

A new parton shower model for Sherpa

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Introduction

The problem

The solution

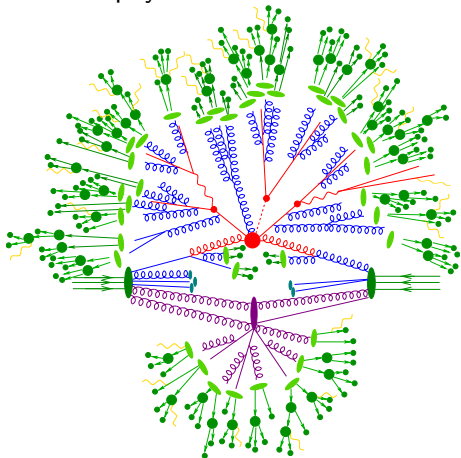
Results

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Introduction

Event Generators crucial for collider physics

Hard Process
Parton Shower
Underlying Event
Hadronization
QED FSR
Hadron Decays



Event Generators crucial for collider physics

Hard Process

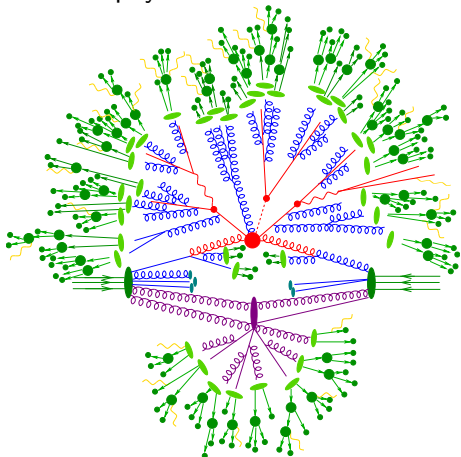
Parton Shower

Underlying Event

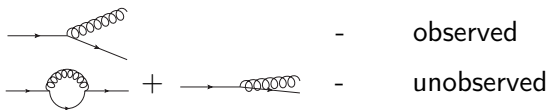
Hadronization

QED FSR

Hadron Decays



Parton branching can occur in two ways

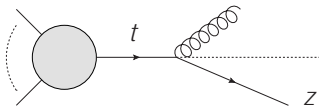


Evolution conserves probability

Decay probability for parton state in collinear limit

$$! \frac{1}{n} \int_t^Z Q^2 d\bar{t} \frac{d}{d\bar{t}} \frac{n+1}{\bar{t}} \times \int_t^Z Q^2 \frac{d\bar{t}}{\bar{t}} \int dz \frac{s}{2} P(z)$$

jets



How to compute probabilities at higher orders Interplay with NLO splitting functions

$$D_{ji}^{(0)}(z;) = ij (1 - z) \quad \leftrightarrow \quad \text{[Diagram: Circle with incoming line } j \text{ and outgoing line } i \text{, split into } z \text{ and } 1-z \text{]} / \text{[Diagram: Circle with incoming line } i \text{ and outgoing line } 1 \text{]}$$

$$D_{ji}^{(1)}(z;) = - \frac{1}{2} P_{ji}^{(0)}(z) \quad \leftrightarrow \quad \text{[Diagram: Circle with incoming line } i \text{ and outgoing line } i \text{, with a gluon emission from the vertex]} / \text{[Diagram: Circle with incoming line } i \text{ and outgoing line } 1 \text{]}$$

$$D_{ji}^{(2)}(z;) = - \frac{1}{2} P_{ji}^{(1)}(z) + \frac{0}{4} P_{ji}^{(0)}(z) + \frac{1}{2} \int_z^1 \frac{dx}{x} P_{jk}^{(0)}(x) P_{ki}^{(0)}(z=x)$$

$$\leftrightarrow \left(\text{[Diagram: Circle with incoming line } i \text{ and outgoing line } i \text{, with a gluon emission from the vertex]} + \text{[Diagram: Circle with incoming line } i \text{ and outgoing line } i \text{, with a gluon emission from the vertex]} \right) / \text{[Diagram: Circle with incoming line } i \text{ and outgoing line } 1 \text{]}$$

Ability to simulate $P_{ji}^{(1)}$ fully differentially prerequisite
to achieve NNLL accuracy for arbitrary observables

Soft and collinear limits schematically identical,
but beware of soft/collinear overlap!

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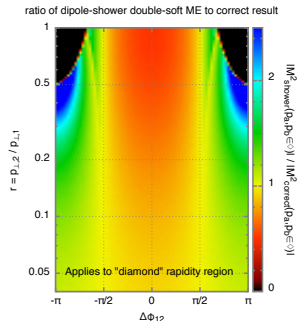
Two problems in commonly used parton showers when compared to analytic NLL resummation:

Angular correlations across multiple emissions due to recoil strategy and choice of evolution variable spoil formal accuracy

Average color charge of initial $q\bar{q}$ dipole after radiation not reflected correctly

Here: address kinematics problem

Transverse recoil should not change hard partons
Angles should be measured in frame of hard event



[Dasgupta, Dreyer, Hamilton, Monni & Salam, JHEP 09 (2018) 033]

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Use mapping where hard directions maintained *exactly*
Guideline: CS identified particle subtraction

[Catani Seymour, Nucl.Phys.B 485 (1997) 291]

Collinear safety criterion

$$p_i \cdot p_j \cdot z \tilde{p}_i ; \quad p_j \cdot p_i \cdot (1 - z) \tilde{p}_i$$

