

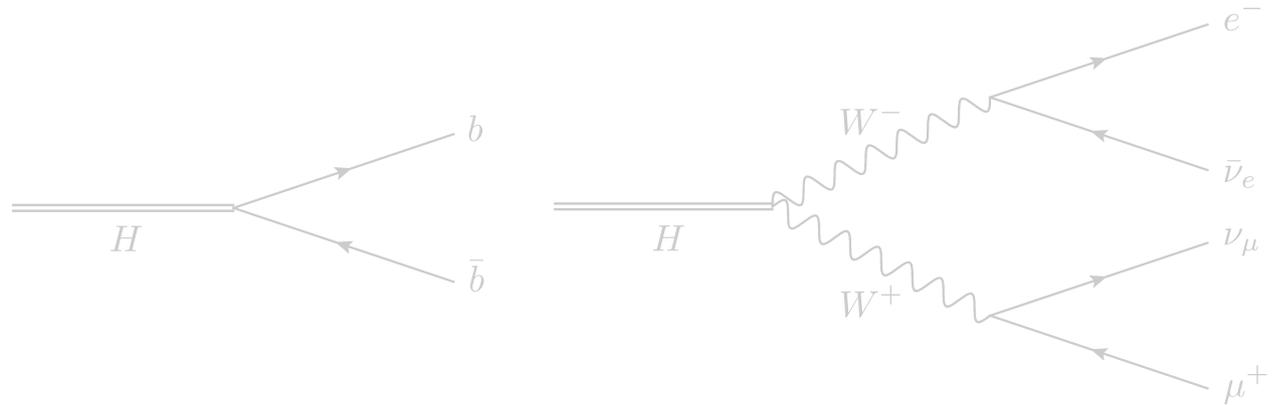
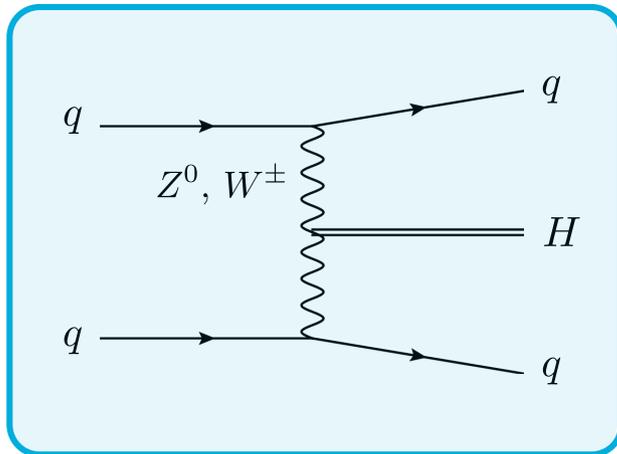


Towards Higgs boson production in weak boson fusion at higher orders with realistic final states

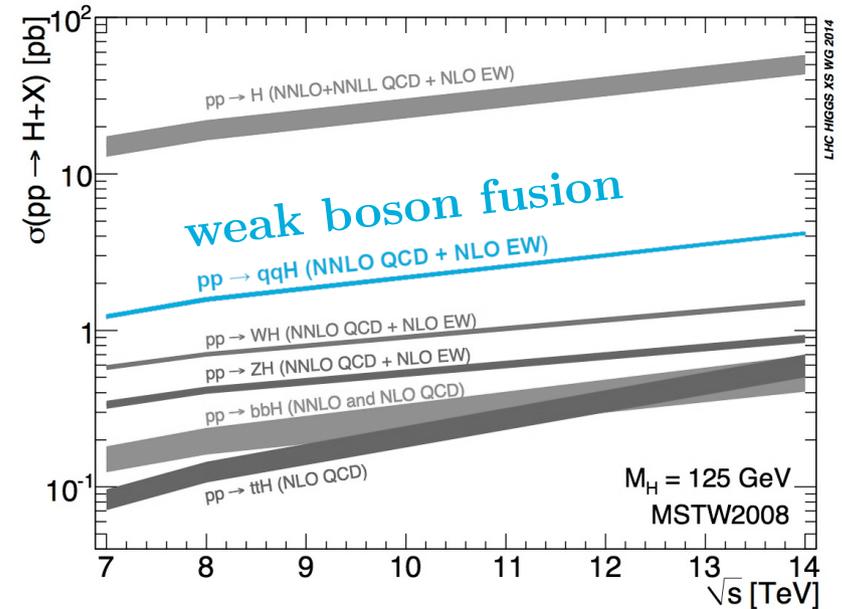
In collaboration with Fabrizio Caola, Kirill Melnikov and Raoul Röntsch
Details in arXiv:2110.02818 / JHEP02(2022)046

Konstantin Asteriadis | 05/12/2022
LoopFest XX, Pittsburgh

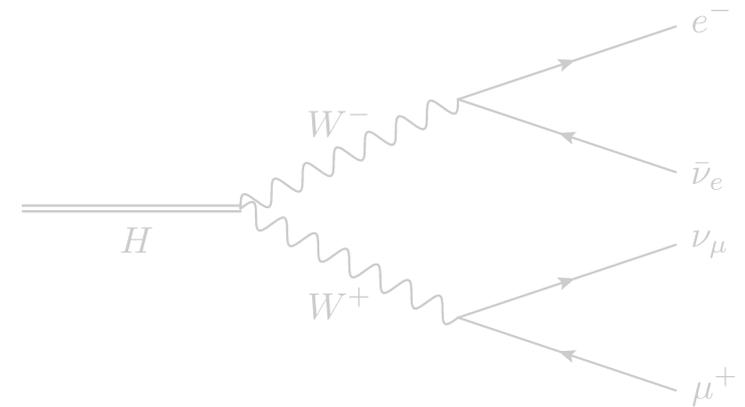
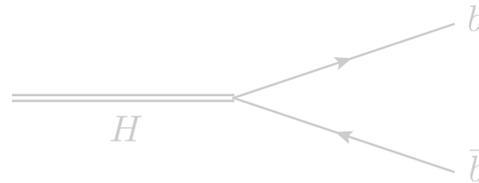
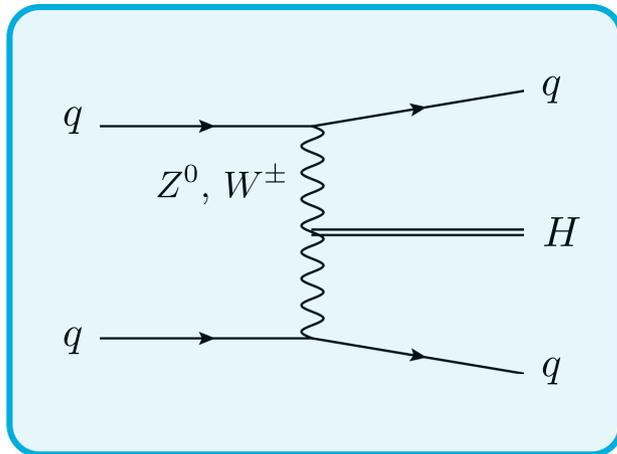
Higgs boson production in vector boson fusion



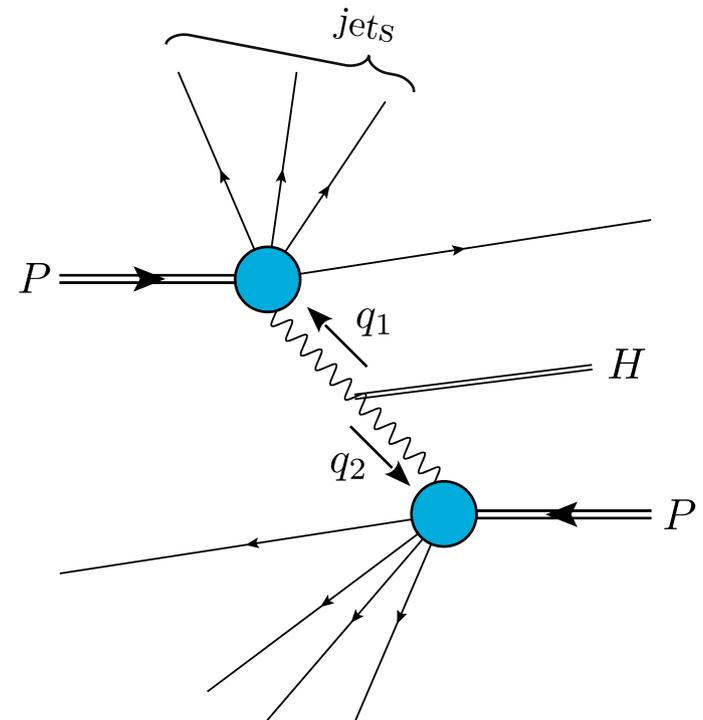
- Production channel with the highest cross section that probes electroweak couplings of the Higgs boson
- Fiducial region: non-trivial jet structure
[Rauch, Zeppenfeld '17; ...] \rightarrow important to control radiative corrections
- State of the art:
 - Inclusive @ $N^3\text{LO}$ QCD [Dreyer, Karlberg '16];
 - Fully differential @ NNLO QCD [Cacciari, Dreyer, Karlberg, Salam, Zanderighi '15; Cruz-Martinez, Glover, Gehrmann, Huss '18, Caola, Melnikov, Roentsch KA '21]



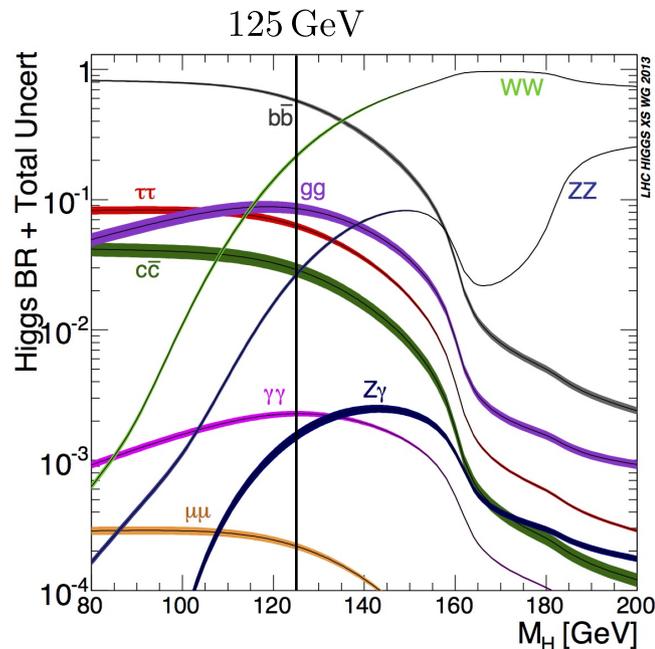
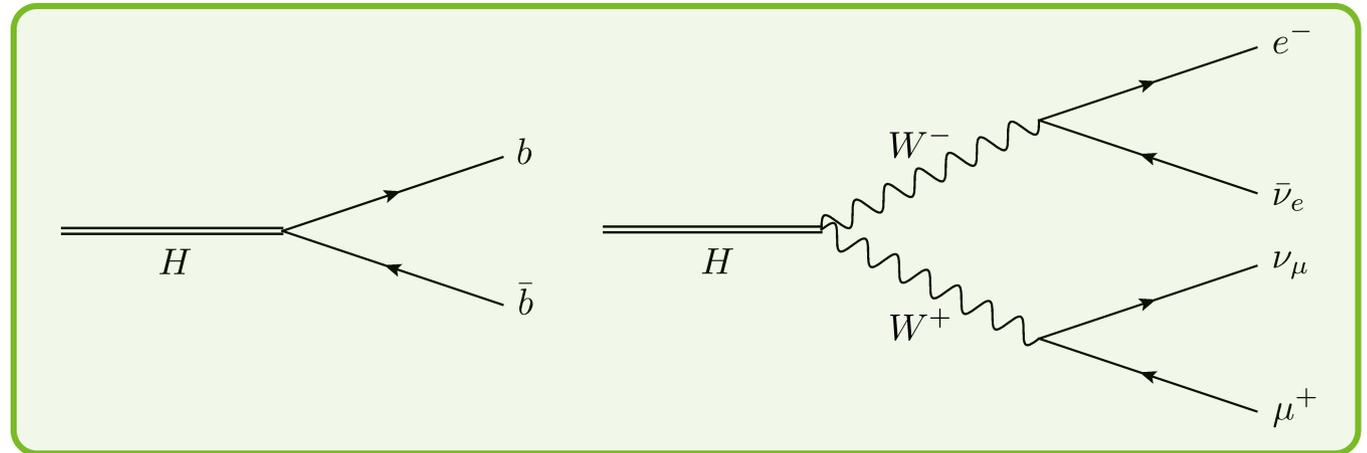
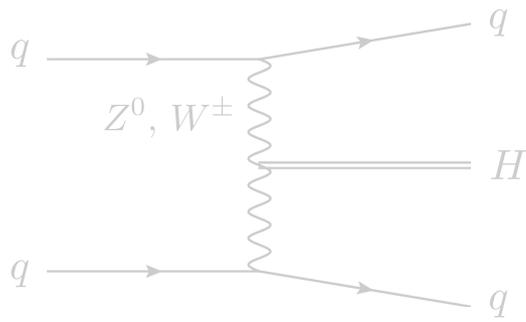
A closer look at QCD corrections



