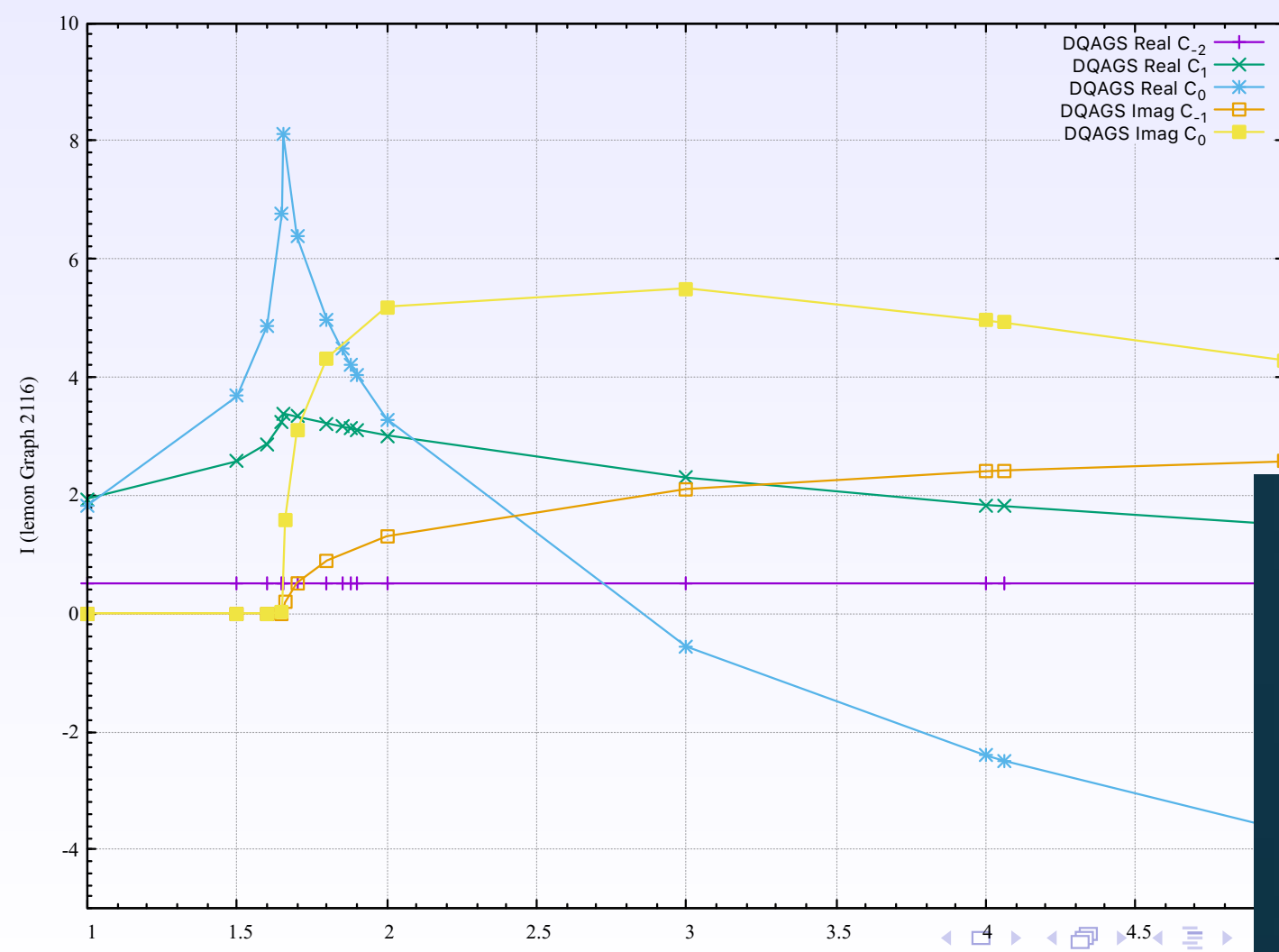
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Numerical approaches

Kinematic distributions

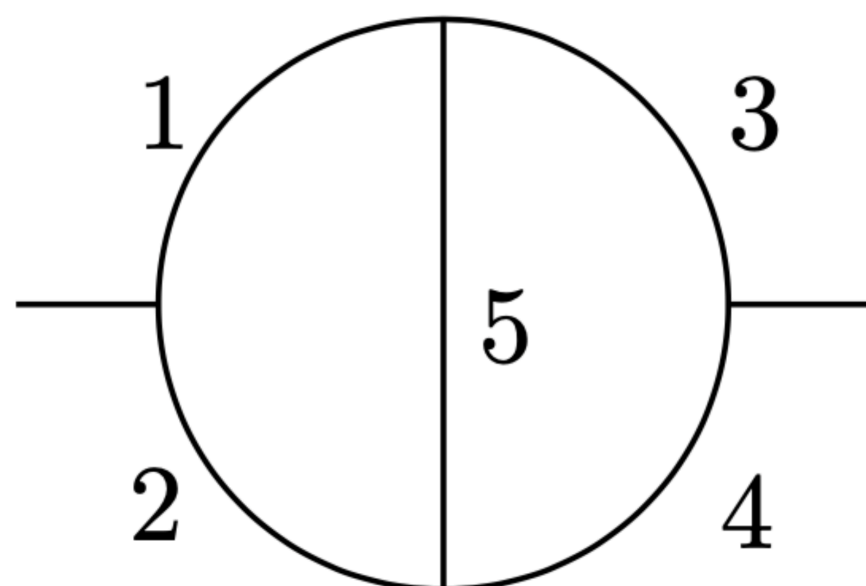
- Introduction - Overview
- Loop Integrals
- Methods
- Results**
- Conclusions

Results Graph 2116: C_{-2}, C_{-1}, C_0 as a function of s



de Doncker, Yuasa, Ishikawai, and Kato

Elise de Doncker

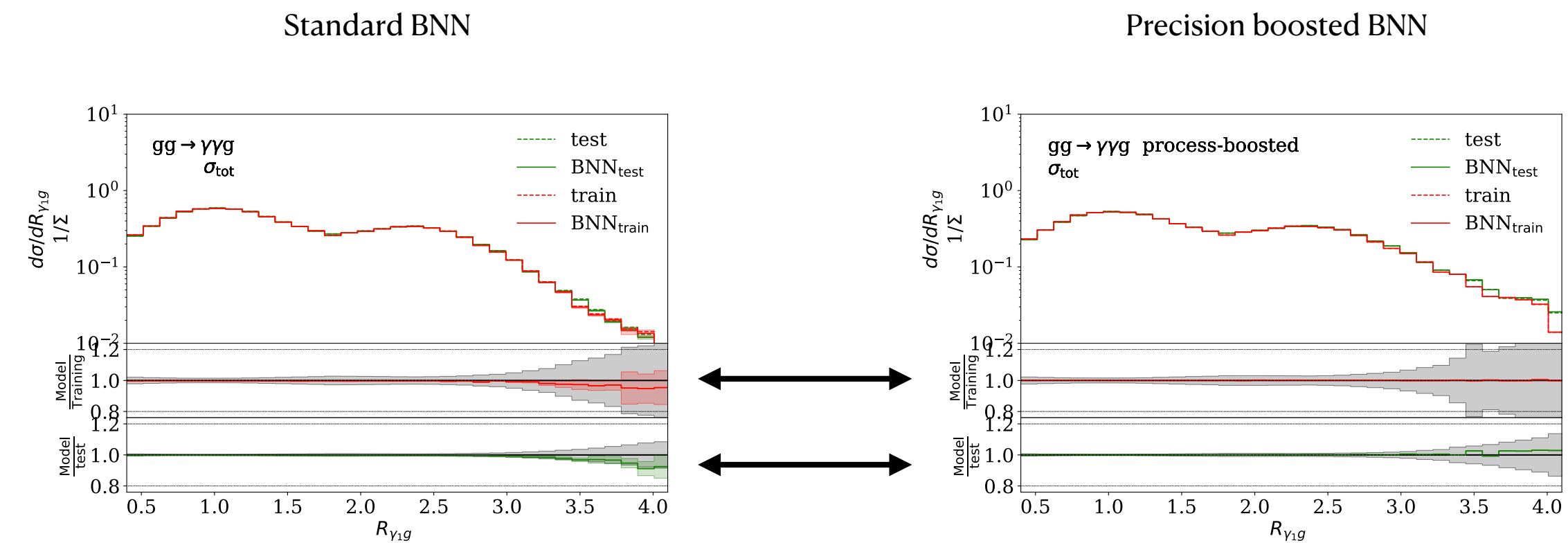


Result

- 100 nodes
- 4 hidden layers
- 4M x 800 = 3.2B PS points

$$p = \log_{10} \left| \frac{e-t}{t} \right|$$

$$L = \text{MSE} \left(f(s_1, \dots, s_m; x_1, \dots, x_k), \frac{d\mathcal{N}(s_1, \dots, s_m; x_1, \dots, x_k)}{dx_1 \dots dx_k} \right)$$

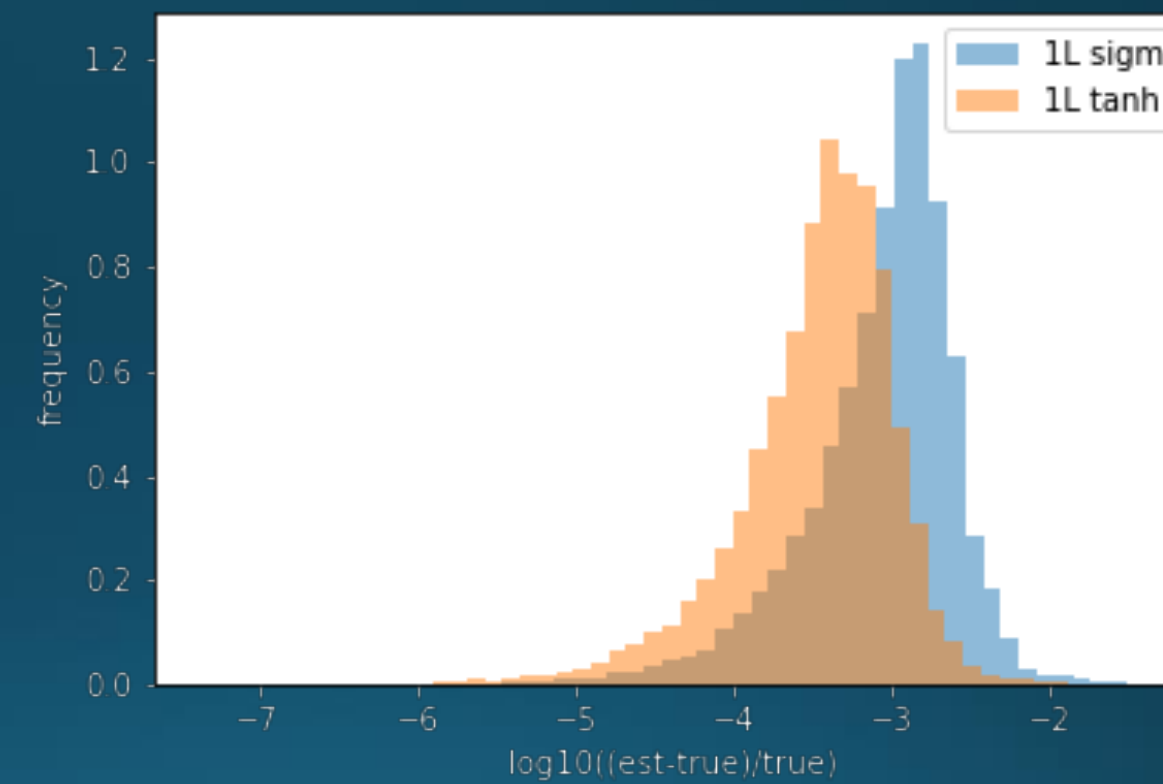


Uncertainties

Gray shades indicate statistical limitation of training data.

