

# The duals of Feynman Integrals

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# Integral reduction

$A_n^{1\text{-loop}} = \mathcal{R}_n + c_{\text{tad}} \text{---} \bigcirc \text{---} + c_{\text{bub}} \text{---} \bigcirc \text{---}$   
 $+ c_{\text{tri}} \text{---} \triangle \text{---} + c_{\text{box}} \text{---} \square \text{---}$

$c_{\text{box}} = \text{---} \square \text{---} = \prod (\text{trees})$  [Britto, Cachazo, Feng 2004]

Generalized unitarity  $\implies$  **Analytic structure** of Feynman integrals (FIs) determined by its **cuts**!

- Residue formulas for sub-topologies hard to find
- Harder when  $d = 4 - 2\epsilon$

**Goal: Systematic formalism for extracting integral coeff in any dimension and for all topologies that maintains the spirit of GU**

# Integral reduction in generic dimension

Perform integral reduction in **generic dimension**

$$\mathcal{A}_n^{1\text{-loop}} = c_{\text{tad}}(\varepsilon) \text{---} \bigcirc \text{---} + c_{\text{bub}}(\varepsilon) \text{---} \bigcirc \text{---} + c_{\text{tri}}(\varepsilon) \triangle + c_{\text{box}}(\varepsilon) \square + c_{\text{pent}}(\varepsilon) \text{pentagon}$$

Want **residue formulas** for  $c_\bullet$  that keep intuition from **unitary cuts**

Traditionally, done via integration by parts (IBP) identities

- Black box
- Squared propagators\*\*
- Cuts are not utilized\*\*

\*\*These problems can be addressed in the IBP framework

# Alternative: Intersection theory

Intersection theory  $\implies$  extraction of  $c_\bullet$  without IBP identities

[Mastrolia, Mizera 2018; Frellesvig, et. al. 2019, 2020; Mizera, AP 2019; Caron-Huot, AP 2021; ...]

FIs **not unique**: shifts by total derivative (IBPs)  $\implies$  **cohomology**

$$|\varphi\rangle \in H^\bullet = \frac{\nabla - \text{closed forms (FIs)}}{\nabla - \text{exact forms (IBPs)}}$$

$\exists$  Poincaré **dual cohomology**  $\langle \varphi^\vee | \in (H^\bullet)^\vee$

**Inner product**  $\langle \varphi^\vee | \varphi \rangle$  integrate the wedge product (intersection number)

$$\langle \varphi_i^\vee | \varphi_j \rangle = \delta_{ij} \quad \implies \quad |\Phi\rangle = \sum_i |\varphi_i\rangle \underbrace{\langle \varphi_i^\vee | \Phi \rangle}_{\text{e.g., } c_{\text{box}}}$$

Computation of  $\langle \bullet | \bullet \rangle$  is **algebraic (residues)**  $\implies$  right direction

## What is the dual space?

What is  $(H^\bullet)^\vee$ ?

# An unsatisfactory answer

$(H^\bullet)^\vee$  well known in math lit for a deformed problem

[Aomoto, Kita, Matsumoto, Yoshida 80's-present]

All propagators raised to **non-integer** powers ( $\varepsilon\alpha \notin \mathbb{Z}$ )

$\implies$  propagators (poles)  $\rightarrow$  branch pts

$$\frac{1}{\ell^2 + m^2} \rightarrow \frac{1}{(\ell^2 + m^2)^{1+\varepsilon\alpha}} \quad (\alpha \rightarrow 0 \text{ before } \varepsilon \rightarrow 0 \text{ at end})$$

Still missing the spirit of generalized unitarity:

- propagators (poles)  $\rightarrow$  branch pts **obscures connection to GU cuts**

# The answer for the undeformed dual space

[Caron-Huot, AP: 2104.06898, 2112.00055]

# The undeformed dual space

Dual forms **localized** to GU cuts:

$$\frac{\text{FIs: } |\varphi\rangle}{\text{propagators/poles}} \leftrightarrow \frac{\text{dual forms: } \langle\varphi^\vee|}{\underbrace{\text{boundaries/zeros (on-shell)}}_{\text{compact support}}}$$