- Largest contribution from **factorizable** corrections [Liu, Melnikov, Penin '19] → focus of this talk
- Effectively DIS scattering of two **protons**
- **However: All current computations are for stable Higgs boson production**

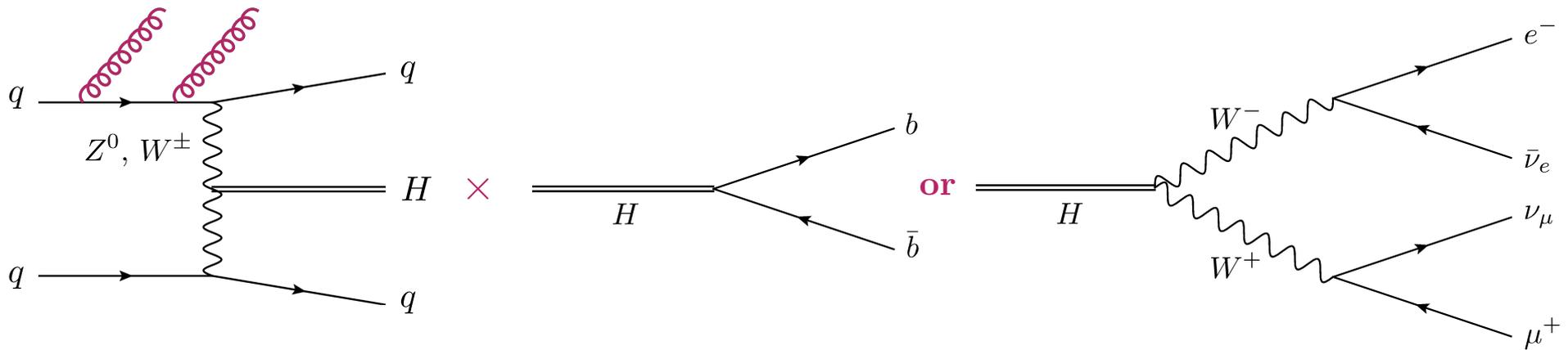


Detecting WBF through realistic final states



- $H \rightarrow b\bar{b}$ and $H \rightarrow WW^* \rightarrow 2l 2\nu$
- Highest branching ratios
- Both studied by ATLAS and CMS
[e.g. Eur. Phys. J. C 81, 537 (2021); Phys. Lett. B 791, 96 (2019)]
- **Doing this at NNLO QCD naively simple, in practice very complicated**

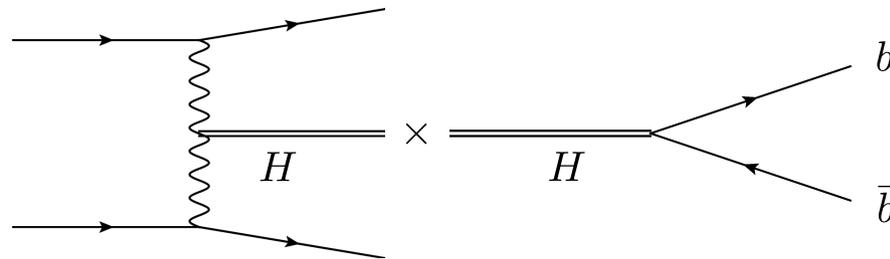
NNLO QCD Higgs boson production + Higgs boson decay



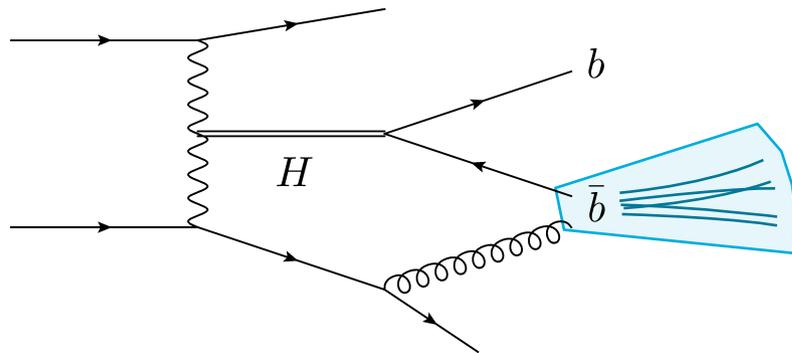
- In this combination, each decay channel comes with its unique challenges:
 - $H \rightarrow b\bar{b}$: non-trivial interplay between partonic jets from production and decay when fiducial cuts are applied
 - $H \rightarrow WW^* \rightarrow 2l 2\nu$: up to 21 dimensional phase space integration that is numerically very challenging
- In what follows: focus on $H \rightarrow b\bar{b}$ decay channel (details on $H \rightarrow WW$ in *JHEP02(2022)046*)
- Consider $H \rightarrow b\bar{b}$ decay only at LO QCD with massless b quarks
- (Side note: Good control on complex final state coming from decay crucial for computing radiative corrections to $H \rightarrow b\bar{b}$ decay channel)

WBF + H \rightarrow $b\bar{b}$ decay

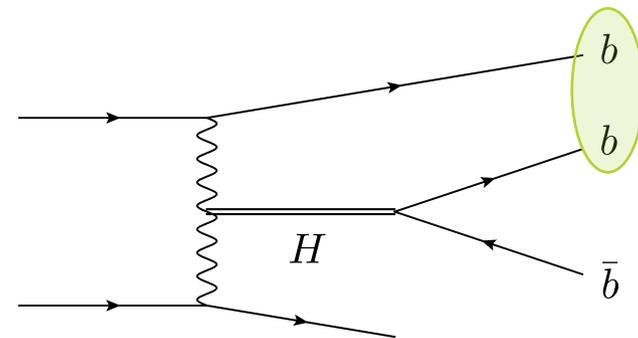
- Narrow width approximation \rightarrow factorization of on-shell Higgs production and on-shell Higgs decay



- Several effects break factorization of production and decay process. For example



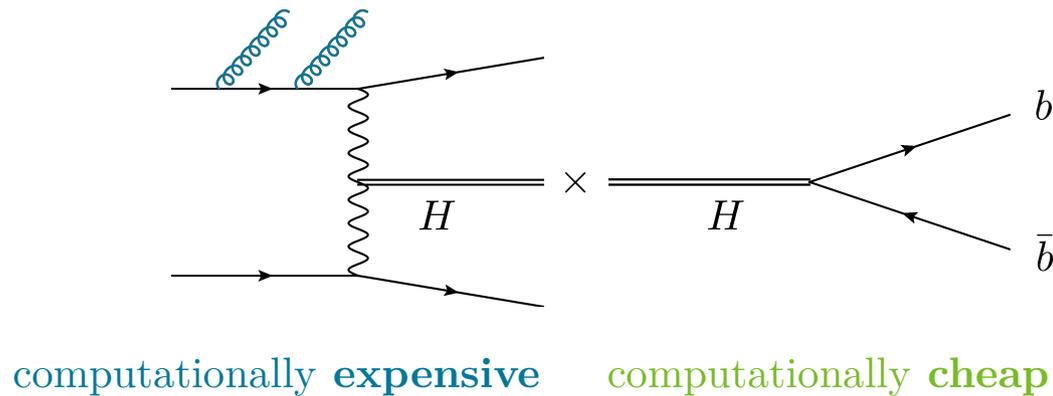
Jet-clustering breaks factorization



B-tagging breaks factorization

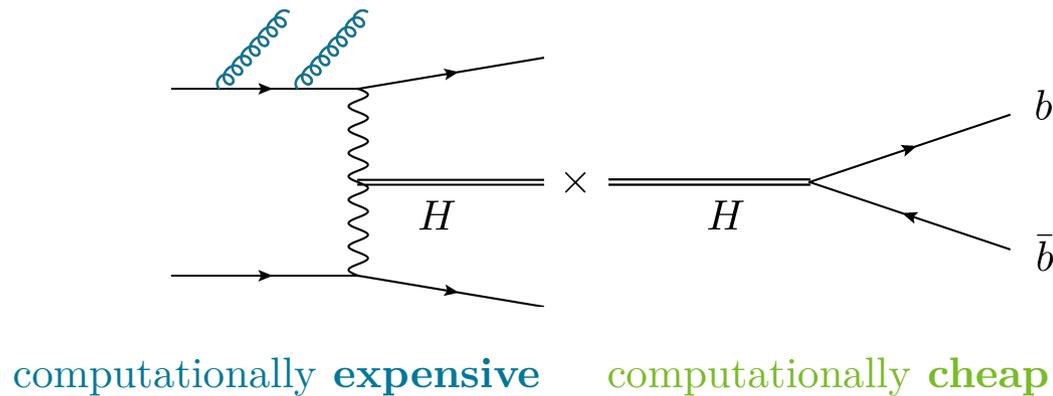
- We don't expect these effects to be very large but it is important to quantify their size
- Finally: cuts on b-jets may change fiducial WBF region

The technical challenge: phase space integration



- Naive approach: simply enlarge the phase space dimension and let the MC integrator do its job
- However, extremely inefficient: it does not take the peculiarities of the situation into account:
 - Huge difference in complexity of production and decay kinematics
 - **Production**: NNLO radiation with already non-trivial Born kinematics
 - **$H \rightarrow b\bar{b}$ @ LO decay**: back-to-back and equally distributed over the angular phase space
- ($H \rightarrow WW$ decay has more complicated kinematics but still far simpler than complexity of WBF production @ NNLO)

