Non-Gaussianity in CMB lensing

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ESA and the Planck Collaboration



Planck Collaboration: Planck 2018 lensing







CMB lensing is weakly non-Gaussian



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CMB lensing is weakly non-Gaussian



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CMB lensing is weakly non-Gaussian

- Lensing is a non-linear process
- Large scale structure growth is non-linear in small scales (gravitational collapse)



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CMB lensing is weakly non-Gaussian

- Lensing is a non-linear process
- Large scale structure growth is non-linear in small scales (gravitational collapse)
- Non-Gaussian introduced from systematics (reconstruction noise, masking, anisotropic beam, foregrounds....)



ESA and the Planck Collaboration





Planck CMB lensing reconstruction



Planck CMB lensing reconstruction



Planck Collaboration: *Planck* 2018 lensing

 σ_T in μ K-arcmin



Planck CMB lensing reconstruction



Planck Collaboration: *Planck* 2018 lensing

The non-Gaussianity in Planck lensing reconstruction is dominated by the noise

 σ_T in μ K-arcmin









decomposing by scales.















 $\chi = #$ disconnected white – #black







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A_{s} , Ω_{cdm} & Σm_{ν} constrain with Planck CMB lensing

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Subtract non-Gaussian information with measurements from Planck like realisations (FFP10)



A_s , Ω_{cdm} & Σm_{ν} constrain with Planck CMB lensing



Subtract non-Gaussian information with measurements from Planck like realisations (FFP10)



Non-Gaussian from lensing itself

Source:

- Non-linear large scale structure growth
- Non-linear lensing process

Method:

- Measure the non-Gaussianity of simulated maps

Simulating CMB lensing by ray-tracing through N-body simulation

Non-Gaussian from lensing itself

• The lensing convergence field $\kappa(\theta)$ in lowest order

$$\kappa(\boldsymbol{\theta}) = \frac{3H_0^2 \Omega_{\rm m}}{2} \int_0^{z_{\rm s}} \frac{dz}{H(z)} \frac{r(z)(r_{\rm s} - r(z))}{a(z)r_{\rm s}} \delta(r\boldsymbol{\theta}, r; z)$$

- r(z) : comoving distance to the redshift z
- $\delta(x;z)$: matter density contrast at position x at z
- *z_s* : Source redshift

Non-Gaussian from lensing itself

• The lensing convergence field $\kappa(\theta)$ in lowest order

$$\kappa(\boldsymbol{\theta}) = \frac{3H_0^2 \Omega_{\rm m}}{2} \int_0^{z_{\rm s}} \frac{dz}{H(z)} \frac{r(z)(r_{\rm s})}{a(z)}$$

$$\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_{k} \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - \frac{1}{2}\right)}{a_k r_k}$$

 $\frac{-r(z)}{z} \delta(r\theta, r; z)$

 $\frac{-r_k}{\delta_k}\delta_k$



A snapshot of N-body at redshift z_k (With periodic boundary)

 $\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_k \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$



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Stacking

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Cutting shell from big snapshot box



A snapshot of N-body at redshift z_k (With periodic boundary)

Stacking

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 $\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_{k} \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$



sphere



A snapshot of N-body at redshift z_k (With periodic boundary)

Stacking

 $\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_{r} \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$

Replace large scale (I>40) with Gaussian realisations



Cutting shell from big snapshot box

Projecting matter distribution to a sphere











 $\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_{k} \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$



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 $\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_{k} \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$



 $\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_{l} \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$

CMB lensing convergence





$$\kappa \approx \frac{3H_0^2 \Omega_{\rm m}}{2} \sum_k \frac{D_k}{H_k} \frac{r_k \left(r_{\rm s} - r_k\right)}{a_k r_{\rm s}} \delta_k$$

Non-Gaussian from full sky simulation

$$\begin{split} S^{(0)} &= \frac{\langle f^3 \rangle_{\rm c}}{\sigma_0^4}, \quad S^{(1)} = \frac{3}{2} \frac{\langle f |\nabla f|^2 \rangle_{\rm c}}{\sigma_0^2 \sigma_1^2}, \\ S^{(2)} &= -3 \frac{\langle |\nabla f|^2 \Delta f \rangle_{\rm c}}{\sigma_1^4}, \end{split}$$

$$\begin{split} K^{(0)} &= \frac{\langle f^4 \rangle_{\rm c}}{\sigma_0^6}, \quad K^{(1)} = 2 \frac{\langle f^2 |\nabla f|^2 \rangle_{\rm c}}{\sigma_0^4 \sigma_1^2}, \\ K_1^{(2)} &= -\frac{4 \langle f |\nabla f|^2 \Delta f \rangle_{\rm c} + \langle |\nabla f|^4 \rangle_{\rm c}}{\sigma_0^2 \sigma_1^4}, \\ K_2^{(2)} &= -\frac{4 \langle f |\nabla f|^2 \Delta f \rangle_{\rm c} + 2 \langle |\nabla f|^4 \rangle_{\rm c}}{\sigma_0^2 \sigma_1^4}, \end{split}$$

$$\begin{split} K^{(0)} &= \frac{\langle f^4 \rangle_{\rm c}}{\sigma_0^6}, \quad K^{(1)} = 2 \frac{\langle f^2 |\nabla f|^2 \rangle_{\rm c}}{\sigma_0^4 \sigma_1^2}, \\ K_1^{(2)} &= -\frac{4 \langle f |\nabla f|^2 \Delta f \rangle_{\rm c} + \langle |\nabla f|^4 \rangle_{\rm c}}{\sigma_0^2 \sigma_1^4}, \\ K_2^{(2)} &= -\frac{4 \langle f |\nabla f|^2 \Delta f \rangle_{\rm c} + 2 \langle |\nabla f|^4 \rangle_{\rm c}}{\sigma_0^2 \sigma_1^4}, \end{split}$$

and first moment of the field

σ_0 and σ_1 are the standard deviation

Non-Gaussian from full sky simulation



Compare with Planck CMB lensing reconstruction



Non-Gaussian information is dramatically dilute by the Gaussian reconstruction noise



Summary

- With decomposition of map, the Minkowski functionals have competitive constrain power compare to power spectrum result.
- For Planck CMB lensing reconstruction, the magnitude of non-Gaussian information in lensing signal is much smaller than the noise.
- The future higher resolution CMB observations like CMB-S4 could give higher signal to noise , which means non-Gaussianity in CMB lensing is still a promising probes for cosmology.