The dual cohomology spanned by holomorphic forms times:

- $\theta(|\text{propagators}| > \epsilon) :$



- $d\theta(|\text{propagators}| > \epsilon) :$



$\theta$  or  $d\theta$  for each propagator



# The undeformed dual space

twist ( $d \rightarrow \nabla$ )

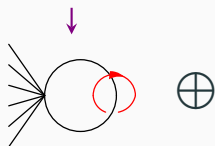
$$d^d l \propto \overbrace{(\ell_{\perp}^2)^{-\varepsilon}}^{\text{branch pts}}$$

$$\overbrace{\frac{d\ell_{\perp}^2 \wedge d^4 l}{\ell_{\perp}^2}}^{\text{5-form}}$$

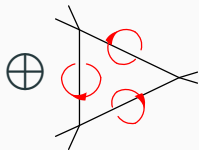
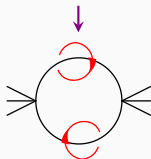
$$\text{red circle} \sim d\theta(\text{propagator})$$

$$\varphi_{1\text{-loop}}^{\nabla} \sim \text{5-forms on 0-cuts} \downarrow \text{0} \oplus$$

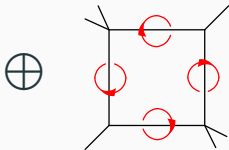
4-forms on 1-cuts



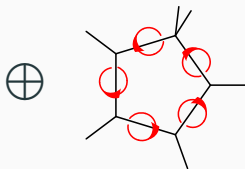
3-forms on 2-cuts



2-forms on 3-cuts



1-forms on 4-cuts



0-forms on 5-cuts

# Streamlining the intersection number

$$\begin{aligned}
 c_{\text{box}}(\varepsilon) &= \left\langle \overbrace{\text{Localized/on-shell dual form}}^{\text{Localized/on-shell dual form}} \mid \overbrace{\text{Any (1-loop) process}}^{\text{Any (1-loop) process}} \right\rangle \\
 &\sim (d\theta)^4 \wedge \left[ \frac{d^1 \ell_{\parallel}}{(\ell_{\perp}^2|_{\text{box}})^1} \right]_{\text{c-map}} \\
 &= \left\langle \varphi_{\text{left-over}}^{\vee}(\ell_{\perp}^2|_{\text{box}}) \mid \underbrace{\text{Res}_{\text{box}} \left[ \text{Any (1-loop) process} \right]}_{\text{Factorizes}} \right\rangle \\
 &= \text{rational fn. of Mandelstams and } \varepsilon
 \end{aligned}$$

Left-over  $\implies$   $d\theta$ 's at branch points  $(\ell_{\perp}^2|_{\text{box}} = 0, \infty)$  from c-map  
 (essential for higher order terms in  $\varepsilon$ )

# Outline

## 1. Workflow and key steps

- DEQs
- Compute intersection number (c-map)