The technical challenge: phase space integration

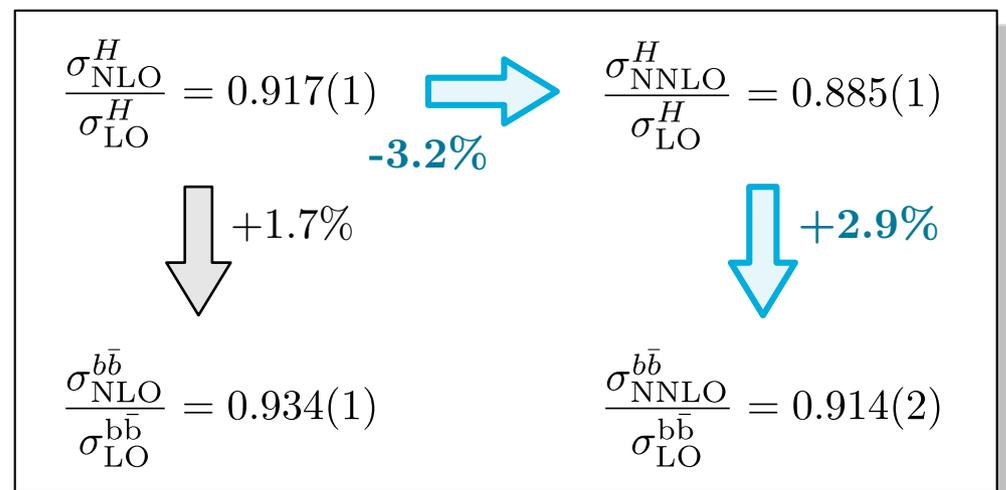
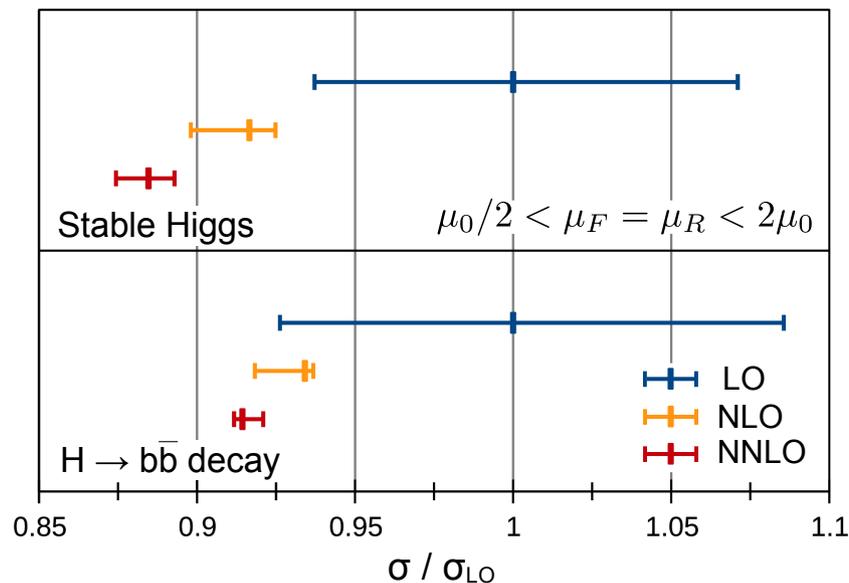


- To take advantage of this situation:
 - 1) Generate high-precision integration grids for stable Higgs boson production
 - 2) Generate production events according to these grids and for each of them generate $O(100)$ decay events
 - 3) When this fails (typically in regions where the NNLO cross section receives large positive and negative contributions) \rightarrow naive “brute-force” integration ($\sim 15\%$ of the time)
- However, $\sim 85\%$ of the time the computation can be done relatively cheaply

Results: fiducial cross section

- Cuts on b-jets; loosely following latest ATLAS measurement: 2 resolved b-jets; $p_{\perp, \text{jb}} > 65 \text{ GeV}$; $|y_{\text{jb}}| < 2.5$ [Eur. Phys. J. C 81, 537 (2021)]
- Cross section for b-jets from production process small \rightarrow at this stage: no flavour tagging in Higgs production process
- Sizable fiducial cross section, $O(100\ 000)$ events with HL-LHC

$$\sigma_{\text{LO}}^{b\bar{b}} = 75.9_{-5.6}^{+6.5} \text{ fb}, \quad \sigma_{\text{NLO}}^{b\bar{b}} = 70.9_{-1.2}^{+0.2} \text{ fb}, \quad \sigma_{\text{NNLO}}^{b\bar{b}} = 69.4_{-0.2}^{+0.5} \text{ fb}$$

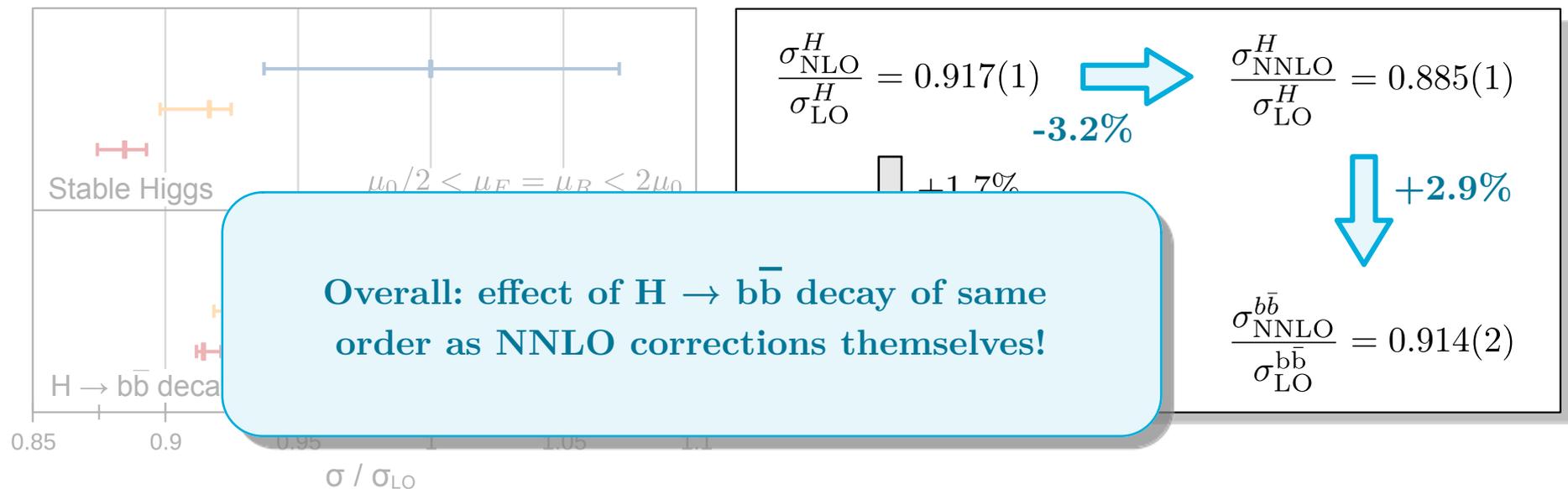


- *Noteworthy features:* smaller residual scale uncertainty and better perturbative convergence compared to stable Higgs production

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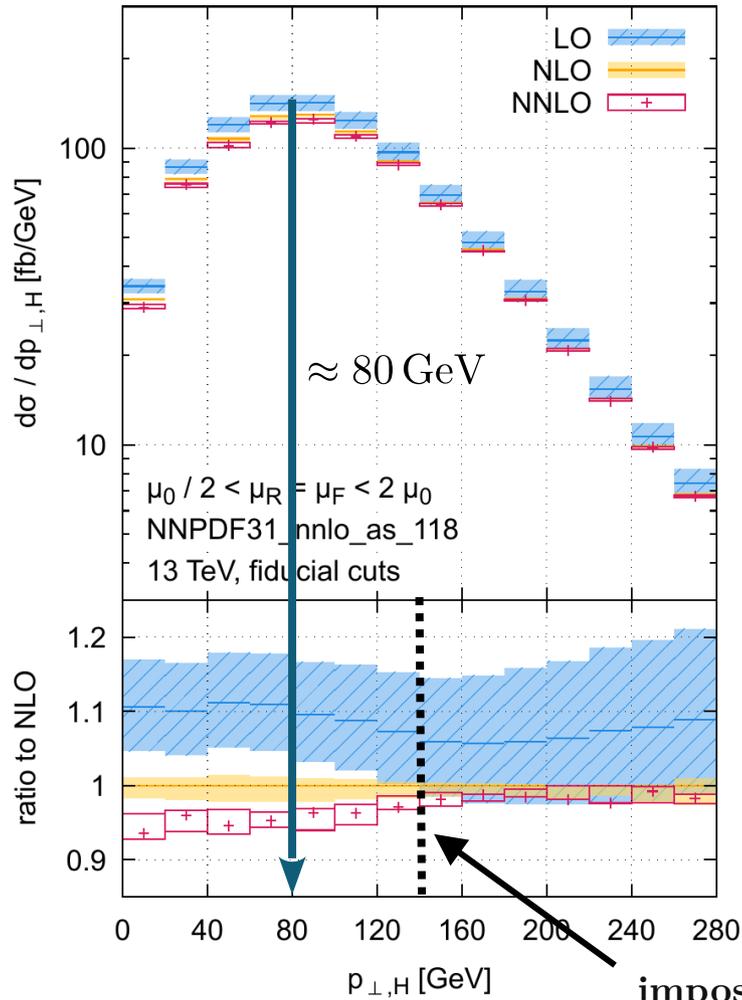
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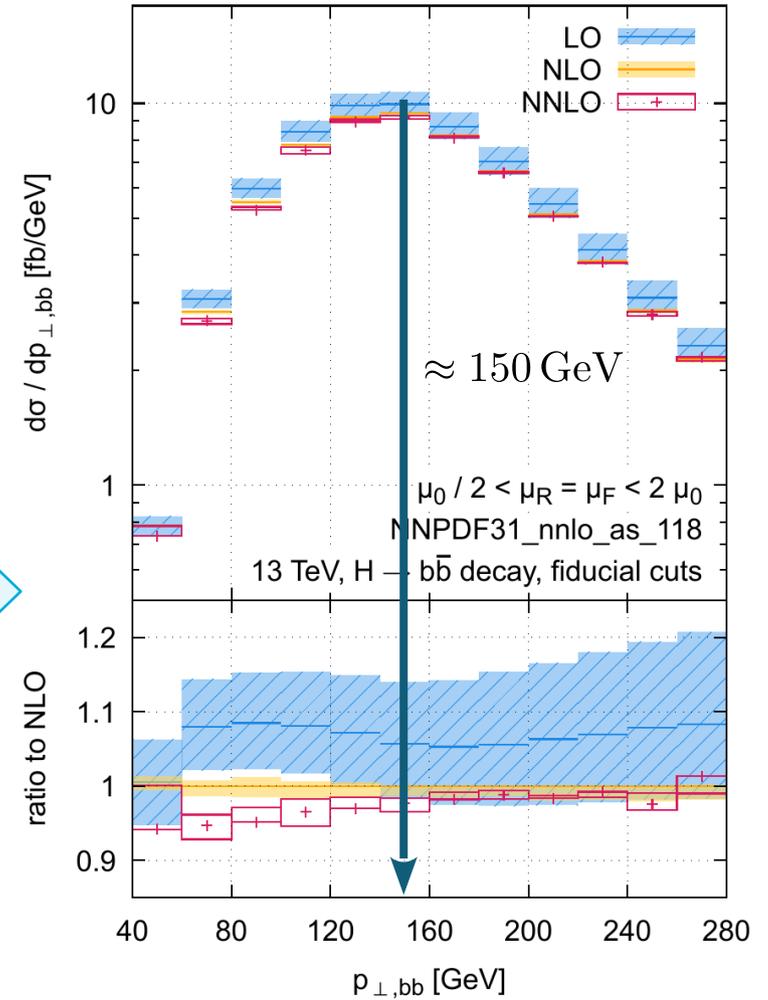
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Results: fiducial cross section