Anja Butter

Daniel Maitre

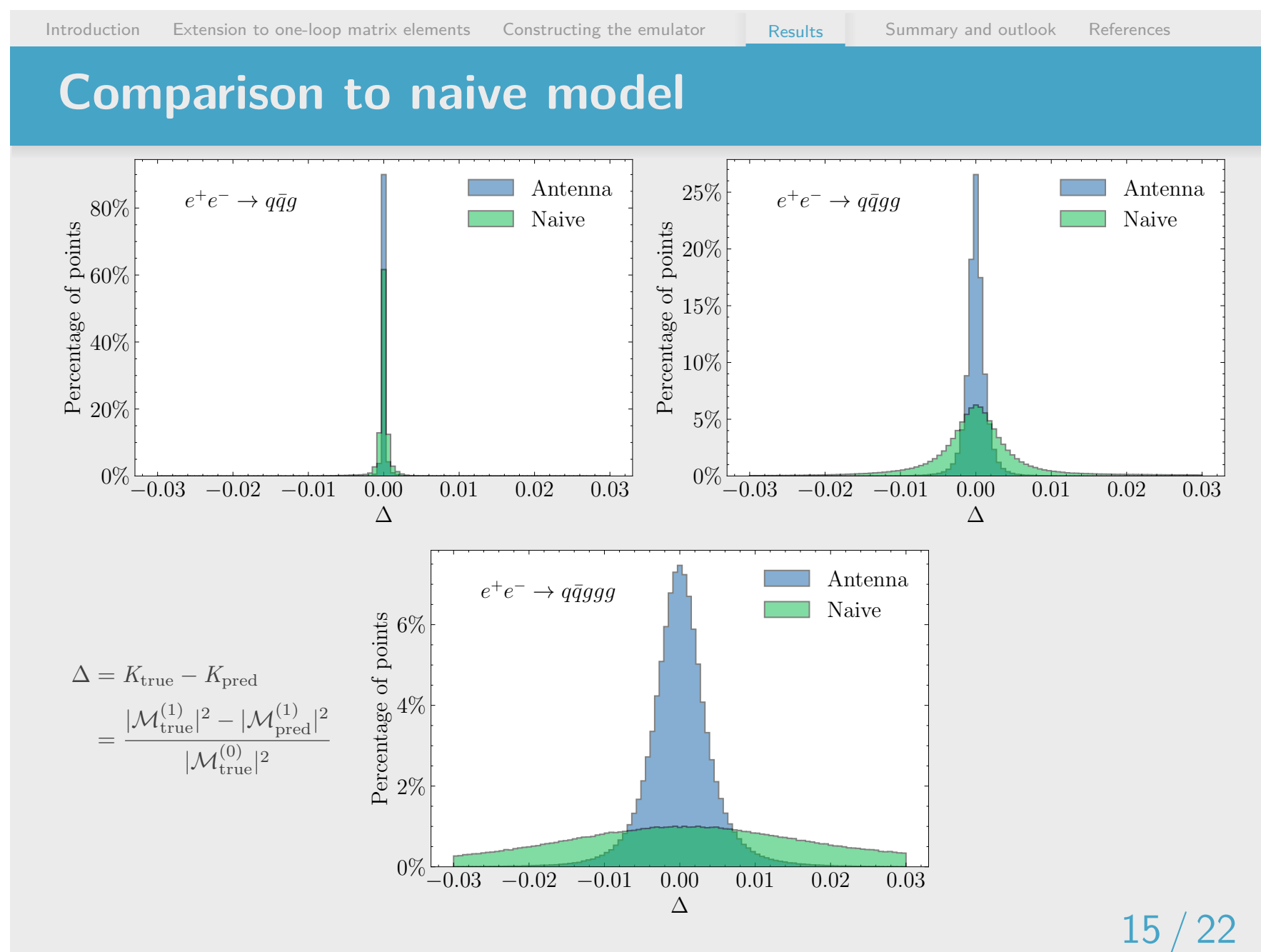


Machine learning the primitive

Numerical approaches - physics informed

$$K_{n+1} = C_0 + \sum_{\{ijk\}} C_{ijk} \frac{X_{ijk}^1}{X_{ijk}^0}$$

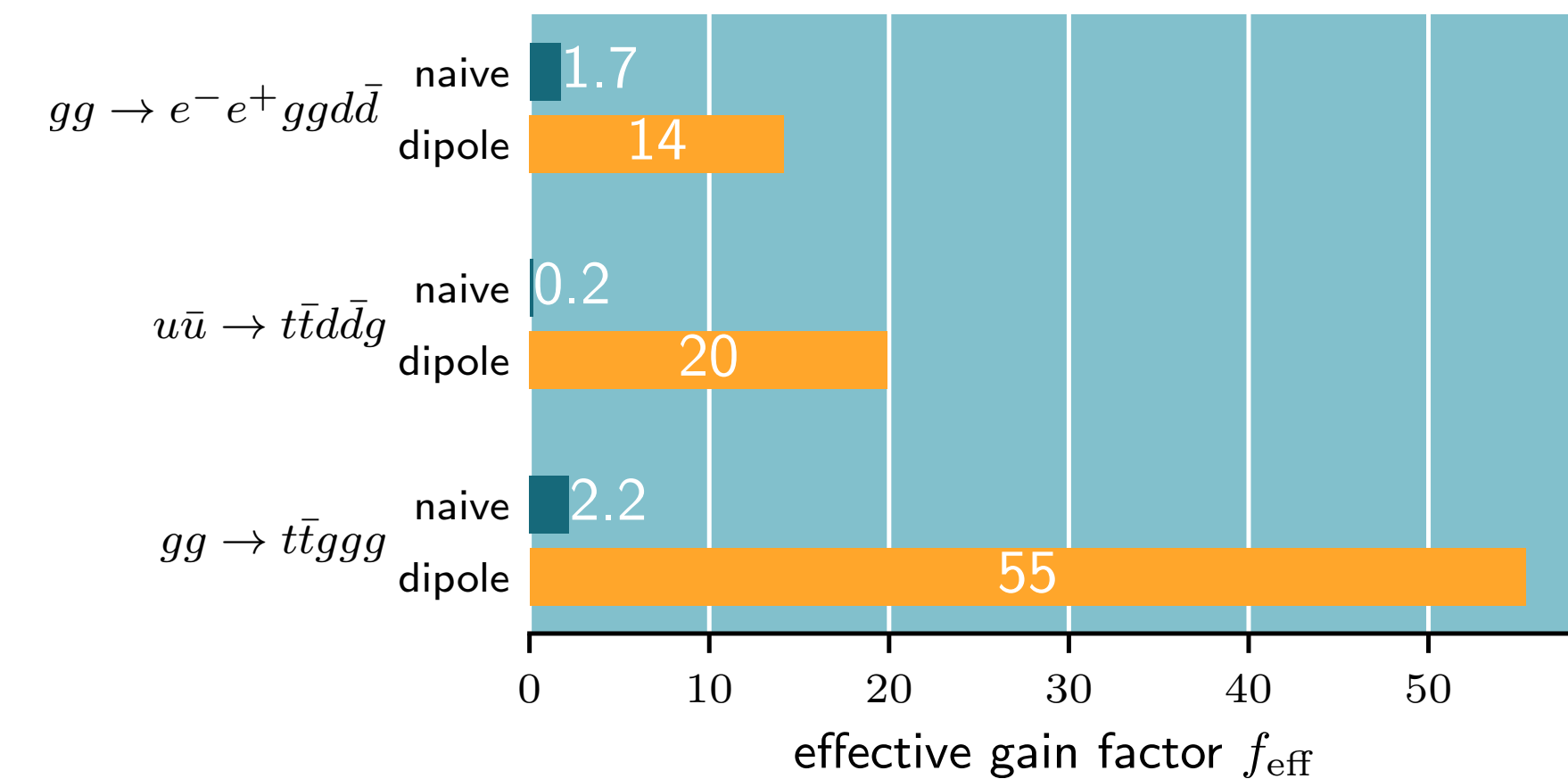
Learning K factors/matrix elements
- coefficients of antenna functions



Results: effective gain factors for LHC multi-jet processes

$$f_{\text{eff}} := \frac{T_{\text{standard}}}{T_{\text{surrogate}}}$$

Using 1M training events:



Analytic approaches

Ryan Moodie

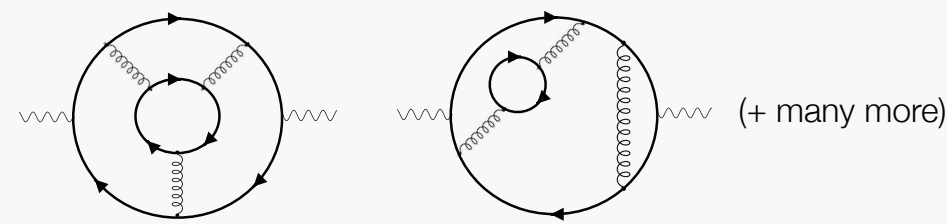
Introduction	Processes	Computation	Finite fields	Reconstruction	Performance	Conclusion
○	○○	○○	○○○	○○○○	○○○○○	○

Timing

Finite fields

Local unitarity

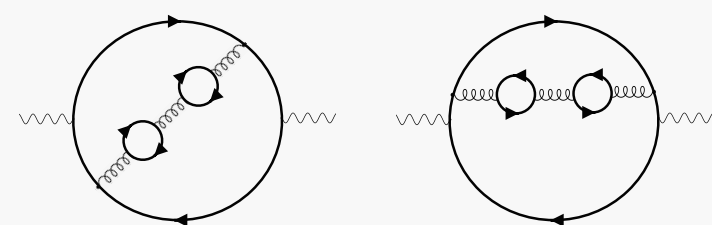
N_f part @N³LO $e^+e^- \rightarrow jj$



$$K_{\alpha_s^3}^{N_f,LU} = -77.1(1.7)$$

$$K_{\alpha_s^3}^{N_f,bm} = -76.8086$$

N_f^2 part @N³LO $e^+e^- \rightarrow jj$



$$K_{\alpha_s^3}^{N_f^2,LU} = -0.35(24)$$

$$K_{\alpha_s^3}^{N_f^2,bm} = -0.331415$$

singlet part @N³LO $e^+e^- \rightarrow jj$



$$K_{\alpha_s^3}^{\text{singlet},LU} = -25.6(1.5)$$

$$K_{\alpha_s^3}^{\text{singlet},bm} = -26.4435$$

Benchmarks:
Herzog, Ruijl, Ueda, Vermaseren, Vogt :
arXiv:1707.01044

Testing {
✓ N3LO IR cancellations
✓ 3-loop UV renormalisation
✓ 1,2,3-loop self-energies

Channel	f64/f64		Evaluation strategy	
	Time (s)	f (%)	Time (s)	f (%)
$gg \rightarrow ggg$	1.39	69	1.89	77
$gg \rightarrow \bar{q}qg$	1.35	91	1.37	91
$qg \rightarrow qgg$	1.34	92	1.57	93
$q\bar{q} \rightarrow ggg$	1.34	93	1.38	93
$\bar{q}Q \rightarrow Q\bar{q}g$	1.14	99	1.16	99
$\bar{q}\bar{Q} \rightarrow \bar{q}\bar{Q}g$	1.36	99	1.39	99
$\bar{q}g \rightarrow \bar{q}Q\bar{Q}$	1.36	99	1.39	99
$\bar{q}q \rightarrow Q\bar{Q}g$	1.14	99	1.14	99
$\bar{q}g \rightarrow \bar{q}q\bar{q}$	1.84	99	1.90	99
$\bar{q}\bar{q} \rightarrow \bar{q}\bar{q}g$	1.82	99	1.94	99
$\bar{q}q \rightarrow q\bar{q}g$	1.71	99	1.77	99
$gg \rightarrow \gamma\gamma g *$	9	99	26	99

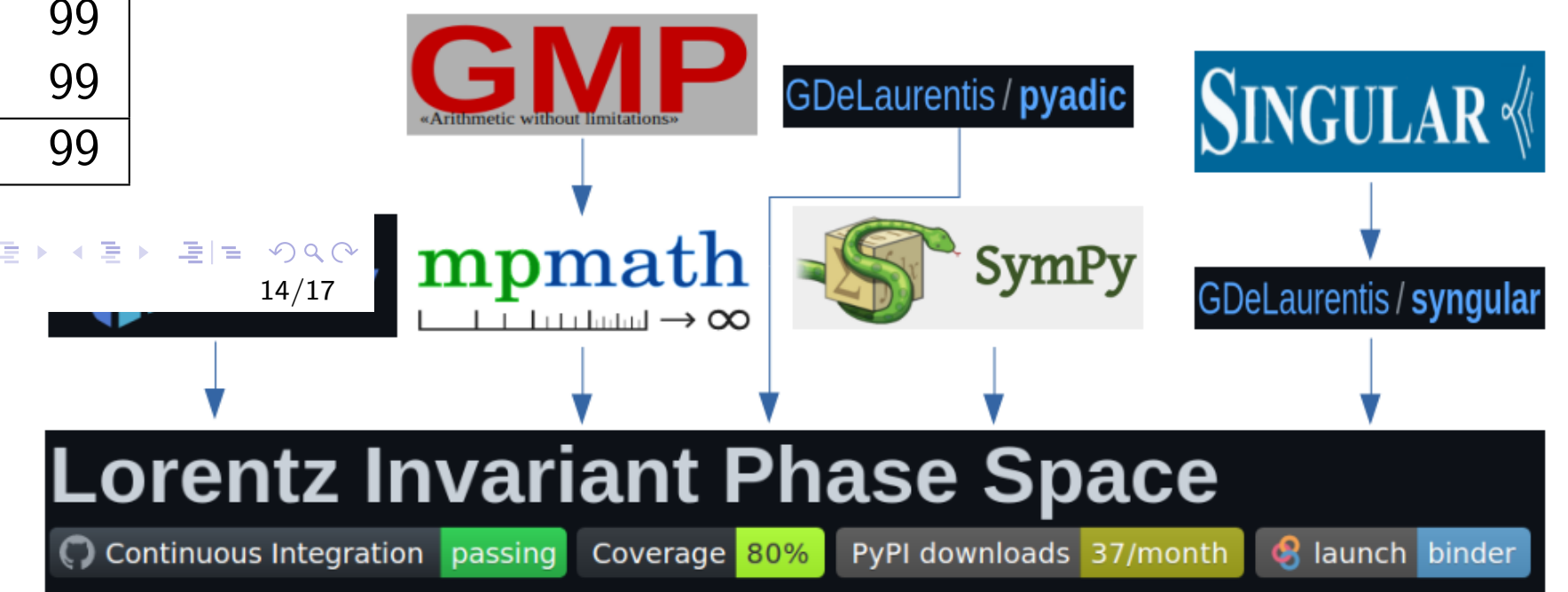
Ryan Moodie (Turin)

Two-loop five-point amplitudes in massless QCD with finite fields

Giuseppe de Laurentis

DEPENDENCIES	3. LIPS: LORENTZ INVARIANT PHASE SPACE	4. CONCLUSION
	○○○○	○

IS GRAPH¹²³⁴



¹Charles R. Harris et al. "Array programming with NumPy". In: *Nature* 585 (2020), pp. 357–362. DOI: 10.1038/s41586-020-2649-2.