Splitting variable defined by

$$z = \frac{p_i \cdot n}{(p_i + p_j) \cdot n}$$

n is an arbitrary auxiliary vector

$$\tilde{p}_i \cdot n \neq 0$$

New vector \tilde{K} , sum of a subset of the momenta

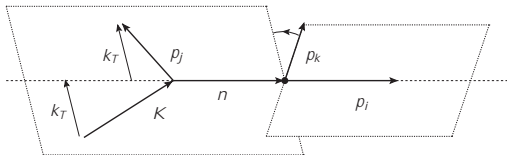
$$\{ \tilde{p}_a; \tilde{p}_a; \tilde{p}_1; \dots; \tilde{p}_{j-1}; \tilde{p}_{j+1}; \dots; \tilde{p}_n \}$$

Shift ensuring momentum conservation

$$p_i = z \tilde{p}_i; \quad n = \tilde{K} + (1 - z) \tilde{p}_i$$

\tilde{K} could be sum of lepton momenta in Drell-Yan

Now: find K and p_j such that $K^2 = \tilde{K}^2$ and collinear safe



Variables

$$v = \frac{p_i p_j}{p_i \tilde{K}} ; \quad = \frac{\tilde{K}^2}{2\tilde{p}_i \tilde{K}}$$

Now we can express p_j and K :

$$p_j = (1 - z) \tilde{p}_i + v \tilde{K} \quad (1 - z + 2) \tilde{p}_i + k_T$$

$$K = \tilde{K} \quad v \tilde{K} \quad (1 - z + 2) \tilde{p}_i \quad k_T$$

! collinear safe since $v \ll E_j(1 - \cos \theta_{ij})$
 n -frame and K -frame agree in soft limit

Angular ordering is approximation to soft coherence

Marchesini & Webber, NPB 310(88)461

Decompose Eikonal

$$W_{ik;j} = \frac{W_{ik;j}}{E_j^2} ; \text{ where } W_{ik;j} = \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(1 - \cos \theta_{jk})}$$

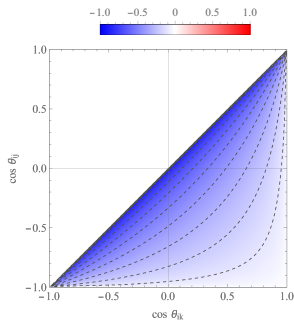
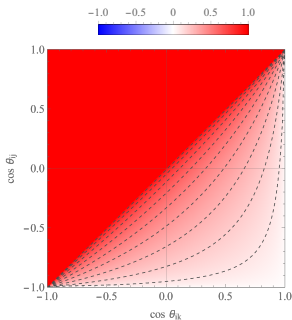
Split into partial radiator functions

$$W_{ik;j} = \tilde{W}_{ik;j}^i + \tilde{W}_{ki;j}^k ;$$

$$\tilde{W}_{ik;j}^i = \frac{1}{2} \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(1 - \cos \theta_{jk})} + \frac{1}{1 - \cos \theta_{ij}} \frac{1}{1 - \cos \theta_{jk}}$$

Azimuthal averaging ! negative and positive contributions cancel outside forward cone

$$\frac{1}{2} \int_0^2 d \cos \theta_{jk} \tilde{W}_{ik;j}^i = \frac{\tilde{l}_{ik;j}^i}{1 - \cos \theta_{jk}} ; \tilde{l}_{ik;j}^i = \begin{cases} 1 & \text{if } j < i \\ 0 & \text{else} \end{cases}$$



Azimuthal modulation needs to be included if observable sensitive to wide-angle soft effects (non-global logarithms)

Alaric approach to soft coherence

A new partial
shower model
for Sherpa

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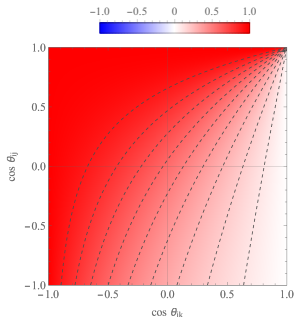
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Capture soft radiation pattern through partial fractioning of angular component of Eikonal

$$\bar{W}_{ik;j}^i = \frac{1 \cos ik}{(1 \cos ij)(2 \cos ij \cos jk)}$$



Advantage: strictly positive radiator, no dead zones

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Full soft PS model given by

$$dP_{ik;j}^i(t; z;) = d_{+1}(f_{pg}; p_j) \frac{s}{2} 2C_i w_{ik;j}^i$$

$$= dt dx \frac{d}{2} \frac{s}{2} \frac{1}{t} 2C_i \frac{(p_i p_k)(p_i n)}{(p_i p_j)(p_k n) + (p_k p_j)(p_i n)}$$

Note: Formalism actually Lorentz invariant

! extendable to initial state?

Extension to initial state works the same as final state
! matches analytic resummation

Currently validating against criteria from

[Dasgupta, Dreyer, Hamilton, Monni, Salam & Soyez, 2002.11114]

$$s(M_Z) = 0 : 005$$

Scheme similar to identified CS subtraction

! matching should be straightforward

ALARIC will become available as part of Sherpa

Include soft/collinear higher-order corrections

[Dulat, Gellersen, Heide & Prestel, 1705.00742, 1805.03757, 2110.05964]