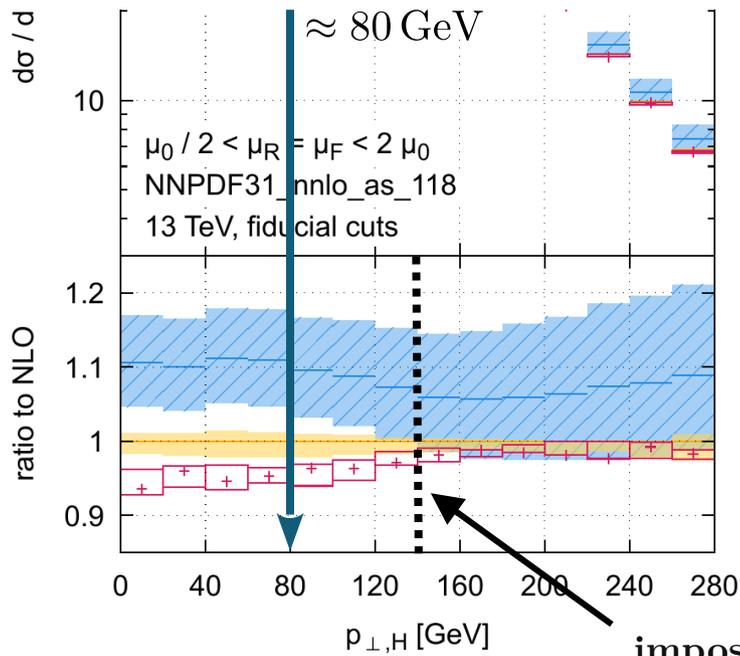
- Simple reason: pt cuts on b-jets ($p_{\perp,j_b} > 65 \text{ GeV}$) preferentially selects events with high Higgs transverse momentum



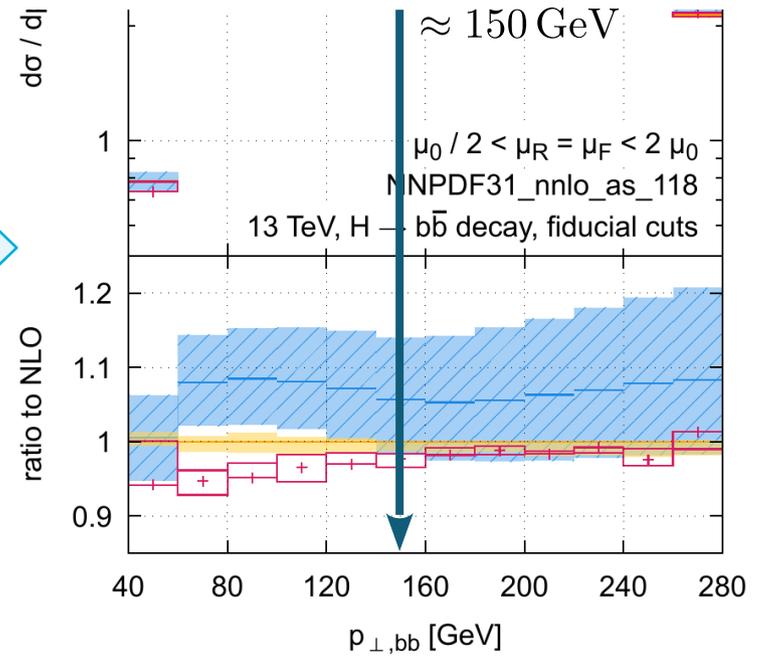
including decay



$$2 p_{\perp,j_b}^{\min} = 130 \text{ GeV}$$



including decay



imposed “soft” pt cut
 $2 p_{\perp, j_b}^{\min} = 130 \text{ GeV}$

- NLO corrections are rather flat → moderate effect
- For $p_t > 130 \text{ GeV}$ NNLO corrections are smaller and within residual scale uncertainty band
- Check: Stable Higgs production with additional pt cut $p_{\perp, H} > 150 \text{ GeV}$

$$\frac{\sigma_{\text{NNLO}}^H}{\sigma_{\text{LO}}^H} = 0.89$$

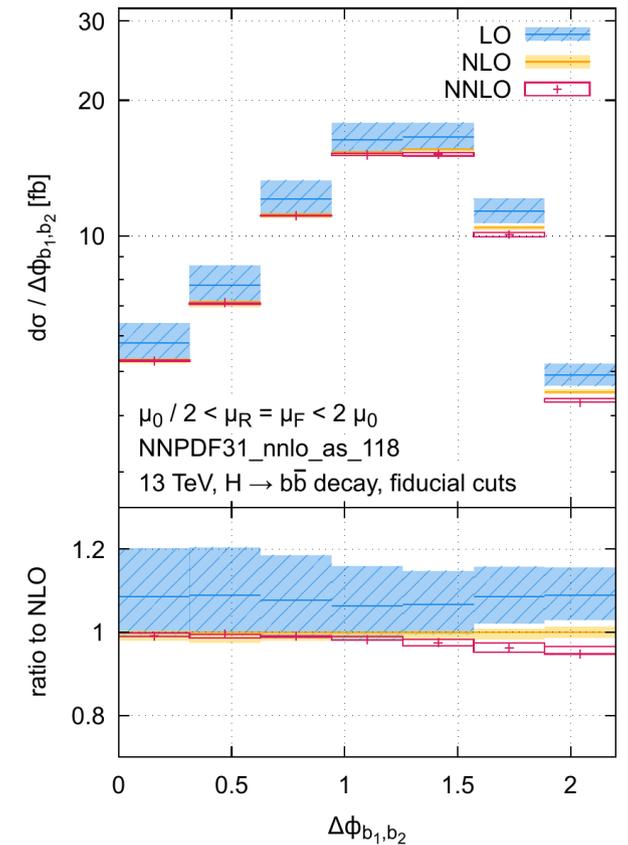
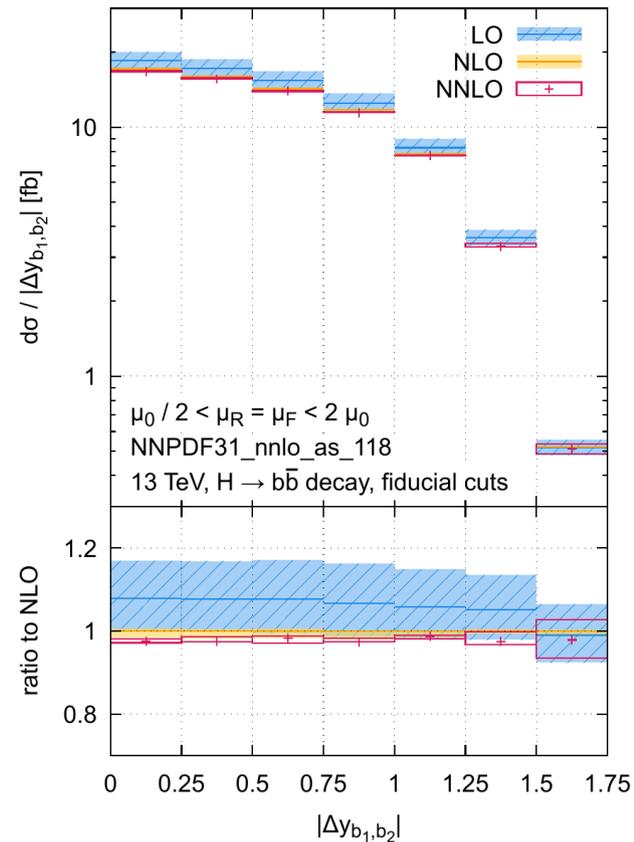
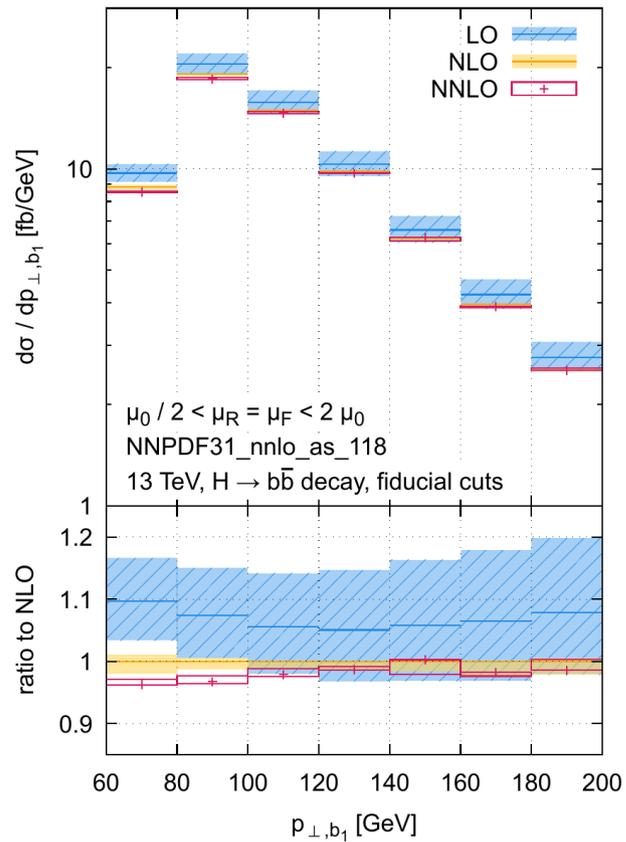
Higgs pt cut

$$\frac{\sigma_{\text{NNLO}}^H}{\sigma_{\text{LO}}^H} = 0.91$$

including decay

$$\frac{\sigma_{\text{NNLO}}^{b\bar{b}}}{\sigma_{\text{LO}}^{b\bar{b}}} = 0.914(2)$$

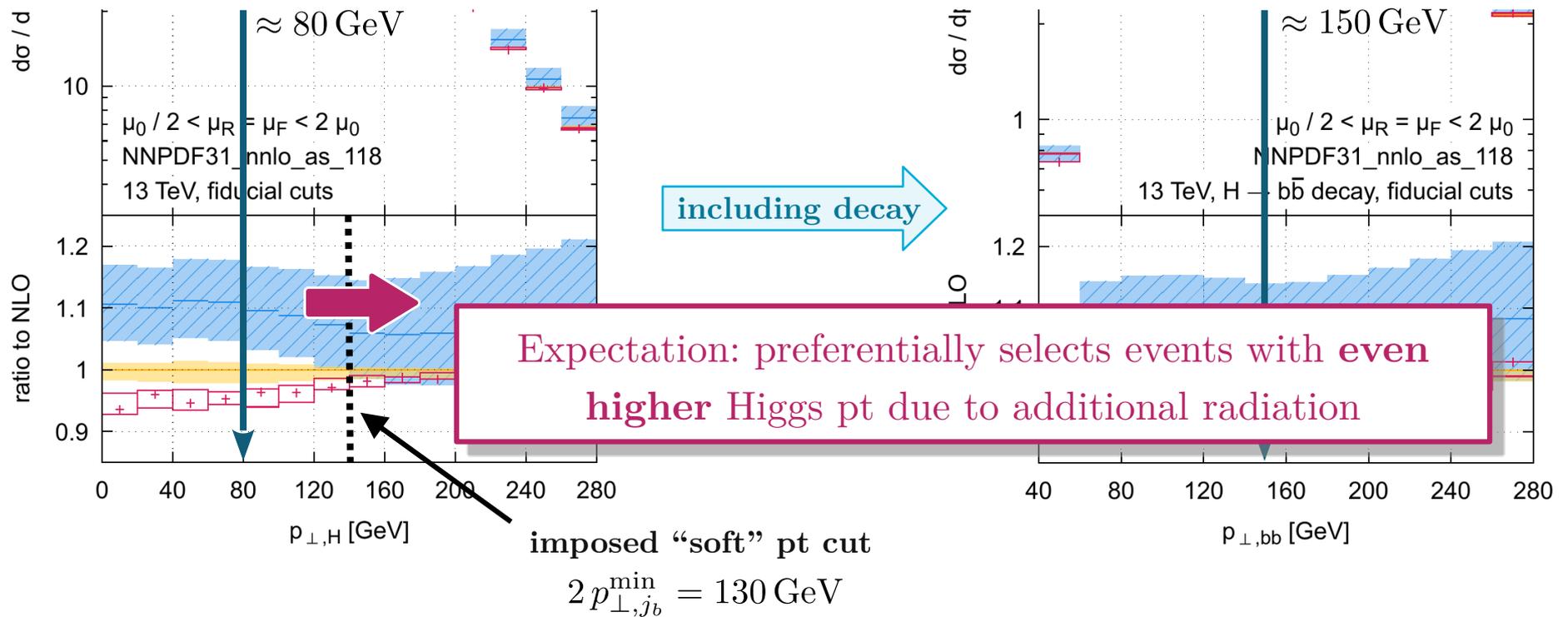
Results: differential cross sections



- Shapes of NLO distributions **not affected** by NNLO corrections
- Simple reweighting possible as long as NNLO/NLO K-factor is computed with a proper cut on the p_t of the stable Higgs boson