²Fredrik Johansson et al. *mpmath: a Python library for arbitrary-precision floating-point arithmetic (version 0.18)*. <http://mpmath.org/>. 2013.

³Aaron Meurer et al. "SymPy: symbolic computing in Python". In: *PeerJ Computer Science* 3 (Jan. 2017), e103. ISSN: 2376-5992.

⁴Wolfram Decker et al. *SINGULAR 4-3-0 — A computer algebra system for polynomial computations*. <http://www.singular.uni-kl.de>. 2022.

Zeno Capatti

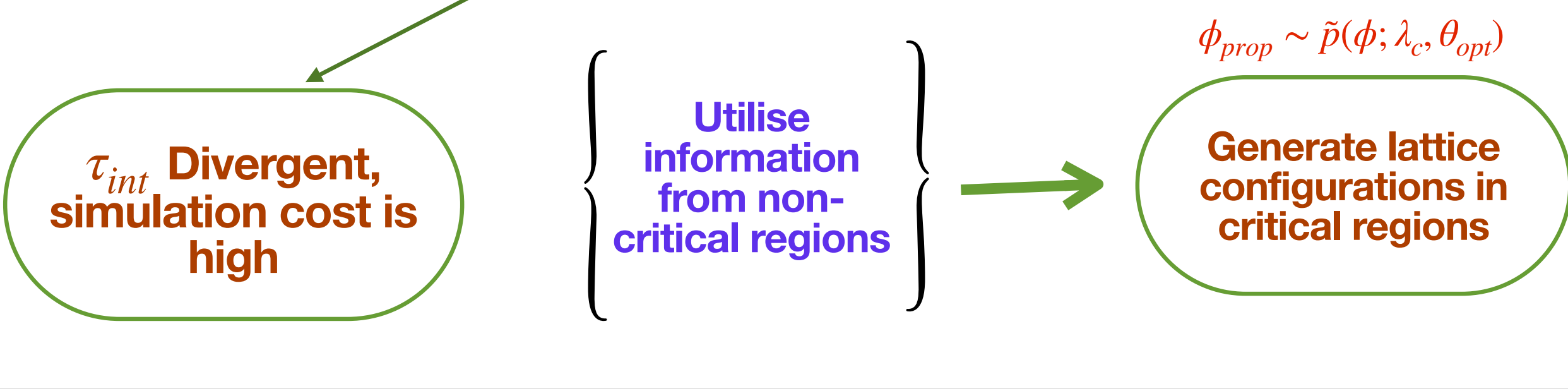
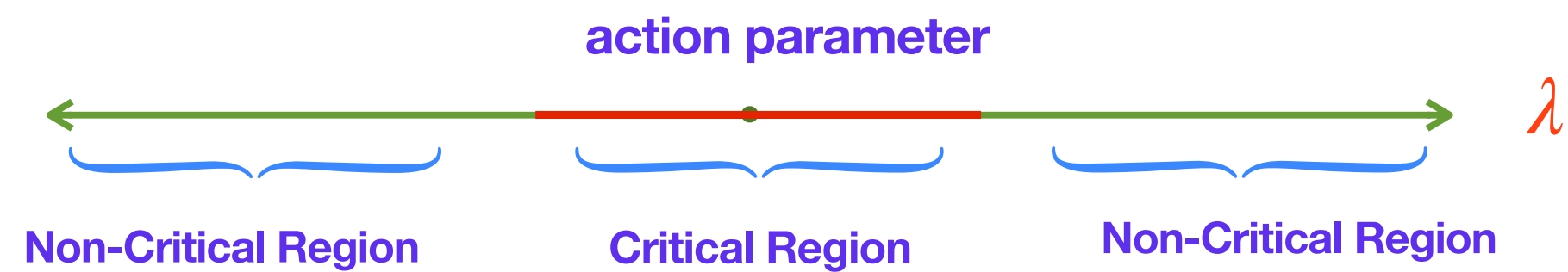
Anke Biekötter for Track 3

Non-perturbative physics

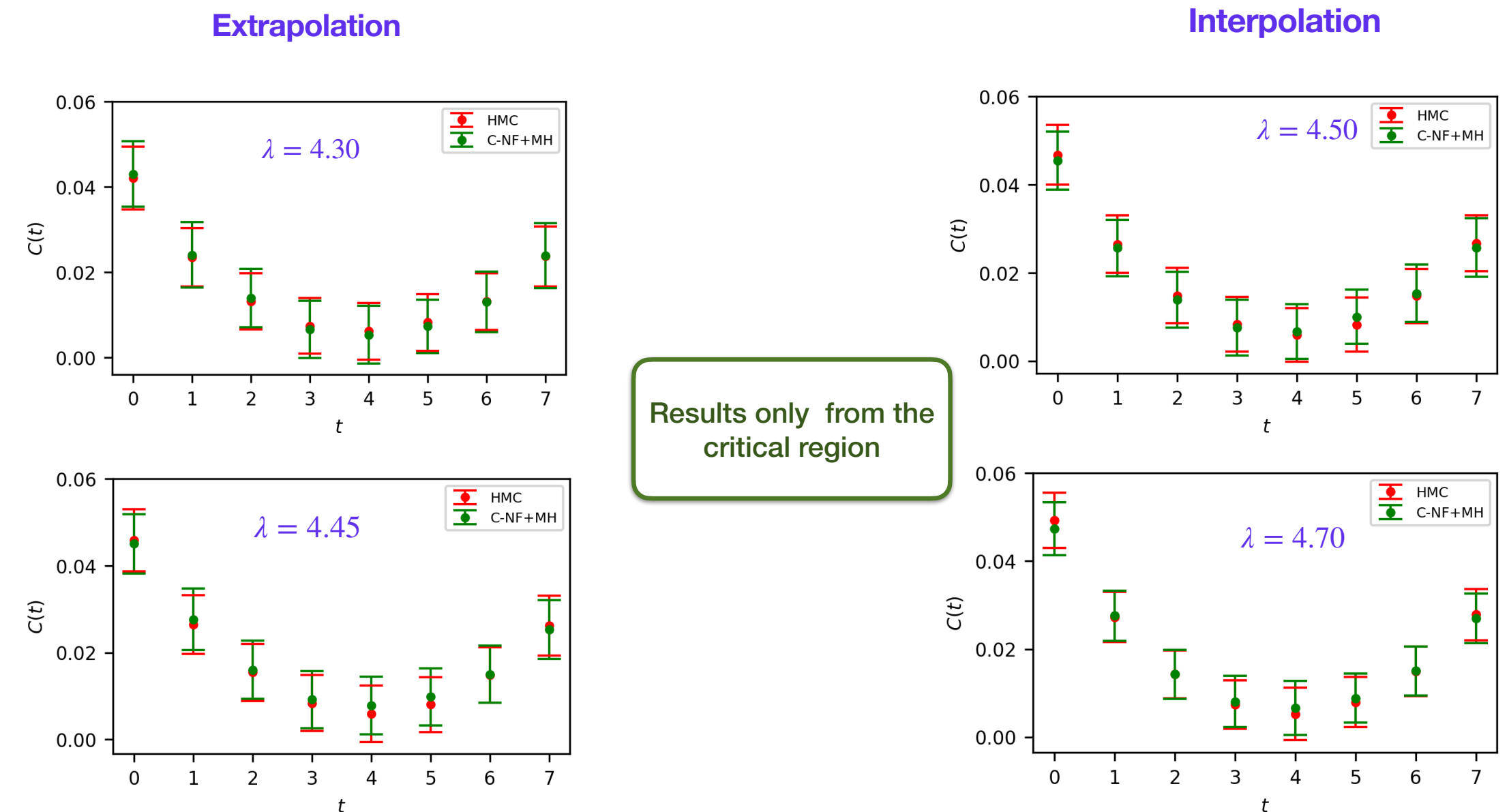
Ankur Singha

Simulation at multiple λ values

For a lattice action : $S(\phi, m_{fixed}, \lambda) \longrightarrow p(\phi | \lambda_i)$