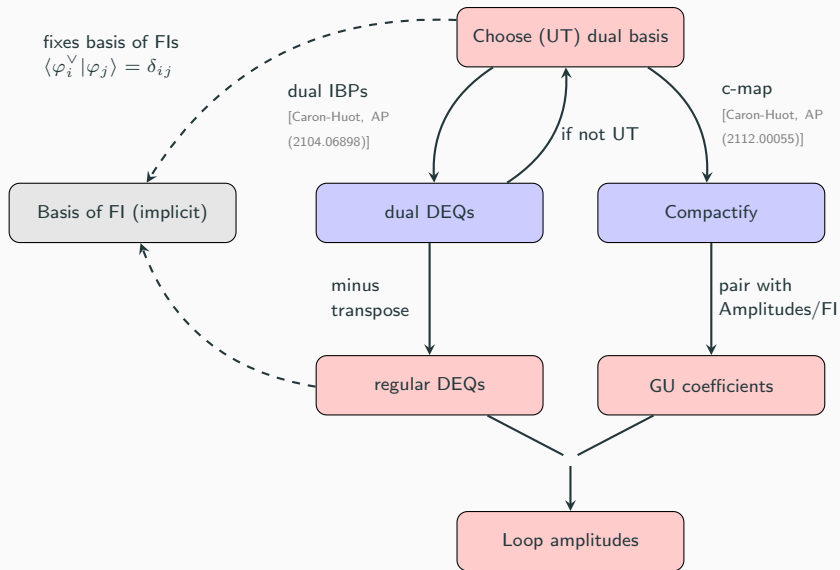
## 2. New technology

- Where are we in the development phase?
- Roadblock: Multi-variate intersections are difficult

## 3. Conclusions and future directions



# Workflow



# Dual DEQs are transposed

$$\nabla^V \left( \theta \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right) = \theta \left( \nabla^V \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right) + d\theta \wedge \left( \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right)$$

Dual IBP-forms are simple: **supported on cuts** and **no squared propagators!**

Obtained simple form for canonical differential equations in **any** dimension for **both** 1-loop dual and Feynman forms

[Caron-Huot, AP (2104.06898); Bourjaily, Gardi, McLeod, Vergu (2020); Abreu, Britto, Duhr, Gardi (2017) Volovich, Spradlin (2011); Schläfli (1860)]

# Compactifying (c-map)

$\varphi_c^\vee$  must have compact support for well defined intersection number

$$\langle \varphi^\vee | \varphi \rangle \propto \int_{\mathbb{C}^5} \varphi_c^\vee \wedge \varphi$$

**Compact support** regulates the singularities of  $\varphi$

Given  $\varphi^\vee$  with partial compact support  $\exists \varphi_c^\vee$  such that  $\varphi^\vee - \varphi_c^\vee = \nabla^\vee(\bullet)$

Construct **local primitives** to remove support

$$\varphi_c^\vee = \varphi^\vee - d\theta(\text{other propagators})\psi^\vee - d\theta(\text{branch pts})\psi^\vee$$
$$\nabla^\vee \psi^\vee \equiv \varphi^\vee \text{ near other propagators and branch pts}$$

as a Laurent series

Trivial for 1-forms (integrate order by order), harder for ( $p > 1$ )-forms

For  $(p > 1)$ -dual forms, there are 2 options:

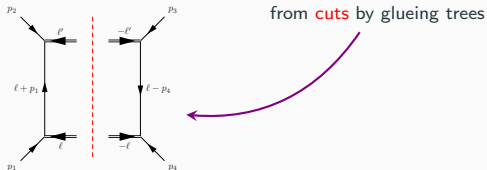
- **Fibration** [Mizera 2019]: 1-variable at a time
  - 😊 Reuse simplicity of 1-form c-map
  - 😞 Matrix/vector valued
  - 😞 Construct  $\varphi_c^\vee$  for dual basis all at once
  - 😞 Singularities at  $\infty$  (from numerators) get progressively worse
  - 😞 Choice: coordinates and order of fibration
- **Combinatoric** [Matsumoto 80's]: deal with the multi-variate nature
  - 😊 Construct  $\varphi_c^\vee$  individual dual forms
  - 😞 Need primitives for all ways of approaching a singularity
  - 😞 Need primitives for the primitives
  - 😞 Doable for  $p = 2$
  - 😞 Choice: coordinates and which primitives are used where