Outlook: Towards a more realistic setup

- $H \rightarrow b\bar{b}$ @LO (and $H \rightarrow WW \rightarrow 2l 2\nu$) as prototypes for $H \rightarrow b\bar{b}$ @ NNLO QCD
- Fully-differential description of $H \rightarrow b\bar{b}$ decay at NNLO QCD (with massive b-quarks) is known
[Bernreuther, Chen, Si '2018; Behring, Bizoń '19]
- Add flavor tagging in WBF Higgs boson production process

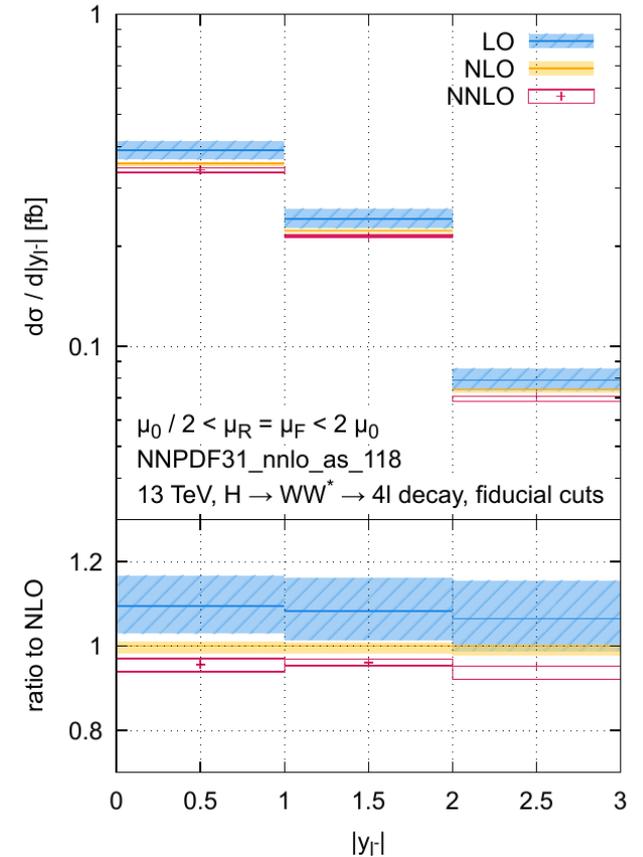
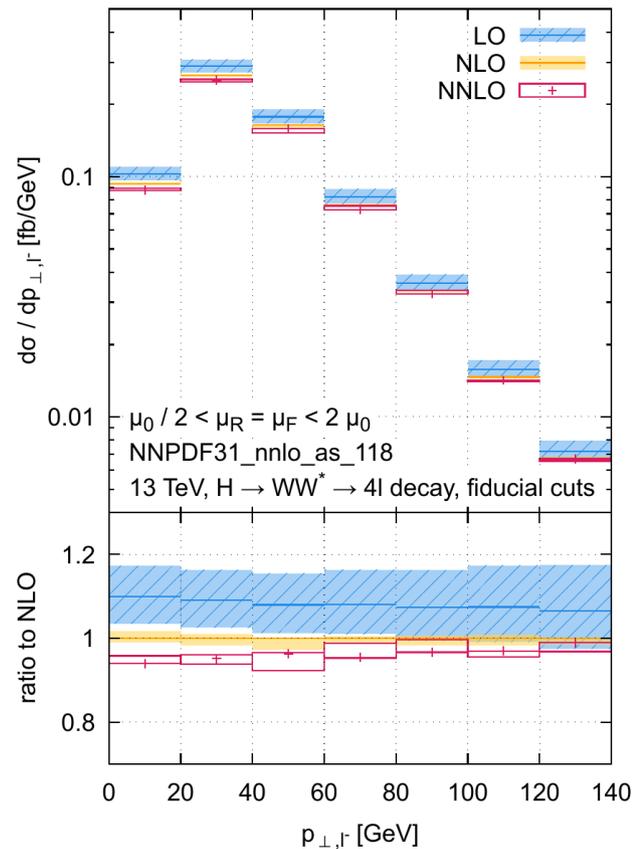
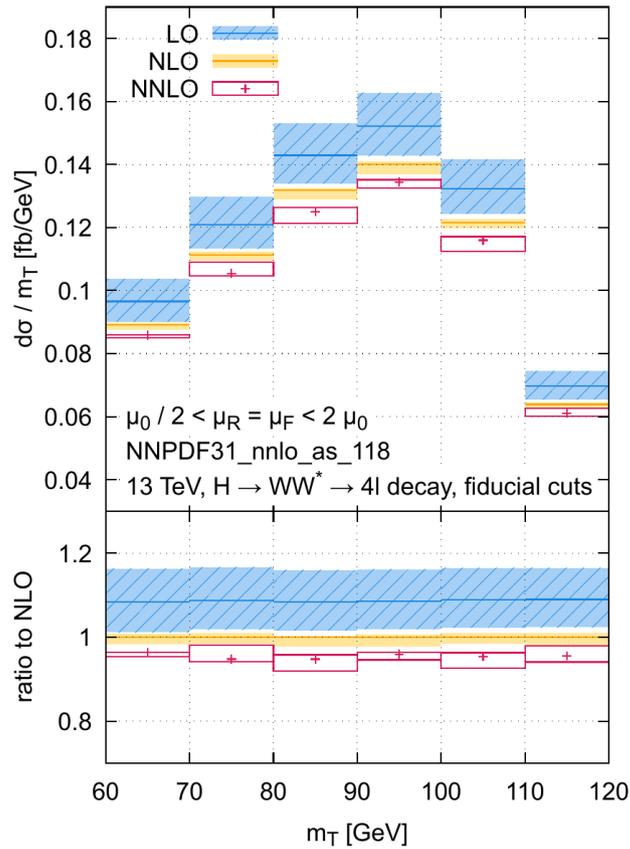


Conclusion

- First NNLO QCD study of Higgs boson production in vector boson fusion that includes the decay of the Higgs boson
- Naively trivial; challenging in practice
- **WBF including $H \rightarrow b\bar{b}$ decay**
 - Non-trivial interplay from jets in production and decay processes
 - Changes in higher order corrections due to cuts on b-jets are comparable to NNLO corrections
 - Smaller residual scale uncertainty / better perturbative convergence
 - **Future work:** Include decay $H \rightarrow b\bar{b}$ massive @ NNLO
[Bernreuther, Chen, Si '18; Behring, Bizoń '19]
- **WBF including $H \rightarrow WW^* \rightarrow 2l 2\nu$ decay**
 - Not presented in this talk but most important result: Effects of decay small and higher order corrections well captured by simple reweighting (with K-factors computed from stable Higgs boson production)