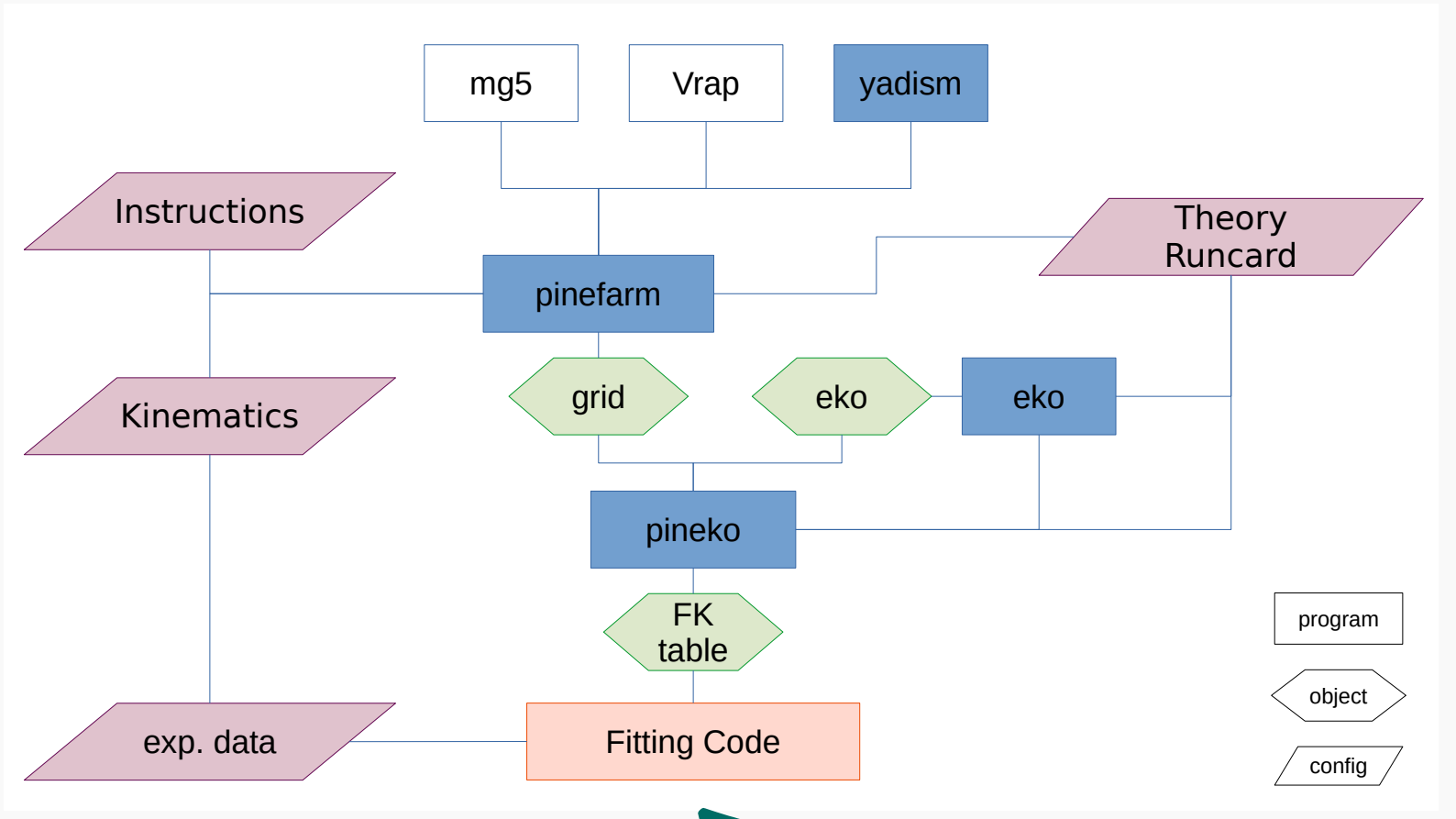


Results: Correlation Functions

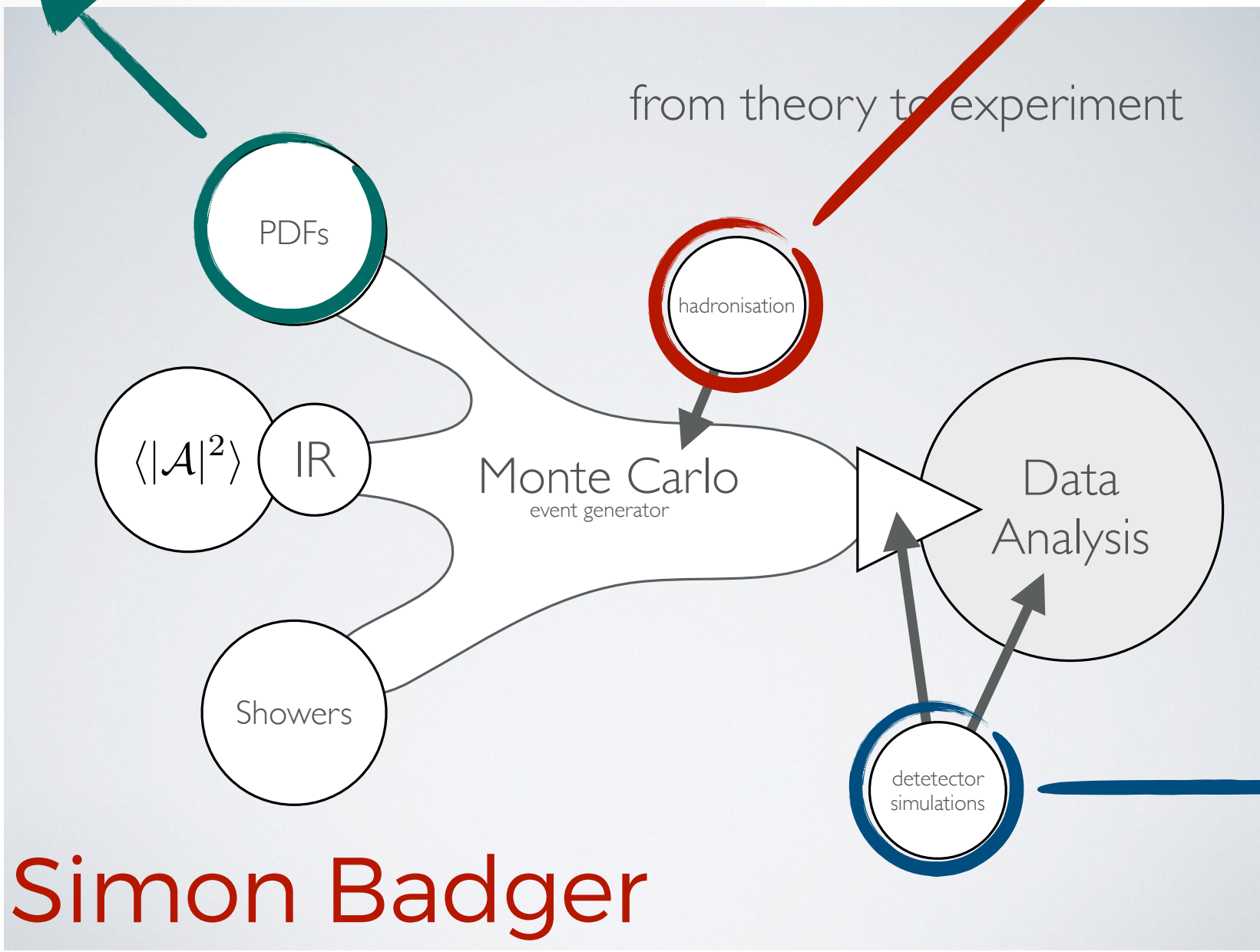


Produce FastKernel (FK) tables!

Felix Hekhorn



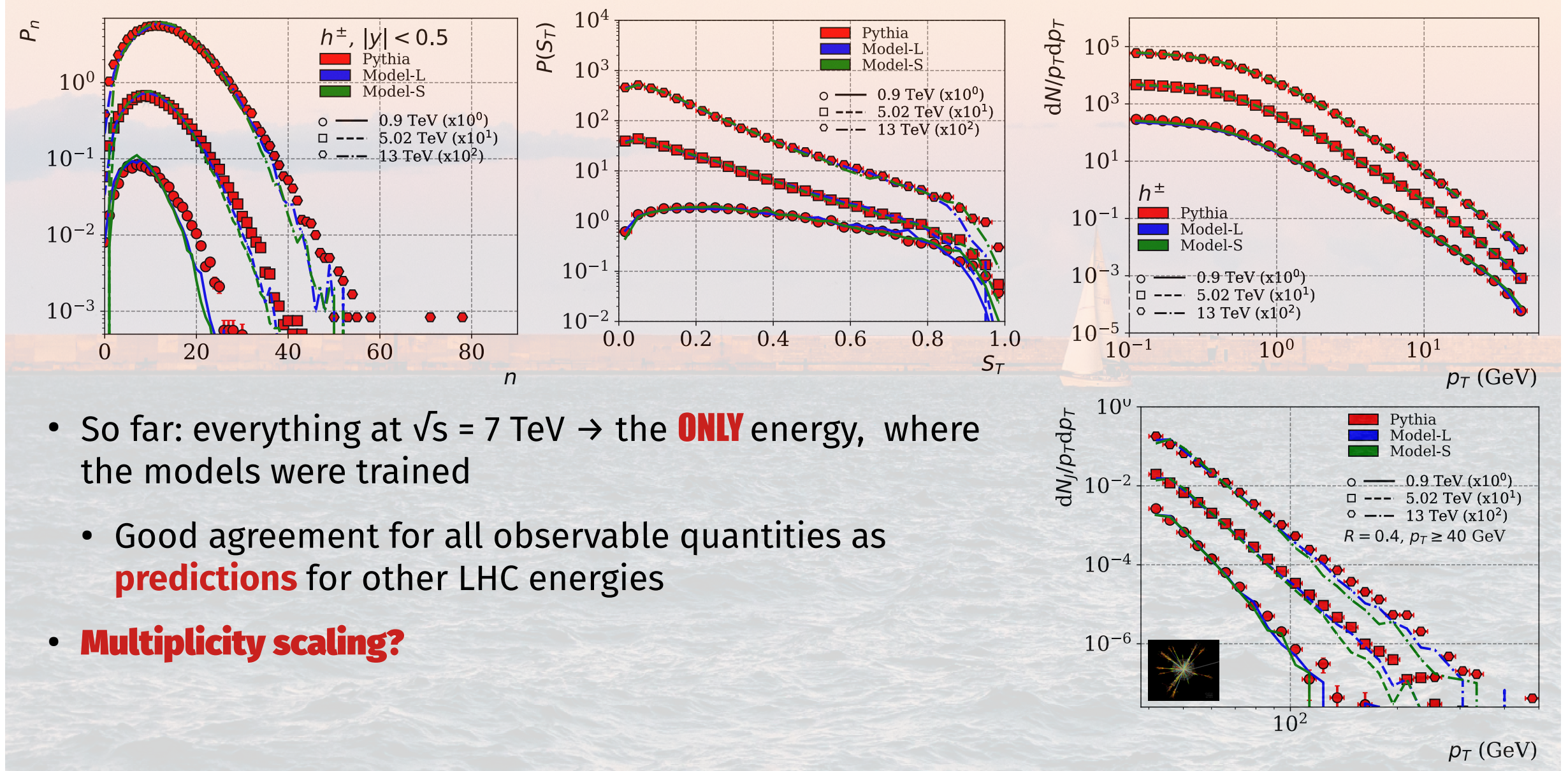
The workhorse in the background



Simon Badger

Alberto Martini

Proton-proton @ 0.9-13 TeV, Predictions

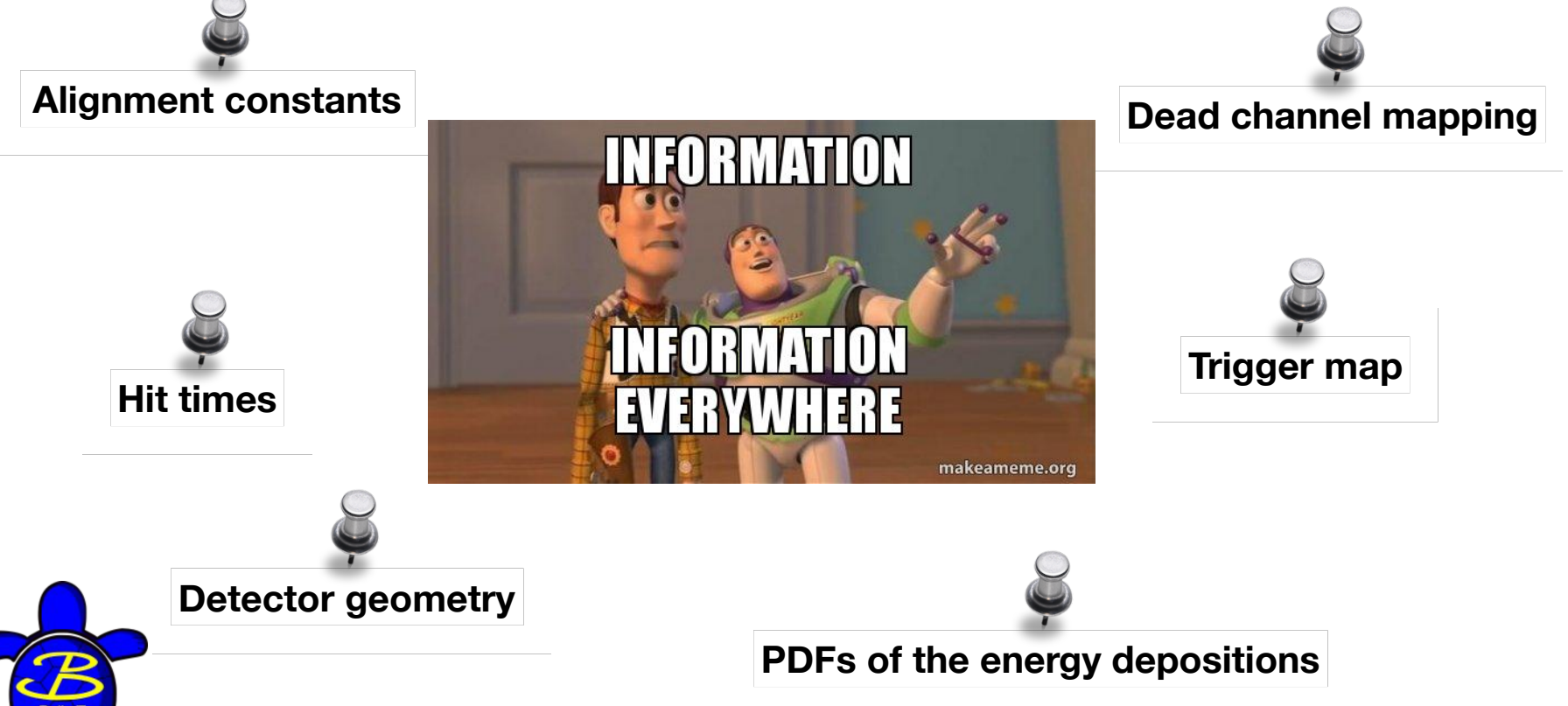


- So far: everything at $\sqrt{s} = 7$ TeV \rightarrow the **ONLY** energy, where the models were trained
- Good agreement for all observable quantities as **predictions** for other LHC energies
- **Multiplicity scaling?**

Preparation of detector configuration

MCrd should have data-driven detector configurations

Manipulate the data calibration constants to provide MC payloads \rightarrow gather together detector related quantities from each sub-detector



Beyond Standard Model

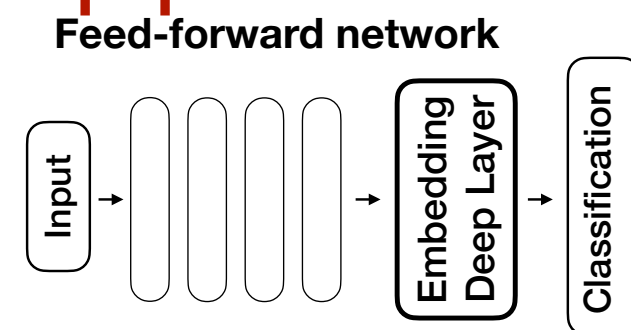
Andrea Wulzer

- Looking for new physics
 - LHC and beyond
- Model independence
- Recasting
- Symmetries

Sven Krippendorf

How to search for symmetries?