# Putting things together: 1-Loop amplitudes

$$\mathcal{A}_4^{h_1 h_2 h_3 h_4} = c_{\text{bub}_s} \left( \text{Diagram 1} \right) + c_{\text{bub}_t} \left( \text{Diagram 2} \right) + c_{\text{box}} \left( \text{Diagram 3} \right)$$

$$\begin{aligned} \mathcal{A}_5^{h_1 h_2 h_3 h_4 h_5} = & \underbrace{\times 5}_{c_{\text{bub}_{s_{12}}}} \left( \text{Diagram 4} \right) + \dots + \underbrace{\times 5}_{c_{\text{box}_{45}}} \left( \text{Diagram 5} \right) \\ & + \dots + \underbrace{\times 1}_{c_{\text{pent}}} \left( \text{Diagram 6} \right) \end{aligned}$$

$$c_{\bullet} = \left\langle d\theta_{\text{cut}}(\text{dual form}) \middle| \text{Feynman form} \right\rangle = \left\langle \text{dual form} \middle| \underbrace{\text{Res}_{\text{cut}}(\text{Feynman form})}_{\text{from cuts by gluing trees}} \right\rangle$$





# Massless 1-Loop amplitudes and 4-dimensional limits

Proof of concept: reproduced 4- and 5-point gluon amplitudes at 1-loop

[Bern, Dixon, Kosower (90's), see 2112.00055 for details]

More interesting, systematic study of 4d-limit

$$\frac{c_s^{(\text{tri})}}{\varepsilon^2} + \frac{c_s^{(\text{bub})}}{\varepsilon} + c_s^{(\text{rat})} \qquad \qquad \qquad \frac{c^{(\text{box})}}{\varepsilon^2}$$
  
$$\mathcal{A}_4 = c_{\text{bub}_s}(\varepsilon) \left( \text{diagram 1} \right) + c_{\text{bub}_t}(\varepsilon) \left( \text{diagram 2} \right) + c_{\text{box}}(\varepsilon) \left( \text{diagram 3} \right)$$
  
$$-2 + 2\varepsilon \log(-s) + \varepsilon^2(\zeta_2 - \log^2(-s)) \qquad \qquad \qquad -\varepsilon^2(\pi^2 + \log^2(t/s))$$

The diagrams are:   
1. A bubble diagram with two external lines on the left and two on the right, connected by two arcs.   
2. A bubble diagram with two external lines on the top and two on the bottom, connected by two arcs.   
3. A box diagram with four external lines, one on each side, connected by four lines forming a square.

# Massless 1-Loop amplitudes

$$c_{\text{bub}_s}[\varphi](\varepsilon) = \sum_{(\ell_{\perp}^2, \ell_{-}, \ell_{+})} \text{Res}_{\ell_{-}} \text{Res}_{\ell_{+}} \text{Res}_{\ell_{\perp}^2} \left[ \underbrace{\psi_{(\ell_{\perp}^2, \ell_{+}, \ell_{-})}^{\vee}}_{\text{triple primitive of bub-dual}} (\text{Res}_{s\text{-channel}} \varphi) \right]$$

		$(0, \infty)$	$(\infty, \infty)$	$(D_{\text{tri}}, \text{soft})$	$(D_{\text{tri}}, \infty)$
$\varepsilon^{-2}$	$\ell_{\perp}^2 = 0$ $\ell_{\perp}^2 = \infty$			✓	
$\varepsilon^{-1}$	$\ell_{\perp}^2 = 0$ $\ell_{\perp}^2 = \infty$		✓		
$\varepsilon^0$	$\ell_{\perp}^2 = 0$ $\ell_{\perp}^2 = \infty$		✓ ✓		✓ ✓

Better coordinates? [Badger 2008]

Next goal: 2-loop rational coefficients

# Conclusions and Future Directions

Recovered generalized unitarity in arbitrary dimension from twisted relative cohomology

1-loop warm up : [S. Caron-Huot, AP (2021); S. Caron-Huot, AP (2021)]

- 1-loop uniform transcendental basis of dual forms
- Obtained DEQ for this basis in any dimension
- Constructed 1-loop 4- and 5-point gluon amplitudes from their cuts (coeffs **all orders** in  $\epsilon$ )
- 4-dimensional limit extract rational terms ( $\epsilon/\epsilon$ )
- Robust checks at intermediate steps

**Next goals:**

- 2-loops (loop-by-loop) [M. Giroux, AP (ongoing)]
- Improve c-map efficiency (what are good choices?)
- Rational terms at 2-loops
- Dual IBP identities are simpler?