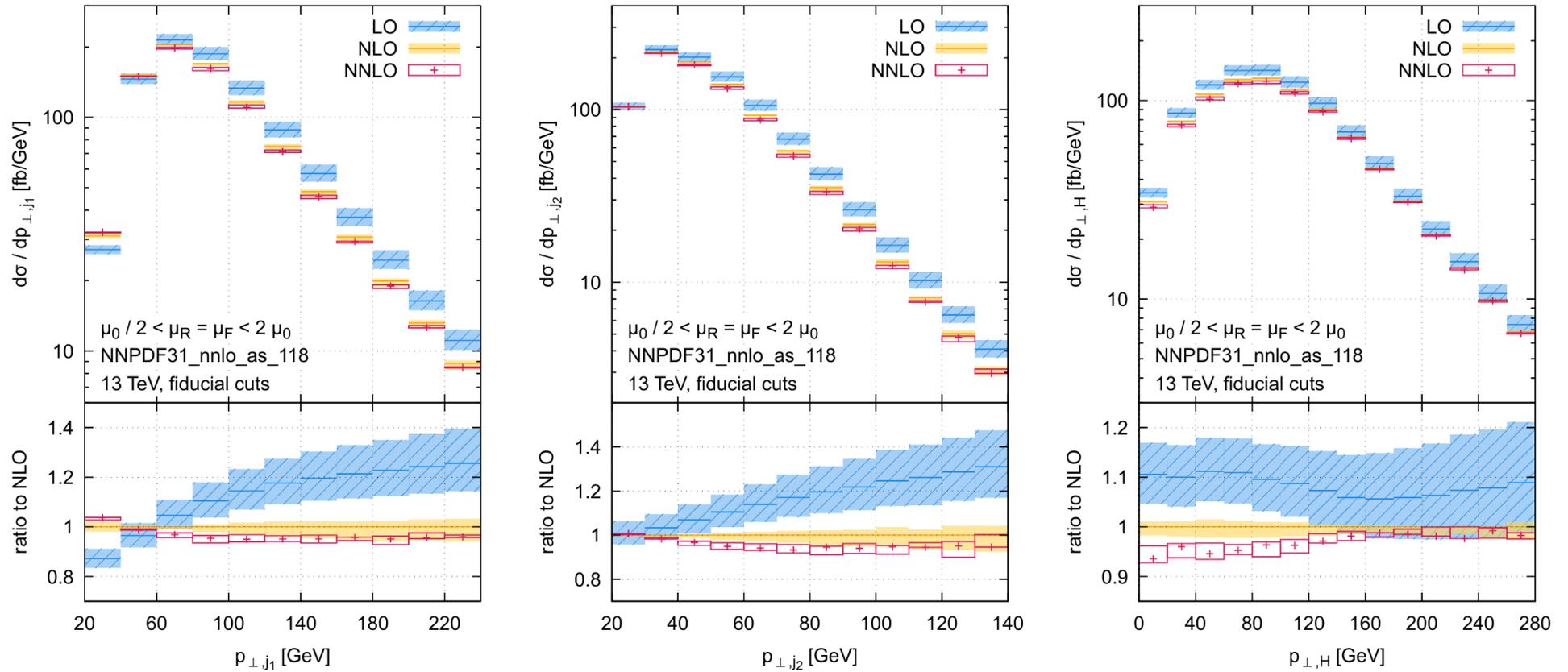
Differential results

Results: differential cross sections

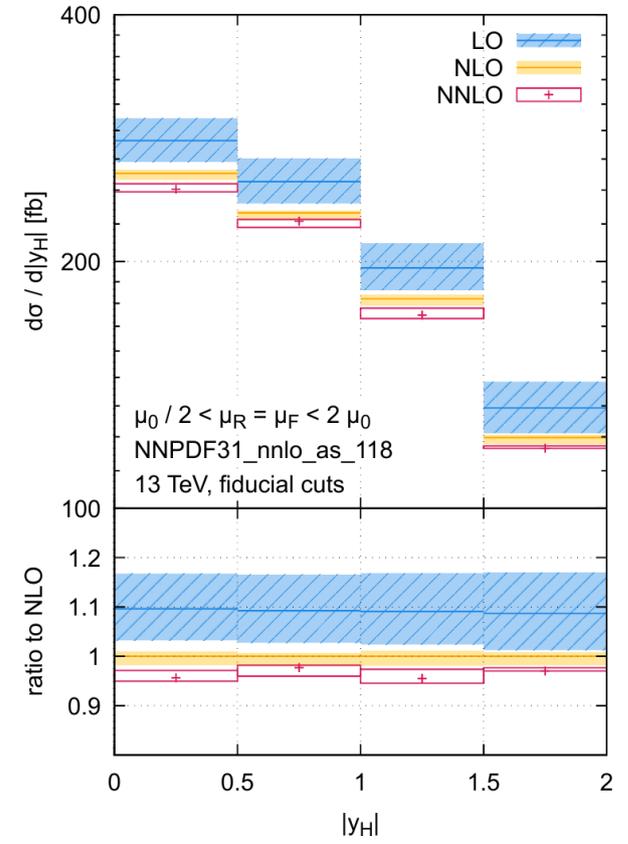
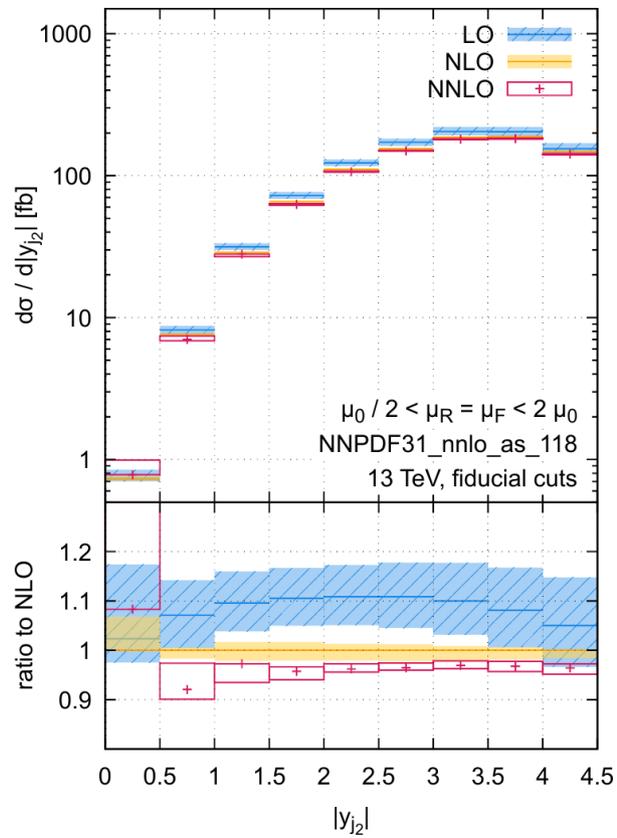
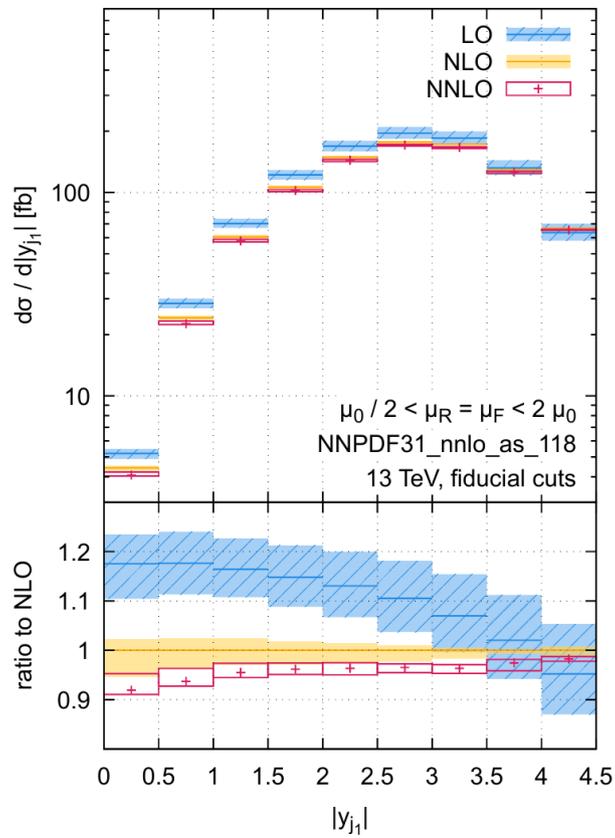


- Differential K-factors rather flat
- NNLO/NLO K-factor computed with stable Higgs is a good approximation within $O(1\%)$ precision

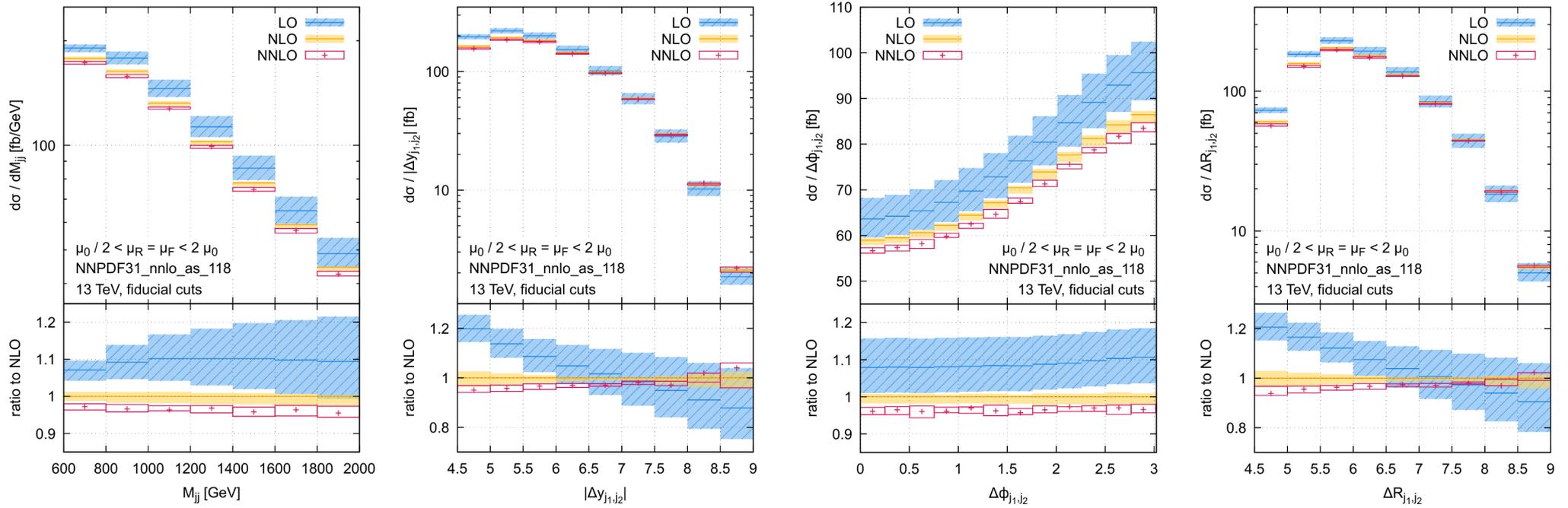
Stable Higgs boson production (I)



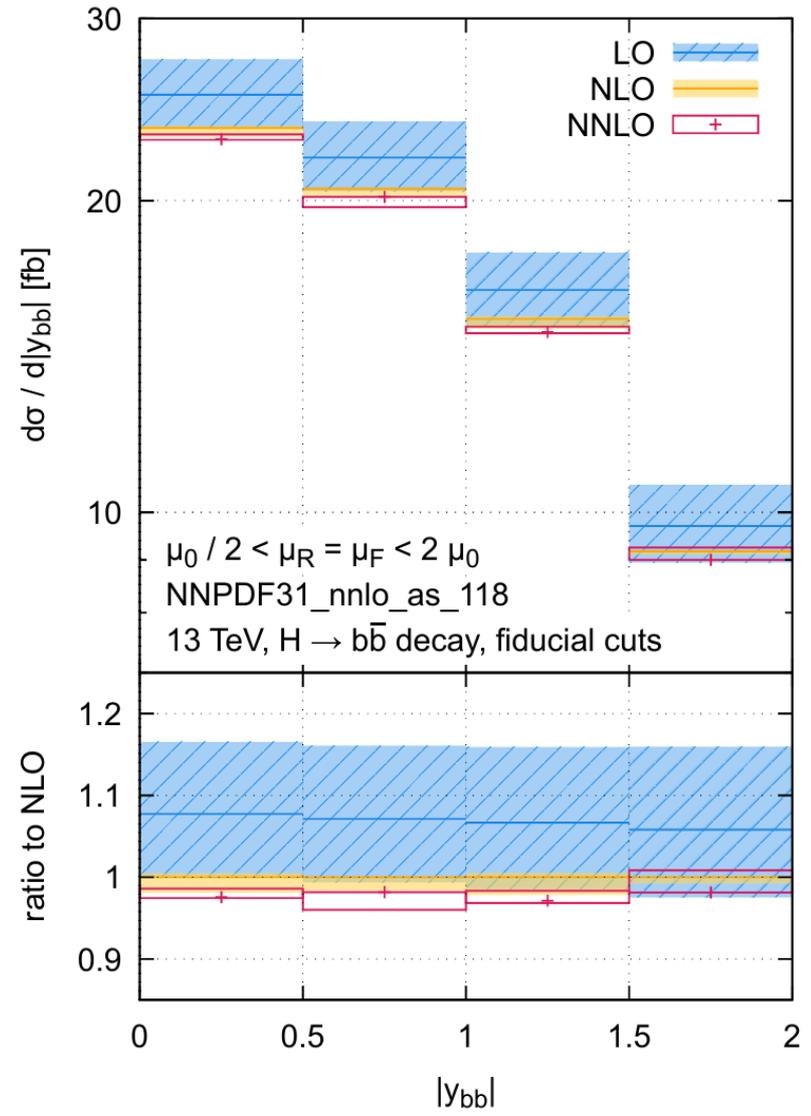
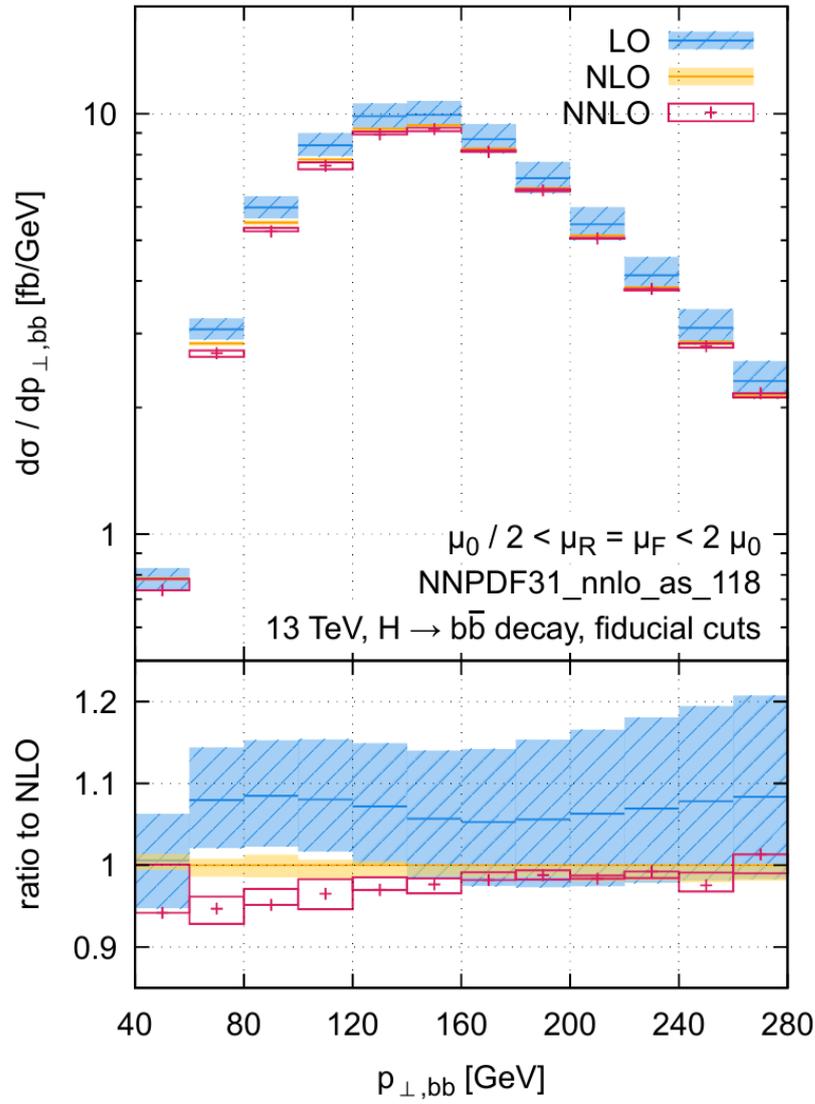
Stable Higgs boson production (II)