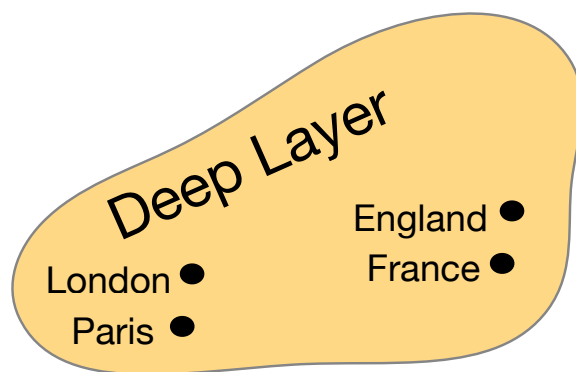
No direct optimisation available: embedding in deep layer



We need: group input with the same meaning together

Word2Vec does it:
(England - London = Paris - France)

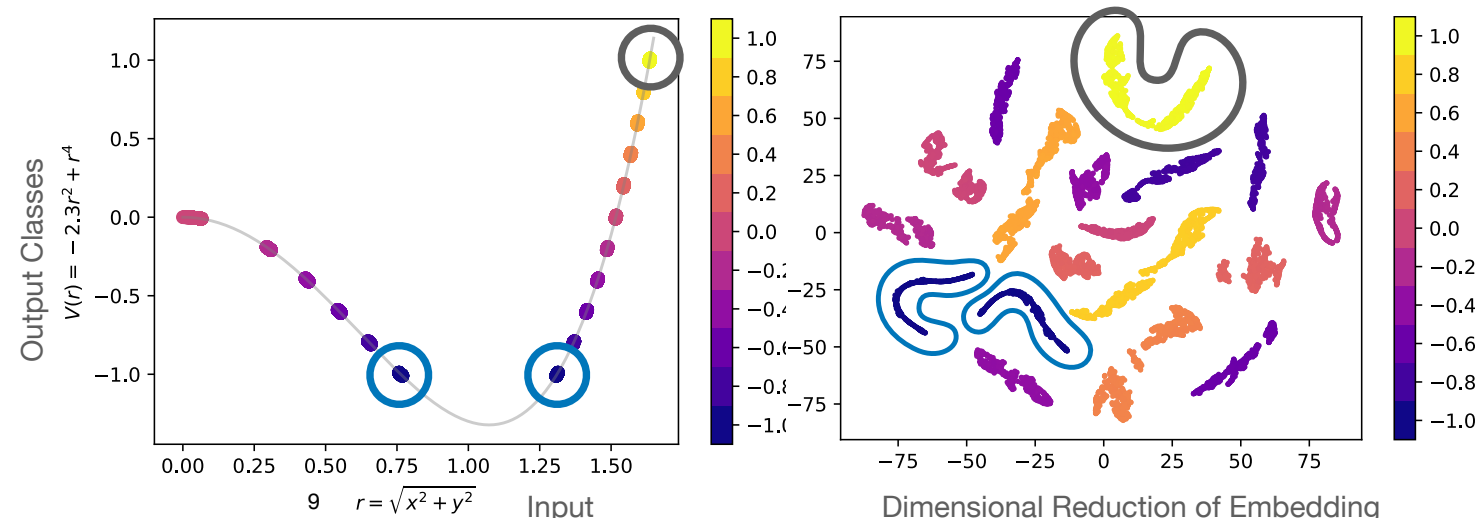
[1301.3781, used for re-discovering periodic table 1807.05617, classifying scents of molecules 1910.10685]



Can we search for symmetries in this way?

Yes!

Examples: SO(2), SU(2), discrete symmetries (CICY)



Krippendorf, Syvaeri 2020

The LHC g.o.f. challenge

By analysing the LHC data, we would like to find evidence of **failure of the SM theory**, suggesting need of **BSM**.

This is a tremendously hard gof problem!

BSM is tiny departure from SM, or large in tiny prob. region
Affecting few (unknown) observables over ∞ many we can measure

Model-dependent H_1 R
BSM searches

- Optimise sensitivity to **one specific BSM model**
- Fail to discover other models. **What if the right theoretical model is not yet formulated?**

Model-independent H_w R
searches

- Could reveal **truly unexpected** new physical laws.
- No hopes to find Optimal strategy. For a Good strategy, we need a **good choice of H_w** .

Beyond Standard Model

Theo Heimesl

The Normalised AutoEncoder

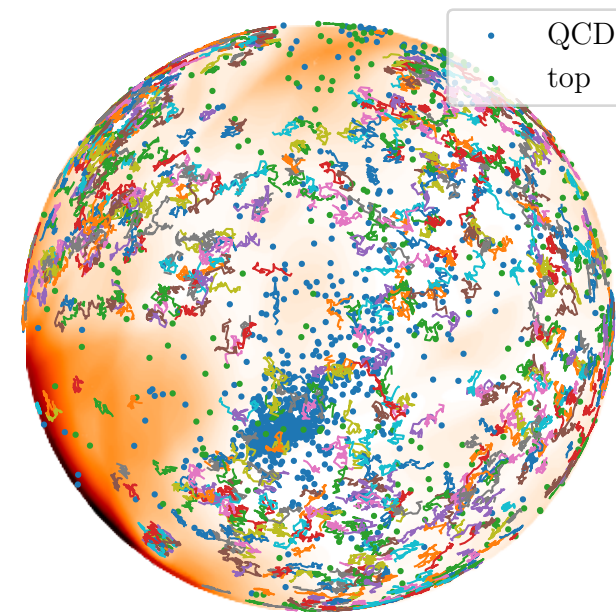
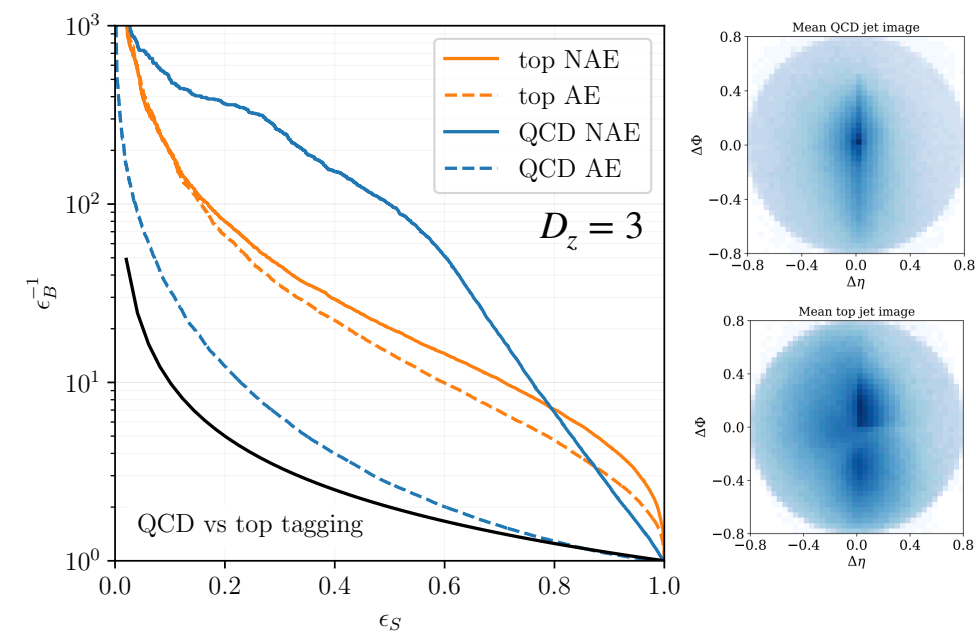
['A normalised autoencoder for LHC triggers' - Dillon, Favaro, Plehn, Sorrenson, Krämer]

No complexity bias!

More robust and reliable anomaly detection

Visualisation of the latent space

Looks like a mess, but very useful for interpreting the results and diagnosing problems with the training!



Barry Dillon

Barry Dillon — Universität Heidelberg — Anomaly searches for new physics at the LHC

Anomaly detection

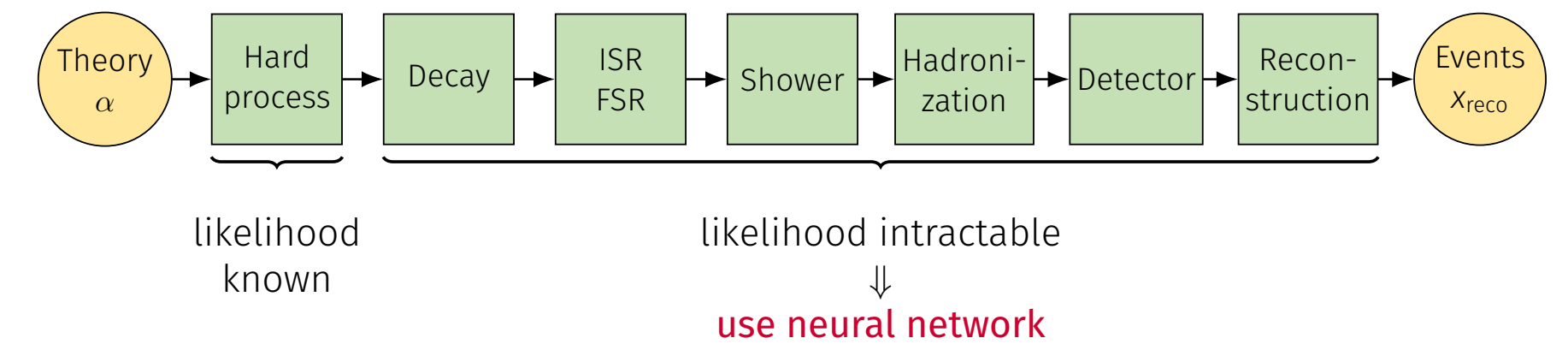
Matrix Element Method

4

- Process with theory parameter α , hard-scattering momenta x_{hard}
- Likelihood at hard-scattering level given by differential cross section

$$p(x_{\text{hard}}|\alpha) = \frac{1}{\sigma(\alpha)} \frac{d\sigma(\alpha)}{dx_{\text{hard}}}$$

- Neyman-Pearson lemma \implies optimal use of information
- Differential cross section only known analytically at hard-scattering level



Beyond Standard Model

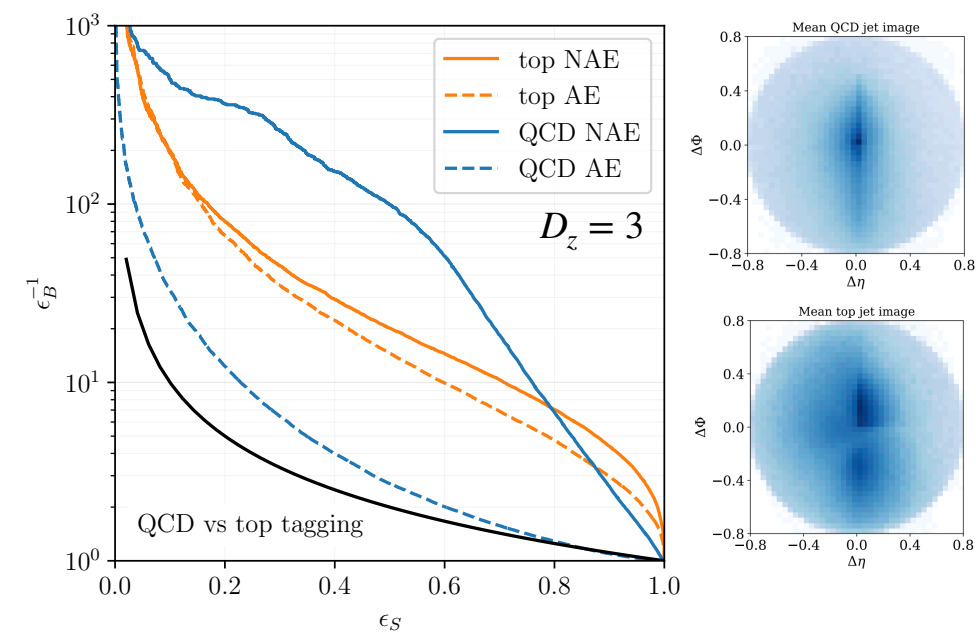
Theo Heime

The Normalised AutoEncoder

['A normalised autoencoder for LHC triggers' - Dillon, Favaro, Plehn, Sorrenson, Krämer]

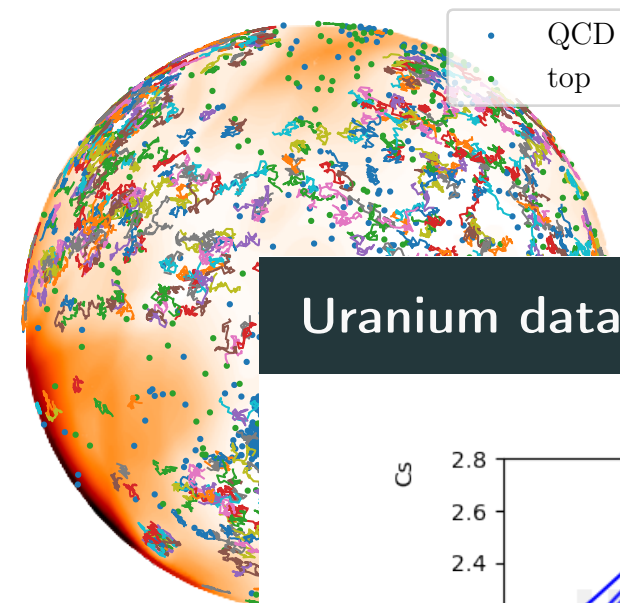
No complexity bias!

More robust and reliable anomaly detection



Visualisation of the latent space

Looks like a mess, but very useful for interpreting the results and diagnosing problems with the training!



Uranium dataset

[Pasquale et al., 2022 in preparation]

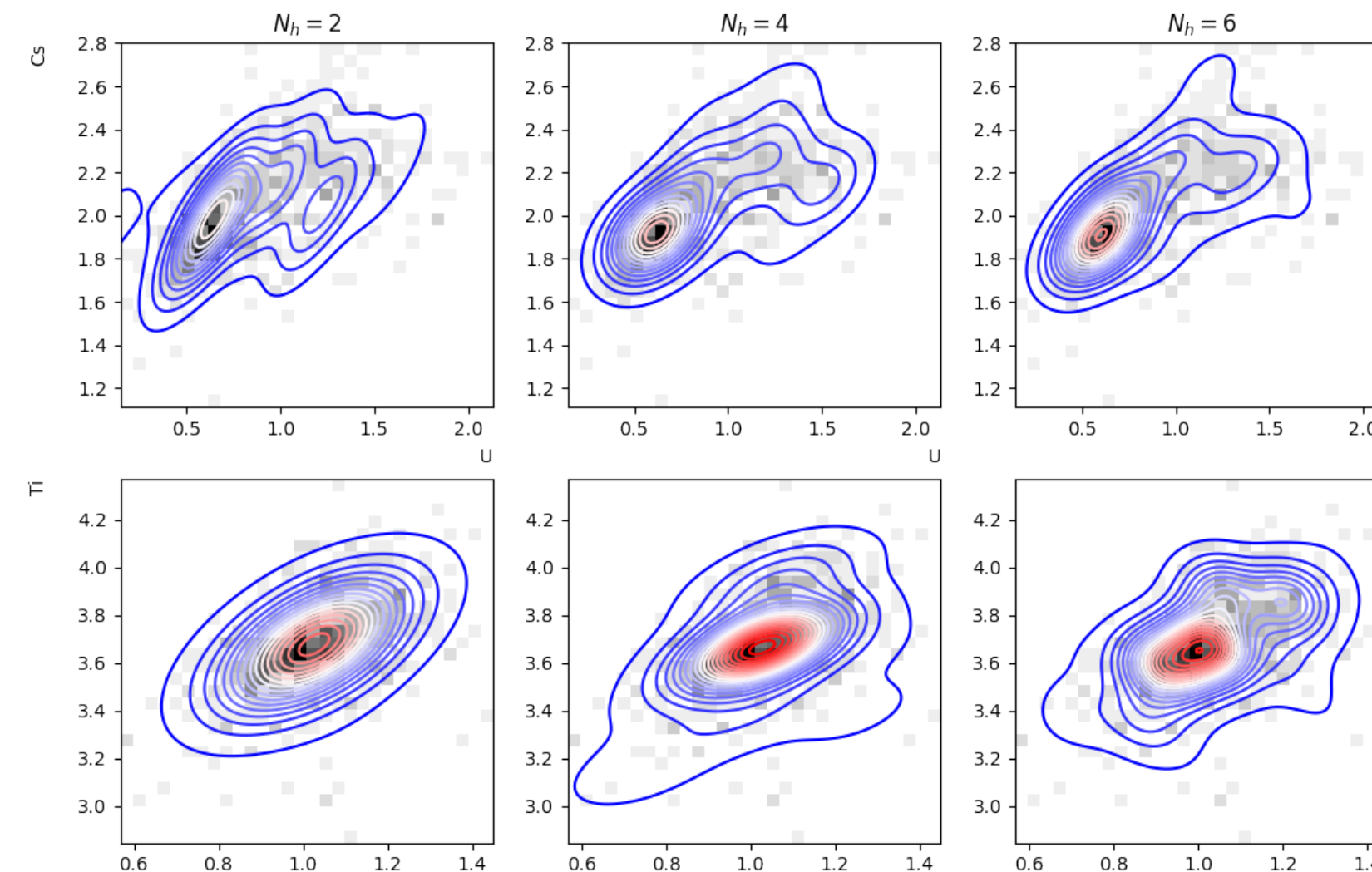


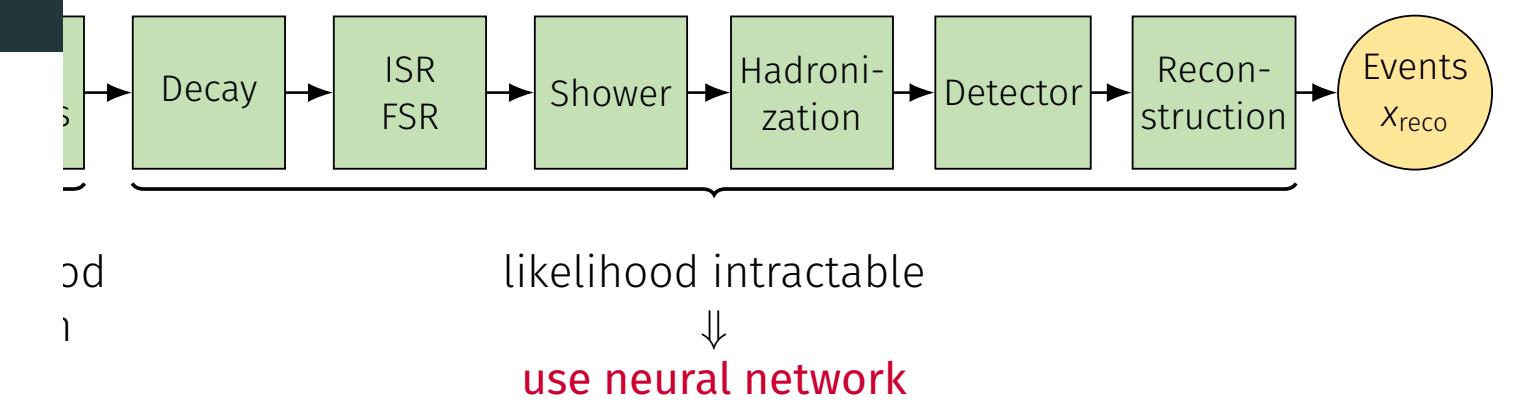
Figure 3: rRTBMs modelling the concentrations of Uranium and Cesium (first row), Cobalt and Titanium (second row) and, Cesium and Scandium (third row) for $N_h = 2, 4, 6$ (left,center,right). The rRTBM contours and histograms of the original data are shown.

Matrix Element Method

- Process with theory parameter α , hard-scattering momenta x_{hard}
- Likelihood at hard-scattering level given by differential cross section

$$p(x_{\text{hard}}|\alpha) = \frac{1}{\sigma(\alpha)} \frac{d\sigma(\alpha)}{dx_{\text{hard}}}$$

- Neyman-Pearson lemma \implies optimal use of information
- Differential cross section only known analytically at hard-scattering level



Barry Dillon

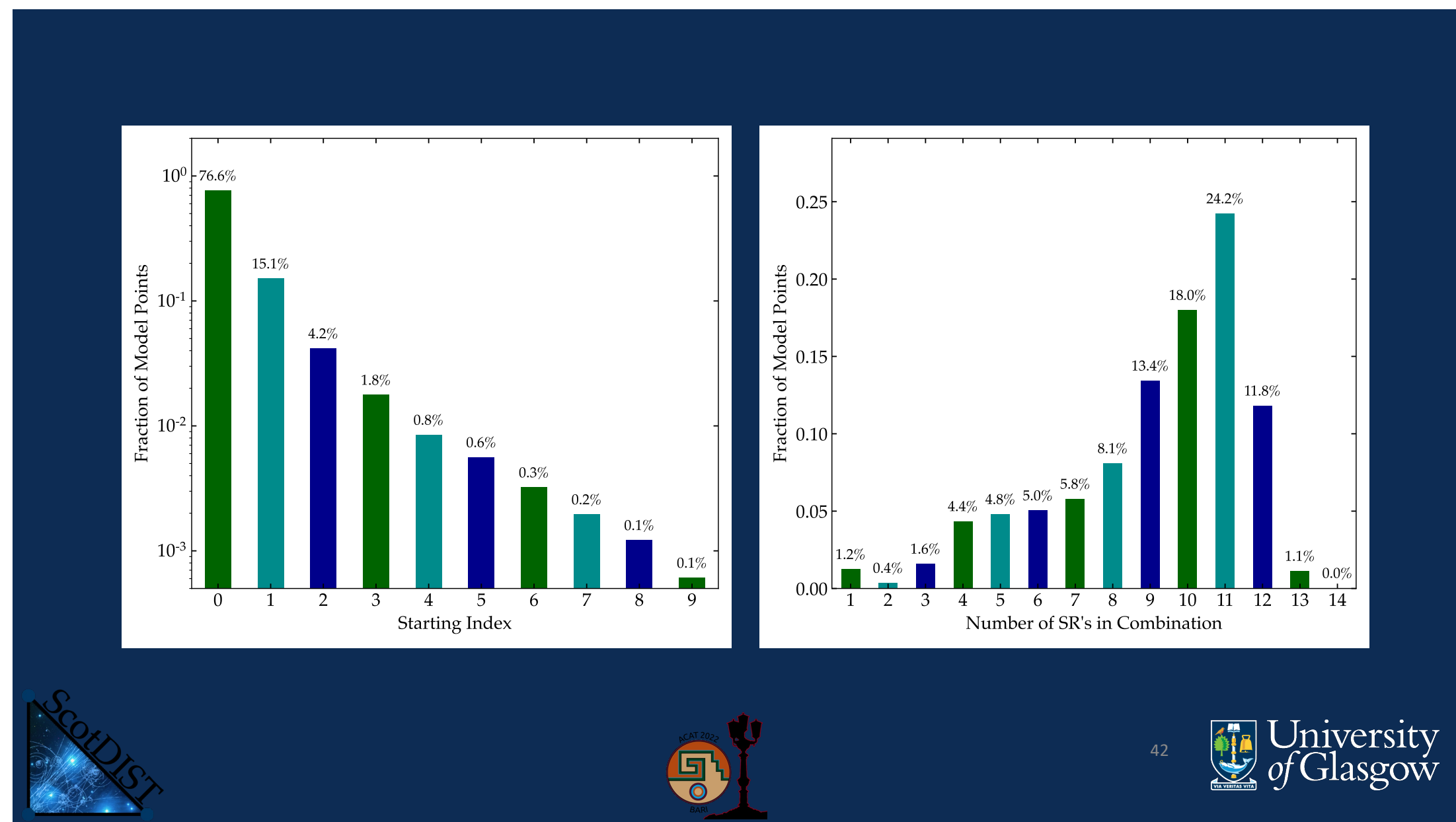
Barry Dillon — Universität Heidelberg — Anomaly searches for new physics at the LHC

Anomaly detection

Andrea Pasquale

Beyond Standard Model

Recasting



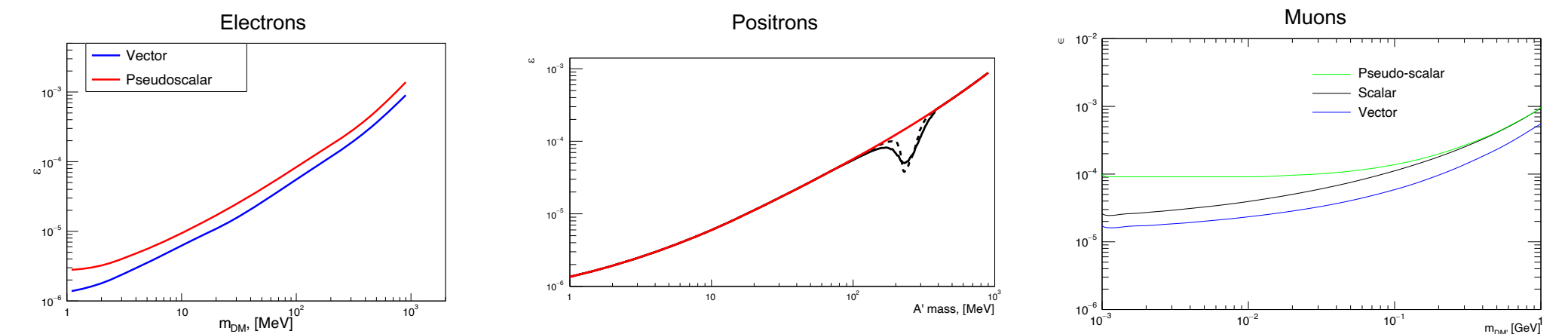
Jamie Yellen

Dark Matter MC event generation

ETH zürich

Beam-dump experiments sensitivity

- Given an experimental set-up and selection cuts, DMG4 enables a realistic study of the **sensitivity** of an experiment with a full set of simulations
- To cope with the extremely low DM production rate, **biasing of the cross-section** can be introduced to generate a reasonable fraction of DM events in the set-up (*BiasSigmaFactor* parameter) → **observe DM** propagation in the set-up and optimize **event selection** (signal to background)
- Final sensitivity is **normalised** w.r.t the biasing factor to retrieve the expected number of DM events