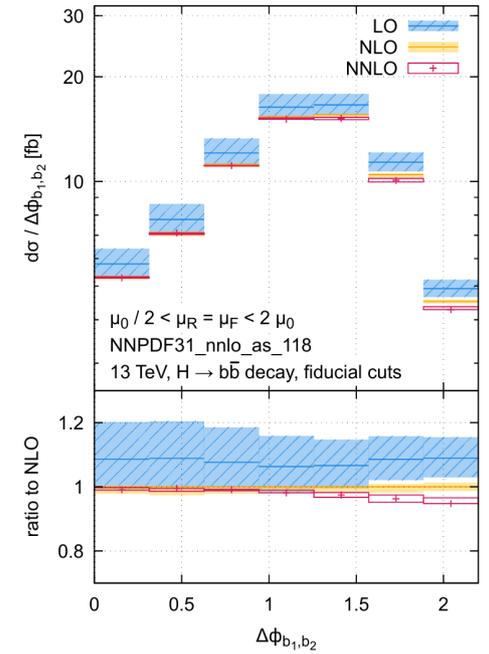
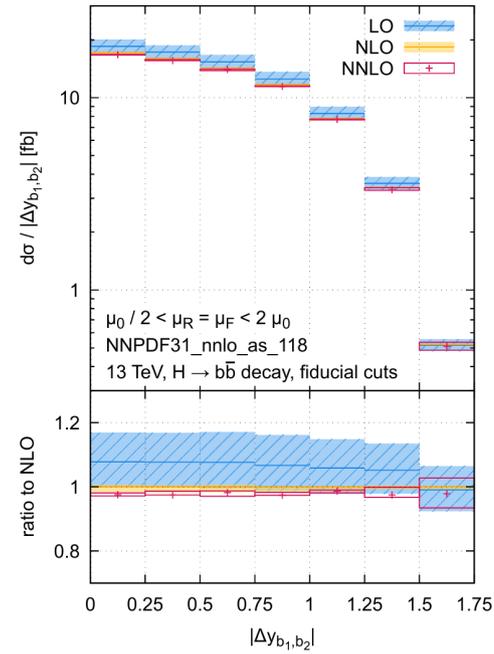
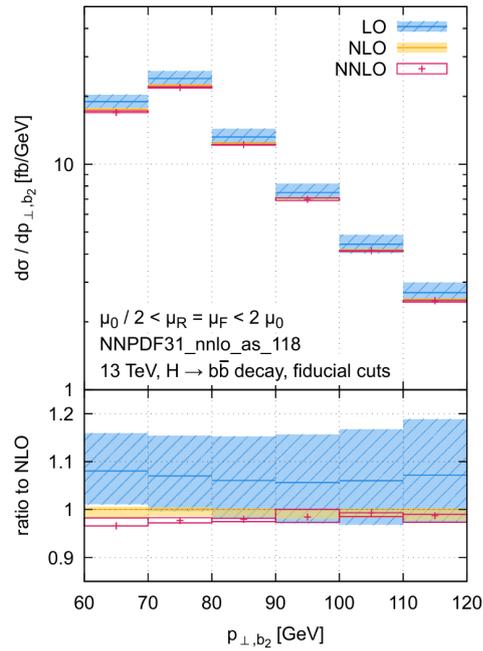
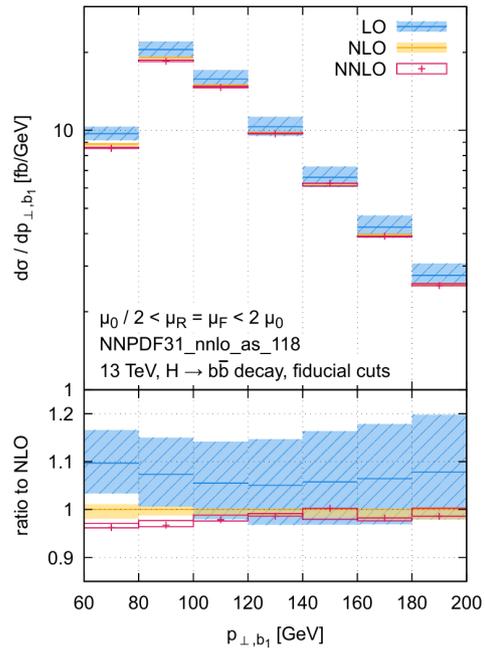
Stable Higgs boson production (III)



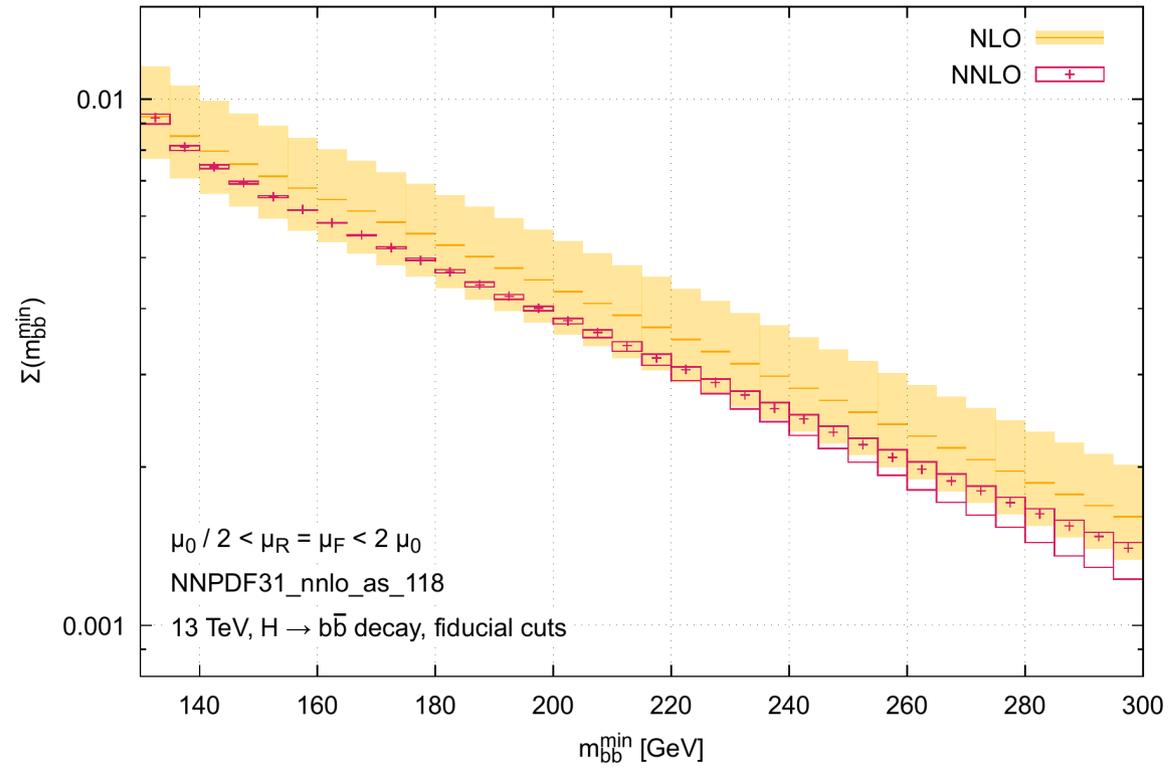
H \rightarrow $b\bar{b}$ decay (I)



H \rightarrow $b\bar{b}$ decay (II)



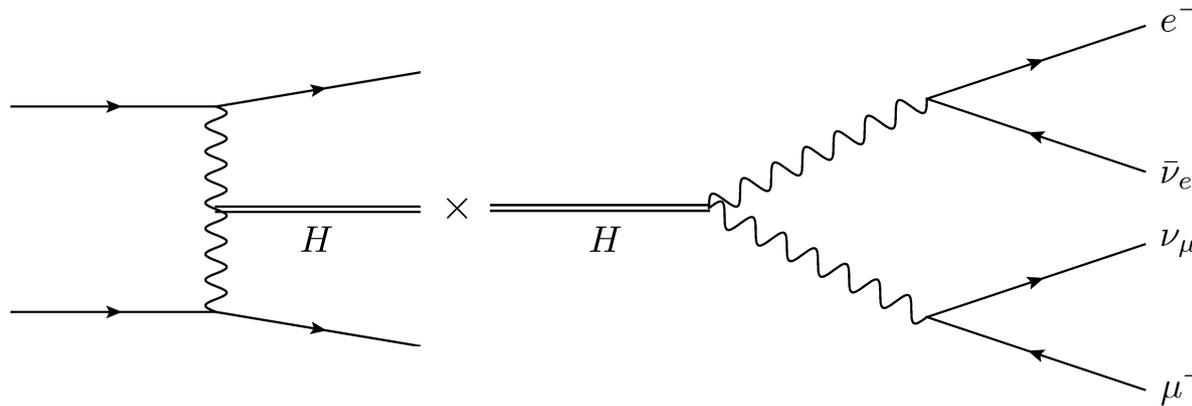
Higgs mass reconstruction



Miscellaneous

WBF + H \rightarrow WW* \rightarrow 2l 2 ν

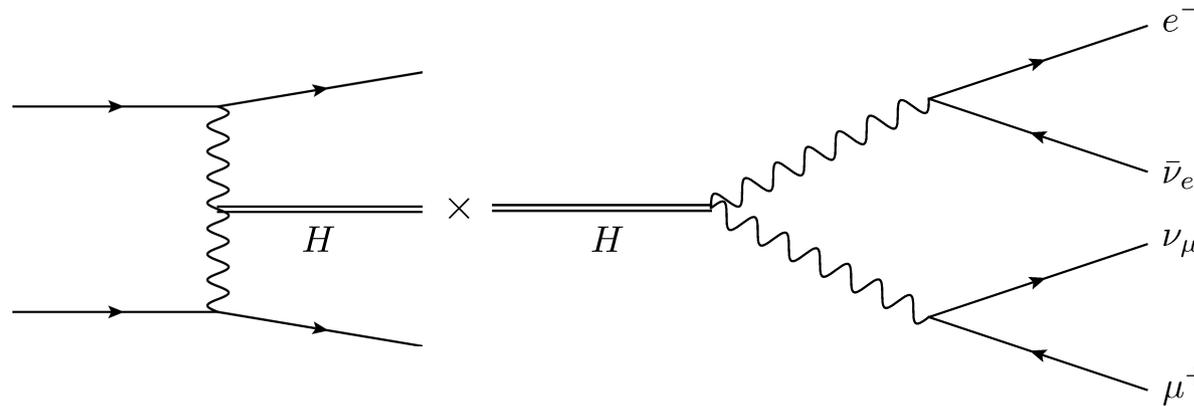
- Integration trick of before enables efficient computations of even more complex final states