Institute for Particle Physics and Astrophysics (IPA)

Henri Sieber on behalf of the DMG4 team | 10/25/22 | 12

Henri Sieber

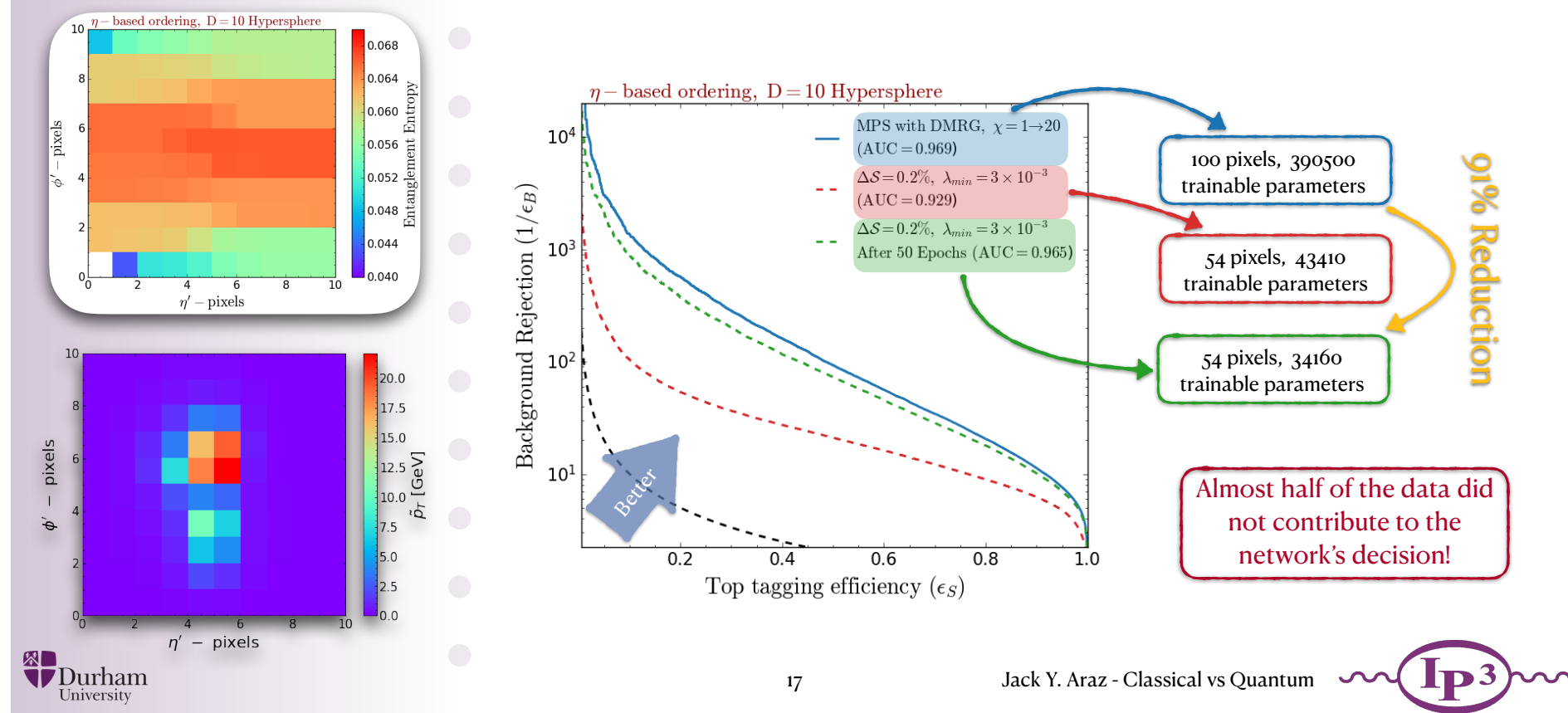
Useful Near-Term Quantum Algorithms

Jungsang Kim

- **Quantum Machine Learning (QML):** Quantum advantages proven for Learning complex patterns w/ quantum feature maps (arXiv:2010.02174)
Exponential gain in predicting certain worst-case error (arXiv:2101.02464)
Quantum correlations used in generative modeling (arXiv:2101.08354)
- **Quantum Chemistry and Materials Studies**
Variational quantum eigensolvers (VQE) for energy estimation
Quantum simulation of dynamics of excitation
Study of quantum many-body phenomena
- **Optimization Problems:** Quantum Approximate Optimization Algorithm

Quantum Machine Learning

Top Tagging through MPS



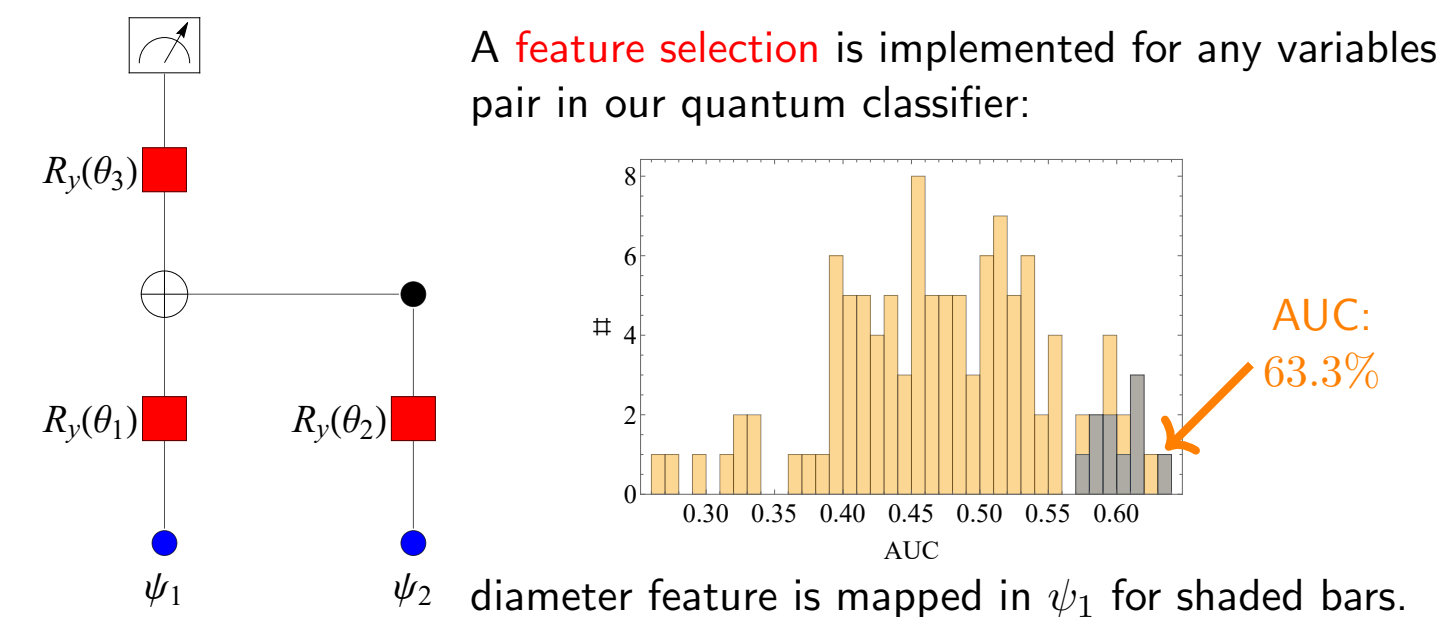
Jack Araz

A hard decision making process

Our best classical (not quantum!) classifier with diameter, grading, histologic type, multifocality, in situ component, PgR:

AUC (%)	Accuracy (%)	Specificity (%)	Sensitivity (%)
70.8 (70.3-71.1)	69.8 (69.3-70.2)	74.8 (72.9-75.1)	61.0 (60.3-61.7)

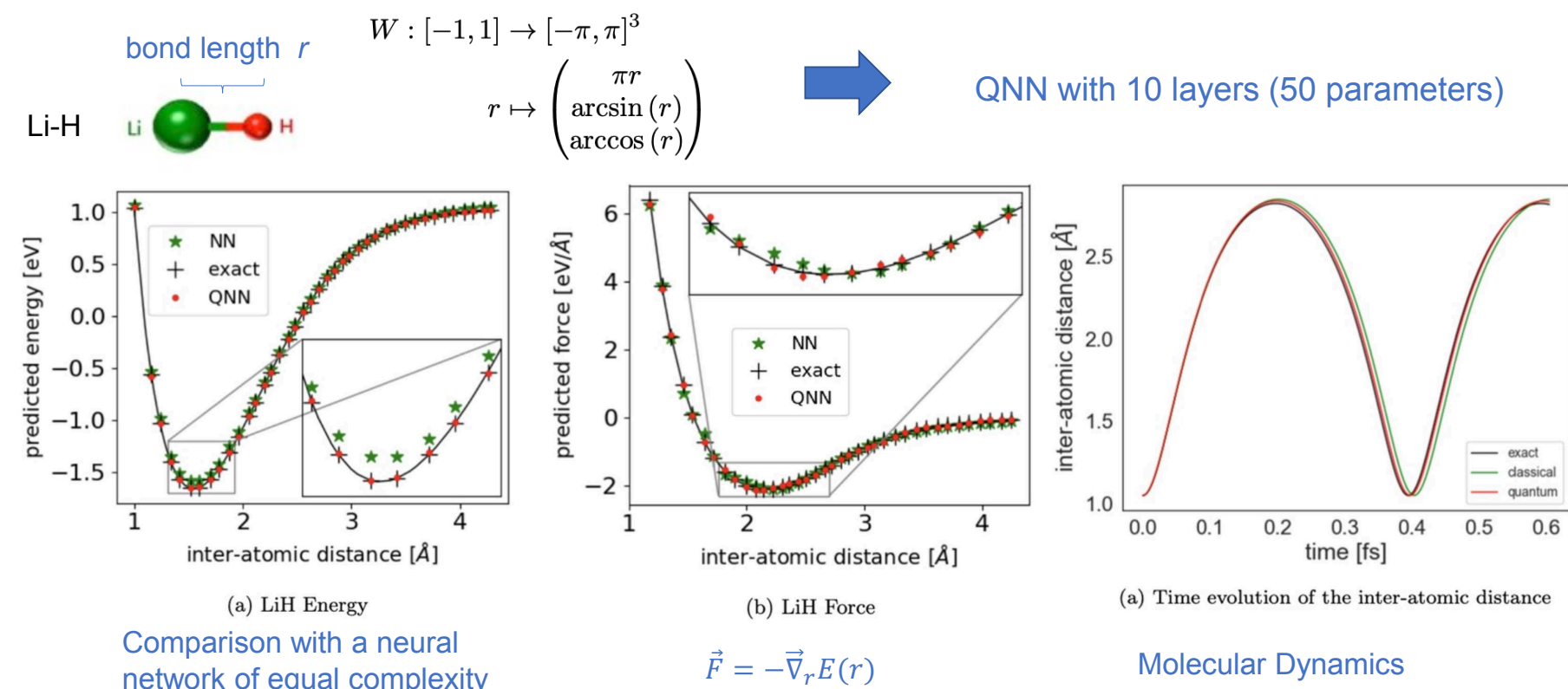
reporting the 1st-3rd interquartile range after 10 ten-fold cross-validations.



Domenico Pomarico

Application to Force Fields (Chemistry)

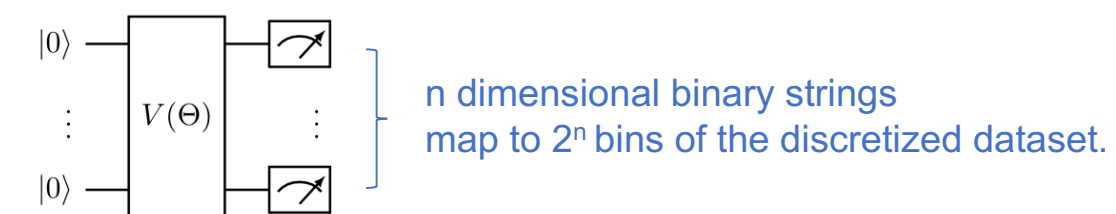
Kiss, Tacchino, et al., *Mach. Learn. Sci. Technol.* 3 035004 (2022)



Oriel Kiss

Quantum Circuit Born machine (QCBM)

1. **Sample** from a variational pure state $|\psi(\theta)\rangle$ by projective measurement with probability given by the **Born rule**: $p_\theta(x) = |\langle x|\psi(\theta)\rangle|^2$.



2. **Training** (Hybrid loop):

- KL divergence Delgado and Hamilton, arXiv:2203.03578 (2022).
- Adversarial (QGAN) Zoufal, et al., *npj Quantum Inf* 5, 103 (2019).
- In the phase space. Kyriienko, et al., arXiv: 2202.08253 (2022).
- Maximum Mean Discrepancy

$$\text{MMD}(P, Q) = \mathbb{E}_{X \sim P} [K(X, Y)] + \mathbb{E}_{Y \sim Q} [K(X, Y)] - 2\mathbb{E}_{X \sim P} [K(X, Y)]$$

3. **Why the Maximum Mean Discrepancy MMD ?**

- Resource efficient for NISQ devices.
- Stable.
- However, empirically less performant than adversarial.

Michele Grossi

Quantum chemistry and fluids

Variational quantum eigensolver

Hardware results (IBMQ)

- Start from the classical solution (warm start).
- Qubit Based Excitation descending UCC Ansatz.
- 10 runs on the 27 qubits IBMQ mumbai chip.

Error mitigation

- **Readout:** individually invert the error matrices

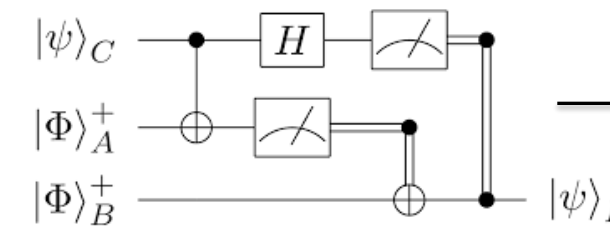
$$S_k = \begin{pmatrix} P_{0,0}^{(k)} & P_{0,1}^{(k)} \\ P_{1,0}^{(k)} & P_{1,1}^{(k)} \end{pmatrix}.$$

- Here, $P_{i,j}^{(k)}$ is the probability of the k -th qubit to be in state $j \in \{0,1\}$ while measured in state $i \in \{0,1\}$.

hardware	No. parameters	No. CNOT	mean	st. deviation	exact	error ratio
ibmq_mumbai raw (g.s.)	9	209	-6.27	0.269	-5.529	13.36%
ibmq_mumbai mitigated (g.s.)	9	209	-5.319	0.24	-5.529	3.81%
ibmq_mumbai raw (1st es)	3	41	-2.907	0.87	-3.420	14.97%
ibmq_mumbai mitigated (1st es)	3	41	-3.424	0.08	-3.420	0.12%

NEXT STEPS

- Build Quantum Circuit for the Collision and Streaming of qLBM
- Implement the Quantum Circuit in the Intel Quantum SDK
- Finally, Solve a simple Fluid Dynamics problem using this Circuit.
- Validate the results.



```

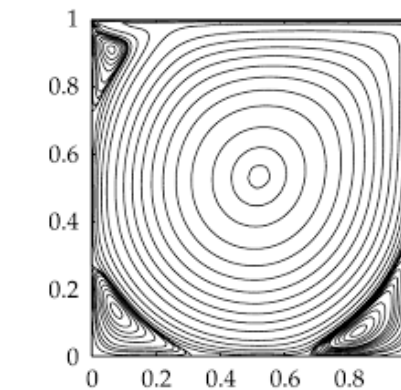
const int N = 5;
/* global array of qubits */
qbit qreg[N];
qbit creg[N];

quantum_kernel void cphase(qbit ctr, qbit tgt, double phase) {
  RZ(ctr, phase / 2);
  CNOT(ctr, tgt);
  RZ(tgt, -1 * phase / 2);
  CNOT(ctr, tgt);
  RZ(tgt, phase / 2);
}

quantum_kernel void prepare() {
  for (int i = 0; i < N; i++) {
    PrepZ(qreg[i]);
  }
}

quantum_kernel void measure() {
  for (int i = 0; i < N; i++) {
    MeasZ(qreg[i], creg[i]);
  }
}

quantum_kernel void QFT() {
  int n = N - 1;
  
```



Oriel Kiss

Tejas Shinde

Optimisation

Quantum annealers

Frameworks

Juan Carlos Criado

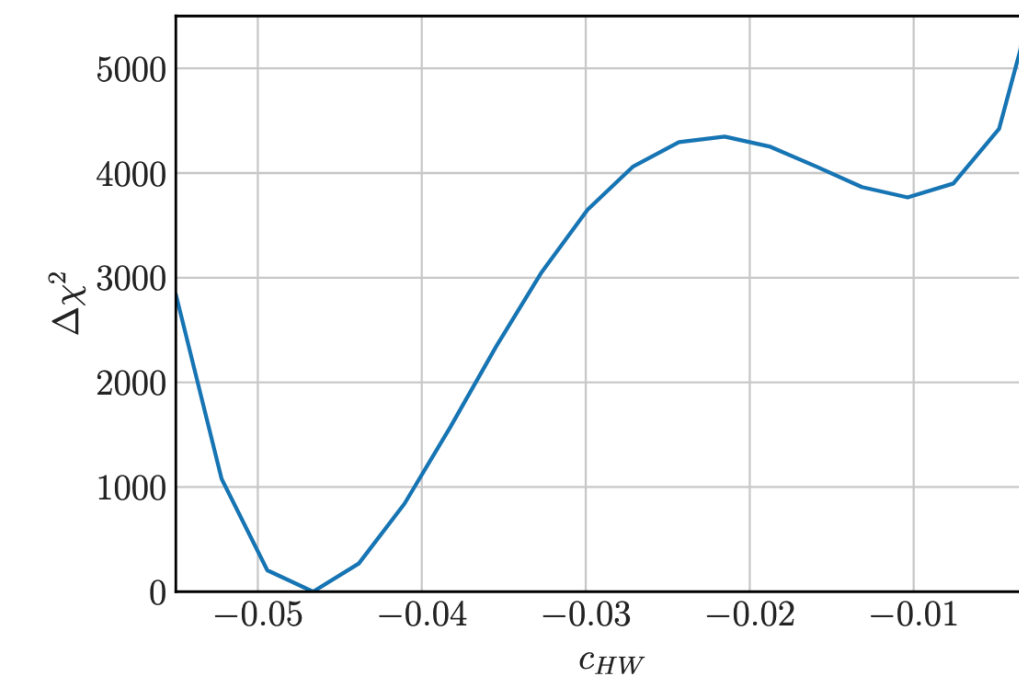
QFitter: EFT Wilson coefficient fits

2207.10088

$$\chi^2 = \sum_{ij} V_a C_{ab}^{-1} V_b, \quad V_a = O_a^{(\text{exp})} - O_a^{(\text{th})}(c)$$

$$O_a^{(\text{th})}(c) = A_a + \sum_i B_{ai} c_i + \sum_{ij} C_{aij} c_i c_j$$

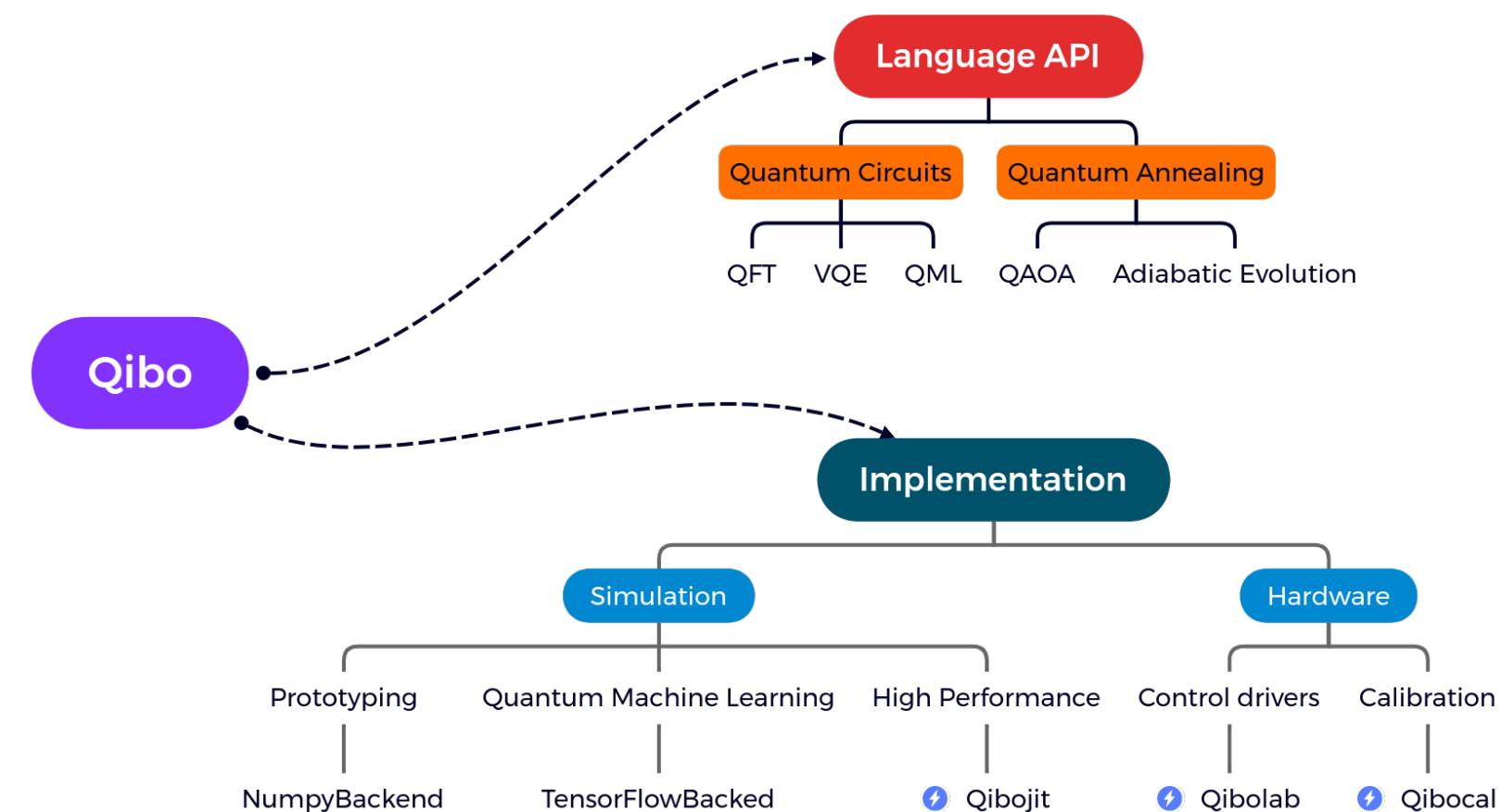
$$\begin{aligned} \mathcal{L} = & \frac{c_{u3yt}}{v^2} (\phi^\dagger \phi) (\bar{q}_L \tilde{\phi} u_R) + \frac{c_{d3yb}}{v^2} (\phi^\dagger \phi) (\bar{q}_L \phi d_R) \\ & + \frac{i c_W g}{2 m_W^2} (\phi^\dagger \sigma^a D^\mu \phi) D^\nu W_{\mu\nu}^a + \frac{c_H}{4 v^2} (\partial_\mu (\phi^\dagger \phi))^2 \\ & + \frac{c_\gamma (g')^2}{2 m_W^2} (\phi^\dagger \phi) B_{\mu\nu} B^{\mu\nu} + \frac{c_g g_S^2}{2 m_W^2} (\phi^\dagger \phi) G_{\mu\nu}^a G^{a\mu\nu} \\ & + \frac{i c_{HW} g}{4 m_W^2} (\phi^\dagger \sigma^a D^\mu \phi) D^\nu W_{\mu\nu}^a \\ & + \frac{i c_{HB} g'}{4 m_W^2} (\phi^\dagger D^\mu \phi) D^\nu B_{\mu\nu} + \text{h.c.} \end{aligned}$$



Qibo

Qibo is an **open-source** full stack API for quantum simulation and quantum hardware control and calibration.

Andrea
Pasquale



<https://github.com/qiboteam/qibo>

The future

- Faster, more precise calculations and event generation
- New physics: model independent searches
- Towards Quantum Computing

There are a lot of exciting ideas for us. Happy coding!

Anke Biekötter - Leonardo Cosmai -
Joshua Davies - Latifa Elouadrhiri