- Cuts loosely following latest CMS measurement [Phys. Lett. B 791, 96 (2019)]
 - Leading lepton $p_{\perp, l_1} > 25 \text{ GeV}$ / Subleading lepton $p_{\perp, l_2} > 13 \text{ GeV}$
 - Rapidity of leptons between rapidity of VBF tagged jets
 - Transverse mass: $60 \text{ GeV} < m_T \equiv \sqrt{2p_T^{l_1 l_2} p_T^{\text{miss}} (1 - \cos \Delta\phi_{l_1 l_2, \vec{p}_T^{\text{miss}}})} < 125 \text{ GeV}$
 - Lepton system: $p_{\perp, l_1 l_2} > 30 \text{ GeV}$, $m_{l_1 l_2} > 12 \text{ GeV}$
 - Missing pt: $p_T^{\text{miss}} > 20 \text{ GeV}$

WBF + H → WW* → 2l 2ν

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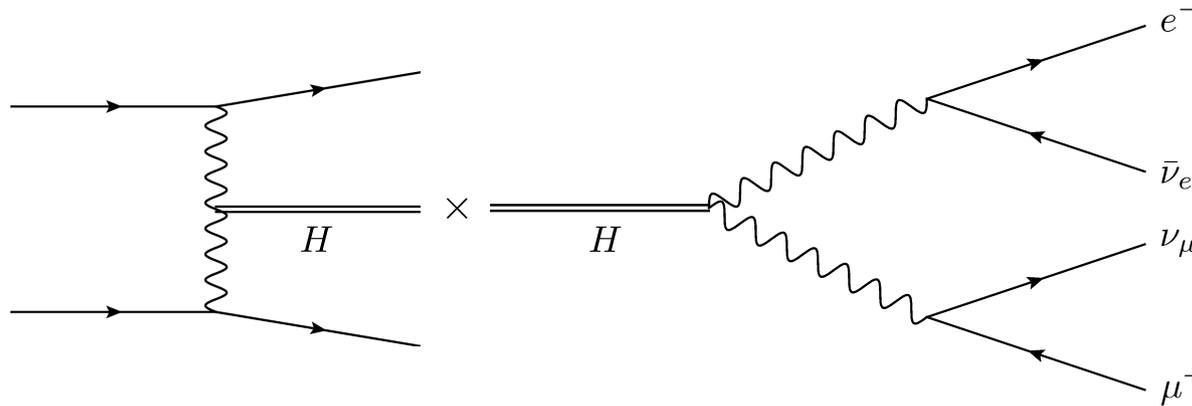
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- Mild cuts on Higgs decay products

→ no significant change in fiducial WBF region expected ...

WBF + H → WW* → 2l 2ν

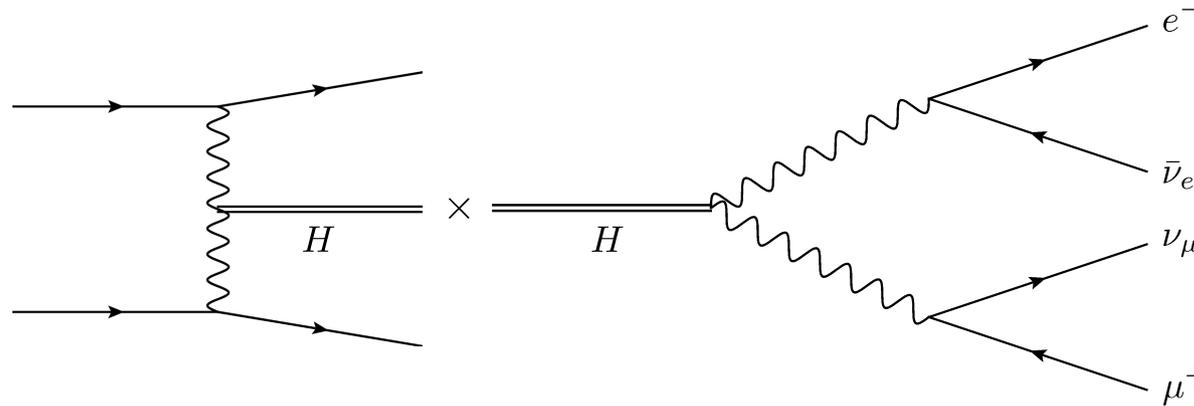
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 - Lepton system: $p_{\perp, l_1 l_2} > 30 \text{ GeV}$, $m_{l_1 l_2} > 12 \text{ GeV}$
 - Missing pt: $p_T^{\text{miss}} > 20 \text{ GeV}$
- Factorization of production and decay »practically« unbroken
 → no huge difference in shapes of decay observables for NLO → NNLO expected ...

WBF + H → WW* → 2l 2ν

- Integration trick of before enables efficient computations of even more complex final states



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 - Leading lepton $p_{\perp, l_1} > 25 \text{ GeV}$ / Subleading lepton $p_{\perp, l_2} > 13 \text{ GeV}$
 - Rapidity of leptons between rapidity of VBF tagged jets

Similar corrections as in case of stable Higgs production found!

$$\sigma_{\text{NNLO}}^{l^- \bar{\nu} l^+ \nu} = \underbrace{0.719 \text{ fb}}_{\text{LO}} + \underbrace{(-0.057) \text{ fb}}_{\Delta\text{NLO}} + \underbrace{(-0.03) \text{ fb}}_{\Delta\text{NNLO}}$$

	-7.9%	-4.2%
(stable Higgs)	-8.3%	-3.2%

Phase space parametrization (partonic center-of-mass frame)

- Consider a single collinear integrated subtraction term (schematically)

$$\int \frac{d^3 \vec{k}_1}{(2\pi)^3 2E_{k_1}} \lim_{k_1 \parallel p_1} \left[\text{Diagram 1} \right]^2 \sim \int^1 dz \left[\frac{1}{1-z} \right]_+ \times \frac{1}{z} \left[\text{Diagram 2} \right]^2 + \dots$$

Diagram 1: A grey circle with an incoming line from the left labeled p_1 and an outgoing line to the right labeled p_2 . A wavy line labeled k_1 enters from the top, and another wavy line labeled k_2 enters from the top. A wavy line labeled *regulated* exits from the bottom.

Diagram 2: A grey circle with an incoming line from the left labeled $z \cdot p_1$ and an outgoing line to the right labeled p_2 . A wavy line labeled k_2 enters from the top, and a wavy line labeled *regulated* exits from the bottom.

- It contains “boosted” and regulated NLO matrix element squared

$$\left[\text{Diagram 3} \right]^2 \approx \left[\text{Diagram 4} \right]^2 - \text{Eikonal}(z \cdot p_1, p_2, k_1) \left[\text{Diagram 5} \right]^2 - \frac{1}{z p_1 \cdot k_2} P_{qq} \left(\frac{z E_1 - E_{k_2}}{z E_1} \right) \left[\text{Diagram 6} \right]^2 - \dots$$

Diagram 3: A grey circle with an incoming line from the left labeled $z \cdot p_1$ and an outgoing line to the right labeled p_2 . A wavy line labeled k_2 enters from the top, and a wavy line labeled *regulated* exits from the bottom.

Diagram 4: A grey circle with an incoming line from the left labeled $z \cdot p_1$ and an outgoing line to the right labeled p_2 . A wavy line labeled k_2 enters from the top, and a wavy line exits from the bottom.

Diagram 5: A grey circle with an incoming line from the left labeled $z \cdot p_1$ and an outgoing line to the right labeled p_2 . A wavy line exits from the bottom.

Diagram 6: A grey circle with an incoming line from the left labeled $x \cdot z \cdot p_1$ and an outgoing line to the right labeled p_2 . A wavy line exits from the bottom.

subtract soft singularity

subtract collinear $k_2 \parallel p_1$ singularity

Phase space parametrization (**hadronic** center-of-mass frame)

- Partonic \rightarrow **hadronic** center-of-mass frame
- Consider a single collinear integrated subtraction term

$$\int dx_1 f(x_1) \int \frac{d^3 \vec{k}_1}{(2\pi)^3 2E_{k_1}} \lim_{k_1 \parallel p_1} \left[\begin{array}{c} \text{diagram with } k_1, k_2, p_1, p_2 \text{ and } x_1 \cdot p_1 \\ \text{regulated} \end{array} \right]^2 \approx \int dx_1 f(x_1) \int^1 dz \left[\frac{1}{1-z} \right]_+ \times \frac{1}{z} \left[\begin{array}{c} \text{diagram with } k_2, p_1, p_2 \text{ and } x_1 z \cdot p_1 \\ \text{regulated} \end{array} \right]^2$$

- Use integration over parton fraction \mathbf{x}_1 to substitute $\mathbf{x}_1 \mathbf{z} \rightarrow \mathbf{y}$ (ignoring finite pieces for $\mathbf{z} < \mathbf{y}$)

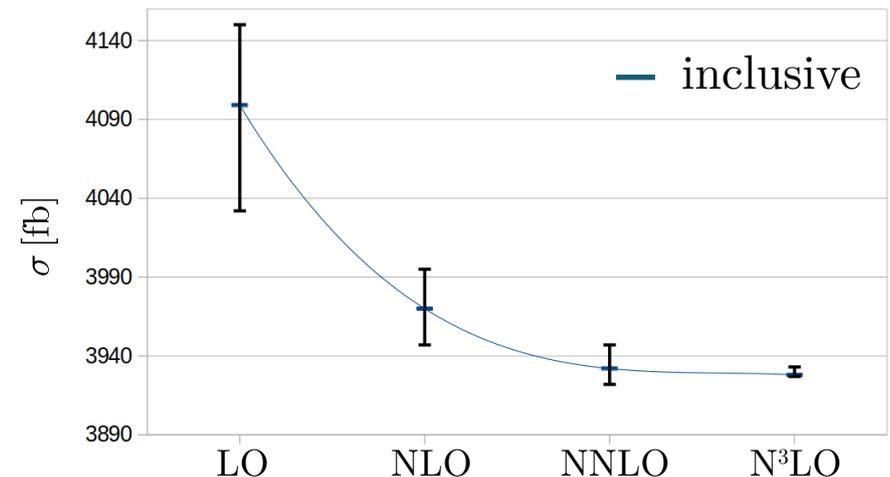
$$\approx \int dy \underbrace{\left[\int_y^1 dz \left[\frac{1}{1-z} \right]_+ \frac{f(y/z)}{z^2} \right]}_{\equiv \tilde{f}(y)} \times \left[\begin{array}{c} \text{diagram with } k_2, p_1, p_2 \text{ and } y \cdot p_1 \\ \text{regulated} \end{array} \right]^2$$

- Integration over “soft singularity” \mathbf{z} factorizes from resolved matrix element
- Use redefined parton distribution functions

State of the art of QCD analysis

- Historically focus on corrections to *factorizable* contributions:
 - **2015 (differential)** “*Fully differential VBF Higgs production at NNLO*” [Matteo Cacciari, Frédéric A. Dreyer, Alexander Karlberg, Gavin P. Salam, Giulia Zanderighi]; Projection-to-Born method; Used for checks of our implementation
 - **2016 (inclusive)** “*Vector-boson fusion Higgs production at N³LO in QCD*” [Dreyer, Karlberg ‘16]
 - **2018 (differential)** “*NNLO corrections to VBF Higgs boson production*” [Juan Cruz-Martinez, Nigel Glover, Thomas Gehrmann, Alexander Huss]; Antenna-subtraction
- Recent growing interest in *non-factorizable* contributions
 - First studies available [Dreyer, Karlberg, Tancredi ‘20; Chen, Figy, Plätzer ‘21]

- **Inclusive:** Nicely converging, N³LO within residual scale uncertainties
- **Fiducial cuts:** NNLO corrections outside of scale uncertainties



State of the art of QCD analysis

[Matteo Cacciari, Frédéric A. Dreyer, Alexander Karlberg, Gavin P. Salam, Giulia Zanderighi]

